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# Passive Q-switching of a diode-side-pumped Nd-doped mixed gadolinium yttrium vanadate bounce laser

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**ABSTRACT** High power passive Q-switching was achieved with a pulse width of 18–32 ns by using a diode-side-pumped Nd:Gd<sub>0.6</sub>Y<sub>0.4</sub>VO<sub>4</sub> bounce amplifier. An average output power of > 8 W was obtained at a pump power of 39 W. The peak power of the Q-switched output was adjusted within 1.9–5.2 kW by changing the Nd concentration. The mixed vanadates showed significantly higher Q-switching performances in comparison with pure Nd:GdVO<sub>4</sub>.

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## 1 Introduction

Passively Q-switched solid-state lasers employing a saturable absorber (SA) such as Cr<sup>4+</sup>:YAG have been intensely investigated because no additional electronic elements are required and they provide robust and inexpensive alternatives to actively Q-switched systems. Passively Q-switched lasers, mainly formed in an end-pumped microchip laser geometry, mostly have the drawback of lower output power due to the residual absorption of the saturable absorbers producing output powers of ~ 4 W [1–3].

A diode-side-pumped bounce amplifier configuration based on neodymium-doped vanadate slabs including Nd:YVO<sub>4</sub> and Nd:GdVO<sub>4</sub>, in which extremely high inversion density is produced below the pump face of the amplifier, is a promising candidate for producing high power outputs at ultra-high efficiency and high beam quality [4–6]. The bounce amplifier geometry will allow an easier power scaling of the passively Q-switched laser systems in comparison with the conventional end-pumping geometry, even with the residual absorption of the saturable absorbers.

However, very large emission cross sections of Nd:YVO<sub>4</sub> and Nd:GdVO<sub>4</sub> frequently prevent stable passive Q-switching, and yield pre-lasing and pulse broadening in the Q-switching operation. Villafana et al. demonstrated a passively Q-switched Nd:YVO<sub>4</sub> bounce laser using a LiF:F<sup>2</sup> saturable absorber, demonstrating ~ 2 W stable Q-switched

output with minimum pulse width of ~ 10 ns [7]. However, the peak power of the Q-switched pulses was limited to ~ 0.5 kW.

Use of *c*-cut Nd-doped vanadates (or *a*-axis-pumped vanadates) is one solution to produce successfully high peak power Q-switching pulses [8]. However, the use of *c*-cut vanadates showing low absorption (~ 11 cm<sup>-1</sup>) for the pump diode wavelength forms a deeper pumped region (~ 3 mm), which would impact total performance of the bounce laser system.

Mixed vanadate technology, in which, for example, a fraction of Gd ions in the GdVO<sub>4</sub> are replaced by Y ions, is capable of customizing the fluorescence spectrum, thereby allowing production of adjustable laser parameters such as stimulated emission cross section and fluorescence lifetime [9–12]. Recently, we demonstrated an actively Q-switched Nd-doped mixed gadolinium yttrium vanadate (Nd:Gd<sub>x</sub>Y<sub>1-x</sub>VO<sub>4</sub>) laser with custom-made pulse repetition frequency (PRF) range of 100–650 kHz and peak power range of 1–3 kW [13].

In this paper, we extend our mixed vanadate bounce laser work to develop a high average power passively Q-switched laser. We demonstrate stable Q-switching operation of a diode-side-pumped mixed gadolinium yttrium vanadate (Nd:Gd<sub>x</sub>Y<sub>1-x</sub>VO<sub>4</sub>) bounce laser at a pulse repetition frequency range of up to 150 kHz and a pulse duration in the range of 18–32 ns. The average power in Q-switched operation was 5.1–8.2 W at a pump power of 40 W. The corresponding peak power of 1.9–5.2 kW was achieved.

## 2 Experiments

### 2.1 Continuous-wave operation

Figure 1 shows the schematic diagram of the bounce amplifier setup. The amplifiers used for the experiment were two *a*-cut, 2 mm × 5 mm × 20 mm slabs, i.e. a 1 at. % Nd-doped Gd-rich mixed vanadate Nd:Gd<sub>0.6</sub>Y<sub>0.4</sub>VO<sub>4</sub> slab and a 1.5 at. % Nd-doped Gd<sub>0.6</sub>Y<sub>0.4</sub>VO<sub>4</sub> slab (Hortek Crystal Co., Ltd). The amplifiers, having two end faces antireflection coated for 1 μm, were wedged to the normal of the pump face to prevent self-lasing within the crystal. A continuous-wave (cw) 808-nm diode array output was delivered by a 12.7 mm cylindrical lens (CL<sub>D</sub>) to be a line focus with dimensions of ~ 0.2 mm × 15 mm on the pump face of the amplifiers.

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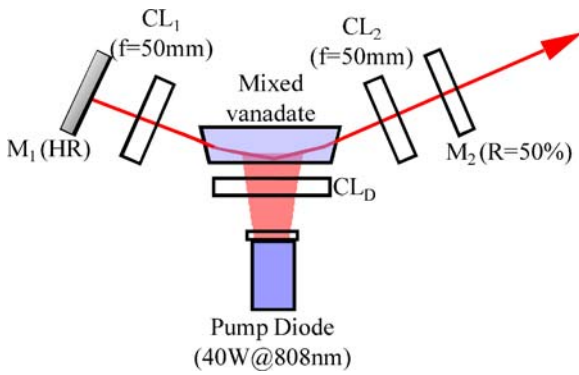


FIGURE 1 Experimental setup for cw operation

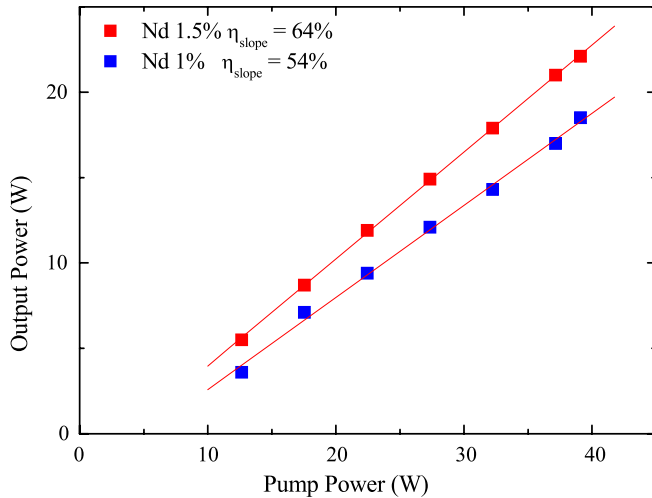


FIGURE 2 Experimental plots of the cw output power as a function of pump power

The laser cavity was the same as that used in our previous publications, and it was composed of a highly reflective flat mirror  $M$  for 1064 nm, a partially reflective flat output coupler OC also for 1064 nm and two cylindrical lenses,  $CL_1$  and  $CL_2$  ( $f = 50$  mm), in the vertical direction for spatially matching the laser mode to the ellipsoidal gain region. The internal incidence angle of the laser mode with respect to the pump face normal was  $\sim 80^\circ$ .

We measured the cw output power from the mixed vanadate lasers using various output couplers. The optimum reflectivity of the output coupler for both the mixed vanadate crystals used was  $\sim 50\%$ . The 1.5 at. % Nd:Gd<sub>0.6</sub>Y<sub>0.4</sub>VO<sub>4</sub> showed slightly lower lasing threshold and higher slope efficiency in comparison with those of the 1 at. % Nd:Gd<sub>0.6</sub>Y<sub>0.4</sub>VO<sub>4</sub>. Maximum output power of 22 W was obtained at the pump power of 39 W (Fig. 2), corresponding to a slope efficiency of 64%. The higher Nd doping produced a relatively shallower pumped region, yielding higher laser gain and resulting in highly efficient laser oscillation with the slope efficiency of  $> 60\%$ .

## 2.2 Passively Q-switched operation

A 4.4-mm-long, 0.3%-doped Cr<sup>4+</sup>:YAG crystal with initial optical density of 0.42 was used as saturable absorber for passive Q-switching. The criterion to achieve stable

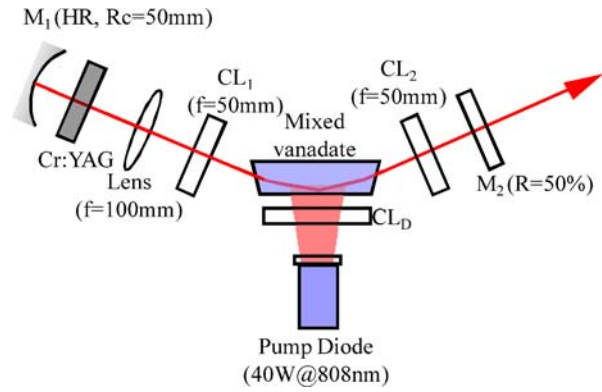


FIGURE 3 Experimental setup for passively Q-switched operation

Q-switching is given by the following expression [14]:

$$\frac{\ln(1/T_0^2)}{\ln(1/T_0^2) + \ln(1/R) + L} \frac{\sigma_a A_e}{\sigma_e A_a} > \frac{1}{1-\beta}, \quad (1)$$

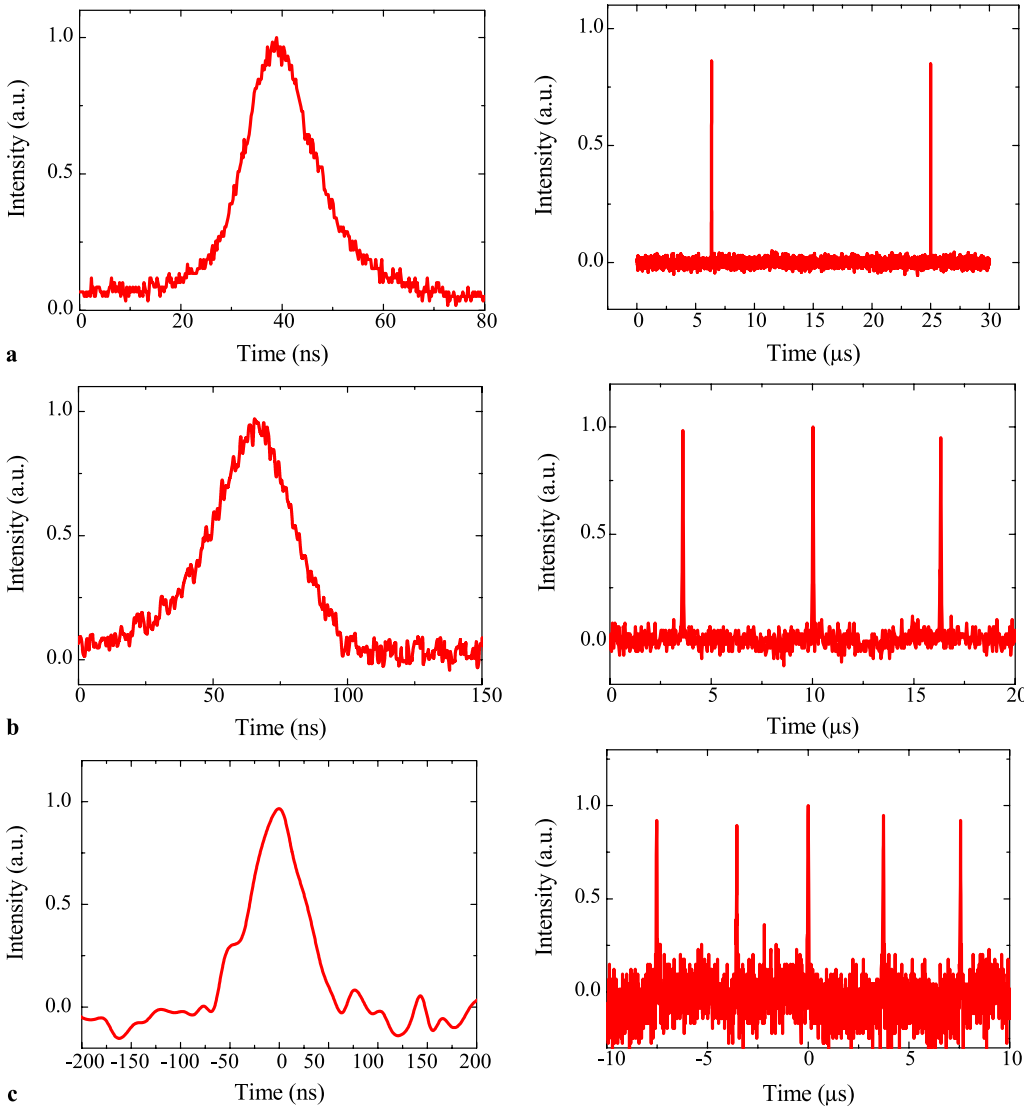
where  $T_0$  ( $\sim 0.45$ ) is the transmission of the unsaturated saturable absorber,  $R$  is the reflectivity of the output coupler,  $L$  (0.02) is the internal loss of the cavity,  $\sigma_a$  ( $3 \times 10^{-18}$  cm<sup>2</sup>) is the ground-state absorption cross section of the saturable absorber,  $\sigma_e$  ( $7 \times 10^{-18}$  cm<sup>2</sup>) is the stimulated emission cross section of the gain medium,  $\beta$  (0.25) is the ratio of the excited-state absorption cross section to the ground-state absorption cross section and  $A_e/A_a$  is the ratio of the laser mode areas in the gain medium and in the saturable absorber. By substituting the physical parameters into (1), stable Q-switching requires the value  $A_e/A_a > 0.8$ . To satisfy the criterion and achieve stable passive Q-switching, the cavity configuration was modified as shown in Fig. 3. The Cr:YAG crystal was located near a confocal point of the focusing lens ( $f = 100$  mm) and a spherical mirror ( $R = 50$  mm). The reflectivity of the output coupler was 10%.

Above the pump power ( $\sim 20$  W) required to achieve threshold, the laser started to operate in Q-switched mode, and its output exhibited a smooth temporal profile without mode beating as shown in Fig. 4a. In the case of the 1 at. % Nd-doped Gd-rich mixed vanadate Nd:Gd<sub>0.6</sub>Y<sub>0.4</sub>VO<sub>4</sub> slab amplifier, the threshold of Q-switched operation was  $\sim 18$  W and a slope efficiency of 20% was obtained. The pulse repetition frequency was 22 kHz near threshold, increasing to 52 kHz at the maximum pump level (39 W). The pulse width was typically  $\sim 20$  ns at all pump levels, resulting in peak powers in the range of 3.7–5.2 kW.

In the case of the 1.5 at. % Nd-doped mixed vanadate slab amplifier, a slope efficiency of 40% was achieved and the maximum output power reached 8.2 W. As shown in Fig. 4b, the output exhibited a slightly broader pulse width ( $\sim 30$  ns) in the PRF region of 45–153 kHz, resulting in the peak power being limited to  $\sim 1.9$  kW. The higher laser gain and the reduced fluorescence lifetime observed in the highly doped mixed vanadate prevent larger energy storage.

To compare the performances of the mixed vanadates with conventional 1.1 at. % Nd:GdVO<sub>4</sub>, we also measured the Q-switched output power from a Nd:GdVO<sub>4</sub> laser (Fig. 4c).

Figure 5 summarizes the Q-switched performance of three vanadate lasers including two mixed vanadates and GdVO<sub>4</sub>.



**FIGURE 4** Temporal evolutions of the Q-switched pulse in (a) 1 at. % Nd:Gd<sub>0.6</sub>Y<sub>0.4</sub>VO<sub>4</sub>, (b) 1.5 at. % Nd:Gd<sub>0.6</sub>Y<sub>0.4</sub>VO<sub>4</sub> and (c) Nd:GdVO<sub>4</sub>

The peak power of the output was limited to  $\sim 0.6$  kW, while the average output power was similar to that of 1.5 at. % Nd-doped mixed vanadate. The pulse width of the output was typically  $\sim 50$  ns. At the maximum pump level, the PRF increased to 252 kHz, which is comparable to the maximum PRF of stable Q-switching expected from the recovery time of saturable absorption ( $\sim 3$   $\mu$ s). These results show that too strong gain in the Nd:GdVO<sub>4</sub> amplifier prevented production of intense Q-switched pulses with high peak power.

The output exhibited a mixed-mode profile as shown in Fig. 6. The corresponding  $M^2$  factor of the output along the horizontal axis was  $\sim 2.5$ , while the  $M^2$  factor along the vertical was 1.2, and it was relatively small in comparison with that ( $\sim 5$ ) of the cw output previously reported by us. This means that the saturable absorber produced a limiting aperture for the laser to operate at lower-order modes.

### 3 Numerical simulations

The rate equations for passive Q-switching with the Cr:YAG crystal can be written as follows [15]:

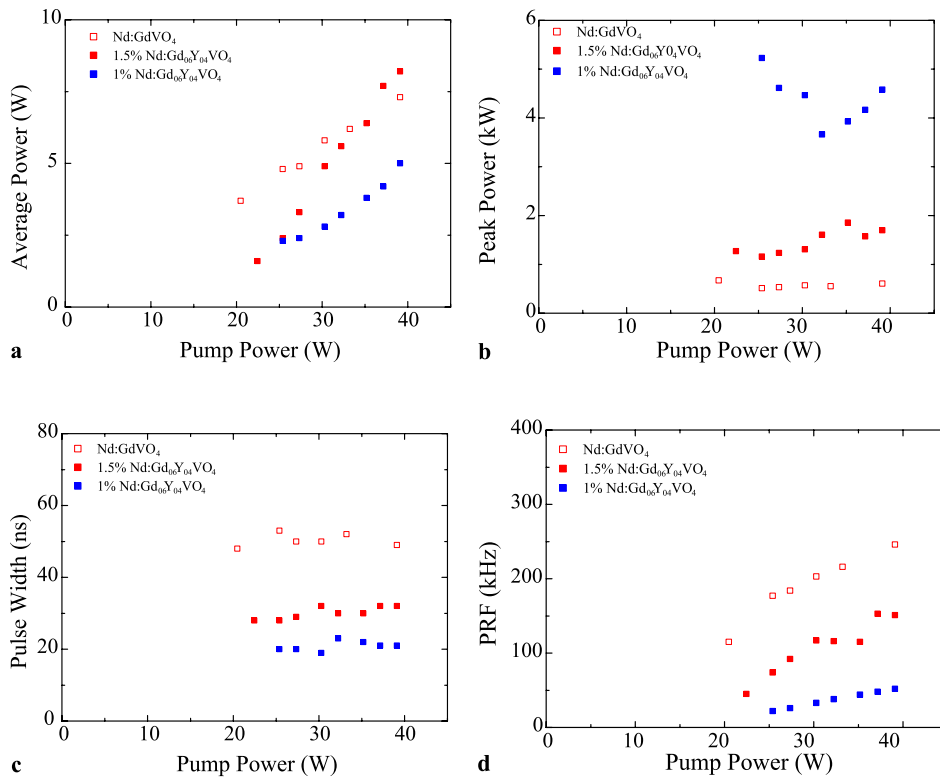
$$\frac{dn}{dt} = \frac{\eta_{\text{pump}} P_{\text{pump}}}{V} \left( \frac{\lambda_p}{hc} \right) - c\sigma_e n \varphi - c\sigma_{\text{esa}} n \varphi - \frac{n}{\tau_f}, \quad (2)$$

$$\frac{d\varphi}{dt} = c\sigma_e n \varphi \frac{l_g}{l} - c\sigma_a n_a \varphi \frac{l_a}{l} \frac{A_e}{A_a} - \frac{\varphi}{\tau_c} + \frac{n}{\tau_f} \Omega, \quad (3)$$

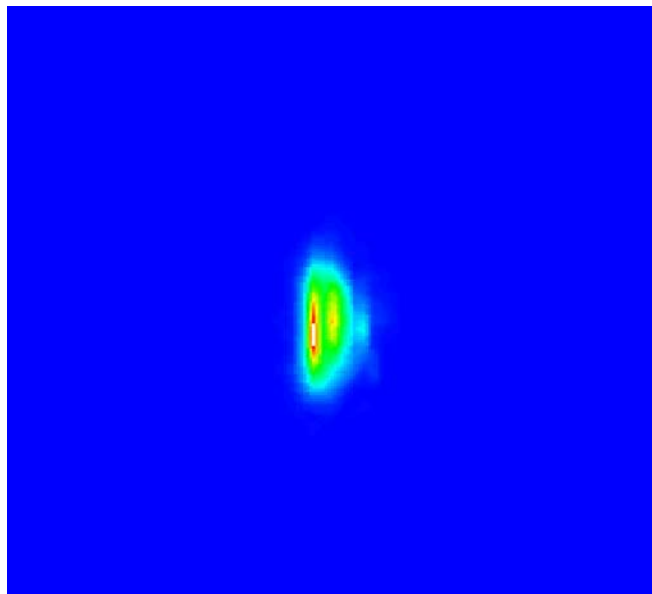
$$\frac{dn_a}{dt} = c\sigma_a (N_0 - n) \varphi \frac{A_e}{A_a} - \frac{n_a}{\tau_a}, \quad (4)$$

$$\frac{1}{\tau_c} = \frac{2l}{c} \left( \ln \left( \frac{1}{R} \right) + L_i \right), \quad (5)$$

where  $n$  is the inversion population density,  $\varphi$  is the photon density,  $\eta_{\text{pump}}$  is the pump quantum efficiency,  $P_{\text{pump}}$  is the pump diode power,  $V$  is the volume of the pumped region,  $hc/\lambda_p$  is the photon energy of the pump diode,  $c$  is the velocity of light,  $\sigma_{\text{esa}}$  is the excited-state absorption cross section of the Cr:YAG crystal,  $\tau_f$  is the fluorescence lifetime of the upper level,  $\tau_a$  is the decay time of the excited level in the Cr:YAG crystal,  $l_g$  is the length of the pumped region,  $l_a$  is the length of the saturable absorber,  $l$  is the cavity length,  $\Omega$  is the factor accounting for the fraction of spontaneous photons contributing to the lasing,  $n_a$  is the population density of the excited state in the Cr:YAG crystal and  $N_0$  is the total



**FIGURE 5** Experimental data in Q-switching operation. (a) Average power, (b) peak power, (c) pulse width and (d) pulse repetition frequency as a function of pump power



**FIGURE 6** Spatial form of the Q-switched output

number of the Cr ions, respectively. By using the values given in Table 1, we numerically simulated the theoretical temporal evolution of Q-switched pulses. The ratio  $A_e/A_a$  of the laser mode areas was difficult to estimate accurately, and it was optimized within the range of 1.3–2. Three curves (red and blue solid curves, and red broken curve) in Fig. 7 show the simulated plots as a function of the pump power. The simulated PRF in the mixed vanadates increases as the pump power increases, and there is good agreement between simulations and experiments. The simulated pulse width of the Q-switched output in the mixed vanadates is also consistent

with the experimental one, though there is a slight mismatch of the absolute value between simulation and experiment.

In the case of GdVO<sub>4</sub>, the simulated pulse width of the Q-switched output increases gradually above the pump power of 30 W, while the experimental one is constant within the whole pump range. The discrepancy between simulation and experiment in Nd:GdVO<sub>4</sub> will be due to uncertainties of the physical parameters.

#### 4 Conclusion

We have successfully demonstrated passive Q-switching of mixed gadolinium yttrium vanadate Nd:Gd<sub>x</sub>Y<sub>1-x</sub>VO<sub>4</sub> slab lasers with a bounce amplifier geometry, for the first time. By changing the Nd doping, we can adjust the peak power (1.9–5.2 kW) and pulse repetition frequency regions (22–145 kHz) of the stable Q-switching operation. We have also compared the Q-switching performance of the mixed vanadate lasers with that of a conventional Nd:GdVO<sub>4</sub> laser. The mixed vanadate slabs are capable of producing higher peak power Q-switched pulses in comparison with the conventional vanadate crystals.

At even higher pumping, severe thermal effects of the saturable absorber will significantly impact the performance of the stable Q-switching operation. For further power scalability, careful cavity design, that would include thermal issues of the Cr:YAG crystal, will be necessary.

We can also extend the system to optimize the product of peak and average powers for nonlinear frequency conversion by varying the initial absorption of the Cr:YAG crystal.

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Parameter, unit	Value
$n_{\text{pump}}$	Pump quantum efficiency 1
$P_{\text{pump}}$	Maximum pump power, W 40
$V$	Laser volume, cm <sup>3</sup> 0.003
$\lambda_p$	Wavelength of pump diode, cm $0.8 \times 10^{-4}$
$\sigma_e$	Stimulated emission cross section, cm <sup>2</sup> $7 \times 10^{-19}$ (Nd:Gd <sub>0.6</sub> Y <sub>0.4</sub> VO <sub>4</sub> ) $12.6 \times 10^{-19}$ (Nd:GdVO <sub>4</sub> )
$\tau_f$	Spontaneous fluorescence lifetime, $\mu\text{s}$ 100 (1 at. % Nd:Gd <sub>0.6</sub> Y <sub>0.4</sub> VO <sub>4</sub> ) 85 (1.5 at. % Nd:Gd <sub>0.6</sub> Y <sub>0.4</sub> VO <sub>4</sub> ) 100 (Nd:GdVO <sub>4</sub> )
$\sigma_a$	Absorption cross section, cm <sup>2</sup> $30 \times 10^{-19}$
$\sigma_{\text{esa}}$	Absorption cross section of excited state, cm <sup>2</sup> $7.5 \times 10^{-19}$
$\tau_a$	Recovery time of saturable absorption, $\mu\text{s}$ 3
$l$	Cavity length, cm 40
$l_e$	Active length, cm 2
$l_a$	Absorption length, cm 4.4
$\Omega$	Spontaneous emission factor $1.0 \times 10^{-5}$
$R$	Reflectivity of output coupler 0.1
$L_i$	Internal dissipative optical losses 0.02
$A_e/A_a$	Ratio of the laser mode areas in the amplifier and absorber 1.9 (1 at. % Nd:Gd <sub>0.6</sub> Y <sub>0.4</sub> VO <sub>4</sub> ) 1.4 (1.5 at. % Nd:Gd <sub>0.6</sub> Y <sub>0.4</sub> VO <sub>4</sub> )

TABLE 1 Physical parameters for numerical simulation

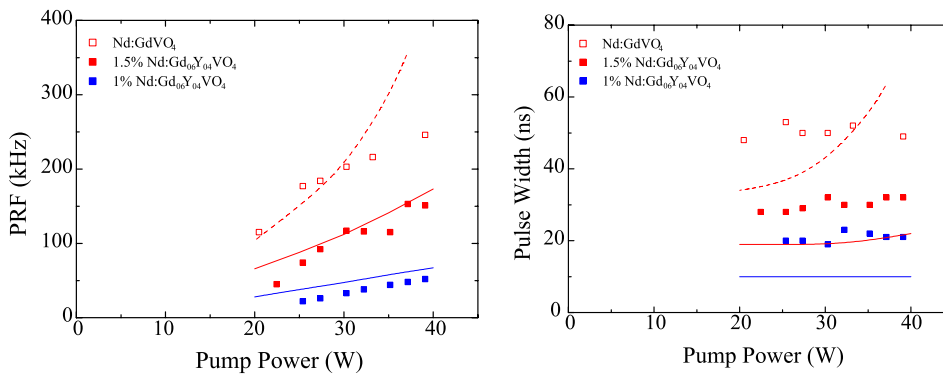


FIGURE 7 PRF and pulse width of Q-switched output from mixed vanadate slab lasers. Red and blue squares show experimental results in 1.5 at. % Nd:Gd<sub>0.6</sub>Y<sub>0.4</sub>VO<sub>4</sub> and 1 at. % Nd:Gd<sub>0.6</sub>Y<sub>0.4</sub>VO<sub>4</sub> lasers, respectively. Red and blue solid curves show the simulated plots. Results for pure Nd:GdVO<sub>4</sub> are shown for comparison

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## REFERENCES

- C. Li, J. Song, D. Shen, N.S. Kim, J. Lu, K. Ueda, Appl. Phys. B **70**, 471 (2000)
- C. Du, J. Liu, Z. Wang, G. Xu, X. Xu, K. Fu, X. Meng, Z. Shao, Opt. Laser Technol. **34**, 699 (2002)
- A. Agnesi, S. Dell'acqua, Appl. Phys. B **76**, 351 (2003)
- J.E. Bernard, A.J. Alcock, Opt. Lett. **19**, 1861 (1994)
- A. Minassian, B. Thompson, M.J. Damzen, Appl. Phys. B **76**, 341 (2003)
- B. Thompson, A. Minassian, M.J. Damzen, 42 W Nd:GdVO<sub>4</sub> bounce laser oscillator, in *Technical Digest of the Conference on Lasers and Electro-Optics (CLEO 2004)*, San Francisco, CA (2004), paper CThJ2
- E.R. Villafana, A.V. Kir'yanov, A. Minassian, M.J. Damzen, Laser Phys. **15**, 1 (2005)
- S. Forget, F. Druon, F. Balembois, P. Georges, N. Landru, J.-P. Fe've, J. Lin, Z. Weng, Opt. Commun. **259**, 816 (2006)
- J. Liu, X. Meng, Z. Shao, M. Jiang, B. Ozygus, A. Ding, H. Weber, Appl. Phys. Lett. **83**, 1289 (2003)
- H. Zhang, J. Wang, C. Wang, L. Zhu, X. Hu, X. Meng, M. Jiang, Y.T. Chow, Opt. Mater. **23**, 449 (2003)
- J. Liu, Z. Wang, X. Meng, Z. Shao, B. Ozygus, A. Ding, H. Weber, Opt. Lett. **28**, 2330 (2003)
- Y.F. Chen, M.L. Ku, L.Y. Tsai, Y.C. Chen, Opt. Lett. **29**, 2279 (2004)
- T. Omatsu, M. Okida, A. Minassian, M.J. Damzen, Opt. Express **14**, 2727 (2006)
- J. Liu, B. Ozygus, S. Yang, J. Erhard, U. Seelig, A. Ding, H. Weber, X. Meng, L. Zhu, L. Qin, C. Du, X. Xu, Z. Shao, J. Opt. Soc. Am. B **20**, 652 (2003)
- W. Chen, Y. Inagawa, T. Omatsu, M. Tateda, N. Takeuchi, Y. Usuki, Opt. Commun. **194**, 401 (2001)