

# Phase transition and changing properties of nanostructured V<sub>2</sub>O<sub>5</sub> thin **flms deposited by spray pyrolysis technique, as a function of tungsten dopant**

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## **Abstract**

To investigate their formation mechanism and basic electronic behaviors, in this study undoped and tungsten (W)-doped vanadium oxide flms were comparatively analyzed by X-ray difraction (XRD), scanning electron microscopy (SEM), atomic force microscopy (AFM), Raman spectroscopy, photoluminescence (PL) and UV–Vis spectroscopy. The XRD measurement showed that the W doping causes a change in crystalline phase, improvement of grain size, a decrease in dislocation and stress with diferent dopant ratios. The SEM and AFM images evidenced that the W-doped flms change nanostructured form from nanorods to composition of nanorods and nanoplates at 400 °C, which is 20% lower than general formation temperature of vanadium pentoxide  $(V_2O_5)$  nanorods, while the undoped film have small-sized nanostructural morphology as an initial stage. The Raman spectrum also indicates signifcant changes of its structural and optical properties. The PL quenching is observed in the PL spectra of the W-doped  $V_2O_5$  films with decreasing deep-level emission. Band gap, Urbach energy, refective index and dielectric constant of all the flms were calculated by data of the UV–Vis and refectance measurements. The results showed that the W-doped  $V_2O_5$  films were more suitable for (opto)electronic device applications such as solar cell, gas sensor, electrochromic electrode and battery.

## **Graphic abstract**



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# **1 Introduction**

Recently, the vanadium oxides as a class of semiconducting oxides, VO,  $V_2O_3$ , VO<sub>2</sub> and  $V_2O_5$  including in the  $V^{2+}$ ,  $V^{3+}$ ,  $V^{4+}$  and  $V^{5+}$  oxidation states, are remarkable due to their typical properties for various technological applications. Foremost among them, the  $V_2O_5$  as a transition-metal oxide

n-type semiconductor is one of the remarkable candidates due to wide band gap  $(2.4–2.8 \text{ eV})$ , high absorption coefficients, thermal stability, multivalence layered structure, abundance, low cost and electrochromic and photocatalytic feature [[1](#page-15-0)[–3](#page-16-0)]. These properties allow this material's entry into the (opto) electronic applications such as solar cells [[4\]](#page-16-1), light emitting diodes [\[5](#page-16-2)], photodiodes [\[6](#page-16-3)], transistors [[7\]](#page-16-4), gas sensors [\[8](#page-16-5)], electrochromic devices [\[9](#page-16-6)] and solid-state batteries [[10\]](#page-16-7).

However, some disadvantages of the  $V_2O_5$  has been reported in the previously studies such as low conductivity, which is most important one of them, for all applications and weak retention capacity in the long-term cycling conditions for storage applications [\[11,](#page-16-8) [12\]](#page-16-9). On the other hand, orthorhombic structure of the  $V_2O_5$  ( $\alpha$ - $V_2O_5$ ) has been generally reported due to most stable form and its suitable properties for the applications [[13,](#page-16-10) [14\]](#page-16-11). However, the activation energy of conductivity is 0.24 and 0.1 eV at near room temperature for α- and  $β$ -V<sub>2</sub>O<sub>5</sub> (tetragonal structure), respec-tively [\[15,](#page-16-12) [16](#page-16-13)]. Therefore, the β-V<sub>2</sub>O<sub>5</sub> is more suitable for conductivity while  $\alpha$ -V<sub>2</sub>O<sub>5</sub> is useful for long-term working conditions. Thus, the mixed phase of  $V_2O_5$  may be solution for the above-mentioned problems.

In addition to the issues of the applications, as reported by many authors,  $V_2O_5$  thin films have generally amorphous structure when deposited at low growth temperatures  $(\leq 300 \degree C)$  [\[17–](#page-16-14)[20\]](#page-16-15). Furthermore, it is emphasized that formation temperature (growth, annealing and post-annealing) of nanostructured  $V_2O_5$  is 400 °C for starting and to be completed about 550 °C [\[21\]](#page-16-16). It is well known that low temperature processes using fabrication of (opto)electronic devices is the desired condition. On the other hand, due to changeable transfer temperature and electrical and optical properties, vanadium oxide has been doped with various dopants such as tungsten (W), molybdenum (Mo), niobium (Nb), fuorine (F), iron (Fe) and chromium (Cr) [\[22–](#page-16-17)[27](#page-16-18)]. For example, Zheng et al. [\[28](#page-16-19)] have reported W-doped  $V_2O_5$ nanobelts have higher specifc capacitance than the previous values of  $V_2O_5$  materials for due to excellent electrochemical property, high specifc capacitance and good rate capability. Chen et al. [[29\]](#page-16-20) have used the Mo to investigation of efect on phase transition temperature of  $VO<sub>2</sub>$  nanopowders and showed decreasing phase transition temperature in the Mo doped VO<sub>2</sub>. On the other hand, F doped  $V_2O_5$  films having tetragonal β-V<sub>2</sub>O<sub>5</sub> phase have been reported by Mousavi et al. [\[25\]](#page-16-21). In this study, while the size of nanostructures decreased with increasing F doping level, the width of nanostructure increased. Also, resistivity, transparency and optical band gap of the flms decreased in the F doped flms. Finally, impurity cations into vanadium oxide lattice can change structural, morphological, electrical and optical properties of the vanadium oxide such as crystalline phase, shape of nanostructure, conductivity, chemical activity and color control of electrochromic electrode.

The nanostructured  $V_2O_5$  thin films can be obtained by several growth techniques such as thermal evaporation [\[30](#page-16-22)], pulsed laser deposition (PLD) [[31](#page-16-23)], magnetron sputtering [[32\]](#page-16-24), sol–gel process [[33\]](#page-16-25), chemical vapor deposition [\[34](#page-16-26)], electron beam evaporation [[35](#page-16-27)], and spray pyrolysis [[36](#page-16-28)]. The chemical spray pyrolysis (CSP) is an advantageous method for preparation of thin flms coated on large area due to being simple and inexpensive process. Furthermore, the CSP deposition technique offers crucial advantages such as control of stoichiometry and thin flm structure [\[37](#page-16-29)].

In this study, undoped (reference sample) and W-doped vanadium oxide thin flms deposited by the CSP at same preparation conditions was investigated. Structural, morphological and optical properties of the thin flms were determined and comparatively discussed.

### **2 Experimental details and measurements**

Microscope glass slides as glass substrate were purchased from Isolab/Germany. Before cleaning and flm growth, the glass slides were cut to smaller pieces of  $1.5 \text{ cm} \times 1.5 \text{ cm}$ . Thickness of the glass substrates is 0.4 cm and the substrates were used without surface polishing. The EDX analysis was conducted to determine of composition of the glass substrates and is given in Fig. [5a](#page-9-0). The chemical salts of vanadium trichloride (VCl<sub>3</sub>) (97% purity) and tungsten hexachloride (WCl<sub>6</sub>) (99.9% purity) were purchased from ACROS Organics and abcr (Gute Chemie), respectively. All the depositions were carried out by compressed purifed air of 2 bar (1.97 atm.) pressure as the carrier gas with 5 ml/min solution flow at vertical position of spray-nozzle. Purified air was used as the carrier gas and the distance between spraynozzle and glass substrate was set 30 cm in all the deposition process. The chemical solvents used were purchased from Sigma-Aldrich and used without post-treatment.

The glass substrates were cleaned as follows: Firstly, substrates in acetone were waited in ultrasonic bath for 15 min and washed with deionized water and then dried with  $N_2$  gas flow. Secondly, they were chemically cleaned in 10% NaOH solution pre-heated up to 55 °C and finally, substrates in deionized water were hold in ultrasonic bath for 15 min and then dried with  $N_2$  gas flow, respectively.

The undoped  $V_2O_5$  thin films were deposited by precursor solutions prepared with 0.05 M (393 mg) VCl<sub>3</sub> in 50 ml deionized water. On the other hand, the doped  $V_2O_5$  thin flms were deposited by the precursor solution including percentage weight ratios (wt%) of 1, 3, 5 and 10 WCl<sub>6</sub>, 4 mg (0.31 mM), 12 mg (0.63 mM), 20 mg (1.57 mM) and 40 mg (3.14 mM), respectively. After the precursor solution was stirred for 60 min at room temperature (RT), the thin flms were deposited on the glass substrates by spray pyrolysis technique at 300 °C for 20 min and naturally cooled to RT.

After then, the undoped and the W-doped  $V_2O_5$  thin films were post-annealed in air at 400 °C for 60 min without heating and cooling step. A photograph of the HOMARC Spray Pyrolysis Equipment and growth control parameters are given in Fig. [1](#page-2-0). After the post-annealing process, the  $V_2O_5$ thin flms were characterized.

The crystallinity of the flms was checked using Bruker D2 PHASER X-ray difraction (XRD) system with CuK*<sup>α</sup>* radiation at  $\lambda = 1.54184$  Å. The images of field emission scanning electron microscopy (FE-SEM) including crosssectional and atomic force microscopy (AFM) to reveal morphological properties were obtained using a system of ZEISS Sigma 300 and Hitachi 5100 N, respectively. Energy dispersive X-ray (EDX) analysis for compositional analysis was conducted with measurement system of ZEISS Sigma 300 equipped with EDX spectroscopy. Raman spectra were recorded using a μ-Raman/PL spectrometer of WITEC Alfa 300R. The excitation source is a He–Ne laser with an excitation wavelength of 531.96 nm. The photoluminescence (PL) spectra were obtained using Horiba in the wavelength range of 350–900 nm. A Perkin-Elmer Lambda2S UV–visible spectrometer was used in order to obtain absorption and calculate the band gaps of the flms. The flm thickness was determined using a surface proflometer (KLA Tencor P7).

## **3 Results and discussion**

#### **3.1 Structural analysis**

Figure [2](#page-3-0) shows the XRD patterns of the undoped flm (a) and the W-doped flms depending on dopant concentration (b–e), respectively. Figure [2](#page-3-0)a shows that the undoped flm as reference sample has mixture crystal structure of monoclinic VO<sub>2</sub>, orthorhombic  $V_2O_5$  and tetragonal  $V_2O_5$ . Besides, high preferential orientations of (200) and (101) are stronger than that of other peaks indicating that the crystallinity



<span id="page-2-0"></span>

of undoped sample is dominantly formed  $\alpha$ - and β-V<sub>2</sub>O<sub>5</sub> phases. It exhibits characteristic 2*θ* values of 12.55, 14.24, 26.79, 31.42 and 42.64 with respective Miller indices (200), (002), (330), (242) and (003), which are consistent with crystalline the  $\beta$ -V<sub>2</sub>O<sub>5</sub> phase and are in good agreement with the JCPDS card No. 045-1074 (crystal system: Tetragonal, space group: PE). The appearance of (101), (201) and (012) crystal plane at 2*θ* values of 20.82, 25.25 and 48.81 illustrates the formation of the  $\alpha$ -V<sub>2</sub>O<sub>5</sub> crystals given in JCPDS (089-2482) (crystal system: Orthorhombic, space group: Pmmn (59)). On the other hand, weak difraction peaks are also seen at 23.80, 25.25, 31.19 and 48.81, which matched with the (201), (110), (111) and (402) planes of the  $VO_2$ given in JCPDS (071-0042) (crystal system: Monoclinic, space group: C2/m (12)). Besides, the XRD patterns clearly indicate the broad background at 2*θ* between 17° and 40° due to the amorphous feature of glass substrate.

Interestingly, the other XRD patterns show that the W-doped thin flms have signifcant structural changes. It can be seen in Fig. [2](#page-3-0)b that the XRD pattern of the thin flm doped at 1 wt% have new difraction peaks with decreased intensity. The weak XRD intensities of the crystallite refected plane directions are due to the nano-crystallinity [[38](#page-16-30)]. However, no reflections attributable to compounds deriving from the dopant elements were observed. It suggests that precursor solution of the  $V_2O_5$  with homogeneously dispersed the W dopant is successfully prepared. For the peak of preferential orientation along (200), the full-width half-maximum (FWHM) value is dramatically decreased from 0.345 to 0.130. This structural change can be attributed that the crystallinity and crystalline structures in the flm is enhanced, leading to the higher lattice order. However, the existence of three crystal phases is not afected and the  $\beta$ -V<sub>2</sub>O<sub>5</sub> phase has become prominent. Furthermore, the increasing dissimilar orientations of phases in the XRD pattern can be attributed to that the dopant induced the change in crystal growth and then reorientation of crystallinity plane. The secondary dominant  $\alpha$ -V<sub>2</sub>O<sub>5</sub> peak disappeared after the doping process. Besides, the predominant peak center signifcantly shifted towards lower difraction angle, from 12.549° to 12.511°.

After doping at 3 wt% (Fig. [2](#page-3-0)c), the (200) β-phase peak still dominated in the XRD spectrum with decreased intensity, indicating that the degree of (200) preferential orientation decreased with increasing doping concentration due to substitution of W ions in  $V_2O_5$  lattice making O ions move to interstitial sites and resulting in a less ordered crystalline phase. In addition to this result, the (400), (241) and (233) β-phase peaks are more sharped with higher intensity, while the intensities of  $\alpha$ -phase and  $VO<sub>2</sub>$  diffraction peaks are getting weaker. The formation of the increasing  $β-V_2O_5$  phase at 3 wt% can be due to **Fig. 1** A photograph of spray pyrolysis system used in this study the improving chemical reactions for various crystalline

<span id="page-3-0"></span>**Fig. 2** XRD patterns of the pure (**a**) and doped thin flms with

the W rate of 1 wt% (**b**), 3 wt%

(**c**), 5 wt% (**d**) and 10 wt% (**e**)



**Fig. 2** (continued)



orientations in the annealing process. It can be evident that the doping at 3 wt% afected crystalline orientation rather than quality because the FWHM value of preferential orientation along (200) was nearly same value of 0.131 while peak center was shifted towards lower difraction angle, from 12.511° to 12.499°. In the theory, original XRD peak positions and peak intensities of a doped structure with other atoms change due to the repeat distances in the lattice structure arising from the diferent size between host and doped atoms and diference in electron density of the doped atom, respectively [[39](#page-16-31)]. Therefore, the changing

intensity, angle and orientation of difraction peaks evidence that our dopant incorporated successfully.

Furthermore, for the flm prepared with 5 wt% W including precursor solution, the XRD pattern in Fig. [2d](#page-3-0) demonstrates a more decreased intensity of (200) peak and shifted 2*θ* position towards 12.463°. Specifcally, the XRD data suggests that the diffraction peaks of monoclinic  $VO<sub>2</sub>$  secondary phase nearly disappeared, except its (201) plane. Besides, the number of difraction planes decreased while intensity generally increased in the XRD pattern. The appeared dominant  $V_2O_5$ phase with increasing dopant level can be due to enhanced

oxygen chemisorption, which is considered as flled oxygen vacancies in the lattice by oxygen through gas phase oxidation. On the other hand, Fig. [2](#page-3-0)e presents the XRD pattern of the flm obtained from 10 wt% W-loaded solution. As seen in Fig. [2](#page-3-0)e, the monoclinic  $VO<sub>2</sub>$  secondary phase fully disappeared and only  $V_2O_5$  phase reflection peaks were observed. This change may be attributed to that tensile stress will be helpful to the phase transition occurring at the lower temperature because its volume expands by 1% of its original volume when vanadium dioxide changes from a semiconducting state to a metal state [\[40](#page-17-0)]. In addition, shifting degree of the  $2\theta$  position at (002) is bigger than the other, as seen in Table [1](#page-5-0). In this regard, the W dopant induced the change in crystal phase, crystalline quality and reorientation planes are clearly evident from their XRD patterns. In addition to the relevant phase changes mainly occurred, microstructural parameters like crystallite size, strain and dislocation density positively were afected. The parameters were calculated from the (200) XRD peak of  $β$ -V<sub>2</sub>O<sub>5</sub> phase and given in Table [1.](#page-5-0)

The grain size along the preferred orientation of  $\beta$ -V<sub>2</sub>O<sub>5</sub> peak was calculated by using the Scherrer's equation:

$$
D = \frac{\text{K}\lambda}{\beta \cos \theta} \tag{1}
$$

where  $D$  is the crystallite size,  $K$  is a dimensionless shape factor and has a typical value of about 0.9, *λ* is the wavelength of the incident radiation (1.5406 Å), *β* is the fullwidth half-maximum (FWHM) and *θ* is the Bragg's angle of difraction. The crystallite size of W-doped flms was found to be  $\geq$  54 nm. The best size is obtained with 1 wt% and 3 wt%, whereas the reference sample has 23.15 nm size. This result indicates that the doped flms are highly crystalline. The *d*-spacing of W-doped  $V_2O_5$  can be calculated by the Bragg formula:  $2d \sin \theta = n\lambda$ , where *d* is the interlayer distance. The increment of interlayer distance of the W-doped thin flms was observed to be blue shifted (200) plane. Additionally, in order to have more information, the dislocation density  $(\delta)$  and strain value were calculated by the formulas, respectively,

$$
\delta = \frac{1}{D^2},\tag{2}
$$

$$
\varepsilon = \frac{\beta \cos \theta}{4}.\tag{3}
$$

It can be seen in Table [1](#page-5-0) that the dislocation density decreases with the W dopant which may be due to decrease in concentration of lattice defects such as oxygen vacancy. However, the dislocation density increased after 3 wt% dopant. Therefore, the value of 3 wt% is critical for dislocations. This result showed that the W dopant plays a constructive role in crystallinity of the flm until the critical value 3 wt%. However, the higher doping rate than the critical value may cause lattice defect due to diference of atomic parameters such as atomic radius. On the other hand, due to the increase in *d*-spacing (*d*: 7.130 Å given in JCPDS (045- 1074)), the strain will be decreased in the structure.

#### **3.2 FE‑SEM and AFM analysis**

Figures [3](#page-6-0) and [4](#page-7-0) show SEM images of the undoped and the W-doped  $V_2O_5$  films prepared at the 1, 3, 5, 10 wt% on the glass substrate and cross-sectional images of the flms, respectively. The average flm thickness of 550 nm was determined by the proflometer. Also, the thicknesses determined from the cross-sectional images are 518, 569, 589, 558 and 564 nm, respectively. As seen in Fig. [3a](#page-6-0), the undoped  $V_2O_5$  film has grown as small seed-nanorod form. However, it can be seen in Fig. [3b](#page-6-0) that the 1 wt% W-doped  $V_2O_5$  obtained by the same growth conditions with the undoped flm have transformed into nearly uniform nanorods and interestingly diameter, length and density of the nanorods have changed by increasing W-doping concentra-tion as seen in Fig. [3c](#page-6-0)–e. The diameter $_{(min-max)}$ /length $_{(min-max)}$ of the nanorods were determined by the FE-SEM system (not given in here) as  $60<sub>min</sub> - 200<sub>max</sub>$  nm/150<sub>min</sub>-430<sub>max</sub> nm,  $90_{\text{min}} - 520_{\text{max}}$  nm/320<sub>min</sub>–2060<sub>max</sub> nm,  $150_{\text{min}} - 630_{\text{max}}$ nm/350<sub>min</sub>–2000<sub>max</sub> nm, 120<sub>min</sub>–720<sub>max</sub> nm/320<sub>min</sub>–1800<sub>max</sub> nm and  $120<sub>min</sub> - 720<sub>max</sub>$  nm/260<sub>min</sub>–1140<sub>max</sub> nm, respectively. This alteration suggests that the form of the W-doped flms still keeps the same shape as the undoped flm while the nanostructure growth is significantly enhanced by the dopant. This enhanced nanostructural morphology provide an advantage as a more suitable sensitive area and bigger

<span id="page-5-0"></span>**Table 1** Determined and calculated parameters from the XRD patterns of all the flms

	Peak position $(2\theta)$	Phases	$d$ -spacing (A)	FWHM	Grain size (nm)	Dislocation den- sity $(\times 10^{-3} \text{ nm}^{-2})$ (line <sup>-2</sup> m <sup>-4</sup> )	Strain
VO.	12.549	$VO_2$ , $\alpha$ - $V_2O_5$ , $\beta$ - $V_2O_5$	7.051	0.345	23.15	1.865	0.085
VO-1\%	12.511	$VO_2$ , $\alpha$ - $V_2O_5$ , $\beta$ - $V_2O_5$	7.072	0.130	61.48	0.265	0.032
$VO-3\%$	12.499	VO <sub>2(reduced)</sub> , $\alpha$ -V <sub>2</sub> O <sub>5</sub> , $\beta$ -V <sub>2</sub> O <sub>5</sub>	7.079	0.131	61.01	0.268	0.032
$VO-5\%$	12.463	VO <sub>2(reduced)</sub> , $\alpha$ -V <sub>2</sub> O <sub>5</sub> , $\beta$ -V <sub>2</sub> O <sub>5</sub>	7.099	0.144	55.50	0.324	0.033
VO-10 $%$	12.452	$\alpha$ -V <sub>2</sub> O <sub>5</sub> , $\beta$ -V <sub>2</sub> O <sub>5</sub>	7.106	0.148	54.01	0.342	0.034



<span id="page-6-0"></span>**Fig. 3** SEM images of the pure (a) and 1, 3, 5 and 10 wt% W-doped  $V_2O_5$  thin films (**b–e**), respectively

interlayer space for device applications such as an efective active layer for gas sensors [\[41](#page-17-1)], a hole transport layer for solar cells [\[42](#page-17-2)], and a cathode material for lithium batteries [[42,](#page-17-2) [43\]](#page-17-3). In addition, it is frequently reported that flm oxidized into  $VO_2$  phase at 300 °C,  $V_6O_{13}$  phase at 350 °C, finally  $V_2O_5$  at 400 °C with starting to grow nanorods structures, and transformed into  $V_2O_5$  nanorods in the higher temperatures [[44](#page-17-4), [45\]](#page-17-5). Therefore, minimum required





<span id="page-7-0"></span>**Fig. 4** Cross-sectional SEM images of the pure (a) and 1, 3, 5 and 10 wt% W-doped V<sub>2</sub>O<sub>5</sub> thin films (**b–e**), respectively

temperature is decreased with the W doping from  $\geq 500$  to 400 °C in this study.

On the other hand, the nanostructural morphology is gradually changed from rod to plate with the increasing dopant amount and then a composition of nanorod- and

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nanoplate-shaped forms occur. Hence, the corresponding approximate ratio of the diameter $_{\text{min}}$ –length $_{\text{max}}$  is undulated as 0.139, 0.044, 0.075, 0.066 and 0.105, respectively. As clearly seen in Figs. [3](#page-6-0) and [4](#page-7-0), in the 3 wt% W-doped flm, the nanorods having high size combine with each other to

form nanoplate, and thus, a decrease occurs from 0.075 to 0.066 in the ratio of diameter $_{\text{min}}$  and length $_{\text{max}}$ . Furthermore, this undulation in the values confrms that as the nanorods combine with each other to form nanoplates, the density of nanostructures is decreased with the increasing W concentration. It may be noticed that the W-doping process strongly infuenced crystallization of the flm. This result also suggests that the surface difusion phenomenon which plays an important role in the growth process of nanorods due to the high dependence on the oxygen absorption is enhanced with the W dopant. It may cause that electron affinity of the W metal atoms (electronegativity: 2.36, electron afnity: 78.6 kJ/mol) is bigger than that of V atoms (electronegativity: 1.63, electron affinity: 50.6 kJ/mol) and then the increased nanorod-size can be obtained with enhanced absorption of  $O<sub>2</sub>$ . In other words, all the specific physical and chemical properties of materials are associated with their atomic structures. The atomic structures of materials are related to atomic parameters such as lattice structure, lattice parameters, radius, growing shape and electron afnity of atoms. Additionally, the above-mentioned properties of materials are strongly based on metal concentration and type of metal in materials [\[46](#page-17-6)]. Figure [5](#page-9-0)b–f shows the EDX spectra of the undoped and W-doped flms, respectively. The EDX analysis carried out to verify the W incorporation shows the presence of the three main elements V, O and W. According to the chemical composition of the doped flms, the percentages of the W are  $0.56$  wt%,  $2.63$  wt%,  $3.78$  wt% and 5.86 wt%, respectively. Although the EDX analysis is used for elemental analyses only, the percentages of the W show decreasing doping efficiency with increasing doping ratio. On the other hand, atomic percentage values of the elements in the undoped and doped films are  $45.51$  and  $54.49$ ; 33.22, 66.67 and 0.11; 32.55, 66.92 and 0.53; 38.78, 60.54 and 0.68; 30.09, 68.83 and 1.08, respectively, for V, O and W.

The surface topology of the nanostructured flms prepared with and without the dopant was analyzed using AFM in the tapping mode. Figure [6](#page-10-0) shows 3D AFM images of the  $V_2O_5$  nanostructured films. It was observed that the SEM and AFM analyses were consistent with the analysis of the nanostructured crystal distributions on the surfaces of the flms. The surface topography images show that the surface morphology and the nanorods size are very sensitive to the dopant ratio, as seen in SEM images. Root mean square (RMS) roughness is between 100 and 400 nm. While the roughness increased with the addition of W dopant due to increased nanorod-size, surface roughness in the flms doped greater than 1 wt% W decreased due to the decreasing angle between the substrate and nanorod, which can be useful for the fabrication of  $V_2O_5$ -based (opto)-electronic devices such as solar cell, photosensor, light emitting diode and gas sensor.

It is well known that performance of these devices depend on interface area. However, carrier difusion length limits thickness of organic or inorganic semiconductor top-layer in fabrication of the devices. On the other hand, a suitable thickness should be in order to completely cover the underlayer to prevent short circuit. Therefore, increasing surface area with decreasing roughness allows a wide range in the selection of top-layer thickness.

#### **3.3 Raman analysis**

Raman spectrum analysis of undoped and W-doped  $V_2O_5$  for all concentrations of the dopant ion is presented in Fig. [7.](#page-11-0) In addition to the other analyses, these spectra suggested that the films have orthorhombic structure of  $V_2O_5$ . Also, applying the W dopant infuenced the crystallinity of the  $V<sub>2</sub>O<sub>5</sub>$  nanorods. The Raman spectrum consists of two groups of peaks, the peaks located at higher frequency region are called internal modes and the ones in low-frequency region are called external modes [\[47\]](#page-17-7). The Raman peaks located at 140.75, 201.62, 260.75, 279.92, 403.54, 481.59, 506.54, 692.81, and 996.67 cm<sup>-1</sup> were indexed to the V<sub>2</sub>O<sub>5</sub> crystal. On the other hand, the signal peak at  $260.75 \text{ cm}^{-1}$  is assigned to the  $VO<sub>2</sub>$  phase with vibrational modes of V–O bonds [\[48](#page-17-8)].

The predominant low-wavenumber band at  $140.75 \text{ cm}^{-1}$ is attributed to the skeleton bent vibration related to B3g symmetry vibration that is an evidence for the layered structure of  $V_2O_5$  phase and nanorods [\[49](#page-17-9), [50\]](#page-17-10), and the peaks at 201.62 and 279.92 cm<sup>-1</sup> are from the bending vibrations of the  $O_C-V-O_B$  bond related to B2g [[51](#page-17-11)]. The peak at 403.54  $\text{cm}^{-1}$  is assigned to the bending vibrations of the V=O bonds related to Ag symmetry vibration. The bending vibration signal of V–O–V doubly coordinated oxygen is observed at 481.59 cm<sup>-1</sup>. The peak at 506.54 cm<sup>-1</sup> is assigned to the triply coordinated oxygen  $(V_3$ –O) stretching modes related to Ag symmetry vibration. The peak at 692.81 cm<sup>-1</sup> is assigned to the doubly coordinated oxygen  $(V<sub>2</sub>–O)$  stretching and bending vibrational modes related to B2g symmetry vibration [\[52](#page-17-12)]. The peak at 996.67 cm<sup>-1</sup> corresponds to the terminal oxygen  $(V<sub>5</sub>-O)$  stretching mode related to Ag symmetry vibration, and also associated to the crystallinity quality and stoichiometry of the layered structure of  $V_2O_5$  nanorods [\[53\]](#page-17-13). In the case of the Raman peaks at 523 and 696 cm<sup>-1</sup>, they are attributed to V<sub>3</sub>–O phonon band and V–O–V modes. Intensity of both these modes was found to be quite small [\[54](#page-17-14)].

Clearly, the Raman analysis is in accord with the XRD patterns. As seen the Raman spectra, in the 5 wt% W-doped thin flm, the weaker mode intensity and broader stretching vibration peaks beyond 400 cm−1 have been observed. It can be attributed that while the mode intensity increases with increasing dominant  $V_2O_5$  diffraction peaks, the dominant



<span id="page-9-0"></span>**Fig. 5** EDS spectra of the glass substrate (**a**), undoped thin flm (**b**) and W-doped thin flms (**c**–**f**), respectively

diffraction peaks of the  $VO<sub>2</sub>$  caused the broad spectrum with weaker intensity.

# **3.4 Optical studies**

Figure [8](#page-11-1) presents the PL spectra recorded at room temperature of the undoped and the W-doped  $V_2O_5$  thin films. In the measurement, an excitation source with wavelength of 325 nm was used for all samples. All the samples show two luminescence mechanisms which are intrinsic and extrinsic transition centralized at the green and the red region of the electromagnetic spectrum, respectively. The broad intrinsic transition peaks around 420 and 570 nm are attributed to recombination of photo-excited pairs that are electron in bottom of conduction band as in lowest split-off V-3d and hole in top of O-2p, respectively [[55](#page-17-15)].



<span id="page-10-0"></span>**Fig. 6** AFM images of the undoped and W-doped  $V_2O_5$  thin films, respectively, and graph of root mean square (RMS) versus the doping rate

The weak peak at 705 nm originates from formation of oxygen vacancies because defects of oxygen vacancy act as recombination center for photo-generated electron. In addition, lower wavelength peaks at 686 nm is derived from the transition from donor defect band to the 0–2p valence band. As seen in Fig. [8](#page-11-1), for band gap values, infrared shift at the near band edge emission (NBE) peak of the W-doped  $V_2O_5$  nanorods is in good agreement with those obtained from UV–Vis analyses. However, PL quenching is observed at the W-loaded flms compared to the undoped  $V_2O_5$  film. With increasing length of the W-doped  $V_2O_5$ nanorods, PL quenching might originate from increasing surface states which are generally formed in the nanosized crystal structures because surface states may act as traps and thereby may decrease excitonic emission. Besides, the infrared shift with the W dopant indicates that excess electrons are localized at the empty orbital of V and W atoms and as a result, the band gap is decreased

<span id="page-11-0"></span>



<span id="page-11-1"></span>**Fig. 8** Photoluminescence spectra of the undoped and doped  $V_2O_5$  thin films

in the W-doped samples with increasing localized states in the band gap. On the other hand, the quenching intensity of the red emission peaks (deep-level emission) is attributable to decreased density of the oxygen vacancy defects in the W-doped  $V_2O_5$  nanorods, especially for doped one at 1 wt%. This situation can be attributed to that the W dopant accelerated the absorption of oxygen at lower postannealing temperature due to high electron negativity than V as the mentioned above.

The optical absorbance and transmittance graphs of the prepared flms in the range of 1.5–3.0 eV energy are depicted in Fig. [9](#page-12-0). The inset in Fig. [9](#page-12-0)a shows the absorption spectra of all the flms in the near UV–vis region. Absorption edges of the W-doped nanostructured flms shifted to lower energy values, and they, respectively, change to the shorter wavelengths for 10, 3, 5 and 1 wt% W, respectively. From the absorbance measurements, using the following equations:



· VO (E<sub>u</sub>:454 meV) ·VO-1% (E<sub>u</sub>:400 meV) **10,6** ·VO-3% (E<sub>u</sub>:369 meV) · VO-5% (E<sub>u</sub>:399 meV) **ln (**α**) (c m -1 ) 10,5** · VO-10% (E<sub>u</sub>:405 meV) **10,4 10,3 (c) 1,95 2,00 2,05 2,10 2,15 2,20 2,25 2,30 2,35 2,40 2,45 Energy (eV)** 35 12  $(%)$ **Avg. Transmittance (%) Avg. Reflectance (%)** 11 **50** 30 10 25 Transmittance (%) **Transmittance (%)** 9 **40** 20 8 15 **30** <sup>0</sup> <sup>1</sup> <sup>2</sup> Dopant ratio (wt. %) <sup>8</sup> <sup>9</sup> <sup>10</sup> <sup>11</sup> **20** VO VO-1% VO-3% **10** VO-5% **(d)** VO-10% **1,6 1,8 2,0 2,2 2,4 2,6 2,8 3,0 Energy (eV)**

<span id="page-12-0"></span>**Fig. 9** The plots of absorption coefficient versus photon energy with inset UV–visible absorption spectra (a);  $(\alpha h\nu)^2$  versus  $h\nu$ , the photon energy (**b**); the corresponding  $ln(a)$  versus photon energy (**c**) and

$$
T = \frac{(1 - R)^2 \exp(-\alpha t)}{(1 - R^2) \exp(-2\alpha t)} = \exp(-A),
$$
\n(4)

$$
\alpha h v = K \left( h v - E_g \right)^{1/2},\tag{5}
$$

absorbance values are converted to absorption coefficients in order to find the band gaps of thin films from the  $\alpha^2$  versus *hν* plots, as seen in Fig. [9b](#page-12-0), where *R* is the refectance,  $\alpha$  is the absorption coefficient, *t* is the film thickness,  $h\nu$ is the incident photon energy and  $K$  is the constant. Using these graphs, the calculated band gap  $(E_{\alpha})$  of the reference flms is 2.45 eV, which is compatible with the reported values of the  $V_2O_5$  film prepared by different processes. On the other hand, band gaps of the W-doped nanostructured flms are narrower than that of undoped flm and values are 2.29 eV, 2.27 eV, 2.24 eV and 2.39 eV according to increasing rate of W-contribution, respectively. Although the band gap values signifcantly reduced in the W-doped  $V_2O_5$  films compared to the undoped  $V_2O_5$  film, the band

optical transmittance (**d**) with inset graph of average transmittance and refectance versus dopant ratio of the thin flms, respectively

gap value of the 10 wt% W-doped  $V_2O_5$  film showed an increase. The increase in the band gap is attributed to the microstructural changes, as seen in the SEM images, and disappearance of the monoclinic  $VO<sub>2</sub>$  secondary phase having lower band gap energy  $(<1 eV)$  [[56](#page-17-16)], as seen in the XRD pattern. Thus, the increased oxygen and electrons in the  $V_2O_5$  lattice cause slight increase in the band gap. On the other hand, despite the  $VO<sub>2</sub>$  secondary phase in the undoped thin flm, the band gap is larger than that in others. It can be attributed to structural and morphological properties of the undoped  $V_2O_5$  film having the small grain size and random small-sized nanorod distribution, as seen in Table [1](#page-5-0). Absorption behavior of semiconductors arises from its crystallinity. Crystalline defects and impurities cause weak and broad absorption region, while strong absorption edge is due to photonic excitation of band to band in the optical band gap. Furthermore, tail region of strong absorption edge inform perturbation of structure and disorder of crystal system [\[57](#page-17-17)]. This region is called Urbach tail and Urbach empirical rule

as relation of the absorption coefficient  $(\alpha)$  and the photon energy (*hv*) can be determined as follows:

$$
\alpha = \alpha_o \exp(hv/E_u),\tag{6}
$$

where  $\alpha$ <sup>*o*</sup> is a constant,  $h\nu$  is the incident photon energy and  $E<sub>u</sub>$  is energy of the band tail (Urbach energy). From Eq. [6,](#page-13-0) the relation between plotting  $ln(a)$  and *hv* can be written as follows:

$$
\ln \alpha = \ln \alpha_o + \left( h v / E_{\rm u} \right). \tag{7}
$$

Figure [9c](#page-12-0) shows the plot of  $ln(a)$  versus *hv* in the range of the Urbach tail. The values of  $E<sub>u</sub>$  were calculated as 454 meV, 400 meV, 369 meV, 399 meV, 405 meV for the pure flm and the 1, 3, 5, 10 wt% W-doped nanostructured flms, respectively. These values are also evidence that the disordered atoms and defects in the structural bonding is decreased with increasing ratio of the W dopant, which is in accord with the XRD analyses, as given in Table [1](#page-5-0). Furthermore, it is also observed in the values of  $E<sub>u</sub>$  that the value of 3 wt% is critical for the increasing disordered atoms and defects in the structural bonding. It can be attributed to the unlike atomic radius and interconnection between nanorods, the gradually increasing dislocation density and strain in the vanadium oxide lattice with the increasing W.

Optical transmittances of the nanostructured flms in the range of 1.5–3.0 eV energy are illustrated in Fig. [9d](#page-12-0). The optical transmittance is also decreased with the W doping. The compatible change of the refectance with the band gap is attributed to efect of "self-shadowing efect" due to enhanced nanostructure in form of the nanorods [\[58\]](#page-17-18). The self-shadowing mechanism leads to voiding within the nanostructured flms due to ratio of the elongated columnar structure and angle of the nanorods. The band gap and average visible transmittance of the flms were calculated, as given in Table [2.](#page-13-1) For the changes clearly seen, graph of average transmittance and refectance versus dopant ratio is given in inset of Fig. [9d](#page-12-0). Interestingly, the average transmittance showed an opposite trend with the band gap of W-doped films. This trend is attributed to the absorption coefficient  $(\alpha)$  having relatively two distinct behaviors of the films in

<span id="page-13-1"></span>**Table 2** Measured and calculated basic optical values of the nanostructured flms

	Band gap $(eV)$	Avg. transmit- tance $(\%)$	Avg. reflectance $(\% )$
V <sub>O</sub>	2.45	34.21	6.94
$VO-1\%$	2.29	15.59	11.48
$VO-3%$	2.27	17.43	9.79
$VO-5%$	2.24	25.58	9.47
$VO-10\%$	2.39	14.50	7.55

the region of under and above of 2.5 eV. For example, while the  $\alpha$  of the 10% W-doped film is smaller than that the other doped films at the energies lower than 2.5 eV energy, the  $\alpha$ at the energies higher than 2.5 eV is the greatest.

<span id="page-13-0"></span>The refectivity spectra obtained by the measurement system of the flms are given in Fig. [10](#page-14-0). As seen in Fig. [10](#page-14-0)a, the refectance of the flms varies between 4 and 14% over visible region. For similar refectance behavior of the flms, it is observed that a strong increment occurred at the lowenergy region in the spectra while refectance between 1.8 and 3.0 eV gradually increased with increasing energy. However, the refectance increased nearly by four times and a broad increment peak occurred in the W-doped nanostructured flms at about 2.5 eV. This energy is very close to the band gap value of the undoped flm and then the changes of the behavior of absorbance and transmittance also can be attributed to increased refectance at the this energy region. Furthermore, a general increment in the refectance can be attributed to the efect of self-shadowing and increasing refection from the nanostructured flm surface.

Refractive index (*n*) is determined from refectance (*R*) and extinction values (*k*) with following relation:

$$
n = \frac{(1+R)}{(1-R)} + \sqrt{\frac{4R}{(1-R)^2} - k^2},
$$
\n(8)

where the extinction coefficient can be calculated from the absorption coefficient, a, by the following equation:  $k = \alpha \lambda/4\pi$ , where  $\lambda$  is the wavelength of the incident photons. The spectra of *n* show an increasing trend with increasing energy of the incident photons and varied in the range of 1.5–2.2 in the near UV–Vis region, as seen in Fig. [10b](#page-14-0). The fundamental electron excitation is given by a complex dielectric constant. It has two components, the real and imaginary, which are connected with *n* and *k* values. Therefore, the complex dielectric function  $(\epsilon_1)$  is related to the complex refractive index by the following equations:

$$
\varepsilon_1 = (n^2 + k^2) \tag{9}
$$

and

$$
\varepsilon_2 = 2nk,\tag{10}
$$

where  $n(\omega)$  and  $k(\omega)$  are frequency-dependent refractive index  $(n)$  and extinction coefficient  $(k)$ , respectively. The variations of the real and imaginary parts of dielectric constant dependent on the energy of incident photons are indicated in Fig. [11.](#page-15-1) As seen in Fig. [11](#page-15-1)a, while the spectral behavior of the refractive index is dominant on the variation of the real dielectric constant, the imaginary dielectric constant has a spectral behavior as a variation of extinction coefficient which is given in inset of Fig.  $11b$ . The minimum and maximum values of  $\varepsilon_1$  varies between 2.2 and 4.9, respectively. The value of  $\varepsilon_2$  has limits of 0.1–0.9 in the visible region.

<span id="page-14-0"></span>



# **4 Conclusion**

In conclusion, the nanostructured films of W-doped  $V<sub>2</sub>O<sub>5</sub>$  were successfully deposited by the spray pyrolysis technique. Besides, formation, morphological variation, structural and optical properties of nanostructured  $V<sub>2</sub>O<sub>5</sub>$  films were explained in detail. The XRD patterns demonstrated mainly three diferent crystal structures of monoclinic VO<sub>2</sub>, orthorhombic  $V_2O_5$  and tetragonal  $V_2O_5$ . The W dopant caused change of the phase with reducing

monoclinic  $VO<sub>2</sub>$  with increasing dopant ratio and mixture phase of α- and  $\beta$ -V<sub>2</sub>O<sub>5</sub> phase is obtained at 10 wt% W. The SEM and the AFM images show that the undoped flm prepared using basic precursor solution at 400 °C annealing temperature formed with small-sized nanostructural morphology, while enhancing nanostructural morphology has been clearly seen in the flms obtained with the different ratios of the W dopant at the same experimental condition. In addition to the SEM and the AFM measurements, the Raman spectra suggested that the surface

<span id="page-15-1"></span>



difusion phenomenon which plays an important role in the growth process of nanorods due to high dependence on the oxygen absorption is enhanced with the W dopant. The PL spectra show PL quenching in the W-doped  $V_2O_5$ nanorods, indicating surface states with increasing length of the nanorods may act as traps and thereby may decrease excitonic emission. Furthermore, the optical absorbance, transmittance and reflectance spectra with the related graphs demonstrated the changed electronic behaviors such as band gap, refective index and dielectric constants.

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