

Radial and non radial solutions for Hardy–Hénon type elliptic systems

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Received: 15 May 2009 / Accepted: 20 September 2009 / Published online: 17 October 2009
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Abstract We discuss existence and non-existence of positive solutions for the following system of Hardy and Hénon type:

$$\begin{cases} -\Delta v = |x|^\alpha u^p, & -\Delta u = |x|^\beta v^q & \text{in } \Omega, \\ u = v = 0 & & \text{on } \partial\Omega, \end{cases}$$

where $\Omega \ni 0$ is a bounded domain in \mathbb{R}^N , $N \geq 3$, $p, q > 1$, and $\alpha, \beta > -N$. We also study symmetry breaking for ground states when Ω is the unit ball in \mathbb{R}^N .

Mathematics Subject Classification (2000) 35J47 · 35J60 · 35J20

1 Introduction

In this paper we consider the following system of superlinear elliptic equations of Hardy and Hénon type:

$$\begin{cases} -\Delta v = |x|^\alpha u^p & \text{in } \Omega, \\ -\Delta u = |x|^\beta v^q & \text{in } \Omega, \\ u > 0 & \text{in } \Omega, \\ v > 0 & \text{in } \Omega, \\ u = v = 0 & \text{on } \partial\Omega, \end{cases} \quad (1)$$

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where Ω is a bounded domain in \mathbb{R}^N with $0 \in \Omega$, $N \geq 3$, $p, q > 1$, and $\alpha, \beta > -N$. In particular, we will investigate existence, multiplicity and qualitative properties (such as radial symmetry in the case Ω a ball) of solutions.

The case of a single equation (especially the case $\alpha > 0$) has been widely studied (see for instance [2–4, 7, 24, 26], and the references therein). Recall that the equation

$$\begin{cases} -\Delta u = |x|^\alpha u^p & \text{in } \Omega \\ u = 0 & \text{on } \partial\Omega, \end{cases} \tag{2}$$

is called of Hardy type if $\alpha < 0$ (because of its relation to the Hardy-Sobolev inequality) and it is of Hénon type if $\alpha > 0$ (this equation was introduced by Hénon in 1973 [15] for the study of stellar systems). One has the following results:

Hardy type: By variational methods one obtains the existence of a nontrivial solution of (2) in $H_0^1(\Omega)$ ($\Omega \subset \mathbb{R}^N$ an arbitrary bounded domain) provided that $2 < p + 1 < \frac{2(N-|\alpha|)}{N-2}$, by an application of the celebrated Caffarelli-Kohn-Nirenberg estimates (CKN, [6]), and due to a generalized Pohozaev type identity one proves non-existence of nontrivial solutions in starshaped domains if $0 \geq \alpha > -N$ and $p + 1 = \frac{2(N-|\alpha|)}{N-2}$.

Hénon type: For $\alpha > 0$ one obtains for Ω arbitrary (bounded) the existence of a solution for $2 < p + 1 < \frac{2N}{N-2}$. On the other hand, if Ω is a ball, then one has the existence of a radial solution in a larger range, namely for $2 < p + 1 < \frac{2(N+\alpha)}{N-2}$ (see [23]), and the non-existence of non-trivial solutions in the range: $p + 1 \geq \frac{2(N+\alpha)}{N-2}$ by a Pohozaev-type identity for radial functions.

Concerning the symmetry of solutions one has the following result: if Ω is a ball, one proves by moving plane techniques [12] that the *minimal energy* solution is positive and *radially symmetric* if $\alpha \leq 0$ (Hardy case). One can pose the question if this symmetry continues to be present also for positive α . In an interesting paper Smets et al. [26] showed that this is not the case: they proved that for $\alpha > 0$ and sufficiently large a *symmetry-breaking* occurs, that is, to the minimal energy level (which is attained for $p + 1 < 2^* = \frac{2N}{N-2}$) corresponds a solution which is *not radially symmetric*. In a related result, Cao and Peng proved in [5] that for $\alpha > 0$ and $p + 1$ sufficiently close to 2^* the ground-state solutions of (2) are not radial since they possess a unique maximum point which tends to $\partial\Omega$ as $p + 1 \rightarrow 2^*$.

Turning to the system (1), we first recall the case $\alpha = \beta = 0$ which has been studied by many authors. Here the natural restriction on the exponents p and q for existence/non-existence of solutions is given by the *critical hyperbola*, that is

$$\frac{N}{p + 1} + \frac{N}{q + 1} = N - 2; \tag{3}$$

this hyperbola was first introduced by Mitidieri [21] who proved non-existence of solutions for (p, q) lying on or above the hyperbola, using a Pohozaev-type identity. Existence of solutions for $(p + 1, q + 1)$ below the critical hyperbola was proved by de Figueiredo and Felmer (see [9]) and by Hulshoff and van der Vorst (see [16, 17]) by using a variational set-up with fractional Sobolev spaces. A different approach, working with Sobolev-Orlicz spaces (which allows a generalization to non-polynomial nonlinearities), can be found in [8, 10].

Recently, the general case $\alpha \neq 0$, and/or $\beta \neq 0$ has been investigated independently by de Figueiredo et al. [11] and Liu and Yang [20]; in both papers an approach via fractional Sobolev spaces is used. As in the scalar case, the presence of the weight functions $|x|^\alpha$ and $|x|^\beta$ affects the range of p and q for which the problem may have solutions. Indeed, in [11] and [20] it is shown that the dividing line between existence and non-existence is given by the following “weighted” critical hyperbola

$$\frac{N + \alpha}{p + 1} + \frac{N + \beta}{q + 1} = N - 2. \tag{4}$$

For future reference, we call the hyperbola (3) the M-hyperbola (for *Mitidieri hyperbola*), and the hyperbola (4) the $\alpha\beta$ -hyperbola).

We remark that systems of type (1) are closely related to the double weighted Hardy–Littlewood–Sobolev inequality (see e.g. Stein and Weiss [27] and Lieb [18]). This becomes clear by the approach we use for system (1), which is different from the one proposed in the previously cited papers. Following for instance Wang [28] we (formally) deduce from the second equation in (1)

$$v = (-\Delta u)^{\frac{1}{q}} |x|^{-\frac{\beta}{q}},$$

and inserting this into the first equation we obtain the following scalar equation for the u -component

$$-\Delta ((-\Delta u)^{1/q} |x|^{-\beta/q}) = |x|^{\alpha} u^p \tag{5}$$

We intend to investigate the rôle played by the weights α and β when dealing with the existence and symmetry of *ground state* (or *minimal energy*) solutions of Eq. 5, that is, minimizers u of the corresponding Raleigh quotient

$$R(u) = \frac{\int_{\Omega} |x|^{-\beta(r-1)} |\Delta u|^r}{\left(\int_{\Omega} |x|^{\alpha} |u|^{p+1} dx\right)^{\frac{r}{p+1}}}, \quad r := \frac{q + 1}{q}, \tag{6}$$

on the weighted Sobolev space

$$W^{2,r}(\Omega, |x|^{-\beta(r-1)} dx) \cap W_0^{1,r}(\Omega)$$

Here we denote with $W^{2,r}(\Omega, |x|^{-\beta(r-1)} dx)$ the set of functions $u \in W_{loc}^{2,1}(\Omega)$ such that

$$\int_{\Omega} (|u|^r + |\nabla u|^r + \sum_{|\xi|=2} |D^{\xi} u|^r |x|^{-\beta(r-1)}) dx < +\infty,$$

endowed with the norm

$$\|u\|_{W^{2,r}(\Omega, |x|^{-\beta/q} dx)} := \left(\int_{\Omega} (|u|^r + |\nabla u|^r + \sum_{|\xi|=2} |D^{\xi} u|^r |x|^{-\beta(r-1)} dx) dx \right)^{1/r};$$

also, we denote with

$$W_{rad}^{2,r}(\Omega, |x|^{-\beta(r-1)} dx)$$

the subspace of $W^{2,r}(\Omega, |x|^{-\beta(r-1)} dx)$ of radial functions.

Furthermore, let

$$W_{D}^{2,r}(\Omega, |x|^{-\beta(r-1)} dx)$$

denote the closure of $\{\phi \in C^{\infty}(\Omega) : \phi = 0 \text{ on } \partial\Omega\}$ in $W^{2,r}(\Omega, |x|^{-\beta(r-1)} dx)$, i.e. the closure of the smooth functions in Ω with Dirichlet boundary conditions, and with

$$W_{D,rad}^{2,r}(\Omega, |x|^{-\beta(r-1)} dx)$$

the corresponding subspace of radial functions (for Ω a ball).

For the definition of this space and related properties we refer the reader to the Appendix. In particular, we prove there the following generalization of the Meyers-Serrin denseness result:

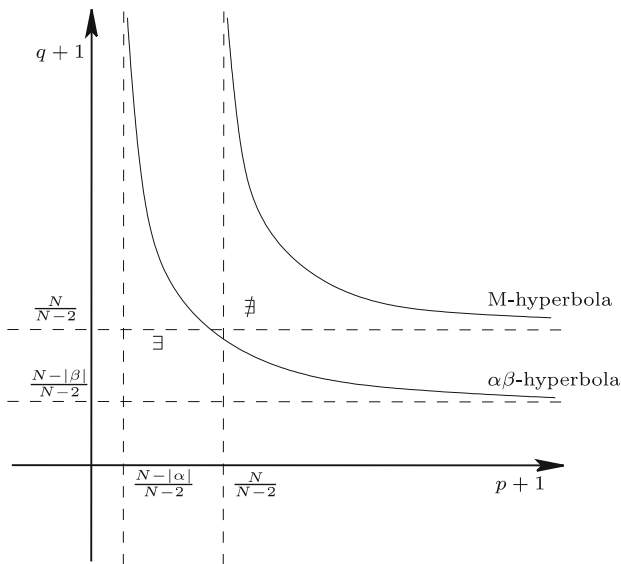


Fig. 1 Hardy type system

Proposition *Let Ω a domain with smooth boundary. Suppose that $\alpha, \beta > -N$ and $p, q > 1$ with $q > \frac{\beta}{N}$. Then*

$$W_D^{2,r}(\Omega, |x|^{-\beta/q} dx) = W^{2,r}(\Omega, |x|^{-\beta/q} dx) \cap W_0^{1,r}(\Omega).$$

It is not difficult to prove that critical points of $R(u)$ on $W_D^{2,r}(\Omega, |x|^{-\beta/q} dx)$ are (up to rescaling) weak solutions of (5), i.e. verifying

$$\begin{cases} \int_{\Omega} (-\Delta u)^{1/q} |x|^{-\beta/q} (-\Delta \varphi) dx = \int_{\Omega} |x|^{\alpha} u^p \varphi dx, \\ \text{for all } \varphi \in W_D^{2,r}(\Omega, |x|^{-\beta/q} dx) \end{cases}$$

and, moreover, if $v = (-\Delta u)^{\frac{1}{q}} |x|^{-\frac{\beta}{q}}$, then $v \in W^{2, \frac{p+1}{p}}(\Omega, |x|^{-\frac{\alpha}{p}} dx) \cap W_0^{1, \frac{p+1}{p}}(\Omega)$. In accordance, by a *strong solution* of the system we mean a couple (u, v) of weak-solutions such that

$$(u, v) \in W^{2,r}(\Omega, |x|^{-\frac{\beta}{q}} dx) \cap W_0^{1,r}(\Omega) \times W^{2, \frac{p+1}{p}}(\Omega, |x|^{-\frac{\alpha}{p}} dx) \cap W_0^{1, \frac{p+1}{p}}(\Omega).$$

In what follows, we will denote by E the space $E = W_D^{2,r}(\Omega, |x|^{-\beta(r-1)} dx) = W^{2,r}(\Omega, |x|^{-\beta/q} dx) \cap W_0^{1,r}(\Omega)$ (if the values β and r are clear from the context), and by E_{rad} the radial component of E .

In this paper we investigate solvability and symmetry properties of the solutions for general exponents α and β . We will see that the solvability and the qualitative properties of the solutions depend on the location of the exponents p, q with respect to the M-hyperbola and the $\alpha\beta$ -hyperbola.

Note that if $\alpha, \beta < 0$, then the $\alpha\beta$ -hyperbola lies below the M-hyperbola, see Fig. 1. In this case, we have the following result:

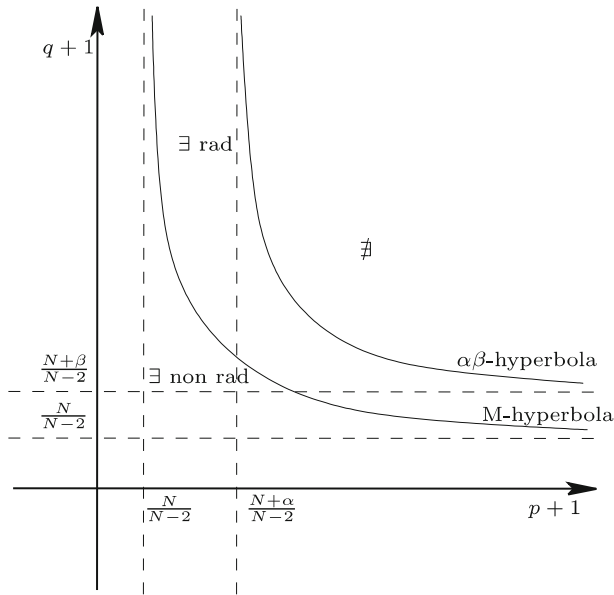


Fig. 2 Hénon type system

Theorem 1 (Hardy-type system) *Let $0 \geq \alpha, \beta > -N$, and $p, q > 1$.*

- (a) *If $\frac{N-|\alpha|}{p+1} + \frac{N-|\beta|}{q+1} = N - 2$ (i.e. on the $\alpha\beta$ -hyperbola), and if Ω is starshaped, then system (1) has no non-trivial solution, and hence $\inf_E R(u)$ is not attained.*
- (b) *If $\frac{N-|\alpha|}{p+1} + \frac{N-|\beta|}{q+1} > N - 2$ (i.e. below the $\alpha\beta$ -hyperbola), then*
 - (b₁) *$\inf_E R(u)$ is attained, and therefore system (1) has a nontrivial solution \bar{u} (the minimal energy solution);*
 - (b₂) *if Ω is a ball, then \bar{u} is radially symmetric*

Next, we consider the case $\alpha, \beta > 0$, i.e. the Hénon-type system. Note that then the $\alpha\beta$ -hyperbola lies above the M-hyperbola, and there are three regions which characterize the behavior of the system, see Fig. 2.

Theorem 2 (Hénon type system) *Let $\alpha, \beta > 0$, and suppose that $p, q > 1$.*

- (a) *If $\frac{N+\alpha}{p+1} + \frac{N+\beta}{q+1} \leq N - 2$ (i.e., on or above the $\alpha\beta$ -hyperbola) and Ω is starshaped, then system (1) has no non-trivial solution;*
- (b) *Suppose that $\Omega = B_1(0)$. If $q > \max\{1, \frac{\beta}{N}\}$ and $\frac{N+\alpha}{p+1} + \frac{N+\beta}{q+1} > N - 2$ (i.e. below the $\alpha\beta$ -hyperbola), then system (1) has a radial solution (not necessarily of minimal energy);*
- (c) *If $\frac{N}{p+1} + \frac{N}{q+1} > N - 2$ (i.e. below the M-hyperbola), then $\inf_E R(u)$ is attained and hence system (1) has a ground state solution; furthermore, if $\alpha > 0$ is sufficiently large, then the ground state solution is not radially symmetric.*

Remarks (1) Note that (c) of the previous theorem can be interpreted as a symmetry breaking: for $\alpha, \beta \leq 0$ we have by Theorem 1 that the ground state solution in a ball is radial, while (c) says that for $\alpha, \beta > 0$ and α large the ground state solution is non radial.

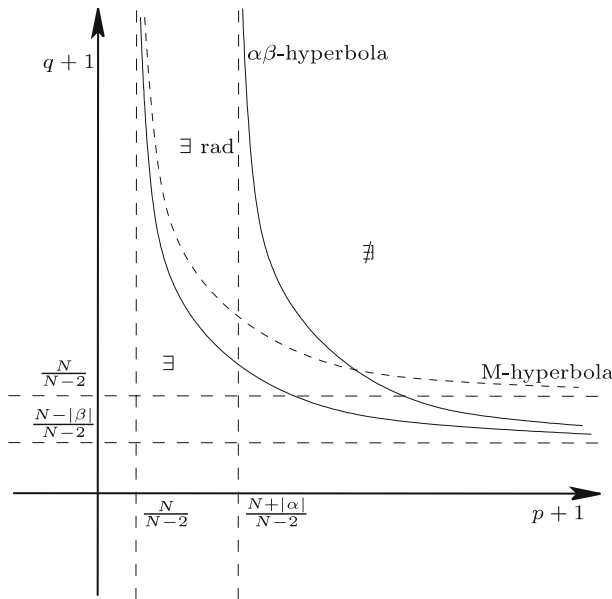


Fig. 3 Hardy-Hénon type system

- (2) By (b) and (c) of Theorem 2 we get the existence of at least two solutions for $(p+1, q+1)$ below the M-hyperbola and α large: one radial solution (obtained as minimum of $R(u)$ on E_{rad}), and the minimal energy solution which is non-radial.

In the Hénon case there is a very recent symmetry breaking result, due to He and Yang ([14]): they prove that if (for q fixed) p goes towards the M-hyperbola, then the ground state solution is non radial; this result extends to systems the corresponding result for the equation by Cao and Peng [5]).

Finally, we consider the case of a mixed *Hénon-Hardy type system*, i.e. one exponent (say α) is positive and the other exponent (i.e. β) is negative. In this case the M-hyperbola and the $\alpha\beta$ -hyperbola intersect, see Fig. 3. We show that in this case a third hyperbola comes into play.

Theorem 3 (Hénon-Hardy type system) *Let $\alpha > 0, 0 > \beta > -N$, and suppose that $p, q > 1$.*

- (a) *If $\frac{N+\alpha}{p+1} + \frac{N-|\beta|}{q+1} = N - 2$ (i.e. on the $\alpha\beta$ -hyperbola) and Ω is starshaped, then system (1) has no solution;*
- (b) *Assume $\Omega = B_1(0)$. If $\frac{N+\alpha}{p+1} + \frac{N-|\beta|}{q+1} > N - 2$ (i.e. below the $\alpha\beta$ -hyperbola), then system (1) has a radial solution (not necessarily of minimal energy);*
- (c) *If (p, q) satisfies $\frac{N}{p+1} + \frac{N-|\beta|}{q+1} > N - 2$, then $\inf_E R(u)$ is attained and hence system (1) has a ground state solution; furthermore, if $\alpha > 0$ is sufficiently large, then the ground state solution is not radially symmetric.*

Remark Suppose that Ω is an arbitrary bounded domain: if $(p + 1, q + 1)$ lies below the $\alpha\beta$ -hyperbola and above the M-hyperbola, then it is *not known* whether $\inf_E R(u)$ is attained.

2 The Hardy case: proof of Theorem 1

Proof of Theorem 1 (a): this is obtained via a generalized identity of Pohozaev-type, see Proposition 7 Sect. 6 below. Note that this case is somewhat delicate due to the singular weights.

Proof of Theorem 1 (b₁): To prove the existence of a positive solution we minimize the Rayleigh quotient $R(u)$ given in (6). By the compactness of the embedding in Lemma 4 (see Sect. 5) the infimum is attained by a positive function, which is a (strong) solution of problem (10).

Proof of Theorem 1 (b₂): for $\alpha = \beta = 0$ and if Ω is a ball, it was proved by X.J. Wang [28] that the ground state of (6) is a radial and radially decreasing positive function. By adapting his argument (moving planes technique and maximum principle) one can extend this result also to the values $\alpha, \beta < 0$ (noting that the weights do not interfere with the moving planes technique). □

3 The Hénon case: proof of Theorem 2

Proof of Theorem 2 (a): this follows again by a Pohozaev-type identity proved in Proposition 6 in Sect. 6 below.

Proof of Theorem 2 (b): The existence of radial solutions under the hypotheses of Theorem 2(b) follows from the embedding result for radial functions in Lemma 9 in Sect. 7 below, by considering again the Rayleigh quotient $R(u)$ on the weighted space $W_{rad}^{2,r}(\Omega, |x|^{-\beta/q} dx)$.

Proof of Theorem 2 (c): We have to show that $m := \inf_E R(u)$ is attained. Let $\{u_n\} \subset E$ be a minimizing sequence. We may assume that

$$\int_{\Omega} |x|^{\alpha} |u_n|^{p+1} = 1, \quad \text{and} \quad \int_{\Omega} |x|^{-\beta/N} |\Delta u_n|^r dx \rightarrow m > 0.$$

Then clearly $\int_{\Omega} |\Delta u_n|^r dx \leq c$, and by the assumption and the Rellich-Kondrachov compactness theorem it follows that $\{u_n\}$ has a convergent subsequence in $L^{p+1}(\Omega)$, and hence also in $L^{p+1}(\Omega, |x|^{\alpha} dx)$. This is sufficient to conclude that $\inf_E R(u)$ is attained. □

Finally, we show that if $\alpha > 0$ is sufficiently large, then the radial ground state level lies above the ground state level: indeed, by Proposition 10 below we have the following lower estimate for the radial ground state level:

$$S_{\alpha,\beta}^{rad} \geq C \alpha^{2r + \frac{r}{p+1} - 1}, \quad \text{for } \alpha \geq \alpha_0$$

On the other hand, for the ground state level the following upper estimate holds (see Proposition 11 below): there exist $C > 0$ and α_0 such that for $\alpha \geq \alpha_0$

$$S_{\alpha,\beta} \leq C \alpha^{2r - N + N \frac{r}{p+1}}$$

From these two inequalities it follows that the ground state is non radial for α sufficiently large, since

$$\frac{r}{p+1} - 1 > -N + N \frac{r}{p+1} \iff \frac{r}{p+1} < 1,$$

which is clearly the case.

4 The mixed case: proof of Theorem 3

Proof of Theorem 3 (a) and (b): as in Theorem 2

Proof of Theorem 3 (c): Let $\{u_n\} \subset E$ be a minimizing sequence for $m = \inf R(u)$. We may assume that

$$\int_{\Omega} |x|^{\alpha} |u_n|^{p+1} = 1, \quad \text{and} \quad \int_{\Omega} |x|^{|\beta|/N} |\Delta u_n|^r dx \rightarrow m.$$

We apply the embedding result in Lemma 4 (see Sect. 5 below) for $\alpha = 0$, i.e. under the hypotheses of Theorem 3 c). By the compactness of the embedding we have that for a subsequence $u_n \rightarrow u$ in $L^{p+1}(\Omega)$, and since $\alpha > 0$ clearly also in $L^{p+1}(\Omega, |x|^{\alpha} dx)$. Thus is sufficient to conclude that $\inf_E R(u)$ is attained.

Proof of Theorem 3 (c): one proves as in Theorem 2 c) that if $\alpha > 0$ is sufficiently large, then the ground state level is non radial. □

5 An embedding result of Caffarelli-Kohn-Nirenberg type

We first prove a preliminary embedding result:

Lemma 4 *Let $\Omega \subset \mathbb{R}^N$ be a bounded smooth domain with $0 \in \Omega$. For given α, β, q let p^* such that*

$$\frac{N - |\alpha|}{p^* + 1} + \frac{N - |\beta|}{q + 1} = N - 2. \tag{7}$$

Then we have the following continuous embedding

$$W^{2,r}(\Omega, |x|^{-\beta(r-1)} dx) \hookrightarrow L^{p+1}(\Omega, |x|^{\alpha}), \quad \text{for } 0 \leq p \leq p^*;$$

furthermore, if

$$\frac{N - |\alpha|}{p + 1} + \frac{N - |\beta|}{q + 1} > N - 2, \quad \text{i.e. } p < p^*, \tag{8}$$

then the embedding is compact.

Proof This follows from the generalization due to Lin (see [19]) of the Caffarelli-Kohn-Nirenberg inequality [6]: if (7) holds, then there exists a constant C such that

$$\left(\int_{\mathbb{R}^N} |x|^{\alpha} |u|^{(p^*+1)} \right)^{\frac{1}{(p^*+1)}} \leq C \left(\int_{\mathbb{R}^N} |x|^{-\beta(r-1)} |D^2 u|^r \right)$$

for all $u \in C_0^{\infty}(\mathbb{R}^N)$.

Via extension theorems (see for instance [1]) we have

$$W^{2,r}(\Omega, |x|^{-\beta(r-1)} dx) \hookrightarrow L^{p^*+1}(\Omega, |x|^{\alpha}),$$

and for any $0 \leq p < p^*$ the embedding is compact. □

Another consequence of the CKN inequalities is the following

Lemma 5 *If $(u, v) \in W_D^{2, \frac{q+1}{q}}(\Omega, |x|^{|\beta|/q} dx) \times W_D^{2, \frac{p+1}{p}}(\Omega, |x|^{|\alpha|/p} dx)$, with p and q as in (7), then $\nabla u \nabla v \in L^1(\Omega)$.*

Proof If $u \in W_D^{2, \frac{q+1}{q}}(\Omega, |x|^{|\beta|/q} dx)$ and $v \in W_D^{2, \frac{p+1}{p}}(\Omega, |x|^{|\alpha|/p} dx)$ then from the CKN inequality applied to ∇u and ∇v , one has $\nabla u \in L^s(\Omega)$ with $s = \frac{N(q+1)}{Nq+|\beta|-(q+1)}$, and $\nabla v \in L^t(\Omega)$ with $t = \frac{N(p+1)}{Np+|\alpha|-(p+1)}$. It is not difficult to prove that $\frac{1}{s} + \frac{1}{t} = 1$. So the assertion follows by the Hölder inequality. \square

6 A generalized identity of Pohozaev-type

In this section we prove, via a generalized identity of Pohozaev-type, the non-existence of “strong” solutions on or above the critical $\alpha\beta$ -hyperbola.

We consider first the Hénon case $\alpha \geq 0, \beta \geq 0$. In this case we can suppose that the solutions (u, v) are of class $C^2(\Omega) \cap C^1(\bar{\Omega})$.

Proposition 6 *Let $\Omega \subset \mathbb{R}^N$ be a bounded, smooth, starshaped domain with respect to $0 \in \mathbb{R}^N$. If $\alpha, \beta \geq 0$ and*

$$\frac{N + \alpha}{p + 1} + \frac{N + \beta}{q + 1} \leq N - 2 \tag{9}$$

then the problem

$$\begin{cases} -\Delta v = |x|^\alpha |u|^{p-1} u & \text{in } \Omega, \\ -\Delta u = |x|^\beta |v|^{q-1} v & \text{in } \Omega, \\ u = v = 0 & \text{on } \partial\Omega, \end{cases} \tag{10}$$

has no nontrivial strong positive solutions.

Proof This follows by adapting the argument of Mitidieri in [21]. Let

$$G(x, u, v) = \frac{1}{p + 1} |x|^\alpha |u|^{p+1} + \frac{1}{q + 1} |x|^\beta |v|^{q+1} \tag{11}$$

so (10) becomes

$$\begin{cases} -\Delta v = \frac{\partial G}{\partial u}(x, u, v) & \text{in } \Omega \\ -\Delta u = \frac{\partial G}{\partial v}(x, u, v) & \text{in } \Omega \\ u = v = 0 & \text{on } \partial\Omega \end{cases} \tag{12}$$

We multiply the equations respectively by $(x \cdot \nabla u)$ and by $(x \cdot \nabla v)$, add the two equations and integrate. By an application of the divergence theorem (see Proposition 2.1 and Corollary 2.1 in [21]) we have for the left sides

$$\begin{aligned} & \int_{\Omega} \{ \Delta u (x \cdot \nabla v) + \Delta v (x \cdot \nabla u) \} dx \\ &= \int_{\partial\Omega} \left\{ \frac{\partial u}{\partial n} \frac{\partial v}{\partial n} (x \cdot n) \right\} ds + (N - 2) \int_{\Omega} (\nabla u \nabla v) dx \end{aligned} \tag{13}$$

(since $u = v = 0$ on $\partial\Omega$, we have $(x \cdot \nabla u) = x \cdot n \frac{\partial u}{\partial n}$ and $(x \cdot \nabla v) = x \cdot n \frac{\partial v}{\partial n}$ on $\partial\Omega$).

For the right sides we have, since

$$\begin{aligned} \frac{\partial G}{\partial u}(x \cdot \nabla u) + \frac{\partial G}{\partial v}(x \cdot \nabla v) &= \operatorname{div}\{xG(x, u, v)\} \\ &\quad - NG(x, u, v) - \alpha \frac{|x|^\alpha}{p+1} |u|^{p+1} - \beta \frac{|x|^\beta}{q+1} |v|^{q+1} \end{aligned} \tag{14}$$

and taking into account that $G(x, u, v) = 0$ on $\partial\Omega$:

$$\begin{aligned} &\int_{\Omega} \left\{ \frac{\partial G}{\partial u}(x \cdot \nabla u) + \frac{\partial G}{\partial v}(x \cdot \nabla v) \right\} dx \\ &= -\frac{N+\alpha}{p+1} \int_{\Omega} |x|^\alpha |u|^{p+1} dx - \frac{N+\beta}{q+1} \int_{\Omega} |x|^\beta |v|^{q+1} dx. \end{aligned} \tag{15}$$

Therefore

$$\begin{aligned} &\int_{\partial\Omega} \left\{ \frac{\partial u}{\partial n} \frac{\partial v}{\partial n} (x \cdot n) \right\} ds + (N-2) \int_{\Omega} (\nabla u \nabla v) dx \\ &= \frac{N+\alpha}{p+1} \int_{\Omega} |x|^\alpha |u|^{p+1} dx + \frac{N+\beta}{q+1} \int_{\Omega} |x|^\beta |v|^{q+1} dx. \end{aligned} \tag{16}$$

Now, multiplying the first equation by u , the second by v and integrating, one obtains

$$\int_{\Omega} (\nabla u \nabla v) dx = \int_{\Omega} |x|^\alpha |u|^{p+1} dx = \int_{\Omega} |x|^\beta |v|^{q+1} dx$$

So (16) becomes (*Pohozaev identity*)

$$\begin{aligned} &\int_{\partial\Omega} \left\{ \frac{\partial u}{\partial n} \frac{\partial v}{\partial n} (x \cdot n) \right\} ds \\ &= \left[-(N-2) + \frac{N+\alpha}{p+1} + \frac{N+\beta}{q+1} \right] \int_{\Omega} |x|^\alpha |u|^{p+1} dx \end{aligned} \tag{17}$$

Since Ω is starshaped and u, v are positive, we have $\int_{\partial\Omega} \left\{ \frac{\partial u}{\partial n} \frac{\partial v}{\partial n} (x \cdot n) \right\} ds > 0$, and hence by (16)

$$\begin{aligned} 0 &< \int_{\partial\Omega} \left\{ \frac{\partial u}{\partial n} \frac{\partial v}{\partial n} (x \cdot n) \right\} ds \\ &= \left[-(N-2) + \frac{N+\alpha}{p+1} + \frac{N+\beta}{q+1} \right] \int_{\Omega} |x|^\alpha |u|^{p+1} dx \end{aligned} \tag{18}$$

which gives a contradiction for the choice of p and q . □

Next, we consider the Hardy case.

Proposition 7 *Let Ω as in Proposition 6 and assume that $0 \geq \alpha, \beta > -N$ and*

$$\frac{N-|\alpha|}{p+1} + \frac{N-|\beta|}{q+1} = N-2; \tag{19}$$

then there exists no positive strong solution (u, v) of (12).

Proof For this case we follow the idea developed by B. Xuan (see Appendix in [29]).

Let (u, v) a positive solution of system (12). Due to the Hardy weights this solution may be singular in the origin, but standard regularity results imply that for every δ small, u and v belong to $C^2(\Omega \setminus B_\delta(0)) \cap C^0(\overline{\Omega} \setminus B_\delta(0))$. We multiply the equations respectively by $(x \cdot \nabla u)$ and by $(x \cdot \nabla v)$, add the two equations and integrate over $\Omega_\delta = \Omega \setminus B_\delta(0)$

$$\begin{aligned}
 & - \int_{\Omega_\delta} \{ \Delta u (x \cdot \nabla v) + \Delta v (x \cdot \nabla u) \} dx \\
 & = \int_{\Omega_\delta} \left\{ \frac{\partial G}{\partial u} (x \cdot \nabla u) + \frac{\partial G}{\partial v} (x \cdot \nabla v) \right\} dx, \tag{20}
 \end{aligned}$$

where $G = G(x, u, v)$ is as in (11).

We apply the Divergence Theorem to xG , so that one has for the right side of (20)

$$\begin{aligned}
 & \int_{\Omega_\delta} \left\{ \frac{\partial G}{\partial u} (x \cdot \nabla u) + \frac{\partial G}{\partial v} (x \cdot \nabla v) \right\} dx \\
 & = - \frac{N - |\alpha|}{p + 1} \int_{\Omega_\delta} \frac{u^{p+1}}{|x|^\alpha} dx - \frac{N - |\beta|}{q + 1} \int_{\Omega_\delta} \frac{v^{q+1}}{|x|^\beta} dx + \int_{|x|=\delta} G(u, v, x)(x \cdot n) ds, \tag{21}
 \end{aligned}$$

while for the left side of (20) (see [21], Corollary 2.1) one has

$$\begin{aligned}
 & \int_{\Omega_\delta} \{ \Delta u (x \cdot \nabla v) + \Delta v (x \cdot \nabla u) \} dx \\
 & = \int_{\partial\Omega_\delta} \left\{ \frac{\partial u}{\partial n} (x \cdot \nabla v) + \frac{\partial v}{\partial n} (x \cdot \nabla u) - (\nabla u \nabla v)(x \cdot n) \right\} ds + (N - 2) \int_{\Omega_\delta} (\nabla u \nabla v) dx \\
 & = \int_{|x|=\delta} \left\{ \frac{\partial u}{\partial n} (x \cdot \nabla v) + \frac{\partial v}{\partial n} (x \cdot \nabla u) - (\nabla u \nabla v)(x \cdot n) \right\} ds \\
 & \quad + \int_{\partial\Omega} \left\{ \frac{\partial u}{\partial n} \frac{\partial v}{\partial n} (x \cdot n) \right\} ds + (N - 2) \int_{\Omega_\delta} (\nabla u \nabla v) dx. \tag{22}
 \end{aligned}$$

Now, multiplying the first equation by u , the second by v and integrating, one obtains

$$\int_{\Omega_\delta} (\nabla u \nabla v) dx - \int_{|x|=\delta} v \nabla u \cdot x = \int_{\Omega_\delta} |x|^\beta |v|^{q+1} dx$$

and

$$\int_{\Omega_\delta} (\nabla u \nabla v) dx - \int_{|x|=\delta} u \nabla v \cdot x = \int_{\Omega_\delta} |x|^\alpha u^{p+1} dx$$

The Pohozaev identity (17) follows if we prove that all the integrals along $\{|x| = \delta\}$ go to zero, at least for a subsequence $\delta_k \rightarrow 0$. But this follows by the mean value theorem, since by the Lemmas in Sect. 5 the integrals

$$\int_{\Omega} G(u, v, x)dx, \int_{\Omega} |u \nabla v| dx, \int_{\Omega} |v \nabla u| dx, \int_{\Omega} |\nabla u \nabla v| dx$$

are finite. Indeed, if ψ is a positive function in $L^1(\Omega)$, then $\epsilon_k = \int_{B_{1/k}(0)} \psi(x)dx \rightarrow 0$ as $k \rightarrow +\infty$. Moreover, if $\psi \in C(\Omega \setminus \{0\})$ then

$$\epsilon_k = \int_{B_{1/k}(0)} \psi(x)dx = \int_0^{1/k} \int_{|x|=\delta} \psi(x)ds d\delta$$

By the Mean Value Theorem there exists $\delta_k \in (0, 1/k)$ such that

$$\epsilon_k = \frac{1}{k} \int_{|x|=\delta_k} \psi(x)ds$$

Therefore

$$\int_{|x|=\delta_k} \psi(x)(x \cdot n)ds = \int_{|x|=\delta_k} \psi(x)\delta_k ds = k\epsilon_k\delta_k \leq \epsilon_k \rightarrow 0.$$

□

Proposition 8 *Let Ω as in Proposition 6, and assume that $\alpha \geq 0, 0 \geq \beta > -N$ and*

$$\frac{N + \alpha}{p + 1} + \frac{N - |\beta|}{q + 1} = N - 2 \tag{23}$$

then there exists no positive strong solution (u, v) of (12).

Proof By combining the previous methods one obtains the result. □

7 An embedding theorem for radial functions

Proposition 9 *Let $\Omega \subset \mathbb{R}^N$ be the ball $\Omega = B_1(0)$. Let $\alpha, \beta > -N$ and let p and q such that $q > \frac{\beta}{N}$ and*

$$\frac{\alpha + N}{p + 1} + \frac{N + \beta}{q + 1} > N - 2. \tag{24}$$

Then the embedding

$$E_{rad} = W_{D,rad}^{2,r}(\Omega, |x|^{-\beta/q} dx) \hookrightarrow L^{p+1}(\Omega, |x|^{\alpha} dx), \quad r = \frac{q + 1}{q}$$

is continuous and compact.

Proof By the density result (Theorem 16 in the Appendix) it is sufficient to prove the assertion for radial $u \in C^\infty(\Omega) \cap E_{rad}$ with $u = 0$ on $\partial\Omega$. For such u we have

$$\Delta u = t^{1-N}(u'(t)t^{N-1})'$$

It is sufficient to prove that there exists a constant C such that

$$\left(\int_0^1 t^\alpha |u(t)|^{p+1} t^{N-1} dt \right)^{\frac{1}{p+1}} \leq C \left(\int_0^1 |(u'(s)s^{N-1})'|^r s^{r-rN+N-1-\beta/q} ds \right)^{1/r} =: C \|u\|_*$$

Set $w(t) = u'(t)t^{N-1}$. Then, since $w(0) = 0$

$$\begin{aligned} |u(t)| &= \left| \int_1^t u'(s)ds \right| = \left| \int_1^t w(s)s^{1-N} ds \right| = \left| \int_t^1 \left[\int_0^s w'(\xi)d\xi \right] s^{1-N} ds \right| \\ &= \left| \int_1^t \left[\int_0^s w'(\xi) \xi^{\frac{r-rN+N-1-\beta/q}{r}} \xi^{-\frac{r-rN+N-1-\beta/q}{r}} d\xi \right] s^{1-N} ds \right| \end{aligned}$$

(by Hölder inequality with exponents $r = 1 + \frac{1}{q}$ and $r' = q + 1$)

$$\begin{aligned} &\leq \int_t^1 \left[\int_0^s |w'(\xi)|^r \xi^{r-rN+N-1-\beta/q} \right]^{1/r} \left[\int_0^s \xi^{-\frac{r-rN+N-1-\beta/q}{r}(q+1)} d\xi \right]^{\frac{1}{q+1}} s^{1-N} ds \\ &\leq \|u\|_* \int_t^1 \left[\int_0^s \xi^{N+\beta-1} d\xi \right]^{\frac{1}{q+1}} s^{1-N} ds \\ &\leq \|u\|_* \int_t^1 s^{\frac{N+\beta}{q+1}+1-N} ds \end{aligned}$$

Now three cases may occur:

Case 1 $\frac{N+\beta}{q+1} > N - 2$ that is $q + 1 < \frac{N+\beta}{N-2}$. In this case we have

$$\int_{\Omega} |x|^\alpha |u(x)|^{p+1} dx \leq C \|u\|_*^{p+1};$$

Case 2 For $q + 1 = \frac{N+\beta}{N-2}$

$$\int_{\Omega} |x|^\alpha |u(x)|^{p+1} dx \leq C \|u\|_* \int_{\Omega} |x|^\alpha |\log(|x|)|^{p+1} dx \leq C \|u\|_*^{p+1},$$

since, for $\alpha > -N$, $|x|^\alpha |\log(|x|)|^{p+1}$ is integrable.

Case 3 Finally for $q + 1 > \frac{N+\beta}{N-2}$

$$\begin{aligned} \int_{\Omega} |x|^\alpha |u(x)|^{p+1} dx &= \omega_{N-1} \int_0^1 t^{\alpha+N-1} |u(t)|^{p+1} dt \\ &\leq C \|u\|_*^{p+1} \int_0^1 t^{\alpha+N-1} t^{(p+1)\left(\frac{N+\beta}{q+1}-N+2\right)} dt \leq C \|u\|_*^{p+1} \end{aligned}$$

for α such that

$$\alpha + N + (p + 1) \left(\frac{N + \beta}{q + 1} - N + 2 \right) > 0,$$

that is

$$\frac{\alpha + N}{p + 1} + \frac{N + \beta}{q + 1} > N - 2.$$

Finally, the proof of the compactness is standard. □

8 Estimates for ground states

8.1 The radial ground state level (β fixed, $\alpha \rightarrow +\infty$)

We give now an estimate from below for the radial level

$$S_{\alpha,\beta}^{rad} = \inf_{u \in E_{rad} \setminus \{0\}} \frac{\int_{\Omega} |x|^{-\beta/q} |\Delta u|^r dx}{\left(\int_{\Omega} |x|^\alpha |u|^{p+1} dx\right)^{\frac{r}{p+1}}}$$

Proposition 10 *There exist $C > 0$ and α_0 such that*

$$S_{\alpha,\beta}^{rad} \geq C \alpha^{2r + \frac{r}{p+1} - 1}, \quad \alpha \geq \alpha_0$$

Proof Let $\varepsilon = \frac{N}{N+\alpha}$ and $u(x) = u(|x|)$ a smooth radial function such that $u = 0$ on $\partial\Omega$. Let $v(\rho) = u(\rho^\varepsilon)$. We have

$$v'(\rho) = \varepsilon u'(\rho^\varepsilon) \rho^{\varepsilon-1} \quad \text{and} \quad v''(\rho) = \varepsilon^2 u''(\rho^\varepsilon) \rho^{2\varepsilon-2} + \varepsilon(\varepsilon - 1) u'(\rho^\varepsilon) \rho^{\varepsilon-2}$$

so that

$$u'(\rho^\varepsilon) = \rho^{1-\varepsilon} \varepsilon^{-1} v'(\rho) \quad \text{and} \quad u''(\rho^\varepsilon) = \varepsilon^{-2} \rho^{2-2\varepsilon} [v''(\rho) - (\varepsilon - 1) \rho^{-1} v'(\rho)]$$

Therefore, by the change of variable $t = \rho^\varepsilon$,

$$\begin{aligned} \int_{\Omega} |x|^{-\beta/q} |\Delta u|^r dx &= \omega_{N-1} \int_0^1 \left| u''(t) + \frac{N-1}{t} u'(t) \right|^r t^{N-1-\beta/q} dt \\ &= \omega_{N-1} \int_0^1 \varepsilon \rho^{\varepsilon N - \frac{\varepsilon\beta}{q} - 1} \left| u''(\rho^\varepsilon) + \frac{N-1}{\rho^\varepsilon} u'(\rho^\varepsilon) \right|^r d\rho \\ &= \omega_{N-1} \int_0^1 \varepsilon \rho^{\varepsilon N - \frac{\varepsilon\beta}{q} - 1} |\varepsilon^{-2} \rho^{2-2\varepsilon} [v''(\rho) - (\varepsilon - 1) \rho^{-1} v'(\rho)] \\ &\quad + \frac{N-1}{\rho^\varepsilon} \varepsilon^{-1} \rho^{1-\varepsilon} v'(\rho)|^r d\rho \\ &= \omega_{N-1} \int_0^1 \varepsilon^{1-2r} \rho^{\varepsilon N - \frac{\varepsilon\beta}{q} - 1} |\rho^{2-2\varepsilon} [v''(\rho) - (\varepsilon - 1) \rho^{-1} v'(\rho)] \\ &\quad + (N-1) \varepsilon \rho^{1-2\varepsilon} v'(\rho)|^r d\rho \\ &= \omega_{N-1} \int_0^1 \varepsilon^{1-2r} \rho^{\varepsilon N - \frac{\varepsilon\beta}{q} - 1 + 2r - 2r\varepsilon} \left| v''(\rho) + \frac{N-2\varepsilon+1}{\rho} v'(\rho) \right|^r d\rho \end{aligned}$$

$$\begin{aligned}
 &= \omega_{N-1} \varepsilon^{1-2r} \int_0^1 \rho^{\varepsilon N - \frac{\varepsilon\beta}{q} - 1 + 2r - 2r\varepsilon} \rho^{-(N-2\varepsilon+1)r} \left| \left(\rho^{N-2\varepsilon+1} v'(\rho) \right)' \right|^r d\rho \\
 &= \omega_{N-1} \varepsilon^{1-2r} \int_0^1 \rho^{\varepsilon N - \frac{\varepsilon\beta}{q} + r - Nr - 1} \left| \left(\rho^{N-2\varepsilon+1} v'(\rho) \right)' \right|^r d\rho \\
 &= \omega_{N-1} \varepsilon^{1-2r} \int_0^1 \rho^{\varepsilon N - \frac{\varepsilon\beta}{q} + r - Nr - 1} \left| (\rho^\gamma v'(\rho))' \right|^r d\rho.
 \end{aligned}$$

where $\gamma = N - 2\varepsilon + 1$. Moreover, by the choice of ε ,

$$\int_\Omega |x|^\alpha |u(x)|^{p+1} dx = \omega_{N-1} \varepsilon \int_0^1 |v(\rho)|^{p+1} \rho^{N-1} d\rho.$$

Thus, we get the following estimate for the radial level:

$$S_{\alpha,\beta}^{rad} = \varepsilon^{-2r - \frac{r}{p+1} + 1} \inf_{v \in E_{rad} \setminus \{0\}} \frac{\int_0^1 \rho^{\varepsilon N - \frac{\varepsilon\beta}{q} + r - Nr - 1} \left| (\rho^\gamma v'(\rho))' \right|^r d\rho}{\left(\int_0^1 |v(\rho)|^{p+1} \rho^{N-1} \right)^{\frac{r}{p+1}}} \tag{25}$$

It is now sufficient to show that there exists $\eta > 0$ such that

$$\inf_{v \in E_{rad} \setminus \{0\}} \frac{\int_0^1 \rho^{\varepsilon N - \frac{\varepsilon\beta}{q} + r - Nr - 1} \left| (\rho^\gamma v'(\rho))' \right|^r d\rho}{\left(\int_0^1 |v(\rho)|^{p+1} \rho^{N-1} \right)^{\frac{r}{p+1}}} \geq \eta > 0 \quad \text{uniformly as } \varepsilon \rightarrow 0$$

We proceed as in the embedding result setting $w(\rho) = v'(\rho)\rho^\gamma$. Then

$$\begin{aligned}
 |v(t)| &= \left| \int_1^t v'(\rho) d\rho \right| = \left| \int_1^t w(\rho) \rho^{-\gamma} d\rho \right| = \left| \int_1^t \rho^{-\gamma} \left(\int_0^\rho w'(s) ds \right) d\rho \right| \\
 &= \left| \int_1^t \rho^{-\gamma} \left[\int_0^\rho w'(s) s^{\frac{\varepsilon N - \frac{\varepsilon\beta}{q} + r - Nr - 1}{r}} s^{-\frac{\varepsilon N - \frac{\varepsilon\beta}{q} + r - Nr - 1}{r}} ds \right] d\rho \right|
 \end{aligned}$$

(Hölder inequality)

$$\begin{aligned}
 &\leq \left| \int_1^t \rho^{-\gamma} \left(\int_0^1 |w'(s)|^r s^{\varepsilon N - \frac{\varepsilon\beta}{q} + r - Nr - 1} ds \right)^{1/r} \left(\int_0^\rho s^{(-\varepsilon N + \frac{\varepsilon\beta}{q} - r + Nr + 1)q} ds \right)^{\frac{1}{q+1}} d\rho \right| \\
 &\leq \left| \int_1^t \rho^{-\gamma} \left(\int_0^1 |w'(s)|^r s^{\varepsilon N - \frac{\varepsilon\beta}{q} + r - Nr - 1} ds \right)^{1/r} \left(\int_0^\rho s^{-\varepsilon Nq + \varepsilon\beta - 1 + N(q+1)} ds \right)^{\frac{1}{q+1}} d\rho \right| \\
 &\leq \left| \int_1^t \rho^{-\gamma} \left(\int_0^1 |w'(s)|^r s^{\varepsilon N - \frac{\varepsilon\beta}{q} + r - Nr - 1} ds \right)^{1/r} \rho^{-\frac{\varepsilon Nq}{q+1} + \frac{\varepsilon\beta}{q+1} + N} d\rho \right| \\
 &= \left(\int_0^1 |w'(s)|^r s^{\varepsilon N - \frac{\varepsilon\beta}{q} + r - Nr - 1} ds \right)^{1/r} \left| \int_1^t \rho^{2\varepsilon - \frac{Nq\varepsilon}{q+1} + \frac{\varepsilon\beta}{q+1} - 1} d\rho \right| \\
 &= \left(\int_0^1 |w'(s)|^r s^{\varepsilon N - \frac{\varepsilon\beta}{q} + r - Nr - 1} ds \right)^{1/r} \left| \int_1^t \rho^{\varepsilon(2-N + \frac{N+\beta}{q+1}) - 1} d\rho \right| =: \mathfrak{S}
 \end{aligned}$$

For $q + 1 \neq \frac{N+\beta}{N-2}$ one has

$$\mathfrak{S} = \left(\int_0^1 |w'(s)|^r s^{\varepsilon N - \frac{\varepsilon\beta}{q} + r - Nr - 1} ds \right)^{1/r} \left| \frac{t^{\varepsilon(2-N + \frac{N+\beta}{q+1})} - 1}{\varepsilon \left(N - 2 - \frac{N+\beta}{q+1} \right)} \right|$$

Therefore

$$\begin{aligned}
 &\left(\int_0^1 |v(t)|^{p+1} t^{N-1} dt \right)^{\frac{r}{p+1}} \\
 &\leq \left(\int_0^1 |w'(s)|^r s^{\varepsilon N - \frac{\varepsilon\beta}{q} + r - Nr - 1} ds \right) \left(\int_0^1 \left| \frac{t^{\varepsilon(2-N + \frac{N+\beta}{q+1})} - 1}{\varepsilon \left(N - 2 - \frac{N+\beta}{q+1} \right)} \right|^{p+1} t^{N-1} dt \right)^{\frac{r}{p+1}}
 \end{aligned}$$

Now we prove that the last term is uniformly bounded as $\varepsilon \rightarrow 0$. Let

$$g_\varepsilon(t) = \left| \frac{t^{\varepsilon(2-N + \frac{N+\beta}{q+1})} - 1}{\varepsilon \left(N - 2 - \frac{N+\beta}{q+1} \right)} \right|^{p+1} t^{N-1}$$

We have that

$$g_\varepsilon(t) \rightarrow (-\log t)^{p+1} t^{N-1} \text{ on } (0, 1), \text{ as } \varepsilon \rightarrow 0$$

and

$$g_\varepsilon(t) \leq (-\log t)^{p+1} t^{N-1} \text{ on } (0, 1).$$

Then, by the Dominated Convergence Theorem

$$\int_0^1 g_\varepsilon(t)dt \rightarrow \int_0^1 (-\log t)^{p+1} t^{N-1} dt$$

which is finite.

The case $q + 1 = \frac{N+\beta}{N-2}$ is easier and left to the reader. This ends the proof. □

8.2 The ground state level

Following the ideas of Smetz, Su and Willem in [26], we give an upper bound for the level

$$S_{\alpha,\beta} = \inf_{W^r(\Omega)\setminus\{0\}} R(u) \tag{26}$$

Proposition 11 *Let p, q as in (4), such that*

$$\frac{N}{p+1} + \frac{N}{q+1} > N - 2. \tag{27}$$

Then there exist $C > 0$ and α_0 such that for $\alpha \geq \alpha_0$

$$S_{\alpha,\beta} \leq C \alpha^{2r-N+N\frac{r}{p+1}} \tag{28}$$

Proof Let ψ a positive smooth function with support in Ω . Let us consider the rescaled function $\psi_\alpha(x) = \psi(\alpha(x - x_\alpha))$, where $x_\alpha = (1 - \frac{1}{\alpha}, 0, \dots, 0)$. Since ψ_α has support in the ball $B(x_\alpha, \frac{1}{\alpha})$, by the change of variable $y = \alpha(x - x_\alpha)$ we obtain for $\beta > 0$

$$\begin{aligned} \int_\Omega |x|^{-\beta/q} |\Delta \psi_\alpha|^r dx &= \int_{B(x_\alpha, \frac{1}{\alpha})} |x|^{-\beta/q} |\Delta \psi_\alpha|^r dx \\ &\leq \alpha^{2r-N} \int_\Omega \left(1 - \frac{2}{\alpha}\right)^{-\beta/q} |\Delta \psi|^r dy, \end{aligned}$$

while for $\beta \leq 0$

$$\begin{aligned} \int_\Omega |x|^{-\beta/q} |\Delta \psi_\alpha|^r dx &= \int_{B(x_\alpha, \frac{1}{\alpha})} |x|^{-\beta/q} |\Delta \psi_\alpha|^r dx \\ &\leq \int_{B(x_\alpha, \frac{1}{\alpha})} |\Delta \psi_\alpha|^r dx = \alpha^{2r-N} \int_\Omega |\Delta \psi|^r dy. \end{aligned}$$

Furthermore,

$$\int_\Omega |x|^\alpha \psi_\alpha^{p+1}(x) dx = \int_{B(x_\alpha, \frac{1}{\alpha})} |x|^\alpha \psi_\alpha^{p+1}(x) dx \geq \left(1 - \frac{2}{\alpha}\right)^\alpha \int_\Omega \alpha^{-N} \psi^{p+1}(y) dy.$$

This implies

$$S_\alpha \leq C \alpha^{2r-N+N\frac{r}{p+1}} \frac{\int_\Omega |\Delta \psi|^r dx}{\left(\int_\Omega \psi^{p+1}(x) dx\right)^{\frac{r}{p+1}}}$$

We remark that $2r - N + N\frac{r}{p+1} > 0$ by (27). □

9 Appendix: Sobolev spaces with A_r weights

9.1 Some definitions

Let $r > 1, \lambda > 0$ be a r -weight, i.e. a function on \mathbb{R}^N such that

$$\lambda > 0 \text{ a.e. on } \mathbb{R}^N, \quad \lambda \quad \text{and} \quad \lambda^{-1/(r-1)} \in L^1_{\text{loc}}(\mathbb{R}^N)$$

Let $\Omega \in \mathbb{R}^N$ a bounded smooth domain. We denote with $L^r(\Omega, \lambda)$ the set of functions $u \in L^1_{\text{loc}}(\Omega)$ such that

$$\int_{\Omega} |u|^r \lambda \, dx < +\infty$$

and with $W^{2,r}(\Omega, \lambda)$ the set of functions $u \in W^{2,1}_{\text{loc}}(\Omega)$ such that

$$\int_{\Omega} (|u|^r + |\nabla u|^r + \sum_{|\xi|=2} |D^{\xi} u|^r \lambda) \, dx < +\infty$$

One can easily prove that endowed with the norm

$$\|u\|_{W^{2,r}(\Omega,\lambda)} := \left(\int_{\Omega} (|u|^r + |\nabla u|^r + \sum_{|\xi|=2} |D^{\xi} u|^r \lambda) \, dx \right)^{1/r}$$

$W^{2,r}(\Omega, \lambda)$ is a Banach Space. We also denote with $\tilde{W}^{2,r}_0(\Omega, \lambda)$ the closure of $\{\phi \in C^{\infty}(\Omega) : \phi = 0 \text{ on } \partial\Omega\}$ in $W^{2,r}(\Omega, \lambda)$. We are interested to study some “density” property of these Sobolev spaces with weights. To this aim we introduce the following

Definition 12 (*Muckenhoupt Class A_r*) We say that a r -weight λ on \mathbb{R}^N is in the Muckenhoupt class A_r if

$$\left(\frac{1}{|B|} \int_B \lambda \, dx \right) \left(\frac{1}{|B|} \int_B \lambda^{-1/(r-1)} \, dx \right)^{r-1} \leq C \tag{29}$$

for every ball B contained in \mathbb{R}^N (here $|B|$ denotes the Lebesgue measure of the ball B).

Example 13 $\lambda(x) = |x|^{\gamma} \in A_r$ iff $-N < \gamma < N(r - 1)$

This class of weights is strictly related to the Hardy-Littlewood maximal function

Definition 14 Let $f \in L^1_{\text{loc}}(\mathbb{R}^N)$. The maximal function of f is defined by

$$(Mf)(x) = \sup_{R>0} \frac{1}{|B_R(x)|} \int_{B_R(x)} |f(y)| \, dy. \tag{30}$$

In fact the “condition” A_r was introduced by B. Muckenhoupt in the following Theorem

Theorem 15 (Muckenhoupt [22]) *Let λ be a r -weight. The following conditions are equivalent:*

(i) *there is a constant C such that*

$$\int_{\mathbb{R}^N} [(Mf)]^r \lambda \, dx \leq C \int_{\mathbb{R}^N} |f|^r \lambda \, dx \quad \forall f \in L^r(\mathbb{R}^N) \tag{31}$$

(ii) $\lambda \in A_r$

9.2 Approximation by smooth functions on Ω

The central part of this section is to prove the following extension of the celebrated Meyer-Serrin result

Theorem 16

$$\tilde{W}_0^{2,r}(\Omega, \lambda) = W^{2,r}(\Omega, \lambda) \cap W_0^{1,r}(\Omega)$$

In order to prove Theorem 16 we need some preliminary results

Theorem 17 *Let $\lambda \in L^1(\Omega)$ a positive measure on Ω . Then $C_0(\Omega)$ is dense in $L^r(\Omega, \lambda)$ ($1 \leq r < +\infty$).*

Proof (Theorem 2.19 in [1]) It is sufficient to prove that for every $\varepsilon > 0$ and a nonnegative function u there exists $\varphi \in C_0(\Omega)$ such that $\|u - \varphi\|_{L^r(\Omega, \lambda)} < \varepsilon$.

For u measurable and nonnegative there exists a monotonically increasing sequence $\{s_n\}$ of nonnegative simple functions converging point-wise to u on Ω and strongly in $L^r(\Omega, \lambda)$ (since $0 \leq s_n(x) \leq u(x)$, we have $s_n \in L^r(\Omega, \lambda)$ and $(u(x) - s_n(x))^r \lambda(x) \leq u(x)^r \lambda(x)$, so that by the Dominated Convergence theorem $s_n \rightarrow u$ in $L^r(\Omega, \lambda)$). Thus there exists $s \in \{s_n\}$ such that $\|u - s\|_{L^r(\Omega, \lambda)} < \varepsilon/2$. By Lusin’s theorem there exists for all $\delta > 0$ a $\varphi \in C_0(\mathbb{R}^N)$ such that

$$|\varphi(x)| \leq \|s\|_\infty$$

and

$$\text{Vol} E < \delta, \quad E = \{x \in \mathbb{R}^N : \varphi(x) \neq s(x)\}.$$

Therefore, by the absolute continuity of the integral, we can choose $\delta = \delta(\varepsilon)$ such that

$$\|s - \varphi\|_{L^r(\Omega, \lambda)} \leq \|s - \varphi\|_\infty \left(\int_E \lambda(x) \, dx \right)^{1/r} < \varepsilon/2$$

□

Lemma 18 [25] *Let J be a nonnegative, real-valued function in $C_0^\infty(\mathbb{R}^N)$ with the following properties*

$$J(x) = 0 \text{ if } |x| \geq 1, \quad \text{and} \quad \int_{\mathbb{R}^N} J(x) = 1.$$

We consider the sequence of “mollifiers” $J_\epsilon(x) = \epsilon^{-N} J(x/\epsilon)$. Then

(i) $J_\epsilon(x) = 0$ if $|x| \geq 1$

(ii) *There exists a positive constant $C = C(N, \sup J)$ such that, if*

$$J_\epsilon * u(x) = \int_{\mathbb{R}^N} J_\epsilon(x - y)u(y)dy ,$$

then

$$|J_\epsilon * u(x)| \leq CM(u)(x), \quad \forall u \in L^1_{\text{loc}}(\mathbb{R}^N).$$

Proof From (i) and the definition of maximal function one has

$$|J_\epsilon * u(x)| \leq \frac{\sup J}{\epsilon^N} \int_{B_\epsilon(x)} |u(y)|dy \leq C \frac{1}{|B_\epsilon(x)|} \int_{B_\epsilon(x)} |u(y)|dy \leq CM(u)(x)$$

□

Theorem 19 ([1]) *Let u be a function which is defined on \mathbb{R}^N and vanishes identically outside Ω . Let λ a r -weight ($1 \leq r < +\infty$) belonging to the Muckenhoupt class A_r .*

- (a) *If $u \in L^1_{\text{loc}}(\mathbb{R}^N)$, then $J_\epsilon * u(x) \in C^\infty(\mathbb{R}^N)$*
- (b) *If $u \in L^1_{\text{loc}}(\Omega)$ and $\text{supp}(u) \subset\subset \Omega$, then, for $\epsilon < \text{dist}(\text{supp}(u), \partial\Omega)$, $J_\epsilon * u(x) \in C^\infty_0(\Omega)$*
- (c) *If $u \in L^r(\Omega, \lambda)$, then $J_\epsilon * u(x) \in L^r(\Omega, \lambda)$. Moreover there exists a positive constant $C = C(N, \sup J)$ such that*

$$\|J_\epsilon * u\|_{L^r(\Omega, \lambda)} \leq C \|u\|_{L^r(\Omega, \lambda)}$$

(d) *If $u \in L^r(\Omega, \lambda)$, then*

$$\|J_\epsilon * u - u\|_{L^r(\Omega, \lambda)} \rightarrow 0, \quad \epsilon \rightarrow 0^+$$

Proof (For (a) and (b) see [1] Theorem 2.29).

If $u \in L^r(\Omega, \lambda)$ then ($u \in L^1_{\text{loc}}(\mathbb{R}^N)$) from Lemma 18 we have

$$|J_\epsilon * u(x)| \leq CM(u)(x).$$

Hence since λ is in the Muckenhoupt class, by (31) (Theorem 15)

$$\int_{\Omega} |J_\epsilon * u(x)|^r \lambda(x)dx \leq C \int_{\Omega} |M(u)|^r(x)\lambda(x) dx \leq C_1 \int_{\Omega} |u(x)|^r \lambda(x) dx.$$

In particular $\|J_\epsilon * u\|_{L^r(\Omega, \lambda)} \leq C \|u\|_{L^r(\Omega, \lambda)}$ (here $C = C(N, \sup J)$). Now, let $\eta > 0$ be given. By Theorem 17 there exists $\varphi \in C^0_0(\Omega)$ such that $\|u - \varphi\|_{L^r(\Omega, \lambda)} < \frac{\eta}{2(C+1)}$.

Now, since $\int_{\mathbb{R}^N} J_\epsilon(y)dy = 1$, by the uniform continuity of φ there exists ϵ_0 such that for all $0 < \epsilon < \epsilon_0$

$$\begin{aligned} |J_\epsilon * \varphi(x) - \varphi(x)| &= \left| \int_{\mathbb{R}^N} J_\epsilon(x - y)(\varphi(y) - \varphi(x))dy \right| \\ &\leq \sup_{|y-x|<\epsilon} |\varphi(y) - \varphi(x)| < \frac{\eta}{2(\int_{\Omega} \lambda(x)dx)^{1/r}} \end{aligned}$$

This is sufficient to obtain

$$\|J_\epsilon * \varphi - \varphi\|_{L^r(\Omega, \lambda)} < \eta/2$$

Finally, from (c) one has

$$\begin{aligned} \|J_\epsilon * u - u\|_{L^r(\Omega, \lambda)} &\leq \|J_\epsilon * u - J_\epsilon * \varphi\|_{L^r(\Omega, \lambda)} + \|J_\epsilon * \varphi - \varphi\|_{L^r(\Omega, \lambda)} + \|u - \varphi\|_{L^r(\Omega, \lambda)} \\ &\leq (1 + C)\|u - \varphi\|_{L^r(\Omega, \lambda)} + \|J_\epsilon * \varphi - \varphi\|_{L^r(\Omega, \lambda)} < \eta \end{aligned}$$

□

Lemma 20 *Let $u \in W^{2,r}(\Omega, \lambda)$. If $\Omega' \subset\subset \Omega$, then $J_\epsilon * u \rightarrow u$ in $W^{2,r}(\Omega', \lambda)$*

Proof Let $\epsilon < \text{dist}(\Omega', \partial\Omega)$, then $D^\alpha J_\epsilon * u = J_\epsilon * D^\alpha u$ in the distributional sense in Ω' (see [1], lemma 3.16). Since $Du \in L^r(\Omega)$, and $D^\alpha u \in L^r(\Omega, \lambda)$ for $|\alpha| = 2$ we have, by Theorem 19(c)

$$\|D^\alpha J_\epsilon * u - D^\alpha u\|_{L^r(\Omega', \lambda)} = \|J_\epsilon * D^\alpha u - D^\alpha u\|_{L^r(\Omega', \lambda)} \rightarrow 0 \text{ as } \epsilon \rightarrow 0^+$$

and

$$\|DJ_\epsilon * u - Du\|_{L^r(\Omega')} = \|J_\epsilon * Du - u\|_{L^r(\Omega')} \rightarrow 0 \text{ as } \epsilon \rightarrow 0^+$$

so that

$$\|J_\epsilon * u - u\|_{W^{2,r}(\Omega', \lambda)} \rightarrow 0 \text{ as } \epsilon \rightarrow 0^+ \tag{32}$$

□

Now we can prove Theorem 16

Proof (we simplify the proof considering the case $\Omega = B_1(0)$) (see also [1] and [13]).

If $u \in W^{2,r}(\Omega, \lambda)$ and $\epsilon > 0$ we prove that there exists $\varphi \in C^\infty(\Omega)$ such that $\|\varphi - u\|_{W_0^{2,r}(\Omega', \lambda)} < \epsilon$. For $k = 1, 2, \dots$ let

$$\Omega_k = \left\{ x \in \Omega : |x| < 1 - \frac{1}{k} \right\}, \quad \Omega_0 = \Omega_{-1} = \emptyset$$

and

$$U_1 = \Omega_2, \quad U_k = \left\{ x : \frac{k-2}{k-1} < |x| < \frac{k}{k+1} \right\}, \quad k = 2, \dots$$

Then

$$\mathcal{O} = \{U_k : k = 1, 2, \dots\}$$

is a collection of open subsets of Ω that covers Ω . Let Ψ be a C^∞ -partition of unity for Ω subordinate to \mathcal{O} . Let ψ_k denote the sum of the finitely many functions $\psi \in \Psi$ whose support are contained in U_k . Then $\psi_k \in C_0^\infty(U_k)$ and $\sum_{k=1}^\infty \psi_k(x) = 1, \forall x \in \Omega$.

Let $0 < \epsilon < \frac{1}{(k+1)(k+2)}$ and $V_k = \{x : \frac{k-3}{k-2} \leq |x| < \frac{k+1}{k+2}, k = 3, \dots\}$, $V_1 = \Omega_3$, $V_2 = \Omega_4$. Then $\text{supp}(J_\epsilon * (\psi_k u)) \subset V_k \subset\subset \Omega$. Since $\psi_k u \in W^{2,r}(\Omega, \lambda)$, by Lemma 20 we may choose $0 < \epsilon_k < \frac{1}{(k+1)(k+2)}$ such that

$$\|J_\epsilon * (\psi_k u) - (\psi_k u)\|_{W^{2,r}(\Omega, \lambda)} = \|J_\epsilon * (\psi_k u) - (\psi_k u)\|_{W^{2,r}(V_k, \lambda)} < \frac{\epsilon}{2k}.$$

Let $\varphi = \sum_{j=1}^{+\infty} J_{\epsilon_k} * (\psi_j u)$. Since on any $\Omega' \subset\subset \Omega$ only finitely many terms in the sum can be nonzero, one has $\varphi \in C^\infty(\Omega)$ and $\varphi = 0$ on $\partial\Omega$. For $x \in \Omega_k$, we have

$$u(x) = \sum_{j=1}^{k+2} \psi_j(x)u(x) \quad \text{and} \quad \varphi(x) = \sum_{j=1}^{k+2} J_{\epsilon_k} * (\psi_j u)(x)$$

Thus

$$\|u - \varphi\|_{W^{2,r}(\Omega_{k,\lambda})} \leq \sum_{j=1}^{k+2} \|J_{\epsilon_k} * (\psi_j u) - \psi_j u\|_{W_0^{2,r}(\Omega,\lambda)} < \varepsilon$$

By the monotone convergence theorem $\|u - \varphi\|_{W_0^{2,r}(\Omega,\lambda)} < \varepsilon$. \square

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