

Multicolor green to orange‑red emission of Tb3+ and Eu3+‑codoped tellurite glasses: Eu3+ concentration and Tb3+→ **Eu3+ energy transfer**

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Abstract

A series of the TeO₂–GeO₂–ZnO glass system was single and double doped with different Tb³⁺/Eu³⁺ ratios. Their luminescent and colorimetric properties were analyzed for possible use as phosphor materials in lighting devices. The characterization by X-ray difraction and Raman spectroscopy verifed the glassy nature of the fabricated samples. The luminescent properties of the doped glasses were analyzed by means of steady-state fluorescence and time-resolved spectroscopy. The Tb³⁺ excitation bands observed in the codoped samples while monitoring the 700 nm emission of Eu^{3+} , as well as the shortening of Tb³⁺ lifetime in presence of europium indicated a Tb³⁺ \rightarrow Eu³⁺ energy transfer, which main interaction type is electric dipole–dipole, according to Inokuti–Hirayama model. The $Eu^{3+} \to Tb^{3+}$ energy transfer also occurs in the samples but with lower efficiency. The CIE1931 chromaticity coordinates, upon different excitation wavelengths, show a multicolor tunning from green to orange-red due to the lanthanide concentration ratio and the $Tb^{3+} \to Eu^{3+}$ energy transfer process.

Keywords Tellurite glasses · Energy transfer · Photoluminescence · Europium · Terbium

1 Introduction

Optoelectronic devices like lasers, fber amplifers, displays devices, and light-emitting diodes (LEDs), among others, have been promoting research towards materials with good physical, stability, and chemical properties adequate to host rare-earth ions for the mentioned applications [[1–](#page-9-0)[6](#page-9-1)]. In several trivalent lanthanide ions, the absorption bands due to 4f–4f or 5d–4f transitions are located in the range of InGaN, GaN, and AlGaN (350–400 nm) near-ultraviolet LEDs (NUV-LEDs) that are compatible with solid-state lighting devices [\[7](#page-9-2)[–9](#page-9-3)]. Depending on the target application, a single ion or a combination of two or three rare-earth ions are introduced in oxide glasses to obtain multicolor or white light emission, following the red (R), green (G), and blue (B) system, tailoring the concentration of the lanthanide ions [\[9](#page-9-3), [10\]](#page-9-4). For example, Sołtys et al. [[11\]](#page-9-5) used a combination

of Dy^{3+} (blue, yellow), Tb^{3+} (green), and Eu^{3+} (red) ions in triply doped $PbO-B_2O_3$ -Al₂O₃-WO₃ glasses_, varying their concentrations, resulting in warm white light emission upon direct excitation of ${}^{6}H_{15/2} \rightarrow {}^{4}K_{17/2}$ transition of Dy³⁺ (390 nm) with the $Dy^{3+} \to Tb^{3+} \to Eu^{3+}$ energy transfer process. In this respect, Tb^{3+} ions with a 4 f^8 configuration and its dominant green emission from ${}^5D_4 \rightarrow {}^7F_5$ transition, upon NUV excitation, have been demonstrated to be a suitable donor for acceptor ions like Sm^{3+} and Eu^{3+} [\[6](#page-9-1), [12](#page-9-6), [13\]](#page-9-7). Jha et al. $[12]$ $[12]$ work on Tb³⁺/Sm³⁺-doped ZnO-P₂O₅-BaO-TiO₂ glasses revealed that the single-doped glasses presented pure green and orange emission. In contrast, the emission of the codoped glasses is possibly tuned from yellowish green to warm white through the energy transfer mechanism from Tb^{3+} to Sm³⁺. Similarly, the combination of trivalent terbium and europium ions has also been studied in several host materials due to the capability to tune the emission over a wide range of visible light, because of the energy transfer process that occurs from terbium to europium ions $[14–16]$ $[14–16]$ $[14–16]$. Even though the emission tunability is a desirable phosphor property for light-emitting diodes or display applications, in the last decade, few reports evaluated the luminescent and colorimetric properties of $\text{Tb}^{3+}/\text{Eu}^{3+}$ doubledoped glasses variating europium concentration. Among the

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glass host materials where the efect of europium content have been studied are the aluminum-zinc-phosphates [\[14](#page-9-8)], fluorophosphate $[16]$ $[16]$, borate $[15, 17]$ $[15, 17]$ $[15, 17]$ $[15, 17]$, fluorozirconate $[17]$ $[17]$, germanate $[18]$ $[18]$ $[18]$, multicomponent $[19]$ $[19]$ glasses, and oxide and oxyfuoride glass–ceramics [[20,](#page-10-3) [21](#page-10-4)]. Furthermore, the energy transfer mechanism has been exploited in these works to tune the global emission tonality. Given the amorphous nature of glasses, presenting short-range order, the spectroscopic characteristics of the trivalent lanthanide ions contained in them are determined by the local structure. According to Bocharova et al. [\[22](#page-10-5)], dopant ions in disordered solids can distribute randomly or with local concentration displacement. In this work, the authors found that the dopant distribution and the Tb³⁺ \rightarrow Eu³⁺ energy transfer depend on the dopant concentration, as well as the glass structure and glass composition.

Gopi et al. [[15\]](#page-9-10) work on Tb^{3+}/Eu^{3+} -doped borate glasses showed that at small increments of Eu^{3+} dopant from 0.1 to 1.0 mol%, under a 464 nm excitation, can shift the emission from greenish-yellow to orange-red. Also, Loos et al. [\[17](#page-10-0)] reported that Tb^{3+}/Eu^{3+} pair exhibit a similar chromatic behavior in borate and fuorozirconate glasses, upon excitation of 376 nm and 486 nm, confrming that the energy transfer process is enhanced as europium concentration increases from 0.1 to 0.5 mol%. Meanwhile, Wileńska et al. [[19\]](#page-10-2) studied Tb/Eu-doped multicomponent glasses with the composition $\text{AlF}_3\text{-TiO}_2\text{-}40\text{PbO-H}_3\text{BO}_3\text{-}Si\text{O}_2$ obtaining green, yellow, and cold-white light emission at different $xEu-(1 - x)Tb$ molar ratios while maintaining the total concentration of dopants below 1 mol%. de Queiroz et al. [\[16](#page-9-9)] increased the molar ratio of Tb^{3+}/Eu^{3+} up to 1.5 mol% in fluorophosphate glasses and studied their luminescent and chromatic properties varying the excitation wavelength between 355 and 360 nm. They evidenced that the emission tonality of the glasses lies in a broad region of the visible light, in the red, green, and blue regions. Raju et al. [[23\]](#page-10-6) only studied the photoluminescent properties of Tb^{3+} and Eu^{3+} single-doped TeO₂–B₂O₃–CdO–Li₂O–AlF₃ glasses up to concentrations of 1.25 mol%. It is worth mentioning that codoping tellurite glasses with Tb^{3+}/Eu^{3+} has not been studied in-depth for high europium molar concentration, performing timeresolved photoluminescence spectroscopy analysis. Regarding oxide glasses, those based on tellurium oxide $(TeO₂)$ and germanium oxide $(GeO₂)$ are amorphous materials with high thermal and chemical stability, mechanical resistance, low melting temperature, mid-IR transparency, and low phonon energies of 700 and 880 cm−1, respectively, that are appropriate as hosts for photoluminescent trivalent rare-earth ions [\[1](#page-9-0), [24–](#page-10-7)[26\]](#page-10-8). GeO₂ is a natural glass former, while TeO₂ is a conditional glass former, which requires to be combined with other compounds like alkaline oxides, sulfdes, fuorides or chlorides to form a stable glass $[26-28]$ $[26-28]$. GeO₂ improves the mechanical resistance of the matrix, as well as their thermal and chemical stability because of the higher bond strength of Ge–O bonds than Te–O bonds [[28](#page-10-9), [29](#page-10-10)]. Lesniak et al. [[30\]](#page-10-11) reported that for TeO_2-GeO_2 : Eu^{3+} glasses, GeO_2 concentration influences the fluorescence intensity ratio of $Eu³⁺$ ions, the refractive index and phonon energy of the glasses. Zinc oxide is often used in combination with $TeO₂$ because it contributes to stability, mechanical resistance, and high concentration acceptance (up to 40 mol% of ZnO) of the glass matrix. The addition of ZnO causes a disruption of the Te–O–Te bonds forming Te–O− terminal bonds, with nonbridging oxygens, depolymerizing the TeO₂ network $[31,$ [32](#page-10-13)]. The non-bridging oxygens, together with tellurium lone electron pairs, have been related to the capability of $TeO₂$ glass to accept and accommodate modifer cations and their optical properties [[33](#page-10-14)]. Thus, we report a comprehensive analysis of the luminescent, photometric, and colorimetric characteristics of TeO₂–GeO₂–ZnO (TGZ) glasses impurifed with trivalent terbium, europium, and both lanthanides, with emphasis on the high molar concentration of europium, while evaluating the effect of the Tb³⁺ \rightarrow Eu³⁺ energy transfer on the color tuning of the global emission of the doubledoped TGZ glasses, covering a wide range of europium concentrations from 0.5 to 2 mol%.

2 Experimental details

TGZ glasses with a composition of $(80-x-y)TeO₂$ –10GeO₂–1 $0ZnO:xTb^{3+}$, yEu^{3+} (mol%) were fabricated according to the previously reported methodology [[32\]](#page-10-13), using reagent grade oxides powders, $Eu(NO_3)_3.5H_2O$ and $Tb(CO_3)_3.2H_2O$ that were acquired from Sigma-Aldrich. Table [1](#page-2-0) presents mol% concentrations of the precursors and dopants used to synthesize the glasses. The obtained samples were cut and polished into $0.6 \times 0.6 \times 0.2$ cm blocks. The vibrational features of the glasses were studied through Raman spectroscopy using a Horiba LabRAM HR spectrometer equipped with a He–Ne laser (λ =632.8 nm). Photoluminescence spectra were registered by a Fluorolog fuorometer using a 450 W ozonefree Xe lamp. The fuorescence decay times of the emission originated from Tb^{3+} and Eu^{3+} were obtained through pulsed excitation using an Opolette HE 355 LD + UVDM (\sim 10 ns, pulse energy $=1.0$ mJ), and the resulting fluorescence signal was analyzed with a Jobin–Yvon Triax 550 monochromator coupled to a Horiba Jobin–Yvon i-Spectrum Two ICCD.

3 Results and discussion

3.1 X‑ray difraction and Raman spectroscopy

Figure [1](#page-2-1)a presents a representative difractogram of the TGZ glasses. Three broad bands located around 2*θ*=22, 27, and

Fig. 1 a Representative difraction pattern of TGZ glasses. **b** Raman spectra of the fabricated glasses, showing their characteristic bands

50 degrees can be distinguished, which are related to shortrange order. The absence of well-defned difraction peaks, within the difractometer resolution, verifes the glassy structure of the sample.

Raman spectra of the fabricated TGZ glasses are displayed in Fig. [1b](#page-2-1), showing the characteristic features of $TeO₂$ -based glasses, which consist of two dominant bands. Band A extends from 300 to 600 cm⁻¹, while band B covers the 600–1000 cm⁻¹ range. Band C around 1550 cm⁻¹, can be observed in the glasses containing europium ions, and its origin is related to the presence of interstitial $O₂$ molecules due to the decomposition of the europium nitrate

hydrate [[26\]](#page-10-8). The band at short wavenumbers is composed of the following vibrational features: Zn–O–Zn and Te–O–Zn bonds (380 cm⁻¹), O–Te–O in TeO₄ units (440 cm⁻¹), and Te–O–Te bridges (490 cm⁻¹) associated with network con-nectivity [[32](#page-10-13), [34](#page-10-15)[–36](#page-10-16)]. Also, around ~450 cm⁻¹, a symmetric stretching vibration of Ge–O–Ge bonds in $GeO₄$ 4-membered rings can be found overlapped with the previously mentioned Te–O and Zn–O bond vibrations [[37](#page-10-17)]. As for band B, three main vibrational features are present, which are the stretching modes of O–Te–O in TeO₄ (660 cm⁻¹), TeO₃₊₁ and $[TeO₃]^{2–} (750 cm⁻¹)$, and short Te–O[–] terminal bonds (820 cm^{-1}) with non-bonding oxygens (NBOs) [\[38](#page-10-18)]. Regarding germanium oxide, the anti-symmetric stretching vibrations of Ge–O–Ge bonds appear around 790 and 870 cm−1, corresponding to the transverse and longitudinal optical split characteristic of $GeO₂$ [[37\]](#page-10-17). It is worth mentioning that the $TeO₂$ vibrational modes dominate the Raman spectra, as it was previously reported for the TeO_2-GeO_2 binary glasses with contents below 20%mol [[28\]](#page-10-9). In addition, there is still no conclusive evidence of the formation of a Te–O–Ge bond that could demonstrate the formation of mixed polyhedral units [[28](#page-10-9)[–30](#page-10-11), [36\]](#page-10-16). However, Kalampou-nias et al. [[28\]](#page-10-9) reported that $GeO₂$ could act as a modifier, increasing the formation of $TeO₃$ units. A more insightful analysis of the glass structure through Raman spectra deconvolution analysis of europium-doped TGZ glasses and the TGZ matrix structural details were previously reported [\[38](#page-10-18)]. However, it was not possible to deconvolute the Raman spectra of TGZ glasses to evidence the bands related to germanium oxide polyhedra because they are completely overlapped with those of $TeO₂$.

3.2 Photoluminescence

The excitation spectrum of europium single-doped glass is presented in Fig. [2](#page-3-0), it was obtained by monitoring the emission at 612 nm, related to ${}^5D_0 \rightarrow {}^7F_2$ transition. It consists of several bands associated to parity-forbidden intra-4f transitions of Eu³⁺: ⁷F₀ \rightarrow ⁵D₄ (365 nm), ⁷F₀ \rightarrow ⁵L₁₀ (367 nm), ${}^{7}F_{0} \rightarrow {}^{5}G_{2,3,4,5}$ (377 nm), ${}^{7}F_{0} \rightarrow {}^{5}L_{8}$ (383 nm), ${}^{7}F_{0} \rightarrow {}^{5}L_{7}$ $(393 \text{ nm}), {}^{7}\text{F}_{1} \rightarrow {}^{5}\text{L}_{7}$ (400 nm), ${}^{7}\text{F}_{0} \rightarrow {}^{5}\text{D}_{3}$, ${}^{5}\text{L}_{6}$ (415 nm),

Fig. 2 Excitation spectrum of TGZE glass ($\lambda_{em} = 612$ nm)

 ${}^{7}F_{3} \rightarrow {}^{5}D_{3}$ (449 nm) ${}^{7}F_{0} \rightarrow {}^{5}D_{2}$ (465 nm), ${}^{7}F_{2} \rightarrow {}^{5}D_{2}$ $(487 \text{ nm}), \, {}^{7}F_{0} \rightarrow {}^{5}D_{1} \text{ (525 nm)}, \, {}^{7}F_{1} \rightarrow {}^{5}D_{1} \text{ (532 nm)},$ ${}^{7}F_{2} \rightarrow {}^{5}D_{1}$ (553 nm), ${}^{7}F_{0} \rightarrow {}^{5}D_{0}$ (578 nm), ${}^{7}F_{1} \rightarrow {}^{5}D_{0}$ (588 nm) [[39](#page-10-19)[–41](#page-10-20)].

The photoluminescence spectrum of TGZE glass, obtained upon excitation at 393 nm (Fig. [3](#page-3-1)), consists of multiple bands assigned to intraconfigurational $4f^6 \rightarrow 4f^6$ transitions of europium ions from 5D_3 , 5D_2 , 5D_1 , and 5D_0 levels to the ${}^{7}F_J$ manifold. The origin of these bands can be explained as follows: during the excitation, electrons are promoted from ${}^{7}F_0$ to ${}^{5}L_6$ level, from where they decay nonradiatively to the ${}^{5}D_3$ emitting level, and subsequently, to

Fig. 3 Emission spectra of TGZE and codoped TGZTE0.5, TGZTE1, TGZTE1.5 and TGZTE2 glasses

the lower ${}^5D_{J=2, 1, 0}$ emitting levels, followed by a radiative decay towards the ${}^{7}F_J$ levels. The same figure presents the emission spectra of TGZTE0.5, TGZTE1, and TGZTE1.5 and TGZTE2 glasses upon the same excitation wavelength.

In all samples, the most intense emission bands are the typical orange-red bands associated to the hypersensitive electric dipole ${}^5D_0 \rightarrow {}^7F_2$ transition, and the magnetic dipole ${}^{5}D_{0} \rightarrow {}^{7}F_{1}$ transition, which is insensitive to the surrounding environment [\[35](#page-10-21), [42,](#page-10-22) [43](#page-10-23)]. It is noticeable that the intensity of the blue-green bands decreases as the content of europium ions increases, while the intensity of the red bands is enhanced $[34, 42, 43]$ $[34, 42, 43]$ $[34, 42, 43]$ $[34, 42, 43]$ $[34, 42, 43]$ $[34, 42, 43]$. This phenomenon is due to an $Eu \rightarrow Eu$ energy transfer, through cross-relaxation processes, given that the energy difference associated to ${}^5D_3 \rightarrow {}^5D_2$ and ${}^5D_2 \rightarrow {}^5D_1$ is resonant with ${}^7F_0 \rightarrow {}^7F_4$ and ${}^7F_1 \rightarrow {}^7F_4$ transitions, according to the following possible mechanisms**:**

(i)
$$
{}^5D_3 + {}^7F_0 \rightarrow {}^5D_2 + {}^7F_4
$$

(ii) ${}^5D_2 + {}^7F_1 \rightarrow {}^5D_1 + {}^7F_4$.

Due to the energy difference between ${}^{5}D_1$ and ${}^{5}D_0$ levels is only 1747 cm⁻¹, a multiphonon relaxation ⁵D₁ → ⁵D₀ occurs, causing that the emissions from ${}^{5}D_1$ level to be negligible compared to those from the ${}^{5}D_0$ level [[34,](#page-10-15) [42,](#page-10-22) [44\]](#page-10-24).

The excitation spectrum of terbium single-doped glass (TGZT) is shown in Fig. [4](#page-3-2). It was obtained by monitoring the emission at 543 nm and consists of multiple bands related to transitions from ${}^{7}F_{6}$ ground state to ${}^{5}G_{2} + {}^{5}L_{7} + {}^{5}L_{8}$ ${}^{5}L_{9} + {}^{5}G_{4}$, ${}^{5}D_{2} + {}^{5}G_{5}$, ${}^{5}L_{10}$, ${}^{5}G_{6} + {}^{5}D_{3}$, and ${}^{5}D_{4}$ excited levels of Tb3+, located at 342, 353, 358, 369, 378 and 485 nm, respectively, being the last one the most intense [[45–](#page-10-25)[48](#page-11-0)].

Fig. 4 Excitation spectrum of TGZT glass obtained by monitoring the emission of Tb^{3+} at 543 nm (blue line), and emission spectra acquired upon excitation of 342 nm (green line)

The emission of the TGZT glass, excited at 342 nm, is displayed in the same fgure. It consists of four low intensity bands related to transitions from the ${}^{5}D_3$ level to ${}^{7}F_6$, ${}^{7}F_5$ ${}^{7}F_{4}$, ${}^{7}F_{3}$ levels, and four bands associated to transitions from ${}^{5}D_4$ level to ${}^{7}F_6$, ${}^{7}F_5$, ${}^{7}F_4$, and ${}^{7}F_3$ levels. The most intense emission band is the one associated to ${}^{5}D_4 \rightarrow {}^{7}F_5$ transition giving origin to the characteristic bright green emission of Tb^{3+} ions.

3.3 Tb3+→ **Eu3+ energy transfer**

The effect of Eu^{3+} content in the codoped glasses was evaluated in a series of samples prepared varying the europium concentration from 0.5 to 2 mol% (TGZTE0.5, TGZTE1, TGZTE1.5, TGZTE2). The emission spectra of this samples, upon excitation of 342 nm, are presented in Fig. [5.](#page-4-0) This wavelength was selected because it does not excite Eu³⁺ ions (see Fig. [2\)](#page-3-0) and it is resonant with ${}^{7}F_6 \rightarrow {}^{5}L_7 + {}^{5}G_2 + {}^{5}L_8$ transitions of terbium ion. The luminescence results in the typical blue-green–red bands of 4f⁸ intra-configurational transitions of Tb³⁺ ion, from 5D_4 to 7F_J manifold, in addition to the characteristic orange-red emissions of Eu^{3+} ions, associated to ${}^5D_0 \rightarrow {}^7F_J$ transitions.

On the other hand, Fig. [6](#page-4-1) depicts the excitation spectra of TGZE, TGZTE0.5, TGZTE1, TGZTE1.5, and TGZTE2 glasses, acquired while monitoring the emission at 700 nm, ${}^{5}D_{0}$ \rightarrow ${}^{5}F_{4}$ transition of europium ions, where no Tb³⁺ emission exists. Two noteworthy features are present in this fgure: first, the excitation bands associated to Eu^{3+} ions are dominant; and second, the presence of additional excitation bands which are related to terbium ions with maxima around 342, 366, 378, and 486 nm. It is worth mentioning

Fig. 6 Excitation spectra of TGZE, TGZTE0.5, TGZTE1, TGZTE1.5, and TGZTE2, monitored at 700 nm

the overlap of the excitation bands of both lanthanide ions at 366 nm and 378 nm; however, the excitation bands around 342 nm and 486 nm only correspond to terbium ion. Also, a decrease in intensity of the terbium excitation bands is observed when the concentration of europium ions increases from 0.5 to 2 mol% in the codoped samples (Fig. [7](#page-4-2)). Considering the spectral overlap between Tb^{3+} emissions $({}_{2}^{5}D_{4} \rightarrow {}_{2}^{7}F_{5} + {}_{5}^{5}D_{4} \rightarrow {}_{7}^{7}F_{4})$ and Eu³⁺ excitation transitions $({}^{7}F_{1} \rightarrow {}^{5}D_{1}, {}^{7}F_{1} \rightarrow {}^{5}D_{0})$ depicted in Fig. [8](#page-5-0), and the results

Fig. 5 Emission spectra of TGZT and codoped glasses excited at 342 nm

Fig. 7 Excitation spectra of TGZT and double-doped TGZ glasses monitoring Tb^{3+} emission at 543 nm

Fig. 8 Spectral overlap of TGZT glass emission spectra, acquired at 342 nm, and excitation spectra of TGZE recorded while monitoring the emission at 612 nm

mentioned above, it is possible to suggest that a $\text{Th}^{3+} \rightarrow \text{Eu}^{3+}$ energy transfer process occurs in the codoped glasses, which could take place through the following three mechanisms, which are illustrated in the partial energy level diagram shown in Fig. [9](#page-5-1) [\[45](#page-10-25), [49](#page-11-1)]:

(1)
$$
{}^{5}D_{3} [Tb^{3+}] + {}^{7}F_{0} [Eu^{3+}] \rightarrow {}^{7}F_{6} [Tb^{3+}] + {}^{5}G_{2,3,4,5} {}^{5}L_{8}
$$

[Eu³⁺],

(2) ${}^{5}D_{3} [Tb^{3+}] + {}^{7}F_{0} [Eu^{3+}] \rightarrow {}^{7}F_{5} [Tb^{3+}] + {}^{5}D_{3} {}^{5}L_{6} [Eu^{3+}]$ and

(3)
$$
{}^{5}D_{4}[Tb^{3+}] + {}^{7}F_{1}[Eu^{3+}] \rightarrow {}^{7}F_{4}[Tb^{3+}] + {}^{5}D_{1}[Eu^{3+}].
$$

Fig. 9 Tb^{3+} and Eu^{3+} energy level diagrams depicting the $Tb^{3+} \rightarrow Eu^{3+}$ energy transfer process

To inquire further about the $\text{Tb}^{3+} \rightarrow \text{Eu}^{3+}$ energy transfer, time decay profiles of terbium ${}^5D_4 \rightarrow {}^7F_5$ transition (543 nm) were measured in the single-doped and codoped glasses. Th^{3+} lifetimes were estimated by ftting the experimental data to an exponential function (Fig. [10a](#page-5-2)). The incorporation of Eu^{3+} in the terbium-doped glasses causes a reduction of the lifetime of the 543 nm emission, and this reduction is greater as the concentration of europium ions increases. This efect suggests a non-radiative energy transfer from Tb^{3+} to Eu³⁺. The efficiency (*η*) and probability (*P*), of $\text{Tb}^{3+} \rightarrow \text{Eu}^{3+}$ energy transfer can be calculated by the following equations, respectively [\[50](#page-11-2)]:

$$
\eta = 1 - \tau_{\text{Tb} \to \text{Eu}} / \tau_{\text{Tb}} \tag{1}
$$

$$
P_{\text{Tb}\to\text{Eu}} = \frac{1}{\tau_{\text{Tb-Eu}}} - \frac{1}{\tau_{\text{Tb}}},\tag{2}
$$

Fig. 10 a Measured decay times profles of the 543 nm emission of Tb³⁺ in the codoped glasses; **b** efficiency of the Tb³⁺ \rightarrow Eu³⁺ energy transfer process as a function of $Eu³⁺$ concentration

Table 2 Lifetimes of Tb^{3+} emission, and calculated values of efficiency and probability of the energy transfer process in TGZT and codoped glasses

Sample	τ (ms)	$\eta(\%)$	$P (s^{-1})$
TGZT	1.086 ± 0.001		
TGZTE0.5	$1.043 + 0.001$	$4.0 + 0.1$	$37.96 + 1.01$
TGZTE1	$0.973 + 0.001$	$10.4 + 0.1$	$106.94 + 1.93$
TGZTE1.5	$0.927 + 0.001$	$14.6 + 0.1$	$157.94 + 2.12$
TGZTE ₂	$0.901 + 0.003$	$17.0 + 0.1$	$189.07 + 3.69$

where τ_{Tb} and $\tau_{\text{Tb-Eu}}$ are the Tb³⁺ lifetimes in the single-doped and codoped TGZ glasses, respectively. The decay times, efficiency, and probability values are enlisted in Table 2 , showing that the efficiency of the process is improved, and its probability is bigger as the concentration of Eu^{3+} increases. The behavior of the efficiency of the energy transfer at different concentrations of Eu^{3+} is shown in Fig. [10b](#page-5-2).

The Inokuti–Hirayama model (I–H model) [\[51](#page-11-3)] is appropriate to describe energy transfer processes where there is not fast energy transfer migration, and can be useful to determine the main mechanism involved in the Tb³⁺ \rightarrow Eu³⁺ non-radiative energy transfer observed in the codoped glass samples. This model was applied to the temporal decay of the terbium emission ${}^{5}D_4 \rightarrow {}^{7}F_5$ at 543 nm, considering multipolar interaction. The non-exponential curve of the temporal decay of the luminescence intensity, *I*(*t*), can be written as follows:

$$
I(t) = I_0 \exp\left[\left(\frac{-t}{\tau_0}\right) - \gamma_S \left(\frac{t}{\tau_0}\right)^{3/S}\right],\tag{3}
$$

where I_0 denotes the intensity at $t=0$, τ_0 is the lifetime value of the donors (Tb³⁺) in the absence of Eu³⁺, γ_s is the energy transfer parameter, while the *S* value can take values of 6 (electric dipole–dipole), 8 (dipole-quadrupole), and 10 (quadrupole–quadrupole), depending on the multipolar interaction between the luminescent centers involved in the energy transfer process. In this way, the temporal behavior of ${}^{5}D_4 \rightarrow {}^{7}F_5$ transition of Tb³⁺ under excitation at 342 nm, was fitted using Eq. [\(3](#page-6-1)) with τ_0 = 1.086 ms, assuming different types of multipolar interactions. For all samples, the best ftting of the experimental data were obtained for $S=6$, which infers that the $\text{Tb}^{3+} \rightarrow \text{Eu}^{3+}$ energy transfer process is dominated mainly by dipole–dipole interactions (Fig. [10a](#page-5-2)). Result that is similar to those observed in other Tb^{3+}/Eu^{3+} codoped glasses [\[11](#page-9-5), [45,](#page-10-25) [52](#page-11-4)]. The estimated energy transfer parameter *γ6* values obtained from the I–H model for TGZTE series of glasses are enlisted in Table [3](#page-6-2). It is observed that the energy transfer parameter varies directly with $Eu³⁺$ concentration.

The critical interaction distance (R_c) for the energy transfer between both lanthanide ions can be calculated using the *γ*⁶ parameter through the following equation [\[53\]](#page-11-5):

$$
R_{\rm c} = \left[\frac{3\gamma_6}{4\pi 1.77\rho_{\rm Eu}}\right]^{1/3},\tag{4}
$$

where ρ_{Eu} is the Eu³⁺ (acceptor) concentration. For the TGZTE glasses the ρ_{En} and R_c values are enlisted in Table [3.](#page-6-2)

The estimated critical interaction distances were found in 6.4–6.7 Å range for all TGZTE samples. Also, considering that ions are evenly distributed in the glass and under the assumption that they are confned in the interior of a sphere with radius *R*, the distance between a pair of randomly distributed ions (D_{random}) can be calculated using Eq. [5](#page-6-3) [[53\]](#page-11-5):

$$
D_{\text{random}} = 2R = 2\left(\frac{3}{4\pi\rho_{\text{T}}}\right)^{1/3},
$$
 (5)

where ρ_T is the total ion concentration (Tb³⁺ plus Eu³⁺). The estimated D_{random} values are between 20 and 18 Å (see Table [3](#page-6-2)). Since the obtained R_c values are considerably shorter than the estimated interaction distance between a randomly distributed pair of ions, it is possible to suggest that the $Tb^{3+} \rightarrow Eu^{3+}$ energy transfer takes place inside clusters, instead of randomly located ions in TGZ glasses.

3.4 Eu3+→**Tb3+ energy transfer**

In the Tb^{3+} excitation spectra shown in Fig. [7](#page-4-2), the presence of two small excitation bands can be distinguished around 393 nm and 465 nm, that are associated to ${}^{7}F_{0} \rightarrow {}^{5}L_{6}$ and ${}^{7}F_{0} \rightarrow {}^{5}D_{2}$ transitions of europium ion, respectively. This fact indicates that the $Eu^{3+} \rightarrow Tb^{3+}$ energy transfer also takes place. To get a better insight about this process, the decay times of the 700 nm emission of Eu^{3+} (${}^{5}D_0 \rightarrow {}^{5}F_4$ transition) in the codoped glasses were compared to the lifetimes of a TGZ series glasses single doped with 0.5, 1.0, 1.5 and 2 mol% of Eu^{3+} : 1.119 ms (TGZE0.5), 1.084 ms (TGZE1.0), 1.054 ms (TGZE1.5) and 1.056 ms (TGZE2.0). According to the estimated results, the $Eu^{3+} \rightarrow Tb^{3+}$ energy transfer in the codoped glasses occurs

Table 3 Multipolar interaction parameter (γ_6) , Eu³⁺ concentration (ρ_{Eu}) , critical interaction distance (R_c) , and distance between two randomly distributed ions (D_{random})

Sample	Y6	ρ_{Eu} (ions/cm ³)	$R_c(\AA)$	$D_{\text{random}}(A)$
TGZTE0.5	0.0567	2.74×10^{19}	6.6	22
TGZTE1	0.1312	5.45×10^{19}	6.7	21
TGZTE1.5	0.1776	8.13×10^{19}	6.5	18
TGZTE ₂	0.2294	1.078×10^{20}	6.4	17

less efficiently, which is demonstrated by small variation of Eu^{3+} decay times, as well as their calculated efficiency and probability (Table [4\)](#page-7-0). As can be noted, the presence of Tb^{3+} in the codoped glasses reduces the lifetime of the 700 nm emission, compared to the value obtained in the TGZE samples. However, this reduction may also be due to a cross-relaxation process between the europium ions, as described above.

Table 4 Decay times, efficiency, and probability values of $Eu^{3+} \rightarrow Tb^{3+}$ energy transfer in the double-doped glasses

$P (s^{-1})$
41.80
20.00
1.80
0.89

3.5 Photometric analysis

The tellurite glasses studied in this paper can fnd application as multicolor phosphor materials, given its wide range of tunability, from green to orange-red, by modifying the lanthanide concentration ratio or the excitation wavelength. To evaluate this tuning property, the fabricated glasses were excited in the UV range using diferent wavelengths and the chromaticity coordinates of their emission were estimated and plotted in the CIE 1931 chromaticity diagram (Fig. [11](#page-7-1)). Additionally, the correlated color temperature was calculated by means of McCamy equation [[54\]](#page-11-6):

$$
CCT = -449n3 + 3525n2 - 6823.8n + 5520.33,
$$
 (6)

where $n = (x - 0.3320)/(y - 0.1858)$, being *x* and *y* the chromaticity coordinates. The global emission of the glasses can be modulated from green to red region, passing through the yellow and orange regions of the spectrum, by simply modifying the Eu^{3+}/Tb^{3+} ratio or the excitation wavelength

Fig. 11 CIE 1931 chromaticity coordinates of the global emission of the fabricated glasses, excited at diferent UV wavelengths using a xenon arc lamp, **a** 342 nm, **b** 366 nm, **c** 378 nm, and d) 393 nm

(342, 366, 378, and 393 nm). The calculated chromaticity coordinates, CCT and color purity (CP) values are summarized in Table [5](#page-8-0).

When the codoped samples are excited upon 342 nm (where europium ions do not have excitation band), the corresponding emission arises from the bands associated with both Tb³⁺ and Eu³⁺ ions, because the Tb³⁺ \rightarrow Eu³⁺ energy transfer process. Besides, the increase of europium concentration and the consequent decrease of terbium emissions causes a shift of the global emission tonality towards the orange region (Fig. [11](#page-7-1)a). In the case of the values shown in Fig. [11](#page-7-1)b, c, the tunability of the global emission arises from the direct excitation of both ions and the energy transfer process. Figure [12](#page-8-1) shows photographs of the glasses under UV-LED light excitation of 365 nm, evidencing the emission color change of the single and codoped glasses from green to reddish-orange. Excitation wavelengths of 366 nm and 378 nm are resonant with ${}^{7}F_6 \rightarrow {}^{5}L_{10}$, ${}^{5}G_6$, ${}^{5}D_3$ transitions of Tb³⁺ and ${}^{7}F_0 \rightarrow {}^{5}$ ${}^{7}F_{6} \rightarrow {}^{5}L_{10}$, ${}^{5}G_{6}$, ${}^{5}D_{3}$ transitions of Tb³⁺ and ${}^{7}F_{0} \rightarrow {}^{5}D_{4}$,
 ${}^{7}F_{0} \rightarrow {}^{5}L_{10}$, ${}^{7}F_{0} \rightarrow {}^{5}G_{2,3,4,5}$ transitions of Eu³⁺. Then, the origin of europium emission bands is due, frst, to direct excitation of the Eu³⁺ and second to the Tb³⁺ \rightarrow Eu³⁺ energy transfer process. Finally, when the glasses are excited at 393 nm, the double-doped glasses emission tonality comes fundamentally from europium ions in the reddish region, as shown in Fig. [11](#page-7-1)d.

Table 5 CIE 1931 chromaticity coordinates, CCT and CP values of the single and codoped TGZ glasses at diferent excitation wavelengths

Under daylight

Under UV-light excitation

4 Conclusions

Tb³⁺, Eu³⁺ single-doped and Tb^{3+}/Eu^{3+} codoped TeO₂–GeO₂–ZnO glasses were successfully prepared by melt quenching technique. The luminescent properties investigated upon UV excitation allowed to identify the presence of a Tb³⁺ \rightarrow Eu³⁺ energy transfer process in the codoped samples. In such samples, the emission from terbium ${}^{5}D_{4}$ level decays faster than in TGZT (Tb^{3+}) -doped glass), revealing that Tb^{3+} can sensitize Eu^{3+} through non-radiative energy transfer. The fitting of the Tb^{3+} luminescence decay profle using the Inokuti-Hirayama model, allow us to infer that the Tb³⁺ \rightarrow Eu³⁺ energy transfer process can be dominated by an electric dipole–dipole interaction in all codoped glasses. Furthermore, an increase of the energy transfer efficiency was observed, coupled with an enhancement of the $Eu³⁺$ emission at the expense of reducing the characteristic emission of Tb^{3+} , as the Eu³⁺ concentration increases. The photometric characterization of Tb^{3+} –Eu³⁺-doped glasses was carried out at diferent excitation wavelengths by evaluating the CIE 1931 chromaticity coordinates, showing a multicolor tuning of the emission from green to the orangered region as a function of europium concentration with high color purity. Thus, the glasses under study can be good candidates as multicolor phosphors for their emission color tuning capability for solid-state devices.

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Data availability The datasets generated during and/or analyzed during the current study are available from the corresponding author on reasonable request.

Declarations

Conflict of interest The authors declare there is no conficts of interest.

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