



# Observed structural relationships between ocean chlorophyll variability and its heating effects on the ENSO

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## Abstract

Ocean chlorophyll (Chl)-induced heating can affect the climate system through the penetration of solar radiation in the upper ocean. Currently, the ocean biology-induced heating (OBH) feedback effects on the climate in the tropical Pacific are still not well understood, and the mechanisms regarding how SST is modulated remain elusive. In this paper, chlorophyll (Chl) data from satellites are combined with physical fields from Argo profiles to estimate OBH-related fields, including the penetration depth ( $H_p$ ) and the ocean mixed-layer (ML) depth ( $H_m$ ). In addition, some directly related heating terms with  $H_m$  and  $H_p$  are diagnosed, including the absorbed solar radiation component within the ML (denoted as  $Q_{abs}$ ), the rate of ML temperature changes that are directly induced by  $Q_{abs}$  (denoted as  $R_{sr} = Q_{abs}/(\rho_0 C_p H_m)$ ), and the portion of solar radiation that penetrates through the bottom of the ML (denoted as  $Q_{pen}$ ). The structural relationships between these related fields are examined to illustrate how these heating terms are affected by  $H_p$  and  $H_m$ . The extent to which  $R_{sr}$  and  $Q_{pen}$  are modulated by  $H_p$  is strikingly different during ENSO cycles. In the western-central equatorial Pacific, inter-annual variations in  $H_p$  tend to be out of phase with those in  $H_m$ . A decrease (increase) in  $Q_{abs}$  from a positive (negative)  $H_p$  anomaly during El Niño (La Niña) tends to be offset by a negative (positive)  $H_m$  anomaly. Thus,  $R_{sr}$  is not closely related with  $H_p$ , even though  $Q_{abs}$  is highly correlated with  $H_p$ , indicating that the direct thermal effect through  $Q_{abs}$  is not a dominant factor that affects the SST. In contrast, the inter-annual variability of  $Q_{pen}$  in the region is significantly enhanced by that of  $H_p$ , with their high positive correlation. The  $H_p$ -induced differential heating in the ML and subsurface layers from the  $Q_{pen}$  and  $Q_{abs}$  terms modifies the thermal contrast, stratification and vertical mixing, which represent a dominant indirect ocean dynamical effect on the SST. The revealed relationships between these related fields provide an observational basis for gaining structural insights into the OBH feedback effects and validating model simulations in the tropical Pacific.

**Keywords** Chlorophyll variability · Ocean biology-induced heating · ENSO modulations · The penetration depth and mixed layer depth · Satellite and Argo data

## 1 Introduction

Interactions between multiple processes are prominent in the tropical Pacific and are responsible for seasonal and inter-annual variabilities. For example, ocean–atmosphere

coupling in the region creates the El Niño Southern Oscillation (ENSO), which is a dominant inter-annual mode in the climate system (e.g. Philander 1983; Zebiak and Cane 1987; Zhang and Levitus 1997; Zhang and Gao 2016; Tang et al. 2018). Additionally, various feedbacks and couplings

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coexist in the tropical Pacific that involve interactions among ocean physics, ocean biology and the climate, which can significantly modulate the ENSO (Zhang and Busalacchi 2008, 2009; Zhang et al. 2012; Zhi et al. 2015; Kang et al. 2017a, b). In particular, ocean biology can influence ocean physics through its modulation of penetrative solar radiation in the upper ocean (Timmermann and Jin 2002; Marzeion et al. 2005; Wetzel et al. 2006; Lengaigne et al. 2007; Gnanadesikan and Anderson 2009; Zhang et al. 2009; Anderson et al. 2009; Jochum et al. 2010; Park et al. 2014). Indeed, recent studies revealed pronounced ocean biology-induced heating feedback on the ENSO over the tropical Pacific (Zhang 2015; Kang et al. 2017a). As a major driver that strongly affects the upwelling and mixing processes in the equatorial Pacific, the ENSO greatly changes the Chl in the region. At the same time, corresponding biological perturbations modulate the penetrative solar radiation in the upper ocean, which produces a feedback on the ENSO.

The solar radiation that arrives at the sea surface penetrates through the upper water column and its shortwave component decays exponentially with depth. Some is absorbed directly within the mixed layer (ML), while some penetrates through the bottom of the ML and into subsurface layers. The ML depth ( $H_m$ ) is a major physical factor that determines the vertical partitioning of solar radiation between the mixed layer (ML) and subsurface layers. Additionally, phytoplankton, detritus and colored dissolved organic matter (CDOM) can absorb solar radiation mainly in the visible spectrum (380–700 nm), thus affecting the vertical penetration of the incoming solar radiation in the upper ocean (Lewis et al. 1990; Morel and Antoine 1994; Siegel et al. 2002; Strutton and Chavez 2004; Kang et al. 2017a). Among these components, chlorophyll is the main factor that attenuates shortwave radiation and consequently affects the ocean's thermal conditions. Correspondingly, a penetration depth ( $H_p$ ) can be defined to represent the extent to which the penetration of solar radiation in the upper ocean is affected by the Chl concentration. Thus, the depth of the mixed layer ( $H_m$ ) and penetration depth ( $H_p$ ) are two factors that affect the distribution of penetrative solar radiation in the ML and subsurface layers; the former ( $H_m$ ) is controlled by physical processes and the latter ( $H_p$ ) is determined by biological processes. Furthermore, bio-effects are realized through a few heating terms that are directly associated with Chl (and the corresponding  $H_p$ ) and  $H_m$ , including the absorbed solar radiation flux within the ML (denoted as  $Q_{\text{abs}}$ ) and its directly related temporal rate of change in the ML temperature (denoted as  $R_{\text{sr}}$ ), alongside the penetrated portion of solar radiation through the bottom of the ML (denoted as  $Q_{\text{pen}}$ ). These heating terms can be used to characterize and quantify the Chl-induced heating feedback effects on the SST because these terms act to modulate

the thermodynamics in the upper ocean, which can further induce indirect dynamical effects on the SST.

Previously, bio-climate interactions were investigated by using coupled ocean–atmosphere models with different degrees of complexity. Large uncertainties exist in model representations of ocean biology-induced heating (OBH) feedback, and its simulated effects strongly depend on models used. Thus, modulations in SST from ocean biology are sensitive to the manner in which bio-effects are represented in models. Currently, models exhibit biases in terms of simulating Chl variability in the tropical Pacific. Additionally, the physical mechanisms that are involved in bio-effects significantly differ among various previous modeling studies because different dominant processes are apparently at work. What are the main factors that control ocean biology-induced effects and the involved physical processes? Why do large intermodal differences exist in bio-effect simulations? Clearly, observations are vital to reveal the physical mechanism for these ocean biology-induced effects and validate model-based simulations and related bio-effects.

Over the past decade, remote sensing has made great progress in providing time series of ocean color data and associated products. For example, chlorophyll (Chl) data are available from satellite measurements and have been used to describe and understand bio-climate interactions (McClain et al. 1998). As observed in the tropical Pacific, for example, the existence and variations of Chl affect surface-layer water turbidity and modify the vertical redistribution of solar radiation in the upper ocean. Correspondingly, the Chl concentration can be used to estimate  $H_p$ . Also, Argo profiles provide unprecedented temperature and salinity data in the upper ocean, which can be used to estimate ocean fields on basin scales, including the depth of the ML ( $H_m$ ). Thus, sufficient observational data of Chl and physical fields have been accumulated in the last decade, which allows us to perform detailed analyses on Chl-induced heating effects in the tropical Pacific.

In this paper, we perform an observation-based analysis for the biological heating effect, with a focus on the inter-annual variability that is associated with the ENSO in the tropical Pacific. A diagnostic analysis method is developed that makes full use of observational data to reveal bio-effects from inter-annual Chl anomalies. More specifically, observed Chl data are combined with other derived fields to characterize the structural relationships of Chl-related heating terms with  $H_m$  and  $H_p$ . For example, observed  $H_m$  and  $H_p$  fields are used to quantify the combined effects on the penetration of solar radiation and related heating effects. The revealed relationships are used to infer the processes that are involved and understand the influence pathways through which the SST is modulated by inter-annual Chl anomalies.

This paper is organized as follows. Section 2 describes the heating terms that are associated with Chl and the

observational data that are used for the analyses. The structural relationships of these heating terms with  $H_m$  and  $H_p$  are analyzed in Sect. 3. Section 4 provides a summary and discussion. An Appendix is provided for the analyses of the annual mean and seasonal variations in some related fields.

## 2 Ocean Chl-induced heating terms and observational data

Comprehensive data are used to diagnose Chl-induced heating terms and bio-effects on the SST, including ocean physical and biological fields that are derived from satellite-based observations and in situ Argo profiles. As shown below, these heating terms are determined by the ocean mixed-layer depth ( $H_m$ ) and penetration depth ( $H_p$ ); the former is determined by physical processes and the latter is determined by biological processes. Observational data are used to quantify these heating terms and structural relationships between physical and biological fields. The effects on these terms can be analyzed using observational data, which are defined below.

### 2.1 Ocean Chl-induced heating terms

The incoming solar radiation that arrives at the sea surface tends to decay very sharply with depth in the upper ocean (a factor that reflects the effect of pure water on penetrative solar radiation); the depth of the mixed layer ( $H_m$ ) is a major factor that determines the vertical partitioning of solar radiation between the mixed layer (ML) and subsurface layers. At the same time, the existence and variations of Chl can affect the penetrative solar radiation in the upper ocean; therefore, the penetration depth ( $H_p$ ) can be introduced to quantify the extent to which solar radiation is affected by biological fields in the upper ocean.

Some heating-term fields are directly related to  $H_p$  and  $H_m$ , including the absorbed solar radiation flux within the ML (denoted as  $Q_{abs}$ ) and its directly related temporal rate of change in the ML temperature (denoted as  $R_{sr}$ ), alongside the portion of solar radiation that penetrates the bottom of the ML (denoted as  $Q_{pen}$ ).

#### 2.1.1 Penetration depth ( $H_p$ )

Ocean biological components (such as phytoplankton, detritus and colored dissolved organic matter (CDOM)) can absorb solar radiation in the visible spectrum (380–700 nm) and further modify the vertical penetration of the incoming solar radiation in the upper ocean (Lewis et al. 1990; Morel and Antoine 1994; Siegel et al. 2002; Strutton and Chavez 2004). The penetration of the solar shortwave radiation in the upper ocean follows the Beer-Lambert Law, which

indicates that shortwave radiation decays exponentially with depth according to attenuation coefficients. Here, chlorophyll is considered to have attenuating effects on shortwave radiation.

Thus, the total attenuation coefficient can be computed as follows (Wang et al. 2008):

$$K_A(z) = K_W + K_C \text{Chl}(z), \quad (1)$$

where  $z$  is the depth;  $K_W = 0.028 \text{ m}^{-1}$  and  $K_C = 0.058 \text{ m}^{-1} (\text{mg chl m}^{-3})^{-1}$  represent the attenuation coefficients for pure water and chlorophyll, respectively; and  $\text{Chl}(z)$  is the chlorophyll concentration. In this study, we only consider the Chl effect on  $K_A$  because Chl plays a dominant role in ocean biology-induced heating effects.

Then, a penetration depth ( $H_p$ ) of the shortwave radiation is defined as the inverse of  $K_A$  calculated from the satellite observations of the Chl field (Murtugudde et al. 2002; Zhang et al. 2011). This  $H_p$  field serves as a linkage between ocean biology and physics and can be used to quantify the ocean biology-induced heating effects on the related heating terms.

#### 2.1.2 Absorbed solar radiation within the ML ( $Q_{abs}$ and $R_{sr}$ )

As expressed in Zhang (2015),  $H_p$  and  $H_m$  are explicitly associated with several ocean biology-induced heating terms.  $Q_{abs}$  denotes the absorbed solar radiation flux within the ML, and  $R_{sr}$  denotes the temporal rate of change in the ML temperature that directly results from the  $Q_{abs}$  effect on SST, both of which are written as follows

$$Q_{abs}(H_m, H_p) = Q_{sr}[1 - \gamma \exp(-H_m/H_p)], \quad (2)$$

$$R_{sr}(H_m, H_p) = Q_{sr}[1 - \gamma \exp(-H_m/H_p)]/(\rho_0 c_p H_m) = Q_{abs}/(\rho_0 c_p H_m), \quad (3)$$

where  $Q_{sr}$  is the incoming solar radiation flux at the sea surface,  $\gamma$  is a constant ( $=0.33$ ) that denotes the fraction of the available radiation to penetrate to depths beyond the first few centimeters of the sea surface,  $C_p$  is the heat capacity, and  $\rho_0$  is the density of sea water.

As expressed above,  $Q_{abs}$ , which is a function of both  $H_m$  and  $H_p$ , is determined by changes in  $H_m$  and  $H_p$ . On the one hand,  $Q_{abs}$  increases exponentially with  $H_m$ . The deeper the ML, the more solar radiation that is directly absorbed within the ML. On the other hand,  $Q_{abs}$  also decreases exponentially with  $H_p$ ; the larger the  $H_p$  value, the less solar radiation that is directly absorbed within the ML, producing more penetration through the bottom of the ML. Therefore, changes in  $H_m$  and  $H_p$  tend to have an opposite effect on  $Q_{abs}$ . That is, a negative (positive) perturbation in  $H_m$  reduces (increases)  $Q_{abs}$ , whereas a negative (positive) perturbation in  $H_p$  increases (reduces)  $Q_{abs}$ .

Furthermore, the absorbed solar radiation within the ML ( $Q_{abs}$ ) directly changes  $R_{sr}$ , thus further changing the

SST, which represents a direct thermal effect.  $R_{sr}$  is proportional to  $Q_{abs}$  (which has an exponential relationship with  $H_m$  and  $H_p$ ) and inversely proportional to  $H_m$  (appearing as a denominator). Thus,  $R_{sr}$  can be affected by  $H_m$  in two fashions, whose effects tend to be of opposite signs. On the one hand,  $H_m$  is exponentially related to  $Q_{abs}$ , so a deepening (shoaling) of the ML (which increases (decreases)  $Q_{abs}$ ) can increase (decrease)  $R_{sr}$ . On the other hand,  $H_m$  is inversely proportional to  $R_{sr}$  (appearing as a denominator), so the deepening (shoaling) of the ML can decrease (increase)  $R_{sr}$ . These two effects on  $R_{sr}$  that are induced by  $H_m$  and  $H_p$  have implications for how this term is modulated by  $H_p$ .

### 2.1.3 Penetrative solar radiation flux through the bottom of the ML ( $Q_{pen}$ )

The solar radiation flux that penetrates the bottom of the ML ( $Q_{pen}$ ) is written as

$$Q_{pen}(H_m, H_p) = Q_{sr}[\gamma \exp(-H_m/H_p)] = Q_{sr} - Q_{abs}. \quad (4)$$

As expressed,  $Q_{pen}$  has exponential relationships with  $H_m$  and  $H_p$ , which tend to have opposite signs relative to  $Q_{abs}$ . For example,  $Q_{pen}$  decreases exponentially with  $H_m$  but increases exponentially with  $H_p$ ; that is, the deeper the ML, the less solar radiation that directly penetrates through the bottom of the ML. Similarly, the larger the  $H_p$  value, the deeper that the solar radiation directly penetrates into the subsurface layer. Therefore, changes in  $H_m$  and  $H_p$  tend to have opposite effects on  $Q_{pen}$ .

## 2.2 Observational data

Various observational data are used to describe the inter-annual variability and relationships among some physical and biological fields. The SST fields are obtained from Reynolds et al. (2002). The shortwave solar radiation data are from the MERRA re-analysis product (Rienecker et al. 2011). The surface chlorophyll datasets are obtained from the GlobColour project from 1998 to 2007, which supplied continuous datasets for merged Level-3 Ocean Color products (including the SeaWiFS, MODIS, MERIS and VIIRS sensors; see details at <http://hermes.acri.fr/index.php>; Maritorena et al. 2010). Then, monthly CHL-1 data (chlorophyll concentration ( $\text{mgm}^{-3}$ ) for case-1 waters) are interpolated from  $0.25^\circ \times 0.25^\circ$  grids to our analysis grids ( $1.0^\circ \times 1.0^\circ$ ).

The Argo-based dataset is provided by the International Pacific Research Center (IPRC)/Asia-Pacific Data-Research Center (APDRC). This product includes the MLD and three-dimensional gridded fields for temperature and salinity with a  $1^\circ$  horizontal resolution at the standard depths. The MLD is defined as the depth at which the density increases from 10 m to a value that is equivalent to a temperature drop of  $0.2^\circ\text{C}$ .

These data during 2005–2015 are averaged to form climatological monthly fields, which are used to calculate their inter-annual anomalies. In addition to these directly observed data, some related heating fields are estimated, including a penetration depth ( $H_p$ ),  $Q_{abs}$ ,  $R_{sr}$ , and  $Q_{pen}$ .

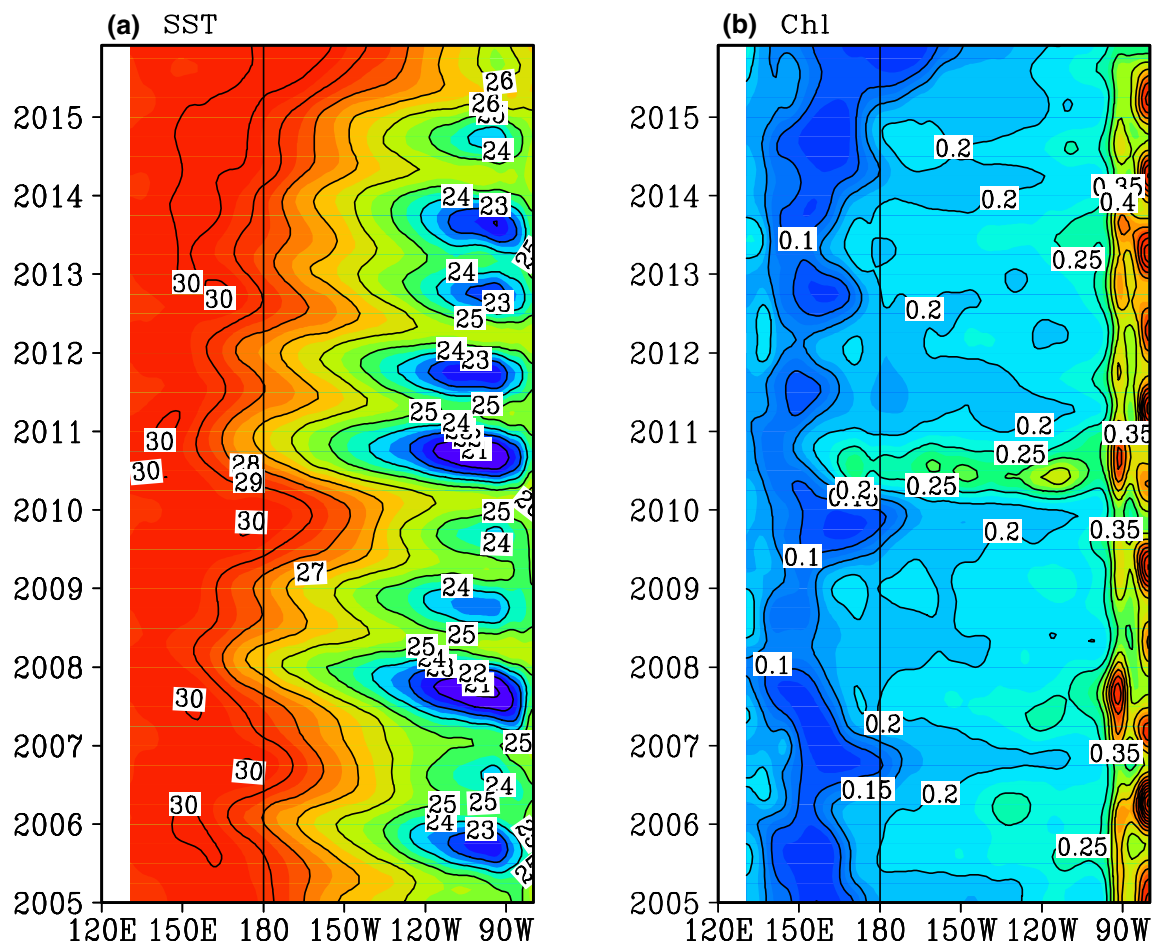
## 3 Structural relationships associated with the bio-effects on ENSO

Figures 1, 2, and 3 display the evolution along the Equator for various related total fields; the corresponding annual-mean and seasonal variations are presented in the Appendix. The inter-annual variabilities of these physical and biological fields are dominated by ENSO signals in the tropical Pacific. Note that full quantity and its variability exhibit differences in their magnitudes between physical and biological fields. For example, seasonal and inter-annual variabilities of physical fields are generally one order of magnitude smaller than its total field. However, the magnitudes of seasonal and inter-annual Chl variabilities are comparable to those of its total value, and the total  $Q_{abs}$  field is one order of magnitude larger than its seasonal and inter-annual variability.

### 3.1 Inter-annual variations in SST and Chl

The total SST field (Fig. 1a) is characterized by a warm pool in the western tropical Pacific and a cold tongue in the eastern equatorial Pacific. The warm pool in the west and cold tongue in the east exhibit large zonal displacements during ENSO cycles. During La Niña, the warm pool retreats to the west, whereas the cold tongue in the east develops and expands westward along the Equator, with the  $25^\circ\text{C}$  SST isotherm located west of  $150^\circ\text{W}$ . During El Niño, the warm waters in the west extend eastward along the Equator (e.g., the  $27^\circ\text{C}$  SST isotherm extends eastward to the east of  $150^\circ\text{W}$ ), while the cold tongue shrinks in the east.

As a major driver, the ENSO induces pronounced perturbations to ocean physical and biological fields in the tropical Pacific (Figs. 1, 2). For example, shortwave radiation exhibited a pronounced inter-annual variation associated with ENSO (Fig. 2a); one interesting feature is the large shift for shortwave radiation around 2010. We speculate that the eastward migration of atmospheric convection center during El Niño is responsible for a decrease in shortwave radiation near the dateline. During 2009–2010, a record-breaking warm sea surface temperature emerged in the central Pacific, which is the strongest Central Pacific type of El Niño in the 21st century (Lee and McPhaden 2010). In terms of biological field, a well-defined pattern of inter-annual Chl variability is seen during El Niño and La Niña cycles. The main characteristics of the mean Chl field and variability have been previously described (e.g. Ballabrera-Poy et al.



**Fig. 1** Time-longitude sections along the Equator during 2005–2015 for the **a** SST and **b** Chl fields from the GlobColour Project. The contour interval is 1 °C in **a** and 0.05 mg m<sup>-3</sup> in **b**

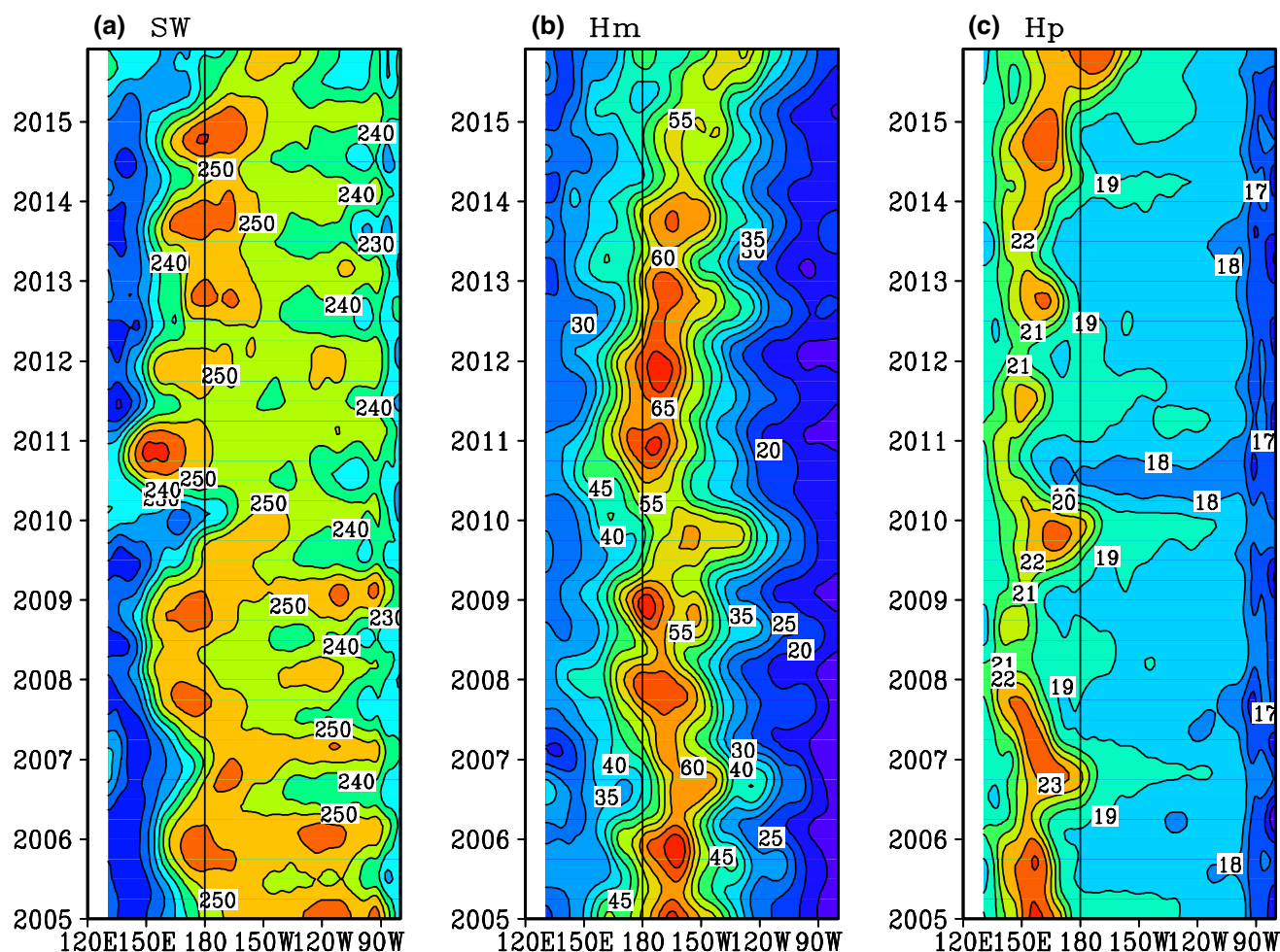
2007). Climatologically (see Fig. 1b and the Appendix), Chl exhibits elevated values from the west to the east across the tropical Pacific: its concentration is low (<0.1 mg m<sup>-3</sup>) in the western equatorial Pacific in association with the warm waters but is high (>0.2 mg m<sup>-3</sup>) in the eastern equatorial region, where the cold tongue develops. Additionally, the Chl concentration is very high in the eastern coastal regions.

Inter-annually, the Chl concentration in the equatorial Pacific significant increases during La Niña events and drops during El Niño events. Detailed examinations indicate that the inter-annual Chl variability exhibits different characteristics in the west and east. In the western-central equatorial Pacific, the Chl field has a clear east–west migration along the Equator during ENSO cycles. Regions with large values extend to the date line during El Niño. In the eastern equatorial Pacific, large variations are associated with local oceanic processes (e.g., upwelling).

The inter-annual variability features of the SST and Chl fields can be more clearly seen in their anomaly fields (Fig. 4). The magnitude of the inter-annual Chl variability

is comparable to that of the total Chl field. ENSO is a major source for inter-annual variability of Chl, with negative anomaly during El Niño events and positive anomaly during La Niña events. The inter-annual Chl variations in the east do not exhibit obvious propagation along the Equator (i.e., they are almost in phase in terms of time), but those in the west have zonal migration along the Equator (Fig. 1). The inter-annual Chl variability is quantified by calculating its standard deviations (Fig. 5a). Two large Chl variability centers are seen: one in the western equatorial Pacific to the west of the date line and another in the eastern equatorial and coastal regions.

Clear relationships exist between the inter-annual variations in SST and Chl; the latter closely follows the former, characterized by their out-of-phase fluctuations during ENSO cycles (Fig. 4). Thus, the ENSO induces bio-responses that are quick and coherent at large scales. In addition, the Chl and SST exhibit spatial shifts in their large variability regions. While the largest SST anomalies occur in the central and eastern equatorial Pacific (Fig. 4a),



**Fig. 2** Time-longitude sections along the Equator during 2005–2015 for the **a** shortwave radiation, **b**  $H_m$ , and **c**  $H_p$  fields. The contour interval is  $10 \text{ W m}^{-2}$  in **a**,  $5 \text{ m}$  in **b**, and  $1 \text{ m}$  in **c**

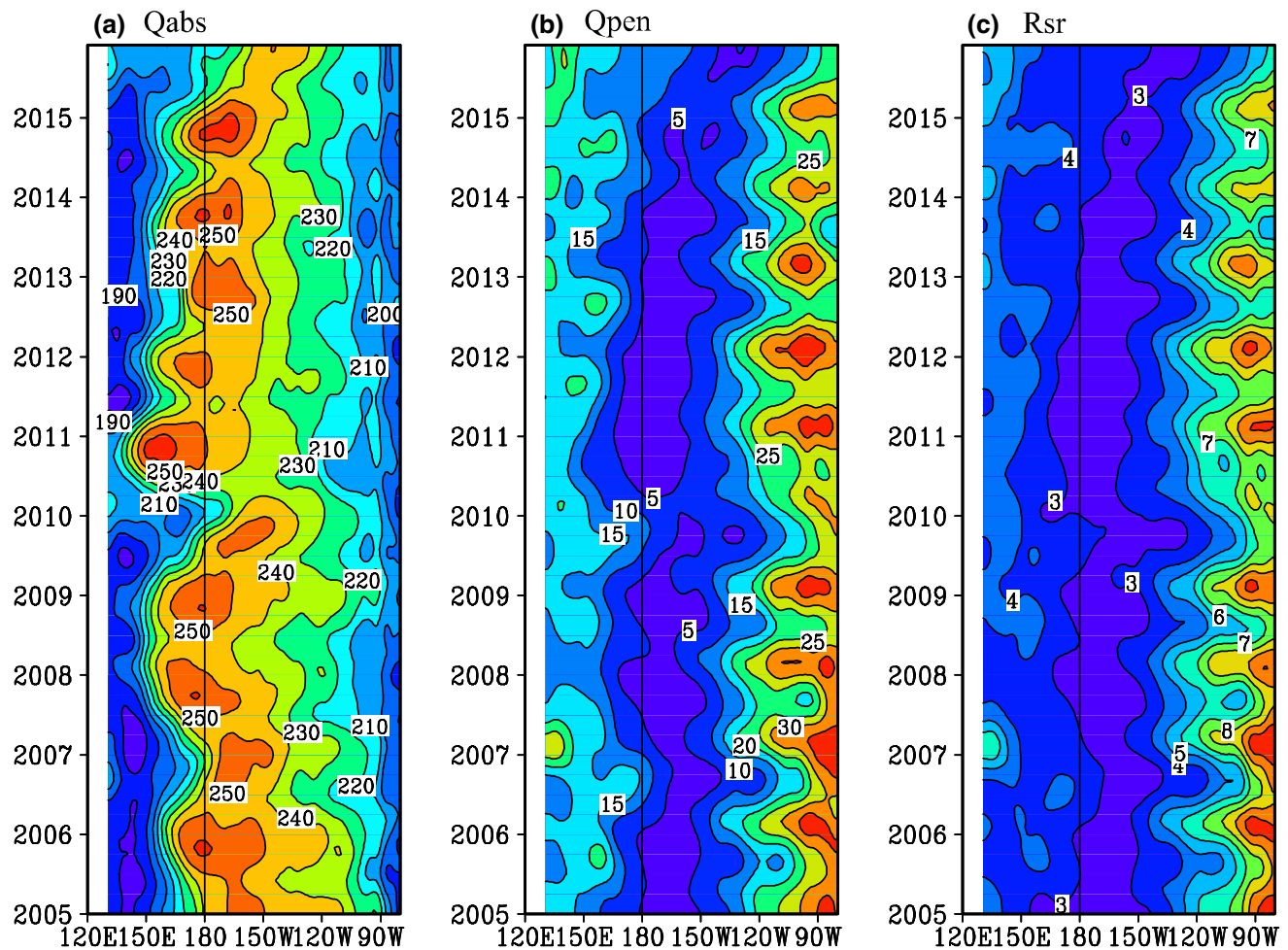
a pronounced inter-annual variability of Chl is located in the western-central equatorial Pacific (Fig. 4b). During La Niña events, for example, a cold SST anomaly is located in the central-eastern equatorial Pacific, but regions with positive Chl concentration are seen in the western-central equatorial region. In contrast, during El Niño events, a warm SST anomaly in the east is associated with a negative Chl concentration in the west. Quantitatively, the correlations between SST and Chl are shown in Fig. 6a. The inter-annual variability of Chl closely follows that of SST, so the inter-annual variations in Chl tend to be negatively correlated with those in the SST over the tropical Pacific, with large negative values in the western equatorial Pacific near the date line.

### 3.2 Inter-annual variations in $H_m$ and $H_p$

The characteristics of the  $H_m$  structure and its variability are seen in Fig. 2b and the Appendix. The ML is relatively shallow ( $< 30 \text{ m}$ ) in the western tropical Pacific because of

the weak surface friction velocity and stabilizing surface buoyancy flux. Regions with values larger than  $50 \text{ m}$  are found in the central basin, which exhibit energetic surface wind and strong surface buoyancy losses. During La Niña, the ML is unusually deep in the western-central equatorial Pacific but shallow in the eastern regions. During El Niño, the ML becomes shallower in the west but deeper in the east. The inter-annual variations in  $H_m$  exhibit a see-saw pattern, with two large anomaly regions in the western and eastern equatorial regions (Fig. 7b) and the zero-crossing line around  $150^\circ\text{W}$ . As such, inter-annual anomalies of  $H_m$  in the west tend to be out of phase with those in the east, a clear see-saw pattern that is associated with the ENSO's evolution.

The above Chl field is used to estimate the penetration depth ( $H_p$ ). As shown in Fig. 2c and the Appendix, the penetration of solar radiation is deep in the western equatorial Pacific but shallow in the east, with a low  $H_p$  value ( $< 19 \text{ m}$ ) in the east and high  $H_p$  value ( $> 20 \text{ m}$ ) in the west. Seasonally



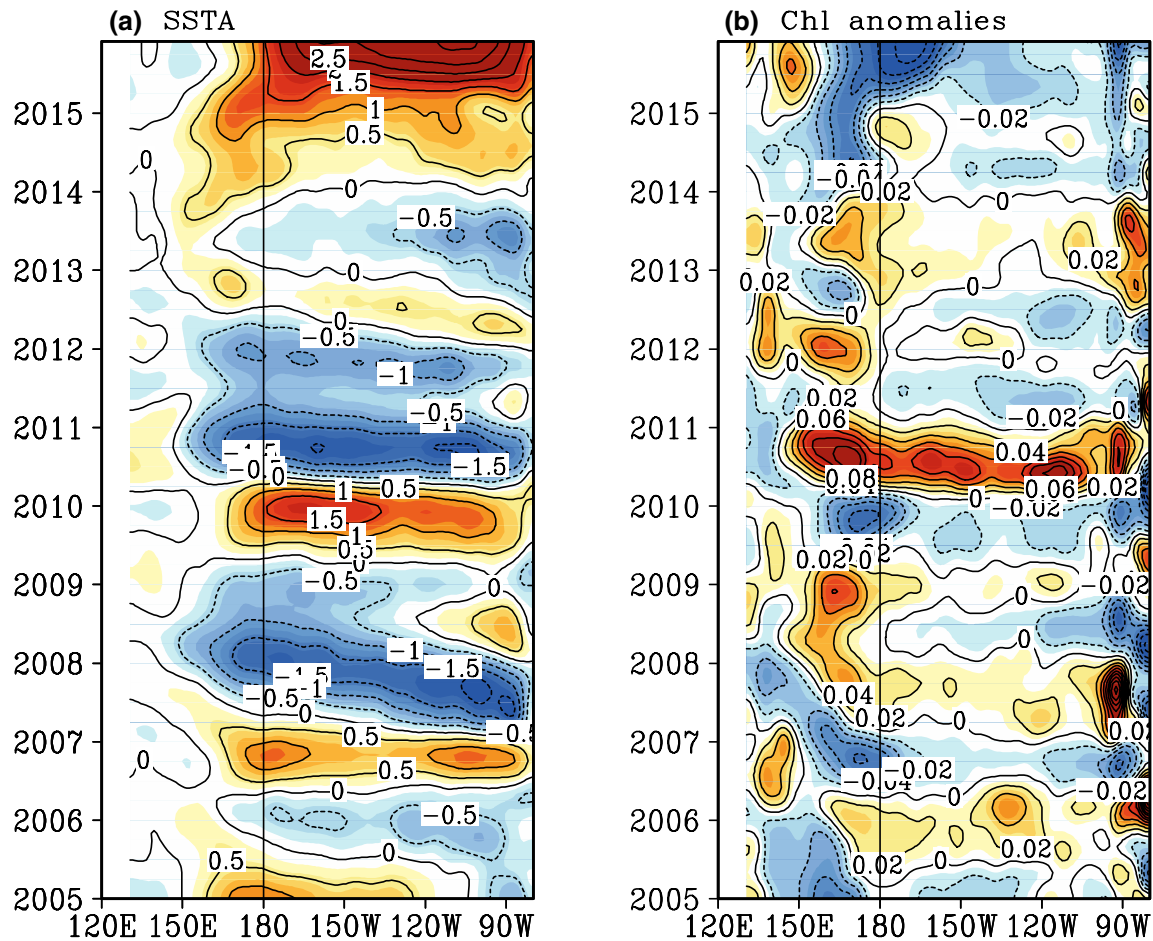
**Fig. 3** Time-longitude sections along the Equator during 2005–2015 for the ocean Chl-related total heating terms: **a**  $Q_{\text{abs}}$ , **b**  $Q_{\text{pen}}$ , and **c**  $R_{\text{sr}}$ . The contour interval is  $10 \text{ W m}^{-2}$  in **a**,  $5 \text{ W m}^{-2}$  in **b** and  $1 \text{ }^{\circ}\text{C month}^{-1}$  in **c**

(see the Appendix), the penetration of solar radiation is deep in spring but shallow in fall. Inter-annually, the penetration depth in the equatorial Pacific exhibits large variations during ENSO cycles (Fig. 7c). Inter-annual  $H_p$  anomalies have a uniform pattern across the equatorial Pacific. For example,  $H_p$  has positive anomalies in the equatorial Pacific during El Niño (corresponding to a reduced Chl concentration with a deeper penetration of solar radiation) but negative anomalies during La Niña (corresponding to an increased Chl concentration with a shallower penetration of solar radiation).

As a response to ENSO cycles, the inter-annual variations in  $H_m$  and  $H_p$  indicate a coherent space–time structure (Figs. 7b, c, 8, 9) clearly display their horizontal patterns during La Niña conditions as represented in August 2010 and during El Niño conditions in August 2015, respectively. One striking feature is that the inter-annual variations in  $H_p$  and  $H_m$  tend to be out of phase in the western-central equatorial Pacific during ENSO cycles but in phase in the east (Figs. 8b, c, 9b, c). That is, in the western equatorial Pacific,

positive (negative)  $H_m$  anomalies are accompanied by negative (positive)  $H_p$  anomalies during La Niña (El Niño). However, in the eastern equatorial Pacific,  $H_p$  anomalies are weak and  $H_m$  has positive and negative anomalies during El Niño and La Niña, respectively.

The standard deviations for the inter-annual anomalies of  $H_p$  and  $H_m$  are shown in Fig. 5b, c. Regions with large inter-annual  $H_p$  variability are seen in the western Pacific west of the date line.  $H_m$  has a large-variability region in the western equatorial region. In the western-central equatorial Pacific near the date line, the magnitude of the inter-annual variability of  $H_p$  is approximately 20% that of  $H_m$ , indicating that the effect of  $H_p$  on the heating terms is not negligible compared to that of  $H_m$ . Quantitatively, the correlations between  $H_m$  and  $H_p$  are shown in Fig. 6c. Corresponding to the negative correlations between  $H_p$  and Chl, positive correlations exist between  $H_p$  and  $H_m$  in the western equatorial Pacific. The regions with large positive correlation near the date line are consistent with those where the amplitude of  $H_p$



**Fig. 4** Time-longitude sections along the Equator during 2005–2015 for inter-annual anomalies of **a** SST and **b** Chl. The contour interval is  $0.5\text{ }^{\circ}\text{C}$  in **a** and  $0.02\text{ mg m}^{-3}$  in **b**

is approximately 20% that of  $H_m$ , indicating that  $H_p$  can have pronounced effects. The effects on the heating terms from the inter-annual variability of  $H_p$  will be analyzed below.

### 3.3 Inter-annual variations in the bio-induced heating terms

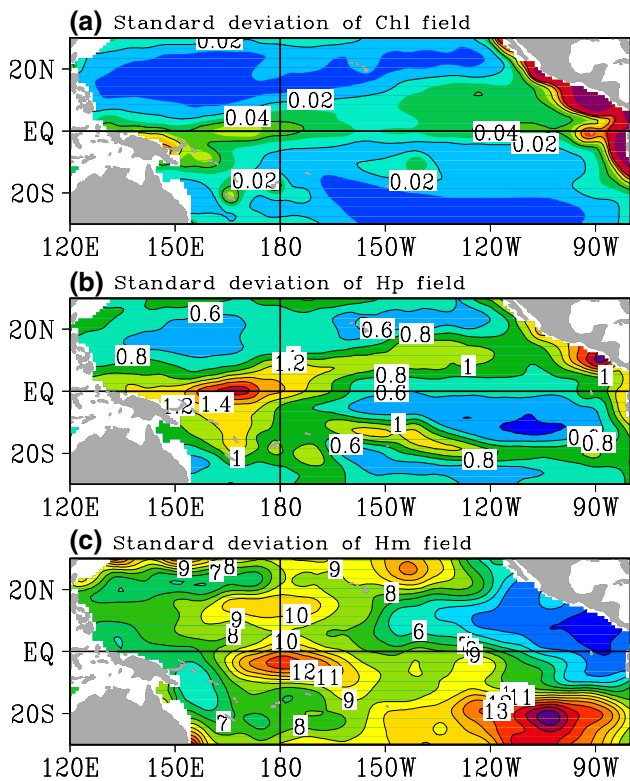
$H_m$  and  $H_p$  are two factors that determine the distribution of solar radiation between the mixed layer and underlying subsurface layers. As mathematically expressed above, the three heating terms are directly associated with  $H_m$  and  $H_p$ . To understand the bio-effects in the tropical Pacific, a diagnostic analysis is performed to quantify the structural relationships of these heating terms with  $H_m$  and  $H_p$ , with emphasis placed on the ENSO. Then, the manner in which these related fields are affected by  $H_p$  can be inferred to reveal how the SST is modulated by the bio-effects and underlying processes that are involved.

Figures 3, 10, 11, and 12 display the space–time evolution of some related fields along the Equator; examples for

the horizontal patterns of inter-annual  $Q_{pen}$  anomalies are demonstrated in Fig. 8 during La Niña conditions and Fig. 9 El Niño conditions, respectively. All these fields exhibit large variability that is associated with the ENSO. The inter-annual variations in  $H_m$  and  $H_p$  exhibit well-defined structures during ENSO cycles, so their effects on the vertical distributions of solar radiation between the ML and subsurface layers are expected, which can be quantified by calculating the modulating effects on these heating terms. For example, the structural relationships of  $Q_{pen}$  with  $H_m$  and  $H_p$  can be clearly seen in the horizontal distributions of their inter-annual anomaly fields during El Niño (Fig. 8) and La Niña (Fig. 9), respectively.

As indicated in the  $H_m$  structure (Fig. 2b), most of the solar radiation is absorbed within the ML ( $Q_{abs}$ ; Fig. 3a), with some penetrating through the bottom of the ML ( $Q_{pen}$ ; Fig. 3b). When the ML changes during ENSO cycles, the amount of solar radiation that is absorbed within the ML ( $Q_{abs}$ ) and penetrates through the bottom of the ML ( $Q_{pen}$ ) also changes. The absorbed component within the ML





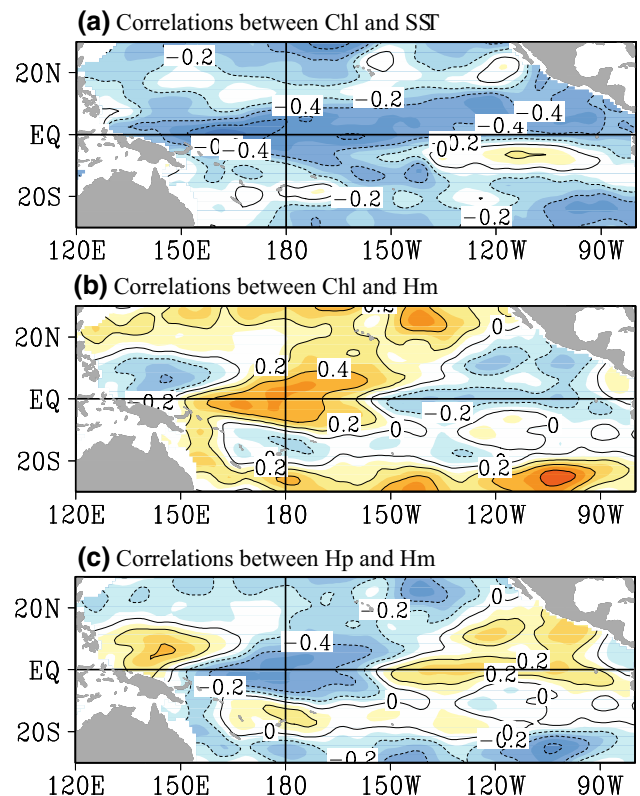
**Fig. 5** Standard deviations for inter-annual anomalies of **a** Chl, **b**  $H_p$  and **c**  $H_m$ . The contour interval is  $0.02 \text{ mg m}^{-3}$  in **a**,  $0.2 \text{ m}$  in **b** and  $1 \text{ m}$  in **c**

directly changes the  $R_{sr}$  value (Fig. 3c), thus modulating the SST, a way to directly change SST. Additionally,  $H_p$  can have a modulating effect on these terms, which is analyzed in detail below.

### 3.3.1 Modulating effect on $Q_{pen}$

The total  $Q_{pen}$  field and corresponding inter-annual  $Q_{pen}$  anomalies along the Equator (Figs. 3b, 10a) indicate their well-defined structure during ENSO cycles. Large  $Q_{pen}$  anomalies are observed in the equatorial regions with a see-saw pattern in the west and east, which is characterized by a positive  $Q_{pen}$  anomaly during El Niño and a negative anomaly during La Niña. In the east, an opposite pattern is seen during ENSO cycles.

$Q_{pen}$  is affected by  $H_m$  and  $H_p$  (Fig. 3), and  $H_m$ 's effect on  $Q_{pen}$  can be clearly seen in the horizontal distributions of the inter-annual anomaly fields during El Niño (Fig. 8) and La Niña (Fig. 9). In the western-central equatorial region, for example, the negative  $Q_{pen}$  anomaly during La Niña is accompanied by a positive  $H_m$  anomaly. An opposite pattern is seen during El Niño (Fig. 9), with a positive  $Q_{pen}$  anomaly accompanied by a negative  $H_m$  anomaly in the western-central region near the date line. The similarity between the

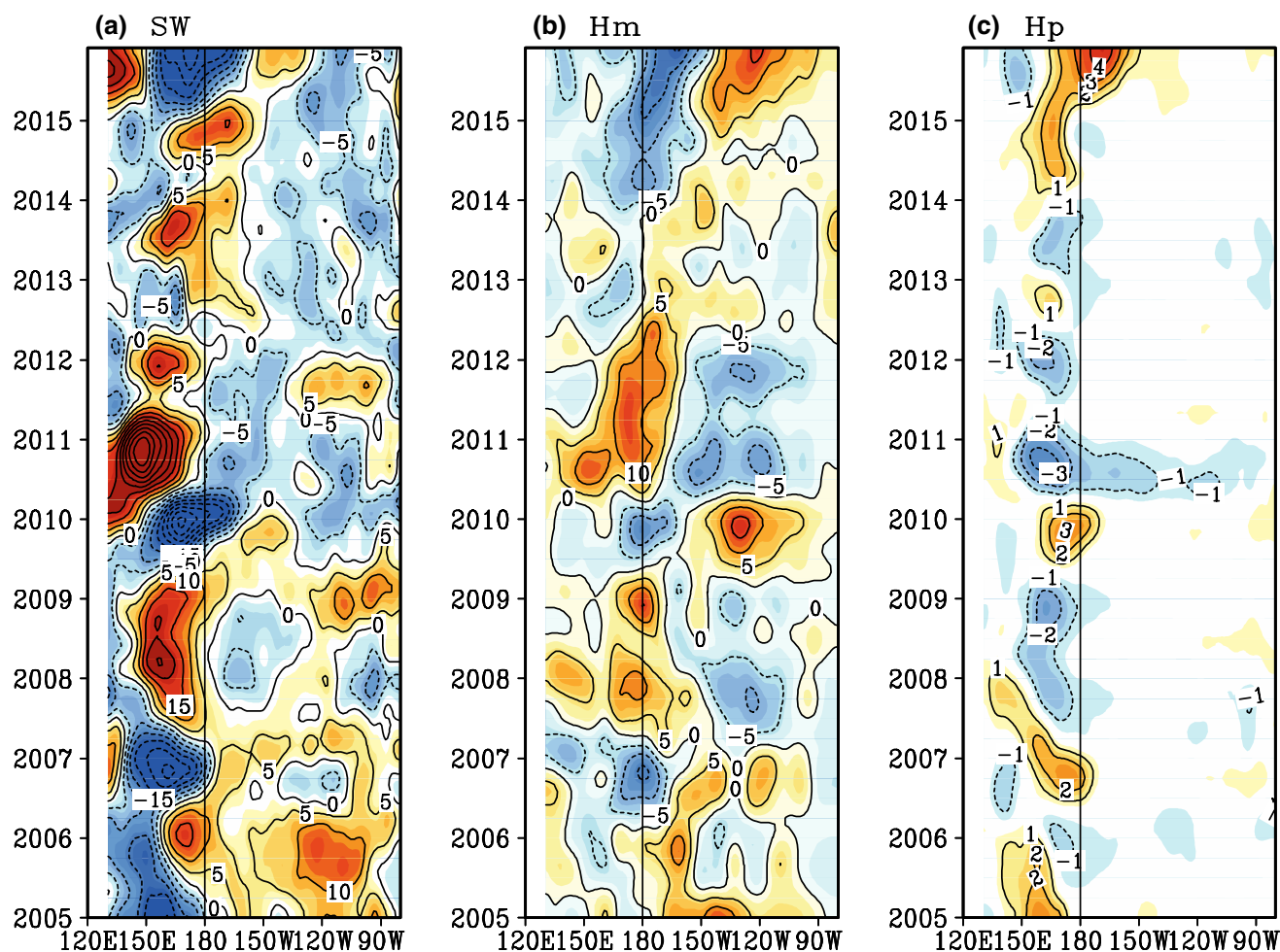


**Fig. 6** Correlations during 2005–2015 between inter-annual anomalies of **a** Chl and SST, **b** Chl and  $H_m$ , and **c**  $H_p$  and  $H_m$ . The contour interval is  $0.2$

space–time evolutions of  $H_m$  (Fig. 7b) and  $Q_{pen}$  (Fig. 10a) indicates that  $H_m$  is a main factor that determines the mean state and variations of  $Q_{pen}$  in the tropical Pacific.

$Q_{pen}$  is also a function of  $H_p$ , so we expect that  $Q_{pen}$  can be modulated by  $H_p$ . As shown in Figs. 7c and 10a, negative and positive  $Q_{pen}$  anomalies have a clear signature of the inter-annual  $H_p$  effects. Thus, a close relationship exists between the inter-annual variations in  $H_p$  and  $Q_{pen}$  in the western-central equatorial basin. A negative  $H_p$  anomaly is observed during La Niña events, which causes less sunlight to penetrate through the bottom of the mixed layer, but become trapped more within the mixed layer. Thus, this negative  $H_p$  anomaly makes the negative  $Q_{pen}$  anomaly during La Niña *more negative*. During El Niño, when  $H_p$  becomes a positive anomaly, sunlight can penetrate deeper into subsurface layers, so the positive  $H_p$  anomaly makes the positive  $Q_{pen}$  anomaly *more positive*.

Clearly, the large observed inter-annual anomalies of  $Q_{pen}$  (e.g., Fig. 10a) represent a combined effect from both  $H_m$  and  $H_p$  during ENSO cycles. Although the inter-annual variability of  $Q_{pen}$  in this region is mainly determined by that of  $H_m$ , the inter-annual variability of  $H_p$  can also make a substantial contribution, enhancing inter-annual anomalies of  $Q_{pen}$  during El Niño–La Niña cycles.



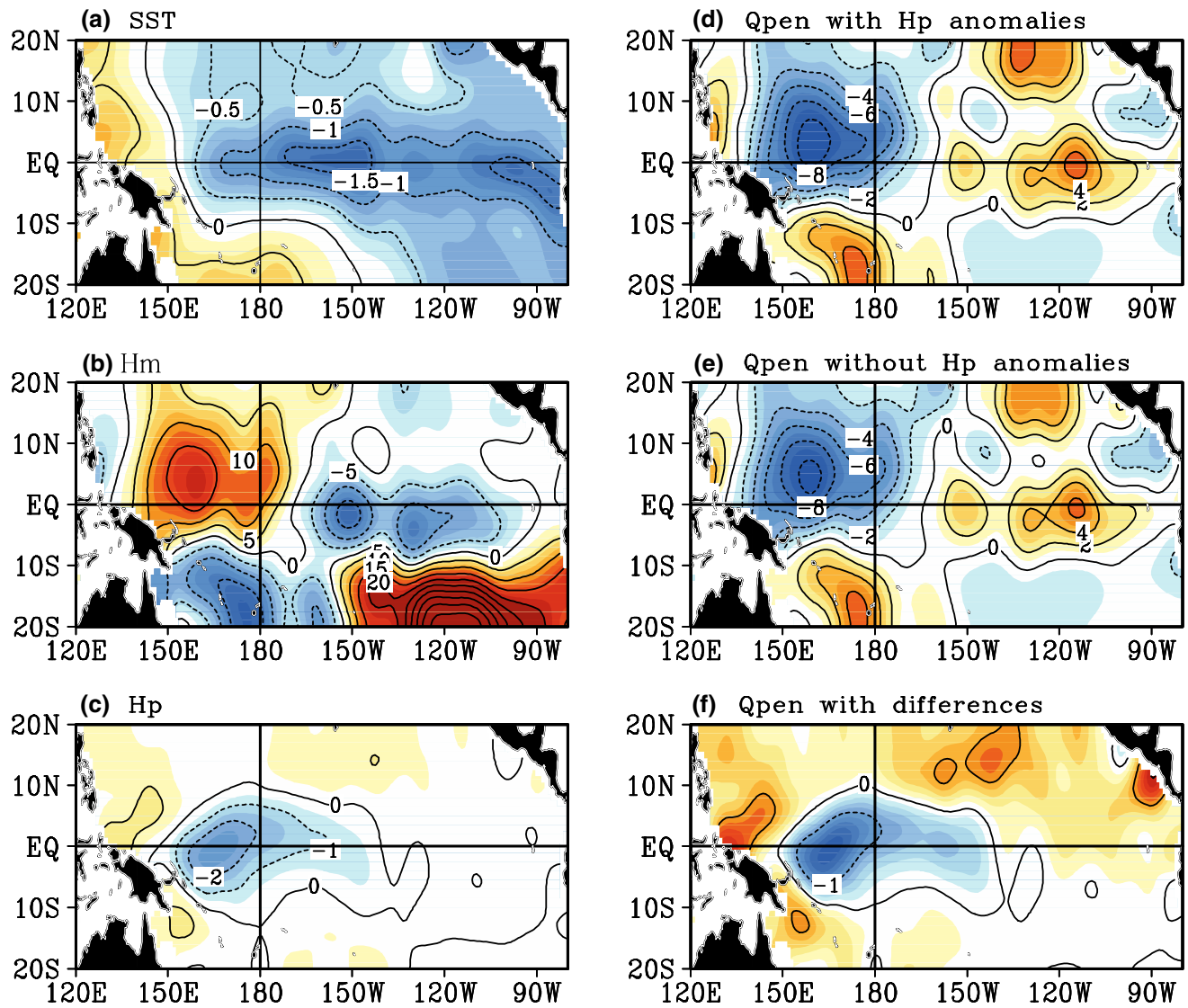
**Fig. 7** Time-longitude sections along the Equator during 2005–2015 for inter-annual anomalies: **a** shortwave radiation, **b**  $H_m$ , and **c**  $H_p$ . The contour interval is  $5 \text{ W m}^{-2}$  in **a**,  $5 \text{ m}$  in **b**, and  $1 \text{ m}$  in **c**

### 3.3.2 Modulating effects on $Q_{\text{abs}}$ and $R_{\text{sr}}$

Next, similar calculations are shown in Figs. 3a and 11 for  $Q_{\text{abs}}$ . Similar to  $Q_{\text{pen}}$ , the inter-annual variability in  $Q_{\text{abs}}$  exhibits a well-defined pattern during ENSO cycles. As defined above, the space–time structural relationships of  $Q_{\text{abs}}$ 's variability with  $H_m$  and  $H_p$  are similar to those of  $Q_{\text{pen}}$  but with the opposite sign. For example, in the western-central equatorial Pacific,  $Q_{\text{abs}}$  exhibits a positive anomaly during La Niña (more gain in solar radiation within the ML) but a negative anomaly during El Niño (loss of solar radiation within the ML), which is determined collectively by  $H_m$  and  $H_p$ . The inter-annual variability of  $Q_{\text{abs}}$  (Fig. 11a) almost mirrors that of  $H_m$  (Fig. 7b), so  $H_m$  is a major factor that controls the structure and variability of  $Q_{\text{abs}}$ . Additionally,  $Q_{\text{abs}}$  can be modulated by  $H_p$  because the former is a function of  $H_p$ . Similar to  $Q_{\text{pen}}$ , a change in  $H_p$  significantly modulates  $Q_{\text{abs}}$  in the western-central equatorial Pacific, and the observed inter-annual variability of  $Q_{\text{abs}}$  also represents the

combined effects of  $H_m$  and  $H_p$ . During La Niña conditions, a positive  $H_m$  anomaly is mainly responsible for a positive  $Q_{\text{abs}}$  anomaly in the region; at the same time,  $H_p$  exhibits a negative anomaly, which increases the absorption of solar radiation within the ML (a corresponding positive  $Q_{\text{abs}}$  anomaly). Thus, the negative  $H_p$  anomaly makes the positive  $Q_{\text{abs}}$  anomaly *more positive*. The combined effects of the positive  $H_m$  anomaly and negative  $H_p$  anomaly enhance the positive  $Q_{\text{abs}}$  anomaly. During El Niño events, the negative  $H_m$  anomaly is responsible for a negative  $Q_{\text{abs}}$  anomaly in the western-central equatorial Pacific; then, the positive  $H_p$  anomaly makes the negative  $Q_{\text{abs}}$  anomaly *more negative*. Thus, the inter-annual variability of  $H_p$  tends to enhance inter-annual anomalies of  $Q_{\text{abs}}$  during ENSO cycles.

Next, we move to the analysis for  $R_{\text{sr}}$  (the rate of the ML temperature change that is directly resulted from the heating effect associated with  $Q_{\text{abs}}$ ). Figures 3c and 12 display the total  $R_{\text{sr}}$  field and its inter-annual variability along the Equator.  $R_{\text{sr}}$  has high values in the west and east but relatively low



**Fig. 8** Horizontal patterns of inter-annual anomaly fields for La Niña conditions as represented in August 2010: **a** SST, **b**  $H_m$ , **c**  $H_p$ , **d**  $Q_{pen}$  when estimated with the inter-annual  $H_p$  effect considered, **e**  $Q_{pen}$  when estimated with the inter-annual  $H_p$  effect excluded, and **f** the

differences (the  $Q_{pen}$  fields estimated with the inter-annual  $H_p$  effect included minus those excluded). The contour interval is  $0.5\text{ }^\circ\text{C}$  in **a**,  $5\text{ m}$  in **b**,  $1\text{ m}$  in **c**,  $2\text{ W m}^{-2}$  in **d** and **e**, and  $1\text{ W m}^{-2}$  in **f**

values in the central basin (Fig. 3c and the Appendix).  $R_{sr}$  also exhibits large inter-annual variability that is associated with ENSO cycles. In the western-central equatorial regions,  $R_{sr}$  has a positive anomaly during El Niño events but a negative anomaly during La Niña events, which is determined by  $H_m$  and  $H_p$ . The inter-annual variability of  $R_{sr}$  (Fig. 12a) follows closely with that of  $H_m$  (Fig. 7b), so  $H_m$  is a major factor that determines the structure and variability of  $R_{sr}$ .  $R_{sr}$  is also a function of  $H_p$ , so we expect that  $R_{sr}$  can be modulated by  $H_p$ , as with  $Q_{abs}$  and  $Q_{pen}$ . However, the manner in which  $R_{sr}$  is affected by  $H_p$  is different from  $Q_{abs}$  and  $Q_{pen}$ , which is explained below.

As expressed by  $R_{sr} = Q_{abs}/(\rho_0 C_p H_m)$ ,  $R_{sr}$  is determined by both  $Q_{abs}$  and  $H_m$ . On the one hand,  $Q_{abs}$ , which is a function of  $H_m$  and  $H_p$ , can be one factor that determines the  $R_{sr}$  value. That is, the absorbed solar radiation within the ML ( $Q_{abs}$ ) directly affects the rate of the ML temperature change, so  $R_{sr}$  and  $Q_{abs}$  can have a coherent variation and can be modulated by  $H_p$  in a similar fashion. In the western-central region, for example, the positive  $H_p$  anomaly during El Niño decreases the absorption of solar radiation within the ML (a corresponding negative  $Q_{abs}$  anomaly), which would produce a negative  $R_{sr}$  anomaly with a negative correlation between  $H_p$  and  $R_{sr}$ . However, this scenario does not occur. Examining the relationship among the inter-annual variations in  $H_p$

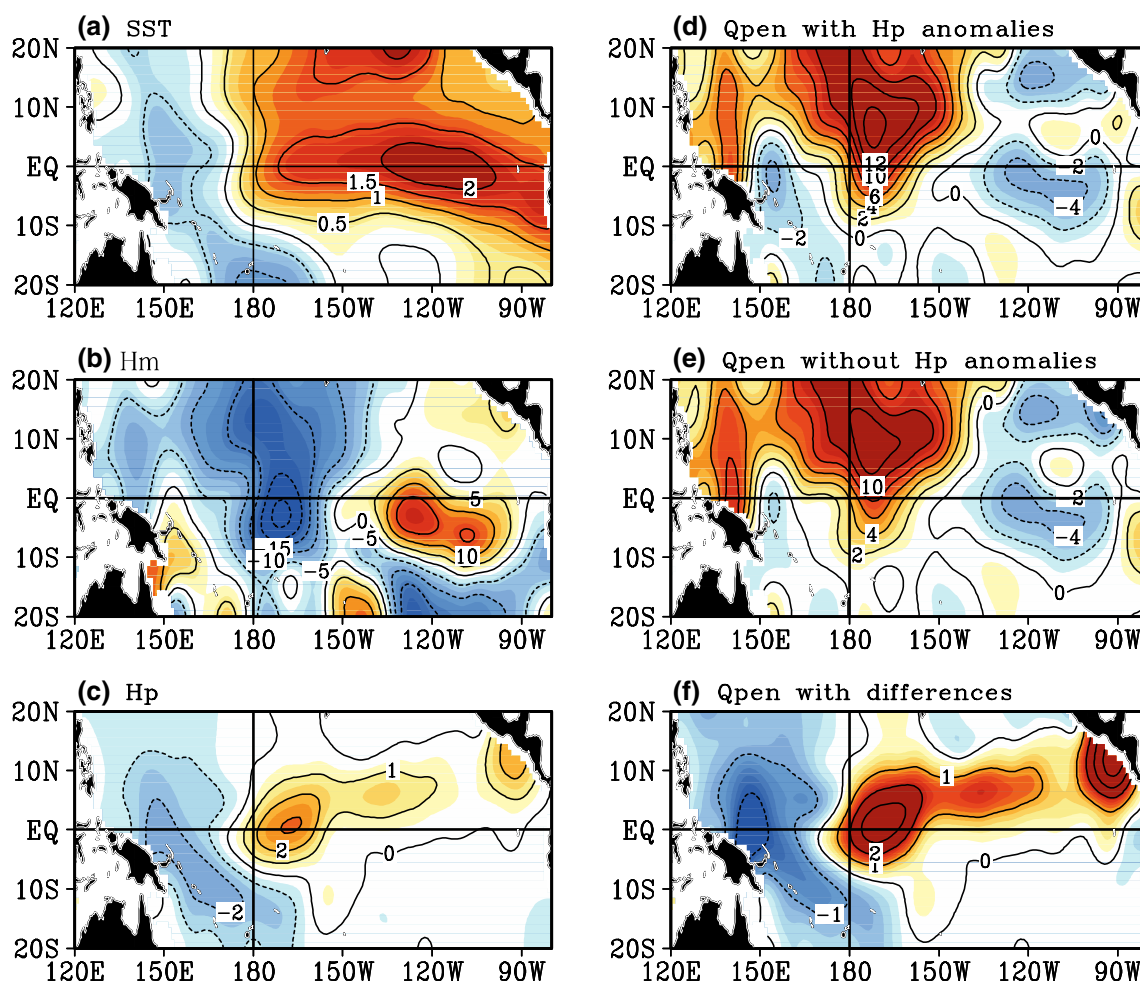


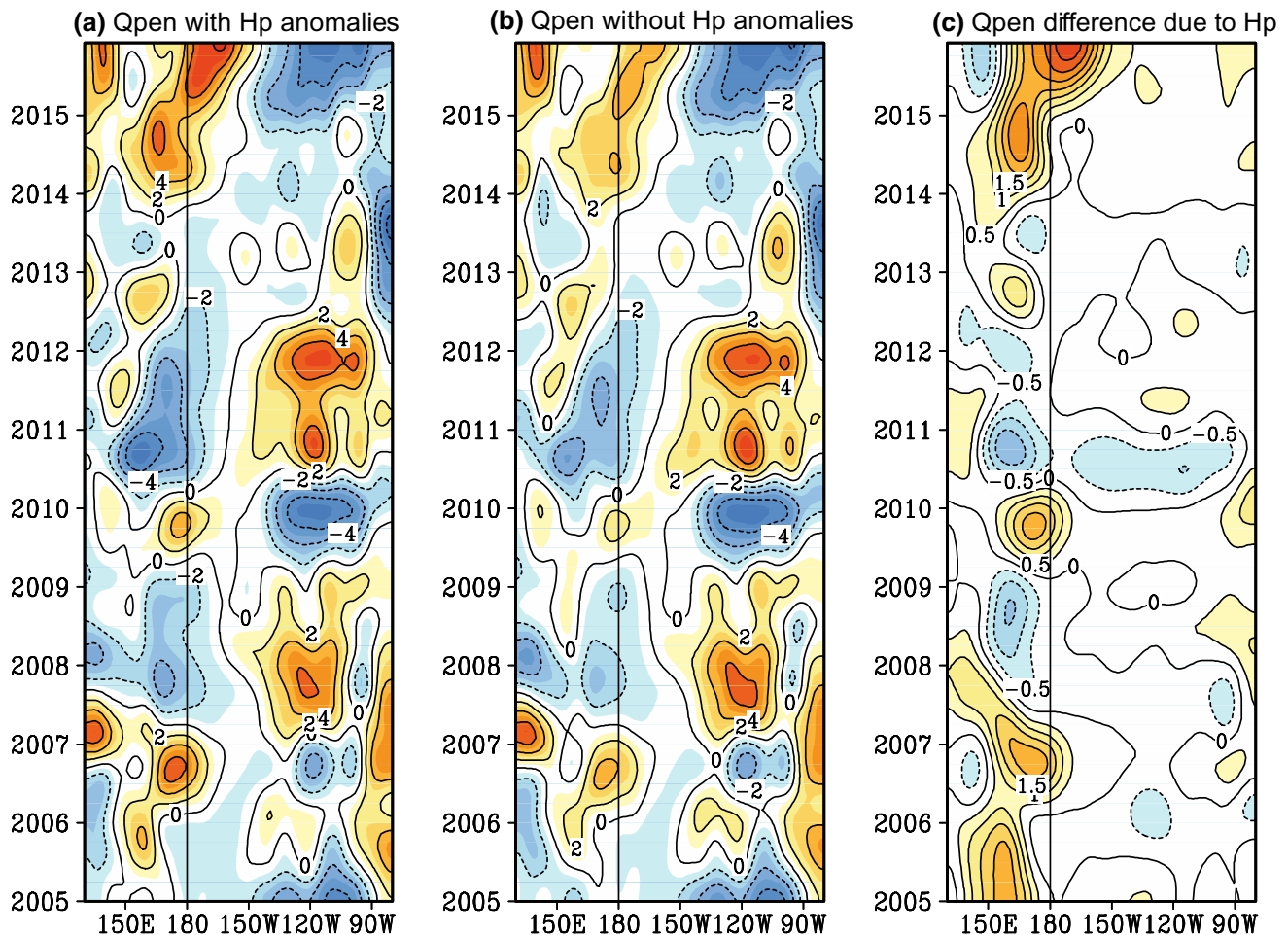
Fig. 9 Same as in Fig. 8 but for El Niño conditions as represented in August 2015

(Fig. 7c),  $Q_{\text{abs}}$  (Fig. 11a), and  $R_{\text{sr}}$  (Fig. 12a) indicates that the inter-annual variations in  $R_{\text{sr}}$  do not follow those in  $Q_{\text{abs}}$ . In fact, their variations in  $Q_{\text{abs}}$  and  $R_{\text{sr}}$  tend to be out of phase in the western-central regions during ENSO cycles. This result indicates that  $R_{\text{sr}}$  is not determined by  $Q_{\text{abs}}$  and that the manner in which  $Q_{\text{abs}}$  is modulated by  $H_p$  is not reflected in  $R_{\text{sr}}$ . This relationship indicates that  $R_{\text{sr}}$  is not affected by  $H_p$  in a coherent fashion.

On the other hand,  $R_{\text{sr}}$  is closely related to  $H_m$  because of their inverse relationship. As such, the manner in which  $R_{\text{sr}}$  is modulated by  $H_p$  can be complicated, which is different from  $Q_{\text{abs}}$ . Indeed, another factor (the inverse relationship of  $R_{\text{sr}}$  with  $H_m$ ) must be considered, which can play an important role in determining the nature of the inter-annual variability of  $R_{\text{sr}}$ . We provide additional explanations below.

The mathematical expressions in Eq. 3 indicate that  $R_{\text{sr}}$  is proportional to  $Q_{\text{abs}}$  (which increases exponentially with  $H_m$  but decreases exponentially with  $H_p$ ); additionally,  $R_{\text{sr}}$  is inversely proportional to  $H_m$  (appearing as a denominator in  $R_{\text{sr}} = Q_{\text{abs}} / (\rho_0 C_p H_m)$ ). As such,  $H_m$  has two effects

on  $R_{\text{sr}}$ : one is exponential through  $Q_{\text{abs}}$  and the other is an inverse relationship with  $H_m$ . Thus, the manner in which  $R_{\text{sr}}$  is affected by  $H_p$  (which is explicitly represented through  $Q_{\text{abs}}$ ) is complicated by the effect of inter-annual  $H_m$  variations. During an El Niño event, for example,  $H_p$  exhibits a positive anomaly in the western-central regions, which decreases the solar radiation within the ML (a decrease in  $Q_{\text{abs}}$ ); at the same time, the ML tends to be anomalously shallow (a negative  $H_m$  anomaly, which decreases  $Q_{\text{abs}}$ ). This decreased  $Q_{\text{abs}}$  field (which is caused by the effect from the positive  $H_p$  anomaly) now acts on this anomalously shallow ML, so the effect on  $R_{\text{sr}}$  ( $= Q_{\text{abs}} / (\rho_0 C_p H_m)$ ) from the reduced  $Q_{\text{abs}}$  is offset by that from the shoaling ML. The combined net effects of the reduced  $Q_{\text{abs}}$  component and negative  $H_m$  anomaly produce an  $R_{\text{sr}}$  value that does not change much. Therefore,  $R_{\text{sr}}$  is not modulated by positive  $H_p$  anomalies as significantly as  $Q_{\text{abs}}$  is during El Niño, and the sign of induced changes in  $R_{\text{sr}}$  is not consistent with what can be expected from the inter-annual variability of  $Q_{\text{abs}}$ .



**Fig. 10** Time-longitude sections along the Equator during 2005-2015 for inter-annual anomalies: **a**  $Q_{pen}$  when estimated with the inter-annual  $H_p$  effect considered, **b**  $Q_{pen}$  when estimated with the inter-

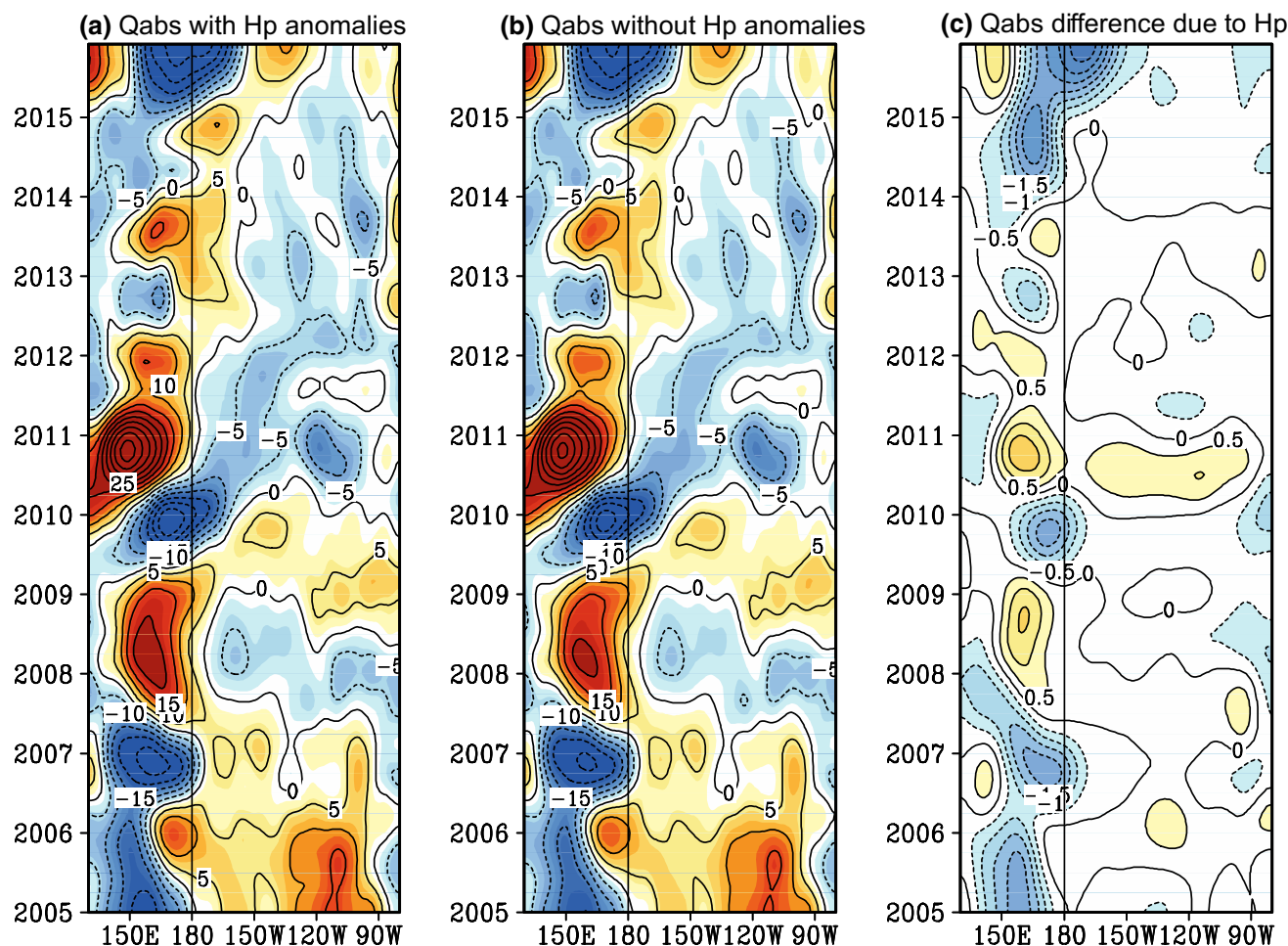
annual  $H_p$  effect excluded, and **c** their differences (the  $Q_{pen}$  fields estimated with the inter-annual  $H_p$  effect included minus those excluded). The contour interval is  $2 \text{ W m}^{-2}$  in **a** and **b** and  $0.5 \text{ W m}^{-2}$  in **c**

The relationships among these fields are further quantified in Fig. 13, which presents the anomaly correlations between  $H_p$  and  $Q_{pen}$  and those between  $H_p$  and  $R_{sr}$ , respectively. The correlation between  $H_p$  and  $Q_{abs}$  is exactly the same as that with  $Q_{pen}$  but with the opposite sign and is not shown here. Evidently, the extent to which  $R_{sr}$  is correlated with  $H_p$  is different from what could be inferred from the manner in which  $Q_{abs}$  is modulated by  $H_p$ , whereas  $H_p$  and  $Q_{pen}$  ( $Q_{abs}$ ) indicate a high positive correlation in the western-central equatorial Pacific. One striking feature of  $H_p$  and  $R_{sr}$  is that these terms are weakly positively correlated in the west, exhibiting the opposite sign to what is expected from the effect of  $H_p$  on  $Q_{abs}$  (i.e., in terms of the relationship with  $Q_{abs}$ ,  $R_{sr}$  would have a negative correlation with  $H_p$ ). These results indicate that the extent to which  $R_{sr}$  is affected by  $H_p$  is different from  $Q_{abs}$  and  $Q_{pen}$ . The fact that  $R_{sr}$  is not significantly modulated by  $H_p$  has implications for the manner in which SST is modulated by bio-effects.

### 3.4 Further attributions to the inter-annual $H_p$ effects

To more clearly understand the bio-effects, further analyses are performed by isolating the contribution from the inter-annual  $H_p$  anomalies. The  $H_p$  field can be separated into its seasonally varying climatological component ( $\overline{H_p}$ ) and inter-annual component ( $H'_p$ ); their direct effects on the three heating terms are then explicitly quantified by re-calculating these heating terms (say,  $Q_{pen}$ ) in two manners, namely, with  $H_p$  taken as its seasonally varying climatology (denoted as  $Q_{pen}(H_m, \overline{H_p})$ ) and with  $H_p$  considered as varying inter-annually (denoted as  $Q_{pen}(H_m, H'_p)$ ). Correspondingly, the effects of the inter-annual  $H_p$  anomalies on  $Q_{pen}$  can be estimated by calculating the differences between  $Q_{pen}(H_m, \overline{H_p})$  and  $Q_{pen}(H_m, H'_p)$ .

Figure 10b, c display the time-longitude sections along the Equator for  $Q_{pen}$ ; the horizontal patterns for the  $Q_{pen}$  difference during El Niño and La Niña are shown in Figs. 8 and



**Fig. 11** Time-longitude sections along the Equator during 2005–2015 for inter-annual anomalies: **a**  $Q_{\text{abs}}$  when estimated with the inter-annual  $H_p$  effect considered, **b**  $Q_{\text{abs}}$  when estimated with the inter-

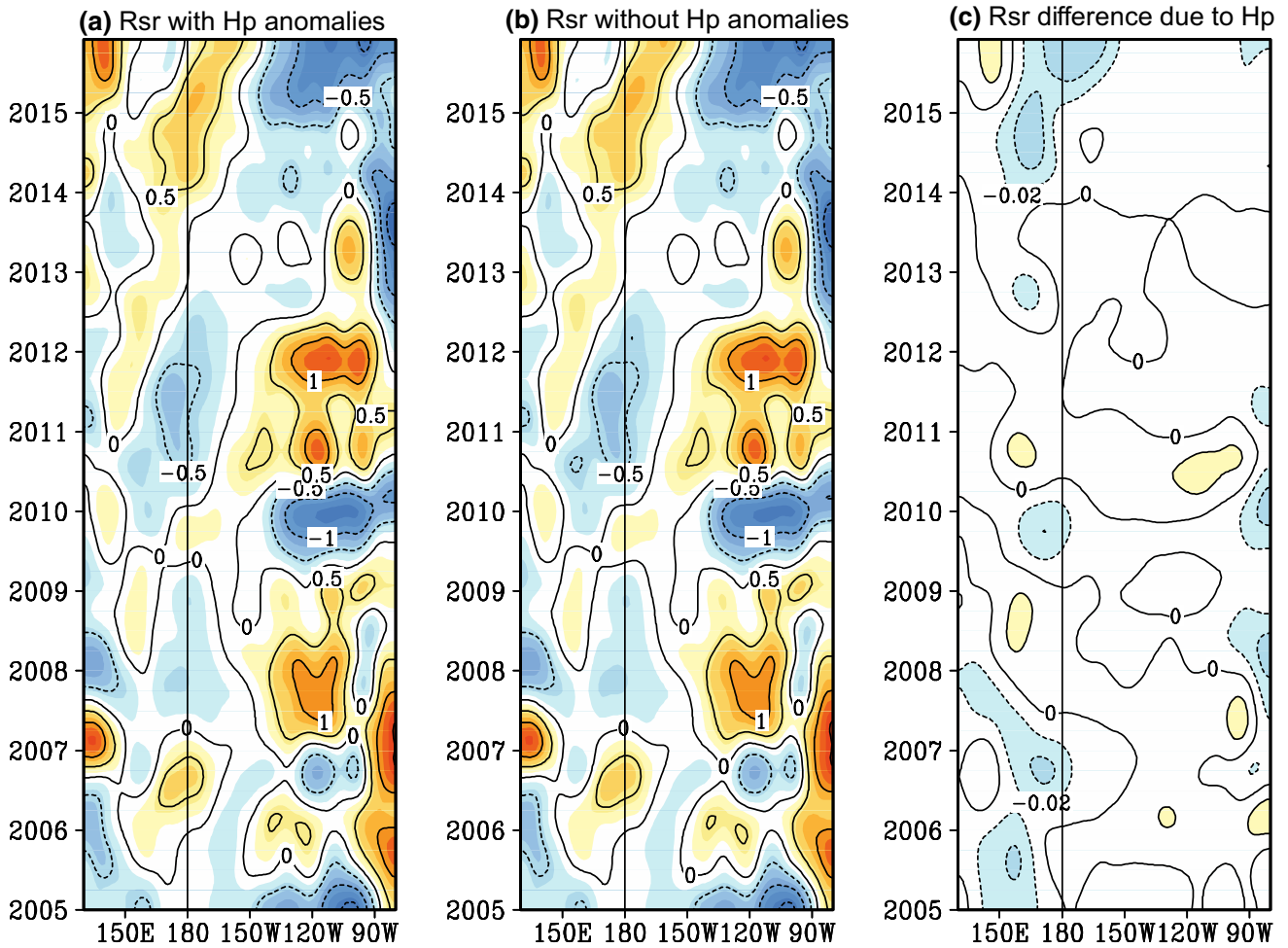
annual  $H_p$  effect excluded, and **c** their differences (the  $Q_{\text{abs}}$  fields estimated with the inter-annual  $H_p$  effect included minus those excluded). The contour interval is  $5 \text{ W m}^{-2}$  in **a** and **b** and  $0.5 \text{ W m}^{-2}$  in **c**

9. The effects of the inter-annual  $H_p$  anomalies on  $Q_{\text{pen}}$  are mostly pronounced in the western-central equatorial Pacific. High similarities in the spatial patterns and temporal variations are seen between the inter-annual variations in  $H_p$  and the  $Q_{\text{pen}}$  difference. Their relationships indicate that  $Q_{\text{pen}}$  is significantly modulated by  $H_p$  in the western-central regions, where a large variability of  $H_p$  exists. The high similarity between  $H_p$  and the  $Q_{\text{pen}}$  difference is consistent with the high positive correlations between the inter-annual variations in  $H_p$  and  $Q_{\text{pen}}$  (Fig. 13).

A similar analysis is performed for  $Q_{\text{abs}}$ , and the corresponding results are shown in Fig. 11b, c. As defined with  $Q_{\text{abs}}$  and  $Q_{\text{pen}}$ , the effects of the inter-annual  $H_p$  anomalies on  $Q_{\text{abs}}$  are the same as for  $Q_{\text{pen}}$  but with the opposite sign. Good agreement exists between the inter-annual variations in  $H_p$  and the  $Q_{\text{abs}}$  difference, indicating that  $Q_{\text{abs}}$  is significantly modulated by  $H_p$  in the western-central equatorial regions.

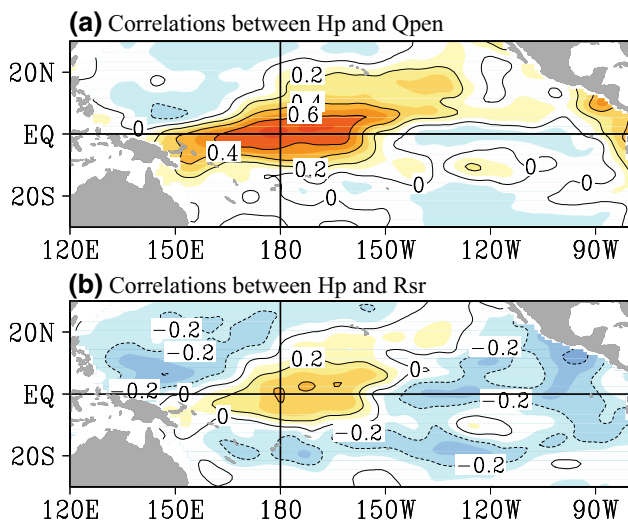
Next, we analyze  $R_{\text{sr}}$ , in which the effects on  $R_{\text{sr}}$  are estimated by taking  $H_p$  as its seasonally varying climatology (denoted as  $R_{\text{sr}}(H_m, \overline{H_p})$ ) and varying inter-annually (denoted as  $R_{\text{sr}}(H_m, H_p)$ ). Figure 12b, c show the time-longitude sections along the Equator for inter-annual anomalies of  $R_{\text{sr}}(H_m, \overline{H_p})$  and the differences. The relationships between  $H_p$  and the difference field in  $R_{\text{sr}}$  are different from those in  $Q_{\text{pen}}$ . A large inter-annual anomaly of  $H_p$  does not produce a correspondingly large difference in  $R_{\text{sr}}$ , indicating that the  $H_p$ -induced modulating effect on  $R_{\text{sr}}$  is weak. The sign of the  $R_{\text{sr}}$  difference is even opposite to what is expected from the inter-annual  $H_p$  effect on  $Q_{\text{abs}}$ . This result is consistent with the low correlation between  $H_p$  and  $R_{\text{sr}}$ .

Table 1 quantifies the modulating effects of the inter-annual  $H_p$  anomalies by calculating the standard deviation of some related variables in the Niño4 region ( $160^\circ\text{E}$ – $150^\circ\text{W}$ ,  $5^\circ\text{S}$ – $5^\circ\text{N}$ ). Although the amplitude of the inter-annual variability of  $H_p$  is approximately 13% that of  $H_m$  in the Niño4



**Fig. 12** Time-longitude sections along the Equator for inter-annual anomalies: **a**  $R_{sr}$  when estimated with the inter-annual  $H_p$  effect considered, **b**  $R_{sr}$  when estimated with the inter-annual  $H_p$  effect

excluded, and **c** their differences (the  $R_{sr}$  fields estimated with the inter-annual  $H_p$  effect included minus those excluded). The contour interval is  $0.5\text{ }^{\circ}\text{C month}^{-1}$  in **a** and **b** and  $0.02\text{ }^{\circ}\text{C month}^{-1}$  in **c**



**Fig. 13** Correlations during 2005–2015 between inter-annual anomalies of **a**  $H_p$  and  $Q_{pen}$  and **b**  $H_p$  and  $R_{sr}$ . The contour interval is 0.2

region, its contribution to the inter-annual variability of  $Q_{pen}$  is approximately 28%, whereas that to  $R_{sr}$  is only 3%. One interesting feature is that the extent to which  $R_{sr}$  is modulated by the inter-annual  $H_p$  variability is different from the extent to which  $Q_{pen}$  and  $Q_{abs}$  are. Quantitatively, the correlation between the inter-annual anomalies of  $H_p$  and  $Q_{pen}$  is highly positive, and that between  $H_p$  and  $R_{sr}$  is weakly positive (Fig. 13), indicating that  $R_{sr}$  is not modulated as significantly as  $Q_{abs}$  and  $Q_{pen}$ . The differences in the extent to which these heating terms are modulated by  $H_p$  have important implications for the manner in which SST is modulated by bio-effects and the mechanism that is involved.

### 3.5 Mechanism by which SST is modulated by the OBH feedback

The effects of ocean biology-induced heating on the penetrative solar radiation are realized through the modulations of

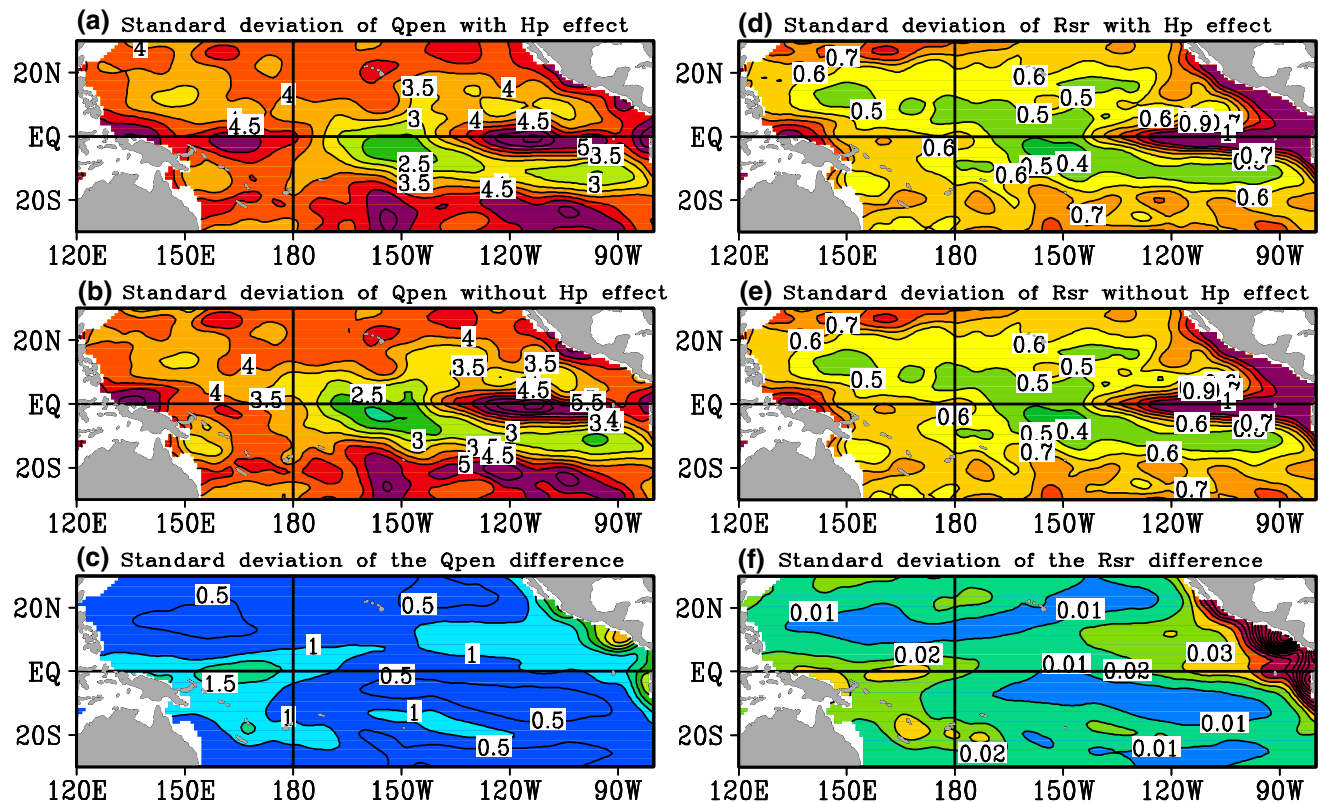
**Table 1** Standard deviation of variables in the Niño4 region (160°E–150°W, 5°S–5°N)

Variables	$H_m$ (m)	$H_p$ (m)	$Q_{pen}$ with $H_p$ effect ( $W m^{-2}$ )	$Q_{pen}$ without $H_p$ effect ( $W m^{-2}$ )	$R_{sr}$ with $H_p$ effect ( $^{\circ}C month^{-1}$ )	$R_{sr}$ without $H_p$ effect ( $^{\circ}C month^{-1}$ )
Standard deviation	6.40	0.81	2.55	1.99	0.32	0.33

the three heating terms, which are explicitly related to  $H_m$  and  $H_p$ . The revealed relationships among these related fields hint at the processes that are involved in the effects of the OBH-related feedback. Two influence pathways are possible by which SST can be modulated by inter-annual variations in  $H_p$  in the equatorial Pacific, depending on which heating term is predominantly affected. One is a direct influence pathway by which SST can be modulated through a direct gain or loss in solar radiation within the ML, which is indicated by  $Q_{abs}$  and  $R_{sr}$ . If a significant change to  $R_{sr}$  is induced from  $Q_{abs}$  in association with inter-annual  $H_p$  anomalies, a direct modulating effect on SST is indicated to be important. In this case, the ML is warmed up or cooled down directly by the  $H_p$ -induced  $Q_{abs}$  contribution. Then, the bio-feedback is considered to be realized primarily through a direct thermal effect on the SST, which should be reflected in  $R_{sr}$ .

Another is an indirect influence pathway, by which  $Q_{pen}$  and  $Q_{abs}$  are significantly modulated by  $H_p$ , but  $R_{sr}$  is not. Then, differential heating is produced by  $H_p$  vertically between the ML ( $Q_{abs}$ ) and subsurface layers ( $Q_{pen}$ ), which modifies the vertical thermal contrast, stratification and vertical mixing in the upper ocean, thus affecting the SST. In this case,  $Q_{pen}$  and  $Q_{abs}$  are significantly modulated by  $H_p$  but  $R_{sr}$  does not show a coherent change with  $Q_{abs}$ , indicating that the direct thermal effect on the SST is not a predominant mechanism for SST modulation by  $H_p$ . As such, an indirect dynamical effect on the SST that is associated with  $Q_{pen}$  and  $Q_{abs}$  is the dominant process. Therefore, the extent to which these heating terms are modulated by  $H_p$  can indicate the importance of a direct thermal effect or indirect dynamical effect on the SST (Fig. 14).

Here, one striking feature revealed from observational data is that inter-annual variations in  $R_{sr}$  are not significantly



**Fig. 14** Standard deviations for inter-annual anomalies of  $Q_{pen}$  (left panels) and  $R_{sr}$  (right panels), which were estimated with the inter-annual  $H_p$  effect (a, d) considered and (b, e) excluded, alongside (e,

f) their differences. The contour interval is  $0.5 W m^{-2}$  in a–c,  $0.1 ^{\circ}C month^{-1}$  in d and e, and  $0.01 ^{\circ}C month^{-1}$  in f



modulated by  $H_p$ . That is,  $R_{sr}$  does not exhibit a corresponding positive anomaly, as would be expected from the effect of inter-annual  $H_p$  anomalies on  $Q_{abs}$  in the western-central regions. The effects of  $H_p$  on  $R_{sr}$  cannot explain the modulating effect on the SST because the sign of the modulating effect on  $R_{sr}$  is opposite to what is expected from the inter-annual  $H_p$  effect on  $Q_{abs}$ . Thus, bio-effects on the SST are not realized through the  $R_{sr}$  term. Other processes that are involved in the effects on  $Q_{pen}$  and  $Q_{abs}$  must play a dominantly important role in modulating the SST. Indeed, the inter-annual variability of  $H_p$  acts to enhance that of  $Q_{abs}$  and  $Q_{pen}$  during ENSO cycles. The  $H_p$ -enhanced anomalies of  $Q_{abs}$  and  $Q_{pen}$  produce large differential heating between the ML and subsurface layers, which modulates the vertical thermal contrast, stratification and vertical mixing, which represent a dominant indirect ocean dynamical effect on the SST.

#### 4 Summary and discussion

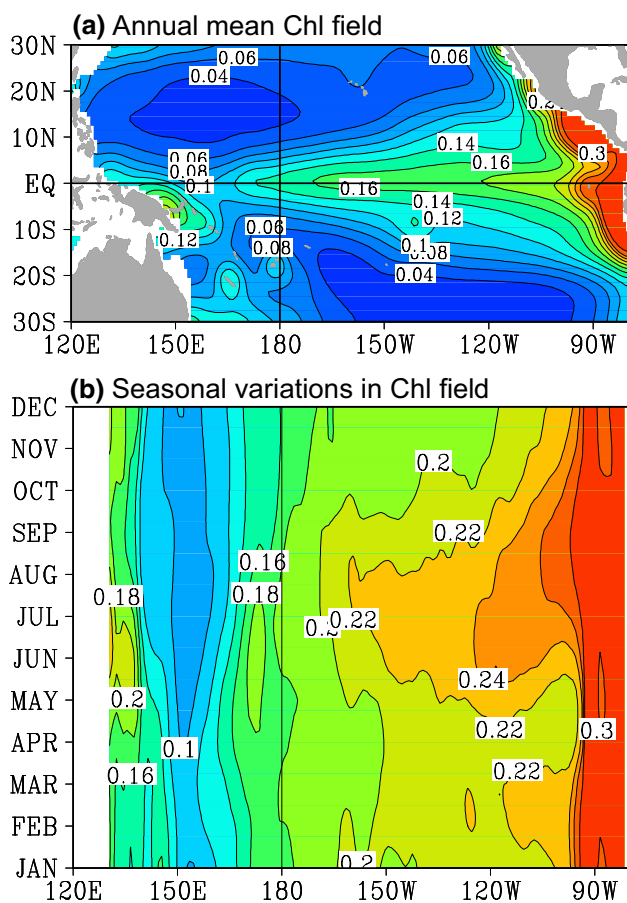
In this paper, observed data were used to characterize the inter-annual variabilities of related physical and biological oceanic fields in the tropical Pacific. In terms of ocean biology, Chl anomalies appeared very quickly and almost simultaneously as a response to ENSO cycles. For example, large inter-annual variations in Chl were concentrated in the western-central equatorial Pacific near the date line: low-Chl-concentration regions extended eastward across the date line during El Niño but retreated westward during La Niña. These inter-annual variations in Chl both represented a response to the ENSO and exhibited feedback onto the ENSO. Here, Chl was considered a major component that affected the vertical penetration of solar radiation in the upper ocean and represented the bio-heating feedback; then, Chl was used to derive the penetration depth ( $H_p$ ) and quantify the effect on the penetrative solar radiation. In terms of physical factors,  $H_m$  was estimated from Argo products. As with  $H_p$ ,  $H_m$  exhibited large inter-annual anomalies in response to ENSO cycles. The incoming shortwave radiation decayed exponentially with depth in the upper ocean, so  $H_m$  was a major factor that controlled the penetration of solar radiation in the upper ocean. Thus,  $H_m$  and  $H_p$  are two factors that affect the distribution of solar radiation between the ML and the underlying subsurface layers.

To quantify the effects from inter-annual  $H_p$  anomalies, three heating terms were explicitly expressed as a function of  $H_p$  and  $H_m$  (Zhang 2015). As previously demonstrated by modeling studies, the biological conditions can affect the mean climate and inter-annual variability over the tropical Pacific through modulating effects on these heating terms. The structural relationships of  $H_m$  and  $H_p$  with the related heating terms were analyzed to gain insight into Chl-induced

heating effects on the ENSO in nature. Furthermore, the revealed relationships from observations could be inferred to determine the processes and mechanisms for ENSO modulations.

The ENSO can induce large perturbations to  $H_p$  and  $H_m$ , whose inter-annual variations exhibited a well-defined structure across the tropical Pacific.  $H_p$  and  $H_m$  directly affected the three heating terms ( $Q_{pen}$ ,  $Q_{abs}$  and  $R_{sr}$ ), which underwent inter-annual variations that were dominated by ENSO signals. A combined modulating effect from  $H_m$  and  $H_p$  was seen on  $Q_{pen}$ . In the western-central equatorial region, the inter-annual variations in  $H_p$  were large and have an out-of-phase relationship with those in  $H_m$  during ENSO cycles. The effects on  $Q_{pen}$  induced by inter-annual anomalies of  $H_p$  tended to be of the same sign as those of  $H_m$ , with a high positive correlation between  $H_p$  and  $Q_{pen}$ . During El Niño, for example,  $Q_{pen}$  exhibited a positive anomaly, which can be attributed to the negative  $H_m$  anomaly. At this time,  $H_p$  was positive. The effect of this positive  $H_p$  anomaly caused the positive  $Q_{pen}$  anomaly during El Niño to become more positive. During La Niña,  $Q_{pen}$  became a negative anomaly in association with a positive  $H_m$  anomaly. At the same time,  $H_p$  became a negative anomaly, which caused the negative  $Q_{pen}$  anomaly to become more negative. So, the inter-annual anomalies of  $H_p$  enhanced those of  $Q_{pen}$ . In terms of  $Q_{abs}$ , the  $H_p$  effect was the same as with  $Q_{pen}$  but with the opposite sign. That is, in the western-central equatorial Pacific, the effects of the positive  $H_p$  anomaly during El Niño made the negative  $Q_{abs}$  more negative, whereas those of the negative  $H_p$  anomaly made the positive  $Q_{abs}$  more positive during La Niña. Thus, the inter-annual  $H_p$  anomalies enhanced the differential heating between the ML ( $Q_{abs}$ ) and subsurface layers ( $Q_{pen}$ ). In terms of the effect on  $R_{sr}$ , the extent to which  $R_{sr}$  was modulated by  $H_p$  was strikingly different from  $Q_{pen}$  and  $Q_{abs}$ . Although  $Q_{abs}$  and  $Q_{pen}$  were significantly modulated by  $H_p$ ,  $R_{sr}$  was not. So, the modulating effects of  $H_p$  on the SST were not realized through the  $R_{sr}$  term and a direct thermal effect from  $H_p$  through  $Q_{abs}$  was not a major process that modulated the SSTs. In contrast,  $Q_{pen}$  and  $Q_{abs}$  were significantly modulated by  $H_p$  in the western-central equatorial Pacific. As such, the  $H_p$ -enhanced anomalies of  $Q_{pen}$  and  $Q_{abs}$  produced large differential heating between the ML and subsurface layers, which modified the vertical thermal contrast, stratification and vertical mixing, thus affecting the SST. Thus, the modulating effects of  $H_p$  on the SST can be traced to the  $Q_{pen}$  and  $Q_{abs}$  terms, and bio-modulations of the ENSO were realized through an indirect dynamical effect on the SST in association with inter-annual  $H_p$  anomalies. Furthermore, the structural relationships indicated a negative feedback on the ENSO.

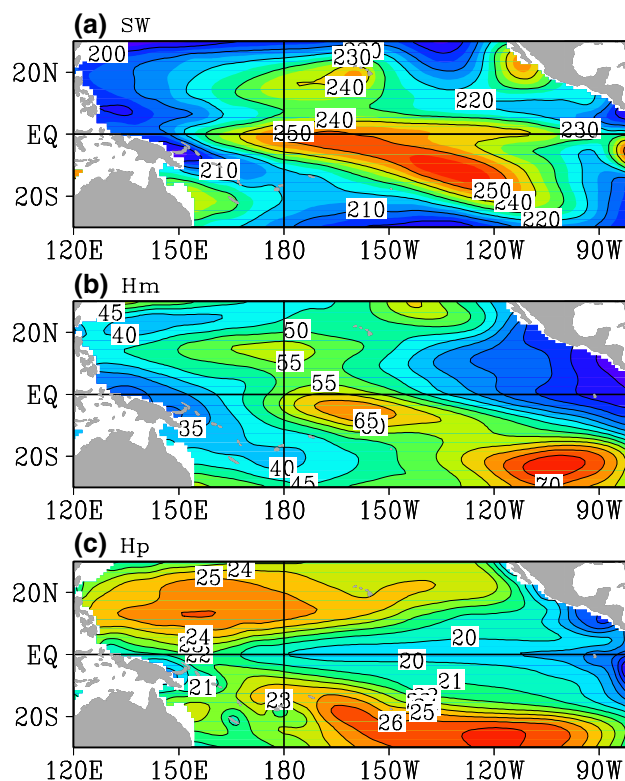
In this study, satellite-based Chl data were combined with Argo data to estimate inter-annual variations in  $H_p$  and  $H_m$ . Furthermore, the derived  $H_p$  and  $H_m$  fields were used to



**Fig. 15** Estimated Chl fields during 2005–2015 from satellite observations: **a** horizontal distribution for the annual mean climatology and **b** seasonal variations along the Equator. The contour interval is  $0.02 \text{ mg m}^{-3}$  in **a** and **b**

reveal their effects on ocean biology-induced heating terms and the modulating effects on the SST. The added values of satellite-Chl observations were clearly demonstrated in association with the use of in situ data. This use of satellite and in situ data elucidated the observed structural relationships among the related fields and processes that were involved in the bio-effects. In particular, an indirect dynamical effect on the SST from the inter-annual Chl variability could explain a negative feedback on the ENSO in the tropical Pacific.

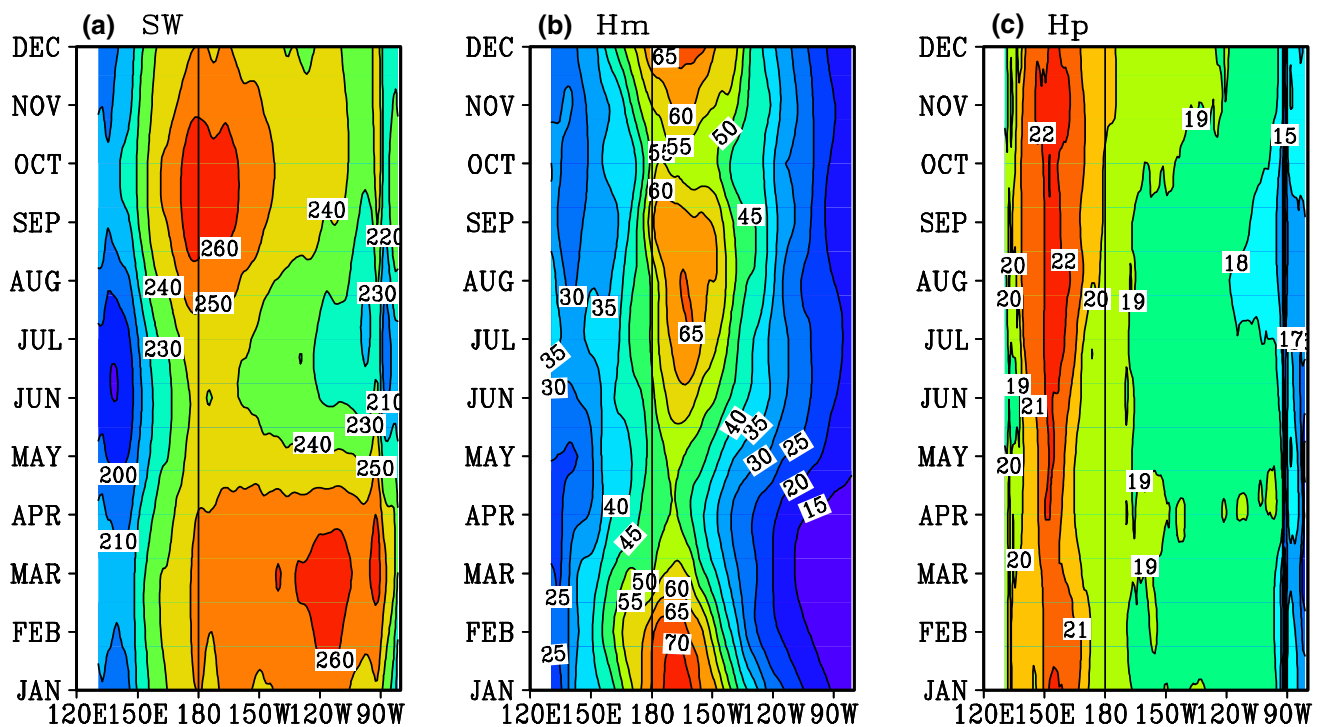
These structural relationships in nature provide an observational basis for validating model simulations (Gnanadesikan and Anderson 2009; Jochum et al. 2010; Kang et al. 2017a; Zhi et al. 2019). For example, in our previous modeling studies, a statistical approach was utilized to derive a model for the inter-annual variations in  $H_p$  from satellite observations of Chl and  $H_p$ . The observational analyses from this study indicated a good relationship between the inter-annual variations in Chl ( $H_p$ ) and the SST, so adopting a statistical approach was justified here. Furthermore, the derived  $H_p$  statistical model was used for coupled ocean–atmosphere



**Fig. 16** Horizontal distributions of the annual-mean fields during 2005–2015: **a** shortwave radiation from MERRA, **b**  $H_m$  from the Argo product, and **c**  $H_p$  from the GlobColour Project. The contour interval is  $10 \text{ W m}^{-2}$  in **a**,  $5 \text{ m}$  in **b**, and  $1 \text{ m}$  in **c**

simulations, and the relationships among these fields and modulating effects on the ENSO in modeling studies were consistent with this observation-based analysis. The results from these observation-based analyses support our previous modeling studies using the statistical  $H_p$  approach. By combining these observations and previous modeling, the modulating effects of ocean biology-induced heating were clearly demonstrated with a negative feedback on the ENSO. Furthermore, an indirect dynamical effect from inter-annual anomalies of  $H_p$  was identified as the predominant mechanism for ENSO modulations. These analyses offer a clear method to trace the influence pathways through which the ENSO is modulated by Chl anomalies, and the methodology that was developed in this study can be used in other models to reveal bio-effects.

In this study, we demonstrated the relationships between inter-annual variability of Chl and ENSO modulations using observational data. A diagnostic analysis is performed to quantify the contribution of interannual  $H_p$  anomalies to the three heating terms in association with that of  $H_m$ . Note that large inter-annual variability of Chl is located in the western-central equatorial Pacific. As seen in this analysis, the ENSO is a clear source for the Chl variability that is strongly affected by the upwelling and mixing processes



**Fig. 17** Seasonal variations along the Equator for climatological fields during 2005–2015: **a** shortwave radiation, **b**  $H_m$ , and **c**  $H_p$ . The contour interval is  $10 \text{ W m}^{-2}$  in **a**,  $5 \text{ m}$  in **b**, and  $1 \text{ m}$  in **c**

in the region. However, the detailed processes that can be responsible for inter-annual variability of Chl have not been analyzed from this observations-based analysis. Also, we present a clear explanation for understanding the feedback onto ENSO induced by inter-annual Chl variability. But, here we only considered the direct impact of  $H_p$  on inter-annual variations of these related heat items using a diagnostic method. In fact, through the modulating effects on stratification and mixing,  $H_p$  is expected to have an impact on  $H_m$ , and thus there exist interactions between inter-annual variations in  $H_p$  and  $H_m$ . These issues cannot be evaluated from this observational analysis, and thus need to be addressed further by modeling studies as shown in (Zhang et al. 2018a, b, 2019).

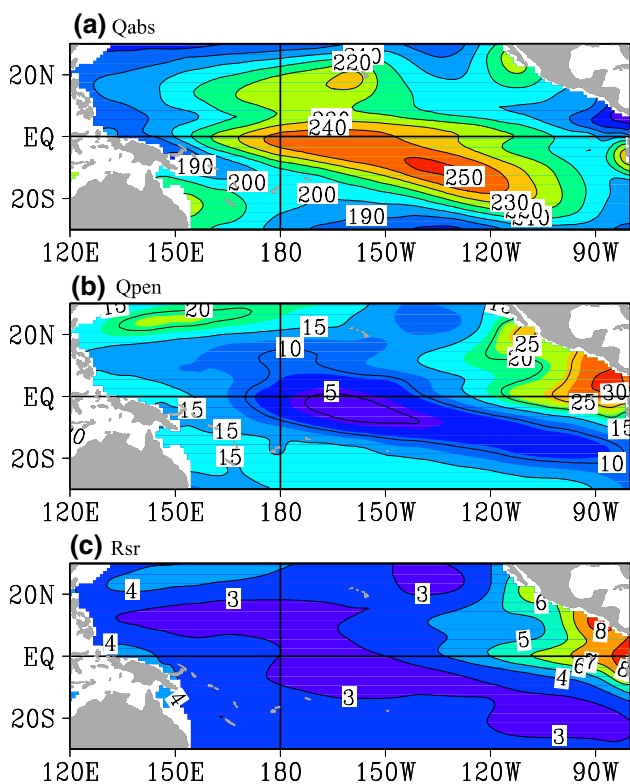
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## Appendix: Mean state and seasonal variations

Observations are used to analyze the structure and variability of ocean biology-related heating effects in the tropical Pacific, including Chl,  $H_p$ ,  $H_m$ , short wave (SW) solar radiation, and the three heating terms. The three heating terms, which were derived from combined satellite and in situ data, are new fields that have not been shown before. In this Appendix, we present the mean state and seasonal variabilities of these related fields to support the inter-annual variability analyses in the main text.

### A1. Chl field

Figure 15 displays the annual-mean field of Chl and its variations in the equatorial Pacific according to satellite measurements. The magnitude of the seasonal and inter-annual Chl variabilities is comparable to that of its total value. The main characteristics of Chl are clearly evident in Fig. 1b and Fig. 15a, as previously described (e.g., Ballabrera-Poy et al. 2007). For example, a pronounced ridge is seen in the equatorial regions, with high Chl values extending from the east to the west along the Equator. The Chl concentration is low ( $< 0.1 \text{ mg m}^{-3}$ ) in the western equatorial Pacific in association with warm waters and high in the central and eastern equatorial regions, where a cold tongue develops.



**Fig. 18** Horizontal distributions of the annual-mean fields during 2005–2015 for ocean Chl-related heating terms: **a**  $Q_{abs}$ , **b**  $Q_{pen}$ , and **c**  $R_{sr}$ . The contour interval is  $10 \text{ W m}^{-2}$  in **a**,  $5 \text{ W m}^{-2}$  in **b** and  $1 \text{ }^\circ\text{C month}^{-1}$  in **c**

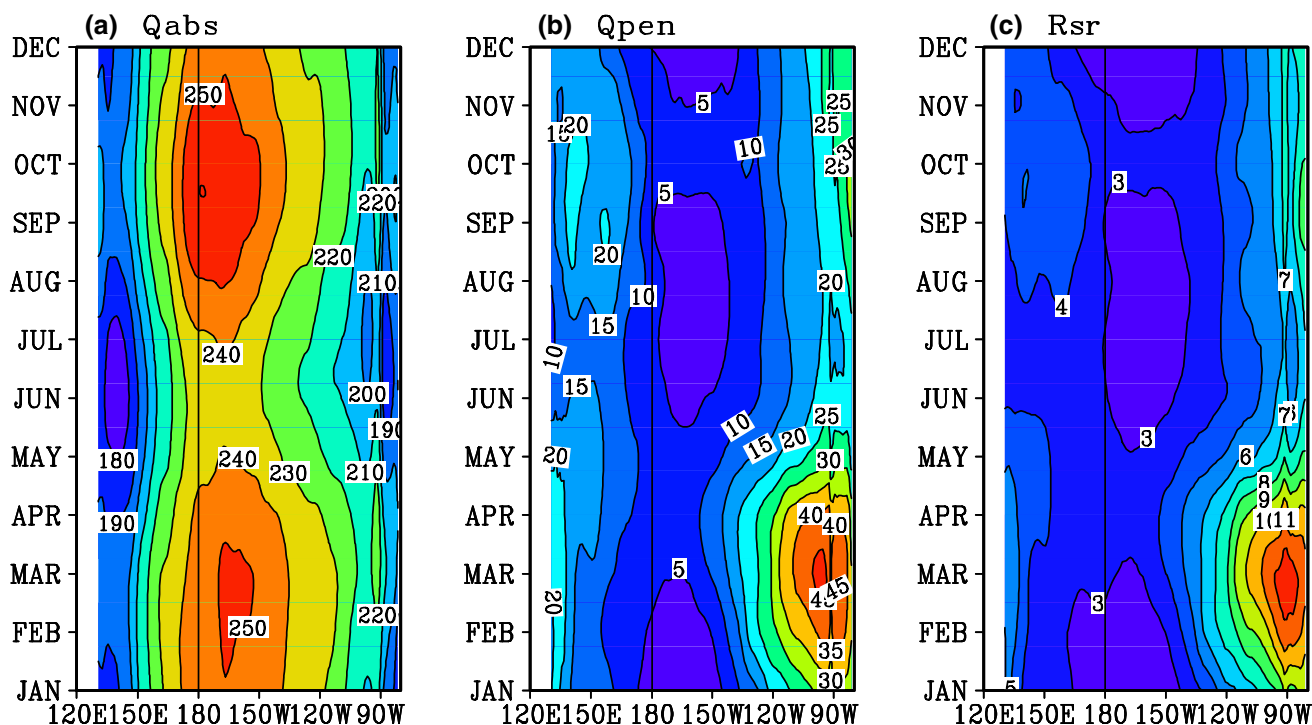
In the far-eastern coastal regions, the Chl concentration is especially high. Correspondingly, large gradients exist in the western equatorial region and far-eastern region.

The annual-mean Chl field and corresponding seasonal variability of Chl along the Equator are shown in Fig. 15. Longitudinally, Chl exhibits elevated values from west to east across the equatorial Pacific. Large seasonal variations in Chl are seen in the western-central equatorial Pacific, with a pronounced peak in summer. In the eastern equatorial region, low values are observed in spring and high values in summer. These seasonal variations in Chl are associated with those in physical conditions in the equatorial Pacific (e.g., upwelling and vertical mixing).

**A2. Short wave (SW) solar radiation,  $H_p$  and  $H_m$**

Figure 16a exhibits the horizontal distributions of the average annual-mean shortwave (SW) solar radiation during 2005–2015; its corresponding seasonal variation is shown in Fig. 17a. Regions with large values are seen in the western-central equatorial Pacific and those with low values are seen in the eastern region. Seasonally, solar radiation that reaches the sea surface has a pronounced semiannual cycle along the Equator (Fig. 17a). The SW radiation penetrates the upper ocean.

The corresponding results for the depth of the ML ( $H_m$ ) are displayed in Figs. 16b and 17b. Regions with a deep ML are located in the western-central equatorial Pacific and



**Fig. 19** Seasonal variations along the Equator for the climatological heating terms during 2005–2015: **a**  $Q_{abs}$ , **b**  $Q_{pen}$ , and **c**  $R_{sr}$ . The contour interval is  $10 \text{ W m}^{-2}$  in **a**,  $5 \text{ W m}^{-2}$  in **b** and  $1 \text{ }^\circ\text{C month}^{-1}$  in **c**

those with a shallow ML are seen in the western and eastern regions. In both the western and eastern sides of the tropical Pacific and the Intertropical Convergence Zone (ITCZ), the ML is relatively shallow because of the weak surface friction velocity and the stabilizing surface buoyancy flux. Values larger than 50 m are found in the central tropical regions, where surface winds are energetic with strong surface buoyancy losses. Seasonal variations are clearly evident. In the eastern equatorial region, the ML is shallow in spring and deep in fall. In the western equatorial basin, the ML is deep in winter but shallow in the late spring.

Chl is used to derive the penetration depth ( $H_p$ ), and its annual-mean structure is shown in Fig. 15a. A pronounced trough is seen in the equatorial regions, with low  $H_p$  values extending from east to west along the Equator. Deep penetration depths with values larger than 20 m are observed in the western equatorial Pacific, while shallow penetration depths are found in the east. Large gradients are located in the western equatorial Pacific and far-eastern regions. Seasonally, the solar radiation exhibits deep penetration in spring and shallow penetration in fall.

### A3. Three heating terms ( $Q_{\text{abs}}$ , $Q_{\text{pen}}$ and $R_{\text{sr}}$ )

As mathematically expressed above, SW,  $H_p$  and  $H_m$  all determine the distributions of penetrative solar radiation between the mixed layer and the underlying subsurface layers. The annual-mean structures of these three heating terms are shown in Fig. 18, and the corresponding seasonal variations are displayed in Fig. 19. The magnitude of the total  $Q_{\text{abs}}$  and  $Q_{\text{pen}}$  fields is one order larger than that of their interannual variabilities. Most of the solar radiation is absorbed within the ML (Fig. 18a), with some penetrating through the bottom of the ML (Fig. 18b). These terms exhibit coherent relationships with SW,  $H_m$  and  $H_p$ . For example, the structure and magnitude of  $Q_{\text{abs}}$  is similar to those of SW;  $Q_{\text{pen}}$  has low values in the western-central equatorial Pacific and high values in the western and eastern equatorial Pacific.  $Q_{\text{abs}}$ ,  $Q_{\text{pen}}$  and  $R_{\text{sr}}$  all have a clear signature for  $H_m$ , so  $H_m$  is major factor that affects the penetration of solar radiation. Seasonally,  $Q_{\text{abs}}$  has a pronounced semiannual cycle along the equator (Fig. 19a), similar to the SW solar radiation reaching the sea surface. Interestingly, seasonal variations in  $Q_{\text{pen}}$  and  $R_{\text{sr}}$  (Fig. 19b, c) exhibit a slight shift in time compared to  $Q_{\text{abs}}$  (Fig. 19c). The seasonal variations in these heating terms indicate that these terms are predominantly determined by  $H_m$ . No clear signature is seen for the effect of  $H_p$  on these heating terms in terms of the mean field and seasonal variations.

Thus,  $H_m$  is a major factor that controls the distribution of solar radiation within the ML and subsurface layers. As such, if the ML is deeper, SW solar radiation is absorbed more within the ML and penetrates less into the subsurface

layers. If the ML is deep enough, all the radiation would be absorbed within the ML, with little penetration through the bottom of the ML. In such a situation,  $H_p$  would barely influence the penetrative solar radiation in the upper ocean.

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