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## Design, realisation and analysis of an apodised, film-loaded acousto-optic tunable filter

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**ABSTRACT** The design and realisation of a SiO<sub>2</sub> film-loaded Ti:LiNbO<sub>3</sub> integrated acousto-optic tunable filter are presented. An optimised Hamming apodisation is achieved through introducing an angular offset between the acoustic and optical waveguides along with carefully placed acoustic absorbers. The realised device has a 20-mm acoustic interaction length and achieves a 3-dB linewidth of 1.8 nm with 19-dB sidelobe suppression. The realised device is highly efficient, requiring a RF drive power of only 10 mW to achieve complete polarisation conversion. The enhanced drive efficiency of the device is analysed using both acoustic and optical mode analysis. Based on this analysis, possible explanations for the improved behaviour are presented and directions that may lead to further enhancements are discussed.

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### 1 Introduction

Integrated acousto-optical tunable filters (IAOTFs) are attractive devices for wavelength-division multiplexing (WDM) optical communication systems [1] and also offer potential for use in bio-medical sensing applications [2]. Recently, efforts to migrate IAOTF implementations from Ti-diffused LiNbO<sub>3</sub> [3] to film-loaded acoustic waveguides, which allow more flexible control of acoustic mode confinement, have been reported [4]. The current state-of-the-art IAOTF devices are implemented for a 0.8-nm channel-spaced WDM system using multi-folded stages of tilted SiO<sub>2</sub> films with 60 wt % InSn acoustic waveguides on LiNbO<sub>3</sub> [5].

A film-loaded IAOTF is depicted in Fig. 1. The application of a RF electric signal to the interdigital transducer electrodes will cause a surface acoustic wave (SAW) to be generated due to the piezoelectric properties of the LiNbO<sub>3</sub> substrate. A strip of reduced acoustic velocity material forms a waveguide for

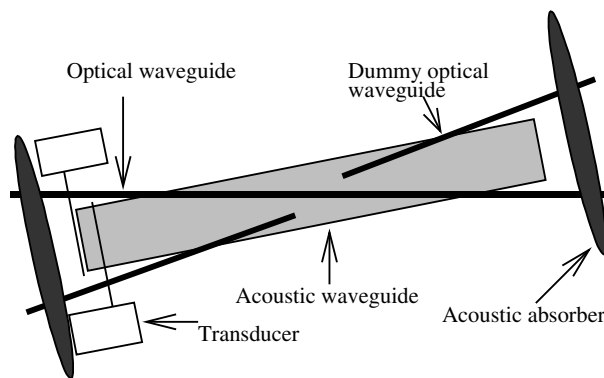


FIGURE 1 Film-loaded integrated acousto-optic tunable filter (IAOTF) structure

the generated SAW along the length of the device. The stress produced by the presence of the SAW on the LiNbO<sub>3</sub> surface causes a perturbation in the optical properties of the material. This perturbation can be used to couple the orthogonal TE and TM polarised optical modes propagating in the titanium-diffused optical waveguide beneath the SAW guide.

Since the stress caused by the SAW is periodic, the optical perturbation is also periodic and thus acts much like an optical grating. Hence, depending on the difference in propagation velocities of the TE and TM polarised optical modes, if the optical wavelength and the acoustic period meet a phase-matching condition, the TE and TM modes will be strongly coupled. For a particular SAW period only a single optical wavelength will meet this phase-matching condition and thus the particular wavelength that has its polarisation converted may be precisely controlled via the RF applied to the transducer. If the input optical signal is of a single polarisation, a polarisation splitter at the output allows the converted wavelength to be dissociated from the other wavelengths.

The resolution and rejection of any filter will largely be determined by the filter apodisation. Many AOTF apodisations have been proposed; however it can be shown that the optimised Hamming apodisation function offers the best compromise between linewidth and sidelobe level [3].

Methods to implement these apodisations in AOTF devices have also been proposed [3, 6]. Some of these techniques can introduce phase variations in the acoustic profile along the path of the optical waveguide, which may result

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in asymmetry in the filter response and poor sidelobe suppression. Apodising the acousto-optic interaction through use of an angular offset between the acoustic and optical waveguides ensures a constant acoustic phase along the optical path and thus is more likely to provide an optimised filter response. When the acoustic waveguide is formed using a Ti-diffused cladding and the optical waveguide is formed using a Ti-diffused core, it can be impractical to arrange angular offset apodisation as it is difficult to define the optical waveguide outside of the undiffused acoustic guiding region. Strip-loaded acoustic waveguides have a clear advantage in this respect.

This investigation thus focuses on the realisation of an optimised Hamming apodised AOTF through use of an angular offset between the acoustic and optical waveguides. First, in Sect. 2, the acoustic properties of several thin films are analysed to determine which film will be most suitable for strip loading. Based on this analysis, test structures of SiO<sub>2</sub> strip-loaded SAW guides are fabricated on a LiNbO<sub>3</sub> substrate to experimentally verify the SAW velocity and attenuation. Section 3 then utilises the experimentally determined SAW guide properties in conjunction with a beam-propagation method (BPM) and coupled-mode-equation models to predict the integrated AOTF response. This modeling enables the design of the SAW guide geometry and relative placement of acoustic and optical waveguides to provide an optimised filter response.

The realisation of the designed filter is described in Sect. 4 along with the measured performance of this device. It is found to offer a reasonably narrow linewidth and good sidelobe suppression, comparing well with predictions. Of particular interest is the enhanced efficiency of the device. Section 5 presents a theoretical analysis of the interaction between the guided acoustic and optical modes in order to better understand the factors contributing to efficient conversion. Drawing on this analysis, several possible avenues for the improvement of the realised device are suggested.

## 2 Film-loaded SAW guide characterisation

The propagation characteristics of the SAW mode depend on the acoustic velocity difference between the bare LiNbO<sub>3</sub> cladding and the film-loaded LiNbO<sub>3</sub> guide. Figure 2 presents the velocity difference modelled for SiO<sub>2</sub>, MgO, ZnO and AlN films on LiNbO<sub>3</sub> using specialised software [6]. This modelling demonstrates that through careful selection of strip material and thickness, a wide range of core acoustic velocities are available. Although films offering high acoustic velocity contrast (such as ZnO) may offer more tightly confined SAW propagation, we have chosen to use SiO<sub>2</sub> for this work. This material offers a moderate reduction in acoustic velocity in the guiding region, very similar to that which can be achieved through Ti diffusion [3, 6], and is routinely produced in our laboratory.

To verify that the acoustic velocity was as predicted in Fig. 2 and to ensure that the acoustic propagation loss was sufficiently low, a SAW test structure was fabricated [7]. The geometry of this test structure is presented in Fig. 3.

The transmitting transducer (T<sub>x</sub>) launches acoustic waves along the strip guide in both directions. These acoustic waves

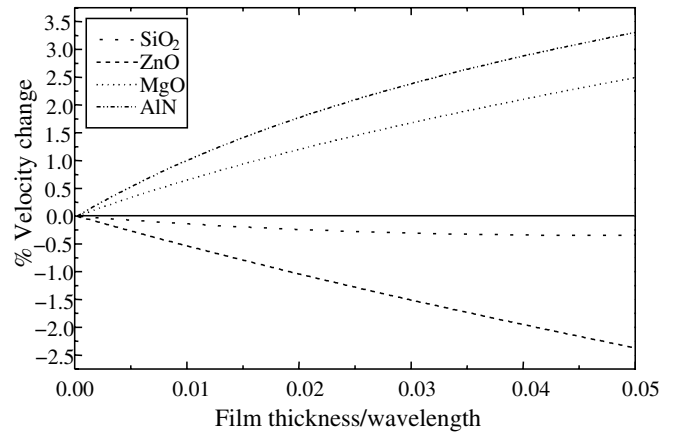


FIGURE 2 SAW velocity change for different films

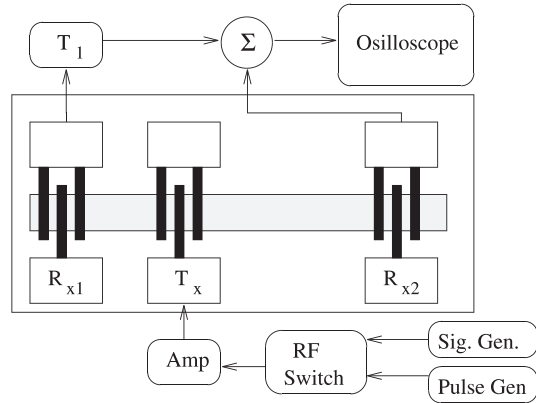


FIGURE 3 SAW velocity-measurement setup

travel two different distances ( $d_1$  and  $d_2$ ) before being detected at two receiver transducers ( $R_{x1}$  and  $R_{x2}$  respectively). The input signal is a pulse-modulated tone and the two outputs are summed and displayed on an oscilloscope.

A time delay ( $T_1$ ) was introduced between the two received signals (as shown in Fig. 3) to approximately compensate for the path difference and then the frequency of the tone was carefully adjusted until the two received signals interfered destructively. The phase velocity  $V_p$  was then found by substituting the frequency of cancellation  $f_{\text{null}}$  into (1):

$$V_p = \frac{2f_{\text{null}}(d_2 - d_1)}{(2i + 1)}, \quad (1)$$

where  $i$  is an integer.

To fabricate the test structure, the SiO<sub>2</sub> film was deposited using a RF magnetron sputter system using a 99.99% purity SiO<sub>2</sub> target. The chamber was evacuated to a base pressure of  $2 \times 10^{-5}$  Torr and the sample temperature was maintained at 140 °C. A gas mixture of 1 : 9 O<sub>2</sub> : Ar at  $1 \times 10^{-2}$  Torr was introduced and excited with a RF power of 160 W for 30 min to produce a 0.38- $\mu\text{m}$ -thick film. A strip of 80- $\mu\text{m}$  width was then patterned on the surface by masking and etching with HF. The transducer consisted of 20 finger pairs with a 20.8- $\mu\text{m}$  period and was fabricated using an evaporated 200 Å : 300 Å : 500 Å Ti : Ni : Au metallisation to achieve low conductor loss and good adhesion. Au wire bonding was used to connect the transducer to a 150-nH inductor and a 3–20-pF tunable capacitor matching network.

Using the technique described above, the phase velocity of the 80- $\mu\text{m}$ -wide, 0.38- $\mu\text{m}$ -thick  $\text{SiO}_2$  strip-loaded SAW guide was found to be  $3685 \text{ ms}^{-1}$  and the  $\text{LiNbO}_3$  velocity to be  $3700 \text{ ms}^{-1}$  at 174.5 MHz. The error in this measurement is estimated to be around  $\pm 1 \text{ ms}^{-1}$ . The SAW attenuation at this frequency was estimated to be  $1.75 \text{ dB cm}^{-1}$  by comparing the amplitudes of the signals received at  $R_{x1}$  and  $R_{x2}$  and the corresponding path-length difference.

### 3 IAOTF response modelling

Having established the propagation characteristics of the  $\text{SiO}_2$  strip-loaded SAW guide in Sect. 2, it was possible to use the experimentally verified data to determine the arrangement of acoustic and optical waveguides that offers the optimum wavelength-filter profile.

The evolution of the acoustic signal was modelled using a beam-propagation method model [8, 9]. The resulting polarisation conversion was modelled by recording the acoustic intensity along the optical waveguide (see Fig. 1) and then using coupled-mode theory [10] to determine the polarisation conversion.

The angular offset between the optical waveguide and the SAW guide characterised in Sect. 2 was optimised to achieve an acoustic intensity profile that approximated an optimised

Hamming weighting function. During this analysis it was found that a second dummy optical waveguide (as shown in Fig. 1) was needed to maintain symmetry of the acoustic amplitude and sustain good symmetric sidelobe suppression.

The acoustic amplitude profile and corresponding filter characteristics predicted for the 80- $\mu\text{m}$ -wide, 20-mm-long SAW guide at several angular offsets are presented in Fig. 4 and Fig. 5 respectively. An input polarisation extinction ratio of 25 dB is assumed.

From Fig. 5 it is evident that an offset of  $0.55^\circ$  offers the narrowest linewidth, whilst maintaining a low sidelobe level.

### 4 Film-loaded IAOTF realisation and results

Based on the predictions of Sect. 3, IAOTF devices were fabricated with a 0.38- $\mu\text{m}$ -thick  $\text{SiO}_2$  film SAW guide of width 80  $\mu\text{m}$  at an angular offset of  $0.55^\circ$  to the optical waveguide. Acoustic absorbers were placed a short distance from either end of the SAW guide (as depicted in Fig. 1) to truncate the acoustic amplitude profile at its extremities.

It is important for good coupling efficiency and extinction that the TE and TM modes of the optical waveguide are of as similar cross-sectional distributions as possible and that they propagate with the same attenuation. Several experiments were conducted by varying the diffusion temperature and the Ti-strip thickness to produce optical waveguides with TE and TM modes of similar characteristics. The final parameters used were 7- $\mu\text{m}$ -wide, 1000  $\text{\AA}$ -thick Ti strip diffused for 8 h at  $1050^\circ\text{C}$ .

An optical carrier, supplied by a tunable laser, was polarised using a fibre polarisation controller and then butt-coupled to the IAOTF device. An amplified RF tone at 171 MHz was applied to the transducer, corresponding to an optical conversion wavelength of 1560 nm. The optical output was imaged through a crossed polariser onto an optical power meter. The wavelength of the tunable laser was varied from 1554 to 1565 nm and the optical power transmitted through crossed polarisers was recorded and is presented in Fig. 6. The unconverted power is also shown in Fig. 6, demonstrating that low cross-talk is achieved.

This 20-mm-long device achieves a 3-dB linewidth of 1.8 nm and a sidelobe level of  $-19 \text{ dB}$ . This response compares well to that predicted in Fig. 5 although some slight asymmetry is evident. This asymmetry is most likely caused by small inhomogeneities in the SAW guide and the optical waveguide along the length of the device. The sidelobe level of  $-19 \text{ dB}$  is respectable but falls short of the ideal  $< -22 \text{ dB}$  level predicted for the optimised Hamming apodisation. This shortcoming may be attributed to the non-ideal polarisation extinction ratio at the device input of only 25 dB. It is believed that improvement of this input polarisation control through the use of an integrated optical polarisation splitter would yield further sidelobe-level reduction.

The RF drive power required for maximum conversion of the central optical wavelength was only 10 mW, which is significantly lower than the  $\geq 50 \text{ mW}$  required by previously reported Ti-diffused geometries [11]. This enhanced conversion efficiency was an important and surprising result. Improving the conversion efficiency is of benefit to system integration for communications and compact spectroscopic applications, in

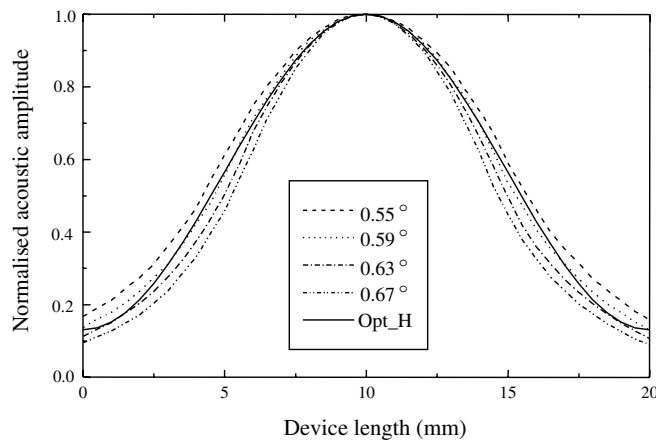


FIGURE 4 Acoustic amplitude for different interaction angles

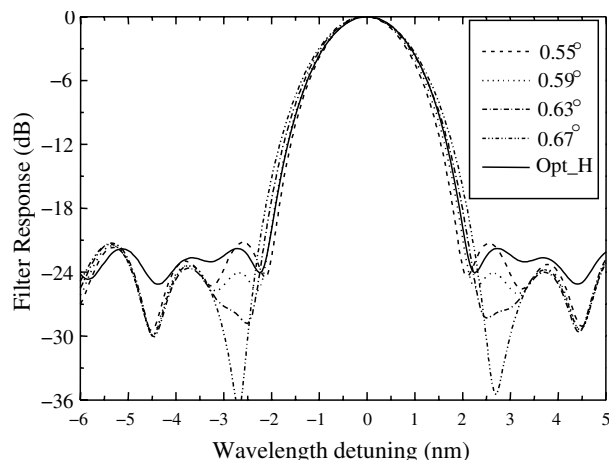
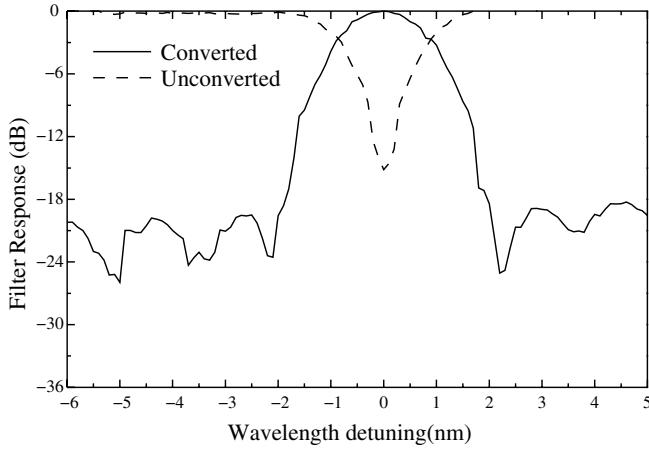


FIGURE 5 Filter response for different interaction angles



**FIGURE 6** Filtered and unfiltered responses at 171 MHz and centre wavelength of 1560 nm

terms of both a reduced impact on the overall power budget for the system and probably more importantly a significant reduction in the heat generated by the device. In order to better understand the observed efficiency, further theoretical analysis of the interaction of the acoustic and optical modes was conducted.

## 5 Discussion of acousto-optic conversion efficiency

It is proposed that the conversion efficiency is determined by three key aspects of the AOTF device: (i) efficient electro-mechanical coupling at the launching transducer; (ii) the confinement of the acoustic guide and its alignment with the optical modes and (iii) the overlap of the TE and TM polarised optical modes. A discussion of each of these effects, whether they contribute to the the observed efficiency of the current device and how they might be utilised to improve the efficiency of future devices, is presented in the following subsections.

### 5.1 Electro-mechanical efficiency at the transducer

The electro-mechanical coupling coefficient,  $K$ , of the transducer is approximately given as [12]

$$K^2 \propto \frac{2(V_f - V_m)}{V_f} \quad (2)$$

Here  $V_f$  and  $V_m$  are the free-material and metallised acoustic wave velocity respectively. According to the results of Sect. 2, the acoustic velocity of a 0.38- $\mu\text{m}$ -thick  $\text{SiO}_2$  film and bare  $\text{LiNbO}_3$  differ by less than 0.3%. The difference between the acoustic velocity of metallised  $\text{SiO}_2$  and  $\text{LiNbO}_3$  should also be small and so the electro-mechanical efficiency should be similar between the strip-loaded and diffused acoustic waveguide transducers. Further, the transducer used in the strip-loaded device reported here was very similar to that reported in [11], having a 20-finger-pair variable capacitor and using a series inductor for impedance matching. Thus it is unlikely that an improvement in efficiency of the order observed could result from enhanced transducer performance.

In this device the transducer was not optimised to compensate for the acoustic walk-off [13]. A slight rotation of

the transducer has been shown to improve the efficiency of launching acoustic waves in diffused SAW guides [14], and it is expected that a similar improvement could result when applied to the strip-loaded guide under investigation. Other possibilities that could be pursued to improve efficiency include implementation of unidirectional transducers and examination of the efficiency with which the transducers launch the fundamental mode of the SAW guide.

### 5.2 Acousto-optic efficiency

Another significant factor contributing to the filter conversion efficiency is the efficiency of the interaction between the optical modes and the acoustic perturbation.

The coupling coefficient in  $X$ -cut  $Y$ -propagating  $\text{LiNbO}_3$  devices can be quite accurately modelled using the following approximation [15]:

$$k \approx k_{xz} = \frac{\omega}{8P_{\text{SAW}}} \int \int (E_x^{\text{TM}})^* \Delta \epsilon_{xz} E_z^{\text{TE}} dx dy \quad (3)$$

Here  $E_x^{\text{TM}}$  and  $E_z^{\text{TE}}$  are the electric field components of the TM and TE optical modes directed along the crystal  $X$  and  $Z$  axes respectively and  $\Delta \epsilon_{xz}$  is the perturbation in the permeability that couples the TE and TM optical modes. This is approximately proportional to the strain and electric field components of the SAW in the  $X$  direction.

To ensure the most efficient polarisation conversion, the  $X$ -directed SAW distribution profile must be a good match to the optical modes. It is thus proposed that a tightly confined acoustic mode would enhance efficiency by causing more acoustic power to be coincident with the propagating optical modes.

The current strip-loaded waveguide does not have a tightly confined acoustic mode, having a similar ratio of guide to cladding acoustic velocity as the diffused device in [16]. The strip-loaded waveguide width of 80  $\mu\text{m}$  is slightly narrower than the 100- $\mu\text{m}$  width of [11]; however this is unlikely to cause the observed improvement in efficiency.

A further enhancement of the efficiency could be sought by significantly increasing the confinement of the acoustic mode so that its width was closer to that of the optical modes. This could be done through the use of an alternative material such as  $\text{ZnO}$  for the strip waveguide as suggested in Sect. 2. Another approach is to further load the  $\text{SiO}_2$  strip waveguide with a thin layer of dense material. Seino et al. [4] have used an  $\text{InSn}$ -loaded  $\text{SiO}_2$  film and significantly narrower strip guide widths to attain a conversion power of only 5 mW. To maintain the same apodisation with a more tightly confined SAW it would be necessary to reduce the angular offset significantly.

### 5.3 TE and TM optical mode overlap

It is evident from (3) that it is not only important for the acoustic mode to match the optical modes, it is also very important that the TE and TM optical modes match each other. Careful experimentation was conducted to ensure that the TE and TM modes were as similar as possible, both in terms of field profile and also in terms of propagation loss to maximise the possible extinction. The measured power-normalised field profiles for the TE and TM optical modes are presented in

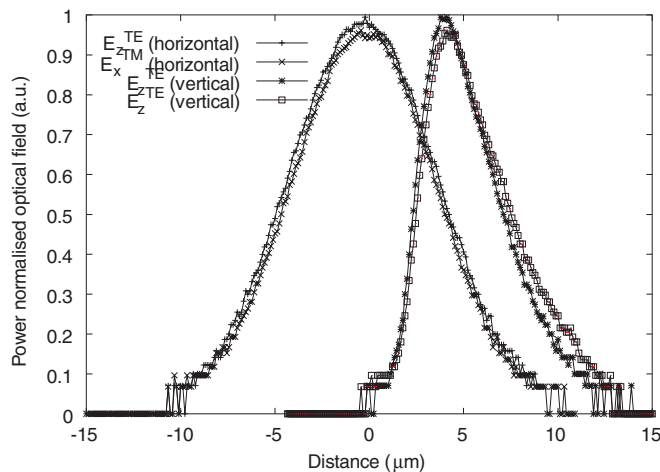


FIGURE 7 Measured power-normalised fields for the TE and TM optical modes in both horizontal and vertical directions

Fig. 7. The field distributions of the two optical modes are evidently very closely matched.

It is thus suggested that most of the improvement in observed conversion efficiency is due to the good overlap between the TE and TM optical modes. It will therefore be worth attempting to implement some of the techniques described above to further improve the conversion efficiency. An investigation of applying these techniques to the present device is currently under way.

## 6 Conclusion

In this investigation we have demonstrated an integrated acousto-optic tunable filter utilising a  $\text{SiO}_2$  film-loaded SAW guide at an angular offset to the optical path to implement an optimised Hamming apodisation. The characteristics of the film-loaded SAW guide in isolation have been accurately modelled and experimentally verified. The strip-loaded SAW guide has been used in the realisation of an apodised  $\text{SiO}_2$  film-loaded integrated acousto-optic tunable filter device and achieves a 19-dB sidelobe suppression and a 3-dB linewidth of 1.8 nm for a 20-mm acoustic interaction length. It is believed that the sidelobe-level suppression and the filter linewidth could be improved with the incorporation of more precise polarisation con-

trol, and extension of the acousto-optic interaction-region length.

The demonstrated device is highly efficient, requiring only 10 mW of RF power for complete optical polarisation conversion. A detailed discussion of the factors determining conversion efficiency in this device has been presented and suggests that this observed efficiency is due to good matching between the TE and TM optical modes. Several alternative avenues for further improvement of the device efficiency have been identified and efforts to implement these proposed techniques are currently under way.

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