Stability and instability of Ricci solitons

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Received: 11 April 2014 / Accepted: 22 May 2014 / Published online: 18 June 2014 © Springer-Verlag Berlin Heidelberg 2014

Abstract We consider the volume-normalized Ricci flow close to compact shrinking Ricci solitons. We show that if a compact Ricci soliton (M, g) is a local maximum of Perelman's shrinker entropy, any normalized Ricci flow starting close to it exists for all time and converges towards a Ricci soliton. If g is not a local maximum of the shrinker entropy, we show that there exists a nontrivial normalized Ricci flow emerging from it. These theorems are analogues of results in the Ricci-flat and in the Einstein case (Haslhofer and Müller, arXiv:1301.3219, 2013; Kröncke, arXiv:1312.2224, 2013).

Mathematics Subject Classification 53C44 · 58E11 · 37C75

1 Introduction

A Riemannian manifold (M, g) is called a Ricci soliton if there exist a vector field $X \in \mathfrak{X}(M)$ and a constant $c \in \mathbb{R}$ such that

$$\operatorname{Ric}_g + L_X g = c \cdot g.$$

The soliton is called gradient, if $X = \operatorname{grad} f$ for some $f \in C^{\infty}(M)$. We call (M, g) expanding, if c < 0, steady, if c = 0 and shrinking, if c > 0. If X = 0, we recover the definition of an Einstein metric with Einstein constant c. If (M, g) is not Einstein, we call the soliton nontrivial.

Ricci solitons were first introduced by Hamilton in the eighties [21]. They appear as self-similar solutions of the Ricci flow

$$\dot{g}(t) = -2\mathrm{Ric}_{g(t)}.$$

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Communicated by M. Struwe.

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More precisely, a Ricci flow starting at a Ricci soliton g evolves by

$$g(t) = (1 - 2ct)\varphi_t^*g$$

where φ_t is a family of diffeomorphisms on *M*. In other words, Ricci solitons are fixed points of the Ricci flow considered as a dynamical system in the space of metrics modulo diffeomorphism and rescaling. Ricci solitons satisfy many interesting geometric properties and have been studied extensively in recent research (see [5] for a survey). They also appear in string theory [17,29].

Compact Ricci solitons satisfy additional properties. In dimensions n = 2, 3, any compact Ricci soliton is of constant curvature [21,28]. By Perelman's work [38], they are always gradient (see also [16, Proposition 3.1]). Moreover, any expanding or steady compact Ricci soliton is necessarily Einstein [5, Proposition 1.1]. In dimensions $n \ge 4$ there exist examples of nontrivial compact Ricci solitons [4,9,14,34,39,43]. All known nontrivial compact Ricci solitons are Kähler.

In this work, we are going to study the Ricci flow as a dynamical system close to compact shrinking Ricci solitons. Previously the behaviour of Ricci flow close to Ricci-flat metrics has been studied by various authors [19,23,26,41]. The Einstein case was studied in [44] and in a recent paper by the author [36].

In the context of Ricci solitons, it is more convenient to deal with the volume-normalized Ricci flow

$$\dot{g}(t) = -2\operatorname{Ric}_{g(t)} + \frac{2}{n} \frac{1}{\operatorname{vol}(M, g(t))} \left(\int_{M} \operatorname{scal}_{g(t)} dV_{g(t)} \right) g(t),$$
(1.1)

where $\operatorname{scal}_{g(t)}$, $dV_{g(t)}$ and $\operatorname{vol}(M, g(t))$ are the scalar curvature, the volume element and the total volume with respect to g(t), respectively.

Definition 1.1 A compact Ricci soliton (M, g) is called dynamically stable if for any C^k -neighbourhood \mathcal{U} of g in the space of metrics (where $k \ge 3$), there exists a C^{k+5} -neighbourhood $\mathcal{V} \subset \mathcal{U}$ such that for any $g_0 \in \mathcal{V}$, the normalized Ricci flow (1.1) starting at g_0 exists for all $t \ge 0$ and converges modulo diffeomorphism to an Einstein metric in \mathcal{U} as $t \to \infty$.

We call a compact Ricci soliton (M, g) dynamically unstable if there exists a nontrivial normalized Ricci flow defined on $(-\infty, 0]$ which converges modulo diffeomorphism to g as $t \to -\infty$.

In his pioneering work [38], Perelman introduced the shrinker entropy ν , which is a functional on the space of metrics and admits precisely the shrinking Ricci solitons as its critical points. Perelman discovered the remarkable property that ν is nondecreasing under the Ricci flow and stays constant only at its critical points. The shrinker entropy is the most important tool in proving our main theorems.

Theorem 1.2 (Dynamical stability) Let (M, g) be a compact shrinking Ricci soliton. If (M, g) is a local maximizer of v, it is dynamically stable.

Note that we do not need to assume that (M, g) is a strict local maximizer in the space of metrics modulo diffeomorphism and rescaling.

Theorem 1.3 (Dynamical instability) Let (M, g) be a compact shrinking Ricci soliton. If (M, g) is not a local maximizer of v, it is dynamically unstable.

Observe that any compact shrinking Ricci soliton must be either dynamically stable or unstable, so we have a complete description of the Ricci flow as a dynamical system close to Ricci solitons. Since non-shrinking compact Ricci solitons are Einstein, they are already covered by previous results.

In the Einstein case, we additionally were able to give geometric stability/instability conditions in terms of the conformal Yamabe invariant and the Laplace spectrum [36]. It would be intersting to find similar geometric conditions in the case of nontrivial Ricci solitons. So far, we can characterize stability in terms of the eigenvalues of a second-order differential operator, supposed that an additional technical condition holds.

Theorem 1.4 Let (M, g) be a compact shrinking Ricci soliton and suppose all infinitesimal solitonic deformations are integrable. Then (M, g) is a local maximum of v if and only if $v'' \leq 0$.

The integrability condition means that all elements in the kernel of v'' can be integrated to curves of Ricci solitons. We show that this condition does not always hold.

Theorem 1.5 The complex projective space ($\mathbb{C}P^n$, g_{fs}) with the Fubini-Study metric admits nonintegrable infinitesimal solitonic deformations.

2 Notation and conventions

For the Riemann curvature tensor, we use the sign convention such that $R_{X,Y}Z = \nabla_{X,Y}^2 Z - \nabla_{Y,X}^2 Z$. Given a fixed metric, we equip the bundle of (r, s)-tensor fields (and any subbundle) with the natural pointwise scalar product induced by the metric. By $S^p M$, we denote the bundle of symmetric (0, p)-tensors. For a given $f \in C^{\infty}(M)$, we introduce some f-weighted differential operators. The f-weighted Laplacian (or Baker-Emery Laplacian) acting on $C^{\infty}(S^p M)$ is

$$\Delta_f h = -\sum_{i=1}^n \nabla_{e_i, e_i}^2 h + \nabla_{\operatorname{grad} f}^2 h.$$

By the sign convention, $\Delta_f = (\nabla_f^*)\nabla$, where ∇_f^* is the adjoint of ∇ with respect to the weighted L^2 -scalar product $\int_M \langle ., . \rangle e^{-f} dV$. The weighted divergence $\delta_f : C^{\infty}(S^p M) \to C^{\infty}(S^{p-1}M)$ and its formal adjoint $\delta_f^* : C^{\infty}(S^{p-1}M) \to C^{\infty}(S^p M)$ with respect to the weighted scalar product are given by

$$\delta_f T(X_1, \dots, X_{p-1}) = -\sum_{i=1}^n \nabla_{e_i} T(e_i, X_1, \dots, X_{p-1}) + T(\operatorname{grad} f, X_1, \dots, X_{p-1}),$$

$$\delta_f^* T(X_1, \dots, X_p) = \frac{1}{p} \sum_{i=0}^{p-1} \nabla_{X_{1+i}} T(X_{2+i}, \dots, X_{p+i}),$$

where the sums 1 + i, ..., p + i are taken modulo p. If f is constant, we recover the usual notions of Laplacian and divergence. In this case, we will drop the f in the notation. For $\omega \in \Omega^1(M)$, we have $\delta_f^* \omega = \frac{1}{2} L_{\omega^{\sharp}} g$ where ω^{\sharp} is the sharp of ω . Thus, $\delta_f^*(\Omega^1(M))$ is the tangent space of the manifold $g \cdot \text{Diff}(M) = \{\varphi^* g | \varphi \in \text{Diff}(M)\}$. Throughout, any manifold M will be compact and any metric considered on M will be smooth, unless the contrary is explicitly asserted.

3 The shrinker entropy

Let g be a Riemannian metric, $f \in C^{\infty}(M)$, $\tau > 0$ and define

$$\mathcal{W}(g, f, \tau) = \frac{1}{(4\pi\tau)^{n/2}} \int_{M} [\tau(|\nabla f|_g^2 + \mathrm{scal}_g) + f - n] e^{-f} \, dV.$$

Let

$$\mu(g,\tau) = \inf \left\{ \mathcal{W}(g,f,\tau) \middle| f \in C^{\infty}(M), \frac{1}{(4\pi\tau)^{n/2}} \int_{M} e^{-f} dV_g = 1 \right\}.$$

For any fixed $\tau > 0$, the infimum is finite and is realized by a smooth function [10, Lemma 6.23 and 6.24]. We define the shrinker entropy as

$$\nu(g) = \inf \{ \mu(g, \tau) \mid \tau > 0 \}$$

Recall also the definition of Perelman's λ -functional

$$\lambda(g) = \inf_{\substack{f \in C^{\infty}(M) \\ \int \\ g = f \ dV_g = 1M}} \int (\operatorname{scal}_g + |\nabla f|_g^2) e^{-f} \ dV_g.$$
(3.1)

If $\lambda(g) > 0$, then $\nu(g)$ is finite and realized by some $\tau_g > 0$ (see [10, Corollary 6.34]). In this case, a pair (f_g, τ_g) realizing $\nu(g)$ satisfies the equations

$$\tau (2\Delta f + |\nabla f|^2 - \text{scal}) - f + n + \nu = 0, \qquad (3.2)$$

$$\frac{1}{(4\pi\tau)^{n/2}} \int_{M} f e^{-f} \, dV = \frac{n}{2} + \nu, \tag{3.3}$$

see e.g. [8, p. 5]. On the other hand, $\nu(g)$ is not finite if $\lambda(g) < 0$ [10, p. 244].

Remark 3.1 Note that \mathcal{W} possesses the symmetries $\mathcal{W}(\varphi^* g, \varphi^* f, \tau) = \mathcal{W}(g, f, \tau)$ for $\varphi \in \text{Diff}(M)$ and $\mathcal{W}(\alpha g, f, \alpha \tau) = \mathcal{W}(g, f, \tau)$ for $\alpha > 0$. Therefore, $\nu(g) = \nu(\alpha \cdot \varphi^* g)$ for any $\varphi \in \text{Diff}(M)$ and $\alpha > 0$.

Remark 3.2 The shrinker entropy is upper semicontinuous with respect to the C^2 -topology if defined. Let g be fixed and (f_g, τ_g) be a minimizing pair. Let $g_i \to g$ in C^2 and $v_i \in C^{\infty}(M)$ such that $e^{-v_i} dV_{g_i} = e^{-f_g} dV_g$. Then,

$$\nu(g) = \mathcal{W}(g, f_g, \tau_g) = \lim_{i \to \infty} \mathcal{W}(g_i, v_i, \tau_g) \ge \limsup_{i \to \infty} \nu(g_i)$$

Proposition 3.3 (First variation of v) Let (M, g) be a Riemannian manifold. Then the first variation of v is given by

$$\nu(g)'(h) = -\frac{1}{(4\pi\tau_g)^{n/2}} \int_M \left\langle \tau_g(\text{Ric} + \nabla^2 f_g) - \frac{1}{2}g, h \right\rangle e^{-f_g} \, dV_g.$$

where (f_g, τ_g) realizes v(g).

Proof See e.g. [8, Lemma 2.2].

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By scale and diffeomorphism invariance, it is easy to see that ν is nondecreasing under the Ricci flow and that ν stays constant if and only if

$$\operatorname{Ric} + \nabla^2 f_g = \frac{1}{2\tau_g} g, \qquad (3.4)$$

i.e. if g is a critical point of v. Observe that in this case, the pair (f_g, τ_g) realizing v(g) is unique.

Remark 3.4 A metric g is a critical point of v if and only if g is a gradient shrinking Ricci soliton. Suppose that (3.4) holds, then g is a gradient shrinking Ricci soliton since $\nabla^2 f_g = \frac{1}{2}L_{\text{grad}}f_g g$ and $\frac{1}{2\tau_g} > 0$. Conversely, any gradient shrinking Ricci soliton g has positive scalar curvature [28, Proposition 1]. Therefore, v(g) is finite, since $\lambda(g) \ge \min \text{scal}_g > 0$. Since v stays constant along the flow starting at g, we necessarily have (3.4).

Proposition 3.5 (Second variation of v) Let (M, g) be a gradient shrinking Ricci soliton. Then the second variation of v at g is given by

$$\nu_g''(h) = \frac{\tau}{(4\pi\tau)^{n/2}} \int_M \langle Nh, h \rangle e^{-f} \, dV \tag{3.5}$$

where (f, τ) is the minimizing pair realizing v. The stability operator N is given by

$$Nh = -\frac{1}{2}\Delta_f h + \mathring{R}h + \delta_f^*(\delta_f(h)) + \frac{1}{2}\nabla^2 v_h - \operatorname{Ric}\frac{\int_M \langle \operatorname{Ric}, h \rangle e^{-f} \, dV}{\int_M \operatorname{scal} e^{-f} \, dV}.$$
(3.6)

Here, $\mathring{R}h(X, Y) = \sum_{i=1}^{n} h(R_{e_i, X}Y, e_i)$ and v_h is the unique solution of

$$\left(-\Delta_f + \frac{1}{2\tau}\right)v_h = \delta_f(\delta_f(h)). \tag{3.7}$$

Proof See [8, Theorem 1.1].

Remark 3.6 In the following, we explain how this variational formula can be substantially simplified. The operator N is formally self-adjoint with respect to the weighted measure because

$$\int_{M} \langle \nabla^2 v_h, k \rangle e^{-f} \, dV = \int_{M} \langle v_h, \delta_f(\delta_f(k)) \rangle e^{-f} \, dV = \int_{M} \langle v_h, (-\Delta_f + \frac{1}{2\tau}) v_k \rangle e^{-f} \, dV$$

and the right hand side is symmetric in h and k. The other summands of N are clearly self-adjoint. Therefore,

$$\left. \frac{d^2}{dsdt} \right|_{s,t=0} \nu(g+th+sk) = \frac{\tau}{(4\pi\tau)^{n/2}} \int_M \langle Nh,k \rangle e^{-f} \, dV.$$

By scale and diffeomorphism invariance of ν , we have, for an arbitrary $k \in C^{\infty}(S^2M)$,

$$\frac{\tau}{(4\pi\tau)^{n/2}} \int_{M} \langle Nh, k \rangle e^{-f} \, dV = 0, \quad \text{if } h \in \mathbb{R} \cdot g \oplus \delta_f^*(\Omega^1(M)) = \mathbb{R} \cdot \text{Ric} \oplus \delta_f^*(\Omega^1(M)).$$

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The equality of the direct sums holds because g is a Ricci soliton. Thus ν'' is only nontrivial on the orthogonal component of Ric $\cdot g \oplus \delta_f^*(\Omega^1(M))$, given by

$$V = \left\{ h \in \Gamma(S^2 M) \mid \delta_f(h) = 0 \text{ and } \int_M \langle \operatorname{Ric}, h \rangle e^{-f} \, dV = 0 \right\},$$

and the stability operator $N|_V : V \to V$ is of the form

$$N = -\frac{1}{2}\Delta_f + \mathring{R}.$$

We also have $\delta_f \operatorname{Ric} = 0$ [8, Lemma 3.1] and thus, $\delta_f^{-1}(0) = \mathbb{R} \cdot \operatorname{Ric} \oplus V$.

Definition 3.7 Let (M, g) be a Ricci soliton. We call the soliton linearly stable if all eigenvalues of N are nonpositive and linearly unstable otherwise. If N is nonpositive and ker $N \cap V \neq \{0\}$, we call the soliton neutrally linearly stable.

Remark 3.8 The round sphere is linearly stable; the complex projective space with the Fubini-Study metric is neutrally linearly stable. Any product of positive Einstein manifolds is unstable. Some more examples are discussed in [6,7,24]. It is conjectured that all compact 4-dimensional nontrivial Ricci solitons are linearly unstable [5, p. 29]. Due to a result by Hall and Murphy [25], any kählerian shrinking Ricci soliton is linearly unstable if dim $H^{1,1}(M) > 1$. This results applies to all known nontrivial Ricci solitons in any dimension.

4 Analyticity and a Lojasiewicz–Simon inequality

A necessary tool in proving our stability and instability theorems is a Lojasiewicz–Simon inequality for ν which will be the main theorem of this section. To prove this, ν needs to be analytic. This is the content of the following

Lemma 4.1 Let (M, g_0) be a gradient shrinking Ricci soliton. Then there exists a $C^{2,\alpha}$ neighbourhood \mathcal{U} of g_0 in the space of metrics such that the minimizing pair (f_g, τ_g) realizing $\nu(g)$ is unique and depends analytically on the metric. Moreover, the map $g \mapsto \nu(g)$ is analytic on \mathcal{U} .

Remark 4.2 Previously, analytic dependence of $\mu(g, 1/2)$ on g in the neighbourhood of a Ricci soliton was shown in [42, Lemma 2.2]. The proof presented here turns out to be similar but some subtleties occur due to the presence of the scale parameter τ .

Proof of Lemma 4.1 We use the implicit function theorem for Banach manifolds. We define a map $H(g, f, \tau) = \tau (2\Delta f + |\nabla f|^2 - \text{scal}) - f + n$. Let

$$C_{g_0}^{k,\alpha}(M) = \left\{ u \in C^{k,\alpha}(M) \mid \int_{M} u e^{-f_{g_0}} dV_{g_0} = 0 \right\}$$

and let $\mathcal{M}^{2,\alpha}$ be the Banach manifold of $C^{2,\alpha}$ -metrics. Define

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$$L: \mathcal{M}^{2,\alpha} \times C^{2,\alpha}(M) \times \mathbb{R}_+ \to C^{0,\alpha}_{g_0}(M) \times \mathbb{R} \times \mathbb{R},$$
$$(g, f, \tau) \mapsto (L_1, L_2, L_3),$$

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where the three components are given by

$$\begin{split} L_1(g, f, \tau) &= H(g, f, \tau) - \frac{1}{(4\pi\tau_{g_0})^{n/2}} \int_M H(g, f, \tau) e^{-f_{g_0}} \, dV_{g_0}, \\ L_2(g, f, \tau) &= \frac{1}{(4\pi\tau)^{n/2}} \int_M f e^{-f} \, dV_g - \frac{n}{2} + \frac{1}{(4\pi\tau)^{n/2}} \int_M H(g, f, \tau) e^{-f_g} \, dV_g, \\ L_3(g, f, \tau) &= \frac{1}{(4\pi\tau)^{n/2}} \int_M e^{-f} \, dV_g - 1. \end{split}$$

This is an analytic map between Banach manifolds. We have $L(g, f, \tau) = (0, 0, 0)$ if and only if there exists a constant $c \in \mathbb{R}$ such that the set of equations

$$\tau (2\Delta f + |\nabla f|^2 - \operatorname{scal}) - f + n = c,$$
(4.1)

$$\frac{1}{(4\pi\tau)^{n/2}} \int_{M} f e^{-f} \, dV - \frac{n}{2} = -c, \tag{4.2}$$

$$\frac{1}{(4\pi\tau)^{n/2}} \int_{M} e^{-f} \, dV = 1 \tag{4.3}$$

is satisfied. Now we compute the differential of L at $(g_0, f_{g_0}, \tau_{g_0})$ restricted to $R = C^{2,\alpha}(M) \times \mathbb{R}$. Let $F_{g_0} = f_{g_0} - \frac{n}{2} - \nu(g_0)$ and

$$V = \left\{ u \in C_{g_0}^{2,\alpha}(M) \mid \int_{M} u \cdot F_{g_0} e^{-f_{g_0}} dV_{g_0} = 0 \right\},$$
$$W = \left\{ u \in C_{g_0}^{0,\alpha}(M) \mid \int_{M} u \cdot F_{g_0} e^{-f_{g_0}} dV_{g_0} = 0 \right\}.$$

By (3.2) and smoothness of f_{g_0} , we have $F_{g_0} \in C_{g_0}^{k,\alpha}(M)$ for all $k \ge 0$. Moreover, we have the decompositions

$$C^{2,\alpha}(M) \cong V \oplus \operatorname{span}(F_{g_0}) \oplus \mathbb{R},$$
$$C^{0,\alpha}_{g_0}(M) \cong W \oplus \operatorname{span}(F_{g_0}),$$

where the last factor in the first decomposition represents the constant functions. Consider the differential of L restricted to R as a linear map

$$dL_{(g_0,f_{g_0},\tau_{g_0})}\Big|_R: V \oplus \operatorname{span}(F_{g_0}) \oplus \mathbb{R} \oplus \mathbb{R} \to W \oplus \operatorname{span}(F_{g_0}) \oplus \mathbb{R} \oplus \mathbb{R}.$$

Straightforward calculations, using the Euler–Lagrange equations (3.2) and (3.2), show that it is equal to

$$dL_{(g_0, f_{g_0}, \tau_{g_0})}\Big|_R = \begin{pmatrix} 2\tau_{g_0}\Delta_{f_{g_0}} - 1 & 0 & 0 & 0\\ 0 & 1 & 0 & \frac{1}{\tau_{g_0}} \\ 0 & -\|F_{g_0}\|_{L^2(v\,dV)}^2 & -\frac{n}{2} - \nu(g_0) & -\frac{n}{2\tau_{g_0}}\left(\frac{n}{2} + \nu(g_0) + 1\right) \\ 0 & 0 & -1 & -\frac{n}{2\tau_{g_0}} \end{pmatrix},$$

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where $L^2(v \, dV)$ is the L^2 norm with respect to the *v*-weighted measure, with $v = \frac{1}{(4\pi \tau_{g_0})^{n/2}} e^{-f_{g_0}}$. In the following, we show that this map is an isomorphism. From (3.2), we conclude that $\Delta_{f_{g_0}} F_{g_0} = \frac{1}{\tau_{g_0}} F_{g_0}$. Thus, the map

$$2\tau_{g_0}\Delta_{f_{g_0}}-1:V\to W$$

is well defined. By [18, Proposition 3.1], the smallest nonzero eigenvalue of $\Delta_{f_{g_0}}$ satisfies $\lambda > \frac{1}{2\tau_{g_0}}$. This implies invertibility of the above map. It remains to consider the lower right 3×3 -block which we denote by A. We have

$$\det(A) = \frac{1}{\tau_{g_0}} \left(\left\| F_{g_0} \right\|_{L^2(v \, dV)}^2 - \frac{n}{2} \right)$$

Since F_{g_0} is an eigenfunction of the weighted Laplacian to the eigenvalue $\frac{1}{\tau_{g_0}}$,

$$\|F_{g_0}\|_{L^2(v\,dV)}^2 = \tau_{g_0}(\Delta_f F_{g_0}, F_{g_0})_{L^2(v\,dV)} = \tau_{g_0} \|\nabla F_{g_0}\|_{L^2(v\,dV)}^2 = \tau_{g_0} \|\nabla f_{g_0}\|_{L^2(v\,dV)}^2.$$

By definition of $v(g_0)$,

$$\begin{aligned} \tau_{g_0} \left\| \nabla f_{g_0} \right\|_{L^2(v \, dV)}^2 &= \frac{1}{(4\pi \, \tau_{g_0})^{n/2}} \int_M \tau_{g_0} |\nabla f_{g_0}|^2 e^{-f_{g_0}} \, dV \\ &= \nu(g_0) - \frac{1}{(4\pi \, \tau_{g_0})^{n/2}} \int_M [\tau_{g_0} \operatorname{scal} + f_{g_0} - n] e^{-f_{g_0}} \, dV \\ &= \frac{n}{2} - \frac{1}{(4\pi \, \tau_{g_0})^{n/2}} \int_M \tau_{g_0} \operatorname{scal} e^{-f_{g_0}} \, dV. \end{aligned}$$

Therefore,

$$\det(A) = -\frac{1}{(4\pi \tau_{g_0})^{n/2}} \int_M \operatorname{scal} e^{-f_{g_0}} dV < 0$$

because the scalar curvature of g_0 is positive [28, Proposition 1]. In summary, we have shown that $dL_{(g_0, f_{g_0}, \tau_{g_0})}|_R$ is invertible. By the implicit function theorem for Banach manifolds, there exists a neighbourhood $\mathcal{U} \subset \mathcal{M}^{2,\alpha}$ of g_0 and an analytic map $P: \mathcal{U} \to C^{2,\alpha}(M) \times \mathbb{R}_+$ such that L(g, P(g)) = 0. Moreover, there exists a neighbourhood $\mathcal{V} \subset C^{2,\alpha}(M) \times \mathbb{R}_+$ such that for any $(g, f, \tau) \in \mathcal{U} \times \mathcal{V}$, we have $L(g, f, \tau) = 0$ if and only if $(f, \tau) = P(g)$.

Now we claim that on a smaller neighbourhood $\mathcal{U}_1 \subset \mathcal{U}$, there is a unique pair of minimizers in the definition of ν and it is equal to P(g). Suppose this is not the case. Then there exist a sequence g_i of metrics such that $g_i \to g_0$ in $C^{2,\alpha}$ and pairs of minimizers (f_{g_i}, τ_{g_i}) such that $P(g_i) \neq (f_i, \tau_{g_i})$ for all $i \in \mathbb{N}$. By substituting $w_{g_i}^2 = e^{-f_{g_i}}$, we see that the pair (w_{g_i}, τ_{g_i}) is a minimizer of the functional

$$\tilde{\mathcal{W}}(g_i, w, \tau) = \frac{1}{(4\pi\tau)^{n/2}} \int_M [\tau(4|\nabla w|^2 + \operatorname{scal}_g w^2) - \log(w^2)w^2 - nw^2] \, dV_g$$

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under the constraint $\frac{1}{(4\pi\tau)^{n/2}}\int_M w^2 dV_{g_i} = 1$. It satisfies the pair of equations

$$-\tau_{g_i}(4\Delta w_{g_i} + \operatorname{scal}_{g_i} w_{g_i}) - 2\log(w_{g_i})w_{g_i} + nw_{g_i} + \nu(g_i)w_{g_i} = 0,$$
(4.4)

$$\frac{1}{(4\pi\,\tau_{g_i})^{n/2}}\int\limits_M w_{g_i}^2\log w_{g_i}^2\,dV_{g_i} = \frac{n}{2} + \nu(g_i). \tag{4.5}$$

By upper semicontinuity, $\nu(g_i) \leq C_1$.

Now we show that there exist constants C_2 , $C_3 > 0$ such that $C_2 \le \tau_{g_i} \le C_3$. Suppose this is not the case. By [10, Lemma 6.30], we have a lower estimate

$$\nu(g_i) = \mu(g_i, \tau_{g_i}) \ge (\tau_{g_i} - 1)\lambda(g_i) - \frac{n}{2}\log\tau_{g_i} - C_4(g_i) \ge (\tau_{g_i} - 1)C_5 - \frac{n}{2}\log\tau_{g_i} - C_6.$$

We have $C_5 > 0$ because $\lambda(g)$ is uniformly bounded on \mathcal{U} and positive. The constant $C_4(g)$ depends on the Sobolev constant and the volume. Now if τ_{g_i} converges to 0 or ∞ , $\nu(g_i)$ diverges, which causes the contradiction. Observe that we also obtained a lower bound on $\nu(g_i)$.

Next, we show that $\|\nabla w_{g_i}\|_{L^2}$ is bounded. Choose $\epsilon > 0$ so small that $2 + 2\epsilon \le \frac{2n}{n-2}$. By Jensen's inequality and the bounds on τ_{g_i} ,

$$\begin{split} \int_{M} w_{g_{i}}^{2} \log w_{g_{i}}^{2} \, dV_{g_{i}} &= \frac{1}{\epsilon} \int_{M} w_{g_{i}}^{2} \log w_{g_{i}}^{2\epsilon} \, dV_{g_{i}} \\ &\leq \frac{1}{\epsilon} \left\| w_{g_{i}} \right\|_{L^{2}}^{2} \log \left(\frac{1}{\left\| w_{g_{i}} \right\|_{L^{2}}^{2}} \int_{M} w_{g_{i}}^{2+2\epsilon} \, dV_{g_{i}} \right) \\ &= \frac{1}{\epsilon} (4\pi \, \tau_{g_{i}})^{n/2} \log \left((4\pi \, \tau_{g_{i}})^{-n/2} \int_{M} w_{g_{i}}^{2+2\epsilon} \, dV_{g_{i}} \right) \\ &\leq C_{7} \log \left(\int_{M} w_{g_{i}}^{2+2\epsilon} \, dV_{g_{i}} \right) + C_{8}. \end{split}$$

By the Sobolev inequality,

$$\int_{M} w_{g_{i}}^{2+2\epsilon} dV_{g_{i}} \leq C_{9} (\left\|\nabla w_{g_{i}}\right\|_{L^{2}}^{2} + \left\|w_{g_{i}}\right\|_{L^{2}}^{2})^{1+\epsilon} \leq C_{9} (\left\|\nabla w_{g_{i}}\right\|_{L^{2}}^{2} + C_{10})^{1+\epsilon}.$$

In summary, we have

$$C_{1} \geq \frac{1}{(4\pi\tau)^{n/2}} \int_{M} [\tau(4|\nabla w_{g_{i}}|^{2} + \operatorname{scal}_{g_{i}}w_{g_{i}}^{2}) - \log(w_{g_{i}}^{2})w_{g_{i}}^{2} - nw^{2}] dV_{g_{i}}$$

$$\geq C_{11} \|\nabla w_{g_{i}}\|_{L^{2}}^{2} - C_{12}\log(\|\nabla w_{g_{i}}\|_{L^{2}}^{2} + C_{10}) - C_{13},$$

which shows that $\|\nabla w_{g_i}\|_{L^2}$ is bounded.

Now we continue with a bootstrap argument. By Sobolev embedding, the bound on $\|w_{g_i}\|_{H^1}$ implies a bound on $\|w_{g_i}\|_{L^{2n/(n-2)}}$. Let p = 2n/(n-2) and choose some q slightly smaller than p. By elliptic regularity and (4.4),

$$\left\|w_{g_{i}}\right\|_{W^{2,q}} \leq C_{14}\left(\left\|w_{g_{i}}\log w_{g_{i}}\right\|_{L^{q}} + \left\|w_{g_{i}}\right\|_{L^{q}}\right).$$

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Since for any $\beta > 1$, $|x \log x| \le |x|^{\beta}$ for |x| large enough, we have

$$\|w_{g_i} \log w_{g_i}\|_{L^q} \le C_{15}(\operatorname{vol}(M, g_i)) + \|w_{g_i}\|_{L^p} \le C_{16} + \|w_{g_i}\|_{L^p}.$$

Thus, $||w_{g_i}||_{W^{2,q}} \leq C(q)$. Using Sobolev embedding, we obtain bounds on $||w_{g_i}||_{L^{p'}}$ for some p' > p. From (4.4) again, we have bounds on $||w_{g_i}||_{W^{2,q'}}$ for any q' < p'. Using these arguments repetitively, we obtain $||w_{g_i}||_{W^{2,q}} \leq C(q)$ for all $q \in (1, \infty)$. Again by elliptic regularity,

$$\|w_{g_i}\|_{C^{2,\alpha}} \le C_{17} \left(\|w_{g_i} \log w_{g_i}\|_{C^{0,\alpha}} + \|w_{g_i}\|_{C^{0,\alpha}} \right) \le C_{18} \left(\left(\|w_{g_i}\|_{C^{0,\alpha}} \right)^{\gamma} + \|w_{g_i}\|_{C^{0,\alpha}} \right)$$

for some $\gamma > 1$. For some sufficiently large q, we have, by Sobolev embedding,

$$\|w_{g_i}\|_{C^{0,\alpha}} \le C_{19} \|w_{g_i}\|_{W^{1,q}} \le C_{19} \cdot C(q).$$

We finally obtained an upper bound on $\|w_{g_i}\|_{C^{2,\alpha}}$. Thus there exists a subsequence, again denoted by (w_{g_i}, τ_{g_i}) , which converges in $C^{2,\alpha'}, \alpha' < \alpha$, to some limit $(w_{\infty}, \tau_{\infty})$. By Remark 3.2,

$$\nu(g_0) \ge \lim_{i \to \infty} \nu(g_i) = \lim_{i \to \infty} \tilde{\mathcal{W}}(g_i, w_{g_i}, \tau_{g_i}) = \tilde{\mathcal{W}}(g_0, w_{\infty}, \tau_{\infty}) \ge \nu(g_0),$$

and therefore, $(w_{\infty}, \tau_{\infty}) = (w_{g_0}, \tau_{g_0})$ because the minimizing pair is unique at g_0 . Moreover, by resubstituting,

$$(f_{g_i}, \tau_{g_i}) \to (f_{\infty}, \tau_{\infty}) = (f_{g_0}, \tau_{g_0})$$

in $C^{2,\alpha'}$. Because the pair (f_{g_i}, τ_{g_i}) satisfies (3.2) and (3.2), $L(g_i, f_{g_i}, \tau_{g_i}) = 0$ and the implicit function argument from above implies that $P(g_i) = (f_{g_i}, \tau_{g_i})$ for large *i*. This proves the claim by contradiction. We also have shown that the constant *c* appearing above is equal to $-\nu(g)$. Since the map $g \mapsto (f_g, \tau_g)$ is analytic, the map

$$g \mapsto v(g) = -\tau_g (2\Delta f_g + |\nabla f_g|^2 - \operatorname{scal}_g) + f_g - n$$

is also analytic. This proves the lemma.

Lemma 4.3 Let (M, g_0) be a gradient shrinking Ricci soliton. Then there exists a $C^{2,\alpha}$ -neighbourhood U of g_0 in the space of metrics and a constant C > 0 such that

$$\left\| \frac{d}{dt} \bigg|_{t=0} f_{g+th} \right\|_{C^{2,\alpha}} \leq C \|h\|_{C^{2,\alpha}}, \quad \left\| \frac{d}{dt} \bigg|_{t=0} f_{g+th} \right\|_{H^{i}} \leq C \|h\|_{H^{i}}, \quad i = 0, 1, 2,$$
$$\left| \frac{d}{dt} \bigg|_{t=0} \tau_{g+th} \right| \leq C \|h\|_{L^{2}}$$

for all $g \in \mathcal{U}$.

Proof Let the maps P, L and the space R be as in the proof of the previous lemma. Note that if U is small enough, the map

$$h \mapsto \left(\frac{d}{dt}\bigg|_{t=0} f_{g+th}, \frac{d}{dt}\bigg|_{t=0} \tau_{g+th}\right)$$

is precisely the differential of P. By the chain rule,

$$dP_g = -\left(dL_{g,f_g,\tau_g}|_R\right)^{-1} \circ dL_{g,f_g,\tau_g}|_{C^{2,\alpha}(S^2M)}.$$
(4.6)

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First, we compute $dL_{g,f,\tau}|_{C^{2,\alpha}(S^2M)}$. We have

$$dH(g,h,f,\tau) := \left. \frac{d}{dt} \right|_{t=0} H(g+th,f,\tau) = 2\tau \Delta' f - \tau h(\operatorname{grad} f,\operatorname{grad} f) - \operatorname{scal}'$$

Then,

$$dL_{1}(h) = dH(g, h, f, \tau) - \frac{1}{(4\pi\tau_{g_{0}})^{n/2}} \int_{M} dH(g, h, f, \tau) e^{-f_{g_{0}}} dV_{g_{0}},$$

$$dL_{2}(h) = \frac{1}{(4\pi\tau)^{n/2}} \frac{1}{2} \int_{M} f e^{-f} \operatorname{tr} h \, dV_{g} + \frac{1}{(4\pi\tau_{g_{0}})^{n/2}} \int_{M} dH(g, h, f, \tau) e^{-f_{g}} \, dV_{g_{0}},$$

$$dL_{3}(h) = \frac{1}{(4\pi\tau)^{n/2}} \frac{1}{2} \int_{M} e^{-f} \operatorname{tr} h \, dV_{g}$$

by the first variation of the volume element. The first variation of the Laplacian and the scalar curvature are

$$\frac{d}{dt}\Big|_{t=0} \Delta_{g+th} f = \langle h, \nabla^2 f \rangle - \left\langle \delta h + \frac{1}{2} \nabla \mathrm{tr} h, \nabla f \right\rangle,$$
$$\frac{d}{dt}\Big|_{t=0} \mathrm{scal}_{g+th} = \Delta_g(\mathrm{tr}_g h) + \delta_g(\delta_g h) - \langle \mathrm{Ric}_g, h \rangle_g,$$

c.f. [2, pp. 62-64]. Therefore, we have the estimates

$$\| dL_{g,f,\tau} \|_{C^{2,\alpha}(S^2M)}(h) \|_{C^{0,\alpha}} \le C_1 \| h \|_{C^{2,\alpha}}, \| dL_{g,f,\tau} \|_{C^{2,\alpha}(S^2M)}(h) \|_{H^{i-2}} \le C_2 \| h \|_{H^i},$$

where i = 0, 1, 2. Now we consider $dL_{g,f,\tau}|_R$. This is essentially an elliptic operator which is invertible at $(g_0, f_{g_0}, \tau_{g_0})$. By continuity, it is also invertible on a small $C^{2,\alpha}$ -neighbourhood of the tuple. By elliptic estimates, we conclude from (4.6) that

$$\|dP_g(h)\|_{C^{2,\alpha}} \le C_3 \|h\|_{C^{2,\alpha}}, \qquad \|dP_g(h)\|_{H^i} \le C_4 \|h\|_{H^i}$$

for g close to g_0 . This finishes the proof of the lemma.

To prove the main result of this section, we need a form of the slice theorem. Recall that Ebin's slice theorem [15] states the following: For any Riemannian metric g on a compact manifold, there exists neighbourhood \mathcal{U} of g and a submanifold $S_g \subset \mathcal{M}$ tangent to $\delta_g^{-1}(0)$ such that any metric in \mathcal{U} is isometric to a unique metric in S_g . We call S_g a slice of the action of Diff(M).

For our purpose, it is more convenient to use a slice tangent to $\delta_{g,f}^{-1}(0)$ where f is the potential function of the Ricci soliton. In fact, we can choose the slice to be affine, i.e.

$$S_{g,f} = \left\{ g + h \mid \delta_{g,f}(h) = 0 \right\}.$$
(4.7)

By [1, Lemma A.5], there exists a $C^{1+k,\alpha}$ -neighbourhood \mathcal{U} of g such that any metric in \mathcal{U} is isometric to a unique metric in $\mathcal{S}_{g,f} \cap \mathcal{U}$.

Theorem 4.4 (Lojasiewicz–Simon inequality) Let (M, g_0) be a gradient shrinking Ricci soliton. Then there exists a $C^{2,\alpha}$ neighbourhood \mathcal{U} of g_0 and constants $\sigma \in [1/2, 1), C > 0$ such that

$$|\nu(g) - \nu(g_0)|^{\sigma} \le C \left\| \tau(\operatorname{Ric}_g + \nabla^2 f_g) - \frac{1}{2}g \right\|_{L^2}$$
 (4.8)

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for all $g \in \mathcal{U}$.

Proof Since both sides are diffeomorphism invariant, it suffices to show the inequality on a slice to the action of the diffeomorphism group. Let

$$S_{g_0, f_{g_0}} = \mathcal{U} \cap \left\{ g_0 + h \mid \delta_{g_0, f_{g_0}}(h) = 0 \right\}.$$

Let $\tilde{\nu}$ be the ν -functional restricted to $S_{g_0, f_{g_0}}$. Obviously, $\tilde{\nu}$ is analytic since ν is. By the first variational formula in Lemma 3.3, the L^2 -gradient of ν is (up to a constant factor) given by $\nabla \nu(g) = [\tau (\operatorname{Ric}_g + \nabla^2 f_g) - \frac{1}{2}g]e^{-f_g}$. It vanishes at g_0 . On the neighbourhood \mathcal{U} , we have the uniform estimate

$$\|\nabla \nu(g_1) - \nabla \nu(g_2)\|_{L^2} \le C \|g_1 - g_2\|_{H^2}, \qquad (4.9)$$

which holds by Taylor expansion and Lemma 4.3. The L^2 -gradient of $\tilde{\nu}_-$ is given by the projection of $\nabla \nu$ to $\delta_{g_0, f_{g_0}}^{-1}(0)$. Therefore, (4.9) also holds for $\nabla \tilde{\nu}$. The linearization of $\tilde{\nu}$ at g_0 vanishes on $\mathbb{R} \cdot \operatorname{Ric}_{g_0}$ and equals

$$\frac{\tau_{g_0}}{(4\pi\,\tau_{g_0})^{n/2}}e^{-f_0}\left(-\frac{1}{2}\Delta_f + \mathring{R}\right)$$

on the $L^2(e^{-f_{g_0}} dV)$ -orthogonal complement of $\mathbb{R} \cdot \operatorname{Ric}_{g_0}$ in $\delta_{g_0, f_{g_0}}^{-1}(0)$, see Remark 3.6. Let us denote this operator by D. By ellipticity,

$$D: (\delta_{g_0, f_{g_0}}^{-1}(0))^{2, \alpha} \to (\delta_{g_0, f_{g_0}}^{-1}(0))^{0, \alpha}$$

is Fredholm. It also satisfies the estimate $||Dh||_{L^2} \leq C ||h||_{H^2}$. By a general Lojasiewicz– Simon inequality [11, Theorem 7.3], there exists a constant $\sigma \in [1/2, 1)$ such that the inequality $|\nu(g) - \nu(g_0)|^{\sigma} \leq ||\nabla \tilde{\nu}_{-}(g)||_{L^2}$ holds for any $g \in S_{g_0, f_{g_0}}$. Since

$$\|\nabla \tilde{\nu}(g)\|_{L^2} \le \|\nabla \nu(g)\|_{L^2} \le C \left\|\tau(\operatorname{Ric}_g + \nabla^2 f_g) - \frac{1}{2}g\right\|_{L^2}$$

(4.8) holds on all $g \in S_{g_0, f_{g_0}}$. By diffeomorphism invariance, it holds on all $g \in \mathcal{U}$.

5 Dynamical stability and instability

In this section, we prove the main theorems of the paper. We consider the τ -flow

$$\dot{g}(t) = -2\operatorname{Ric}_{g(t)} + \frac{1}{\tau_{g(t)}}g(t)$$
 (5.1)

which is well-defined in a neighbourhood of a gradient shrinking Ricci soliton. Observe that ν is nondecreasing under the τ -flow. We also construct a modified τ -flow as follows: Let $\varphi_t \in \text{Diff}(M), t \ge 1$ be the family of diffeomorphisms generated by $X(t) = -\text{grad}_{g(t)} f_{g(t)}$ and define $\tilde{g}(t) := \varphi_t^* g(t)$, where g(t) is a solution of (5.1). Then we have

$$\frac{d}{dt}\tilde{g}(t) = -2(\operatorname{Ric}_{\tilde{g}(t)} + \nabla^2 f_{\tilde{g}(t)}) + \frac{1}{\tau_{\tilde{g}(t)}}\tilde{g}(t).$$
(5.2)

This is the gradient flow of τ with respect to the weighted L^2 -measure.

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Lemma 5.1 Let (M, g) be a gradient shrinking Ricci soliton and let $k \ge 3$. Then for each C^k -neighbourhood \mathcal{U} of g, there exists a C^{k+5} -neighbourhood \mathcal{V} of g such that the modified τ -flow (5.2), starting at any metric in \mathcal{V} , stays in \mathcal{U} for all $t \in [0, 1]$.

Proof Let us denote the ϵ -ball with respect to the C_g^k -norm by \mathcal{B}_{ϵ}^k . Without loss of generality, we may assume that \mathcal{U} is of the form $\mathcal{U} = \mathcal{B}_{\epsilon}^k$ for some $\epsilon > 0$. Throughout the proof, let us assume that we are in a neighbourhood of g such that Lemmas 4.1 and 4.3 hold. All covariant derivatives, Laplacians and norms in this proof are taken with respect to g(t) (resp. $\tilde{g}(t)$). Along the (unmodified) τ -flow, we have the evolution equations

$$\partial_t R = -\Delta R + R * R + \frac{2}{\tau}R,$$

$$\partial_t \operatorname{Ric} = -\Delta \operatorname{Ric} + R * \operatorname{Ric},$$

$$\partial_t \frac{1}{2\tau}g = -\frac{\partial_t \tau}{2\tau^2}g + \frac{1}{2\tau}\left(-2\operatorname{Ric} + \frac{1}{\tau}g\right)$$

$$\partial_t \nabla^2 f = \nabla \operatorname{Ric} * \nabla f + \nabla^2 \partial_t f,$$

where * is Hamilton's notation for a combination of tensor products and contractions. The first two formulas follow from rescaling the evolution equations for the standard Ricci flow [3, pp. 26–28]. The third formula is clear and the last one follows from the first variation of the Hessian [36, Lemma A.2]. The evolution equation for the Riemann tensor yields the evolution inequality

$$\partial_t |\nabla^i R|^2 \le -\Delta |\nabla^i R|^2 + \sum_{j=1}^{i-1} C_{ij} |\nabla^j R| |\nabla^{i-j} R| |\nabla^i R| + C_{i0} \left(|R| + \frac{1}{\tau} \right) |\nabla^i R|^2.$$

From the maximum principle, we obtain the following. Suppose we have a τ -flow g(t) defined on $[0, T], T \le 1$ and the bounds

$$\sup_{p \in M} |R_{g(t)}|_{g(t)} \le K, \qquad \frac{1}{\tau_{g(t)}} \le K, \qquad \sup_{p \in M} |\nabla^{i} R_{g(0)}|_{g(0)} \le K$$
(5.3)

for all $t \in [0, T]$ and $i \le k + 3$. Then there exists a constant $\tilde{K}(K, n, k)$ such that

$$\sup_{p \in M} |\nabla^i R_{g(t)}|_{g(t)} \le \tilde{K}$$
(5.4)

for all $t \in [0, T]$ and $i \le k + 3$. Furthermore, we have the evolution inequality

$$\begin{aligned} \partial_t |\operatorname{Ric} + \nabla^2 f - 1/2\tau \cdot g|^2 &\leq -\Delta |\operatorname{Ric} + \nabla^2 f - 1/2\tau \cdot g|^2 \\ &+ 2|\operatorname{Ric} + \nabla^2 f - 1/2\tau \cdot g||(*)|, \end{aligned}$$

where (*) is given by

$$(*) = R * \operatorname{Ric} + \frac{\partial_t \tau}{2\tau^2} g - \frac{1}{2\tau} \left(-2\operatorname{Ric} + \frac{1}{\tau} g \right) + \nabla \operatorname{Ric} * \nabla f + \nabla^2 \partial_t f + \Delta \nabla^2 f.$$

We obtain a bound

$$|(*)| \le C(||f||_{C^4}, ||R||_{C^1}, ||\partial_t f||_{C^2}),$$

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where we used the estimate $|\partial_t \tau| \le C \|\operatorname{Ric} - \frac{1}{2\tau}g\|_{L^2}$. Similarly, we have

$$\begin{aligned} \partial_t |\nabla^i (\operatorname{Ric} + \nabla^2 f - 1/2\tau \cdot g)|^2 &\leq -\Delta |\nabla^i (\operatorname{Ric} + \nabla^2 f - 1/2\tau \cdot g)|^2 \\ &+ 2 |\nabla^i (\operatorname{Ric} + \nabla^2 f - 1/2\tau \cdot g)||(**)|, \end{aligned}$$

where we have the bound

$$|(**)| \le C(\|f\|_{C^{i+4}}, \|R\|_{C^{i+1}}, \|\partial_t f\|_{C^{i+2}}).$$
(5.5)

It remains to control the norms of f and $\partial_t f$. By the differential equation

$$\tau (2\Delta f + |\nabla f|^2 - \text{scal}) - f + n + \nu = 0$$
(5.6)

and elliptic regularity,

$$\|f\|_{C^{i,\alpha}} \le C(\|f\|_{C^{i-1,\alpha}} + \|\operatorname{scal}\|_{C^{i-2,\alpha}} + |\nu + n|).$$
(5.7)

From differentiating (5.6), we obtain

$$\|\partial_t f\|_{C^{i,\alpha}} \le C(\|\operatorname{Ric} - 1/2\tau \cdot g\|_{C^{i,\alpha}}, \|f\|_{C^{i,\alpha}}, \|R\|_{C^{i-2,\alpha}}).$$
(5.8)

Suppose now again that (5.3) (and therefore, also 5.4) holds. Using an iteration argument in (5.7) and descending to C^i -norms we obtain that

$$\|f\|_{C^{i+4}} \le C(\|f\|_{C^{2,\alpha}}, \|R\|_{C^{i+3}}), \qquad \|\partial_t f\|_{C^{i+2}} \le C(\|f\|_{C^{2,\alpha}}, \|R\|_{C^{i+3}}).$$

For all $i \leq k$, we thus have, by (5.4) and (5.5),

$$\partial_t |\nabla^i (\operatorname{Ric} + \nabla^2 f - 1/2\tau \cdot g)|^2 \le -\Delta |\nabla^i (\operatorname{Ric} + \nabla^2 f - 1/2\tau \cdot g)|^2 + C(||f||_{C^{2,\alpha}}, \tilde{K}) |\nabla^i (\operatorname{Ric} + \nabla^2 f - 1/2\tau \cdot g)|$$
(5.9)

along the τ flow. By diffeomorphism invariance, conclusion (5.4) and the inequality (5.9) also hold for the modified τ -flow. Let us denote the modified τ -flow by \tilde{g} . Choose the C^{k+5} -neighbourhood \mathcal{V} so small that we have the bounds (5.3) for some constant K > 0 and any modified τ -flow starting in \mathcal{V} as long as the flow stays in \mathcal{U} . Then (5.4) holds. Let $\epsilon_1 > 0$. By (5.9) and the maximum principle, we can choose $\delta_1 = \delta_1(\epsilon_1, \mathcal{U}, \mathcal{V}) > 0$ so small that if

$$|\nabla^{i}(\operatorname{Ric}_{\tilde{g}(0)} + \nabla^{2} f_{\tilde{g}(0)} - 1/2\tau_{\tilde{g}(0)} \cdot \tilde{g}(0))| \le \delta_{1},$$

then

$$|\nabla^{i}(\operatorname{Ric}_{\tilde{g}(t)} + \nabla^{2} f_{\tilde{g}(t)} - 1/2\tau_{\tilde{g}(t)} \cdot \tilde{g}(t))| \leq \epsilon_{1}$$

for $i \le k$ and $t \in [0, 1]$ as long as the flow stays in \mathcal{U} . Let T > 0 be the maximal time such that the modified τ -flow stays in \mathcal{U} . Suppose that $T \le 1$. By integration,

$$\begin{split} \|\tilde{g}(T) - g\|_{C_g^k} &\leq \|\tilde{g}(0) - g\|_{C_g^k} + \int_0^T \frac{d}{dt} \|\tilde{g}(t) - \tilde{g}(0)\|_{C_g^k} dt \\ &\leq \delta(\mathcal{V}) + C(\mathcal{U}) \int_0^T \left\| \dot{\tilde{g}}(t) \right\|_{C_{\tilde{g}(t)}^k} dt \\ &\leq \delta(\mathcal{V}) + C(\mathcal{U}) \cdot k \cdot \epsilon_1 \leq \frac{\epsilon}{2}, \end{split}$$

provided that we have chosen ϵ_1 and $\delta(\mathcal{V})$ small enough. This contradicts the maximality of T and proves the lemma.

Lemma 5.2 (Shi estimates for the τ -flow) Let g(t), $t \in [0, T]$ be a solution of the τ -flow (5.1) and suppose that

$$\sup_{p \in M} |R_{g(t)}|_{g(t)} + \frac{1}{\tau_{g(t)}} \le T^{-1} \quad \forall t \in [0, T].$$

Then for each $k \ge 1$, there exists a constant C(k) such that

$$\sup_{p \in M} |\nabla^k R_{g(t)}|_{g(t)} \le C(k) \cdot T^{-1} t^{-k/2} \quad \forall t \in (0, T].$$

Proof See [37, Lemma 6.5.6].

Estimates of that type are well-known for the standard Ricci flow [22, Theorem 7.1]. Note that we do not make assumptions on the derivatives of the curvature at t = 0.

Theorem 5.3 (Dynamical stability) Let (M, g) be a gradient shrinking Ricci soliton and let $k \ge 3$. Suppose that g is a local maximizer of v. Then for every C^k -neighbourhood U of g, there exists a C^{k+5} -neighbourhood V such that the following holds:

For any metric $g_0 \in \mathcal{V}$, there exists a 1-parameter family of diffeomorphisms φ_t such that for the τ -flow (5.1) starting at g_0 , the modified flow $\varphi_t^*g(t)$ stays in \mathcal{U} for all time and converges to a gradient shrinking Ricci soliton g_∞ in \mathcal{U} as $t \to \infty$. The convergence is of polynomial rate, i.e. there exist constants $C, \alpha > 0$ such that

$$\left\|\varphi_t^*g(t) - g_\infty\right\|_{C^k} \le C(t+1)^{-\alpha}.$$

Proof Without loss of generality, we may assume that $\mathcal{U} = \mathcal{B}_{\epsilon}^k$ and that $\epsilon > 0$ is so small that Theorem 4.4 holds on \mathcal{U} .

By Lemma 5.1, we can choose a small neighbourhood \mathcal{V} such that the modified τ -flow (5.2), starting at any metric $g \in \mathcal{V}$ stays in $\mathcal{B}_{\epsilon/4}^k$ up to time 1. Let $T \ge 1$ be the maximal time such that any solution of (5.2), starting in \mathcal{V} , stays in \mathcal{U} . By definition of T, we have uniform bounds

$$\sup_{p \in M} |R_{\tilde{g}(t)}|_{\tilde{g}(t)} \le C_1 \quad \forall t \in [0, T),$$
$$|\tau_{\tilde{g}(t)}| \le C_2 \quad \forall t \in [0, T).$$

By Lemma 5.2 and diffeomorphism invariance, we have

$$\sup_{p \in M} |\nabla^l R_{\tilde{g}(t)}|_{\tilde{g}(t)} \le C(l) \quad \forall t \in [1, T).$$
(5.10)

Because $f_{\tilde{g}(t)}$ satisfies the equation $\tau (2\Delta f + |\nabla f|^2 - \text{scal}) - f + n + \nu = 0$, we obtain

$$\sup_{p \in M} |\nabla^l f_{\tilde{g}(t)}|_{\tilde{g}(t)} \le \tilde{C}(l) \quad \forall t \in [1, T).$$
(5.11)

We have

$$\|\tilde{g}(T) - g\|_{C^{k}} \le \|\tilde{g}(1) - g\|_{C^{k}} + \int_{1}^{T} \left\|\dot{\tilde{g}}(t)\right\|_{C^{k}} dt \le \frac{\epsilon}{4} + \int_{1}^{T} \left\|\dot{\tilde{g}}(t)\right\|_{C^{k}} dt.$$

By interpolation (c.f. [20, Corollary 12.7]), (5.10) and (5.11), we have

$$\left\|\dot{\tilde{g}}(t)\right\|_{C^{k}} \le C_{3} \left\|\dot{\tilde{g}}(t)\right\|_{L^{2}}^{1-\eta}$$

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for $t \in [1, T)$ and for η as small as we want. In particular, we can assume that $\theta := 1 - \sigma(1 + \eta) > 0$, where σ is the constant appearing in the Lojasiewicz–Simon inequality 4.4. By the first variation of ν ,

$$\frac{d}{dt}\nu(\tilde{g}(t)) \ge C_4 \left\| \dot{\tilde{g}}(t) \right\|_{L^2}^{1+\eta} \left\| \dot{\tilde{g}}(t) \right\|_{L^2}^{1-\eta}.$$

By Theorem 4.4,

$$\begin{aligned} -\frac{d}{dt} |v(\tilde{g}(t)) - v(g)|^{\theta} &= \theta |v(\tilde{g}(t)) - v(g)|^{\theta - 1} \frac{d}{dt} v(\tilde{g}(t)) \\ &\geq C_5 |v(\tilde{g}(t)) - v(g)|^{-\sigma(1+\eta)} \left\| \dot{\tilde{g}}(t) \right\|_{L^2}^{1+\eta} \left\| \dot{\tilde{g}}(t) \right\|_{L^2}^{1-\eta} \\ &\geq C_6 \left\| \dot{\tilde{g}}(t) \right\|_{C^k}. \end{aligned}$$

for $t \in [1, T)$. Hence by integration,

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$$\int_{1}^{c} \left\| \dot{\tilde{g}}(t) \right\|_{C^{k}} dt \le \frac{1}{C_{6}} |\nu(\tilde{g}(1)) - \nu(g)|^{\theta} \le \frac{1}{C_{6}} |\nu(\tilde{g}(0)) - \nu(g)|^{\theta} \le \frac{\epsilon}{4}$$

provided that \mathcal{V} is small enough. Thus, $T = \infty$ and $\tilde{g}(t)$ converges to some limit metric $g_{\infty} \in \mathcal{U}$ as $t \to \infty$. By the Lojasiewicz–Simon inequality, we have

$$\frac{d}{dt}|\nu(\tilde{g}(t))-\nu(g)|^{1-2\sigma}\geq C_7,$$

which implies

$$|\nu(\tilde{g}(t)) - \nu(g)| \le C_8(t+1)^{-\frac{1}{2\sigma-1}}$$

Therefore, $\nu(g_{\infty}) = \nu(g)$, so g_{∞} is a gradient shrinking Ricci soliton, since it is also a local maximum of ν . The convergence is of polynomial rate, since for $t_1 < t_2$,

$$\|\tilde{g}(t_1) - \tilde{g}(t_2)\|_{C^k} \le C_9 |\nu(\tilde{g}(t_1)) - \nu(g)|^{\theta} \le C_{10}(t_1 + 1)^{-\frac{\nu}{2\sigma-1}}.$$

The assertion follows from $t_2 \to \infty$. By the above arguments, one also sees that $\tilde{g}(t)$ converges in any C^k -norm and therefore, the limit metric is smooth.

Theorem 5.4 (Dynamical instability) Let (M, g) be a gradient shrinking Ricci soliton that is not a local maximizer of v. Then there exists a nontrivial ancient τ -flow $g(t), t \in (-\infty, 0]$ and a 1-parameter family of diffeomorphisms $\varphi_t, t \in (-\infty, 0]$ such that $\varphi_t^*g(t) \to g$ as $t \to \infty$.

Proof Let $g_i \to g$ in C^k and suppose that $\nu(g_i) > \nu(g)$ for all *i*. Let $\tilde{g}_i(t)$ be a solution of (5.2) starting at g_i . Then by Lemma 5.1, $\bar{g}_i = g_i(1)$ converges to g in C^{k-5} and by monotonicity, $\nu(\bar{g}_i) > \nu(g)$ as well. Let $\epsilon > 0$ be so small that Theorem 4.4 holds on $\mathcal{B}_{2\epsilon}^{k-5}$. Then we have the differential inequality

$$\frac{d}{dt}(\nu(\tilde{g}_i(t)) - \nu(g))^{1-2\sigma} \ge -C_1,$$

from which we obtain

$$[(\nu(\tilde{g}_i(t)) - \nu(g))^{1-2\sigma} - C_1(s-t)]^{-\frac{1}{2\sigma-1}} \le (\nu(\tilde{g}_i(s)) - \nu(g)),$$

as long as $\tilde{g}_i(t)$ stays in $\mathcal{B}_{2\epsilon}^{k-5}$. Thus, there exists a t_i such that

$$\|\tilde{g}_i(t_i) - g\|_{C^{k-5}} = \epsilon,$$

and $t_i \to \infty$. If $\{t_i\}$ was bounded, $\tilde{g}_i(t_i) \to g$ in C^{k-5} . By interpolation,

$$\left\|\dot{\tilde{g}}_i(t)\right\|_{C^{k-5}} \le C_2 \left\|\dot{\tilde{g}}_i(t)\right\|_{L^2}^{1-r}$$

for $\eta > 0$ as small as we want. We may assume that $\theta = 1 - \sigma (1 + \eta) > 0$. By Theorem 4.4, we have the differential inequality

$$\frac{d}{dt}(\nu(\tilde{g}_i(t))-\nu(g))^{\theta}\geq C_3\left\|\dot{\tilde{g}}_i(t)\right\|_{L^2}^{1-\eta},$$

if $\nu(\tilde{g}_i(t)) > \nu(g)$. Thus,

$$\epsilon = \|\tilde{g}_i(t_i) - g\|_{C^{k-5}} \le \|\bar{g}_i - g\|_{C^{k-5}} + C_4(\nu(\tilde{g}_i(t_i)) - \nu(g))^{\theta}.$$
(5.12)

Now put $\tilde{g}_i^s(t) := \tilde{g}_i(t+t_i), t \in [T_i, 0]$, where $T_i = 1 - t_i \to -\infty$. We have

$$\begin{aligned} \left\| \tilde{g}_i^s(t) - g \right\|_{C^{k-5}} &\leq \epsilon \quad \forall t \in [T_i, 0] \\ \tilde{g}_i^s(T_i) \to g \text{ in } C^{k-5}. \end{aligned}$$

Because the embedding $C^{k-6}(M) \subset C^{k-5}(M)$ is compact, we can choose a subsequence of the \tilde{g}_i^s , converging in $C_{loc}^{k-6}(M \times (-\infty, 0])$ to an ancient flow $\tilde{g}(t), t \in (-\infty, 0]$, which satisfies the differential equation

$$\dot{\tilde{g}}(t) = -2\left(\operatorname{Ric}_{\tilde{g}(t)} - \frac{1}{2\tau_{\tilde{g}(t)}}\tilde{g}(t) + \nabla^2 f_{\tilde{g}(t)}\right).$$

From taking the limit $i \to \infty$ in (5.12), we have $\epsilon \le C_4(\nu(\tilde{g}(0)) - \nu(g))^{\beta/2}$ which shows that the Ricci flow is nontrivial. For $T_i \le t$, the Lojasiewicz–Simon inequality implies

$$\begin{split} \left\| \tilde{g}_{i}^{s}(T_{i}) - \tilde{g}_{i}^{s}(t) \right\|_{C^{k-6}} &\leq C_{4}(\nu(\tilde{g}_{i}(t+t_{i})) - \nu(g))^{\theta} \\ &\leq C_{4}[-C_{1}t + (\nu(\tilde{g}_{i}(t_{i})) - \nu(g))^{1-2\sigma}]^{-\frac{\theta}{2\sigma-1}} \\ &\leq [-C_{5}t + C_{6}]^{-\frac{\theta}{2\sigma-1}}. \end{split}$$

Thus,

$$\|g - \tilde{g}(t)\|_{C^{k-6}} \le \|g - \tilde{g}_i^s(T_i)\|_{C^{k-6}} + [-C_5t + C_6]^{-\frac{\vartheta}{2\sigma-1}} + \|\tilde{g}_i^s(t) - \tilde{g}(t)\|_{C^{k-6}}.$$

It follows that $\|g - \tilde{g}(t)\|_{C^{k-6}} \to 0$ as $t \to -\infty$.

Remark 5.5 For any τ -flow g(t) we obtain a solution $\hat{g}(t)$ of the normalized Ricci flow (1.1) by projecting to the subset \mathcal{M}_c of metrics of volume c and by rescaling the time parameter suitably. Now if g(t) converges to some Ricci soliton as $t \to \pm \infty$, the same holds for $\hat{g}(t)$. This shows that Theorems 5.3 and 5.4 also hold when replacing the τ -flow by the volume-normalized Ricci flow. In this way, we obtain the theorems as stated in the introduction.

Remark 5.6 All known nontrivial Ricci solitons are dynamically unstable (c.f. Remark 3.8).

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6 The integrability condition

Let g_t a C^1 -curve of Ricci solitons through g and suppose that all g_t are of the same volume c. By projecting to the affine slice, we obtain a C^1 -curve of Ricci solitons

$$\tilde{g}_t \subset \mathcal{M}_c \cap \left\{ g + k \mid \delta_{g, f_g}(k) = 0 \right\}$$

where \mathcal{M}_c is the space of metrics of volume c. We have

$$\frac{d}{dt}\Big|_{t=0}\tilde{g}_t = h \in V = \left\{h \in \Gamma(S^2M) \middle| \delta_f(h) = 0 \text{ and } \int_M \langle \operatorname{Ric}, h \rangle e^{-f} \, dV = 0\right\}$$

and $\nu''(h) = 0$ because ν is constant along \tilde{g}_t . This motivates the following definition:

Definition 6.1 Let (M, g) be a gradient shrinking Ricci soliton and let N be the stability operator of Proposition 3.5. We call $h \in C^{\infty}(S^2M)$ an infinitesimal solitonic deformation if $h \in V$ and N(h) = 0. An infinitesimal solitonic deformation is called integrable if there exists a curve of Ricci solitons g_t through $g = g_0$ such that $\frac{d}{dt}|_{t=0}g_t = h$.

If g is Einstein, this generalizes the notion of infinitesimal Einstein deformations (IED). Recall that an IED is a trace-free and divergence-free (TT) tensor which lies in the kernel of the Einstein operator $\Delta_E = \Delta - 2\mathring{R}$ [2, p. 347]. An IED is called integrable if there exists a curve of Einstein metrics tangent to it.

Lemma 6.2 Let (M, g) be an Einstein manifold with Einstein constant μ . Let IED be the space of infinitesimal Einstein deformations and ISD the space of infinitesimal solitonic deformations. Then we have

$$ISD = IED \oplus \left\{ \mu v \cdot g + \nabla^2 v | v \in C^{\infty}(M), \, \Delta v = 2\mu v \right\}.$$

Proof Since g is Einstein, we have $\Delta_E = -\frac{1}{2}N$ on V, where V is as above. The space V splits as

$$V = TT \oplus \left\{ F(v) = (\Delta v - \mu v)g + \nabla^2 v \mid \int_M v \, dV = 0 \right\}$$

The kernel of Δ_E on TT-tensors is *IED* by definition. On the second component, we have $\Delta_E(F(v)) = F((\Delta - 2\mu)v)$ [6, p. 7]. See also [37, p. 106] for details.

IED's and their integrability were studied by Koiso [30–33], see also [2, Chapter 12]. Recently, Podestà and Spiro [40] generalized some of Koiso's results to the case of Ricci solitons. For a given Ricci soliton g, they constructed a slice S_{g,f_g} in the space of metrics tangent to $\delta_{g,f_g}^{-1}(0)$ such that the following holds: For any $s \ge [\frac{n}{2}]+3$, there exists a H^s -neighbourhood $\mathcal{U} \subset S_{g,f_g}$ of g and a finite-dimensional submanifold $\mathcal{Z} \subset \mathcal{U}$ such that $T_g \mathcal{P} = \ker(N|_V) \oplus \mathbb{R} \cdot \operatorname{Ric}_g$ and the set of Ricci solitons $\mathcal{P} \subset \mathcal{Z}$ is a real analytic subset of \mathcal{Z} [40, Theorem 3.4]. Any metric which is H^s -close to g is isometric to a unique metric in S_{g,f_g} .

For the main theorem in this section, we impose the technical condition that all infinitesimal solitonic deformations are integrable. If this condition holds, $\mathcal{P} = \mathcal{Z}$ provided that \mathcal{U} is small enough. In particular, the set of Ricci solitons in \mathcal{U} is a finite-dimensional manifold.

The slice used in [40] is constructed via the exponential map of the weak Riemannian structure on \mathcal{M} . Hence it differs from the affine slice $\tilde{\mathcal{S}}_{g,f_g}$ that we use in this paper. However,

we can identify the Ricci solitons in S_{g,f_g} and \tilde{S}_{g,f_g} via the map $\Psi : \mathcal{P} \to \tilde{S}_{g,f_g}$ which associates to $g_1 \in S_{g,f_g}$ the unique metric $\tilde{g}_1 \in \tilde{S}_{g,f_g}$ isometric to g_1 . Since \mathcal{P} is finitedimensional, it also consists of all Ricci solitons in a suitable $C^{2,\alpha}$ -neighbourhood of g. By [1, Lemma A.5], Ψ is a diffeomorphism onto its image and the set $\tilde{\mathcal{P}} = \Psi(\mathcal{P})$, consisting of all Ricci solitons in \tilde{S}_{g,f_g} , is a finite-dimensional manifold.

Lemma 6.3 Let (M, g_0) be a gradient shrinking Ricci soliton. Then there exists a $C^{2,\alpha}$ -neighbourhood U of g_0 and a constant C > 0 such that

$$\left\| \frac{d^2}{dtds} \right\|_{t,s=0} f_{g+sk+th} \right\|_{H^1} \le C \|k\|_{C^{2,\alpha}} \|h\|_{H^1},$$
$$\left| \frac{d^2}{dtds} \right|_{t,s=0} \tau_{g+sk+th} \le C \|k\|_{C^{2,\alpha}} \|h\|_{H^1},$$

for all $g \in \mathcal{U}$.

Proof Let the maps *P*, *L* and the space *R* be as in the proof of the Lemma 4.1. We then have

$$\left(\frac{d^2}{dtds}\Big|_{t,s=0}f_{g+sk+th}, \frac{d^2}{dtds}\Big|_{t,s=0}\tau_{g+sk+th}\right) = \frac{d}{ds}\Big|_{s=0}dP_{g+sk}(h).$$

Let us denote $A_g = dL_{g, f_g, \tau_g}|_R$ and $B_g = dL_{g, f_g, \tau_g}|_{C^{2,\alpha}(S^2M)}$, so that we can rewrite (4.6) as

$$dP_g = -A_g^{-1} \circ B_g. ag{6.1}$$

By differentiating,

$$\frac{d}{ds}\Big|_{s=0}dP_{g+sk} = -A_g^{-1}\circ\left(\frac{d}{ds}\Big|_{t=0}A_{g+sk}\right)\circ A_g^{-1}\circ B_g - A_g^{-1}\circ\left(\frac{d}{ds}\Big|_{s=0}B_{g+sk}\right).$$

The derivatives of A_g and B_g contain the second variation of the Laplacian and the scalar curvature, which can be schematically written as

$$\begin{aligned} \frac{d^2}{dsdt} \bigg|_{s,t=0} \Delta_{g+sk+th} f &= \nabla k * h * \nabla f + k * \nabla h * \nabla f, \\ \frac{d^2}{dsdt} \bigg|_{s,t=0} \operatorname{scal}_{g+sk+th} &= \nabla^2 k * h + k * \nabla^2 h + \nabla k * \nabla h + R * k * h, \end{aligned}$$

see [36, Lemma A.3]. Thus by straightforward calculation and standard estimates,

$$\left\| \frac{d}{ds} \right|_{t=0} A_{g+sk}(f) \right\|_{H^{-1}} \le C \|k\|_{C^{2,\alpha}} \|f\|_{H^{1}},$$
$$\left\| \frac{d}{ds} \right|_{t=0} B_{g+sk}(h) \right\|_{H^{-1}} \le C \|k\|_{C^{2,\alpha}} \|h\|_{H^{1}},$$

where $f \in R$ and $h \in C^{2,\alpha}(S^2M)$. By combining these estimates with those for A_g and B_g in the proof of Lemma 4.3, we obtain the result.

Proposition 6.4 (Estimates of the third variation of v) Let (M, g_0) be a gradient shrinking Ricci soliton. There exists a $C^{2,\alpha}$ -neighbourhood U of g_0 such that

$$\left|\frac{d^3}{dt^3}\right|_{t=0} \nu(g+th) \le C \|h\|_{H^1}^2 \|h\|_{C^{2,\alpha}}$$

for all $g \in U$ and some constant C > 0.

Proof We put $v = \frac{e^{-f}}{(4\pi\tau)^{n/2}}$ and $\nabla v = \tau (\operatorname{Ric} + \nabla^2 f) - \frac{1}{2}g$. Then

$$\begin{split} \frac{d^3}{dt^3}\Big|_{t=0} \nu(g+th) &= -\frac{d^2}{dt^2}\Big|_{t=0} \int_M \langle \nabla \nu, h \rangle \nu \, dV \\ &= -\int_M \langle (\nabla \nu)'', h \rangle \nu \, dV - 6 \int_M \langle \nabla \nu, h \circ h \circ h \rangle \nu \, dV - \int_M \langle \nabla \nu, h \rangle (\nu \, dV)'' \\ &+ 2 \int_M \langle (\nabla \nu)', h \circ h \rangle \nu \, dV + 2 \int_M \langle \nabla \nu, h \circ h \rangle (\nu \, dV)' \\ &- \int_M \langle (\nabla \nu)', h \rangle (\nu \, dV)'. \end{split}$$

Here, \circ denotes the composition of symmetric (0, 2)-tensors, considered as endomorphisms on TM. We have

$$\frac{d}{dt}\Big|_{t=0} \nabla_{g+th}^2 f = \nabla h * \nabla f,$$
$$\frac{d}{dt}\Big|_{t=0} \operatorname{Ric}_{g+th} = (\nabla^2 * h) + (R * h)$$

see [36, Lemma A.1 and Lemma A.2]. The second variation of the Ricci tensor and the Hessian appearing in $(\nabla v)''$ are of the form

$$\frac{d}{ds}\frac{d}{dt}\Big|_{s,t=0}\nabla^2_{g+sk+th}f = k*\nabla h*\nabla f + \nabla k*h*\nabla f,$$
$$\frac{d}{ds}\frac{d}{dt}\Big|_{s,t=0}\operatorname{Ric}_{g+sk+th} = k*\nabla^2 h + \nabla^2 k*h + \nabla k*\nabla h + R*k*h,$$

see [36, Lemma A.3]. Now, standard estimates and the Lemmas 4.3 and 6.3 yield an upper bound of the form $C \|h\|_{H^1}^2 \|h\|_{C^{2,\alpha}}$ for each of these terms.

Theorem 6.5 Let (M, g_0) be a shrinking gradient Ricci soliton and suppose that all infinitesimal solitonic deformations are integrable. If g_0 is linearly stable, then there exists a small $C^{2,\alpha}$ -neighbourhood $\mathcal{U} \subset \mathcal{M}$ of g_0 such that $v(g) \leq v(g_0)$ for all $g \in \mathcal{U}$.

Remark 6.6 Observe that the converse implication always holds, even if we drop the integrability assumption.

Proof By diffeomorphism invariance of v, it suffices to prove the theorem on a slice in the space of metrics. Let

$$\mathcal{S}_{g_0,f_{g_0}} = \mathcal{U} \cap \left\{ g_0 + h \mid h \in \delta_{g_0,f_{g_0}}^{-1}(0) \right\}$$

and let

$$\mathcal{P} = \left\{ g \in \mathcal{S}_{g_0, f_{g_0}} \middle| \operatorname{Ric}_g + \nabla^2 f_g = \frac{1}{2\tau_g} \cdot g \right\}$$

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be the set of gradient shrinking Ricci solitons in the slice. By integrability, \mathcal{P} is a finitedimensional manifold with tangent space

$$T_{g_0}\mathcal{P} = \mathbb{R} \cdot \operatorname{Ric} \oplus \ker(N|_{\delta_{f_{g_0}}^{-1}(0)})$$

where *N* is the stability operator in Proposition 3.5. Let *W* be the $L^2(e^{-f_{g_0}} dV)$ -orthogonal complement of $T_{g_0}\mathcal{P}$ in $\delta_{f_{g_0}}^{-1}(0)$. Let us pass to $C^{2,\alpha}$ -spaces for the rest of the proof. By the inverse function theorem for Banach manifolds, any $(C^{2,\alpha})$ -metric *g* in $\mathcal{S}_{g_0,f_{g_0}}$ can be written as $g = \bar{g} + h$, where $\bar{g} \in \mathcal{P}$ and $h \in W$. By Taylor expansion,

$$\nu(\bar{g}+h) = \nu(\bar{g}) + \frac{1}{2} \frac{d^2}{dt^2} \Big|_{t=0} \nu(\bar{g}+th) + R(\bar{g},h),$$
$$R(\bar{g},h) = \int_0^1 \left(\frac{1}{2} - t + \frac{1}{2}t^2\right) \frac{d^3}{dt^3} \nu(\bar{g}+th) dt.$$

By assumption, $v_{g_0}^{''}$ is negative definite on V and therefore,

$$\nu_{g_0}^{\prime\prime}(h) = \frac{\tau}{(4\pi\tau)^{n/2}} \left[-\epsilon \int_M |\nabla h| e^{-f} \, dV + \int_M \left\langle \left[\left(\epsilon - \frac{1}{2} \right) \Delta_f + \mathring{R} \right] h, h \right\rangle e^{-f} \, dV \right] \right]$$

$$\leq -C_1 \left\| h \right\|_{H^1}^2 \tag{6.2}$$

for all $h \in V$. By Taylor expansion again,

$$|\nu_{\bar{g}}''(h) - \nu_{g_0}''(h)| \le C_2 \|\bar{g} - g_0\|_{C^{2,\alpha}} \|h\|_{H^1}.$$
(6.3)

More precisely,

$$\frac{d}{ds}\Big|_{s=0} N_{g+sk}h = \nabla^2 k * h + \nabla k * \nabla h + k * \nabla^2 h + R * h * k + \nabla k * \nabla v_h + \nabla^2 v + \nabla f * \nabla k * h + \nabla f * k * \nabla h + \nabla^2 f' * h + \nabla f' * h$$

by straightforward calculation. By elliptic regularity,

$$\|v_h\|_{H^1} \le C_3 \|h\|_{H^1}$$

and by differentiating (3.7) and Lemma 4.3,

$$\|v'\|_{H^1} \le C_4 \|h\|_{H^1} \|k\|_{C^{2,\alpha}}$$

Together with Lemma 4.3 again, this proves (6.3). As a consequence,

$$\nu_{\bar{g}}^{\prime\prime}(h) \le -C_1 \, \|h\|_{H^1}^2 \tag{6.4}$$

for all $h \in V$ and for all \overline{g} in a sufficiently small neighbourhood of g_0 . Thus by Proposition 6.4,

$$\nu(\bar{g}+h) \leq \nu(\bar{g}) - (C_5 - C_6 \|h\|_{C^{2,\alpha}}) \|h\|_{H^1}^2$$

which shows that any $\nu(g) \leq \nu(g_0)$ for any $g \in S_{g_0, f_{g_0}}$ close enough to g_0 in $C^{2,\alpha}$. By [1, Lemma A.5], any metric in \mathcal{U} is isometric to some metric in $S_{g_0, f_{g_0}}$, so that this inequality holds for all $g \in \mathcal{U}$. In particular, it holds for smooth metrics.

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Corollary 6.7 Let (M, g_0) be a gradient shrinking Ricci soliton and suppose that all infinitesimal solitonic deformations are integrable. Then, linear stability and dynamical stability are equivalent.

Remark 6.8 It would be interesting to find curvature conditions which ensure linear stability of a given Ricci soliton. This would generalize previous results obtained in the Einstein case, see [12, 13, 27, 30, 33, 35]. A detailed study of the stability operator N will be the content of our future investigations.

Proof of Theorem 1.5 The Fubini-Study metric is Einstein and its Einstein constant is given by $1/2\tau$. By [36, pp. 25–26], there exists an eigenfunction $v \in C^{\infty}(M)$ of the Laplacian with eigenvalue $1/\tau$ such that $\int_{M} v^3 dV \neq 0$. By Lemma 6.2, $h = v \cdot g_{fs} + 2\tau \nabla^2 v$ is an infinitesimal solitonic deformation. Suppose there exists a curve of Ricci solitons g_t such that $\frac{d}{dt}|_{t=0}g_t = h$. Since g_{fs} is a critical point of v and h lies in the kernel of its linearization, $\frac{d}{dt}|_{t=0}v(g_t) = \frac{d^2}{dt^2}|_{t=0}v(g_t) = 0$. The third variation equals

$$\frac{d^3}{dt^3}\Big|_{t=0}v(g_t) = v'''(h) = v'''(v \cdot g_{fs}) = \frac{3n-4}{\operatorname{vol}(M, g_{fs})} \int_M v^3 \, dV \neq 0.$$

Here we used the diffeomorphism invariance of v and the third variational formula in [36, Proposition 9.1]. This contradicts the fact that v must be constant along g_t and thus, h is not integrable.

Remark 6.9 It seems likely that $(\mathbb{C}P^n, g_{fs})$ is isolated in the moduli space of Ricci solitons, because all infinitesimal solitonic deformations arise from conformal deformations.

Remark 6.10 The space ($\mathbb{C}P^n$, g_{fs}) is neutrally linearly stable but dynamically unstable [36, Corollary 1.11].

Acknowledgments The author would like to thank Christian Bär, Stuart Hall and Thomas Murphy for their support and helpful discussions. Moreover, the author thanks Sonderforschungsbereich 647 funded by Deutsche Forschungsgemeinschaft for financial support.

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