Excited-state absorption and stimulated emission of tetrahedral Co2⁺ ion in LiGa₅O₈

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Received: 5 June 1998/Revised version: 20 October 1998/Published online: 19 May 1999

Abstract. Differential absorption spectra measurements of tetrahedral Co^{2+} -doped LiGa₅O₈ have been made using a ps pump-probe technique. The stimulated emission from the ${}^{4}T_{1}({}^{4}P)$ level to the ${}^{4}A_{2}({}^{4}F)$ and ${}^{4}T_{2}({}^{4}F)$ states is observed to be overlapped by the excited-state absorption. The excitedstate absorption is tentatively assigned to transition from the ${}^{4}T_{1}({}^{4}P)$ level of the Co²⁺ ion to the conduction band of LiGa₅O₈. For Co²⁺:LiGa₅O₈, laser operation is expected to be possible in the vicinity of 720 nm and in the 860–970 nm spectral region.

PACS: 42.65; 42.70

 $Co²⁺$ -doped crystals in which the $Co²⁺$ ions are in a tetrahedrally coordinated lattice site have been under investigation [1–7]. This is due to the fact that crystals doped with tetrahedral Co^{2+} ions are attractive for laser applications. The tetrahedral Co^{2+} in $Y_3Sc_2Ga_3O_{12}$ (YSGG) and $Y_3Al_5O_{12}$ (YAG) has been used as saturable absorber for passive Q-switching of the Er:glass laser [8, 9]. Co^{2+} :ZnSe has been used as saturable absorber Q-switch for the Er^{3+} : Yb^{3+} :glass laser [10]. On the other hand, the tetrahedral $Co²⁺$ ion in a number of spinels such as $MgAl₂O₄$ [1], LiGa₅O₈ [2], and $ZnGa₂O₄$ [3] exhibits strong broad luminescence bands in the visible and near-infrared spectral regions with sub-µs lifetimes. This allows us to consider such crystals as possible candidates for tunable solid-state lasers in the visible and near-infrared. In this paper, excited-state absorption and stimulated emission of Co^{2+} :LiGa₅O₈ are studied under ps excitation in the spectral region 460–970 nm.

 $LiGa₅O₈$ has an inverse spinel structure (space group $O⁶$ -P4332), and contains four formula units in the cubic cell [3]. The lattice constant is $a = 0.833$ nm. One half of the Ga³ ions occupy tetrahedral sites while the other half of the Ga^{3+} ions, together with the $Li⁺$ ions, occupy octahedral sites. It has been shown [5] that the Co^{2+} ions substitute for tetrahedrally coordinated Ga^{3+} ions in the LiGa₅O₈ lattice and occupy sites of C_3 point-group symmetry. The thickness of the Co^{2+} :LiGa₅O₈ sample studied was 1.1 mm.

1 Experiment

Excited-state absorption and stimulated emission measurements were carried out by use of a ps pump-probe technique. The differential absorption spectra were measured as

$$
\Delta OD(\lambda) = -\lg(T/T_0),\tag{1}
$$

where T and T_0 are probe beam transmittance of pumped and unpumped crystal, respectively. The 15-ps-duration pulses at 0.54 µm from a frequency-doubled passively mode-locked Nd^{3+} :YAlO₃ laser were used as a pump beam. The ps whitelight continuum, which was produced in a $D₂O$ cell, was used as a probe beam. An optical multichannel analyzer with two photodiode arrays combined with a spectrometer was used as a recording system. All the measurements were performed at room temperature, and with a spectral resolution ≈ 1 nm. The differential absorption ∆OD of the probe beam after all the excited ions have accumulated in a level whose lifetime is much longer compared to the pulse duration, can be expressed as $[11-13]$

$$
\Delta OD(\lambda) = \Delta NL(\sigma_{\text{ESA}} - \sigma_{\text{GSA}} - \sigma_{\text{SE}}), \qquad (2)
$$

where σ_{ESA} , σ_{GSA} , and σ_{SE} are the excited-state absorption (ESA), the ground-state absorption (GSA), and the stimulated emission (SE) cross sections, respectively, ∆*N* is the concentration of ions in the excited state, and *L* is the sample length. Notice that $\Delta OD > 0$ means $\sigma_{\text{ESA}} > \sigma_{\text{GSA}} + \sigma_{\text{SE}}$, and Δ OD < 0 signifies that $\sigma_{\text{ESA}} < \sigma_{\text{GSA}} + \sigma_{\text{SE}}$. The relative shape of the $\sigma_{\text{SE}}(\lambda)$ was obtained by measuring the fluorescence spectrum $I(\lambda)$ which in bulk media is related to the $\sigma_{\text{SE}}(\lambda)$ by $\sigma_{\text{SE}}(\lambda) \propto \lambda^5 I(\lambda)$ [14].

2 Results

The GSA spectrum of Co^{2+} :LiGa₅O₈ shown in Fig. 1 is similar to that of Co^{2+} ions in a tetrahedrally coordinated site in crystals with spinel, garnet, and fluorite structure $[1-7]$. In accordance with a Tanabe–Sugano diagram for a d^7 ion

Fig. 1. The absorption spectra of the Co^{2+} :LiGa₅O₈ and undoped LiGa₅O₈ (*dashed curve*) crystals

in a tetrahedral crystal field (Fig. 2), the visible absorption band centered at 16600 cm^{-1} (≈ 600 nm) is assigned to the ${}^4A_2({}^4F) \rightarrow {}^4T_1({}^4P)$ transition, and the near-infrared band around 6700 cm⁻¹ ($\approx 1.5 \,\mu$ m) to ⁴A₂(⁴F) \rightarrow ⁴T₁(⁴F) transition [2]. The absorption band due to the ⁴A₂(⁴F) → ⁴T₂(⁴F) transition expected to lie in the vicinity of 4000 cm^{−1} $(\approx 2.5 \,\mu\text{m})$ is not observed because of its low oscillator strength. In a pure tetrahedral field this transition is electricdipole forbidden. In a lower-symmetry crystal field (the $Co²⁺$ ions in $LiGa₅O₈$ occupy distorted tetrahedrally coordinated sites with C_3 point-group symmetry [2]) both the ⁴T₁(⁴P)

and ${}^{4}T_{1}({}^{4}F)$ energy levels split into two components, a singlet and a doublet, which are exhibited by two GSA bands with peaks at $\approx 16900 \text{ cm}^{-1}$ and $\approx 15700 \text{ cm}^{-1}$ originated from the ⁴A₂(⁴F) \rightarrow ⁴T₁(⁴P) transition, and two GSA bands
at \approx 6150 cm⁻¹ and \approx 7000 cm⁻¹ related to the ⁴A₁(⁴F) \rightarrow ${}^{4}T_{1}({}^{4}F)$ transition [2]. The absorption feature at 18 000 cm⁻¹, which is observed as a shoulder on the high-energy site of the ${}^4A_2({}^4F) \rightarrow {}^4T_1({}^4P)$ transition, is most likely due to one of the doublet levels arising from the ${}^{2}G$ free-ion level [2, 6]. The weak peak at $\approx 20600 \text{ cm}^{-1}$ may be assigned to the ${}^4A_2({}^4F) \rightarrow {}^2T_1({}^2P)$ transition by comparison with the spectral position of this transition observed for the tetrahedrally coordinated Co^{2+} ion in MgAl₂O₄ and ZnAl₂O₄ [1–4]. The GSA cross sections were determined from an expression $\sigma_{\text{GSA}}(\lambda) = \alpha(\lambda)/N_0$, where $\alpha(\lambda)$ is the measured ground-state absorption coefficient, N_0 is the Co²⁺ concentration in the $LiGa₅O₈$ crystal. From the known concentration of cobalt in the raw material $(0.1 \text{ wt. } %)$, the concentration N_0 was estimated to be 5.7×10^{19} cm⁻³. The obtained values $\sigma_{\text{GSA}}(\lambda)$ are shown in Fig. 1. The band-to-band absorption of $LiGa₅O₈$ (Fig. 1) begins at $\approx 32000 \text{ cm}^{-1}$ (312 nm). The absorption edge of the Co^{2+} :LiGa₅O₈ crystal rising from \approx 27 000 cm⁻¹ (370 nm) can be assigned to the transition from die impurity level 4A_2 to the conduction band.

Figure 3 shows the differential absorption spectra of the $Co^{2+}:\text{LiGa}_5\text{O}_8$ crystal in the wavelength region of 460–970 nm. The sharp feature in this spectrum is a marker for the position of the pump. At this wavelength some light from the pump laser is scattered in the probe direction. GSA bleaching, ESA, and SE are observed simultaneously. Clearly observed is an ESA band in the 460–550 nm and 740–860 nm regions as well as a SE in the range of 860–970 nm (inset in Fig. 3). One can see a SE around 675 nm which is expected

Fig. 2. Energy-level diagram of the Co^{2+} ion in LiGa₅O₈ in T_d symmetry. Ground-state absorption and stimulated-emission transitions are indicated by arrows. Note that doublet levels are omitted. CB, conduction band

Fig. 3. The differential absorption spectrum $\Delta OD = -\lg(T/T_0)$ and corresponding ground-state absorption (*solid curve*) and stimulated-emission (*dashed curves*) spectra for Co^{2+} :LiGa₅O₈. The pump wavelength is 540 nm. The delay time between pump and probe pulses is 30 ps. The *inset* shows an expanded view of the ∆OD signal indicating lower values of that. The peak values of $\sigma_{\text{SE}}(\lambda)$ for the SE bands around 695 and 950 nm should not be compared

Fig. 4a–c. The differential absorption $\Delta OD = -\lg(T/T_0)$ of $Co^{2+}:\text{LiGa}_5O_8$ as a function of pump-probe delay time at different wavelength of 675 nm (**a**), 603 nm (**b**), and 480 nm (**c**)

to overlap with the GSA bleaching and ESA. Figure 4 shows the temporal dependence of differential absorption ∆OD for three wavelengths from the spectral regions in which the GSA bleaching, ESA, and SE are observed. The GSA bleaching, ESA, and SE appear with an instrument-limited rise time of \approx 20 ps, and there is its negligible decay after our maximum delay time of 230 ps.

3 Discussion

The pump wavelength at 540 nm excites the ${}^{4}A_{2}({}^{4}F) \rightarrow {}^{4}T_{1}({}^{4}P)$ transition (Fig. 2). For the LiGa₅O₈ crystal, the lifetime of the ${}^{4}T_{1}({}^{4}P)$ level has been found to be 200 ns [2] which is much longer than the 15-ps pulse duration in our experiments. This indicates that only transitions from the ${}^{4}T_{1}({}^{4}P)$ level can appear in our measured ΔOD spectra as contribution from ESA and (or) SE. The pump pulse depletes the ground-state level ⁴A₂(⁴F) resulting in the ⁴A₂(⁴F) \rightarrow the ground-state level ⁴A₂(⁴F) resulting in the ⁴A₂(⁴F) \rightarrow ⁴T₁(⁴P) transition bleaching. Since no energetically higher lying quartet levels exist (than the ${}^{4}T_{1}({}^{4}P)$ one) there are no spin-allowed ESA transitions from the ${}^{4}T_{1}({}^{4}P)$ level.

The ESA is tentatively assigned to the transition from the impurity energy level ${}^{4}T_{1}({}^{4}P)$ to the conduction band (Fig. 2).

This is supported by the relative energy position of the ${}^4A_2 \rightarrow$ CB absorption edge which extends from 27 000 cm−¹ and the ${}^{4}A_{2}({}^{4}F) \rightarrow {}^{4}T_{1}({}^{4}P)$ transitions which are located at 15700 and 16900 cm^{-1} (Fig. 1). So the absorption related to the transition from the ${}^{4}T_{1}({}^{4}P)$ excited level of the tetrahedral $Co²⁺$ ion to the conduction band of the LiGa₅O₈ crystal may be expected to rise from $\approx 11000 \text{ cm}^{-1}$ ($\approx 900 \text{ nm}$). It can not be excluded that, in accordance with a Tanabe–Sugano diagram for a d^7 electronic configuration in a tetrahedral crystal field, the spin-allowed transitions from the thermally populated ²E(² \hat{G}) state to doublet levels arising from the ²F freeion level contribute to the ESA in the 740–860 nm region. In spectral region studied here the Co^{2+} :LiGa₅O₈ crystal exhibits two SE bands at around 695 nm and 950 nm (Fig. 3). These SE bands are attributed to the ${}^{4}T_{1}({}^{4}P) \rightarrow {}^{4}A_{2}({}^{4}F)$ and ${}^{4}T_{1}({}^{4}P) \rightarrow {}^{4}T_{2}({}^{4}F)$ transitions, respectively (Fig. 2). The ${}^{4}T_{1}({}^{4}P) \rightarrow {}^{4}T_{2}({}^{4}F)$ SE band is affected by the ESA at the short-wavelength edge of the SE band (inset in Fig. 3). The spectral region in which σ_{SE} exceeds σ_{ESA} ranges from 860 to 970 nm. Taking into account that there is no GSA in this region we can conclude that the region 860–970 nm is promising for laser operation. The ${}^4T_1({}^4P) \rightarrow {}^4A_2({}^4F)$ SE clearly overlaps with the GSA bleaching and the ESA. The wavelength region around 720 nm in which σ_{SE} exceeds σ_{ESA} , and σ_{GSA} seems to be low compared with $\sigma_{\text{SE}} - \sigma_{\text{ESA}}$ (the reabsorption by GSA is low) appears suitable for laser operation. In order to estimate the effective stimulated-emission cross sections $\sigma_{\text{EFF}} = (\sigma_{\text{SE}} - \sigma_{\text{ESA}})$, the $\Delta \text{OD}(\lambda)$ spectrum must be corrected for the GSA bleaching, see (2). The unknown constant ∆*NL* can be determined from (2) if a specific wavelength exists, for which the ESA and SE are negligible. Under the assumption, that the ESA and SE are small in the vicinity of 640 nm, the cross sections σ_{EFF} can be estimated to be 4.3×10^{-19} cm² at 700 nm and 2.2×10^{-20} cm² between 880 nm and 900 nm.

4 Conclusion

The differential absorption spectra measurements of tetrahedral Co^{2+} -doped LiGa₅O₈ crystal were carried out with a ps pump-probe technique. The stimulated emission bands due to the transitions from the ⁴T₁⁽⁴P) level to the ⁴A₂⁽⁴F) and ${}^{4}T_{2}({}^{4}F)$ levels as well as the ESA in the 460–550 nm and 740–860 nm regions were clearly observed. The ESA is tentatively assigned to the transition from the ${}^{4}T_{1}({}^{4}P)$ level of $Co²⁺$ ion to the conduction band of the LiGa₅O₈ crystal. For tetrahedral Co^{2+} -doped LiGa₅O₈ crystal, laser operation is expected in the 860–970 nm spectral region and in the vicinity of 720 nm.

References

- 1. N.V. Kuleshov, V.P. Mikhailov, V.G. Scherbitsky, P.V. Prokoshin, K.V. Yumashev: J. Lumin. **55**, 265 (1993)
- 2. J.F. Donegan, F.G. Anderson, F.J. Bergin, T.J. Glynn, G.F. Imbusch: Phys. Rev. B **45**, 563 (1992)
- 3. T. Abritta, F.H. Blak: J. Lumin. **48**/**49**, 558 (1991)
- 4. J. Ferguson, D.L. Wood, L.G. Van Uitert: J. Chem. Phys. **51**, 2904 (1969)
- 5. R. Pappalardo, D.L. Wood, R.C. Linares: J. Chem. Phys. **35**, 2041 (1961)
- 6. D.L. Wood, J.P. Remeika: J. Chem. Phys. **46**, 3595 (1967)
- 7. V.M. Orera, R. Merino, R. Cases, R. Alcala: J. Phys. Condens. Matter. **5**, 3717 (1993)
- 8. M.B. Camargo, R.D. Stultz, M. Birnbaum, M. Kokta: Opt. Lett. **20**, 339 (1995)
- 9. R.D. Stultz, M.B. Camargo, M. Birnbaum, M. Kokta: OSA Proc. on Advanced Solid-State Lasers, Vol. 24, 1995, ed. by B.H.T. Chai, S. Payne, pp. 460–464
- 10. M. Birnbaum, M.B. Camargo, S. Lee, F.Unlu. R.D. Stultz: OSA TOPS on Advanced Solid-State Lasers, Vol. 10, 1997, ed. by C.R. Pollock, W.R. Bosenberg, pp 148–151
- 11. N.T. Melamed, F. de S.Barros, P.J. Viecaro, J.O. Artman: Phys. Rev. B **5**, 3377 (1972)
- 12. R.M. MacFarlane, J.C. Vial: Phys. Rev. B **34**, 1 (1986)
- 13. H. Manaa, C. Pedrini, R. Moncorge: OSA Proc. on Advanced Solid-
- State Lasers, Vol. 10, 1991, ed. by G. Dube, L. Chase, pp. 371–375
- 14. B.F. Aull, H.P. Jensen: IEEE J. Qauntum Electron. **QE-18**, 925 (1982)