

# Output properties of diode-pumped passively Q-switched 1.06 $\mu\text{m}$ Nd:GdVO<sub>4</sub> laser using a [100]-cut Cr<sup>4+</sup>:YAG crystal

Y.F. Ma · X. Yu · F.K. Tittel · R.P. Yan · X.D. Li ·  
C. Wang · J.H. Yu

Received: 8 February 2012 / Published online: 30 March 2012  
© Springer-Verlag 2012

**Abstract** A passively Q-switched 1.06  $\mu\text{m}$  Nd:GdVO<sub>4</sub> laser with a [100]-cut Cr<sup>4+</sup>:YAG saturable absorber was demonstrated. The output characteristics were investigated when the anisotropic transmission of Cr<sup>4+</sup>:YAG crystal and the incident pump power level were considered. The experimental results showed that it was feasible to generate laser with narrower pulse width ( $\tau_p$ ), higher pulse energy and peak power when the polarization direction of laser was parallel to the [001], [010], [00 $\bar{1}$ ], and [0 $\bar{1}$ 0] orientations of the Cr<sup>4+</sup>:YAG crystal. The different changes of  $\tau_p$  as a function of incident pump power was observed due to the anisotropy of transmission of Cr<sup>4+</sup>:YAG and the different gain levels (pump power levels). If the Cr<sup>4+</sup>:YAG was fully bleached as a result of high cavity gain or due to the laser polarization direction was parallel to the [001], [010], [00 $\bar{1}$ ], and [0 $\bar{1}$ 0] orientations,  $\tau_p$  was constant, otherwise  $\tau_p$  decreased when the gain increased.

## 1 Introduction

Q-switched lasers are used in many applications, such as laser lidar, remote sensing, micro-machining, and microsurgery [1, 2]. Compared with active Q-switching, the passive techniques have advantages such as lower cost, compactness, simplicity in set-up, and operation since they do

not require external control. Cr<sup>4+</sup>:YAG crystal as a saturable absorber has advantages of improved thermo-mechanical properties, large absorption cross section, low saturable intensity, and high damage threshold [3, 4].

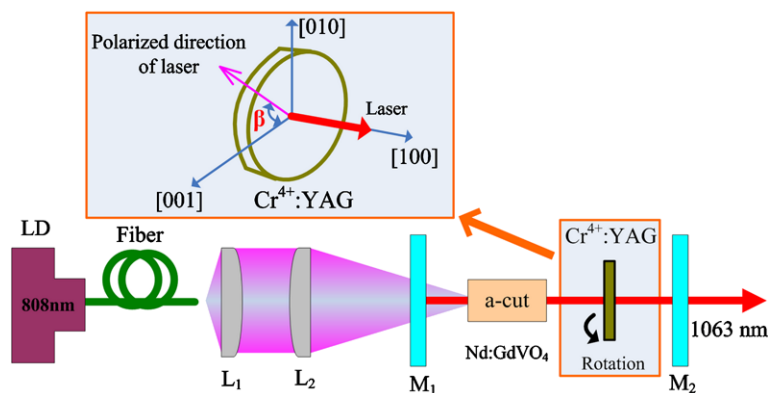
The output performance, such as output energy,  $E_p$ , pulse duration,  $\tau_p$ , pulse peak power,  $P_p$ , and pulse repetition rate,  $f$ , are important for passively Q-switched lasers. Many theoretical models and experimental results about passively Q-switched lasers have appeared in the literature [5–12]. But some of the theoretical opinions and experimental observations differed in these publications. For example, Degnan [5] reported that  $\tau_p$  was not affected by the pump power level, but was determined by other parameters, such as the cavity length. This was confirmed by experimental results reported in [6–8]. However, Lu et al. [9] reported that  $\tau_p$  decreases when the incident pump power increases, which was supported by experimental investigations [10–12]. In our characterization studies of a Cr<sup>4+</sup>:YAG passively Q-switched Nd:GdVO<sub>4</sub> laser, both these two different experimental phenomenon were observed for the first time in the same experimental configuration, which we attributed them to the anisotropy of the Cr<sup>4+</sup>:YAG transmission and the different laser cavity gain. The anisotropic transmission of Cr<sup>4+</sup>:YAG crystal was previously reported in [13, 14], but to the best of our knowledge, there was no researcher to report the above two changes of  $\tau_p$  simultaneously, not to mention giving an explanation.

In this paper, we demonstrated a passively Q-switched 1.06  $\mu\text{m}$  Nd:GdVO<sub>4</sub> laser with a [100]-cut Cr<sup>4+</sup>:YAG crystal by 808 nm laser diode pumping. The output characteristics were investigated when the anisotropic transmission of Cr<sup>4+</sup>:YAG crystal and the different gain levels were considered. The experimental results showed that we could obtain a pulsed laser output with narrower  $\tau_p$ , higher  $E_p$  and  $P_p$  when the polarization direction of laser was paral-

Y.F. Ma (✉) · X. Yu · R.P. Yan · X.D. Li · C. Wang · J.H. Yu  
National Key Laboratory of Science and Technology on Tunable  
Laser, Harbin Institute of Technology, Harbin 150001, China  
e-mail: Yufei.Ma@rice.edu

Y.F. Ma · F.K. Tittel  
Department of Electrical and Computer Engineering, Rice  
University, 6100 Main Street, Houston, TX 77005, USA

**Fig. 1** Experimental setup of a passively Q-switched Nd:GdVO<sub>4</sub> laser and the schematic of crystallographic axis of the Cr<sup>4+</sup>:YAG, polarized laser, and a rotation angle  $\beta$



rel to the  $[001]$ ,  $[010]$ ,  $[00\bar{1}]$ , and  $[0\bar{1}0]$  orientations of the Cr<sup>4+</sup>:YAG crystal, respectively. The different changes of  $\tau_p$  as a function of incident pump power were observed due to the anisotropic transmission of Cr<sup>4+</sup>:YAG crystal and the different cavity gain levels.

## 2 Experimental setup

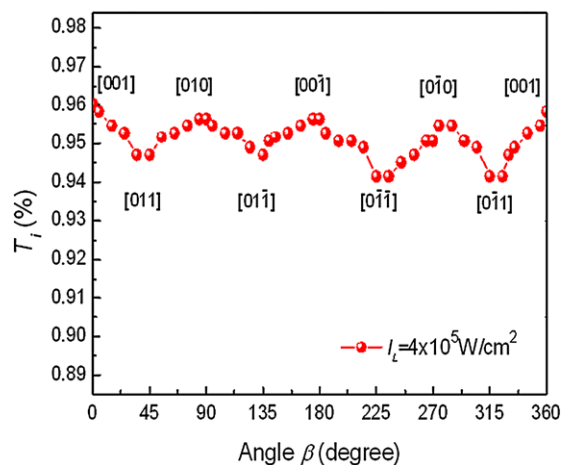
The experimental configurations of a  $[100]$ -cut Cr<sup>4+</sup>:YAG passively Q-switched Nd:GdVO<sub>4</sub> laser using 808 nm laser-diode end-pumping is shown in Fig. 1.

The 808 nm pump source was a fiber coupled laser diode. L<sub>1</sub> and L<sub>2</sub> were a set of collimating and focusing lenses. The a-cut Nd:GdVO<sub>4</sub> crystal with a 0.3 at.% Nd<sup>3+</sup> concentration had dimensions of  $3 \times 3 \times 8$  mm<sup>3</sup>. The crystal was wrapped with indium foil and placed into water-cooled copper heat sink. The input mirror M<sub>1</sub> was a flat mirror with an antireflection coating at 808 nm and a high reflectivity at 1063 nm. The output coupler M<sub>2</sub> had a transmission of 30 % at 1063 nm. The  $[100]$ -cut Cr<sup>4+</sup>:YAG crystal with initial transmission  $T_0$  of 90 % acted as a saturable absorber. The cavity length was about 160 mm.

## 3 Experimental results and discussions

The dependence of energy transmission  $T_i$  of the Cr<sup>4+</sup>:YAG crystal with the angle between the crystallographic axis and the polarization direction of laser was investigated firstly. The crystal was excited with a linearly polarized, AO Q-switched 1063 nm Nd:GdVO<sub>4</sub> laser propagating along its growth direction. The schematic of crystallographic axis of Cr<sup>4+</sup>:YAG, polarized laser and rotation angle  $\beta$  is same as shown in Fig. 1. The laser was focused by a lens to increase the power density  $I_L$ . The crystal was fixed on a precise rotatable mount in order to measure the rotation dependent  $T_i$ . The experimental results are shown in Fig. 2.

From Fig. 2, we can see that at a power density of  $I_L = 4 \times 10^5$  W/cm<sup>2</sup>, there is an obvious anisotropic energy transmission.  $T_i$  varied with the angle  $\beta$ , and the almost same  $T_i$

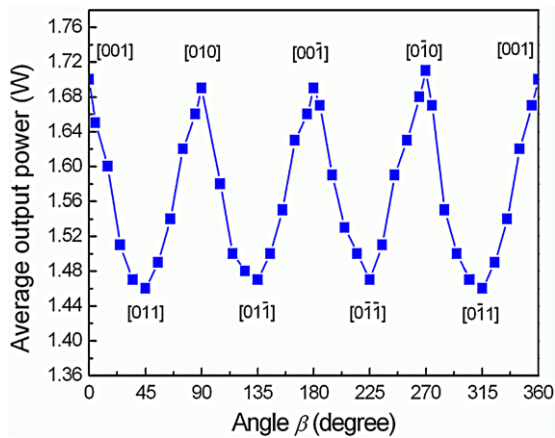


**Fig. 2**  $T_i$  of the  $[100]$ -cut Cr<sup>4+</sup>:YAG crystal as a function of rotation angle

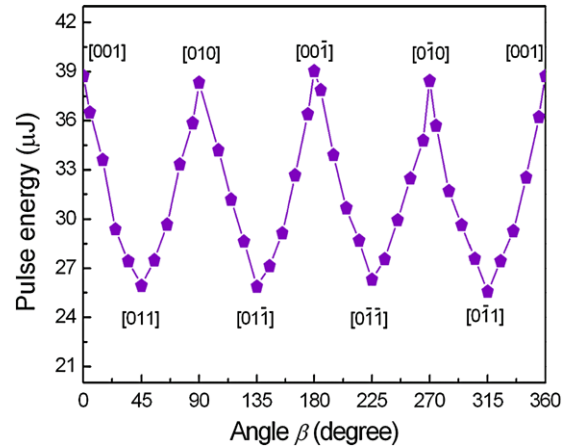
peaks appeared every 90°. When  $\beta$  was 0°, 90°, 180°, and 270°, the polarization direction of laser was parallel to the  $[001]$ ,  $[010]$ ,  $[00\bar{1}]$ , and  $[0\bar{1}0]$  orientations of Cr<sup>4+</sup>:YAG, respectively, and  $T_i$  had the maximum value. When the polarization direction of laser was parallel to the  $[011]$ ,  $[0\bar{1}1]$ ,  $[01\bar{1}]$ ,  $[0\bar{1}\bar{1}]$ , orientations ( $\beta = 45^\circ$ ,  $135^\circ$ ,  $225^\circ$ , and  $275^\circ$ , respectively),  $T_i$  was minimum.

The output characteristics of passively Q-switched Nd:GdVO<sub>4</sub> laser with the  $[100]$ -cut Cr<sup>4+</sup>:YAG crystal were investigated. The experimental configuration is shown in Fig. 1. The incident pump power was 7.5 W and the Cr<sup>4+</sup>:YAG crystal was rotated in a full 360°. The measured results are shown in Figs. 3–7.

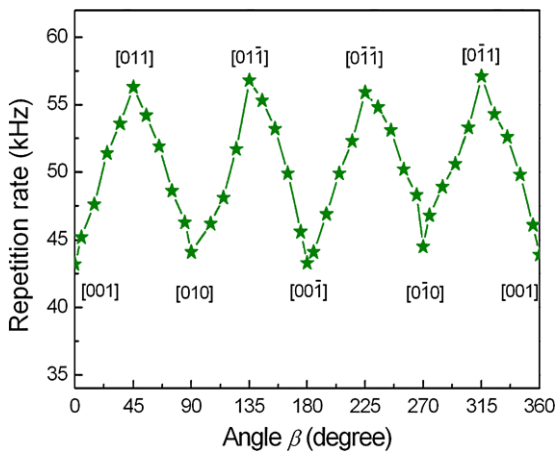
From Fig. 3, we can see that the maximum average output power was obtained when  $\beta$  was 0°, 90°, 180°, and 270°, respectively. This is because of when the polarization direction of laser was parallel to the  $[001]$ ,  $[010]$ ,  $[00\bar{1}]$ , and  $[0\bar{1}0]$  orientations, the  $T_i$  of Cr<sup>4+</sup>:YAG crystal had a maximum value. The loss was smaller and the laser was easier to oscillate compared with the directions of  $[011]$ ,  $[0\bar{1}1]$ ,  $[01\bar{1}]$ , and  $[0\bar{1}\bar{1}]$ .



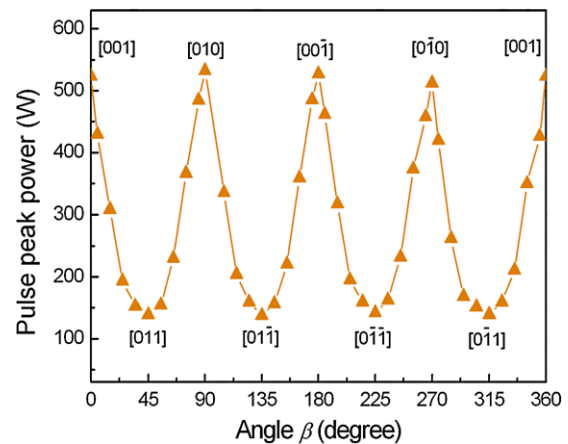
**Fig. 3** Relation between the average output power and the rotation angle  $\beta$



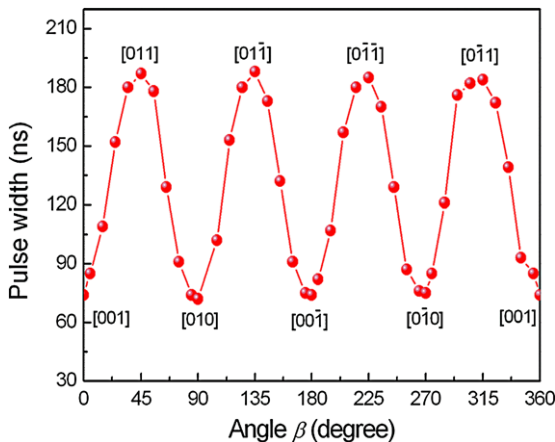
**Fig. 6** Relation between the pulse energy and the rotation angle  $\beta$



**Fig. 4** Relation between the repetition rate and the rotation angle  $\beta$



**Fig. 7** Relation between the pulse peak power and the rotation angle  $\beta$



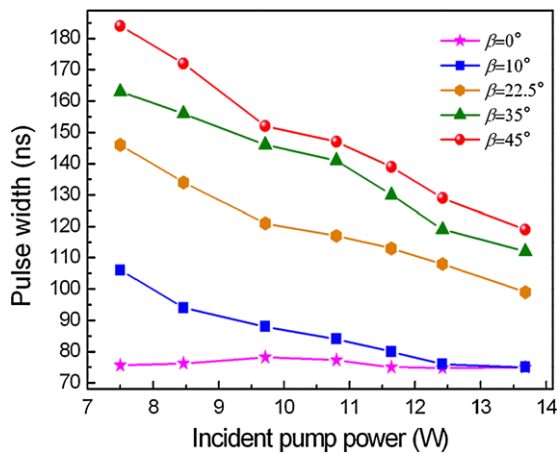
**Fig. 5** Relation between the pulse width and the rotation angle  $\beta$

The modulation depth  $Q_o$  of Cr<sup>4+</sup>:YAG crystal is approximately directly proportional to the difference of  $T_i$  and  $T_o$ .  $Q_o$  peaks emerged when  $\beta$  was 0°, 90°, 180°, and 270°,

respectively. Therefore, at these  $\beta$  values, the repetition rate and the pulse width were at a minimum and the pulse energy and the pulse peak power were maximum, as shown in Figs. 4, 5, 6 and 7.

There are four sinusoidal modulations of  $T_i$  in a full 360°. We studied the pulse width as a function of incident pump power at different rotation degree within 45° to illustrate the law of the full 360°. The measured results are depicted in Fig. 8.

From Fig. 8, we can see that at the same incident pump power, the narrowest  $\tau_p$  was achieved when  $\beta$  was 0° and the largest could be obtained when  $\beta$  was 45°. Due to the anisotropic transmission, at  $\beta$  of 0° and in the investigated pump power range (from 7.5 W to 13.7 W), the Cr<sup>4+</sup>:YAG crystal was fully bleached, and the pulse width was constant and ~76 ns. At other angles ( $\beta = 10^\circ, 22.5^\circ, 35^\circ, \text{ and } 45^\circ$ , respectively),  $\tau_p$  was a decreasing function of the incident pump power, and it would decrease to a constant value of 76 ns ( $\beta = 0^\circ$ ) when the incident pump power increased. Because the gain of the laser oscillator is a function of the incident pump power level, as shown in Fig. 8, the gain impacts



**Fig. 8** Pulse width as a function of incident pump power at different rotation degree

the pulse duration. If the  $\text{Cr}^{4+}$ :YAG was fully bleached as the result of high gain or due to the laser polarization direction was parallel to the  $[001]$ ,  $[010]$ ,  $[00\bar{1}]$ , and  $[0\bar{1}0]$  orientations,  $\tau_p$  was constant, otherwise  $\tau_p$  decreased when the gain increased.

#### 4 Conclusions

In conclusion, we have demonstrated a passively Q-switched  $1.06 \mu\text{m}$  Nd:GdVO<sub>4</sub> laser with a  $[100]$ -cut  $\text{Cr}^{4+}$ :YAG saturable absorber by 808 nm laser diode pumping. The output characteristics were investigated when the anisotropic transmission of  $\text{Cr}^{4+}$ :YAG crystal and the different gain levels were considered. The experimental results showed that a pulsed laser output with narrower  $\tau_p$ , higher  $E_p$  and  $P_p$

can be achieved when the laser polarization direction was parallel to the  $[001]$ ,  $[010]$ ,  $[00\bar{1}]$ , and  $[0\bar{1}0]$  orientations of  $\text{Cr}^{4+}$ :YAG, respectively. Two different effects of constant  $\tau_p$  and varying  $\tau_p$  were observed due to the anisotropy of transmission of  $\text{Cr}^{4+}$ :YAG crystal and the gain level. If the  $\text{Cr}^{4+}$ :YAG was fully bleached as the result of high gain or due to the laser polarization direction was parallel to the  $[001]$ ,  $[010]$ ,  $[00\bar{1}]$ , and  $[0\bar{1}0]$  orientations of  $\text{Cr}^{4+}$ :YAG, which was more easily to be complete bleached,  $\tau_p$  was constant, otherwise  $\tau_p$  decreased when the gain increased.

#### References

1. H.-H. Wu, S.-F. Chen, C.-H. Cheng, *J. Opt. A, Pure Appl. Opt.* **9**, 376 (2007)
2. Y. Ma, Y. Zhang, X. Yu, X. Li, F. Chen, R. Yan, *Opt. Commun.* **284**, 2569 (2011)
3. J. Liu, C. Wang, S. Liu, W. Tian, L. Li, S. Liu, M. Liu, *J. Mod. Opt.* **55**, 1971 (2008)
4. Y. Ma, X. Yu, X. Li, R. Fan, J. Yu, *Appl. Opt.* **50**, 3854 (2011)
5. J.J. Degnan, *IEEE J. Quantum Electron.* **31**, 1890 (1995)
6. F.Q. Liu, H.R. Xia, S.D. Pan, W.L. Gao, D.G. Ran, S.Q. Sun, Z.C. Ling, H.J. Zhang, S.R. Zhao, J.Y. Wang, *Opt. Laser Technol.* **39**, 1449 (2007)
7. G. Xiao, M. Bass, *IEEE J. Quantum Electron.* **33**, 41 (1997)
8. M. Montes, C. delas Heras, D. Jaque, *Opt. Mater.* **28**, 408 (2006)
9. M. Lu, C.R. Chatwin, R.C.D. Young, P.M. Birch, *Opt. Lasers Eng.* **47**, 617 (2009)
10. J. Liu, J. Yang, J. He, *Opt. Laser Technol.* **35**, 4314 (2003)
11. C. Du, J. Liu, Z. Wang, G. Xu, X. Xu, K. Fu, X. Meng, Z. Shao, *Opt. Laser Technol.* **34**, 699 (2002)
12. Y.F. Ma, R.W. Fan, X. Yu, X.D. Li, D.Y. Chen, J.H. Yu, *Laser Phys.* **21**, 1570 (2011)
13. H. Eilers, K.R. Hoffman, W.M. Dennis, S.M. Jacobsen, W.M. Yen, *Appl. Phys. Lett.* **61**, 2958 (1992)
14. X. Zhang, S. Zhao, Q. Wang, S. Zhang, L. Sun, X. Liu, S. Zhang, H. Chen, *J. Opt. Soc. Am. B* **18**, 770 (2001)