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Comment on "New exact solution of Euler's equations (rigid body dynamics) in the case of rotation over the fixed point"

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Abstract The analysis of the paper "New exact solution of Euler's equations (rigid body dynamics) in the case of rotation over the fixed point" is made. It is shown that the solution is reduced to a well-known subcase of the Euler's case: steady rotation.

Keywords Euler's equations (rigid body dynamics) · Steady rotation

1 The consideration of the problem in the absolute system of coordinates

In a recent article [\[1\]](#page-2-0) Ershkov reported a new exact solution of Euler's equation. Thorough analysis of this work showed that the proposed "generalization" of the Euler's case is just a special case of the Euler's case. To verify this, let us consider the problem in the absolute system of coordinates *OXYZ*. Denote by **K***OXYZ* the components of the angular momentum vector with respect to the fixed point of the rigid body *O*, the gravity force by P_{OXYZ} , $\gamma_{OXYZ} = -P_{OXYZ}/|P_{OXYZ}| = -P_{OXYZ}/P$, the point of gravity force application r_{OXYZ} .

The main assumption of the article [\[1\]](#page-2-0) is the integral of motion (2.1) $(\mathbf{K}_{OXYZ})^2 = C_0^2$. The equality (2.1), the first equation in (1.3) $(\gamma_{OXYZ})^2 = 1$ and the second equation in (1.3) $\mathbf{K}_{OXYZ} \cdot \gamma_{OXYZ} = C_0 = \text{const}$ imply the colinearity of the vectors $\mathbf{K}_{OXYZ} \|\gamma_{OXYZ}$ (this holds in any axes set, in particular, in the frame of reference fixed in the rotating body (2.2)). This fact follows from:

$$
(\mathbf{K}_{OXYZ})^2 = (\mathbf{K}_{OXYZ} \cdot \gamma_{OXYZ})^2 + (\mathbf{K}_{OXYZ} \times \gamma_{OXYZ})^2.
$$
 (1)

So we obtain that

$$
\mathbf{K}_{OXYZ} = C_0 \gamma_{OXYZ} = -C_0 \mathbf{P}_{OXYZ} / P = \text{const}_{OXYZ}.
$$
 (2)

By the angular momentum theorem [\[2\]](#page-2-1), we have that

$$
\frac{\mathrm{d}K_{OXYZ}}{\mathrm{d}t} = \mathbf{r}_{OXYZ} \times \mathbf{P}_{OXYZ} = 0,\tag{3}
$$

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and, therefore,

$$
\mathbf{r}_{OXYZ} \|\mathbf{P}_{OXYZ}\| \mathbf{K}_{OXYZ},\tag{4}
$$

i.e. the momentum of external forces is identically zero. Therefore, the cases considered in the paper (paragraphs 3 and 4) cannot be beyond the classical Euler's case. Let us show that these cases are just subcases of the Euler's case.

As we consider a rigid body $|\mathbf{r}_{OXYZ}|$ = const and taking into account [\(4\)](#page-1-0), we obtain that

$$
\mathbf{r}_{OXYZ} = \mathbf{const}_{OXYZ},\tag{5}
$$

which may only hold in the following two cases

- 1. $\mathbf{r}_{OXYZ} = 0$ —this is a trivial case,
- 2. $r_{OXYZ} \neq 0$, but the motion is steady rotation about one of the principle axes, and the gravity center lies on this axis and the gravity force vector is parallel to this axis. This case is easily understood using the Poinsot geometrical interpretation [\[2](#page-2-1)[,3\]](#page-2-2).

2 The consideration of the problem in a frame of reference fixed in the rotating body

Now, following the author's paper, we deduce equation [\(4\)](#page-1-0) in the frame of reference fixed in the rotating body *Oxyz*, preserving the notations of the variables. (2.2) implies that $\mathbf{K}_{Oxyz} \|\mathbf{P}_{Oxyz}$. Substituting (2.2) in (1.1) and considering (1.2) (1.1) and (1.2) and (1.2) and the case of consequenting and in constitution in (1.1) and considering (1.2) $((1.1)$ and (1.2) are the same and agree with corresponding equations in [\[3](#page-2-2)]), we get that

$$
\begin{cases}\n0 = P (\gamma_2 c - \gamma_3 b) \\
0 = P (\gamma_3 a - \gamma_1 c) \\
0 = P (\gamma_1 b - \gamma_2 a)\n\end{cases}
$$
\n(6)

that is $\mathbf{r}_{Oxyz} \times \mathbf{P}_{Oxyz} = 0$, which implies [\(4\)](#page-1-0) in a moving set of axes

$$
\mathbf{r}_{Oxyz} \|\mathbf{P}_{Oxyz}\| \mathbf{K}_{Oxyz},\tag{7}
$$

which we needed to deduce.

3 The hidden triviality of paragraph 3 in paper [\[1\]](#page-2-0)

Let us show that proposition (3.1) in the case $a = 0$ in paragraph 3 implies that the rigid body is at rest. Indeed, in the case $a = 0$ statement [\(7\)](#page-1-1) corresponds to: firstly,

$$
\Omega_1 = 0 \tag{8}
$$

and, secondly,

$$
\frac{I_2\Omega_2}{b} = \frac{I_3\Omega_3}{c}.
$$
\n(9)

Substituting (8) in the first equation (1.1) and considering the first equation in (6) , we obtain that

$$
(I_3 - I_2) \Omega_2 \Omega_3 = 0. \tag{10}
$$

Considering the author's inequalities $I_2 > I_1 > I_3$, $c \neq 0$ and (3.1), from [\(9\)](#page-1-4) we obtain that $\Omega_2 = \Omega_3 = 0$, thus completing our proof.

4 The stationary rotation in paragraph 4 in paper [\[1\]](#page-2-0)

Let us show that the case $a = c = 0$ of paragraph 4 implies that $\Omega_1 = \Omega_3 = 0$, $\Omega_2 = const$.

Indeed, the equalities $a = c = 0$ and [\(7\)](#page-1-1) imply that $\Omega_1 = \Omega_3 = 0$. Considering also (2.1) we obtain that Ω_2 = const, thus completing the proof.

Thus, in this case, the body is rotating about the principle axis with constant angular speed $\Omega_2 = \text{const}$ (steady rotation). In equation (4.3), $f_1 = f_2 = 0$, and thus in this equation, we have zero divided by zero.

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