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## Oxidation during metasomatism in ultramafic xenoliths from the Wesselton kimberlite, South Africa: implications for the survival of diamond

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**Abstract** Garnets in xenoliths from the Wesselton kimberlite show significant zoning in major and trace elements. The garnets were studied using room temperature Mössbauer spectroscopy with high spatial resolution, and show an increase in  $\text{Fe}^{3+}/\Sigma\text{Fe}$  from core to secondary rim. Temperatures and pressures were determined using the garnet–olivine, garnet–orthopyroxene and Ni in garnet formulations, and indicate conditions close to 1,000 °C and 37 kbar for most of the garnets. Oxygen fugacities calculated using the garnet–olivine–orthopyroxene oxybarometer show an increase of approximately one log-bar unit from garnet core to secondary rim, relative to the quartz–fayalite–magnetite buffer curve. Combined with reanalysis of literature data from unaltered material from the same locality, there was an increase in relative oxygen fugacity of approximately two log-bar units during the course of metasomatism. Existing data from other South African garnet peridotites were recalculated using the same thermobarometers and oxybarometers, and indicate relative oxygen fugacities that lie at least two log-bar units below the diamond–carbonate equilibrium in peridotitic systems, which defines the maximum limit of diamond

stability in peridotite. Diamond would hence be preserved during the initial stages of metasomatism, but in later stages fluid would react with the diamonds, leading to their resorption and eventual destruction.

### Introduction

The study of zoning profiles in minerals can provide information regarding chemical processes that occurred during the formation of a mineral assemblage. The electron, proton and ion microprobes provide information on major and trace element chemical compositions, but, except in a small number of cases, provide no information on iron oxidation state (e.g. Canil and O'Neill 1996; McCammon 1999; Sobolev et al. 1999). Such information can be used in conjunction with oxygen barometers to determine oxygen fugacity during mineral formation, which can have significant effects on many processes including element partitioning, diffusivity, partial melting reactions and diamond stability. Mössbauer spectroscopy is an excellent method to distinguish between  $\text{Fe}^{2+}$  and  $\text{Fe}^{3+}$ , and can provide an estimate of their relative abundance in individual phases. Normally, experiments are performed on polycrystalline samples of ~1 cm diameter, but a simple technique (McCammon et al. 1991; McCammon 1994) enables routine Mössbauer measurements to be performed on absorbers more than two orders of magnitude smaller.

Zoned garnets from phlogopite–harzburgite xenoliths in the Wesselton kimberlite, Kimberly, South Africa, were studied by Griffin et al. (1999) using electron and proton microprobes, and the data were used to infer the history of a multistage metasomatic alteration. Modelling of zoning profiles suggested a three-part process involving infiltration of fluids with different compositions over several different time-scales:

1. Slow diffusion of Ca, Zr and Y from rim to core over periods of 10–30,000 years, from a fluid depleted in Ti, Ga and Y.

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2. Formation of garnet overgrowths high in Ca, Zr, Y and Ti, followed by annealing over periods of  $\sim 10^3$  years.
3. Formation of secondary rims on very short time-scales prior to eruption.

Preliminary work suggested differences in  $\text{Fe}^{3+}$  concentration between the various garnet generations (McCammon et al. 1995), implying a change in oxygen fugacity conditions.

In this paper we extend the work described by McCammon et al. (1995) and report results of a Mössbauer study of zoned garnets from the Wesselton kimberlite, and calculate oxygen fugacity using the garnet–olivine–orthopyroxene oxybarometer (Luth et al. 1990; Gudmundsson and Wood 1995). Results are used to infer redox conditions during the different processes recorded in the zoning patterns and their implications for diamond stability.

## Experimental procedure

Mineral fragments from ultramafic xenoliths in coarse concentrate from the Wesselton kimberlite were mounted in epoxy disks and polished for preliminary analysis using the electron and proton microprobes. Garnets showing significant heterogeneity were prepared as polished sections to determine zoning patterns using the electron microprobe at CSIRO and Macquarie University (Australia) and the proton microprobe at CSIRO (Australia). Coexisting olivine and orthopyroxene were found to be homogeneous within each grain. X-ray maps of the garnets were prepared and used to locate interesting areas for Mössbauer analysis. Descriptions of the four garnets selected for this study are listed in Table 1, and chemical compositions are listed in Griffin et al. (1999).

Samples were prepared for Mössbauer spectroscopy by grinding the disks to a thickness of  $\sim 300$   $\mu\text{m}$ , which gives an iron density of  $\sim 5$   $\text{mg Fe/cm}^2$  based on the chemical composition of the garnet. Note that because Mössbauer spectroscopy is a non-destructive technique, the top polished surface of the disks are preserved for further analysis. To isolate different parts of the sample, a piece of 25- $\mu\text{m}$  thick Ta foil (absorbs 99.9% of 14.4 keV gamma rays) drilled with a 400- $\mu\text{m}$  diameter hole was positioned over the area to be analysed. Ca X-ray maps showing the locations where Mössbauer spectra were collected are illustrated in Fig. 1.

A simple modification to a conventional Mössbauer spectrometer allows spectra of absorbers nearly two orders of magnitude smaller diameter to be collected (McCammon et al. 1991; McCammon 1994). To obtain adequate count rates, a conventional Mössbauer source (typical specific activity 100  $\text{mCi/cm}^2$ ) is

replaced by a point source (specific activity  $\geq 2,000$   $\text{mCi/cm}^2$ ), which can be obtained commercially at a cost similar to conventional sources. The gamma rays are collimated to the selected absorber diameter using a Pb shield, and the source-absorber distance is reduced to  $< 5$  mm. The latter results in a solid angle similar to conventional experiments, and hence a similar count rate. Because the signal quality depends on absorber density (measured in  $\text{mg Fe/cm}^2$ ) and not the total amount of iron in the absorber, the reduction in size has no effect on the effective thickness of the absorber. When non-resonant absorption caused by heavier elements is low and the point source is relatively new ( $< 1$  year old), high quality Mössbauer spectra (comparable to conventional measurements) can be recorded on absorbers with diameters as small as 50  $\mu\text{m}$ .

Mössbauer spectra were recorded at room temperature in transmission mode using a constant acceleration Mössbauer spectrometer with a nominal 20  $\text{mCi } ^{57}\text{Co}$  high specific activity source (2  $\text{Ci/cm}^2$ ) in a 12- $\mu\text{m}$  Rh matrix. Spectra were collected over time periods ranging from 1–5 days, depending on the signal quality. The velocity scale was calibrated relative to 25  $\mu\text{m}$   $\alpha$ -Fe foil using the positions certified for National Bureau of Standards standard reference material no. 1541; line widths of 0.42 mm/s for the outer lines of  $\alpha$ -Fe were obtained at room temperature. The spectra were fitted to Lorentzian and Voigt line-shapes using the commercially available fitting program NORMOS written by R.A. Brandt (distributed by Wissenschaftliche Elektronik GmbH, Germany).

## Results from Mössbauer spectroscopy

The Mössbauer spectra of garnet were fitted to two quadrupole doublets, corresponding to  $\text{Fe}^{2+}$  and  $\text{Fe}^{3+}$ . Spectra were initially fit to Lorentzian doublets with conventional constraints (equal component widths and areas), but residuals showed large deviations. Final fits were therefore obtained using Voigt doublets for  $\text{Fe}^{2+}$  absorption, where areas and Lorentzian line-widths of individual components were allowed to vary, and Lorentzian doublets with conventional constraints for  $\text{Fe}^{3+}$  absorption. Area asymmetry reflects recoil-free fraction anisotropy (Gol'danskii–Karyagin effect; Geiger et al. 1992), whereas line-width asymmetry probably reflects differences in next-nearest neighbour configurations (e.g. Angel et al. 1998). Although line-width and area asymmetry parameters are correlated and cannot be determined unambiguously, the high resolution of  $\text{Fe}^{2+}$  and  $\text{Fe}^{3+}$  absorption enables robust values for other hyperfine parameters (centre shift, quadrupole splitting and relative area) to be determined.

Hyperfine parameters for  $\text{Fe}^{2+}$  (Table 2) correspond to those reported by Amthauer et al. (1976) for the

**Table 1** Description of zoned garnets studied using Mössbauer spectroscopy

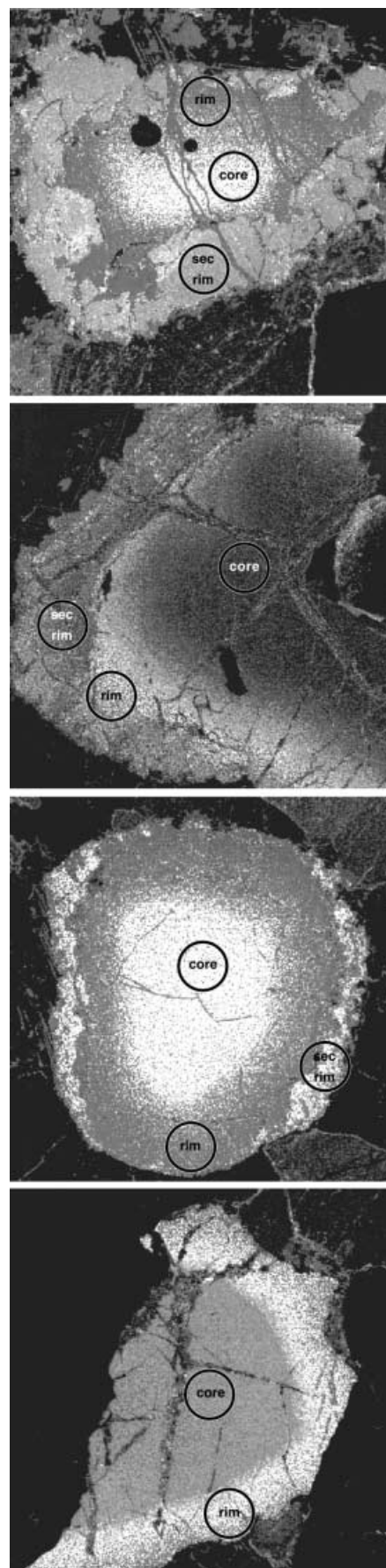
Sample	Description
933	A single grain $\sim 1.6 \times 2.5$ mm, surrounded by olivine and orthopyroxene. Composition varies somewhat irregularly from core to rim, and the grain is surrounded by a secondary garnet rim that varies from 100–150 $\mu\text{m}$ in thickness
937	A single ovoid grain, of which an area $\sim 1.7$ mm across is available for analysis. It is surrounded by olivine. There is pronounced zoning from core to rim, and a secondary replacement rim of up to 500 $\mu\text{m}$ in thickness
940	A single grain $\sim 3.5$ mm in the long diameter, surrounded by olivine and orthopyroxene. The primary garnet grain shows broad symmetrical zoning from core to rim, and on one side there is a 50–75 $\mu\text{m}$ zone of secondary replacement garnet
951	An incomplete grain $\sim 3$ mm in diameter, surrounded by olivine and orthopyroxene. The relatively homogeneous core is surrounded by a rim of distinctly different composition

dodecahedral site, and agree within experimental error with values reported by Luth et al. (1990) for mantle-derived garnets of similar composition. Hyperfine parameters of  $\text{Fe}^{3+}$  correspond to those reported in the above references for the octahedral site, and there is no evidence for tetrahedrally coordinated  $\text{Fe}^{3+}$ . Line-widths for both  $\text{Fe}^{2+}$  and  $\text{Fe}^{3+}$  doublets are somewhat greater than values commonly observed for mantle garnets, which is likely related to microstresses in the single crystals and would normally be removed during grinding. The increased line-widths do not have a significant effect on the values determined for  $\text{Fe}^{3+}/\Sigma\text{Fe}$ .

$\text{Fe}^{3+}$  contributions are well resolved in all spectra consisting primarily of garnet, and relative areas could be determined with high precision. De Grave and Van Alboom (1991) have shown that recoil-free fractions of sites containing  $\text{Fe}^{3+}$  tend to be larger than for sites containing  $\text{Fe}^{2+}$ , which agrees with observations for garnet (e.g. Lyubutin et al. 1970; Lyubutin and Dodokin 1971; Amthauer et al. 1976). Accordingly, we corrected the relative areas of  $\text{Fe}^{3+}$  components to determine  $\text{Fe}^{3+}/\Sigma\text{Fe}$  as described by Rancourt et al. (1994) using values for the Mössbauer Debye temperatures of  $\Theta_{\text{M}}(\text{Fe}^{2+}) = 340$  K and  $\Theta_{\text{M}}(\text{Fe}^{3+}) = 500$  K, which are similar to values reported by De Grave and Van Alboom (1991) for other silicates. We did not correct the relative areas for thickness effects because these are minimal considering the large intrinsic line-widths combined with the small effective thicknesses of the garnets (e.g. Rancourt et al. 1993).  $\text{Fe}^{3+}/\Sigma\text{Fe}$  values vary in a consistent manner in all garnets studied, where values increase from core to secondary rim (Table 2). Figure 2 illustrates Mössbauer spectra taken from different regions of the garnets.

We collected Mössbauer spectra for minerals coexisting with garnet (Fig. 3). Mössbauer spectra for olivine give hyperfine values that agree with those in the literature (e.g. Shinno 1981) and show no evidence for  $\text{Fe}^{3+}$  within the detection limit ( $\sim 2\%$   $\text{Fe}^{3+}/\Sigma\text{Fe}$ ). Mössbauer spectra for orthopyroxene show unequal component areas because of preferred orientation (the absorbers are single crystals) and the non-cubic point group symmetry of the iron sites. Hyperfine parameters indicate that most  $\text{Fe}^{2+}$  is partitioned into the M2 site, and values are consistent with those observed for other mantle orthopyroxenes of similar composition (e.g. Luth and Canil 1993). Values for  $\text{Fe}^{3+}/\Sigma\text{Fe}$  in olivine and orthopyroxene are consistent with results from other Mössbauer studies of mantle xenoliths (e.g. Canil and O'Neill 1996).

Several spectra collected from garnet rims showed additional absorption (Fig. 3, bottom), and compari-



**Fig. 1** Electron microprobe X-ray maps of garnet grains indicating distribution of Ca (data from Griffin et al. 1999). Samples are ordered in the sequence (top to bottom): 933; 937; 940; 951. Transmission Mössbauer spectra were recorded of the regions marked by circles (400  $\mu\text{m}$  diameter), where labels correspond to those given in Table 2

**Table 2** Room temperature hyperfine parameters of garnets and coexisting minerals. *CS* Centre shift (relative to  $\alpha$ -Fe); *QS* quadrupole splitting;  $\Gamma$  Lorentzian full width at half maximum of low-velocity component;  $\sigma$  Gaussian standard deviation;  $A_{21}$  component area asymmetry;  $\Gamma_{21}$  component line-width asymmetry.

Estimated standard deviations for garnet spectra are  $\pm 0.005$  mm/s (*CS* and *QS*);  $\pm 0.01$  mm/s ( $\Gamma$  and  $\sigma$ );  $\pm 0.05$  ( $A_{21}$  and  $\Gamma_{21}$ ). Estimated standard deviations for opx and ol spectra are  $\pm 0.01$  mm/s (*CS* and *QS*);  $\pm 0.01$  mm/s ( $\Gamma$ ). Parameter values in italics were held fixed during the fitting process

Sample	Fe <sup>2+</sup>						Fe <sup>3+</sup>			Phase	
	CS (mm/s)	QS (mm/s)	$\Gamma$ (mm/s)	$\sigma$ (mm/s)	$A_{21}$	$\Gamma_{21}$	CS (mm/s)	QS (mm/s)	$\Gamma$ (mm/s)	Relative area (%)	Fe <sup>3+</sup> / $\Sigma$ Fe <sup>a</sup> (%)
Garnet 933											
Core	1.288	3.555	0.36	0.1	0.98	0.87	0.347	0.337	0.31	100	6 (2)
Rim	1.321	3.583	0.48	–	<i>0.95</i>	0.81	<i>0.354</i>	0.373	<i>0.45</i>	55	12 (9)
Sec rim	1.292	3.563	0.32	0.18	0.93	0.79	0.385	0.000	0.61	95	10 (3)
Garnet 937											
Core	1.292	3.557	0.3	0.27	0.95	0.77	<i>0.354</i>	0.262	<i>0.50</i>	100	5 (2)
Rim	1.284	3.522	0.22	0.27	0.96	0.98	<i>0.354</i>	0.365	<i>0.50</i>	92	9 (3)
Sec rim	1.316	3.584	0.47	–	<i>0.95</i>	0.82	<i>0.354</i>	0.346	<i>0.45</i>	53	10 (7)
Garnet 940											
Core	1.309	3.606	0.41	0.07	0.96	0.82	0.344	0.247	0.46	100	5 (2)
Rim	1.314	3.585	0.47	0.06	0.89	0.78	0.404	0.265	0.31	98	4 (2)
Sec rim	1.321	3.587	0.46	–	0.95	0.83	<i>0.354</i>	0.340	<i>0.45</i>	44	7 (6)
Garnet 951											
Core	1.285	3.578	0.39	0.14	0.99	0.82	<i>0.354</i>	0.405	0.48	100	5 (2)
Rim	1.300	3.541	0.4	0.15	0.95	0.76	<i>0.354</i>	0.000	0.69	91	6 (3)
Coexisting minerals											
Orthopyroxene 1.15 933		2.13	0.45							100	0 (2)
Orthopyroxene 1.13 951		2.13	0.41				0.24	0.82	<i>0.41</i>	100	2 (2)
Olivine 937	1.14 <sup>b</sup>	2.98 <sup>b</sup>	0.41 <sup>b</sup>							100	0 (2)

<sup>a</sup>Corrected for differential recoil-free fractions assuming  $\Theta_M(\text{Fe}^{2+}) = 340$  K;  $\Theta_M(\text{Fe}^{3+}) = 500$  K

<sup>b</sup>Weighted mean values for all Fe<sup>2+</sup> absorption

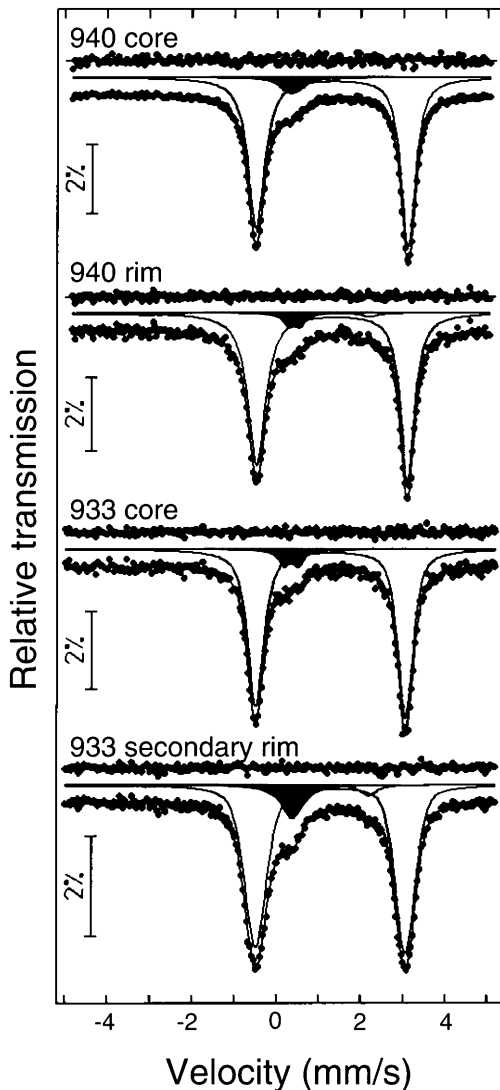
son of hyperfine parameters with literature values suggests the presence of spinel and phlogopite. Griffin et al. (1999) reported grain-boundary replacement of garnet by phlogopite, with inclusions of chrome spinel. This was confirmed in the present samples using an optical microscope. The iron content of the spinel is significantly higher than that of the phlogopite, which is reflected in the relative intensities of Mössbauer absorption. To fit garnet spectra containing these additional phases, we added doublets corresponding to Fe<sup>2+</sup> and Fe<sup>3+</sup> in each phase, and constrained hyperfine values (centre shift, quadrupole splitting and line-width) to those reported in the literature for spinel and phlogopite with similar composition (spinel: Osborne et al. 1981; phlogopite: Canil and O'Neill 1996). The larger uncertainties in Fe<sup>3+</sup>/ $\Sigma$ Fe for garnet (Table 2) reflect primarily uncertainties in the assumption of equal component areas for phlogopite. Mössbauer sites in spinel have cubic point group symmetry, hence component areas will be equal even for single crystals (Gibb 1978), but, in the case of phlogopite, component areas vary with orientation. Nevertheless, component areas of phlogopite are relatively well constrained by the high resolution of the Mössbauer spectrum in the relevant velocity range, and appear to be close to 1:1. This is consistent with petrographic observations that phlogopite is fine-grained compared with the garnet, and would imply

that there is no significant preferred orientation. In support of our fitting approach, the few garnet spectra that contain absorption from spinel and phlogopite gave Fe<sup>3+</sup>/ $\Sigma$ Fe garnet values that were consistent with results from garnet-only spectra.

## Thermobarometry

The garnets in the harzburgitic xenoliths that we studied showed marked chemical zoning over several hundred microns, whereas coexisting olivine and orthopyroxene are homogeneous based on electron microprobe data. Olivine and orthopyroxene dominate the assemblage (62 and 27 vol%, respectively, according to the Wesselton sample studied by Carswell and Dawson 1970), compared with the minor amount of garnet (3 vol%). Because olivine and orthopyroxene account for ~90% of Fe and Mg in the rock and partitioning of Fe<sup>2+</sup> and Mg between these phases is not significantly temperature dependant, Fe<sup>2+</sup>/Mg in both phases is not likely to change during metasomatism unless there was a large amount of fluid that was significantly out of equilibrium with regard to Fe<sup>2+</sup>/Mg.

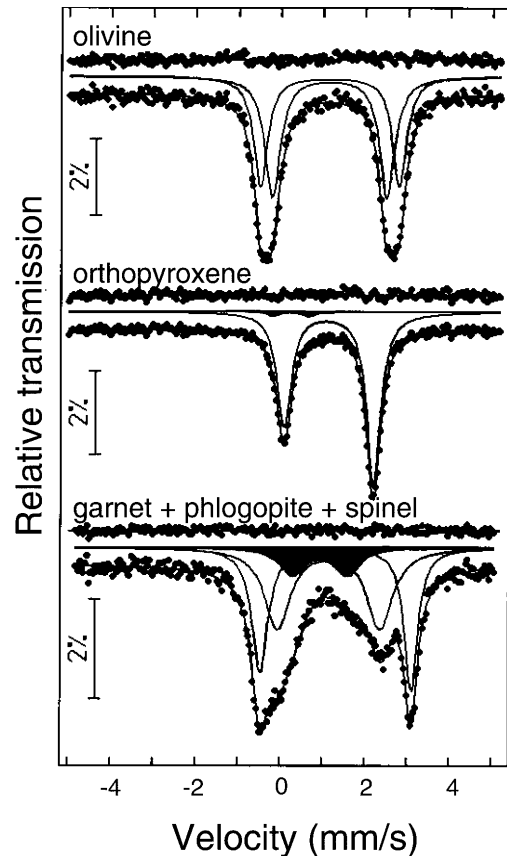
We therefore assumed that Fe<sup>2+</sup>/Mg of olivine and orthopyroxene remained constant throughout all stages of metasomatism, whereas the garnet Fe/Mg composition changed in response to the different conditions,



**Fig. 2** Room temperature Mössbauer spectra of zoned garnets illustrating the variation in  $\text{Fe}^{3+}$  absorption (shaded black) between core and rim (garnet 940), and core and secondary rim (garnet 933)

allowing temperatures to be estimated based on either the garnet–orthopyroxene or garnet–olivine thermometers. Pressures were estimated from the Al in orthopyroxene (equilibrated with garnet) barometer, but could not be determined for all stages of metasomatism. Hence, we calculated pressures based on the composition of the garnet rims equilibrated with orthopyroxene, and assumed that pressure remained relatively constant throughout the metasomatism. It should be emphasised, however, that all of the above assumptions regarding thermobarometry contribute a relatively minor uncertainty to the determination of oxygen fugacity (see the following section).

The garnet–olivine thermometer of O'Neill and Wood (1979, 1980; OW79) and the garnet–orthopyroxene thermometer of Harley (1984a; H84a) give temperatures that generally agree within  $\pm 50^\circ$  of one another (Table 3). All of the garnet cores and rims, as well as one



**Fig. 3** Room temperature Mössbauer spectra of minerals coexisting with garnet taken from samples 937 (top); 951 (middle); and 937 (bottom). Subspectra are shaded as follows: black ( $\text{Fe}^{3+}$ ); light grey ( $\text{Fe}^{2+}$  in phlogopite); dark grey ( $\text{Fe}^{2+}$  in spinel)

of the secondary rims, show temperatures of approximately 1,000 °C, whereas the other two secondary rims show temperatures approximately 200 °C higher. The Ni-in-garnet thermometer (Griffin et al. 1989; Ryan et al. 1996; G89) provides an independent temperature determination based on the Ni content of garnet, assumed to have equilibrated with olivine. Temperatures from the Ni-in-garnet thermometer are also near 1,000 °C, with two exceptions showing temperatures up to 300 °C higher (Table 3). Similar results are obtained if the Ni-in-garnet thermometer calibration of Canil (1999; C99) is used. Pressures were calculated to be  $\sim 37$  kbar using the Al in orthopyroxene barometer of Brey and Köhler (1990; BK90), where orthopyroxene was assumed to be in equilibrium with the garnet rims. Similar pressures were obtained using the formulation of Harley (1984b; H84b).

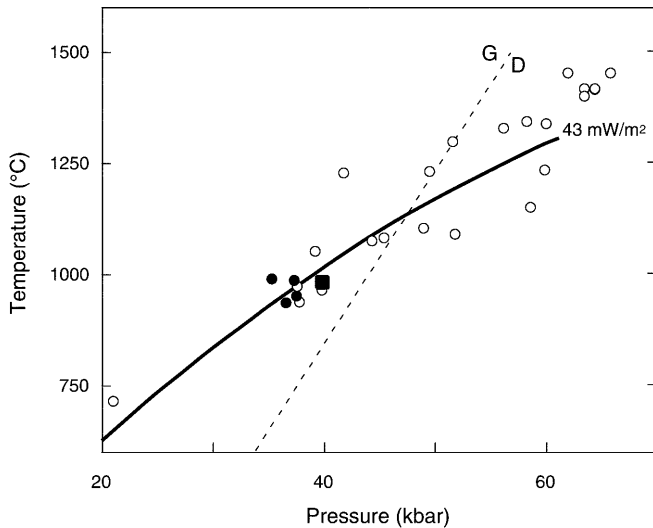
The P,T estimates for the metasomatised garnet harzburgites from the Wesselton kimberlite are consistent with results from other garnet lherzolites and harzburgites from South Africa and neighbouring localities (Table 4 and Fig. 4). We used the analyses reported by Luth et al. (1990) and Canil and O'Neill (1996) to determine P,T conditions according to either (1) the two-pyroxene thermometer of Brey and Köhler

**Table 3** Calculated equilibrium conditions and oxygen fugacities of zoned garnet xenoliths from Wesselton kimberlite

	933			937			940			951	
	Core	Rim	Sec rim	Core	Rim	Sec rim	Core	Rim	Sec rim	Core	Rim
<b>Thermobarometry</b>											
$T_{\text{gt-ol}}(\text{OW79}), ^\circ\text{C}$	986	1,022	1,238	951	1,000	1,161	935	968	998	990	957
+ $P_{\text{gt-opx}}(\text{BK90}),$ kbar		37.3			37.5			36.6			35.3
$T_{\text{gt-opx}}(\text{H84a}), ^\circ\text{C}$	998	1,005	1,113	997	987	1,080	961	964	998	976	932
+ $P_{\text{gt-opx}}(\text{H84b}),$ kbar		38.1			37.5			38.2			34.5
$T_{\text{Ni}}(\text{G89}), ^\circ\text{C}$	1,025	1,233	1,012	1,038	1,018	1,305	971	992	1,025	985	985
$T_{\text{Ni}}(\text{C99}), ^\circ\text{C}$	1,063	1,188	1,055	1,075	1,063	1,233	1,027	1,041	1,062	1,039	1,039
$P$ used, kbar	37.3	37.3	37.3	37.5	37.5	37.5	36.6	36.6	36.6	35.3	35.3
$T$ used, $^\circ\text{C}$	986	1,022	1,238	951	1,000	1,161	935	968	998	990	957
<b>Olivine/orthopyroxene compositions</b>											
$x_{\text{Fa}}$	0.069	0.069	0.069	0.063	0.063	0.063	0.059	0.059	0.059	0.058	0.058
$x_{\text{Fs}}$	0.061	0.061	0.061	0.056	0.056	0.056	0.052	0.052	0.052	0.050	0.050
<b>Cations in garnet (based on 12 O)</b>											
Si	2.991	2.977	3.018	3.001	3.027	2.981	2.979	2.970	3.004	3.026	3.023
Ti	0.016	0.022	0.031	0.005	0.019	0.028	0.002	0.010	0.009	0.001	0.021
Al	1.521	1.526	1.538	1.551	1.568	1.524	1.492	1.520	1.546	1.605	1.599
Cr	0.444	0.449	0.396	0.425	0.371	0.403	0.481	0.441	0.398	0.380	0.338
$\text{Fe}^{2+}$	0.353	0.321	0.298	0.337	0.308	0.292	0.326	0.322	0.305	0.293	0.303
$\text{Fe}^{3+}$	0.021	0.044	0.031	0.016	0.030	0.031	0.019	0.016	0.022	0.017	0.018
Mn	0.021	0.021	0.014	0.016	0.018	0.010	0.009	0.019	0.013	0.017	0.015
Ni	0.000	0.000	0.000	0.000	0.000	0.001	0.000	0.000	0.000	0.000	0.000
Mg	2.341	2.272	2.443	2.512	2.282	2.538	2.457	2.376	2.427	2.474	2.306
Ca	0.297	0.376	0.213	0.138	0.357	0.219	0.267	0.361	0.288	0.165	0.360
Na	0.008	0.015	0.001	0.008	0.007	0.004	0.003	0.010	0.004	0.004	0.012
<b>Oxygen fugacity</b>											
$\log f_{\text{O}_2}(\text{FMQ}), \text{bar}$	-2.2(5)	-1.2(10)	-1.6(4)	-2.4(6)	-1.5(5)	-1.4(9)	-1.9(5)	-2.3(7)	-1.7(11)	-1.9(5)	-1.9(7)

**Table 4** Oxygen fugacities of South African garnet peridotites recalculated using the Gudmundsson and Wood (1996) oxybarometer

Sample	Locality	$P$ (kbar)	$T$ ( $^\circ\text{C}$ )	$x_{\text{Fa}}$	$x_{\text{Fs}}$	gt $\text{Fe}^{3+}/\Sigma\text{Fe}$ (%)	$\log f_{\text{O}_2}$ (FMQ) (bar)
<b>Luth et al. (1990)</b>							
FRB838	Bultfontein	38	937	0.084	0.075	3.3	-2.9
FRB135	Mothae	38	973	0.069	0.060	5.7	-2.4
PHN1917	Mothae	44	1,075	0.068	0.059	4.8	-3.2
FRB131	Mothae	45	1,081	0.073	0.064	5.4	-2.9
FRB1033	Jagersfontein	56	1328	0.091	0.078	11.6	-2.9
PHN1611	Thaba Putsoa	62	1,452	0.121	0.101	12.2	-3.5
PHN5549	Gibeon Town-lands #1	42	1,228	0.088	0.081	5.4	-3.3
PHN1925	Mothae	66	1,453	0.093	0.084	11.8	-3.7
PHN5267	Premier	64	1,415	0.087	0.073	12.1	-3.9
BD2501	Mothae	52	1,298	0.081	0.072	11.2	-2.6
FRB76	Frank Smith	60	1,338	0.088	0.076	11.8	-3.4
FRB140	Mothae	58	1,344	0.093	0.081	10.7	-3.3
<b>Canil and O'Neill (1996)</b>							
BD1140	Bultfontein	40	965	0.082	0.073	5.1	-2.3
BD1150	Bultfontein	49	1,104	0.097	0.083	8.3	-2.5
BD1201	Wesselton	40	981	0.083	0.087	3.3	-3.4
BD1354	Matsoku	60	1,234	0.173	0.147	9.8	-3.7
PHN5239	Finsch	64	1,418	0.085	0.073	10.8	-3.5
PHN5273	Finsch	50	1,231	0.078	0.080	10.6	-2.9
FRB909	Premier	64	1,400	0.091	0.082	11.5	-3.7
FRB921	Premier	39	1,052	0.078	0.070	4.4	-2.8
PHN5267	Premier	64	1,418	0.087	0.073	12.0	-3.9
F865	Premier	52	1,090	0.071	0.061	7.3	-3.3
F556	Premier	59	1,151	0.089	0.072	7.2	-4.0
FRB1350	Premier	21	714	0.083	0.074	3.4	-1.1

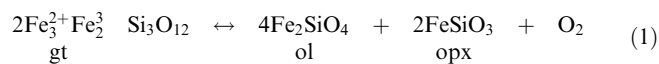


**Fig. 4** Pressure–temperature plot for samples from this study (*solid circles*), garnet lherzolite from Wesselton (Carswell and Dawson 1970) (*square*) and garnet peridotites from South Africa (Luth et al. 1990; Canil and O'Neill 1996) (*open circles*). Also shown are *curves* for the graphite–diamond phase boundary (Kennedy and Kennedy 1976) (*dotted line*) and a 43 mW/m<sup>2</sup> model conductive geotherm (Pollack and Chapman 1977) (*solid line*)

(1990) combined with the BK90 garnet–orthopyroxene barometer, or (2) the garnet–olivine thermometer (OW79) combined with the garnet–orthopyroxene barometer (BK90) for assemblages where no clinopyroxene was present. Results from a garnet lherzolite collected from the Wesselton mine (originally described by Carswell and Dawson 1970) gave conditions of 40 kbar and 981 °C, which are nearly identical to the values determined for the metasomatised samples from Wesselton described in the present study.

### Oxygen fugacity

Oxygen fugacity was estimated based on the garnet–orthopyroxene–olivine oxybarometer (Luth et al. 1990), which was re-calibrated by Gudmundsson and Wood (1995) using the following reaction:



where  $\text{Fe}_3^{2+}\text{Fe}_2^3\text{Si}_3\text{O}_{12}$  is the skiaigite component in the garnet. We used the corrected expression for the skiaigite activity coefficient reported by Woodland and Peltonen (1999).

Oxygen fugacities corresponding to the different regions of the four garnets are listed in Table 3, where values are given relative to the fayalite–magnetite–quartz (FMQ) buffer of O'Neill (1987) corrected for pressure (Ballhaus et al. 1991). Conditions tend to more oxidising going from core to rim to secondary rim. The calculated relative oxygen fugacities do not vary

significantly with temperature or pressure, where a variation of 200 °C or 2 kbar produces a difference in  $\Delta\log f_{\text{O}_2}$  of only 0.2 log-bar units. Therefore, a heating event accompanying metasomatism does not have a significant effect on oxygen fugacities. Uncertainties in the compositions of olivine and orthopyroxene also contribute small errors:  $\leq 0.3$  log-bar units for a variation in  $x_{\text{Fe}}$  of  $\pm 0.005$ . To explain the higher  $\text{Fe}^{3+}/\Sigma\text{Fe}$  values in the secondary rims compared with the cores by changes in temperature or pressure alone would require the secondary rims to have formed below 600 °C, or at pressures more than 47 kbar, and so, in the absence of drastic changes to temperature or pressure during the formation of these garnets, there is a clear indication of an increase in relative oxygen fugacity from core to secondary rim.

### Kinetics of oxidation and diffusion

The extent to which differences in oxidation state will be preserved in the garnet depends to varying degrees on (1) the relevant cation and anion diffusion rates in the garnet structure, and (2) the kinetics of the oxidation/reduction reaction. When oxidised fluids infiltrate the mineral assemblage, oxygen in the fluid will react with  $\text{Fe}^{2+}$  in the garnet to form  $\text{Fe}^{3+}$ . Oxidation studies of fayalite (Mackwell 1992) and measurements of oxygen self-diffusion in diopside (Farver 1989) indicate that oxygen ions are relatively immobile compared with iron cations, which is probably true for garnet also. The oxidation reaction therefore occurs at the surface of the garnet, and proceeds through diffusion of iron cations through the lattice. The rate of iron diffusion probably provides a minimum estimate for the kinetics of garnet oxidation, based on studies of fayalite oxidation (Mackwell 1992). In upper mantle garnets that are low in Ti,  $\text{Fe}^{2+}$  occupies only the dodecahedral site, whereas  $\text{Fe}^{3+}$  occupies only the octahedral site (e.g. Amthauer et al. 1976), implying that exchange of  $\text{Fe}^{2+}$  and  $\text{Fe}^{3+}$  in garnet must involve cation exchange over different crystallographic sites, and would be probably of similar magnitude to diffusion rates of other cations in garnet. This would imply that zoning profiles of  $\text{Fe}^{3+}$  would be preserved on time-scales similar to those of the major and trace elements.

Volume diffusion of iron in garnet is probably controlled by vacancies, hence the diffusivity of cations will vary as a function of oxygen fugacity (e.g. Buening and Buseck 1973).  $\text{Fe}^{2+}/\text{Fe}^{3+}$  diffusivity is probably more sensitive to oxygen fugacity than other major and trace elements, hence increasing oxygen fugacity would increase the rate at which  $\text{Fe}^{2+}$  is oxidised to  $\text{Fe}^{3+}$  in the garnet structure. Griffin et al. (1999) noted that the original major and trace element compositions of the garnet cores were probably not preserved, hence it is likely that the original  $\text{Fe}^{3+}/\Sigma\text{Fe}$  of the garnet was also not preserved. The garnet cores therefore reflect a higher

oxygen fugacity than that existing prior to metasomatism.

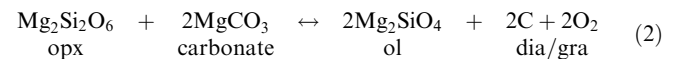
### Implications for diamond preservation

The oxygen fugacities determined from the metasomatised garnet harzburgites of the Wesselton kimberlite can be compared with conditions in other parts of the Kapvaal craton. We used the values for  $\text{Fe}^{3+}/\Sigma\text{Fe}$  in garnet determined by Luth et al. (1990) and Canil and O'Neill (1996), combined with the chemical analyses reported for coexisting olivine, orthopyroxene and garnet, to calculate the relative oxygen fugacity according to the Gudmundsson and Wood (1995) garnet–orthopyroxene–olivine oxybarometer. Pressures and temperatures were determined as described above. Results

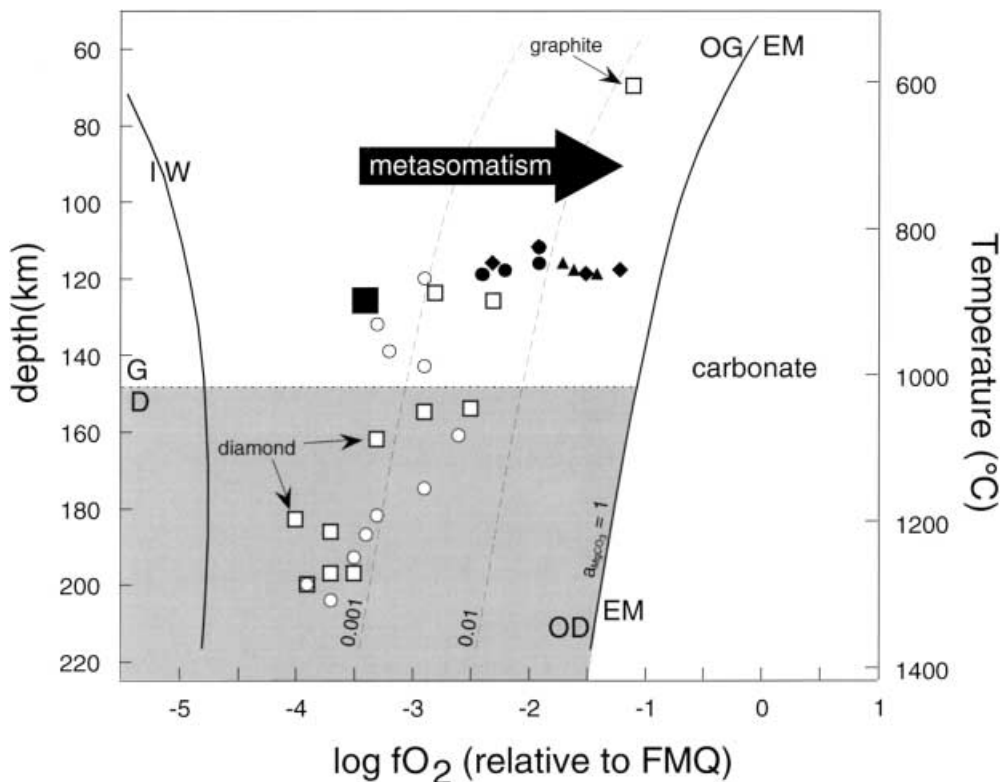
**Fig. 5** Oxygen fugacity (relative to FMQ) of garnet peridotites as a function of depth. Zoned garnets from Wesselton (this study) are indicated by *solid circles* (cores), *solid diamonds* (rims) and *solid triangles* (secondary rims), while an unzoned Wesselton sample (Carswell and Dawson 1970; Canil and O'Neill 1996) is indicated by the large solid square. Oxygen fugacities calculated from data of Luth et al. (1990) (*open circles*) and Canil and O'Neill (1996) (*open squares*) are also shown. Assemblages containing graphite and diamond are indicated by *arrows*. The graphite–diamond transition (Kennedy and Kennedy 1976) and the EMOG/EMOD reactions (Eggler and Baker 1982) define the maximum stability region of diamond (*shaded grey*), where *dashed lines* indicate iso-activity curves for carbonate. The iron–wüstite buffer ( $P,T$  formulation of Ballhaus et al. 1991) is indicated for reference, and temperatures are given according to the 43 mW/m<sup>2</sup> model conductive geotherm (Pollack and Chapman 1977)

show that for nearly all samples, relative oxygen fugacities are lower than those found in the present study (Table 4). In particular, the relative oxygen fugacity calculated for the garnet lherzolite from the Wesselton mine described by Carswell and Dawson (1970) is more than one log-bar unit below the smallest value for the metasomatised assemblage. This is consistent with the conclusion by Griffin et al. (1999) that the garnet cores have also been metasomatised.

The increase in relative oxygen fugacity during metasomatism has some relevance to diamond preservation. The upper oxygen fugacity limit of graphite or diamond stability in harzburgites is defined by the EMOG/EMOD reaction:



(Eggler and Baker 1982). Together with the graphite–diamond univariant equilibrium, this defines the maximum stability region of diamond as a function of pressure, temperature and oxygen fugacity in the peridotite system (Fig. 5). The activity of carbonate in the system increases with increasing oxygen fugacity according to reaction (2) until the EMOG boundary is reached. For diamond-bearing assemblages that undergo metasomatism, the degree of diamond preservation is probably influenced by the activity of carbonate in the metasomatic fluid. Increasing carbonate activity would increase the degree of diamond reaction with the fluid, hence causing resorption of the diamond even at oxygen fugacities below the EMOG boundary. Figure 5 illustrates two iso-activity curves for carbonate. Although





carbonate activities of 0.001 and 0.01 would not be expected to cause reaction of diamond with the fluid, higher activities could be expected to produce an observable reaction.

The Wesselton mine harzburgites show a progressive oxidation of approximately two log-bar units from the original assemblage to the final stage of metasomatism, nearly reaching the EMOG curve (Fig. 5). This is consistent with suggestions by Griffin et al. (1999) that the fluid in stage (3) was derived from the kimberlite itself, which is generally believed to be a hostile environment to diamonds (Woermann and Rosenhauer 1985 and references therein). Other garnet peridotites from South Africa show significantly lower relative oxygen fugacities, and many lie within the diamond stability field (Fig. 5). Diamond would hence be preserved during the initial stages of metasomatism, up to an increase in relative oxygen fugacity of 1–2 log-bar units. Beyond that point, however, the fluid would react with the diamonds, leading to their oxidation and resorption.

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