# **Mathematical Physics**



# **An Abstract Birkhoff Normal Form Theorem and Exponential Type Stability of the 1d NLS**

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**Abstract:** We study stability times for a family of parameter dependent nonlinear Schrödinger equations on the circle, close to the origin. Imposing a suitable Diophantine condition (first introduced by Bourgain), we prove a rather flexible Birkhoff Normal Form theorem, which implies, e.g., exponential and sub-exponential time estimates in the Sobolev and Gevrey class respectively.

# **Contents**



#### <span id="page-1-0"></span>**1. Introduction and Main Results**

We consider families of NLS equations on the circle with external parameters of the form:

<span id="page-1-3"></span>
$$
i u_t + u_{xx} - V * u + f(x, |u|^2)u = 0,
$$
\n(1.1)

where  $i = \sqrt{-1}$  and  $V*$  is a Fourier multiplier

$$
V * u = \sum_{j \in \mathbb{Z}} V_j u_j e^{ijx}, \quad (V_j)_{j \in \mathbb{Z}} \in w_q^{\infty},
$$

living in the weighted  $\ell^{\infty}$  space

$$
w_q^{\infty} := \{ V = (V_j)_{j \in \mathbb{Z}} \in \ell^{\infty} \mid |V|_q := \sup_{j \in \mathbb{Z}} |V_j| \langle j \rangle^q < \infty \}, \quad q \ge 0,
$$

where  $\langle i \rangle := \max\{|i|, 1\}$ , while  $f(x, y)$  is  $2\pi$  periodic and real analytic in *x* and is real analytic in *y* in a neighborhood of  $y = 0$ . We shall assume that  $f(x, y)$  has a zero in  $y = 0$ . By analyticity, for some a,  $R > 0$  we have

<span id="page-1-2"></span>
$$
f(x, y) = \sum_{d=1}^{\infty} f^{(d)}(x) y^d, \quad |f|_{a, R} := \sum_{d=1}^{\infty} |f^{(d)}|_{\mathbb{T}_a} R^d < \infty,
$$
 (1.2)

where, given a real analytic function  $g(x) = \sum_{j \in \mathbb{Z}} g_j e^{ijx}$ , we set<sup>[1](#page-1-1)</sup>  $|g|_{\mathbb{T}_a}^2 := \sum_{j \in \mathbb{Z}} g_j e^{ijx}$  $|g_j|^2 e^{2a|j|}$ . Note that if *f* is independent of *x* [\(1.2\)](#page-1-2) reduces to

<span id="page-1-5"></span>
$$
|f|_{R} := \sum_{d=1}^{\infty} |f^{(d)}| R^d < \infty.
$$
 (1.3)

Equation [\(1.1\)](#page-1-3) is at least locally well-posed (say in a neighborhood of  $u = 0$  in  $H^1$ , see e.g. Lemma [5.4\)](#page-23-0) and has an elliptic fixed point at  $u = 0$ , so that an extremely natural question is to understand *stability times* for small initial data. One can informally state the problem as follows: let  $E \text{ }\subset H^1$  be some Banach space and consider [\(1.1\)](#page-1-3) with initial datum  $u_0$  such that  $|u_0|_E \le \delta \ll 1$ . By local well posedness, the solution  $u(t, x)$ of  $(1.1)$  with such initial datum exists and is in  $H<sup>1</sup>$ .

We call *stability time*  $T = T(\delta)$  the supremum of the times *t* such that for all  $|u_0|_E \leq \delta$  one has  $u(t, \cdot) \in E$  with  $|u(t, \cdot)|_E \leq 2\delta$ .

Computing the stability time  $T(\delta)$  is out of reach, so the goal is to give lower (and possibly upper) bounds.

A good comparison is with the case of a finite dimensional Hamiltonian system with a non-degenerate elliptic fixed point, which in the standard complex symplectic coordinates  $u_j = \frac{1}{\sqrt{2}}$  $\frac{1}{2}(q_j + i p_j)$  is described by the Hamiltonian

<span id="page-1-4"></span>
$$
\sum_{j=1}^{n} \omega_j |u_j|^2 + O(u^3), \quad \text{where } \omega_j \in \mathbb{R} \text{ are the linear frequencies.} \tag{1.4}
$$

<span id="page-1-1"></span><sup>&</sup>lt;sup>1</sup> Namely *g* is a holomorphic function on the domain  $\mathbb{T}_a := \{x \in \mathbb{C}/2\pi\mathbb{Z} : |\text{Im } x| < a\}$  with  $L^2$ -trace on the boundary.

Here if the frequencies  $\omega$  are sufficiently non degenerate, say diophantine,<sup>2</sup> then one can prove exponential lower bounds on  $T(\delta)$  and, if the nonlinearity satisfies some suitable hypothesis (e.g. convexity or steepness ), even super-exponential ones. This was proved in [\[MG95](#page-64-0)] (see also the recent paper [\[BFN15\]](#page-63-0) and references therein).

The strategy for obtaining exponential bounds is made of two main steps. The first one consists in the so-called Birkhoff normal form procedure: after  $N \geq 1$  steps the Hamiltonian [\(1.4\)](#page-1-4) is transformed into

<span id="page-2-1"></span>
$$
\sum_{j=1}^{n} \omega_j |u_j|^2 + Z + R , \qquad (1.5)
$$

where *Z* depends only on the actions  $(|u_i|^2)_{i=1}^n$  while  $R = O(|u|^{2N+3})$  contains terms of order the contains that this proceedure contributions of of order at least  $2N + 3$  in |*u*|. It is well known that this procedure generically diverges in N, so the second step consists in finding  $N = N(\delta)$  which minimizes the size of the remainder *R*.

The problem of *long-time* stability for equations [\(1.1\)](#page-1-3) has been studied by many authors. In the context of infinite chains with a finite range coupling, we mention [\[BFG88](#page-63-1)]. Regarding applications to PDEs (and particularly the NLS) the first results were given in [\[Bou96a](#page-63-2)] by Bourgain, who proved polynomial bounds for the stability times in the following terms: for any N there exists  $p = p(N)$  such that initial data which are  $\delta$ -small in the  $H^{p'+p}$  norm stay small in the  $H^{p'}$  norm, for times of order  $\delta^{-N}$ . Afterwards, Bambusi in [\[Bam99b](#page-63-3)] proved that superanalytic initial data stay small in analytic norm, for times of order  $e^{\ln(\frac{1}{\delta})^{1+b}}$ , where  $b > 0$ .

Following the strategy proposed in [\[Bam03\]](#page-63-4) for the Klein–Gordon equation Bambusi and Grébert in [\[BG03\]](#page-63-5) first considered Eq.  $(1.1)$  on  $\mathbb{T}^d$  and then, in [\[BG06\]](#page-63-6), proved polynomial bounds for a class of *tame-modulus* PDEs, which includes [\(1.1\)](#page-1-3). Their main result is that for any  $N \gg 1$  there exists  $p(N)$  (tending to infinity as  $N \to \infty$ ) such that for all  $p \ge p(N)$  and all  $\delta$ -small initial data in  $H^p$  one has  $T \ge C(N, p)\delta^{-N}$ , provided  $\delta < \delta_0(N, p)$ . Similar results were also proved for the Klein–Gordon equation on tori and Zoll manifolds in [\[DS04](#page-63-7)[,DS06](#page-63-8)[,BDGS07\]](#page-63-9). Successively Faou and Grébert in [\[FG13](#page-63-10)] considered the case of analytic initial data and proved subexponential bounds of the form  $T \ge e^{\ln(\frac{1}{\delta})^{1+b}}$ ,  $b > 0$ , for classes of NLS equations in  $\mathbb{T}^d$  (which include  $(1.1)$  by taking  $d = 1$ ). Regarding derivative NLS equations, the first results were in [\[YZ14](#page-64-1)] for the semilinear case. Recently, Feola and Iandoli in [\[FI\]](#page-63-11) prove polynomial lower bounds for the stability times of reversible NLS equations with two derivatives in the nonlinearity.

A closely related topic is the study of orbital stability times close to periodic or quasi-periodic solutions of [\(1.1\)](#page-1-3). In the case  $E = H^1$ , Bambusi in [\[Bam99a](#page-63-12)] proved a lower bound of the form  $T \ge e^{\delta^{-b}}$ ,  $b > 0$ , for perturbations of the integrable cubic NLS close to a quasi-periodic solution. Regarding higher Sobolev norms, most results are in the periodic case. See [\[FGL13](#page-63-13)] (polynomial bounds for Sobolev initial data) and the preprint [\[MSW18](#page-64-2)] (subexponential bounds for Gevrey initial data).

A dual point of view is to construct special orbits for which the Sobolev norms grow as fast as possible (thus giving an upper bound on the stability times). As far as we are aware such results are mostly on  $\mathbb{T}^2$  and in parameterless cases (for instance [\[CKS+10](#page-63-14)[,GK15](#page-63-15),[GHP16\]](#page-63-16)) and the time scales involved are much longer than our stability

<span id="page-2-0"></span><sup>&</sup>lt;sup>2</sup> A vector  $\omega \in \mathbb{R}^n$  is called diophantine when it is badly approximated by rationals, i.e. it satisfies, for some  $\gamma$ ,  $\tau > 0$ ,  $|k \cdot \omega| \ge \gamma |k|^{-\tau}$ ,  $\forall k \in \mathbb{Z}^n \setminus \{0\}$ .

times (see [\[Gua14](#page-64-3)] for the instability of  $(1.1)$  on  $\mathbb{T}^2$  and [\[Han14](#page-64-4)] for the instability of the plane wave in  $H^p$  with  $p < 1$ ).

In this paper we propose an abstract Birkhoff normal form result (see Theorem [1.3\)](#page-10-0) on weighted sequence spaces (based on  $\ell^2$ ) and deduce from it stability estimates for initial data in analytic, Gevrey and Sobolev class. An important difference of our approach with respect to the aforementioned papers and one of the main motivations of our work is that we use a different diophantine non-resonance condition on the linear frequencies, originally introduced in [\[Bou05](#page-63-17)] in the context of almost-periodic solutions. More precisely set

$$
\Omega_q := \left\{ \omega = \left( \omega_j \right)_{j \in \mathbb{Z}} \in \mathbb{R}^{\mathbb{Z}}, \quad \sup_j |\omega_j - j^2| \langle j \rangle^q < 1/2 \right\} \tag{1.6}
$$

and, for  $\gamma > 0$ , define the set of "good frequencies" as

<span id="page-3-0"></span>
$$
D_{\gamma,q} := \left\{ \omega \in \Omega_q : |\omega \cdot \ell| > \gamma \prod_{n \in \mathbb{Z}} \frac{1}{(1 + |\ell_n|^2 \langle n \rangle^{2+q})} . \quad \forall \ell \in \mathbb{Z}^{\mathbb{Z}} : |\ell| < \infty \right\} .
$$
\n(1.7)

It is known that  $D_{\gamma,q}$  is large with respect to a natural probability product measure on  $q_a$  (for a proof see [\[Bou05\]](#page-63-17) or Lemma [4.1](#page-16-0) in the present paper). It turns out that such diophantine conditions are very natural and easy to use in the context of PDEs on the circle with a superlinear dispersion law. Then from now on we shall fix  $\gamma > 0$ ,  $q > 0$ and assume that  $\omega \in D_{\nu,q}$ .

*Remark 1.1.* We note that some non-resonance condition on the frequencies is inevitable if one wants to prove *long-time* stability, indeed if one takes  $V = 0$  and  $f(x, |u|^2) = |u|^4$ then one can exhibit orbits in which the Sobolev norm is unstable in times of order  $\delta^{-4}$ , see [\[GT12](#page-63-18)[,HP17\]](#page-64-5).

At the formal level our BNF scheme is identical to the one used in finite dimensional systems, see formula  $(1.5)$ . The fact that such a scheme may be applied in an infinite dimensional context follows from introducing a suitable norm (see Definition [1.2](#page-9-0) and the comments thereafter); it turns out that our norm has explicit (and for us quite surprising) immersion properties (see Proposition [3.1\)](#page-14-0) and allows good bounds on the solution of the *homological equation* (see Lemma [4.2\)](#page-16-1). The gist of these properties is that they ensure that any vector field mapping a (neighborhood of) given Hilbert space in itself also maps (smaller neighborhoods of) *more regular* Hilbert spaces in themselves. Analogously also the vector field solving the homological equation maps *sufficiently more regular* Hilbert spaces in themselves.

To show that our procedure works in significative cases, we have computed stability times for various regularity classes. More precisely we improve the results in [\[FG13\]](#page-63-10) on *analytic and Gevrey* initial data, see Theorem [1.1.](#page-4-1) Moreover we recover [\[BG06\]](#page-63-6) on *Sobolev* initial data, giving an explicit control on the dependence of the stability time and of the smallness condition on the regularity, see Proposition [1.1](#page-5-0) and the improved estimates of Theorem [1.2.](#page-6-0)

**Comments on possible generalizations**. In this paper we have considered the simplest possible example of dispersive PDE on the circle. One can easily see that the same strategy can be followed word by word in more general cases provided that the non-linearity does not contain derivatives and that the dispersion law is superlinear. A much more challenging question is to consider NLS models with derivatives in the non-linearity. As we have mentioned a semilinear case was discussed by [\[CMW\]](#page-63-19). A very promising approach to Birkhoff normal form for quasilinear PDEs is the one of [\[BD18,](#page-63-20)[BDG10](#page-63-21)[,BDGS07,](#page-63-9)[BFG88](#page-63-1)[,BFG18](#page-63-22)[,BFN15](#page-63-0),[BG03](#page-63-5),[BG06,](#page-63-6)[Bou96a,](#page-63-2) [Bou96b](#page-63-23)[,Bou05](#page-63-17)[,CKS+10](#page-63-14)[,CLSY](#page-63-24)[,CMW](#page-63-19)[,Del12\]](#page-63-25) which was applied to fully-nonlinear reversible NLS equations in [\[FI](#page-63-11)]. It seems very plausible (at least in the reversible case) that one can adapt their methods (based on paralinearizations and paradifferential calculus) to our setting.

A natural generalization would be the extension to higher dimensions. While the immersion properties would work essentially in the same way, the diophantine condition should be adapted, for instance one could use the condition in [\[FG13](#page-63-10)].

Equation [\(1.1\)](#page-1-3) contains infinitely many external parameters. Of course one would like to consider parameterless equations as in the very interesting recent preprint [\[BFG18](#page-63-22)]. In this direction a natural question would be to understand if one could impose similar diophantine conditions by tuning only one parameter such as the mass in the beam or wave equations (see, e.g., [\[Bam03](#page-63-4), BD18]).

Before explaining the abstract BNF procedure in detail let us describe our stability results.

<span id="page-4-0"></span>*1.1. Stability results.* **Analytic and Gevrey initial data**. Our result is similar to [\[FG13\]](#page-63-10) in the sense that we also prove *subexponential bounds* on the time. We mention however that in [\[FG13](#page-63-10)] the control of the Sobolev norm in time is in a lower regularity space w.r.t. the initial datum. Recently we have been made aware of a preprint by Cong, Mi and Wang [\[CMW\]](#page-63-19) in which the authors give subexponential bounds for Gevrey initial data of a model like [\(1.1\)](#page-1-3), very similar to ours. A difference is that in their case the non linearity contains a derivative (see the comments after Theorem [1.1\)](#page-4-1) but satisfies momentum conservation. The two results were obtained independently and contemporarily, anyway, the overall strategies of proofs are quite different. In particular our result is a consequence of the general Birkhoff Normal Form Theorem [1.3](#page-10-0) and the non-resonance conditions are different (recall [\(1.7\)](#page-3-0)).

To state our result, let us fix  $0 < \theta < 1$ , and define the function space<sup>3</sup>

<span id="page-4-3"></span>
$$
\mathcal{H}_{p,s,a} := \left\{ u(x) = \sum_{j \in \mathbb{Z}} u_j e^{ijx} \in L^2 : |u|_{p,s,a}^2 := \sum_{j \in \mathbb{Z}} |u_j|^2 \langle j \rangle^{2p} e^{2a|j| + 2s\langle j \rangle^{\theta}} < \infty \right\}.
$$
\n(1.8)

with the assumption  $a \geq 0$ ,  $s > 0$ ,  $p > 1/2$ . We remark that if  $a > 0$  this is a space of analytic functions, while if  $a = 0$  the functions have Gevrey regularity. Note that for technical reasons connected to the way in which we control the small divisors, we *cannot* deal with the purely analytic case  $\theta = 1$ , see Lemmas [6.1,](#page-28-0) [7.1.](#page-32-1) For this reason we denote this result as G (Gevrey case).

<span id="page-4-1"></span>Our result, stated below, depends on some constants  $\delta_{\text{G}}$ ,  $T_{\text{G}}$ , *explicitely defined* in Subsection A, and depending only on  $\gamma$ , *q*, a, *R*,  $|f|_{a,R}$ , *p*, *s*, *a*,  $\theta$ .

<span id="page-4-2"></span><sup>&</sup>lt;sup>3</sup> Actually  $H_{p,s,a}$  also depends on  $\theta$ , however, since we think  $\theta$  fixed, we omit to write explicitly the dependence on it.

**Theorem 1.1** (Gevrey Stability). *Fix any a* > 0, *s* > 0 *such that a* + *s* < a *and any*  $p > 1/2$ *. For any*  $0 < \delta \leq \delta_G$  *and any*  $u_0$  *such that* 

$$
|u_0|_{p,s,a}\leq \delta,
$$

*the solution u(t) of* [\(1.1\)](#page-1-3) *with initial datum u(0) =*  $u_0$  *exists and satisfies* 

$$
|u(t)|_{p,s,a} \leq 2\delta \text{ for all times } |t| \leq \frac{T_G}{\delta^2} e^{\left(\ln \frac{\delta_G}{\delta}\right)^{1+\theta/4}}.
$$

*Remark 1.2.* Some comments on Theorem [1.1](#page-4-1) are in order.

1. The main point in the proof is to verify that the abstract Birkhoff Normal Form Theorem [1.3](#page-10-0) is applicable. Then we put the Hamiltonian of the NLS in Birkhoff normal form:

<span id="page-5-3"></span>
$$
\sum_{j\in\mathbb{Z}}\omega_j|u_j|^2 + Z + R\,,\tag{1.9}
$$

where *Z* depends only on the actions  $(|u_i|^2)_{i \in \mathbb{Z}}$  while  $R = O(|u|^{2N+3})$  is analytic in a ball centered at zero of  $h_{p,s,a}$  and has a zero of order at least  $2N + 3$  in  $u = 0$ . Then we find  $N = N(\delta)$  which minimizes the size of the remainder *R*.

2. We did not make an effort to maximize the exponent  $1 + \theta/4$  in the stability time. In fact, by trivially modifying the proof, one could get  $1 + \theta/(2^+)$ . We remark that in [\[CMW](#page-63-19)], in which  $\theta = 1/2$ , the exponent is better, i.e. it is  $1 + 1/(2^+)$ .

**Sobolev initial data**. Here our first goal was to recover by our methods the result of [\[BG06](#page-63-6)], computing explicitly all the constants in the estimates. In particular it is fundamental to have a good control on the dependence of the stabiliy time *T* on the the regularity *p*. Indeed there are two natural ways of taking a small ball around zero: reducing the size  $\delta$  or increasing the regularity  $p$ . A crucial point is that, in the case of Sobolev regularity, the number of BNF steps that one may perform is (apparently unavoidably) tied to the regularity *p*. This is clearly seen in [\[BG06\]](#page-63-6), where the number of steps is <sup>∼</sup> <sup>√</sup>*p*. It seemed an interesting point to verify how our approach worked in such a case, and wether we would see the same phenomenon.

<span id="page-5-0"></span>As before, our estimates depend on some constants, denoted by  $\tau_S$ ,  $\delta_S$ ,  $k_S$ ,  $T_S$ , *explicitly defined* in "Appendix A". These constants depend only on  $\gamma$ , *q*, a, *R*,  $|f|_{a,R}$ .

**Proposition 1.1** (A quantitative version of [\[BG06\]](#page-63-6)). *Consider Eq.*[\(1.1\)](#page-1-3)*with f satisfying*  $(1.2)$  *for* a,  $R > 0$ *. For any*  $p \geq 3\tau_s + 1$  *and any initial datum*  $u(0) = u_0$  *satisfying* 

<span id="page-5-1"></span>
$$
|u_0|_{H^p} := |u_0|_{L^2} + |\partial_x^p u_0|_{L^2} \le \delta \le \delta_S(\mathbf{k}_S p)^{-3p}
$$
 (1.10)

*the solution u(t) of* [\(1.1\)](#page-1-3) *with initial datum u(0)* =  $u_0$  *exists and satisfies* 

<span id="page-5-2"></span>
$$
|u(t)|_{H_p} \le 4\delta \text{ for all times } |t| \le \mathcal{T}_\mathcal{S} p^{-5p} \left(\frac{\delta_\mathcal{S}}{\delta}\right)^{\frac{2(p-1)}{\tau_\mathcal{S}}}.
$$
 (1.11)

*Remark [1.3.](#page-10-0)* Also in this result we just have to verify the hypotheses of Theorem 1.3. However as it happens in  $[BG06]$  the maximum number N of steps of BNF we can perform depends on *p*, in particular  $N = \left[\frac{p-1}{\tau_s}\right]$ . This is in fact slightly better than the previously cited paper (N  $\sim p$  instead of  $\sqrt{p}$ ). On the other hand it is not difficult to show that the bound  $\delta \leq \delta_S(k_S p)^{-3p}$  is essentially optimal (see Remark [10.1\)](#page-40-1).

Looking at the proof of the Theorem or even constructing other finite-dimensional models, one can see that in the traslation invariant case, the very restrictive smallness condition in  $(1.10)$  is only due to interactions between the modes 0, 1,  $-1$  and all the others. It then seems natural to consider initial data for which the energy on such modes is smaller, namely  $|u_0|_{L^2} \leq 2^{-p}\delta$ . We refer to this case as M, the relevant constants can be found in "Appendix A"

<span id="page-6-0"></span>**Theorem 1.2.** *Consider Eq.* [\(1.1\)](#page-1-3) with f independent of x and satisfying [\(1.3\)](#page-1-5) for  $R > 0$ . *For any p*  $> 3\tau_M + 1$  *and for any initial datum u*(0) = *u*<sub>0</sub> *satisfying* 

<span id="page-6-1"></span>
$$
|u_0|_{H^p} \le \delta \le \frac{\delta_{\mathcal{M}}}{\sqrt{p}}, \quad |u_0|_{L^2} \le 2^{-p}\delta \tag{1.12}
$$

*the solution u*(*t*) *of* [\(1.1\)](#page-1-3) *exists and satisfies*

<span id="page-6-4"></span>
$$
|u(t)|_{H^p} \le 8\delta \text{ for all times } |t| \le \mathcal{T}_{\mathbb{M}} \left( \frac{8\delta_{\mathbb{M}}^2}{(p-1)\delta^2} \right)^{\frac{p-1}{\tau_{\mathbb{M}}}}. \tag{1.13}
$$

*Remark 1.4.* Note that, since the  $L^2$  norm is a constant of motion, one trivially has  $|u(t)|_{L^2} < 2^{-p}\delta$ . Comparing with [\(1.11\)](#page-5-2), we see that the time estimate is more or less the same but now it holds in a *much bigger neighborhood* of zero ( $\delta \leq p^{-1/2}$  instead of <sup>≤</sup> *<sup>p</sup>*−3*p*).

If one requires a stronger condition on the  $L^2$  norm, i.e.,  $|u_0|_{L^2} \leq 3^{-p}\delta$ , it turns out that the size of the perturbation is *exponentially decreasing in p* and, therefore, keeping δ fixed and sending *p* to infinity one immediately obtains stability.

The main difference between the Gevrey and Sobolev cases is that in the latter the number of BNF steps N depends on the regularity, while in the former it is independent. Thus in the Sobolev case we cannot fix both  $\delta$  and  $p$  and optimize in N. What we can do is to fix δ and find an optimal regularity *p*(δ), which maximizes the stability time. It turns out that the two cases S and M behave differently. Indeed the weaker smallness condition [\(1.12\)](#page-6-1) allows us to take much bigger  $p(\delta)$ , obtaining much longer stability times. As before our statements depend on some constants, denoted by  $\delta_{\rm S}$ ,  $\delta_{\rm M}$  *explicitly defined* in Subsection A.

<span id="page-6-2"></span>**Corollary 1.1** (Sobolev stability: optimization). (S) *For any*  $0 < \delta \leq \delta_S$  *and any*  $u_0$  *such that* 

<span id="page-6-5"></span>
$$
|u_0|_{H^p} \le \delta, \quad p = p(\delta) := 1 + \frac{\ln(\delta_S/\delta)}{6 \ln \ln(\delta_S/\delta)}, \tag{1.14}
$$

*the solution u(t) of* [\(1.1\)](#page-1-3) *with initial datum u(0) =*  $u_0$  *exists and satisfies* 

<span id="page-6-3"></span>
$$
|u(t)|_{H^p} \le 4\delta \quad \text{for all times} \quad |t| \le \mathcal{T}_\mathcal{S} e^{\frac{\ln^2(\delta_\mathcal{S}/\delta)}{4\tau_\mathcal{S}\ln\ln(\delta_\mathcal{S}/\delta)}}. \tag{1.15}
$$

(M) Assume that f in [\(1.1\)](#page-1-3) is independent of x. For any  $0 < \delta \leq \delta_M$  and

<span id="page-6-6"></span>
$$
\forall p \ge p(\delta) := \frac{\delta_{M}^{2}}{\delta^{2}}, \quad \forall u_{0} \quad s.t. \quad |u_{0}|_{H^{p}} \le \delta, \quad |u_{0}|_{L^{2}} \le 2^{-p}\delta,
$$
 (1.16)

*the solution u(t) of* [\(1.1\)](#page-1-3) *with initial datum u(0)* =  $u_0$  *exists and satisfies* 

<span id="page-6-7"></span>
$$
|u(t)|_{H^p} \le 8\delta \text{ for all times } |t| \le \mathbb{T}_M e^{\frac{\delta_M^2}{\tau_M \delta^2}}.
$$
 (1.17)

*Remark 1.5.* Some remarks on Corollary [1.1](#page-6-2) are in order.

Note that  $(1.15)$  is the stability time computed in [\[BFG88](#page-63-1)] for short range couplings.

- 1. We will prove the case M only for  $p = p(\delta)$ , the general case being analogous<sup>[4](#page-7-1)</sup> (with the same constants!) also if  $p \ge p(\delta)$ .
- 2. One can easily restate Corollary [1.1](#page-6-2) in terms of the Sobolev exponent *p*, instead of δ, since the map  $\delta \to p(\delta)$  is injective.

*Remark 1.6* (finite dimensional examples). It is interesting to compute the stability times predicted by our theorems for initial data supported on a finite number of modes. To this purpose consider an initial datum  $u^{(0)}$  uniformly distributed over the modes 1, ..., *j*:

$$
|u_i^{(0)}| = \varepsilon, \quad \forall i = 1, \dots, j.
$$

Theorem [1.1](#page-4-1) with  $a = 0$ ,  $p = 1$  states that if  $\varepsilon \le \varepsilon_G := \delta_G e^{-2j\theta}$  then  $u(t)$  stays stable, in Gevrey norm, for times of order  $e^{(\ln \frac{\varepsilon_G}{\varepsilon})^{1+\theta/4}}$ .

Now if  $\varepsilon \leq \varepsilon_M(p) := \delta_M j^{-p-1} / \sqrt{p}$  we have  $|u_0|_{H_p} < \delta$  and  $|u_0|_{L^2} < 2^{-p}\delta$ ; then by Theorem [1.2](#page-6-0) the solution  $u(t)$  stays stable, in  $\hat{H}^p$  norm, for times of order  $T \sim \left(\frac{\varepsilon_M(p)}{\varepsilon}\right)^{2(p-1)/\tau_M}$ . Maximizing the time in *p* with fixed  $\varepsilon$  we get

<span id="page-7-2"></span>
$$
p \sim \frac{\ln(\frac{\delta_{\rm M}^2}{\varepsilon})}{2\ln j}, \quad T \sim e^{\frac{(\ln \varepsilon^{-1})^2}{\ln j}} \tag{1.18}
$$

provided that  $\varepsilon \lesssim j^{-7\tau_{\text{M}}}$ . Explicitly we get a weaker constraint on  $\varepsilon$  and a better time estimate. Of course one could play the same game directly with the estimate of Proposition [1.1.](#page-5-0) As it should be expected the time estimate is more or less the same as [1.18](#page-7-2) but the smallness condition is much stronger, i.e. of the type  $\varepsilon \lesssim e^{-2j^{\theta}}$ .

<span id="page-7-0"></span>*1.2. The abstract Birkhoff Normal Form.* We start by setting our functional framework. The main point is to introduce a weighted majorant norm which penalizes the terms in the Hamiltonian which do not preserve momentum, see Definition [1.1.](#page-8-0)

Let us pass to the Fourier side via the identification

<span id="page-7-4"></span>
$$
u(x) = \sum_{j \in \mathbb{Z}} u_j e^{ijx} \ \mapsto \ u = (u_j)_{j \in \mathbb{Z}}, \tag{1.19}
$$

where *u* belongs to some complete subspace of  $\ell^2$ . Fix the symplectic structure to be

$$
i\sum_{j} du_j \wedge d\bar{u}_j. \tag{1.20}
$$

In this framework the Hamiltonian of [\(1.1\)](#page-1-3) is

<span id="page-7-3"></span>
$$
H_{\text{NLS}}(u) := D_{\omega} + P, \text{ where}
$$
  
\n
$$
D_{\omega} := \sum_{j \in \mathbb{Z}} \omega_j |u_j|^2,
$$
  
\n
$$
P := \int_{\mathbb{T}} F(x, |u(x)|^2) dx, \quad F(x, y) := \int_0^y f(x, s) ds. \quad (1.21)
$$

<span id="page-7-1"></span><sup>4</sup> Indeed, thanks to the immersion property of our norms (see Proposition [6.1](#page-26-1) below) the canonical transformation putting the system in Birkhoff Normal Form (see Theorem [8.1](#page-37-0) below) in the *p*-case is simply the restriction to  $H^p$  of the one of the  $p(\delta)$ -case.

We shall always work with quite regular solutions; given a real sequence  $w =$  $(w_i)_{i \in \mathbb{Z}}$ , with  $w_i \geq 1$  let us set the Hilbert space<sup>5</sup>

$$
h_w := \left\{ u := (u_j)_{j \in \mathbb{Z}} \in \ell^2(\mathbb{C}) : \ |u|_w^2 := \sum_{j \in \mathbb{Z}} w_j^2 |u_j|^2 < \infty \right\}.
$$
 (1.22)

As examples of  $h_w$  we consider:

- G) (Gevrey case)  $w(p, s, a) := \left(\langle j \rangle^p e^{a|j|+s\langle j \rangle^\theta}\right)$ *j*∈Z, which is isometrically isomorphic, by Fourier transform, to  $H_{p,s,a}$  defined in [\(1.8\)](#page-4-3).
- S) (Sobolev case) w(p) := w(p, 0, 0) =  $(\langle j \rangle^p)_{i \in \mathbb{Z}}$ , which is isometrically isomorphic, by Fourier transform, to  $H_{p,0,0}$  defined in [\(1.8\)](#page-4-3) and is equivalent to  $H^p$  equipped with the norm  $|\cdot|_{L^2} + |\partial_x^p \cdot|_{L^2}$  with equivalence constants independent of *p* (see [\(5.28\)](#page-24-1))
- M) (Modified-Sobolev case)  $w_j = \lfloor j \rfloor^p$ , where  $\lfloor j \rfloor := \max\{\lfloor j \rfloor, 2\}$ ; this space is equivalent to *H*<sup>*p*</sup> equipped with the norm  $2^p | \cdot |_{L^2} + |\partial_x^p \cdot |_{L^2}$  with equivalence constants independent of  $p$  (see  $(5.30)$ )

Here and in the following, given  $r > 0$ , by  $B_r(\mathbf{h}_w)$  we mean the closed ball of radius *r* centered at the origin of  $h_w$ .

In the following we always consider Hamiltonians  $H : B_r(h_w) \to \mathbb{R}$  such that there exists a pointwise absolutely convergent power series expansion<sup>6</sup>

$$
H(u) = \sum_{\substack{\boldsymbol{\alpha}, \boldsymbol{\beta} \in \mathbb{N}^{\mathbb{Z}}, \\ |\boldsymbol{\alpha}| + |\boldsymbol{\beta}| < \infty}} H_{\boldsymbol{\alpha},\boldsymbol{\beta}} u^{\boldsymbol{\alpha}} \bar{u}^{\boldsymbol{\beta}}, \qquad u^{\boldsymbol{\alpha}} := \prod_{j \in \mathbb{Z}} u_j^{\boldsymbol{\alpha}_j}
$$

with the following properties:

(i) Reality condition:

<span id="page-8-4"></span>
$$
H_{\alpha,\beta} = \overline{H}_{\beta,\alpha},\tag{1.23}
$$

this means that *H* is real analytic in the real and imaginary part of  $u$  (see section [2\)](#page-12-1); (ii) Mass conservation:

$$
H_{\alpha,\beta} = 0 \quad \text{if } |\alpha| \neq |\beta|,\tag{1.24}
$$

namely the Hamiltonian Poisson commutes with the mass  $\sum_{j\in\mathbb{Z}} |u_j|^2$ ;

The Hamiltonian functions being defined modulo <sup>a</sup> constant term, we shall assume without loss of generality that  $H(0) = 0$ .

We say that a Hamiltonian *H* as above preserves momentum when

<span id="page-8-3"></span>
$$
H_{\boldsymbol{\alpha},\boldsymbol{\beta}} = 0 \quad \text{if} \quad \pi(\boldsymbol{\alpha} - \boldsymbol{\beta}) := \sum_{j \in \mathbb{Z}} j(\boldsymbol{\alpha}_j - \boldsymbol{\beta}_j) \neq 0,
$$
 (1.25)

namely the Hamiltonian *H* Poisson commutes with  $\sum_{j\in\mathbb{Z}} j|u_j|$ <sup>2</sup>. Note that if the nonlinearity *f* in Eq. [\(1.1\)](#page-1-3) does not depend on the variable *x*, then the Hamiltonian *P* in [\(1.21\)](#page-7-3) preserves momentum.

<span id="page-8-0"></span><sup>&</sup>lt;sup>5</sup> Endowed with the scalar product  $(u, v)_{h_w} := \sum_{j \in \mathbb{Z}} w_j^2 u_j \bar{v}_j$ .

<span id="page-8-2"></span><span id="page-8-1"></span><sup>6</sup> As usual given a vector  $k \in \mathbb{Z}^{\mathbb{Z}}$ ,  $|k| := \sum_{j \in \mathbb{Z}} |k_j|$ .

**Definition 1.1** (*η*-majorant analytic Hamiltonians). For  $\eta > 0$ ,  $r > 0$  let  $\mathcal{A}_{r,n}(\mathbf{h}_{\mathbf{w}})$  be the space of Hamiltonians as above such that the  $\eta$ -majorant

<span id="page-9-1"></span>
$$
\underline{H}_{\eta}(u) := \sum_{\alpha,\beta \in \mathbb{N}^{\mathbb{Z}}} \left| H_{\alpha,\beta} \right| e^{\eta |\pi(\alpha-\beta)|} u^{\alpha} \bar{u}^{\beta} \tag{1.26}
$$

is point-wise absolutely convergent on  $B_r(h_w)$ . If we take  $\eta = 0$  we denote  $H_0(u) =$  $H(u)$  as the majorant of  $H$ .

The exponential weight  $e^{\eta |\pi(\alpha - \beta)|}$  is added in order to ensure that the monomials which do not preserve momentum have an exponentially small coefficient.

We will say that a Hamiltonian  $H(u) \in A_{r,\eta}(\mathbf{h}_w)$  is  $\eta$ -regular if  $X_{H_n}: B_r(\mathbf{h}_w) \to \mathbf{h}_w$ and is uniformly bounded, where  $X_{H_n}$  is the vector field associated to the  $\eta$ -majorant Hamiltonian in [\(1.26\)](#page-9-1). More precisely we give the following

<span id="page-9-0"></span>**Definition 1.2** (*η*-regular Hamiltonians). For  $\eta \ge 0$ ,  $r > 0$  let  $\mathcal{H}_{r,n}(\mathbf{h}_{\mathbf{w}})$  be the subspace of  $A_{r,n}$ (h<sub>w</sub>) of those Hamiltonians *H* such that

$$
|H|_{\mathcal{H}_{r,\eta}(\mathbf{h}_{\mathbf{w}})} = |H|_{r,\eta,\mathbf{w}} := r^{-1} \left( \sup_{|u|_{\mathbf{h}_{\mathbf{w}} \leq r}} \left| X_{\underline{H}_{\eta}} \right|_{\mathbf{h}_{\mathbf{w}}} \right) < \infty.
$$

We shall show in Sect. [2](#page-12-1) that this guarantees that the Hamiltonian flow of *H* exists at least locally and generates a symplectic transformation on  $h_w$ .

*Remark 1.7.* Definition [1.2](#page-9-0) with  $\eta = 0$ , i.e. the idea of controlling an analytic function through the sup of its *Cauchy majorant*, dates back to Cauchy-Kovalevskaya. In the context of analytic functions on Hilbert spaces, this class of functions is defined and studied, with a slightly different approach, in [\[Nik86\]](#page-64-6) and [\[KP10](#page-64-7)], where it is referred to as "normally analytic" functions.

Regarding the idea of introducing a weight which penalizes monomials which do not preserve momentum, this was used already in [\[Bam03](#page-63-4)].

In our work the crucial point is that all the dependence on the parameters  $r, \eta, w$  of the norm in Definition [1.2](#page-9-0) can be *encoded* in the coefficients

<span id="page-9-3"></span>
$$
c_{r,\eta,w}^{(j)}(\boldsymbol{\alpha},\boldsymbol{\beta}) := r^{|\boldsymbol{\alpha}|+|\boldsymbol{\beta}|-2} e^{\eta|\pi(\boldsymbol{\alpha}-\boldsymbol{\beta})|} \frac{\mathbf{w}_j^2}{\mathbf{w}^{\boldsymbol{\alpha}+\boldsymbol{\beta}}}, \qquad \mathbf{w}^{\boldsymbol{\alpha}+\boldsymbol{\beta}} = \prod_{j \in \mathbb{Z}} \mathbf{w}_j^{\boldsymbol{\alpha}_j+\boldsymbol{\beta}_j} \tag{1.27}
$$

defined for any  $\alpha$ ,  $\beta \in \mathbb{N}^{\mathbb{Z}}$  and  $j \in \mathbb{Z}$  (see formula [\(3.1\)](#page-14-1) and Lemma [3.1\)](#page-14-2). This allows us to give a simple and explicit condition which guarantees the *immersion*  $\mathcal{H}_{r,\eta}(\mathbf{h}_{w}) \subseteq$  $\mathcal{H}_{r',\eta'}(\mathbf{h}_{w'})$  in terms of the ratio of the coefficients  $c_{r,\eta,w}^{(j)}(\boldsymbol{\alpha},\boldsymbol{\beta}), c_{r',w}^{(j)}$ <sup>(*J*)</sup>, <sub>*r'*, *η'*, <sub>w</sub>'</sub> (α, *β*), see Propo-sition [3.1.](#page-14-0)

As it is well known a Birkhoff Normal Form is achieved by an iterative procedure. Let us describe the general step. Given a Hamiltonian

<span id="page-9-2"></span>
$$
H = \sum_{j \in \mathbb{Z}} \omega_j |u_j|^2 + Z + R,\tag{1.28}
$$

where *Z* is a normal form and *R* has a zero of degree say  $2d+2$  (with  $d \ge 1$ ) at  $u = 0$ , we look for a change of variables, which conjugates *H* to a Hamiltonian  $\sum_{j\in\mathbb{Z}} \omega_j |u_j|^2 + Z'$ +

 $R'$  so that now *R* has a zero of degree at least  $2d + 4$ . The desired change of variables is generated by the time one flow of a Hamiltonian *S* which solves the homological equation<sup>[7](#page-10-1)</sup>

$$
\{\sum_{j\in\mathbb{Z}}\omega_j|u_j|^2, S\}=R.
$$

As for the immersion properties, given<sup>[8](#page-10-2)</sup>  $r' \le r, \eta' \le \eta$  and  $w' \ge w$  such that  $\mathcal{H}_{r,n}(\mathbf{h}_w) \subseteq$  $\mathcal{H}_{r',\eta'}(\mathbf{h}_{w'})$ , in Proposition [4.2](#page-16-1) and Lemma [5.2](#page-18-0) we give a simple and explicit condition -in terms of the ratio of the coefficients  $c_{r,\eta,w}^{(j)}(\boldsymbol{\alpha},\boldsymbol{\beta})$ ,  $c_{r',w}^{(j)}$  $r'$ ,  $\eta'$ ,  $w'$  ( $\alpha$ ,  $\beta$ ) - which ensures that if  $R \in H_{r,\eta}(\mathbf{h}_{w})$  is appropriately small, then *S* is well defined and generates a close to identity change of variables  $B_{r'}(h_{w'}) \to h_{w'}$ . With this procedure we start in some phase space  $h_w$  and then show the existence of the Birkhoff change of variables on a ball which not only has a smaller radius but is taken in the stronger toplogy  $h_{w'}$ . Note that this is not a smoothing change of variables: it is defined *from the smaller space to itself*.

<span id="page-10-3"></span>Starting with a Hamiltonian as in [\(1.28\)](#page-9-2) with a zero of order 4, in order to reach the form [\(1.9\)](#page-5-3) we need to perform N steps of BNF. To this purpose we make the following

**Assumption 1.** We say that  $\eta \ge 0$  and two weights  $w_0 \le w$  satisfy the Birkhoff assumption at step  $N \geq 1$  if the following holds. The exists a sequence of weights  $w_0 \leq w_1 \leq \cdots \leq w_N = w$  such that

<span id="page-10-5"></span>
$$
\mathfrak{C} := \max \left\{ 1, \sup_{0 \le n < \mathbb{N}} \sup_{\substack{j, \boldsymbol{\alpha}, \boldsymbol{\beta} \\ \boldsymbol{\alpha}_j + \boldsymbol{\beta}_j \neq 0}} \frac{c_{\varrho_n^*, \eta_{n+1}, \mathbb{W}_{n+1}}^{(j)}(\boldsymbol{\alpha}, \boldsymbol{\beta})}{c_{\varrho_n, \eta_n, \mathbb{W}_n}^{(j)}(\boldsymbol{\alpha}, \boldsymbol{\beta})} \right\} < \infty,
$$
\n
$$
\hat{\mathfrak{K}} := \max \left\{ 1, \sup_{0 \le n < \mathbb{N}} \sup_{\substack{j, \boldsymbol{\alpha}, \boldsymbol{\beta} \\ \boldsymbol{\alpha}_j + \boldsymbol{\beta}_j \neq 0}} \frac{c_{\varrho_n^*, \eta_{n+1}, \mathbb{W}_{n+1}}^{(j)}(\boldsymbol{\alpha}, \boldsymbol{\beta})}{c_{\varrho_n, \eta_n, \mathbb{W}_n}^{(j)}(\boldsymbol{\alpha}, \boldsymbol{\beta})|\omega \cdot (\boldsymbol{\alpha} - \boldsymbol{\beta})|} \right\} < \infty,
$$
\n
$$
\hat{\mathfrak{K}}^{\sharp} := \max \left\{ 1, \sup_{0 \le n < \mathbb{N}} \sup_{\substack{j, \boldsymbol{\alpha}, \boldsymbol{\beta} \\ \boldsymbol{\alpha}_j + \boldsymbol{\beta}_j \neq 0}} \frac{c_{\varrho_n^*, \eta_{\mathbb{N}}, \mathbb{W}_n}^{(j)}(\boldsymbol{\alpha}, \boldsymbol{\beta})|\omega \cdot (\boldsymbol{\alpha} - \boldsymbol{\beta})|}{c_{\varrho_n^*, \eta_n, \mathbb{W}_n}^{(j)}(\boldsymbol{\alpha}, \boldsymbol{\beta})|\omega \cdot (\boldsymbol{\alpha} - \boldsymbol{\beta})|} \right\} < \infty, \quad (1.29)
$$

where

$$
\varrho_n = (2 - \frac{n}{N}), \quad \eta_n = (1 - \frac{n}{N})\eta, \quad 0 \le n \le N, \quad \varrho_n^* = \frac{\varrho_{n+1} + \varrho_n}{2}, \quad 0 \le n < N.
$$

Informally speaking  $\mathfrak{C} < \infty$  guarantees the immersion properties at each step, while  $\mathfrak{K} < \infty$  guarantees that one can solve the homological equation at each step. Finally  $\mathfrak{K}^{\sharp} < \infty$  guarantees that the composition of the changes of variables of all steps is well defined and close to identity on some ball  $B_r(h_{w_N})$ .

Let

<span id="page-10-4"></span>
$$
\mathcal{K}_r(\mathbf{h}_w) := \{ H \in \mathcal{H}_{r,0}(\mathbf{h}_w) \mid H = \sum_{\alpha \in \mathbb{N}^\mathbb{Z}} H_{\alpha,\alpha} |u|^{2\alpha} \}
$$
(1.30)

<span id="page-10-0"></span>be the subspace of normal form Hamiltonians.

<sup>&</sup>lt;sup>7</sup> Since  $\omega \in D_{\gamma,q}$ , w.l.o.g. we may assume that *R* is in the range of the operator  $\{\sum_{j\in\mathbb{Z}}\omega_j|u_j|^2,\cdot\}$ .

<span id="page-10-2"></span><span id="page-10-1"></span><sup>&</sup>lt;sup>8</sup> As usual  $w \leq w'$  means that  $w_j \leq w'_j$  for every  $j \in \mathbb{Z}$ .

**Theorem 1.3** (Abstract Birkhoff Normal Form). *Consider a Hamiltonian of the form*

<span id="page-11-2"></span>
$$
H = D_{\omega} + G, \quad D_{\omega} = \sum_{j} \omega_{j} |u|_{j}^{2}
$$
 (1.31)

*with*  $\omega \in D_{\gamma,q}$  *and*  $G \in \mathcal{H}_{\bar{r},\eta}(\mathbf{h}_{w_0}),$  *for some*  $\bar{r} > 0$ ,  $\eta \geq 0$ *. Assume moreover that* G *has a zero of order at least* 4 *at u* = 0*. Consider*  $N \ge 1$  *and*  $w \ge w_0$  *such that*  $\eta$ ,  $w_0$ , w *satisfy the Birkhoff assumption at step* N. *Set*

<span id="page-11-3"></span>
$$
\widehat{\mathbf{r}} := \min\left\{ \frac{r_{\star}}{\sqrt{\mathrm{N}\max\{\mathfrak{C}\mathfrak{K}, \mathfrak{K}^{\sharp}\}}}, \frac{\bar{r}}{2} \right\}, \quad \text{where} \quad r_{\star} := \bar{r}\sqrt{\frac{\gamma}{2^{11}e|G|_{\bar{r}, \eta, w_0}}}.\quad (1.32)
$$

*Then for all*  $0 < r \leq \hat{\tau}$  *there exists an invertible*<sup>[9](#page-11-0)</sup> *symplectic change of variables* 

<span id="page-11-4"></span>
$$
\Psi: B_r(\mathbf{h}_w) \mapsto B_{2r}(\mathbf{h}_w),
$$
  
\n
$$
\sup_{u \in B_r(\mathbf{h}_w)} |\Psi(u) - u|_w \leq \widehat{C}_1 r^3 \leq \frac{r}{8}, \quad \widehat{C}_1 := \frac{\mathfrak{K}^{\sharp}}{2^7 er_{\star}^2},
$$
\n(1.33)

*such that in the new coordinates*

$$
H \circ \Psi = D_{\omega} + Z + R, \quad Z \in \mathcal{K}_r(\mathbf{h}_{\mathbf{w}}),
$$

*where*

<span id="page-11-1"></span>
$$
|Z|_{r,0,w} \leq \widehat{C}_2 r^2, \quad |R|_{r,0,w} \leq \widehat{C}_3 r^{2(N+1)}, \quad with
$$
  

$$
\widehat{C}_2 := \frac{8|G|_{\bar{r},\eta,w_0}}{\bar{r}^2}, \quad \widehat{C}_3 := \frac{\gamma}{2^9 e r_{\star}^2} \left(\frac{\mathfrak{C} \mathfrak{K} N}{4r_{\star}^2}\right)^N.
$$
 (1.34)

The theorem follows by a straightforward iteration, see Sect. [5.](#page-17-0)

As it is well known the bounds [\(1.34\)](#page-11-1) imply a lower bound on the stability time; we discuss this in Corollary [5.1](#page-24-2) where we show that the solution  $u(t)$  of the Hamiltonian flow of [\(1.31\)](#page-11-2) with initial datum  $u(0) = u_0$  such that  $|u_0|_w \leq \frac{3r}{8}$  exists and satisfies

$$
|u(t)|_{w} \le r \quad \text{ for all times} \quad |t| \le \frac{1}{8\widehat{C}_3 r^{2(N+1)}}.
$$

By Theorem [1.3](#page-10-0) and Corollary [5.1,](#page-24-2) in order to prove the stability results we only need to define suitable sequence spaces verifying Assumption [1.](#page-10-3) In particular we consider the three applications G, S, M introduced at page 7. Another interesting example (suggested to us by Z. Hani) could be the space

$$
\Big\{(u_j)\in L^2: \ \ |u|^2:=\sum_j |u_j|^2 e^{\ln(|j|)^2}<\infty\Big\},\
$$

where  $\lfloor j \rfloor = \max\{|j|, 2\}$ . In this case one may get  $T \approx \delta^{\ln(\ln(1/\delta))}$ .

A preliminary version of these results was announced in [\[BMP19\]](#page-63-26).

<span id="page-11-0"></span><sup>&</sup>lt;sup>9</sup> In the sense that there exists a symplectic change of variables  $\Phi : B_r(h_w) \mapsto B_{2r}(h_w)$  such that  $\Psi \circ \Phi u = \Phi \circ \Psi u = u, \forall u \in B_{\frac{7}{8}r}(\mathbf{h}_{w}).$ 

#### <span id="page-12-1"></span><span id="page-12-0"></span>**Part 1. An Abstract Framework for Birkhoff Normal Form on Sequences Spaces**

### **2. Symplectic Structure and Hamiltonian Flows**

**Spaces of Hamiltonians**. As explained in the Introduction our wheighted spaces  $h_w$  are contained in  $\ell^2(\mathbb{C})$ , so we endow them with the standard symplectic structure coming from the Hermitian product on  $\ell^2(\mathbb{C})$ .

We identify  $\ell^2(\mathbb{C})$  with  $\ell^2(\mathbb{R}) \times \ell^2(\mathbb{R})$  through  $u_j = (x_j + iy_j)/\sqrt{2}$  and induce on  $\ell^2(\mathbb{C})$  the structure of a real symplectic Hilbert space<sup>10</sup> by setting, for any  $(u^{(1)}, u^{(2)}) \in$  $\ell^2(\mathbb{C}) \times \ell^2(\mathbb{C}),$ 

$$
\langle u^{(1)}, u^{(2)} \rangle = \sum_j \left( x_j^{(1)} x_j^{(2)} + y_j^{(1)} y_j^{(2)} \right), \quad \omega(u^{(1)}, u^{(2)}) = \sum_j \left( y_j^{(1)} x_j^{(2)} - x_j^{(1)} y_j^{(2)} \right),
$$

which are the standard scalar product and symplectic form  $\Omega = \sum_j dy_j \wedge dx_j$ . For convenience and to keep track of the complex structure, one often writes the vector fields and the differential forms in complex notation, that is

$$
\Omega = \mathbf{i} \sum_{j} du_{j} \wedge d\bar{u}_{j}, \quad X_{H}^{(j)} = \mathbf{i} \frac{\partial}{\partial \bar{u}_{j}} H
$$

where the one form and vector field are defined through the identification between C and  $\mathbb{R}^2$ , given by

$$
du_j = \frac{1}{\sqrt{2}}(dx_j + idy_j), \quad d\bar{u}_j = \frac{1}{\sqrt{2}}(dx_j - idy_j),
$$

$$
\frac{\partial}{\partial u_j} = \frac{1}{\sqrt{2}}\left(\frac{\partial}{\partial x_j} - i\frac{\partial}{\partial y_j}\right), \quad \frac{\partial}{\partial \bar{u}_j} = \frac{1}{\sqrt{2}}\left(\frac{\partial}{\partial x_j} + i\frac{\partial}{\partial y_j}\right).
$$

<span id="page-12-3"></span>*Remark 2.1.* By mass conservation and since  $H(0) = 0$ , it is straightforward to prove that the norm  $|\cdot|_{r,n,w}$  is increasing in the radius parameter *r* (see also Propositio[n3.1\)](#page-14-0).

Note that if  $|H|_{r,\eta,w} < \infty$  then *H* admits an analytic extension *H*, that is

$$
(u_+, u_-) \in B_r(\ell^2(\mathbb{C})) \times B_r(\ell^2(\mathbb{C})) \to \widehat{H}(u_+, u_-): \qquad H(u) = \widehat{H}(u, \bar{u}),
$$

whose Taylor series expansion is

$$
\widehat{H}(u_+, u_-) = \sum_{\alpha, \beta \in \mathbb{N}^{\mathbb{Z}}}^* H_{\alpha, \beta} u_+^{\alpha} u_-^{\beta}
$$

where we denote by  $\sum^*$  the sum restricted to those  $\alpha, \beta : |\alpha| = |\beta| < \infty$ . One can see that

$$
\frac{\partial}{\partial \bar{u}_j} H(u) = \frac{\partial \widehat{H}(u_+, u_-)}{\partial u_{-,j}} \Big|_{u_+ = \bar{u}_- = u}.
$$

**Poisson structure and hamiltonian flows.** The scale  $\{\mathcal{H}_{r,n}(\mathbf{h}_{\mathbf{w}})\}_{r>0}$  is a Banach-Poisson algebra in the following sense

$$
\langle u, v \rangle = 2\text{Re}(u, v), \quad \omega(u, v) = 2\text{Im}(u, v).
$$

<span id="page-12-2"></span><sup>&</sup>lt;sup>10</sup> We recall that given a complex Hilbert space *H* with a Hermitian product  $(·, ·)$ , its realification is a real symplectic Hilbert space with scalar product and symplectic form given by

**Proposition 2.1.** *For*  $0 < \rho < r$  *and*  $\eta > 0$  *we have* 

<span id="page-13-8"></span>
$$
|\{F, G\}|_{r, \eta, w} \le 4\left(1 + \frac{r}{\rho}\right)|F|_{r + \rho, \eta, w}|G|_{r + \rho, \eta, w}.
$$
\n(2.1)

*Proof.* It is essentially contained in [\[BBP13\]](#page-63-27). See in particular Lemma 2.16 of [BBP13] with  $n = 0$  (no action variables here) and no *s* and  $s'$  (no actions variable here). Note that the constant in Lemma 2.16 is 8, instead of 4 in the present paper, because of the presence there of action variables which scale different from the cartesian ones (namely  $(2r)^2$  instead of 2*r*). Recall also the required properties of the space *E* (named  $h_w$  in the present paper) mentioned after Definition 2.5.  $\Box$ 

<span id="page-13-1"></span>The following Lemma is a simple corollary and its proof is postponed to the appendix. **Lemma 2.1** (Hamiltonian flow). Let  $0 < \rho < r$ , and  $S \in \mathcal{H}_{r+\rho, n}(\mathbf{h}_{\mathbf{w}})$  with

<span id="page-13-2"></span>
$$
|S|_{r+\rho,\eta,w} \le \delta := \frac{\rho}{8e(r+\rho)}.\tag{2.2}
$$

*Then the time* 1*-Hamiltonian flow*  $\Phi_S^1$  :  $B_r(h_w) \to B_{r+\rho}(h_w)$  *is well defined, analytic, symplectic with*

<span id="page-13-4"></span>
$$
\sup_{u \in B_r(\mathrm{h}_w)} \left| \Phi_S^1(u) - u \right|_{\mathrm{h}_w} \le (r + \rho) |S|_{r + \rho, \eta, w} \le \frac{\rho}{8e}.
$$
\n(2.3)

*For any*  $H \in \mathcal{H}_{r+\rho,\eta}(\mathbf{h}_{\mathbf{w}})$  *we have that*  $H \circ \Phi_S^1 = e^{\{S,\cdot\}}H \in \mathcal{H}_{r,\eta}(\mathbf{h}_{\mathbf{w}})$  *and* 

<span id="page-13-7"></span><span id="page-13-6"></span><span id="page-13-5"></span>
$$
\left| e^{\{S,\cdot\}} H \right|_{r,\eta,\mathbf{w}} \le 2|H|_{r+\rho,\eta,\mathbf{w}},\tag{2.4}
$$

$$
\left| \left( e^{\{S,\cdot\}} - \mathrm{id} \right) H \right|_{r,\eta,\mathbf{w}} \le \delta^{-1} |S|_{r+\rho,\eta,\mathbf{w}} |H|_{r+\rho,\eta,\mathbf{w}},\tag{2.5}
$$

$$
\left| \left( e^{\{S,\cdot\}} - \mathrm{id} - \{S,\cdot\} \right) H \right|_{r,\eta,\mathrm{w}} \le \frac{1}{2} \delta^{-2} |S|_{r+\rho,\eta,\mathrm{w}}^2 |H|_{r+\rho,\eta,\mathrm{w}} \tag{2.6}
$$

*More generally for any h* ∈ N *and any sequence*  $(c_k)_{k \in \mathbb{N}}$  *with*  $|c_k| \leq 1/k!$ *, we have* 

<span id="page-13-3"></span>
$$
\left|\sum_{k\geq h} c_k \operatorname{ad}_S^k(H)\right|_{r,\eta,\mathbf{w}} \leq 2|H|_{r+\rho,\eta,\mathbf{w}}\big(|S|_{r+\rho,\eta,\mathbf{w}}/2\delta\big)^h,\tag{2.7}
$$

*where*  $\text{ad}_S(\cdot) := \{S, \cdot\}.$ 

# <span id="page-13-0"></span>**3. Immersions for Spaces of Hamiltonians**

Given two positive sequences  $w = (w_j)_{j \in \mathbb{Z}}$ ,  $w' = (w'_j)$ we write that  $w \leq w'$  if the inequality holds point wise, namely

$$
\mathbf{w} \leq \mathbf{w}' \quad : \iff \quad \mathbf{w}_j \leq \mathbf{w}'_j, \quad \forall j \in \mathbb{Z}.
$$

In this way if  $r' \le r$  and  $w \le w'$  then  $B_{r'}(h_{w'}) \subseteq B_r(h_w)$ . Consequently if  $r' \le r, \eta' \le \eta$ and  $w \leq w'$  then  $\mathcal{A}_{r,\eta}(\mathbf{h}_w) \subseteq \mathcal{A}_{r',\eta'}(\mathbf{h}_{w'}).$ 

We thus wish to study conditions on  $(r, \eta, w)$ ,  $(r^*, \eta', w')$  (with  $r^* \le r$ ) which ensure that  $\mathcal{H}_{r,n}(\mathbf{h}_{\mathbf{w}}) \subseteq \mathcal{H}_{r^*,n'}(\mathbf{h}_{\mathbf{w}'})$ . Note that this is not obvious at all, since we are asking that an Hamiltonian vector field of  $X_H \in \mathcal{H}_{r,n}(\mathbf{h}_w)$ , when restricted to the smaller domain  $B_{r^*}(\mathbf{h}_{w'})$  belongs to the smaller space  $\mathbf{h}_{w'}$ .

**The coefficients**  $c_{r,\eta,w}^{(j)}(\alpha,\beta)$ . Let us start by rewriting the norm  $|\cdot|_{r,\eta,w}$  in a more *adimensional* way. In this way all the dependence on the parameters *r*, η, w of the norm  $|\cdot|_{r,n,w}$  is *encoded* in the coefficients [\(1.27\)](#page-9-3).

**Definition 3.1.** For any  $H \in \mathcal{H}_{r,n}(\mathbf{h}_{w})$  we define a map

$$
B_1(\ell^2) \to \ell^2
$$
,  $y = (y_j)_{j \in \mathbb{Z}} \mapsto (Y_H^{(j)}(y; r, \eta, w))_{j \in \mathbb{Z}}$ 

by setting

<span id="page-14-1"></span>
$$
Y_H^{(j)}(y; r, \eta, w) := \sum_* |H_{\alpha, \beta}| \frac{(\alpha_j + \beta_j)}{2} c_{r, \eta, w}^{(j)}(\alpha, \beta) y^{\alpha + \beta - e_j}
$$
(3.1)

<span id="page-14-2"></span>.

where  $e_j$  is the *j*-th basis vector in  $\mathbb{N}^{\mathbb{Z}}$ , while the coefficient  $c_{r,\eta,\mathbf{w}}^{(j)}(\boldsymbol{\alpha},\boldsymbol{\beta})$  was defined in [\(1.27\)](#page-9-3). For brevity, we set

$$
\sum_*:=\sum_{\alpha,\beta:|\alpha|=|\beta|}
$$

The momentum  $\pi(\cdot)$  was defined in [\(1.25\)](#page-8-3).

The vector field  $Y_H$  is a majorant analytic function on  $\ell^2$  which has the *same norm as H*. Since the majorant analytic functions on a given space have a natural ordering this gives us a natural criterion for immersions, as formalized in the following Lemma.

**Lemma 3.1.** *Let r, r*<sup>\*</sup> > 0,  $\eta$ ,  $\eta' \ge 0$ , w, w'  $\in \mathbb{R}^{\mathbb{Z}}_+$ *. The following properties hold.* 

*(i) The norm of H can be expressed as*

$$
|H|_{r,\eta,w} = \sup_{|y|_{\ell^2} \le 1} |Y_H(y; r, \eta, w)|_{\ell^2}
$$
 (3.2)

*(ii) Given*  $H^{(1)}$  ∈  $\mathcal{H}_{r^*,\eta',w'}$  and  $H^{(2)}$  ∈  $\mathcal{H}_{r,\eta,w}$ , *such that for all*  $\alpha, \beta \in \mathbb{N}^{\mathbb{Z}}$  *and*  $j \in \mathbb{Z}$  *with*  $\alpha_j + \beta_j \neq 0$  *one has* 

$$
|H_{\boldsymbol{\alpha},\boldsymbol{\beta}}^{(1)}|c_{r^*,\eta',w'}^{(j)}(\boldsymbol{\alpha},\boldsymbol{\beta})\leq c|H_{\boldsymbol{\alpha},\boldsymbol{\beta}}^{(2)}|c_{r,\eta,w}^{(j)}(\boldsymbol{\alpha},\boldsymbol{\beta}),
$$

*for some*  $c > 0$ *, then* 

$$
|H^{(1)}|_{r^*,\eta',w'} \le c |H^{(2)}|_{r,\eta,w}.
$$

*Proof.* See "Appendix B". □

<span id="page-14-0"></span>As a corollary we get the following "immersion theorem" for spaces of Hamiltonians

**Proposition 3.1** *(Immersion). Let*  $r, r^* > 0, \eta, \eta' \ge 0, w, w' \in \mathbb{R}_+^{\mathbb{Z}}$ . *If* 

<span id="page-15-4"></span>
$$
C := \sup_{\substack{j \in \mathbb{Z}, \boldsymbol{\alpha}, \boldsymbol{\beta} \in \mathbb{N}^{\mathbb{Z}} \\ \boldsymbol{\alpha}_j + \boldsymbol{\beta}_j \neq 0}} \frac{c_{r^*, \eta', w'}^{(j)}(\boldsymbol{\alpha}, \boldsymbol{\beta})}{c_{r, \eta, w}^{(j)}(\boldsymbol{\alpha}, \boldsymbol{\beta})} < \infty,
$$
\n(3.3)

*then*  $\mathcal{H}_{r,n}(\mathbf{h}_{\mathbf{w}}) \subseteq \mathcal{H}_{r^*,n'}(\mathbf{h}_{\mathbf{w}'})$ *, with* 

<span id="page-15-1"></span>
$$
|H|_{r^*, \eta', w'} \le C |H|_{r, \eta, w}.\tag{3.4}
$$

*In particular*  $|\cdot|_{r,n,w}$  *is increasing in r and η, namely if*  $r^* \leq r$  *and*  $\eta' \leq \eta$  *then* 

$$
|H|_{r^*,\eta',\mathbf{w}} \leq |H|_{r,\eta,\mathbf{w}}.
$$

*Moreover, if*  $r^* \le r$ ,  $w \le w'$  *and*  $H \in \mathcal{K}_{r,n}(\mathbf{h}_w)$  *then* 

<span id="page-15-2"></span>
$$
|H|_{r^*, \eta', \mathbf{w}'} \le |H|_{r, \eta, \mathbf{w}}.\tag{3.5}
$$

*Furthermore, if H preserves momentum then*

$$
|H|_{r^*, \eta', \mathbf{w}'} \le C_0 |H|_{r, \eta, \mathbf{w}},\tag{3.6}
$$

*where*

$$
C_0 := \sup_{\substack{j \in \mathbb{Z}, \, \boldsymbol{\alpha}, \boldsymbol{\beta} \in \mathbb{N}^{\mathbb{Z}}, \\ \boldsymbol{\alpha}_j + \boldsymbol{\beta}_j \neq 0, \\ \sum_i i(\boldsymbol{\alpha}_i - \boldsymbol{\beta}_i) = 0}} \frac{c_{r^*, \eta', w'}^{(j)}(\boldsymbol{\alpha}, \boldsymbol{\beta})}{c_{r, \eta, w}^{(j)}(\boldsymbol{\alpha}, \boldsymbol{\beta})} < \infty, \tag{3.7}
$$

*Proof.* Inequality [\(3.4\)](#page-15-1) directly follows from Lemma [3.1](#page-14-2) (ii), while [\(3.5\)](#page-15-2) follows directly by [\(1.27\)](#page-9-3) since in the kernel  $\alpha_j + \beta_j \neq 0$  implies  $\alpha_j + \beta_j \geq 2$ . The momentum preserving case follows analogously.  $\Box$ case follows analogously.

*Remark 3.1.* The above immersion properties, with different norm and in a different context, were implicitly used by Bourgain in [\[Bou05](#page-63-17)].

# <span id="page-15-0"></span>**4. Small Divisors and Homological Equation**

Let us consider the set of frequencies

$$
\Omega_q := \left\{ \omega = (\omega_j)_{j \in \mathbb{Z}} \in \mathbb{R}^{\mathbb{Z}}, \quad \sup_j |\omega_j - j^2| \langle j \rangle^q < 1/2 \right\};\tag{4.1}
$$

this set is isomorphic to  $[-1/2, 1/2]^{\mathbb{Z}}$  via the identification

$$
\xi \mapsto \omega(\xi)
$$
, where  $\omega_j(\xi) = j^2 + \frac{\xi_j}{\langle j \rangle^q}$ . (4.2)

We endow  $\Omega_a$  with the probability measure  $\mu$  induced<sup>11</sup> by the product measure on  $[-1/2, 1/2]^{\mathbb{Z}}$ .

We now define the set of Diophantine frequencies, the following definition is a slight generalization of the one given by Bourgain in [\[Bou05\]](#page-63-17).

<span id="page-15-3"></span><sup>&</sup>lt;sup>11</sup> Denoting by  $\mu$  the measure in  $\Omega_q$  and by  $\nu$  the product measure on  $[-1/2, 1/2]^{\mathbb{Z}}$ , then  $\mu(A)$  $\nu(\omega^{(-1)}(A))$  for all sets  $A \subset \Omega_q$  such that  $\omega^{(-1)}(A)$  is v-measurable.

**Definition 4.1.** Given  $\gamma > 0$  and  $q \ge 0$ , we denote by  $D_{\gamma,q} \equiv D_{\gamma,q}^{\mu_1,\mu_2}$  the set of  $\mu_1, \mu_2, \gamma$ -Diophantine frequencies

$$
D_{\gamma,q}^{\mu_1,\mu_2} :=
$$
\n
$$
\left\{\omega \in \Omega_q : |\omega \cdot \ell| > \gamma \prod_{n \in \mathbb{Z}} \frac{1}{(1 + |\ell_n|^{\mu_1} \langle n \rangle^{\mu_2 + q})}, \quad \forall \ell \in \mathbb{Z}^{\mathbb{Z}} : 0 < |\ell| < \infty \right\}.
$$
\n(4.3)

Now we have that

**Lemma 4.1.** *For*  $\mu_1, \mu_2 > 1$  *the exists a positive constant*  $C_{\text{meas}}(\mu_1, \mu_2)$  *such that* 

<span id="page-16-0"></span>
$$
\mu\big(\Omega_q\setminus D^{\mu_1,\mu_2}_{\gamma,q}\big)\leq C_{\text{meas}}(\mu_1,\mu_2)\gamma.
$$

*Proof.* In "Appendix C" □

This means that, for all  $\mu_1, \mu_2 > 1$ , Diophantine frequencies are typical in  $\Omega_q$  in the sense that they have full measure. Here and in the following we shall always assume that

<span id="page-16-4"></span>
$$
0 < \gamma \le 1, \qquad \omega \in \mathbb{D}^{2,2}_{\gamma,q} = \mathbb{D}_{\gamma,q}.\tag{4.4}
$$

In the remaining part of this section, on appropriate source and target spaces, we will study the invertibility of the "Lie derivative" operator

$$
L_{\omega}: H \mapsto L_{\omega}H := \sum_{*} \mathbf{i}(\omega \cdot (\boldsymbol{\alpha} - \boldsymbol{\beta}))H_{\boldsymbol{\alpha},\boldsymbol{\beta}}u^{\boldsymbol{\alpha}}\bar{u}^{\boldsymbol{\beta}}, \qquad (4.5)
$$

which is nothing but the action of the Poisson bracket  $\left\{\sum_j \omega_j | u_j \right\}$  $\left\{ \frac{2}{n}, \cdot \right\}$  on *H*.

Recalling the definition of  $K_r(h_w)$  in [\(1.30\)](#page-10-4) we give the following

# **Definition 4.2.** Let

$$
\mathcal{R}_{r,\eta}(\mathbf{h}_{\mathbf{w}}) := \{ H \in \mathcal{H}_{r,\eta}(\mathbf{h}_{\mathbf{w}}) \mid H = \sum_{\alpha \neq \beta} H_{\alpha,\beta} u^{\alpha} \bar{u}^{\beta} \}.
$$
 (4.6)

Then we have the decomposition  $\mathcal{H}_{r,n}(\mathbf{h}_w) = \mathcal{R}_{r,n}(\mathbf{h}_w) \oplus \mathcal{K}_r(\mathbf{h}_w)$  and the continuous projections $12$ 

<span id="page-16-3"></span>
$$
|\Pi_{\mathcal{K}}H|_{r,\eta,\mathbf{w}}, |\Pi_{\mathcal{R}}H|_{r,\eta,\mathbf{w}} \leq |H|_{r,\eta,\mathbf{w}}.
$$
\n(4.7)

Obviously for diophantine frequency  $\mathcal{R}_{r,n}(\mathbf{h}_{\mathbf{w}})$  and  $\mathcal{K}_{r}(\mathbf{h}_{\mathbf{w}})$  represent the range and kernel of *L*ω.

For any  $\overline{r}$ ,  $\eta$ , w and  $\alpha$ ,  $\beta \in \mathbb{N}^{\mathbb{Z}}$  recall the coefficient defined in [\(1.27\)](#page-9-3)

$$
c_{r,\eta,\mathsf{w}}^{(j)}(\boldsymbol{\alpha},\boldsymbol{\beta}):=r^{|\boldsymbol{\alpha}|+|\boldsymbol{\beta}|-2}e^{\eta|\pi(\boldsymbol{\alpha}-\boldsymbol{\beta})|}\frac{\mathsf{w}_j^2}{\mathsf{w}^{\boldsymbol{\alpha}+\boldsymbol{\beta}}}.
$$

<span id="page-16-1"></span>In following Lemma we consider  $R \in \mathcal{R}_{r,n}(\mathbf{h}_{\mathbf{w}})$  and state sufficient conditions which ensure that  $L_{\omega}^{-1}R \in \mathcal{R}_{r^*,\eta'}(\mathbf{h}_{w'}).$ 

<span id="page-16-2"></span><sup>12</sup> Explicitely  $\Pi_{\mathcal{K}} H := \sum_{\alpha=\beta} H_{\alpha,\beta} u^{\alpha} \bar{u}^{\beta}, \ \Pi_{\mathcal{R}} H := \sum_{\alpha \neq \beta} H_{\alpha,\beta} u^{\alpha} \bar{u}^{\beta}.$ 

**Lemma 4.2** (Homological equation). *Fix*  $\omega \in D_{\gamma}$  *q. Consider two ordered weights* 0 <  $r^* \le r, 0 \le \eta' \le \eta, w' \ge w$ , *such that* 

<span id="page-17-1"></span>
$$
K := \gamma \sup_{\substack{j \in \mathbb{Z}, \, \boldsymbol{\alpha} \neq \boldsymbol{\beta} \in \mathbb{N}^{\mathbb{Z}} \\ \boldsymbol{\alpha}_j + \boldsymbol{\beta}_j \neq 0}} \frac{c_{r^*, \eta', w'}^{(j)}(\boldsymbol{\alpha}, \boldsymbol{\beta})}{c_{r, \eta, w}^{(j)}(\boldsymbol{\alpha}, \boldsymbol{\beta}) |\omega \cdot (\boldsymbol{\alpha} - \boldsymbol{\beta})|} < \infty,
$$
\n(4.8)

*then for any*  $R \in \mathcal{R}_{r,n}(\mathbf{h}_{\mathbf{w}})$  *the homological equation* 

$$
L_{\omega}S=R
$$

*has a unique solution*  $S = L_{\omega}^{-1}R$  *in*  $\mathcal{R}_{r^*,\eta'}(\mathbf{h}_{\mathsf{w}'})$ *, which satisfies* 

<span id="page-17-2"></span>
$$
\left| L_{\omega}^{-1} R \right|_{r^*, \eta', \mathbf{w}'} \leq \gamma^{-1} K |R|_{r, \eta, \mathbf{w}}.
$$
\n(4.9)

*Similarly, if R preserves momentum, assuming only*

<span id="page-17-4"></span>
$$
K_0 := \gamma \sup_{\substack{j \in \mathbb{Z}, \, \boldsymbol{\alpha} \neq \boldsymbol{\beta} \in \mathbb{N}^{\mathbb{Z}} \\ \boldsymbol{\alpha}_j + \boldsymbol{\beta}_j \neq 0}} \frac{c_{r^*, \eta', \omega'}^{(j)}(\boldsymbol{\alpha}, \boldsymbol{\beta})}{c_{r, \eta, \omega}^{(j)}(\boldsymbol{\alpha}, \boldsymbol{\beta})|\omega \cdot (\boldsymbol{\alpha} - \boldsymbol{\beta})|} < \infty, \tag{4.10}
$$
\n
$$
\sum_{i} i(\boldsymbol{\alpha}_i - \boldsymbol{\beta}_i) = 0
$$

*we have that S also preserves momentum and*

$$
\left| L_{\omega}^{-1} R \right|_{r^*, \eta', \mathbf{w}'} \le \gamma^{-1} K_0 |R|_{r, \eta, \mathbf{w}}.
$$
\n(4.11)

*Proof.* Given any Hamiltonian  $R \in \mathcal{R}$ , the formal solution of  $L_S = R$  is given by

$$
L_{\omega}^{-1}R = \sum_{|\alpha| = |\beta|, \alpha \neq \beta} \frac{1}{i(\omega \cdot (\alpha - \beta))} R_{\alpha,\beta} u^{\alpha} \bar{u}^{\beta},
$$
(4.12)

where *u* ∈ *B*<sub>*r*<sup>∗</sup>(h<sub>w</sub>')</sub>. By Lemma [3.1](#page-14-2) (ii) (applied to  $H^{(1)} = L_{\omega}^{-1}R$  and  $H^{(2)} = R$ ) and  $(4.8)$ , we get  $(4.9)$ . The momentum preserving case is analogous.  $\square$ 

#### <span id="page-17-0"></span>**5. Abstract Birkhoff Normal Form**

In this section we prove the abstract Birkoff normal form Theorem [1.3.](#page-10-0) We start by defining a degree decomposition which endows  $\mathcal{H}_{r,\eta}(\mathbf{h}_{w})$  with a graded Poisson algebra structure.

**Definition 5.1** (minimal scaling degree). We say that *H* has *minimal* scaling degree  $d = d(H)$  (at zero) if

<span id="page-17-3"></span>
$$
H_{\alpha,\beta} = 0, \quad \forall \alpha, \beta : \quad |\alpha| = |\beta| \leq d,
$$
  

$$
H_{\alpha,\beta} \neq 0, \quad \text{for some } \alpha, \beta : \quad |\alpha| = |\beta| = d+1.
$$

We say that  $d(0) = +\infty$ .

Essentially *H* has scaling degree d if and only if it has a zero of order  $2d + 2$  at zero, we prefer this notation because we find it more intrinsic, it produces a graded Poisson algebra structure and one has the following

**Lemma 5.1** *If*  $H \in \mathcal{H}_{r,n}(\mathbf{h}_{\mathbf{w}})$  *with*  $d(H) \geq d$ *, then for all*  $r^* \leq r$  *one has* 

$$
|H|_{r^*, \eta, w} \le \left(\frac{r^*}{r}\right)^{2d} |H|_{r, \eta, w}.
$$

*Proof.* Recalling  $(1.27)$ , we have

$$
\frac{c_{r^*,\eta,\mathsf{w}}^{(j)}(\boldsymbol{\alpha},\boldsymbol{\beta})}{c_{r,\eta,\mathsf{w}}^{(j)}(\boldsymbol{\alpha},\boldsymbol{\beta})}=\left(\frac{r^*}{r}\right)^{|\boldsymbol{\alpha}|+|\boldsymbol{\beta}|-2}.
$$

Since  $|\alpha| + |\beta| - 2 \ge 2d$ , the inequality follows by Proposition [3.1.](#page-14-0)

The normal form will be proved iteratively by means of the following Lemma, which constitutes the main step of the procedure.

Basically we start with a Hamiltonian  $H = D_{\omega} + Z + R$  with  $Z \in \mathcal{K}_r(\mathbf{h}_{\omega})$  in normal form and  $R \in \mathcal{R}_{r,n}(\mathbf{h}_{\mathbf{w}})$  of minimal degree d, and we consider  $r' \leq r, \eta' \leq \eta, \mathbf{w}' \geq \mathbf{w}$ so that  $\mathcal{H}_{r,\eta}(\mathbf{h}_{w}) \subseteq \mathcal{H}_{r',\eta'}(\mathbf{h}_{w'})$ . Then we give a sufficient condition which ensures the existence of a change of variables  $\Phi : B_{r'}(h_{w'}) \rightarrow B_r(h_{w'})$  such that

$$
H\circ\Phi=D_{\omega}+Z'+R',
$$

<span id="page-18-0"></span>with  $Z'$ ,  $R' \in \mathcal{H}_{r',\eta'}(\mathbf{h}_{w'})$  and  $R'$  of minimal degree  $d+1$ .

**Lemma 5.2.** *Fix*  $\omega \in D_{\gamma,q}$ . *Let*  $r > r' > 0$ ,  $\eta \geq \eta' \geq 0$ ,  $\omega \leq \omega'$ . *Consider* 

$$
H = D_{\omega} + Z + R, \quad Z \in \mathcal{K}_r(\mathbf{h}_{\mathbf{w}}), \quad R \in \mathcal{R}_{r,\eta}(\mathbf{h}_{\mathbf{w}}), \quad d(Z) \ge 1, \quad d(R) \ge d \ge 1.
$$

*Assume that* [\(3.3\)](#page-15-4) *and* [\(4.8\)](#page-17-1) *hold and that*[13](#page-18-1)

<span id="page-18-3"></span>
$$
|R|_{r,\eta,\mathbf{w}} \le \frac{\gamma \delta}{K}, \quad \text{with} \quad \delta := \frac{r - r'}{16er}.
$$

*Then there exists a change of variables*

<span id="page-18-6"></span>
$$
\Phi : B_{r'}(\mathbf{h}_{w'}) \rightarrow B_r(\mathbf{h}_{w'}), \qquad (5.2)
$$

*such that*

$$
H \circ \Phi = D_{\omega} + Z' + R', \quad Z' \in \mathcal{K}_{r', \eta'}(\mathbf{h}_{\mathbf{w}'}), R' \in \mathcal{R}_{r', \eta'}(\mathbf{h}_{\mathbf{w}'}), \quad d(Z') \ge 1, \quad d(R') \ge d+1.
$$

*Moreover*[14](#page-18-2)

<span id="page-18-4"></span>
$$
|Z'|_{r',\eta',\mathbf{w}'} \leq |Z|_{r,\eta,\mathbf{w}} + (\gamma \delta)^{-1} K |R|_{r,\eta,\mathbf{w}} (C|R|_{r,\eta,\mathbf{w}} + |Z|_{r,\eta,\mathbf{w}}),
$$
  

$$
|R'|_{r',\eta',\mathbf{w}'} \leq (\gamma \delta)^{-1} K |R|_{r,\eta,\mathbf{w}} (C|R|_{r,\eta,\mathbf{w}} + |Z|_{r,\eta,\mathbf{w}}).
$$
 (5.3)

*Finally, for*  $w^{\sharp} \geq w'$ , *assume the further conditions* 

<span id="page-18-5"></span>
$$
\gamma \sup_{\substack{j\in\mathbb{Z}, \, \boldsymbol{\alpha}\neq\boldsymbol{\beta}\in\mathbb{N}^{\mathbb{Z}} \\ \boldsymbol{\alpha}_j+\boldsymbol{\beta}_j\neq 0}} \frac{c_{r^*,\eta',\text{w}^\sharp}^{(j)}(\boldsymbol{\alpha},\boldsymbol{\beta})}{c_{r,\eta,\text{w}}^{(j)}(\boldsymbol{\alpha},\boldsymbol{\beta})|\omega\cdot(\boldsymbol{\alpha}-\boldsymbol{\beta})|} =: K^\sharp < \infty, \qquad r^* := \frac{r'+r}{2} \tag{5.4}
$$

<sup>13</sup> *K* is the constant in  $(4.8)$ .

<span id="page-18-2"></span><span id="page-18-1"></span><sup>14</sup> *C* is defined in  $(3.3)$ .

*and*

<span id="page-19-0"></span>
$$
|R|_{r,\eta,\mathbf{w}} \le \frac{\gamma \delta}{K^{\sharp}}.\tag{5.5}
$$

*Then*

<span id="page-19-1"></span>
$$
\Phi|_{B_{r'}(\mathbf{h}_{\mathbf{w}^{\sharp}})} : B_{r'}(\mathbf{h}_{\mathbf{w}^{\sharp}}) \to B_r(\mathbf{h}_{\mathbf{w}^{\sharp}}),
$$
  
\n
$$
\sup_{u \in B_{r'}(\mathbf{h}_{\mathbf{w}^{\sharp}})} |\Phi(u) - u|_{\mathbf{h}_{\mathbf{w}^{\sharp}}} \leq r\gamma^{-1} K^{\sharp} |R|_{r, \eta, \mathbf{w}}.
$$
\n(5.6)

*Moreover if R preserves momentum, assuming only that*

$$
K_0^{\sharp} := \gamma \sup_{\substack{j \in \mathbb{Z}, \, \boldsymbol{\alpha} \neq \boldsymbol{\beta} \in \mathbb{N}^{\mathbb{Z}} \\ \alpha_j + \beta_j \neq 0, \\ \sum_i i(\boldsymbol{\alpha}_i - \boldsymbol{\beta}_i) = 0}} \frac{c_{r^*, \eta', \text{w}^{\sharp}}^{(j)}(\boldsymbol{\alpha}, \boldsymbol{\beta})}{c_{r, \eta, \text{w}}^{(j)}(\boldsymbol{\alpha}, \boldsymbol{\beta})|\omega \cdot (\boldsymbol{\alpha} - \boldsymbol{\beta})|} < \infty \tag{5.7}
$$

*and that* [\(5.1\)](#page-18-3), [\(5.5\)](#page-19-0) *hold with*  $K_0$ ,  $K_0^{\sharp}$  *instead of*  $K$ ,  $K^{\sharp}$  *we have that*  $R'$  *preserves momentum and* [\(5.6\)](#page-19-1) *holds with*  $K_0^{\sharp}$  *instead of*  $K^{\sharp}$ .

*Proof.* By Lemma [4.2](#page-16-1) let  $S = L_{\omega}^{-1} R$  in  $\mathcal{R}_{r^*,\eta'}(\mathbf{h}_{\mathbf{w}'})$  be the unique solution of the homological equation  $L_{\omega} S = R$  on  $B_{r^*}(\mathbf{h}_{w'})$ . Note that  $d(S) \ge d$ . We have

<span id="page-19-2"></span>
$$
|S|_{r^*,\eta',\mathbf{w}'} \le \gamma^{-1} K |R|_{r,\eta,\mathbf{w}}.\tag{5.8}
$$

We now apply Lemma [2.1](#page-13-1) with  $(r, \eta, w) \rightsquigarrow (r', \eta', w')$  and  $\rho := r^* - r'$ . Note that [\(5.1\)](#page-18-3) and [\(5.8\)](#page-19-2) imply [\(2.2\)](#page-13-2). We define  $\Phi := \Phi_S^1$  and compute

$$
H' := H \circ \Phi = D_{\omega} + Z + (e^{\{S,\cdot\}} - \mathrm{id} - \{S,\cdot\})D_{\omega} + (e^{\{S,\cdot\}} - \mathrm{id})(Z + R)
$$
  
=  $D_{\omega} + Z - \sum_{j=2}^{\infty} \frac{(\mathrm{ad}S)^{j-1}}{j!}R + (e^{\{S,\cdot\}} - \mathrm{id})(Z + R).$ 

We now set

$$
Z' = \Pi_{\mathcal{K}} H' - D_{\omega}, \quad R' = \Pi_{\mathcal{R}} H'.
$$

Since the scaling degree is additive w.r.t. Poisson brackets, we have that  $d(Z') \ge 1$  and  $d(R') \geq d + 1$ . By [\(2.7\)](#page-13-3)

$$
|Z'|_{r',\eta',\mathbf{w}'} \leq |Z|_{r',\eta',\mathbf{w}'} + (\gamma \delta)^{-1} K |R|_{r,\eta,\mathbf{w}} (|R|_{r^*,\eta',\mathbf{w}'} + |Z|_{r^*,\eta',\mathbf{w}'}),
$$
  

$$
|R'|_{r',\eta',\mathbf{w}'} \leq (\gamma \delta)^{-1} K |R|_{r,\eta,\mathbf{w}} (|R|_{r^*,\eta',\mathbf{w}'} + |Z|_{r^*,\eta',\mathbf{w}'}).
$$

Since  $(4.8)$  holds we can apply Proposition [3.1:](#page-14-0) by  $(3.4)$  and  $(3.5)$  we get

$$
|R|_{r^*,\eta',w'}\leq C|R|_{r,\eta,w},\qquad |Z|_{r^*,\eta',w'}\leq |Z|_{r,\eta,w}.
$$

 $(5.3)$  follows.

Finally assume [\(5.5\)](#page-19-0) and [\(5.4\)](#page-18-5). By Lemma [4.2](#page-16-1) let  $S^{\sharp} = L_{\omega}^{-1} R$  in  $\mathcal{R}_{r^*,\eta'}(\mathbf{h}_{w^{\sharp}})$  be the solution of the homological equation  $L_{\omega}S^{\sharp} = R$  on  $B_{r^*}(\mathbf{h}_{\omega}) \subset B_{r^*}(\mathbf{h}_{\omega})$ . Since *S* and *S*<sup>‡</sup> solve the same linear equation on  $B_r$ <sup>∗</sup> ( $h_{w<sup>\sharp</sup>}$ ), we have that

$$
S^{\sharp}=S_{\left|B_{r^{*}}(\mathbf{h}_{\mathbf{w}^{\sharp}})\right|}.
$$

By  $(4.9)$  we get

<span id="page-20-0"></span>
$$
|S|_{r^*,\eta',w^\sharp} \le \gamma^{-1} K^\sharp |R|_{r,\eta,w}.\tag{5.9}
$$

We now apply Lemma [2.1](#page-13-1) with  $(r, \eta, w) \rightsquigarrow (r', \eta', w^{\sharp})$  and  $\rho := r^* - r'$ . Note that [\(5.5\)](#page-19-0) and  $(5.9)$  imply  $(2.2)$ . Then  $(5.6)$  follows by  $(2.3)$  and  $(5.9)$ .

The momentum preserving case is analogous.

Theorem [1.3](#page-10-0) follows Given  $\eta \ge 0$  and a sequence of weights  $w_0 \le w_1 \le \cdots \le w_N =$ w. For any given  $r > 0$  we set

<span id="page-20-5"></span>
$$
r_n = (2 - \frac{n}{N})r, \quad \eta_n = (1 - \frac{n}{N})\eta, \quad 0 \le n \le N, \qquad r_n^* = \frac{r_{n+1} + r_n}{2}, \quad 0 \le n < N. \tag{5.10}
$$

From Assumption [1](#page-10-3) and  $(1.27)$  we have<sup>15</sup>

<span id="page-20-2"></span>
$$
\max\left\{1,\sup_{0\leq n<\mathbb{N}}\sup_{\substack{j,\boldsymbol{\alpha},\boldsymbol{\beta}\\\boldsymbol{\alpha}_j+\boldsymbol{\beta}_j\neq 0}}\frac{c_{r_n^*,\eta_{n+1},w_{n+1}}^{(j)}(\boldsymbol{\alpha},\boldsymbol{\beta})}{c_{r_n,\eta_n,w_n}^{(j)}(\boldsymbol{\alpha},\boldsymbol{\beta})}\right\}=\mathfrak{C}<\infty,
$$
\n(5.11)

$$
\max\left\{1,\sup_{0\leq n<\mathbb{N}}\sup_{\substack{j,\boldsymbol{\alpha},\boldsymbol{\beta}\\\boldsymbol{\alpha}_j+\boldsymbol{\beta}_j\neq 0}}\frac{c_{r^*_{j^*},\eta_{n+1},\mathbb{W}_{n+1}}^{(j)}(\boldsymbol{\alpha},\boldsymbol{\beta})}{c_{r_n,\eta_n,\mathbb{W}_n}^{(j)}(\boldsymbol{\alpha},\boldsymbol{\beta})|\boldsymbol{\omega}\cdot(\boldsymbol{\alpha}-\boldsymbol{\beta})|}\right\}=\mathfrak{K}<\infty,
$$
(5.12)

$$
\max\left\{1,\sup_{0\leq n<\mathbb{N}}\sup_{\substack{j,\boldsymbol{\alpha},\boldsymbol{\beta}\\\boldsymbol{\alpha}_j+\boldsymbol{\beta}_j\neq 0}}\frac{c_{r_n^*,\eta_N,\mathbb{W}_N}^{(J)}(\boldsymbol{\alpha},\boldsymbol{\beta})}{c_{r_n,\eta_n,\mathbb{W}_n}^{(J)}(\boldsymbol{\alpha},\boldsymbol{\beta})|\omega\cdot(\boldsymbol{\alpha}-\boldsymbol{\beta})|}\right\}=\mathfrak{K}^{\sharp}<\infty.
$$
 (5.13)

(*j*)

For brevity we set

$$
h_n := h_{w_n}, \quad \mathcal{H}_n := \mathcal{H}_{r_n, \eta_n}(h_n), \quad 0 \le n \le N, \mathcal{H}_{n,*} := \mathcal{H}_{r_n^*, \eta_{n+1}}(h_{n+1}), \quad 0 \le n < N,
$$
\n
$$
(5.14)
$$

and, correspondingly,  $\mathcal{R}_n$ ,  $\mathcal{K}_n$ ,  $\mathcal{R}_n$ ,  $\mathcal{K}_{n,*}$  and

<span id="page-20-6"></span>
$$
|\cdot|_{n} := |\cdot|_{r_n, \eta_n, w_n}, \qquad |\cdot|_{n,*} := |\cdot|_{r_n^*, \eta_{n+1}, w_{n+1}}.
$$
 (5.15)

**Lemma 5.3.** *By Assumption* [\(5.11\)](#page-20-2) *we have the immersion properties*

<span id="page-20-4"></span>
$$
\mathcal{H}_0 \subseteq \mathcal{H}_{0,*} \subseteq \cdots \subseteq \mathcal{H}_n \subseteq \mathcal{H}_{n,*} \subseteq \mathcal{H}_{n+1} \subseteq \cdots \subseteq \mathcal{H}_N, \tag{5.16}
$$

*with estimates*

<span id="page-20-3"></span>
$$
H \in \mathcal{H}_n \quad \Longrightarrow \quad |H|_{n,*} \leq C|H|_n, \qquad 0 \leq n \leq i \leq N-1
$$
\n
$$
H \in \mathcal{K}_n \quad \Longrightarrow \quad |H|_{n,*} \leq |H|_n, \qquad 0 \leq n \leq i \leq N-1. \tag{5.17}
$$

<span id="page-20-1"></span><sup>&</sup>lt;sup>15</sup> We are just using the fact that the ratio  $c_{r,\eta,w}^{(j)}(\boldsymbol{\alpha},\boldsymbol{\beta})/c_{r',\eta',w'}^{(j)}(\boldsymbol{\alpha},\boldsymbol{\beta})$  depends on *r*, *r'* only through their ratio.

*Proof.* We apply Proposition [3.1](#page-14-0) with

$$
r, \eta, w \rightsquigarrow r_n, \eta_n, w_n, r^*, \eta', w' \rightsquigarrow r_n^*, \eta_{n+1}, w_{n+1},
$$

by noting that the bound  $(3.3)$  follows from  $(5.11)$ . The bounds in  $(5.17)$  follow form  $(3.4)$  and  $(3.5)$ . The chain of inclusions  $(5.16)$  follows.  $\Box$ 

*Proof of Theorem [1.3.](#page-10-0)* We will prove the thesis inductively. Let us start by noticing that

$$
\widehat{r} = \min \left\{ \frac{\bar{r}}{8\sqrt{|G|_{\bar{r}, \eta, w_0}}} \sqrt{\frac{\gamma \hat{\delta}}{\max\{\mathfrak{CR}, \mathfrak{R}^{\sharp}\}}}, \frac{\bar{r}}{2} \right\}, \quad \hat{\delta} := \frac{1}{32eN}
$$

and, for all  $0 < r \leq \hat{r}$ , let us set

$$
\varepsilon := \gamma^{-1} \left(\frac{2r}{\bar{r}}\right)^2 |G|_{\bar{r}, \eta, w_0} = \frac{1}{2^9 e} \left(\frac{r}{r_\star}\right)^2.
$$

From definition  $(1.32)$  we thus deduce that

<span id="page-21-0"></span>
$$
8\,\varepsilon \max\{\mathfrak{CR}, \mathfrak{R}^{\sharp}\}\hat{\delta}^{-1} \le 1. \tag{5.18}
$$

Recalling the notations introduced in  $(5.10)$ – $(5.15)$ , by Lemma  $(5.1)$  we have

$$
\gamma^{-1}|G|_0\leq \varepsilon,
$$

hence, setting  $Z^{(0)} := \Pi_{\mathcal{K}} G$  and  $R^{(0)} := \Pi_{\mathcal{R}} G$ , from [\(4.7\)](#page-16-3) it follows that

$$
\gamma^{-1} |Z^{(0)}|_0, \ \gamma^{-1} |R^{(0)}|_0 \leq \varepsilon.
$$

We perform an iterative procedure producing a sequence of Hamiltonians, for  $n =$  $0,\ldots,N$ 

<span id="page-21-1"></span>
$$
H^{(n)} = D_{\omega} + Z^{(n)} + R^{(n)},
$$
  
\n
$$
Z^{(n)} \in \mathcal{K}_n, \quad R^{(n)} \in \mathcal{R}_n, \quad d(Z^{(n)}) \ge 1, \quad d(R^{(n)}) \ge n+1,
$$
  
\n
$$
\gamma^{-1} |Z^{(n)}|_n \le \varepsilon \sum_{h=0}^n 2^{-h}, \quad \gamma^{-1} |R^{(n)}|_n \le \varepsilon^{n+1} \left( 4 \mathfrak{C} \mathfrak{K} \hat{\delta}^{-1} \right)^{n} \stackrel{(5.18)}{\le} 2^{-n} \varepsilon. \quad (5.19)
$$

Fix any  $k < N$ . Let us assume that we have constructed  $H^{(0)}, \ldots, H^{(k)}$  satisfying [\(5.19\)](#page-21-1) for all  $0 \le n \le k$ . We want to apply Lemma [5.2](#page-18-0) with

$$
H, r, \eta, w \rightsquigarrow H^{(k)}, r_k, \eta_k, w_k \quad \text{and} \quad r', \eta', w', w^{\sharp}, d \rightsquigarrow r_{k+1}, \eta_{k+1}, w_{k+1}, w_{N}, k+1.
$$

By construction the bounds [\(3.3\)](#page-15-4), [\(4.8\)](#page-17-1) and [\(5.4\)](#page-18-5) hold since  $C \leq \mathfrak{C}$ ,  $K \leq \mathfrak{K}$ ,  $K^{\sharp} \leq \mathfrak{K}^{\sharp}$ , where  $\widehat{C}$ ,  $\widehat{K}$ ,  $\widehat{K}^{\sharp}$  were defined in [\(5.11\)](#page-20-2), [\(5.12\)](#page-20-2), [\(5.13\)](#page-20-2). We just have to verify that [\(5.1\)](#page-18-3) holds, namely

$$
|R^{(k)}|_k \leq \frac{\gamma}{\mathfrak{K}} \frac{r_k - r_{k+1}}{16er_k}.
$$

In fact, by applying the inductive hypothesis  $(5.19)$  and the smallness condition  $(5.18)$ , we get

$$
|R^{(k)}|_k \leq \gamma \left(4\mathfrak{C}\mathfrak{K}\hat{\delta}^{-1}\right)^k \varepsilon^{k+1} \leq \frac{\gamma \varepsilon}{2^k} \leq \frac{\gamma}{16\varepsilon \mathfrak{K}(2N-k)} = \frac{\gamma}{\mathfrak{K}} \frac{r_k - r_{k+1}}{16er_k}.
$$

The verification of  $(5.5)$  is completely analogous.

So, by applying Lemma [5.2](#page-18-0) we construct a change of variable  $\Phi_k$  as in [\(5.2\)](#page-18-6) with

$$
\Phi_k : B_{r_{k+1}}(\mathbf{h}_{w_{k+1}}) \to B_{r_k}(\mathbf{h}_{w_{k+1}}).
$$

Let us now set

$$
H^{(k+1)} = D_{\omega} + Z^{(k+1)} + R^{(k+1)} := H_k \circ \Phi_k
$$

with  $Z^{(k+1)} \in \mathcal{K}_{k+1}, R^{(k+1)} \in \mathcal{R}_{k+1}$  and  $d(Z^{(k+1)}) \ge 1$ ,  $d(R^{(k+1)}) \ge k+2$ . It remains to prove the bounds in the second line of  $(5.19)$  (with  $n = k + 1$ ). By  $(5.3)$  we have

$$
|Z^{(k+1)}|_{k+1} \le |Z^{(k)}|_{k} + (\gamma \hat{\delta})^{-1} \mathfrak{K} |R^{(k)}|_{k} (\mathfrak{C} |R^{(k)}|_{k} + |Z^{(k)}|_{k}),
$$
  

$$
|R^{(k+1)}|_{k+1} \le (\gamma \hat{\delta})^{-1} \mathfrak{K} |R^{(k)}|_{k} (\mathfrak{C} |R^{(k)}|_{k} + |Z^{(k)}|_{k}).
$$
 (5.20)

By substituting the inductive hypothesis  $(5.19)$ , we have the following chain of inequalities

$$
\gamma^{-1} |R^{(k+1)}|_{k+1} \leq \hat{\delta}^{-1} \varepsilon^2 \mathfrak{K} (4 \mathfrak{C} \mathfrak{K} \hat{\delta}^{-1} \varepsilon)^k (\mathfrak{C} (4 \mathfrak{C} \mathfrak{K} \hat{\delta}^{-1} \varepsilon)^k + 2)
$$
  
\n
$$
\leq \delta^{-1} \varepsilon^2 \mathfrak{K} (4 \mathfrak{C} \mathfrak{K} \hat{\delta}^{-1} \varepsilon)^k (\mathfrak{C} + 2)
$$
  
\n
$$
\leq (4 \mathfrak{C} \mathfrak{K} \hat{\delta}^{-1})^{k+1} \varepsilon^{k+2} = (4 \mathfrak{C} \mathfrak{K} \hat{\delta}^{-1} \varepsilon)^{k+1} \varepsilon,
$$

which proves the bound on  $R^{(n)}$  in [\(5.19\)](#page-21-1) for any *n*.

En passant, we note that

<span id="page-22-1"></span>
$$
\gamma \varepsilon \left( 4 \mathfrak{C} \mathfrak{K} \hat{\delta}^{-1} \varepsilon \right)^{\mathcal{N}} = \frac{\gamma}{2^9 e r_{\star}^2} \left( \frac{\mathfrak{C} \mathfrak{K} \mathcal{N}}{4 r_{\star}^2} \right)^{\mathcal{N}} r^{2(\mathcal{N}+1)}.
$$
 (5.21)

Finally, using the same strategy as above, we also get

$$
\gamma^{-1}|Z^{(k+1)}|_{k+1} \leq \varepsilon \left(\sum_{h=0}^k 2^{-h} + (4\mathfrak{C}\mathfrak{K}\hat{\delta}^{-1})^{k+1} \varepsilon^{k+1}\right) \stackrel{(5.18)}{\leq} \varepsilon \sum_{h=0}^{k+1} 2^{-h},
$$

which completes the proof of the inductive hypothesis  $(5.19)$ , and remark that

<span id="page-22-2"></span>
$$
\varepsilon \sum_{h=0}^{N} 2^{-h} = \frac{r^2}{2^8 er_\star^2} \Big( 1 - 2^{-N-1} \Big). \tag{5.22}
$$

By  $(5.6)$  we have

<span id="page-22-0"></span>
$$
\Phi_k : B_{r_{k+1}}(\mathbf{h}_{w_N}) \to B_{r_k}(\mathbf{h}_{w_N}),
$$
  
\n
$$
\sup_{u \in B_{r_{k+1}}(\mathbf{h}_{w_N})} |\Phi_k(u) - u|_{w_N} \le r_k \gamma^{-1} \mathfrak{K}^{\sharp} |R^{(k)}|_k.
$$
\n(5.23)

In conclusion we define

$$
\Psi := \Phi_0 \circ \Phi_1 \circ \cdots \circ \Phi_{N-1} : B_r(h_N) \to B_{2r}(h_N).
$$

Since we have

$$
\Phi_0 \circ \Phi_1 \circ \cdots \circ \Phi_{N-1} - id
$$
  
=  $(\Phi_0 - id) \circ \Phi_1 \circ \cdots \circ \Phi_{N-1} + (\Phi_1 - id) \circ \Phi_2 \circ \cdots \circ \Phi_{N-1} + \dots \Phi_{N-1} - id.$ 

 $\Box$ 

By  $(5.23)$  we get

$$
\sup_{u\in B_r(h_{w_N})}|\Psi(u)-u|_{w_N}\leq \sum_{k=0}^{N-1}r_k\gamma^{-1}\mathfrak{K}^{\sharp}|R^{(k)}|_{k}\stackrel{(5.19)}{\leq}2r\varepsilon\mathfrak{K}^{\sharp}\sum_{k=0}^{N-1}2^{-k}\leq 4r\mathfrak{K}^{\sharp}\varepsilon,
$$

proving the first bound in [\(1.33\)](#page-11-4). The second bound in (1.33) can be written as  $8\widehat{C}_1 r^2 \le 1$ ,<br>which follows from  $r < \widehat{x}$ . We finally set  $Z = Z_1$ ,  $P = P_2$  and the estimates (1.34) which follows from  $r \leq \hat{r}$ . We finally set  $Z = Z_N$ ,  $R = R_N$  and the estimates [\(1.34\)](#page-11-1) follow by [\(5.21\)](#page-22-1)–[\(5.22\)](#page-22-2). Of course the same reasoniong can be applied in order to construct the inverse, i.e. a symplectic change of variables  $\Phi : B_r(h_w) \mapsto B_{2r}(h_w)$  such that

<span id="page-23-1"></span>
$$
\Psi \circ \Phi u = \Phi \circ \Psi u = u, \quad \forall u \in B_{\frac{7}{8}r}(\mathbf{h}_w). \tag{5.24}
$$

When the nonlinearity *G* preserves momentum Theorem [1.3](#page-10-0) can be reformulated under slightly weaker assumptions. More precisely, setting  $\eta = 0$ 

<span id="page-23-2"></span>
$$
\mathfrak{C}_{0} := \max \left\{ 1, \sup_{0 \leq n < \mathbb{N}} \sup_{\substack{j,\boldsymbol{\alpha},\boldsymbol{\beta}, \\ \boldsymbol{\alpha}_{j}+\boldsymbol{\beta}_{j}\neq 0, \\ \pi(\boldsymbol{\alpha}-\boldsymbol{\beta})=0}} \frac{c_{\varrho_{n}^{*},0,w_{n+1}}^{(j)}(\boldsymbol{\alpha},\boldsymbol{\beta})}{c_{\varrho_{n},0,w_{n}}^{(j)}(\boldsymbol{\alpha},\boldsymbol{\beta})} \right\} < \infty,
$$
\n
$$
\mathfrak{K}_{0} := \max \left\{ 1, \sup_{0 \leq n < \mathbb{N}} \sup_{\substack{j,\boldsymbol{\alpha},\boldsymbol{\beta}, \\ \boldsymbol{\alpha}_{j}+\boldsymbol{\beta}_{j}\neq 0, \\ \pi(\boldsymbol{\alpha}-\boldsymbol{\beta})=0}} \frac{c_{\varrho_{n}^{*},0,w_{n+1}}^{(j)}(\boldsymbol{\alpha},\boldsymbol{\beta})}{c_{\varrho_{n},0,w_{n}}^{(j)}(\boldsymbol{\alpha},\boldsymbol{\beta})|\omega \cdot (\boldsymbol{\alpha}-\boldsymbol{\beta})|} \right\} < \infty,
$$
\n
$$
\mathfrak{K}_{0}^{\sharp} := \max \left\{ 1, \sup_{0 \leq n < \mathbb{N}} \sup_{\substack{j,\boldsymbol{\alpha},\boldsymbol{\beta}, \\ \boldsymbol{\alpha}_{j}+\boldsymbol{\beta}_{j}\neq 0, \\ \pi(\boldsymbol{\alpha}-\boldsymbol{\beta})=0}} \frac{c_{\varrho_{n}^{*},0,w_{n}}^{(j)}(\boldsymbol{\alpha},\boldsymbol{\beta})|\omega \cdot (\boldsymbol{\alpha}-\boldsymbol{\beta})|}{c_{\varrho_{n}^{*},0,w_{n}}^{(j)}(\boldsymbol{\alpha},\boldsymbol{\beta})|\omega \cdot (\boldsymbol{\alpha}-\boldsymbol{\beta})|} \right\} < \infty, \quad (5.25)
$$

the following holds.

**Proposition 5.1.** *If G preserves momentum Theorem [1.3](#page-10-0) holds word by word with*  $\mathfrak{C}_0$ ,  $\mathfrak{K}_0$ ,  $\mathfrak{K}_0^{\sharp}$  instead of  $\mathfrak{C}$ ,  $\mathfrak{K}$ ,  $\mathfrak{K}^{\sharp}$ . Moreover also the new perturbation R preserves mo*mentum.*

We note that in the case that *G* preserves momentum, the same result holds with  $\mathfrak{C}_0$ ,  $\mathfrak{K}_0$ ,  $\mathfrak{K}_0^{\sharp}$  instead of  $\mathfrak{C}$ ,  $\mathfrak{K}$ ,  $\mathfrak{K}^{\sharp}$ ; moreover also *R* preserves momentum.

<span id="page-23-0"></span>We finally give the following abstract stability result, whose proof is postponed to the "Appendix B".

Lemma 5.4. On the Hilbert space  $h_w$  consider the dynamical system

$$
\dot{v} = X_{\mathcal{N}} + X_R, \quad v(0) = v_0, \quad |v_0|_{\mathbf{w}} \leq \frac{3}{4}r,
$$

*where*  $N \in A_{r,0}(\mathbf{h}_{\mathbf{w}})$  *and*  $R \in \mathcal{H}_{r,\eta}(\mathbf{h}_{\mathbf{w}})$  *for some*  $r > 0$ ,  $\eta \geq 0$ . *Assume that* 

$$
\operatorname{Re}(X_{\mathcal{N}},v)_{h_w}=0.
$$

*Then*

<span id="page-24-3"></span>
$$
\left| |v(t)|_{w} - |v_0|_{w} \right| < \frac{r}{8}, \qquad \forall |t| \le \frac{1}{8|R|_{r,\eta,w}}.\tag{5.26}
$$

<span id="page-24-2"></span>**Corollary 5.1.** *Under the same assumptions of Theorem [1.3,](#page-10-0) the solution u*(*t*) *of the Hamiltonian flow of* [\(1.31\)](#page-11-2) *with initial datum*  $u(0) = u_0$  *such that*  $|u_0|_w \leq \frac{3r}{8}$  *exists and satisfies*

$$
|u(t)|_{\mathbf{w}} \le r \quad \text{for all times} \quad |t| \le \frac{1}{8\widehat{\mathbf{C}}_3 r^{2(\mathbf{N}+1)}}. \tag{5.27}
$$

*Proof.* Let us consider Hamiltonian [\(1.31\)](#page-11-2), take an initial datum  $|u_0|_w := r < \frac{3}{8}\hat{r}$  and apply the change of vartiables of Theorem 1.3. Denoting by  $v(0) = \Psi(u_0)$  we are under apply the change of vartiables of Theorem [1.3.](#page-10-0) Denoting by  $v(0) = \Psi(u_0)$  we are under the hypotheses of Lemma [5.4](#page-23-0) with  $\eta = 0$  and we conclude

$$
|v(t)|_w \leq \frac{7}{8}r
$$
,  $\forall |t| \leq \frac{1}{8|R|_{r,0,w}}$ .

Now we can apply [\(5.24\)](#page-23-1) in order to return to the original variables and deduce that  $u(t) = \Phi v(t)$  satisfies

$$
|u(t)|_{w} \leq r, \qquad \forall |t| \leq \frac{r^{-2(N+1)}}{8\widehat{C}_3} \leq \frac{1}{8|R|_{r,0,w}}.
$$

#### <span id="page-24-0"></span>**Part 2. Applications to Gevrey and Sobolev Cases**

In Part 2 we show how to apply the abstract BNF to Gevrey and Sobolev cases. Following the notations given in the introduction we work in the three sequence spaces defined for the applications G, S, M, see page 7. As explained in the introduction, in order to prove the estimates on the stability times we just need to verify that Assumption [1](#page-10-3) holds. This is the content of the next sections.

Let us start by setting some notations.

**Case** G) In the case  $w(p, s, a) = \left(\langle j \rangle^p e^{s \langle j \rangle^{\theta} + a|j|}\right)$ we denote  $h_{w(p,s,a)} = h_{p,s,a}$ ,<br> $\in \mathbb{Z}$ ,  $\in \mathbb{Z}$ , same notation for the norm of vectors  $|\cdot|_{p,s,a}$ . Regarding the norm of Hamiltonians we write  $|\cdot|_{r,\eta,w(p,s,a)}$ , consistently with Definition [1.2.](#page-9-0) Of course, for any  $0 \le p \le p', 0 \le$  $s \leq s', 0 \leq a \leq a'$  we have

$$
h_{p',s',a'} \subseteq h_{p,s,a}, \qquad |v|_{p,s,a} \le |v|_{p',s',a'}, \quad \forall v \in h_{p',s',a'}.
$$

**Case** S) If  $a = s = 0$  we denote  $h_{p,0,0} = h_p$ , same notation for the norm of vectors  $|\cdot|_p$  and hamiltonians  $|\cdot|_{r,n,w(p)}$ .

*Remark 5.1.* Note that, via the usual Fourier identification one has:

<span id="page-24-1"></span>
$$
|u|_p \le |u(x)|_{L^2} + |\partial_x^p u(x)|_{L^2} \le 2|u|_p. \tag{5.28}
$$

 $\Box$ 

**Case** M) In the case  $w_i = |j|^p$  where

$$
\lfloor j \rfloor := \max\{|j|, 2\}
$$

we denote the norm of vectors as

$$
||u||_p^2 = ||u||_w^2 := \sum_{j \in \mathbb{Z}} |j|^{2p} |u_j|^2.
$$
 (5.29)

*Remark 5.2.* Note that  $h_w$  in M) and  $h_p$  are the same vector space endowed with two equivalent norms. Moreover one has

<span id="page-25-0"></span>
$$
||u||_p \le 2^p |u(x)|_{L^2} + |\partial_x^p u(x)|_{L^2} \le 2||u||_p.
$$
\n(5.30)

**Definition 5.2** (momentum preserving regular Hamiltonians). Given  $r > 0$ ,  $p \ge 0$  let  $\mathcal{H}^{r,p}$  be the space of point-wise absolutely convergent Hamiltonians on  $||u||_p \leq r$  which preserves momentum and such that

$$
||H||_{r,p} := r^{-1} \left( \sup_{||u||_p \le r} ||X_{\underline{H}}||_p \right) < \infty,
$$
\n(5.31)

namely.<sup>16</sup>

<span id="page-25-4"></span>
$$
\|\cdot\|_{r,p} = |\cdot|\mathcal{H}_{r,0}(\mathbf{h}_{w}), \qquad \mathbf{w}_{j} = \lfloor j \rfloor^{p}
$$

We now verify that the nonlinearities in [\(1.1\)](#page-1-3) are bounded in the norm  $|\cdot|_{r,n,w}$  in the cases S, M, G.

**Proposition 5.2.** *Consider the correction term*  $P = \int_{\mathbb{T}} F(x, |u|^2) dx$  *in the NLS Hamiltonian* [\(1.21\)](#page-7-3)*, where the argument f in F satisfies*(1.2*). Let p > 1/2.* 

*(i)* For any  $a, s, \eta \ge 0$  *such that*  $a + \eta < a$  *and any*  $r > 0$  *such that*<sup>[17](#page-25-2)</sup>  $(C_{a} \log(p)r)^2 \le R$ , *we have*

<span id="page-25-5"></span>
$$
|P|_{r,\eta,w(p,s,a)} \leq C_{\text{Nem}}(p,s,a-a-\eta) \frac{(C_{\text{alg}}(p)r)^2}{R} |f|_{a,R} < \infty \qquad (5.32)
$$

*where f and*  $|f|_{a,R}$  *are defined in* [1.2.](#page-1-2) *(ii) If F is independent of*<sup>[18](#page-25-3)</sup> *x, for*  $(C_{\text{alg,M}}(p)r)^2 \leq R$  *we have* 

<span id="page-25-6"></span>
$$
||P||_{r,p} \le 2^p \frac{(C_{\text{alg,M}}(p)r)^2}{R} |f|_R < \infty.
$$
 (5.33)

This Proposition follows directly from the fact that the corresponding sequence spaces  $h_w$  are closed w.r.to convolution.

Let  $\star$  : h<sub>p,s,a</sub>  $\times$  h<sub>p,s,a</sub>  $\rightarrow$  h<sub>p,s,a</sub> be the convolution operation defined as

$$
(f,g)\mapsto f\star g:=\left(\sum_{j_1,j_2\in\mathbb{Z},\ j_1+j_2=j}f_{j_1}g_{j_2}\right)_{j\in\mathbb{Z}}.
$$

<span id="page-25-7"></span>The map  $\star$  :  $(f, g) \mapsto f \star g$  is continuous in the following sense:

<sup>&</sup>lt;sup>16</sup> Note that on the preserving momentum subspace  $\mathcal{H}_{r,n}(\mathbf{h}_{\mathbf{w}})$  coincides with  $\mathcal{H}_{r,0}(\mathbf{h}_{\mathbf{w}})$  for every  $\eta$ .

<span id="page-25-1"></span><sup>&</sup>lt;sup>17</sup> *R* is defined in [\(1.2\)](#page-1-2) and the constants in "Appendix A".

<span id="page-25-3"></span><span id="page-25-2"></span><sup>18</sup> i.e. *P* preserves momentum and we are assuming [\(1.3\)](#page-1-5).

**Lemma 5.5.** *For*  $p > 1/2$  *we have* 

$$
|f \star g|_{p,s,a} \leq C_{\text{alg}}(p)|f|_{p,s,a}|g|_{p,s,a}, \quad \|f \star g\|_{p} \leq C_{\text{alg,M}}(p)\|f\|_{p}\|g\|_{p}.
$$
\n(5.34)

The proof is given in "Appendix B".

*Proof of Proposition* [5.2.](#page-25-4) By definition (recall [\(1.2\)](#page-1-2) and [\(1.21\)](#page-7-3))

<span id="page-26-2"></span>
$$
F(x, y) = \int_0^y f(x, s)ds = \sum_{d=2}^\infty \frac{f^{(d-1)}(x)}{d}y^d =: \sum_{d=2}^\infty F^{(d)}(x)y^d \qquad (5.35)
$$

therefore we have

$$
P = \int_{\mathbb{T}} F(x, |u|^2) dx = \sum_{d \ge 2} \left( F^{(d)} \star \underbrace{u \star \cdots \star u}_{d \text{ times}} \star \underbrace{\bar{u} \star \cdots \star \bar{u}}_{d \text{ times}} \right)_0.
$$

To each analytic function  $F^{(d)}(x)$  we associate its Fourier coefficients; we have  $\left(F^{(d)}_j\right)$ *<sup>j</sup>*∈<sup>Z</sup>  $\in$  h<sub>p,s,a0</sub> for  $a_0 := a + \eta < a$  and  $s, p \ge 0$ . Indeed

$$
|F^{(d)}|_{p,s,a_0}^2 := \sum_j e^{2a_0|j|+2s\langle j\rangle^{\theta}} \langle j\rangle^{2p} |F_j^{(d)}|^2 \stackrel{(5.35)}{=} \sum_j e^{2a_0|j|+2s\langle j\rangle^{\theta}} \langle j\rangle^{2p} \frac{|f_j^{(d-1)}|^2}{d^2}
$$
  

$$
\leq \frac{c^2(p,s,a-a_0)}{d^2} \sum_j e^{2a|j|} |f_j^{(d-1)}|^2 = \frac{c^2(a-a_0,s,p)}{d^2} |f^{(d-1)}|_{\mathbb{T}_a}^2
$$

with

$$
c(p, s, t) := e^s + \sup_{x \ge 1} x^p e^{-tx + sx^\theta}.
$$

Now condition  $(1.2)$  ensures that  $(B.12)$  holds and our claim follows, by Lemma [B.2,](#page-50-0) setting  $a_0 = a + n$ . (ii) Follows from  $(B.14)$ .  $\square$ 

# <span id="page-26-0"></span>**6. Immersions**

The following proposition gathers the immersion properties of the norm  $|\cdot|_{r,\eta,w(p,s,a)}$ with respect to the parameters *p*,*s*, *a*.

**Proposition 6.1.** *The following inequalities hold:*

*(1)* Variations w.r.t. the paramater *p.* For any  $0 < \rho < r$ ,  $0 < \sigma < \eta$  and  $p_1 > 0$  we *have*

<span id="page-26-1"></span>
$$
|H|_{r-\rho,\eta-\sigma,w(p+p_1,s,a)} \leq C_{\text{mon}}(r/\rho,\sigma,p_1)|H|_{r,\eta,w(p,s,a)}.
$$

*(2)* Variation w.r.t. the parameter *s.* For any  $0 < \sigma < \eta$  we have

<span id="page-26-3"></span>
$$
|H|_{r,\eta-\sigma,w(p,s+\sigma,a)} \leq |H|_{r,\eta,w(p,s,a)}.\tag{6.1}
$$

*(3)* Variation w.r.t. the parameter *a. For any*  $0 < \sigma < \eta$ 

$$
|H|_{e^{-\sigma}r,\eta-\sigma,w(p,s,a+\sigma)} \leq e^{2\sigma}|H|_{r,\eta,w(p,s,a)}.\tag{6.2}
$$

*Remark 6.1.* All the items in the previous Proposition describe immersion properties of  $\mathcal{H}_{r,n}(\mathbf{h}_{p,s,a})$  w.r.t variations of the parameters.

In item (1) we say that if  $H \in H_{r,n}(\mathbf{h}_{p,s,a})$  (i.e. if its vector field maps  $B_r(\mathbf{h}_{p,s,a}) \to$  $h_{p,s,a}$ ) then it is also in  $\mathcal{H}_{r-\rho,\eta-\sigma}$  ( $h_{p+p_1,s,a}$ ) for any  $\rho, \sigma, p_1 > 0$ . Note however that the norm of  $H$  in the latter space is in general much larger, we denote this constant by  $C_{\text{mon}}$ .

In item (3) we have essentially the same phenomenon, only in order to increase the analiticity parameter  $a \rightarrow a + \sigma$ , we need to decrease the radius to  $e^{-\sigma}r$ .

Item (2) gives the best bound, indeed not only  $\mathcal{H}_{r,n-\sigma}(\mathbf{h}_{p,s+\sigma,a}) \subseteq \mathcal{H}_{r,n}(\mathbf{h}_{p,s,a})$  but the norm of *H* in the latter space does not increase.

To prove this Proposition we show that the hypotheses of Proposition [3.1](#page-14-0) hold. In order to prove this, in turn we strongly rely on some notation and results introduced by Bourgain in [\[Bou05](#page-63-17)] and extended later on by Cong–Li–Shi–Yuan in [\[CLSY](#page-63-24)] (Definition [6.1](#page-27-0) and Lemma [6.1](#page-28-0) below). The definitions and lemmata given below are the key technical arguments. Many of the ideas come from Bourgain in [\[Bou05](#page-63-17)] in the case of Gevrey regularity and for momentum preserving Hamiltonians, here we give a detailed presentation adapted to our more general setting and covering also the case of Sobolev regularity.

<span id="page-27-0"></span>**Definition 6.1.** Given a vector  $v = (v_i)_{i \in \mathbb{Z}} v_i \in \mathbb{N}$ ,  $|v| < \infty$  we denote by  $\hat{n} = \hat{n}(v)$  the vector  $(\hat{n}_l)_{l \in I}$  (where  $I \subset \mathbb{N}$  is finite) which is the decreasing rearrangement of

{N  $\ni$  *h* > 1 repeated  $v_h$  +  $v_{-h}$  times} ∪ {1 repeated  $v_1$  +  $v_{-1}$  +  $v_0$  times}.

*Remark 6.2.* A good way of envisioning this list is as follows. Given  $v = (v_i)_{i \in \mathbb{Z}}$  consider the monomial  $\bar{x}^v := \prod_i^v x_i^{v_i}$ . We can write uniquely

$$
x^v = \prod_i x_i^{v_i} = x_{j_1} x_{j_2} \cdots x_{j_{|v|}}
$$

then  $\hat{n}(v)$  is the decreasing rearrangement of the list  $(\langle j_1 \rangle, \dots, \langle j_{|v|} \rangle)$ .<br>As an example, consider the case  $v \neq 0$ . Then, by construction then

As an example, consider the case  $v \neq 0$ . Then, by construction there exists a unique *J* ≥ 0 such that  $v_j = 0$  for all  $|j| > J$  and  $v_j + v_{-j} \neq 0$  hence

$$
v = (\ldots, 0, v_{-J}, \ldots, v_0, \ldots, v_J, 0 \ldots).
$$

If  $J = 0$  then

$$
\widehat{n} = (\underbrace{1, \ldots, 1}_{v_0 \text{ times}})
$$

otherwise we have

$$
\widehat{n} = (\underbrace{J, \ldots, J}_{v_J+v_{-J} \text{ times}}, \underbrace{J-1, \ldots, J-1}_{v_{J-1}+v_{-J+1} \text{ times}}, \underbrace{1, \ldots, 1}_{v_1+v_{-1}+v_0 \text{ times}}).
$$

Given  $\alpha, \beta \in \mathbb{N}^{\mathbb{Z}}$  with  $1 < |\alpha| = |\beta| < \infty$ , from now on we define

$$
\widehat{n}=\widehat{n}(\boldsymbol{\alpha}+\boldsymbol{\beta}).
$$

We set the even number

$$
N:=|\alpha|+|\beta|,
$$

which is the cardinality of  $\hat{n}$ . We observe that, given

$$
\pi = \sum_{i \in \mathbb{Z}} i(\boldsymbol{\alpha}_i - \boldsymbol{\beta}_i) = \sum_{h > 0} h(\boldsymbol{\alpha}_h - \boldsymbol{\beta}_h - \boldsymbol{\alpha}_{-h} + \boldsymbol{\beta}_{-h}),
$$

there exists a choice of  $\sigma_i = \pm 1$ , 0 such that

<span id="page-28-1"></span>
$$
\pi = \sum_{l} \sigma_{l} \widehat{n}_{l} \tag{6.3}
$$

with  $\sigma_l \neq 0$  if  $\hat{n}_l \neq 1$ . Hence,

<span id="page-28-3"></span><span id="page-28-0"></span>
$$
\widehat{n}_1 \le |\pi| + \sum_{l \ge 2} \widehat{n}_l. \tag{6.4}
$$

Indeed, if  $\sigma_1 = \pm 1$ , the inequality follows directly from [\(6.3\)](#page-28-1); if  $\sigma_1 = 0$ , then  $\hat{n}_1 = 1$ and consequently  $\hat{n}_l = 1 \forall l$ . Since the mass is conserved, the list  $\hat{n}$  has at least two elements, and the inequality is achieved.

**Lemma 6.1.** *Given*  $\alpha$ ,  $\beta$  *such that*  $\sum_i i(\alpha_i - \beta_i) = \pi \in \mathbb{Z}$ , we have that setting  $\widehat{n} = \widehat{n}(\alpha + \beta)$  $\widehat{n}(\alpha + \beta)$ 

<span id="page-28-5"></span>
$$
\sum_{i} \langle i \rangle^{\theta} (\alpha_i + \beta_i) = \sum_{l \ge 1} \widehat{n}_l^{\theta} \ge 2\widehat{n}_1^{\theta} + (2 - 2^{\theta}) \sum_{l \ge 3} \widehat{n}_l^{\theta} - \theta |\pi|.
$$
 (6.5)

*Proof.* In "Appendix C".

The lemma above was proved in the simpler case of momentum preserving Hamiltonians in [\[Bou05\]](#page-63-17) for  $\theta = \frac{1}{2}$  and for general  $\hat{\theta}$  in [\[CLSY](#page-63-24)]. It is fundamental in discussing the properties of  $\mathcal{H}_{r,n}(\overline{h_{p,s,a}})$  with  $s > 0$ , indeed it implies

<span id="page-28-2"></span>
$$
\sum_{i} \langle i \rangle^{\theta} (\alpha_i + \beta_i) - 2 \langle j \rangle^{\theta} + |\pi(\alpha - \beta)| \ge (1 - \theta) \left( \sum_{l \ge 3} \widehat{n}_l^{\theta} + |\pi| \right) \ge 0 \quad (6.6)
$$

for all  $\alpha$ ,  $\beta$  such that  $\alpha$  *j* +  $\beta$  *j*  $\neq$  0.

*Proof of Proposition* [6.1.](#page-26-1) In all that follows we shall use systematically the fact that our Hamiltonians preserve the mass and are zero at the origin. These facts imply that  $|\alpha| = |\beta| > 1$ .

Let us start by proving **Item** (2), which is the simplest case. We need to show that

<span id="page-28-4"></span>
$$
\frac{c_{r,\eta-\sigma,w(p,s+\sigma,a)}^{(j)}(\boldsymbol{\alpha},\boldsymbol{\beta})}{c_{r,\eta,w(p,s,a)}^{(j)}(\boldsymbol{\alpha},\boldsymbol{\beta})} = \exp(-\sigma \left(\sum_{i} \langle i \rangle^{\theta} (\boldsymbol{\alpha}_{i} + \boldsymbol{\beta}_{i}) - 2 \langle j \rangle^{\theta} + |\pi(\boldsymbol{\alpha} - \boldsymbol{\beta})|) \le 1(6.7)
$$

The last inequality follows by  $(6.6)$  of Lemma  $6.1$ 

$$
\Box
$$

**Item** (1) First we assume that  $\rho \le r/2$ . By Proposition [3.1](#page-14-0) for any  $0 \le \rho \le r/2$ ,  $0 < \sigma < \eta$  and  $p_1 > 0$  we need to compute

<span id="page-29-1"></span>
$$
C_{\text{mon}} := \sup_{\substack{j,\boldsymbol{\alpha},\boldsymbol{\beta} \\ \boldsymbol{\alpha}_j + \boldsymbol{\beta}_j \neq 0}} \frac{c_{r-\rho,\eta-\sigma,w(p+p_1,s,a)}^{(j)}(\boldsymbol{\alpha},\boldsymbol{\beta})}{c_{r,\eta,w(p,s,a)}^{(j)}(\boldsymbol{\alpha},\boldsymbol{\beta})}
$$
  
= 
$$
\sup_{\substack{j,\boldsymbol{\alpha},\boldsymbol{\beta} \\ \boldsymbol{\alpha}_j + \boldsymbol{\beta}_j \neq 0}} \left(\frac{\langle j \rangle^2}{\prod_i \langle i \rangle^{\boldsymbol{\alpha}_i + \boldsymbol{\beta}_i}}\right)^{p_1} e^{-\sigma|\pi|} \left(\frac{r-\rho}{r}\right)^{|\boldsymbol{\alpha}|+|\boldsymbol{\beta}|-2}.
$$
 (6.8)

We use the notations of Definition [6.1,](#page-27-0) with  $\hat{n}(\alpha + \beta) \equiv \hat{n}$ . Since  $\alpha_j + \beta_j \neq 0$  we have that  $\langle j \rangle \leq \hat{n}_1$ . Note that

<span id="page-29-3"></span>
$$
\prod_{i} \langle i \rangle^{\alpha_i + \beta_i} = \prod_{l \ge 1} \widehat{n}_l. \tag{6.9}
$$

Hence

$$
\frac{\langle j \rangle^2}{\prod_i \langle i \rangle^{\alpha_i + \beta_i}} \leq \frac{\widehat{n}_1}{\prod_{l \geq 2} \widehat{n}_l}.
$$

Let us call  $N = |\alpha| + |\beta| \ge 2$ . By [\(6.4\)](#page-28-3) we have that

<span id="page-29-0"></span>
$$
\sup_{j,\alpha,\beta} \frac{\langle j \rangle^2}{\prod_i \langle i \rangle^{\alpha_i + \beta_i}} \leq \frac{\widehat{n}_1}{\prod_{l \geq 2} \widehat{n}_l} \leq \frac{\sum_{l=2}^N \widehat{n}_l + |\pi|}{\prod_{l=2}^N \widehat{n}_l}
$$

$$
\leq \frac{(N-1)\widehat{n}_2 + |\pi|}{\prod_{l=2}^N \widehat{n}_l} \leq \frac{N + |\pi|}{\prod_{l=3}^N \widehat{n}_l}.
$$
(6.10)

We have shown that

<span id="page-29-2"></span>
$$
\sup_{\substack{j,\boldsymbol{\alpha},\boldsymbol{\beta} \\ \boldsymbol{\alpha}_j+\boldsymbol{\beta}_j\neq 0}}\frac{\langle j\rangle^2}{\prod_i \langle i\rangle^{\boldsymbol{\alpha}_i+\boldsymbol{\beta}_i}} \leq N+|\pi|.
$$

Since  $(N + |\pi|)^{p_1} \leq 2^{p_1}(N^{p_1} + |\pi|^{p_1})$ , denoting  $L := \ln (r/r - \rho)$  we repeatedly use Lemma [C.1](#page-45-1) in order to control

$$
\sup_{N \ge 2, \pi \in \mathbb{Z}} (N + |\pi|)^{p_1} e^{-\sigma |\pi|} \left(\frac{r - \rho}{r}\right)^{N - 2}
$$
\n
$$
\le 2^{p_1} \left( \sup_{N \ge 2, \pi \in \mathbb{Z}} N^{p_1} e^{-\sigma |\pi| - L(N - 2)} + \sup_{N \ge 2, \pi \in \mathbb{Z}} |\pi|^{p_1} e^{-\sigma |\pi| - L(N - 2)} \right)
$$
\n
$$
\le 2^{p_1} \left( \max \left\{ \left(\frac{p_1}{L}\right)^{p_1}, 1 \right\} + \left(\frac{p_1}{\sigma}\right)^{p_1} \right) \le 2^{p_1 + 1} \max \left\{ \left(\frac{p_1}{L}\right)^{p_1}, \left(\frac{p_1}{\sigma}\right)^{p_1}, 1 \right\}
$$
\n
$$
\le 2^{p_1 + 1} p_1^{p_1} \max \left\{ \left(\frac{2r}{\rho}\right)^{p_1}, \left(\frac{1}{\sigma}\right)^{p_1}, 1 \right\} = C_{\text{mon}}, \tag{6.11}
$$

using that

$$
L \geq \ln(1+\rho/r) \geq 2\ln(3/2)\rho/r \geq \rho/2r,
$$

which holds since we are in the case  $\rho \le r/2$ . This completes the proof in the case  $\rho \leq r/2$ .

Consider now the case  $r/2 < \rho < r$ . Using the monotonicity of the norm w.r.t. *r* and the already proved case with  $\rho = r/2$ , we have

$$
|H|_{r-\rho,\eta-\sigma,w(p+p_1,s,a)} \leq |H|_{r/2,\eta-\sigma,w(p+p_1,s,a)}
$$
  
\n
$$
\leq 2^{p_1+1} \max\left\{ (4p_1)^{p_1}, \left( \frac{p_1}{e\sigma} \right)^{p_1}, 1 \right\} |H|_{r,\eta,w(p,s,a)}
$$
  
\n
$$
\leq 2^{p_1+1} p_1^{p_1} \max\left\{ \left( \frac{2r}{\rho} \right)^{p_1}, \left( \frac{1}{\sigma} \right)^{p_1}, 1 \right\} |H|_{r,\eta,w(p,s,a)},
$$

proving (1) also in the case  $r/2 < \rho < r$ . **Item** (3) We proceed as in item  $(1) - (2)$ ,

<span id="page-30-1"></span>
$$
\frac{c_{e^{-\sigma}r,\eta-\sigma,w(p,s,a+\sigma)}^{(j)}(\alpha,\beta)}{c_{r,\eta,w(p,s,a)}^{(j)}(\alpha,\beta)} = \exp(-\sigma(\sum_{i} \langle i \rangle^{\theta}(\alpha_i + \beta_i) - 2\langle j \rangle^{\theta} + |\pi(\alpha - \beta)| - (|\alpha| + |\beta| - 2)) \le e^{2\sigma}.
$$
 (6.12)

our claim follows since, by formula [\(6.4\)](#page-28-3), one has

<span id="page-30-0"></span>
$$
\sum_{i} (\boldsymbol{\alpha}_i + \boldsymbol{\beta}_i)|i| - 2|j| + |\pi| \ge \sum_{l \ge 2} \widehat{n}_l - \widehat{n}_1 + |\pi| - |\boldsymbol{\alpha}_0 + \boldsymbol{\beta}_0| \ge -(|\boldsymbol{\alpha}| + |\boldsymbol{\beta}|)(6.13)
$$

*Remark 6.3.* Note that a key point in Items (1) and (2) are the estimates [\(6.10\)](#page-29-0) and [\(6.6\)](#page-28-2) where we control the ratio of the coefficient [\(1.27\)](#page-9-3) in terms of  ${\hat{n}}_l$ <sub>*l*l≥3</sub> (namely uniformly with respect to  $\hat{n}_1$  and  $\hat{n}_2$ ). This means that if  $\hat{n}_3$  is "big", then the norm of the Hamiltonian is correspondingly small: polynomially in the Sobolev case and subexponentially in the Gevrey one. This is a seminal property which appears in different flavors thoughout the literature; in Proposition [6.1](#page-26-1) we do not really need to exploit it. Instead, it will be heavily used for a sharp control on the small divisors appearing in the Homological equation (see proof of Proposition [7.1\)](#page-32-2).

<span id="page-30-2"></span>Incidentally we note that norm  $|\cdot|_{r,n,w(p,s,a)}$  possesses the tameness property.

#### **Proposition 6.2.**

$$
\sup_{|u|_{p_0,s,a} \le r-\rho} \frac{|X_{\underline{H}}|_{p,s,a}}{|u|_{p,s,a}} \le C_{\text{tame}}(\rho, \eta, p)|H|_{r, \eta, w(p_0,s,a)}.
$$

*Proof.* In "Appendix B". □

**Proposition 6.3.** *The norm*  $\|\cdot\|_{r,p}$  *is monotone decreasing in p, namely*  $\|\cdot\|_{r,p+p_1} \le$  $\| \cdot \|_{r,p}$  *for any p*<sub>1</sub> > 0*.* 

*Proof.* For the norm  $\|\cdot\|_{r,p}$  the quantity in [\(1.27\)](#page-9-3) becomes (recall that in the norm of a momentum preserving hamiltonian there is need of introducing the parameter  $\eta$ )

$$
c_{r,p}^{(j)}(\boldsymbol{\alpha},\boldsymbol{\beta}) := r^{|\boldsymbol{\alpha}|+|\boldsymbol{\beta}|-2} \left(\frac{\lfloor j\rfloor^2}{\prod_{i\in\mathbb{Z}}\lfloor i\rfloor^{(\boldsymbol{\alpha}_i+\boldsymbol{\beta}_i)}}\right)^p.
$$
 (6.14)

By Lemma [3.1](#page-14-2) item (ii) we only need to show that

<span id="page-31-2"></span>
$$
c_{r,p+p_1}^{(j)}(\boldsymbol{\alpha},\boldsymbol{\beta}) \le c_{r,p}^{(j)}(\boldsymbol{\alpha},\boldsymbol{\beta})
$$
\n(6.15)

for all  $j$ ,  $\alpha$ ,  $\beta$  with  $|\alpha| = |\beta| \ge 1$  and  $\alpha_j + \beta_j \ge 1$  (recall the momentum conservation), namely we have to prove that

<span id="page-31-0"></span>
$$
\sup_{\substack{j,\boldsymbol{\alpha},\boldsymbol{\beta} \\ \boldsymbol{\alpha}_j+\boldsymbol{\beta}_j\geq 1}}\frac{|j|^2}{\prod_i |j|^{\boldsymbol{\alpha}_i+\boldsymbol{\beta}_i}}\leq 1.
$$
\n(6.16)

We first show that the inequality holds in the case  $j = 0, \pm 1$ . Indeed we have

$$
\prod_i \lfloor i \rfloor^{\alpha_i + \beta_i} \ge \prod_i 2^{\alpha_i + \beta_i} = 2^{\sum_i \alpha_i + \beta_i} \ge 4
$$

since  $\sum_i \alpha_i + \beta_i \ge 2$  (by the fact that  $|\alpha| = |\beta| \ge 1$ ). Consider now the case  $|j| = \lfloor j \rfloor \ge 2$ . Since  $\alpha_j + \beta_j \ge 1$ , inequality [\(6.16\)](#page-31-0) follows by

<span id="page-31-1"></span>
$$
\sup_{j,\alpha,\beta} \frac{|j|}{\prod_{i \neq j} \lfloor i \rfloor^{\alpha_i + \beta_i}} \le 1.
$$
\n(6.17)

By momentum conservation we have

$$
|j| \leq \sum_{i \neq j} |i| (\boldsymbol{\alpha}_i + \boldsymbol{\beta}_i) \leq \sum_{i \neq j} |i| (\boldsymbol{\alpha}_i + \boldsymbol{\beta}_i)
$$
 (6.18)

and [\(6.17\)](#page-31-1) follows if we show that

$$
\sup_{j,\boldsymbol{\alpha},\boldsymbol{\beta}} \frac{\sum_{i \neq j} \lfloor i \rfloor (\boldsymbol{\alpha}_i + \boldsymbol{\beta}_i)}{\prod_{i \neq j} \lfloor i \rfloor^{\boldsymbol{\alpha}_i + \boldsymbol{\beta}_i}} \le 1,
$$
\n(6.19)

where we can restrict the sum and the product to the indexes *i* such that  $\alpha_i + \beta_i \geq 1$ . This last estimates follows by the fact that given  $x_k \geq 1$ 

$$
\frac{\sum_{2\leq k\leq n}kx_k}{\prod_{2\leq k\leq n}k^{x_k}}\leq 1,
$$

as it can be easly proved by induction over *n* (noting that  $n^x \ge nx$  for  $n \ge 2$ , and any  $x > 1$ ).  $\Box$ 

#### <span id="page-32-0"></span>**7. Homological Equation**

Now we give estimates on the solution of the homological equation

<span id="page-32-4"></span><span id="page-32-3"></span>
$$
L_{\omega}S := \{D_{\omega}, S\} = R
$$

<span id="page-32-2"></span>The constants  $C_1$ ,  $C_2$ ( $r$ ,  $\sigma$ ,  $t$ ) are defined in "Appendix A". Note that  $C_1$  depends only on θ.

**Proposition 7.1.** Let  $\omega \in D_{\gamma,q}$  and let  $0 < \sigma < \eta$ ,  $0 < \rho < r/2$ . For any  $R \in$  $\mathcal{R}_{r,\eta}(\mathbf{h}_{p,s,a})$ , the Homological equation  $L_{\omega}S = R$  has a unique solution  $S = L_{\omega}^{-1}R$ , *which satisfies the following two bounds:*

$$
\left| L_{\omega}^{-1} R \right|_{r, \eta - \sigma, w(p, s + \sigma, a)} \leq \gamma^{-1} e^{\mathcal{C}_1 \sigma^{-\frac{3}{\theta}}} |R|_{r, \eta, w(p, s, a)}
$$
(G)

$$
\left| L_{\omega}^{-1} R \right|_{r-\rho, \eta-\sigma, \mathbf{w}(p+\tau, s, a)} \leq \gamma^{-1} C_2(r/\rho, \sigma, \tau) |R|_{r, \eta, \mathbf{w}(p, s, a)} \tag{S}
$$

hence  $L_{\omega}^{-1}R \in \mathcal{R}_{r,\eta-\sigma}(\mathbf{h}_{p,s+\sigma,a}) \cap \mathcal{R}_{r-\rho,\eta-\sigma}(\mathbf{h}_{p+\tau,s,a}).$ 

If *R* preserves momentum  $R \in \mathcal{R}_{r,0}(h_w)$ , with  $w_i = |j|^p$ , the unique solution of the Homological equation preserves momentum and satisfies

<span id="page-32-5"></span>
$$
||L_{\omega}^{-1}R||_{r,p+\tau_1} \leq \gamma^{-1}6^{\tau_1}(4^6e^{27})^{2+q}||R||_{r,p},
$$
\n(M)

so 
$$
S = L_{\omega}^{-1} R \in \mathcal{R}_{r,0}(\mathbf{h}_{\mathbf{w}'})
$$
, with  $\mathbf{w}'_j = \lfloor j \rfloor^{p+\tau_1}$ .

*Remark 7.1.* As in the abstract case we assume that  $X_R$  maps  $B_r(h_{p,s,a}) \to h_{p,s,a}$  and then show that *S* maps some smaller ball (because it has smaller radius or is in a stronger topology) to itself. This can be done in two ways: if we increase the Gevrey regularity index  $s \rightarrow s + \sigma$  (case G) then the increase can be arbitrarily small, at the price of an exponential increase in the bound.

If we want to keep *s* fixed (say that we start with  $s = a = 0$  and want to stay in the Sobolev class) then we have to increase the regularity *p by a fixed amount.* The main difference between the cases S and M is that in the first case one has to decrease *r*, η and the bound on *S* diverges as  $\rho, \sigma \to 0$ . In the second case, instead we have to increase the regularity *p* by a slightly larger amount but then we get a uniform bound for *S*.

Note that, differently from Proposition [6.1,](#page-26-1) we cannot consider the purely analytic case  $(s, p$  fixed say to 0, 1). This is due to the fact that in  $(6.13)$  we have a much weaker bound for the ratio of the coefficients in  $(6.12)$ , w.r.t. the one afforded by  $(6.6)$  and  $(6.10)$  for the Gevrey and Sobolev cases.

The following Lemma is the key point in the control of the small divisors appearing in the solution of the Homological equation. Here we strongly use the fact that we are working with a dispersive PDE on the circle with superlinear dispersion law.

<span id="page-32-1"></span>**Lemma 7.1.** *Consider*  $\alpha$ ,  $\beta \in \mathbb{N}^{\mathbb{Z}}$  *with*  $1 \leq |\alpha| = |\beta| < \infty$ . If

<span id="page-32-6"></span>
$$
\left|\sum_{i} (\boldsymbol{\alpha}_{i} - \boldsymbol{\beta}_{i})i^{2}\right| \le 10 \sum_{i} |\boldsymbol{\alpha}_{i} - \boldsymbol{\beta}_{i}|,
$$
\n(7.1)

*then for all j such that*  $\alpha_j + \beta_j \neq 0$  *one has* 

<span id="page-33-2"></span>
$$
\sum_{i} |\alpha_{i} - \beta_{i}| \langle i \rangle^{\theta/2} \leq C_{*} \left( \sum_{i} (\alpha_{i} + \beta_{i}) \langle i \rangle^{\theta} - 2 \langle j \rangle^{\theta} + |\pi| \right), \qquad C_{*} = \frac{13}{1 - \theta}
$$
\n(7.2)

$$
\prod_{i} (1 + |\boldsymbol{\alpha}_i - \boldsymbol{\beta}_i| \langle i \rangle) \le e^{27} (1 + |\pi|)^3 N^6 \prod_{l=3}^N \widehat{n}_l^{\tau_0}
$$
\n(7.3)

*where*  $N = |\boldsymbol{\alpha}| + |\boldsymbol{\beta}|$  *and*  $\pi = \sum_i i(\boldsymbol{\alpha}_i - \boldsymbol{\beta}_i)$  (recall [\(1.25\)](#page-8-3). *Proof.* In "Appendix C"

Note that

<span id="page-33-0"></span>
$$
\left|\sum_{i} (\alpha_i - \beta_i) i^2\right| \ge 10 \sum_{i} |\alpha_i - \beta_i| \quad \implies \quad |\omega \cdot (\alpha - \beta)| \ge 1. \tag{7.4}
$$

Indeed denoting  $\omega_j = j^2 + \xi_j \langle j \rangle^{-q}$  with  $|\xi_j| \leq \frac{1}{2}$ ,

$$
|\omega \cdot (\boldsymbol{\alpha} - \boldsymbol{\beta})| \ge 10 \sum_j |\boldsymbol{\alpha}_j - \boldsymbol{\beta}_j| - \frac{1}{2} \sum_j |\boldsymbol{\alpha}_j - \boldsymbol{\beta}_j| \ge 1.
$$

*Proof.* In the following, we will compute for each item the corresponding  $K$ ,  $K_0$  defined in [\(4.8\)](#page-17-1) and [\(4.10\)](#page-17-4), and show their finiteness in order to apply Lemma [4.2](#page-16-1) and give the explicit upper bounds entailed in Proposition [7.1](#page-32-2) ([G](#page-32-3))–([S](#page-32-4))– $(M)$  $(M)$  $(M)$ .

Item G) In this case by  $(6.7)$ 

$$
K = \gamma \sup_{j:\boldsymbol{\alpha}_j+\boldsymbol{\beta}_j\neq 0} \frac{e^{-\sigma\left(\sum_i \langle i \rangle^{\theta}(\boldsymbol{\alpha}_i+\boldsymbol{\beta}_i)-2\langle j \rangle^{\theta}+|\pi|\right)}}{|\omega \cdot (\boldsymbol{\alpha}-\boldsymbol{\beta})|}.
$$

There are two cases.

If [\(7.1\)](#page-32-6) does not hold, then by [\(7.4\)](#page-33-0)  $|\omega \cdot (\alpha - \beta)| \ge 1$  and by [\(6.5\)](#page-28-5) and [\(4.4\)](#page-16-4) we get

$$
\gamma \frac{e^{-\sigma(\sum_i \langle i \rangle^{\theta}(\boldsymbol{\alpha}_i+\boldsymbol{\beta}_i)-2\langle j \rangle^{\theta}+|\pi|)} }{|\omega \cdot (\boldsymbol{\alpha}-\boldsymbol{\beta})|} \leq 1
$$

and the bound is trivially achieved.

Otherwise, let us consider the case in which [\(7.1\)](#page-32-6) holds. By applying Lemma [7.1,](#page-32-1) since  $\omega \in D_{\gamma,q}$  we get:

<span id="page-33-1"></span>
$$
\gamma \frac{e^{-\sigma(\sum_{i} \langle i \rangle^{\theta}(\boldsymbol{\alpha}_{i}+\boldsymbol{\beta}_{i})-2\langle j \rangle^{\theta}+|\pi|)} }{|\omega \cdot (\boldsymbol{\alpha}-\boldsymbol{\beta})|}
$$
  
\n
$$
\leq e^{-\frac{\sigma}{C_{*}}\sum_{i} |\boldsymbol{\alpha}_{i}-\boldsymbol{\beta}_{i}| \langle i \rangle^{\frac{\theta}{2}} \prod_{i} \left(1+(\boldsymbol{\alpha}_{i}-\boldsymbol{\beta}_{i})^{2} \langle i \rangle^{2+q}\right)}
$$
  
\n
$$
\leq \exp \sum_{i} \left[-\frac{\sigma}{C_{*}} |\boldsymbol{\alpha}_{i}-\boldsymbol{\beta}_{i}| \langle i \rangle^{\frac{\theta}{2}} + \ln \left(1+(\boldsymbol{\alpha}_{i}-\boldsymbol{\beta}_{i})^{2} \langle i \rangle^{2+q}\right)\right]
$$
  
\n
$$
= \exp \sum_{i} f_{i} (|\boldsymbol{\alpha}_{i}-\boldsymbol{\beta}_{i}|) \qquad (7.5)
$$

$$
\overline{a}
$$

where, for  $0 < \sigma < 1$ ,  $i \in \mathbb{Z}$  and  $x > 0$ , we defined

$$
f_i(x) := -\frac{\sigma}{C_*} x \langle i \rangle^{\frac{\theta}{2}} + \ln \left( 1 + x^2 \langle i \rangle^{2+q} \right).
$$

In order to bound [\(7.5\)](#page-33-1), we need the following lemma, whose proof is postponed to "Appendix C".

# <span id="page-34-2"></span>**Lemma 7.2.** *Setting*

$$
i_{\sharp} := \left(\frac{8C_*(q+3)}{\sigma\theta} \ln \frac{4C_*(q+3)}{\sigma\theta}\right)^{\frac{2}{\theta}},
$$

*we get*

<span id="page-34-0"></span>
$$
\sum_{i} f_{i}(|\ell_{i}|) \le 7(q+3)i_{\sharp} \ln i_{\sharp} - \frac{\sigma}{2C_{*}}(\widehat{n}_{1}(\ell))^{\frac{\theta}{2}}
$$
(7.6)

*for every*  $\ell \in \mathbb{Z}^{\mathbb{Z}}$  *with*  $|\ell| < \infty$ .

The inequality ([G](#page-32-3)) follows from plugging  $(7.6)$  into  $(7.5)$  and evaluating the constant. Item S) In this case  $K$  in  $(4.8)$  is (recall  $(6.8)$ )

<span id="page-34-1"></span>
$$
K = \gamma \sup_{j:\alpha_j+\beta_j\neq 0} \left(1-\frac{\rho}{r}\right)^{N-2} \left(\frac{\langle j\rangle^2}{\prod_i \langle i\rangle^{\alpha_i+\beta_i}}\right)^{\tau} \frac{e^{-\sigma|\pi|}}{|\omega \cdot (\alpha-\beta)|},\tag{7.7}
$$

where  $N = |\alpha| + |\beta|$ .

As before we consider two cases.

If  $(7.1)$  is not satisfied then[\(7.4\)](#page-33-0) holds and the right hand side of  $(7.7)$  is bounded by the quantity in [\(6.8\)](#page-29-1) and it is estimated analogusly.

If  $(7.1)$  holds instead, by applying formula  $(6.10)$ , Lemma [7.1](#page-32-1) and the fact that  $\omega \in D_{\nu,q}$  we get:

$$
\left(\frac{\langle j\rangle^2}{\prod_i \langle i\rangle^{\alpha_i+\beta_i}}\right)^{\tau} \frac{1}{|\omega \cdot (\alpha - \beta)|} \leq \left(\frac{\langle j\rangle^2}{\prod_i \langle i\rangle^{\alpha_i+\beta_i}}\right)^{\tau} \prod_i \left(1 + |\alpha_i - \beta_i|^2 \langle i\rangle^{2+q}\right)
$$
\n
$$
\leq \left(\frac{N+|\pi|}{\prod_{l=3}^N \widehat{n}_l}\right)^{\tau} \left(e^{27}(1+|\pi|)^3 N^6 \prod_{l\geq 3} \widehat{n}_l^{r_0}\right)^{2+q}
$$
\n
$$
\leq e^{27(2+q)}(N+|\pi|)^{\tau+9(2+q)} \leq e^{27(2+q)}(N+|\pi|)^{3\tau}.
$$

By using Lemma [C.1](#page-45-1) (just like explained in detail in formula [\(6.11\)](#page-29-2) with  $p_1 = 3\tau$ ), *K* in  $(7.7)$  is bounded by

$$
e^{27(2+q)}(N+|\pi|)^{3\tau} \left(1-\frac{\rho}{r}\right)^{N-2} e^{-\sigma|\pi|}
$$
  
 
$$
\leq e^{27(2+q)} 2^{3\tau+1} (3\tau)^{3\tau} \max \left\{ \left(\frac{2r}{\rho}\right)^{3\tau}, \left(\frac{1}{\sigma}\right)^{3\tau}, 1 \right\}
$$

Item  $M$ ) Note that in this case the constant in  $(4.8)$  amounts to

$$
K_0 = \gamma \sup_{\substack{j \in \mathbb{Z}, \, \boldsymbol{\alpha} \neq \boldsymbol{\beta} \in \mathbb{N}^{\mathbb{Z}} \\ \boldsymbol{\alpha}_j + \boldsymbol{\beta}_j \neq 0, \, \sum_i i(\boldsymbol{\alpha}_i - \boldsymbol{\beta}_i) = 0}} \left( \frac{\lfloor j \rfloor^2}{\prod_i \lfloor i \rfloor^{\boldsymbol{\alpha}_i + \boldsymbol{\beta}_i}} \right)^{\tau_1} \frac{\gamma}{|\omega \cdot (\boldsymbol{\alpha} - \boldsymbol{\beta})|}.
$$

We have two cases. If [\(7.4\)](#page-33-0) holds  $K_0 \le \gamma$  by [\(6.16\)](#page-31-0). Otherwise [\(7.1\)](#page-32-6) holds and, therefore, [\(7.3\)](#page-33-2) (note that here  $\pi = 0$ ) applies, giving

$$
K_0 \le \sup \left( \frac{|j|^2}{\prod_i |i|^{|\alpha_i + \beta_i}} \right)^{\tau_1} \prod_i \left( 1 + |\alpha_i - \beta_i|^2 \langle i \rangle^{2+q} \right)
$$
  

$$
\le \sup \left( \frac{|j|^2}{\prod_i |i|^{|\alpha_i + \beta_i}} \right)^{\tau_1} e^{27(2+q)} N^{6(2+q)} \prod_{l=3}^N \widehat{n}_l^{\tau_0(2+q)}
$$

since  $\omega \in D_{\gamma,q}$ . We claim that

$$
N \le 4 \prod_{l=3}^{N} [\widehat{n}_l]^{\frac{1}{4 \ln 2}}.
$$
 (7.8)

Indeed if  $N = 2$ , the inequality is trivial. Since N is even we have to consider only the case  $N \geq 4$ , which follows by Lemma [C.1.](#page-45-1) Recalling [\(6.9\)](#page-29-3) we have

$$
\prod_{i} \lfloor i \rfloor^{\alpha_i + \beta_i} = \prod_{l \ge 1} \lfloor \widehat{n}_l \rfloor. \tag{7.9}
$$

Then

$$
\sup_{\substack{j,\boldsymbol{\alpha},\boldsymbol{\beta} \\ \boldsymbol{\alpha}_j+\boldsymbol{\beta}_j\geq 1}}\frac{|j|^2}{\prod_i |i|^{\boldsymbol{\alpha}_i+\boldsymbol{\beta}_i}} \leq \frac{|\hat{n}_1|^2}{\prod_{l\geq 1}|\widehat{n}_l|} = \frac{|\hat{n}_1|}{\prod_{l\geq 2}|\widehat{n}_l|} \leq \frac{\sum_{l\geq 2}|\widehat{n}_l|}{\prod_{l\geq 2}|\widehat{n}_l|} = \frac{1}{\prod_{l\geq 3}|\widehat{n}_l|} + \frac{\sum_{l\geq 3}|\widehat{n}_l|}{\prod_{l\geq 2}|\widehat{n}_l|},
$$

where the last inequality holds by momentum conservation. Then<sup>19</sup>

$$
K_0 \leq 2^{\tau_1 - 1} \left( \frac{1}{\prod_{l \geq 3} [\widehat{n}_l]^{\tau_1}} + \frac{(\sum_{l \geq 3} [\widehat{n}_l])^{\tau_1}}{\prod_{l \geq 2} [\widehat{n}_l]^{\tau_1}} \right) (4^6 e^{27})^{2+q} \prod_{l \geq 3} [\widehat{n}_l]^{\tau_1/2}
$$
  

$$
\leq 2^{\tau_1 - 1} (4^6 e^{27})^{2+q} \left( 1 + \frac{(\sum_{l \geq 3} [\widehat{n}_l])^{\tau_1}}{[\widehat{n}_2]^{\tau_1} \prod_{l \geq 3} [\widehat{n}_l]^{\tau_1/2}} \right)
$$
  

$$
\leq 2^{\tau_1 - 1} (4^6 e^{27})^{2+q} \left( 1 + \frac{([\widehat{n}_3]^{1/2} + 4)^{\tau_1}}{[\widehat{n}_2]^{\tau_1}} \right)
$$

by Lemma [C.2](#page-50-0) with  $a = 1/2$ . The estimate on  $K_0$ , hence inequality ([M](#page-32-5)) follows.  $\square$ 

<span id="page-35-0"></span><sup>19</sup> Using that  $(a + b)^{\tau_1} \leq 2^{\tau_1 - 1} (a^{\tau_1} + b^{\tau_1})$  for  $a, b \geq 0, \tau_1 \geq 1$ .

### <span id="page-36-0"></span>**8. Birkhoff Normal Form**

<span id="page-36-4"></span>We are now ready to apply Theorem [1.3](#page-10-0) to the three applications G, S, M, defined in page 7. We start by verifying the assumptions.

#### **Lemma 8.1.** *The following holds*

G*)* Let  $s > 0$ ,  $p > 1/2$  and  $a > 0$ . Then for all  $N > 1$ ,  $0 < n < s$ ,  $w := w(p, s, a)$  $and w_0 := w(p, s - p, a)$  *satisfy the Birkhoff assumption at step* N *and in* [\(1.29\)](#page-10-5) *we can take*

$$
\mathfrak{C}=1,\quad \mathfrak{K},\ \mathfrak{K}^{\sharp}\leq e^{\mathcal{C}_1\left(\frac{N}{\eta}\right)^{\frac{3}{\theta}}}.
$$

S*)* Let  $\tau_S = \tau$ ,  $s, a \ge 0$ ,  $p \ge 3\tau_S + 1$  and set<sup>[20](#page-36-1)</sup> N :=  $[\frac{p-1}{\tau_S}]$ . Then  $\eta > 0$ , w :=  $w(p, s, a)$  *and*<sup>[21](#page-36-2)</sup>  $w_0 := w(p - N\tau_s, s, a)$  *satisfy the Birkhoff assumption at step* N *and in* [\(1.29\)](#page-10-5) *we can take*

$$
\mathfrak{C}\leq C_{\text{mon}}(4N,\eta/N,\tau_{\text{S}}),\quad \mathfrak{K}\leq C_{2}(4N,\eta/N,\tau_{\text{S}}),\quad \mathfrak{K}^{\sharp}\leq C_{2}(4N,\eta/N,N\tau_{\text{S}}).
$$

 $(M)$  Let  $\tau_M = \tau_1$ ,  $p \geq 3\tau_M + 1$  and<sup>[22](#page-36-3)</sup> set  $N := [\frac{p-1}{\tau_M}]$ . Then  $\eta = 0$ ,  $w := (\lfloor j \rfloor^p)_{j \in \mathbb{Z}}$  and *w*<sub>0</sub> := ( $\lfloor j \rfloor^{p-Nt_M}$ )<sub>*j*∈Z</sub> *satisfy the "momentum preserving" Birkhoff assumption at step*<br>
and in the step N *and in* [\(5.25\)](#page-23-2) *we can take*

$$
\mathfrak{C}_0 = 1, \quad \mathfrak{K}_0, \; \mathfrak{K}_0^{\sharp} \leq 6^{\tau_{\mathbb{M}}} (4^6 e^{27})^{2+q}.
$$

*Proof.* G) Set

$$
\mathsf{w}_{n,j}:=\mathsf{w}_{0,j}e^{\frac{n\eta}{\mathsf{N}}\langle j\rangle^{\theta}},\qquad\forall n=1,\ldots,\mathsf{N}.
$$

The computation of  $\mathfrak C$  follows from [\(6.1\)](#page-26-3); the ones of  $\mathfrak K$ ,  $\mathfrak K^{\sharp}$  from Proposition [7.1.](#page-32-2) S) Set

$$
w_{n,j} := w_{0,j} \langle j \rangle^{n\tau_{S}}, \quad \forall n = 1, \ldots, N.
$$

The computation of  $\mathfrak C$  follows from [\(6.8\)](#page-29-1); the ones of  $\mathfrak K$ ,  $\mathfrak K^{\sharp}$  from Proposition [7.1.](#page-32-2) M) Set

$$
w_{n,j} := w_{0,j} \lfloor j \rfloor^{n \tau_M}, \quad \forall n = 1, \ldots, N.
$$

The computation of  $\mathfrak{C}_0$  follows from [\(6.15\)](#page-31-2); the ones of  $\mathfrak{K}_0$ ,  $\mathfrak{K}_0^{\sharp}$  again from Proposition  $7.1.$ 

<span id="page-36-1"></span>

<sup>&</sup>lt;sup>20</sup> [·] is the integer part.<br><sup>21</sup> Note that  $1 \le p - N\tau_S < 1 + \tau_S$ .

<span id="page-36-3"></span><span id="page-36-2"></span><sup>&</sup>lt;sup>22</sup> Note that  $1 \leq p - N\tau_M < 1 + \tau_M$ .

We now state the Birkhoff Normal Form Theorem  $(1.3)$  for the Hamiltonian in  $(1.21)$ in the usual three cases. First we define

<span id="page-37-2"></span>
$$
r(G) := \min \left\{ \frac{\delta_G}{\sqrt{N}e^{\frac{1}{2}C_1\left(\frac{N}{\eta_G}\right)^{\frac{3}{2}}}}, \frac{\sqrt{R}}{2C_{a1g}(p)} \right\},\newline
$$
  
\n
$$
\delta_G := \frac{\sqrt{\gamma R}}{C_{a1g}(p)\sqrt{2^{11}eC_{Nem}(p, s - \eta_G, a - a - \eta_G)|f|_{a,R}}}
$$
  
\n
$$
C_1(G) := \frac{e^{C_1\left(\frac{N}{\eta_G}\right)^{\frac{3}{2}}}}{2^7e\delta_G^2},\newline
$$
  
\n
$$
C_2(G) := \frac{\gamma}{2^9e\delta_G^2}\left(\frac{Ne^{C_1\left(\frac{N}{\eta_G}\right)^{\frac{3}{2}}}}{4\delta_G^2}\right)^N.
$$
  
\n
$$
r(S) := \min \left\{\frac{d_S}{\sqrt{NC_2(4N, a/2N, N\tau_S)}}, \frac{\sqrt{R}}{5 \cdot 2^{\tau_S+2}}\right\}, \text{ where}
$$
  
\n
$$
d_S := \frac{\sqrt{\gamma R}}{\sqrt{2^{17}C_{Nem}(p - N\tau_S, 0, a/2)|f|_{a,R}}},\newline
$$
  
\n
$$
C_1(S) := \frac{c_2(4N, a/2N, N\tau_S)}{2^7e d_S^2}, \quad C_2(S) := \frac{\gamma}{2^8e d_S^2},\newline
$$
  
\n
$$
C_3(S) := \frac{2^8C_{Nem}(p - N\tau_S, 0, a/2)|f|_{a,R}}{eR}\left(\frac{NC_{mon}(4N, a/2N, \tau_S)C_2(4N, a/2N, \tau_S)}{4d_S^2}\right)^N,
$$
  
\n
$$
r(M) := \min \left\{\frac{\delta_M}{\sqrt{N}}, \frac{R}{\sqrt{\frac{R}{2^{\tau_M+6}}}}\right\},\newline
$$
  
\n
$$
\delta_M := \frac{\sqrt{\gamma R}}{\sqrt{2^{17}e12^{\tau_M}(4^6e^{27})^{2+\frac{q}{q}}|f|_R}}{R},\newline
$$
  
\n
$$
C_1(M) := \frac{1}{2^9\delta_M^3}, \quad C_2(M) := \frac{2^{\tau_M+
$$

<span id="page-37-0"></span>**Theorem 8.1** (Birkhoff Normal Form). *Under the same assumptions of Lemma [8.1](#page-36-4) the following holds. Consider the Hamiltonian* [\(1.21\)](#page-7-3)*, assuming, only in the case* M, *that f does not depend on x (momentum conservation). Then for any*  $0 < r \le r$  *there exists two close to identity invertible symplectic change of variables*

<span id="page-37-1"></span>
$$
\Psi, \Psi^{-1}: \quad B_r(\mathbf{h}_w) \mapsto \mathbf{h}_w, \quad \sup_{|u|_w \le r} |\Psi^{\pm 1}(u) - u|_w \le C_1 r^3 \le \frac{1}{8} r, \Psi \circ \Psi^{-1} u = \Psi^{-1} \circ \Psi u = u, \quad \forall u \in B_{\frac{7}{8}r}(\mathbf{h}_w)
$$
\n(8.3)

*such that in the new coordinates*

$$
H\circ\Psi=D_{\omega}+Z+R,
$$

*for suitable majorant analytic Hamiltonians Z,*  $R \in A_r(h_w)$ *, Z*  $\in K$ *, <i>satisfying the estimate*

<span id="page-38-2"></span>
$$
\sup_{|u|_w \le r} |X_{\underline{Z}}|_w \le C_2 r^3, \quad \sup_{|u|_w \le r} |X_{\underline{R}}|_w \le C_3 r^{2N+3}, \tag{8.4}
$$

*XZ (resp. XR), being the hamiltonian vector field generated by the the majorant of Z (resp. R). Moreover, in the case* M*, R preserves momentum.*

*Proof.* We use Theorem [1.3](#page-10-0) with  $G \rightarrow P$ .

G) Setting

<span id="page-38-1"></span>
$$
\eta = \eta_{\mathcal{G}} := \min\left\{\frac{a-a}{2}, s\right\}, \quad \bar{r} := \frac{\sqrt{R}}{C_{\text{alg}}(p)},\tag{8.5}
$$

we have that

$$
|P|_{\bar{r},\eta,\mathbf{w}_0} = |P|_{\bar{r},\eta,\mathbf{w}(p,s-\eta,a)} \stackrel{(5.32)}{\leq} C_{\text{Nem}}(p,s-\eta,a-a-\eta)|f|_{\mathbf{a},R}. \tag{8.6}
$$

By [\(1.32\)](#page-11-3)  $r_{\star} \ge \delta_{\rm G} \ge \delta_{\rm G}$  (see "Appendix A"). Then, recalling (1.32) and Lemma [8.1,](#page-36-4) one can verify that

$$
\widehat{r} \ge r(G), \widehat{C}_1 \le C_1(G), \widehat{C}_2 \le C_2(G), \widehat{C}_3 \le C_3(G)
$$

S) Set

$$
\eta := a/2, \quad \bar{r} := \frac{\sqrt{R}}{C_{\text{alg}}(p - \text{N}\tau_{\text{S}})}.\tag{8.7}
$$

Then Assumption [1](#page-10-3) is satisfied by Lemma [8.1](#page-36-4) with the same choice of N,  $w_0$ , w. We have that

$$
|P|_{\bar{r}, \eta, w_0} = |P|_{\bar{r}, a/2, w(p - N\tau_S)} \stackrel{(5.32)}{\leq} C_{\text{Nem}}(p - N\tau_S, 0, a/2)|f|_{a, R}. \tag{8.8}
$$

For the various constants we refer to "Appendix A". Recalling  $\tau_s = \tau$ , we note that

<span id="page-38-3"></span>
$$
C_{\text{mon}}(4N, \eta/N, \tau) = 2^{2\tau + 1} \tau^{\tau} N^{\tau} \max \{4, (1/2\eta)\}^{\tau} = 2(4\tau \max \{4, (1/2\eta)\} N)^{\tau},
$$
  
\n
$$
C_2(4N, \eta/N, \tau) = 2e^{27(2+q)} (12\tau \max \{4, (1/2\eta)\} N)^{3\tau}
$$
  
\n
$$
C_2(4N, \eta/N, N\tau) = 2e^{27(2+q)} (6N\tau)^{3N\tau} \max \{ (8N)^{3N\tau}, (N/\eta)^{3N\tau} \}
$$
  
\n
$$
= 2e^{27(2+q)} (12\tau \max \{4, (2\eta)^{-1}\} N^2)^{3N\tau},
$$
\n(8.9)

we have that for  $N \geq 3$ 

<span id="page-38-0"></span>
$$
C_{\text{mon}}(4N, \eta/N, \tau)C_2(4N, \eta/N, \tau) \le \sqrt{C_2(4N, \eta/N, N\tau)}.
$$
 (8.10)

By [\(1.32\)](#page-11-3)

$$
r_{\star} \geq \frac{\sqrt{\gamma R}}{C_{\text{alg}}(p - \text{N}\tau_{\text{S}})\sqrt{2^{11}eC_{\text{Nem}}(p - \text{N}\tau_{\text{S}}, 0, \text{a}/2)|f|_{\text{a},R}}} \geq d_{\text{S}}.
$$

Then, recalling [\(1.32\)](#page-11-3) and [\(8.10\)](#page-38-0) one has  $\hat{\tau} > \tau(s)$ . Moreover (recall [\(1.34\)](#page-11-1))

<span id="page-39-4"></span>
$$
\widehat{C}_1 \le C_1(S), \qquad \widehat{C}_2 \le C_2(S), \qquad \widehat{C}_3 \le C_3(S). \tag{8.11}
$$

Finally the last inequality in  $(8.3)$  follows from the second bound in  $(1.33)$ . M) Set

$$
\eta := 0, \quad \bar{r} := \frac{\sqrt{R}}{2^{\tau_M/2} C_{\text{alg,M}} (p - \tau_M N)} \ge \sqrt{\frac{R}{10 \cdot 2^{\tau_M}}}.
$$
\n(8.12)

Then Assumption [1](#page-10-3) is satisfied by Lemma [8.1,](#page-36-4) case M, with the same choice of  $N, w_0, w$ . We have that

$$
|P|_{\bar{r},0,w_0} = ||P||_{\bar{r},p-\tau_M N} \stackrel{(5.33)}{\leq} 2|f|_R. \tag{8.13}
$$

By [\(1.32\)](#page-11-3)

$$
r_{\star} \geq \frac{\sqrt{\gamma R}}{\sqrt{2^{\tau_{\text{M}}+17} |f|_R}}.
$$

Then

$$
\widehat{r} \ge r(M), \quad \widehat{C}_1 \le C_1(M), \quad \widehat{C}_2 \le C_2(M), \quad \widehat{C}_3 \le C_3(M).
$$

 $\Box$ 

### <span id="page-39-0"></span>**9. Gevrey Stability. Proof of Theorem [1.1](#page-4-1)**

Actually we prove of Theorem [1.1](#page-4-1) for the slightly longer stability time  $|t| \leq \frac{2^4 e \delta_0^2}{\gamma \delta^2}$  $e^{\left(\ln \frac{\delta_G}{\delta}\right)^{1+\theta/4}}$ , where  $\delta_{\rm G} > \delta_{\rm G}$  (recall "Appendix A"). We set

 $r := 2\delta$ 

and choose

<span id="page-39-2"></span>
$$
N(r) := \left[ \left( 2 \ln \frac{2 \delta_G}{r} \right)^{\theta/4} \right] = \left[ \left( 2 \ln \frac{\delta_G}{\delta} \right)^{\theta/4} \right].
$$
 (9.1)

Recalling [\(8.5\)](#page-38-1) by Corollary [5.2](#page-25-4) solutions of the PDE [\(1.1\)](#page-1-3) in the space  $h_{p,s,a}$ , correspond, by Fourier identification  $(1.19)$ , to orbits of the Hamiltonian System  $(1.21)$  in the space

$$
h_w \quad \text{with} \quad w_j = e^{a|j|+s\langle j\rangle^\theta} \langle j\rangle^p.
$$

An initial datum  $u_0$  satisfying  $|u_0|_{p,s,a} \leq \delta$  corresponds to<sup>23</sup>  $u_0 \in h_w$  with  $|u_0|_w \leq \delta$ . We claim that  $r \leq 2\delta_{\rm G}$  implies

<span id="page-39-3"></span>
$$
\frac{r\mathrm{N}e^{C_1\left(\frac{\mathrm{N}}{\eta_{\mathrm{G}}}\right)^{\frac{3}{\theta}}}}{2\delta_{\mathrm{G}}}\leq 1.\tag{9.2}
$$

<span id="page-39-1"></span><sup>&</sup>lt;sup>23</sup> We still denote it by  $u_0$ .

Indeed we have

$$
N(r) \ge N_G := \max\left\{\frac{16(4C_1)^{\theta}}{\eta_G^3}, 2^{\frac{2\theta+4}{4-\theta}}\right\}
$$

and by [\(9.1\)](#page-39-2)  $r \le 2\delta_{\rm G}e^{-\frac{1}{2}(\mathbb{N}(r)/2)^{4/\theta}}$  and [\(9.2\)](#page-39-3) follows if we show that the function

$$
N \rightarrow e^{-\frac{1}{2}(N/2)^{4/\theta}}Ne^{\mathcal{C}_1\left(\frac{N}{\eta_G}\right)^{\frac{3}{\theta}}}
$$

is  $\leq 1$  for  $N \geq N_G$ . This is true since the function is decreasing for  $N \geq N_G$  and is  $\leq 1$ for  $N = N_G$ . This proves the claim [\(9.2\)](#page-39-3).

Then we apply Theorem  $8.1$  in the case G. Recalling  $(8.1)$ , by  $(8.4)$  and  $(9.2)$ 

$$
C_3(G)r^{2(N+1)} \leq \frac{\gamma r^2}{2^9 e \delta_G^2} \left(\frac{r}{2\delta_G}\right)^{N(r)} = \frac{\gamma \delta^2}{2^7 e \delta_G^2} \left(\frac{\delta}{\delta_G}\right)^{N(r)} \leq \frac{\gamma \delta^2}{2^7 e \delta_G^2} e^{-\left(\ln \frac{\delta_G}{\delta}\right)^{1+\theta/4}},
$$

since  $N(r) \geq \left(\ln \frac{2\delta_G}{r}\right)^{\theta/4} = \left(\ln \frac{\delta_G}{\delta}\right)^{\theta/4}$ . We deduce the stability time by applying Lemma 5.1

#### <span id="page-40-0"></span>**10. Sobolev Stability**

<span id="page-40-1"></span>Before proving Proposition [1.1](#page-5-0) we add a comment on the optimality of condition [\(1.10\)](#page-5-1).

*Remark 10.1.* We construct a finite dimensional Hamiltonian, which is a reduction of [\(1.1\)](#page-1-3) to a finite number of Fourier indices and which exhibits fast drift in a time of order 1. For instance, consider

$$
H(u_1, u_j) := (1 + V_1)|u_1|^2 + (j^2 + V_j)|u_j|^2 + e^{-a_j} \operatorname{Re}(|u_1|^2 u_1 \bar{u}_j),
$$

which is a finite dimensional model for [\(1.1\)](#page-1-3) with  $f(x, |u|^2) = e^{-a} \cos((\jmath-1)x)|u|^2$ . Consider now the initial datum  $u(0) = (u_1(0), u_j(0)) = (\delta/4, |j|^{-p}\delta/4)$ , which clearly has  $H_p$  norm  $\lt \delta$ . A direct computation shows that in a time *T* of order 1, the Sobolev norm of  $u(T)$  is of order

 $\delta^3 e^{-a}$  *j*  $\mu$ <sup>*j*</sup> hence greater than 4 $\delta$  if  $\delta^2 e^{-a}$  *j p* is large. Maximizing on *j* we get a constraint of the form  $\delta^2 e^{-p} (a^{-1} p)^p < 1$ .

Of course this pathological "fast diffusion" phenomenon comes from the fact that *f* is NOT traslation invariant (and hence *H* does not preserve momentum). Actually, restricting to translation invariant Hamiltonians would not result in signficantly weaker constraints on the smallness of  $\delta$  w.r.t. *p*. This can be seen in the following example. Consider the familiy of Hamiltonians (in three degrees of freedom)

$$
K^{(j)} := V_0|u_0|^2 + (1+V_1)|u_1|^2 + (j^2+V_j)|u_j|^2 + \text{Re}(\bar{u}_0^{j-1}u_1^j\bar{u}_j)
$$

with the constants of motion

$$
L = |u_0|^2 + |u_1|^2 + |u_j|^2, \quad M = |u_1|^2 + j|u_j|^2.
$$

Following the same approach as in the previous example one shows that  $|u_j|^2$  can have a drift of order  $j^{-p}\delta^{2j}$  in a time *T* of order 1. This means that the Sobolev norm of *u*(*T*) is of order

 $\delta^{2j} j^p$ . Maximizing on *j* we get a constraint of the form  $\delta e^{p^{1-}} < 1$ .

*Proof of Proposition [1.1.](#page-5-0)* As before we set  $r := 2\delta$ . An initial datum  $u_0$  satisfying  $|u_0|_{L^2} + |\partial_x^p u_0|_{L^2} \le \delta$  corresponds to  $24 u_0 \in h_{w(p)}$  with  $|u_0|_p \le \delta$  by [\(5.28\)](#page-24-1). We apply the Birkhoff Normal Form Theorem [8.1](#page-37-0) in the case S (recall that  $N = \left\lfloor \frac{p-1}{\tau_s} \right\rfloor$ ). Recalling the definition of  $r(S)$  in [\(8.2\)](#page-37-2), we verify that, for any  $N \ge 1$ 

<span id="page-41-1"></span>
$$
\delta_{\rm S}(\mathbf{k}_{\rm S}p)^{-3p} \le \frac{\mathbf{d}_{\rm S}}{2\sqrt{\rm N}\mathcal{C}_2(\mathbf{4}\mathbf{N}, \mathbf{a}/2\mathbf{N}, \mathbf{N}\tau)}.\tag{10.1}
$$

Indeed

$$
\frac{d_S}{2\sqrt{NC_2(4N, a/2N, NT)}} \stackrel{(8.9)}{=} \frac{d_S}{2\sqrt{2}e^{27(2+q)/2}} \frac{1}{\sqrt{N}} \frac{1}{(\sqrt{12\tau}N \max\{2, a^{-1/2}\})^{3N\tau}}
$$

$$
\geq \frac{d_S\sqrt{\tau}}{2\sqrt{2}e^{27(2+q)/2}} \left(\sqrt{\frac{12}{\tau}} \max\{2, a^{-1/2}\}\right)^{-3(p-1)} (p-1)^{-3(p-1)-1/2}
$$

$$
= \delta_S(k_S)^{-3p}(p-1)^{-3(p-1)-1/2}
$$

setting

$$
\delta_{\rm S} = \frac{\mathrm{d}_{\rm S}\sqrt{\tau}}{2\sqrt{2}e^{27(2+q)/2}}\mathrm{k}_{\rm S}^3,
$$

[\(10.1\)](#page-41-1) follows by verifying that  $\delta_S \leq \delta_S$  and noting that  $p^{-3p} < (p-1)^{-3(p-1)-1/2}$ for  $p > 1$ .

By  $(8.11)$  and  $(8.9)$ 

$$
C_{3}(S)(2\delta)^{2(N+1)}
$$
\n
$$
= \delta^{2} \frac{2^{10}C_{\text{Nem}}(p - \text{N}\tau, 0, \alpha/2)|f|_{a,R}}{eR}
$$
\n
$$
\left(\frac{\text{N}C_{\text{mon}}(4\text{N}, \alpha/2\text{N}, \tau)C_{2}(4\text{N}, \alpha/2\text{N}, \tau)\delta^{2}}{d_{S}^{2}}\right)^{N}
$$
\n
$$
= \delta^{2} \frac{2^{10}C_{\text{Nem}}(p - \text{N}\tau, 0, \alpha/2)|f|_{a,R}}{eR}
$$
\n
$$
\left(\frac{4e^{27(2+q)}3^{3\tau}(4\max\{4, (1/a)\})^{4\tau}\delta^{2}}{\tau d_{S}^{2}}\right)^{N}(\tau\text{N})^{\frac{4\tau+1}{\tau}(\tau\text{N})}
$$
\n
$$
\leq \delta^{2} \frac{2^{10}C_{\text{Nem}}(p - \text{N}\tau, 0, \alpha/2)|f|_{a,R}}{eR}
$$
\n
$$
\left(\frac{4e^{27(2+q)}3^{3\tau}(4\max\{4, (1/a)\})^{4\tau}\delta^{2}}{\tau d_{S}^{2}}\right)^{\frac{p-1}{\tau}-1}(p-1)^{\frac{4\tau+1}{\tau}(p-1)}
$$

(remember that  $N = [(p - 1)/\tau]$ ). Then, noting that  $(p - 1)^{\frac{4\tau + 1}{\tau}(p-1)} < p^{5p}$  for  $p > 1$ (recall that  $\tau \geq 15$ ), we get

$$
C_3(S)(2\delta)^{2(N+1)} < \frac{1}{8\mathbf{T}_S} p^{5p} \left(\frac{\delta}{\delta_S}\right)^{2\frac{p-1}{\tau}}.
$$

<span id="page-41-0"></span><sup>&</sup>lt;sup>24</sup> We still denote it by  $u_0$ .

We conclude by applying Lemma [5.1](#page-24-2) and  $(5.28)$ 

$$
|u(t)|_{L^2}+|\partial_x^p u(t)|_{L^2}\leq 2|u(t)|_p\leq 4\delta,\qquad \forall |t|\leq \mathrm{T}_{\mathrm{S}}p^{-5p}\left(\frac{\delta_{\mathrm{S}}}{\delta}\right)^{\frac{2(p-1)}{\mathrm{T}_{\mathrm{S}}}},
$$

proving  $(1.11)$ .  $\Box$ 

*Proof of Theorem [1.2.](#page-6-0)* It is similar to the previous case but now we consider

$$
h_w \quad \text{with} \quad w_j = \lfloor j \rfloor^p \quad (\text{and } |\cdot|_w = ||\cdot||_p).
$$

We set  $r = 4\delta$ , an initial datum  $u_0$  satisfying  $2^p |u_0|_{L^2}$ ,  $|u_0|_{L^2} + |\partial_x^p u_0|_{L^2} \le \delta$  corresponds to  $u_0 \in h_w$  with  $||u_0||_p \le 2\delta$  by [\(5.30\)](#page-25-0). Now we can apply the Birkhoff Normal Form Theorem [8.1](#page-37-0) with  $N = \left[\frac{p-1}{\tau_1}\right]$ 

$$
4\delta \le 4 \frac{\delta_{\mathbb{M}}}{\sqrt{p}} \le \mathbf{r}(\mathbb{M}) = \min \left\{ \frac{\delta_{\mathbb{M}}}{\sqrt{\mathbb{N}}}, \frac{\sqrt{R}}{\sqrt{2^{\tau_{\mathbb{M}}+4}}} \right\}.
$$
 (10.2)

Proceeding as in the case S and noting that now

$$
8C_3(M)(4\delta)^{2(N+1)} = 2^{\tau_M + 12} \frac{|f|_R}{R} \left(\frac{N\delta^2}{2\delta_M^2}\right)^N \delta^2
$$
  

$$
\leq 2^{\tau_M + 13} \delta_M^2 \frac{|f|_R}{R} \left(\frac{p-1}{\tau_1}\right)^{\frac{p-1}{\tau_1}} \left(\frac{\delta^2}{2\delta_M^2}\right)^{\frac{p-1}{\tau_1}}
$$
  

$$
\leq \frac{1}{T_M} \left(\frac{(p-1)\delta^2}{8\delta_M^2}\right)^{\frac{p-1}{\tau_M}}.
$$

Finally by Corollary [5.1](#page-24-2) and [\(5.30\)](#page-25-0) we get

$$
|u(t)|_{L^2}+|\partial_x^p u(t)|_{L^2}\leq 2||u(t)||_p\leq 8\delta,\qquad \forall |t|\leq \mathrm{T}_M\left(\frac{8\delta_M^2}{(p-1)\delta^2}\right)^{\frac{p-1}{\tau_1}},
$$

proving  $(1.13)$ .  $\Box$ 

*Proof of Corollary* [1.1.](#page-6-2) In case S we start by noticing that for  $3p \ln(k_S p) \leq \ln(\delta_S/\delta)$ the function  $\frac{T_S}{\delta^2}(p)^{-5p} \left(\frac{\delta_S}{\delta}\right)^{\frac{2(p-1)}{\tau_S}}$  is increasing in *p*.

Let us check that  $p(\delta)$  defined in [\(1.14\)](#page-6-5) satisfies [\(1.10\)](#page-5-1) and is  $\geq 3\tau_s + 1$  namely, passing to the logarithms and setting  $y := \ln(\delta_s/\delta)$ , we have to check that  $\frac{y}{\ln(y)} > 6\tau_s$ and  $3p \ln(k_S p) \leq y$ . The first bound follows from the definition of  $\delta_S$ . For the second, we have

$$
3p \ln(\mathbf{k}_{\rm S} p) \le 3\left(1 + \frac{1}{6} \frac{y}{\ln(y)}\right) \left(\ln(\mathbf{k}_{\rm S}) + \ln(1 + \frac{1}{6} \frac{y}{\ln(y)})\right) \le y
$$

provided that $25$ 

$$
y \ge \max\{k_{\rm S}, 40\}.
$$

Now we have to show that

$$
\mathrm{T}_{\mathrm{S}}e^{-\frac{\ln^2(\delta_{\mathrm{S}}/\delta)}{4\tau\ln\ln(\delta_{\mathrm{S}}/\delta)}} \leq \mathrm{T}_{\mathrm{S}}(p)^{-5p} \left(\frac{\delta_{\mathrm{S}}}{\delta}\right)^{\frac{2(p-1)}{\tau}}
$$

wich amounts to

$$
e^{\frac{y^2}{4\tau \ln y}}(p)^{5p}e^{-\frac{2(p-1)}{\tau}y} \le 1
$$

or equivalently

$$
\frac{y^2}{4\tau \ln y} + 5p \ln(p) - \frac{2(p-1)}{\tau} y \le 0.
$$

Assuming  $\frac{y}{\ln y} > 6$ , we have  $1 + \frac{y}{6 \ln y} - \frac{7}{6} < p < \frac{y}{3 \ln y}$  we get

$$
\frac{y^2}{4\tau \ln y} + 5p \ln(\mathbf{K}_{\mathbf{S}}p) - \frac{2(p-1)}{\tau}y \le \frac{y^2}{4\tau \ln y} \n+ \frac{5}{3} \frac{y}{\ln(y)} \ln(\frac{y}{3\ln(y)}) - 2y(\frac{1}{6\tau} \frac{y}{\ln(y)} - \frac{1}{6}) \n\le -\frac{y^2}{12\tau \ln y} + 2y < 0
$$

if  $\frac{y}{\ln(y)} > 24\tau > 6$ . Note that the last inequality holds if  $y \ge 24\tau^2$  (recall that  $\tau \ge 15$ ). Recollecting the condition that *y* has to satisfy is

$$
y \ge \max\{k_{\rm S},\, 24\tau^2\},\
$$

namely  $\delta \leq \delta_{\rm S}$ .

M) Since we are assuming  $\delta \leq \delta_M$  we have that p defined in [\(1.16\)](#page-6-6) satisfies  $p > 1+3\tau_M$ . Moreover by  $(1.16)$ , the bound  $(1.12)$  holds. Then Theorem [1.2](#page-6-0) applies and  $(1.17)$  follows directly by  $(1.13)$ .

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<span id="page-43-0"></span><sup>25</sup> Note that the function

$$
y \mapsto y - 3\left(1 + \frac{1}{6}\frac{y}{\ln(y)}\right)\left(\ln y + \ln(1 + \frac{1}{6}\frac{y}{\ln(y)})\right)
$$

is positive for  $y > 40$ .

# <span id="page-44-0"></span>**Part 3. Appendices**

# <span id="page-44-1"></span>**Appendix A. Constants**

In this subsection are listed all the constants appearing along the paper. We first introduce some auxiliary constants. Given  $t, \sigma, \zeta > 0$ ,  $p > 1/2$ ,  $0 < \theta < 1$ ,  $s, q > 0$ , we set<sup>26</sup>

$$
C_{\text{alg}}(p) := 2^{p} \Big( \sum_{i \in \mathbb{Z}} \langle i \rangle^{-2p} \Big)^{1/2},
$$
  
\n
$$
C_{\text{alg},M}(p) := \sqrt{2} \sqrt{2 + \frac{2p+1}{2p-1}},
$$
  
\n
$$
C_{\text{Nem}}(p, s, t) := C_{\text{alg}}(p) \Big( e^{s} + \sup_{x \geq 1} x^{p} e^{-tx + sx^{\theta}} \Big),
$$
  
\n
$$
C_{\text{mon}}(t, \sigma, p) := 2^{p+1} p^{p} \max \Big\{ (2t)^{p}, \sigma^{-p}, 1 \Big\},
$$
  
\n
$$
C_{1} := 28 \theta^{-1} (q + 3) \Big( \frac{2^{3} \cdot 13(q + 3)}{\theta (1 - \theta)} \Big)^{\frac{2}{\theta}}
$$
  
\n
$$
\Big( \ln \Big( \frac{2^{3} \cdot 13(q + 3)}{\theta (1 - \theta)} \Big)^{\frac{2}{\theta} + 1}, \quad C_{*} = 13/(1 - \theta),
$$
  
\n
$$
C_{2}(t, \sigma, \zeta) := e^{27(2+q)} C_{\text{mon}}(t, \sigma, 3\zeta),
$$
  
\n
$$
\tau := \tau_{0}(2 + q), \quad \tau_{0} := 15/2,
$$
  
\n
$$
\tau_{1} := 2 \Big( \tau_{0} + \frac{3}{2 \ln 2} \Big) (2 + q).
$$

Here are the constants appearing in Theorem [1.1:](#page-4-1)

$$
\delta_{\mathbf{G}} := \min \left\{ \frac{\sqrt{R}}{4C_{\text{alg}}(p)}, \ \delta_{\mathbf{G}}e^{-\left(\max \left\{16(4C_1)^{\theta} \eta_{\mathbf{G}}^{-3}, 2^{\frac{2\theta+4}{4-\theta}}\right\}\right)^{4/\theta}}\right\},
$$
\n
$$
\mathbf{T}_{\mathbf{G}} := \frac{2^4 e \delta_{\mathbf{G}}^2}{\gamma}, \quad \text{where}
$$
\n
$$
\delta_{\mathbf{G}} := \frac{\sqrt{\gamma R}}{C_{\text{alg}}(p)\sqrt{2^{11}eC_{\text{Nem}}(p, s - \eta_{\mathbf{G}}, \mathbf{a} - a - \eta_{\mathbf{G}})|f|_{\mathbf{a},R}}},
$$
\n
$$
\eta_{\mathbf{G}} := \min \left\{\frac{\mathbf{a} - a}{2}, s\right\}.
$$

<span id="page-44-2"></span> $\overline{^{26}$  Regarding  $C_{\text{Nem}}$  note that

$$
\sup_{x\geq 1} x^p e^{-tx+sx^{\theta}} \leq \exp\left( (1-\theta) \left( \frac{s}{t^{\theta}} \right)^{\frac{1}{1-\theta}} \right) \max\left\{ \frac{p}{e(1-\theta)t}, e^{-\frac{t(1-\theta)}{p}} \right\}^p.
$$

Here are the constants appearing in Proposition [1.1](#page-5-0)

$$
\tau_{\rm S} := \tau = \frac{15}{2} (2+q),
$$
\n
$$
\delta_{\rm S} = \min \left\{ \frac{\sqrt{3^{\tau} \tau \gamma R}}{2^{10} e^{2\tau} (k_{\rm S} \tau)^{4\tau} \sqrt{C_{\rm Nem}(p - N\tau_{\rm S}, 0, a/2) |f|_{a,R}}}, \frac{\sqrt{R}}{2^{\tau+5}} \right\}
$$
\n
$$
k_{\rm S} := \sqrt{\frac{12}{\tau}} \max \left\{ 2, a^{-1/2} \right\},
$$
\n
$$
\tau_{\rm S} = \frac{2^6 e^{2\tau} 3^{3\tau} (4 \max \{4, (1/a)\})^{4\tau}}{\tau \gamma}.
$$

Here are the constants in Theorem [1.2](#page-6-0)

$$
\tau_{\mathbb{M}} := \tau_1 = \left(15 + \frac{3}{\ln 2}\right)(2 + q), \qquad \delta_{\mathbb{M}} := \min\left\{\frac{\sqrt{\tau_1} \delta_{\mathbb{M}}}{4}, \sqrt{\frac{R}{2^{\tau_1 + 8}}}\right\},\newline
$$

$$
T_{\mathbb{M}} := \frac{R}{2^{\tau_1 + 13} \delta_{\mathbb{M}}^2 |f|_R} = \frac{2^3 e^{6\tau_1} (4^6 e^{27})^{2 + q}}{\gamma};
$$

where, recalling [8.2,](#page-37-2)

$$
\delta_{\rm M} = \frac{\sqrt{\gamma R}}{\sqrt{2^{17} e 12^{\tau_1} (4^6 e^{27})^{2+q} |f|_R}}.
$$

Here are the constants appearing in Corollary [1.1:](#page-6-2)

$$
\bar{\delta}_{\rm S} := \delta_{\rm S} e^{-\max\{k_{\rm S}, 24\tau^2\}}, \quad \bar{\delta}_{\rm M} := \frac{\delta_{\rm M}}{4\tau_1}.
$$

#### <span id="page-45-0"></span>**Appendix B. Proofs of the Main Properties of the Norms**

<span id="page-45-1"></span>**Lemma B.1.** *Let*  $0 < r_1 < r$ *. Let*  $E$  *be a Banach space endowed with the norm*  $|\cdot|_E$ *. Let*  $X : B_r \to E$  *a vector field satisfying* 

$$
\sup_{B_r} |X|_E \le \delta_0.
$$

*Then the flow*  $\Phi(u, t)$  *of the vector field*<sup>[27](#page-45-2)</sup> *is well defined for every* 

$$
|t| \leq T := \frac{r - r_1}{\delta_0}
$$

*and*  $u \in B_{r_1}$  *with estimate* 

$$
|\Phi(u,t) - u|_E \le \delta_0 |t|, \quad \forall |t| \le T.
$$

<span id="page-45-2"></span><sup>&</sup>lt;sup>27</sup> Namely the solution of the equation  $\partial_t \Phi(u, t) = X(\Phi(u, t))$  with initial datum  $\Phi(u, 0) = u$ .

*Proof.* Fix  $u \in B_{r_1}$ . Let us first prove that  $\Phi(u, t)$  exists  $\forall |t| \leq T$ . Otherwise there exists a time<sup>28</sup> 0 < *t*<sub>0</sub> < *T* such that  $|\Phi(u, t)|_E$  < *r* for every  $0 \le t < t_0$  but  $|\Phi(u, t_0)|_E = r$ . Then, by the fundamental theorem of calculus

$$
\Phi(u, t_0) - u = \int_0^{t_0} X(\Phi(u, \tau))d\tau.
$$
 (B.1)

Therefore

$$
r - r_1 \le |\Phi(u, t_0)|_E - |u|_E \le |\Phi(u, t_0) - u|_E \le \int_0^{t_0} |X(\Phi(u, \tau))|_E d\tau \le \delta_0 t_0
$$
  
<  $\delta_0 T = r - r_1$ ,

which is a contradiction. Finally, for every  $|t| < T$ ,

$$
|\Phi(u,t)-u|_E \leq \left| \int_0^t |X(\Phi(u,\tau))|_E d\tau \right| \leq \delta_0 |t|.
$$

*Proof of Lemma [2.1.](#page-13-1)* For brevity we set, for every  $r' > 0$ 

$$
|\cdot|_{r'}:=|\cdot|_{r',\eta,w}.
$$

We use Lemma [B.1,](#page-45-1) with  $E \to h_w$ ,  $X \to X_S$ ,  $\delta_0 \to (r+\rho)|S|_{r+\rho}$ ,  $r \to r+\rho$ ,  $r_1 \to r$ ,  $T \to 8e$ . Then the fact that the time 1-Hamiltonian flow  $\Phi_S^1 : B_r(h_w) \to B_{r+\rho}(h_w)$  is well defined, analytic, symplectic follows, since for any  $\eta \geq 0$ 

$$
\sup_{u\in B_{r+\rho}(\Omega_w)}|X_S|_{\mathcal{H}_w}\leq (r+\rho)|S|_{r+\rho}<\frac{\rho}{8e}.
$$

Regarding the estimate  $(2.3)$ , again by Lemma [B.1](#page-45-1) (choosing  $t = 1$ ), we get

$$
\sup_{u\in B_r(h_w)}\left|\Phi_S^1(u)-u\right|_{h_w}\leq (r+\rho)|S|_{r+\rho}<\frac{\rho}{8e}.
$$

Estimates [\(2.4\)](#page-13-5), [\(2.5\)](#page-13-6), [\(2.6\)](#page-13-7) directly follow by [\(2.7\)](#page-13-3) with  $h = 0, 1, 2$ , respectively and  $c_k = 1/k!$ , recalling that by Lie series

$$
H \circ \Phi_S^1 = e^{ad_S} H = \sum_{k=0}^{\infty} \frac{ad_S^k H}{k!} = \sum_{k=0}^{\infty} \frac{H^{(k)}}{k!},
$$

where  $H^{(i)} := \text{ad}_S^i(H) = \text{ad}_S(H^{(i-1)}), H^{(0)} := H.$ Let us prove [\(2.7\)](#page-13-3). Fix  $k \in \mathbb{N}$ ,  $k > 0$  and set

$$
r_i := r + \rho (1 - \frac{i}{k}), \quad i = 0, ..., k.
$$

Note that, by the immersion properties of the norm (recall Remark [2.1\)](#page-12-3)

<span id="page-46-1"></span>
$$
|S|_{r_i} \le |S|_{r+\rho}, \qquad \forall i = 0, \dots, k. \tag{B.2}
$$

 $\Box$ 

<span id="page-46-0"></span><sup>&</sup>lt;sup>28</sup> We assume  $t_0$  positive, the negative case is analogous.

Noting that

<span id="page-47-0"></span>
$$
1 + \frac{kr_i}{\rho} \le k \left( 1 + \frac{r}{\rho} \right), \quad \forall i = 0, \dots, k,
$$
 (B.3)

by using  $k$  times  $(2.1)$  we have

$$
|H^{(k)}|_{r} = |\{S, H^{(k-1)}\}|_{r} \le 4(1 + \frac{kr}{\rho})|H^{(k-1)}|_{r_{k-1}}|S|_{r_{k-1}}
$$
  

$$
\stackrel{\text{(B.2)}}{\le} |H|_{r+\rho}|S|_{r+\rho}^{k} 4^{k} \prod_{i=1}^{k} (1 + \frac{kr_{i}}{\rho})^{\frac{(\mathbf{B.3})}{\leq}} |H|_{r+\rho} \left(4k\left(1 + \frac{r}{\rho}\right)|S|_{r+\rho}\right)^{k}.
$$

Then, using  $k^k \leq e^k k!$ , we get

$$
\left| \sum_{k \ge h} c_k H^{(k)} \right|_r \le \sum_{k \ge h} |c_k| |H^{(k)}|_r \le |H|_{r+\rho} \sum_{k \ge h} \left( 4e \left( 1 + \frac{r}{\rho} \right) |S|_{r+\rho} \right)^k
$$
  
= |H|\_{r+\rho} \sum\_{k \ge h} (|S|\_{r+\rho}/2\delta)^k \stackrel{(2.2)}{\le} 2|H|\_{r+\rho} (|S|\_{r+\rho}/2\delta)^h.

Finally, if *S* and *H* satisfy mass conservation so does each  $\text{ad}_{S}^{k}H$ ,  $k \geq 1$ , hence  $H \circ \Phi_{S}^{1}$ too.

*Proof of Lemma [3.1.](#page-14-2)* We first prove (i). It is easily seen that:

$$
X_{\underline{H}_{\eta}}^{(j)}(u) = \mathbf{i} \sum_{\alpha,\beta \in \mathbb{N}^{\mathbb{Z}}} |H_{\alpha,\beta}|\beta_j e^{\eta |\pi(\alpha-\beta)|}u^{\alpha}\bar{u}^{\beta-\epsilon_j}.
$$

Now

$$
|X_{\underline{H}_{\eta}}(u)|_{w} \leq |X_{\underline{H}_{\eta}}(\underline{u})|_{w}, \quad \underline{u} = (|u_j|)_{j \in \mathbb{Z}}
$$

hence, in evaluating the supremum of  $|X_{\underline{H}_n}|_w$  over  $|u|_w \leq r$  we ca restrict to the case in which  $u = (u_j)_{j \in \mathbb{Z}}$  has all real positive components. Hence

$$
|H|_{r,\eta,w}=r^{-1}\sup_{|u|_{w}\leq r}\left|\left(\sum_{\alpha=1}^{k}|H_{\alpha,\beta}|\beta_j e^{\eta|\pi(\alpha-\beta)|}|u|^{\alpha+\beta-\epsilon_j}\right)_{j\in\mathbb{Z}}\right|_{w}.
$$

Then

$$
|H|_{r,\eta,w} = \frac{1}{2r} \sup_{|u|_{w} \le r} \left| \left( W_{\eta}^{(j)}(u) \right)_{j \in \mathbb{Z}} \right|_{w}
$$
 (B.4)

where

$$
W_{\eta}^{(j)}(u) = \sum_{j=1}^{N} |H_{\alpha,\beta}| (\alpha_j + \beta_j) e^{\eta |\pi(\alpha - \beta)|} u^{\alpha + \beta - \epsilon_j},
$$

since, by the reality condition [1.23,](#page-8-4) we have

$$
\sum^* |H_{\alpha,\beta}|\beta_j e^{\eta |\pi(\alpha-\beta)|}u^{\alpha+\beta-\epsilon_j} = \sum^* |H_{\alpha,\beta}|\alpha_j e^{\eta |\pi(\alpha-\beta)|}u^{\alpha+\beta-\epsilon_j} = \frac{1}{2}W_{\eta}^{(j)}(u).
$$

By the linear map

$$
L_{r,w}: \ell^2 \to h_w, \qquad y_j \mapsto \frac{r}{w_j} y_j = u_j,
$$

the ball of radius 1 in  $\ell^2$  is isomorphic to the the ball of radius *r* in  $h_w$ , namely  $L_{r,w}(B_1(\ell^2)) = B_r(\mathbf{h}_w)$ . We have

$$
Y_H^{(j)}(y; r, \eta, w) = \frac{1}{2} W_{\eta}^{(j)}(L_{r,w} y).
$$

Then (i) follows.

In order to prove item (ii) we rely on the fact that, since we are using the  $\eta$ -majorant norm, the supremum over *y* in the norm is achieved on the real positive cone. Moreover, given *u*,  $v \in \ell^2$ , if

$$
|u_j|\leq |v_j|,\quad \forall j\in\mathbb{Z}
$$

then  $|u|_{\ell^2} \leq |v|_{\ell}$ 2.  $\Box$ 

*Proof of Lemma* [5.4.](#page-23-0) Let us look at the time evolution of  $|v(t)|_w^2$ . By construction and Cauchy-Schwarz inequality

$$
2|v(t)|_w \left| \frac{d}{dt} |v(t)|_w \right| = \left| \frac{d}{dt} |v(t)|_w^2 \right| = 2| \operatorname{Re}(v, \dot{v})_{h_w} | = 2| \operatorname{Re}(v, X_R)_{h_w} | \leq 2|v(t)|_w |X_{\underline{R}}|_w \leq 2r |v(t)|_w |R|_{r, \eta, w}
$$

as long as  $|v(t)|_w \leq r$ ; namely

<span id="page-48-1"></span>
$$
\left| \frac{d}{dt} |v(t)|_{w} \right| \le r |R|_{r, \eta, w} \tag{B.5}
$$

as long as  $|v(t)|_w \leq r$ .

Assume by contradiction that there exists a time<sup>29</sup>

$$
0
$$

such that

<span id="page-48-2"></span>
$$
\left| |v(t)|_{w} - |v_0|_{w} \right| < \frac{r}{8}, \quad \forall \, 0 \le t < T_0, \qquad \text{but} \quad \left| |v(T_0)|_{w} - |v_0|_{w} \right| = \frac{r}{8}. \tag{B.6}
$$

Then

$$
|v(t)|_w \le |v_0|_w + \frac{r}{8} < r
$$
  $\forall 0 \le t \le T_0.$ 

By  $(B.5)$  we get

$$
\left| |v(T_0)|_{w} - |v_0|_{w} \right| \le r |R|_{r, \eta, w} T_0 < \frac{r}{8},
$$

which contradicts [\(B.6\)](#page-48-2), proving [\(5.26\)](#page-24-3).  $\Box$ 

<span id="page-48-0"></span><sup>&</sup>lt;sup>29</sup> The case  $T_0 < 0$  is analogous.

*Proof of Lemma* [5.5.](#page-25-7) We first note that (see, e.g. Lemma 17 of  $\overline{[BDG10]}$  $\overline{[BDG10]}$  $\overline{[BDG10]}$ ) for  $p > 1/2$ and every sequence  $\{x_i\}_{i \in \mathbb{Z}}$ ,  $x_i \geq 0$ ,

$$
\left(\sum_{i\in\mathbb{Z}}x_i\right)^2\leq c\sum_{i\in\mathbb{Z}}\left(\frac{\langle i\rangle^p\langle j-i\rangle^p}{\langle j\rangle^p}x_i\right)^2,
$$

with  $c := 4^p \sum_{i \in \mathbb{Z}} \langle i \rangle^{-2p} = (C_{\text{alg}}(p))^2$ . Then

$$
\begin{split} |f \star g|_{p,s,a}^2 &\leq \sum_{j} e^{2s\langle j \rangle^{\theta}} e^{2a|j|} \langle j \rangle^{2p} \Big( \sum_{i} |f_i| |g_{j-i}| \Big)^2 \\ &\leq c \sum_{j} e^{2s\langle j \rangle^{\theta}} e^{2a|j|} \sum_{i} \langle i \rangle^{2p} \langle j - i \rangle^{2p} |f_i|^2 |g_{j-i}|^2 \\ &= c \sum_{i} e^{2s\langle i \rangle^{\theta}} e^{2a|i|} \langle i \rangle^{2p} |f_i|^2 \sum_{j} \langle j - i \rangle^{2p} e^{2s\langle j - i \rangle^{\theta}} e^{2a|j - i|} |g_{j-i}|^2 \\ &= c |f|_{p,s,a}^2 |g|_{p,s,a}^2. \end{split}
$$

Regarding the second estimate, we set

$$
\phi(i, j) := \frac{\lfloor j \rfloor}{\lfloor i \rfloor \lfloor j - i \rfloor}, \qquad \forall i, j \in \mathbb{Z}.
$$

Note that

<span id="page-49-0"></span>
$$
\phi(i, j) = \phi(j, i) = \phi(-i, -j).
$$
 (B.7)

We claim that

<span id="page-49-1"></span>
$$
\phi(i,j) \le 1. \tag{B.8}
$$

Indeed by [\(B.7\)](#page-49-0) we can consider only the case  $j \ge 0$ . Since  $\phi(-|i|, j) \le \phi(|i|, j)$  we can consider only the case  $i \ge 0$ . Again by [\(B.7\)](#page-49-0) we can assume  $j \ge i$ . In particular we can take  $j > i > 0$ , [\(B.8\)](#page-49-1) being trivial in the cases  $j = i$ ,  $i = 0$ . We have

$$
\phi(i+1, i) = \frac{i+1}{2\lfloor i \rfloor} \le \frac{3}{4}, \qquad \phi(j, 1) = \frac{j}{2(j-1)} \le 1.
$$

Then it remains also to discuss the case  $j - 2 \ge i \ge 2$ ; we have

$$
\phi(i, j) = \frac{j}{i(j - i)} = \frac{1}{i} + \frac{1}{j - i} \le 1,
$$

proving [\(B.8\)](#page-49-1).

For  $q \geq 0$  set

<span id="page-49-2"></span>
$$
c_q := \sup_{j \in \mathbb{Z}} \sum_{i \in \mathbb{Z}} (\phi(i, j))^q = \sup_{j \ge 0} \sum_{i \in \mathbb{Z}} (\phi(i, j))^q.
$$
 (B.9)

We claim that

<span id="page-49-3"></span>
$$
c_q \le 4 + 2\frac{q+1}{q-1} < \infty, \quad \forall q > 1.
$$
 (B.10)

Indeed, since  $\lfloor j \rfloor / \lfloor j + 1 \rfloor$  ≤ 1 and  $\lfloor j \rfloor / \lfloor j - 1 \rfloor$  ≤ 3/2 for *j* > 0, we have<sup>30</sup>

$$
c_q = \sup_{j\geq 0} \left( \frac{\lfloor j \rfloor^q}{2^{q-1} \lfloor j+1 \rfloor^q} + \frac{1}{2^{q-1}} + \frac{\lfloor j \rfloor^q}{2^{q-1} \lfloor j-1 \rfloor^q} + \sum_{i \leq -2, 2 \leq i \leq j-2, i \geq j+2} (\phi(i, j))^q \right)
$$
  

$$
\leq 2^{3-q} + \sup_{j\geq 0} \left( \sum_{i\geq 2} \frac{\lfloor j \rfloor^q}{i^q (j+i)^q} + \sum_{2 \leq i \leq j-2} \left( \frac{1}{i} + \frac{1}{(j-i)} \right)^q + \sum_{i\geq j+2} \frac{\lfloor j \rfloor^q}{i^q (i-j)^q} \right)
$$
  

$$
\leq 2^{3-q} + \sum_{i\geq 2} \frac{1}{i^q} + 2^{q-1} \sum_{2 \leq i \leq j-2} \left( \frac{1}{i^q} + \frac{1}{(j-i)^q} \right) + \sum_{i\geq j+2} \frac{1}{(i-j)^q}
$$
  

$$
\leq 4 + 2\frac{q+1}{q-1},
$$

using that  $(x + y)^q \le 2^{q-1}(x^q + y^q)$  for *x*,  $y \ge 0$  and that<sup>[31](#page-50-2)</sup>

$$
\sum_{i\geq 2} i^{-q} \leq \frac{q+1}{2^q(q-1)}.
$$

Note that for every  $q, q_0 \geq 0$  we have

<span id="page-50-3"></span>
$$
c_{q_0+q} \le c_{q_0} \tag{B.11}
$$

since

$$
c_{q_0+q} := \sup_{j \in \mathbb{Z}} \sum_{i \in \mathbb{Z}} (\phi(i, j))^{q_0} (\phi(i, j))^{q} \stackrel{\text{(B.8)}}{\leq} \sup_{j \in \mathbb{Z}} \sum_{i \in \mathbb{Z}} (\phi(i, j))^{q_0} = c_{q_0}.
$$

We now note that for  $p > 1/2$ ,  $j \in \mathbb{Z}$  and every sequence  $\{x_i\}_{i \in \mathbb{Z}}$ ,  $x_i \geq 0$ , we have by Cauchy-Schwarz inequality

$$
\left(\sum_{i\in\mathbb{Z}}x_i\right)^2=\left(\sum_{i\in\mathbb{Z}}(\phi(i,j))^p(\phi(i,j))^{-p}x_i\right)^2\leq c_{2p}\sum_{i\in\mathbb{Z}}\left((\phi(i,j))^{-p}x_i\right)^2,
$$

with  $c_{2p}$  defined in [\(B.9\)](#page-49-2). Using the above inequality we get

$$
||f \star g||_p^2 \le \sum_j |j|^{2p} \Big( \sum_i |f_i||g_{j-i}| \Big)^2
$$
  
\n
$$
\le c_{2p} \sum_j \sum_i |j|^{2p} |j-i|^{2p} |f_i|^2 |g_{j-i}|^2
$$
  
\n
$$
= c_{2p} \sum_i |j|^{2p} |f_i|^2 \sum_j |j-i|^{2p} |g_{j-i}|^2
$$
  
\n
$$
= c_{2p} ||f||_p^2 ||g||_p^2.
$$

The proof ends recalling  $(B.10)$ .  $\square$ 

<span id="page-50-2"></span><span id="page-50-1"></span><span id="page-50-0"></span><sup>30</sup> Note that the term  $\left(\frac{1}{i} + \frac{1}{(j-i)}\right)^q$  for  $j = 4$  and  $i = 2$  is 1 for every *q*. 31  $\sum_{i\geq 2} 1^{-q} \leq 2^{-q} + \int_2^{\infty} x^{-q} dx$ .

**Lemma B.2** (Nemitskii operators). *Let*  $p > 1/2$ . *(i) Fix s*  $> 0$ ,  $a_0 > 0$ . *Consider a sequence*  $F^{(d)} = (F_j^{(d)})$  $j ∈ \mathbb{Z}$  ∈  $h_{p,s,a_0}$ ,  $d ≥ 1$ , *such that* 

<span id="page-51-0"></span>
$$
\sum_{d=1}^{\infty} d|F^{(d)}|_{p,s,a_0} R^d < \infty
$$
 (B.12)

*for some*  $R > 0$ *. For*  $u = (u_j)_{j \in \mathbb{Z}}$  *let*  $\bar{u} = (\overline{u_{-j}})_{j \in \mathbb{Z}}$  and consider the Hamiltonian

$$
H(u) = \sum_{d=1}^{\infty} \left( F^{(d)} \star \underbrace{u \star \cdots \star u}_{d \text{ times}} \star \underbrace{\overline{u} \star \cdots \star \overline{u}}_{d \text{ times}} \right)_0.
$$

*For all*  $(\eta, a, r)$  *such that*  $\eta + a \le a_0$  *and*  $(C_{\text{alg}}(p)r)^2 \le R$ *, we have that*  $H \in$  $\mathcal{H}_{r,\eta}(\mathbf{h}_{p,s,a})$  *and* 

$$
|H|_{r,\eta,w(p,s,a)} \le r^{-1} \sum_{d=1}^{\infty} d|F^{(d)}|_{p,s,a_0} (C_{\text{alg}}(p)r)^{2d-1} < \infty.
$$

*(ii) Analogously if*  $F^{(d)}$  *are constants satisfying* 

<span id="page-51-2"></span>
$$
\sum_{d=1}^{\infty} d|F^{(d)}|R^d < \infty \tag{B.13}
$$

 $and$   $(C_{a \log M}(p)r)^2 \leq R$ , *then*  $H \in \mathcal{H}^{r,p}$  *with* 

<span id="page-51-1"></span>
$$
||H||_{r,p} \le 2^p r^{-1} \sum_{d=1}^{\infty} d|F^{(d)}|(C_{\text{alg,M}}(p)r)^{2d-1} < \infty.
$$
 (B.14)

*Proof.* (i) By definition the  $\eta$ -majorant Hamiltonian is

$$
\underline{H}_{\eta} = \sum_{d} \sum_{\substack{j_0, j_1, \dots, j_{2d} \\ j_0 + \sum_{i=1}^{2d} (-1)^i j_i = 0}} e^{\eta |\pi_{j_1, \dots, j_{2d}}|} |F_{j_0}^{(d)}| \overline{u_{j_1}} u_{j_2} \overline{u_{j_3}} \dots u_{j_{2d}}
$$

where

$$
\pi_{j_1,\dots,j_{2d}} = \sum_{i=1}^{2d} (-1)^i j_i = -j_0,
$$

hence

$$
\underline{H}_{\eta} = \sum_{d} \left( \underline{F}_{\eta}^{(d)} \star \underline{u \star \cdots \star u}_{d \text{ times}} \star \underline{\bar{u} \star \cdots \star \bar{u}}_{d \text{ times}} \right)_{0}, \quad \underline{F}_{\eta}^{(d)} := \left( e^{\eta |j|} |F_{j}^{(d)}| \right)_{j \in \mathbb{Z}}
$$

consequently

$$
X_{\underline{H}_{\eta}}^{(j)} = \sum_{d} d \left( \underbrace{F_{\eta}^{(d)} \star \underbrace{u \star \cdots \star u}_{d \text{ times}} \star \underbrace{\overline{u} \star \cdots \star \overline{u}}_{d-1 \text{ times}} \right)_{j}.
$$

Moreover

$$
|X_{\underline{H}_{\eta}}|_{p,s,a} \leq \sum_{d} d(C_{\text{alg}}(p))^{2d-1} |\underline{F}_{\eta}^{(d)}|_{p,s,a} (|u|_{p,s,a})^{2d-1}.
$$

Since

$$
|\underline{F}_{\eta}^{(d)}|_{p,s,a} = |F^{(d)}|_{s,a+\eta,p} \leq |F^{(d)}|_{p,s,a_0}
$$

we get

$$
|X_{\underline{H}_{\eta}}|_{p,s,a} \leq \sum_{d} d(C_{\text{alg}}(p))^{2d-1} |F^{(d)}|_{p,s,a_0} (|u|_{p,s,a})^{2d-1}.
$$

Therefore

$$
|H|_{r,p,\eta}^{(p,s,a,0)} = r^{-1} \left( \sup_{|u|_{p,s,a} < r} \left| X_{\frac{H}{\eta}} \right|_{p,s,a} \right) \le r^{-1} \sum_{d} d |F^{(d)}|_{p,s,a_0} (C_{\text{alg}}(p)r)^{2d-1} < \infty.
$$

(ii) The proof is analogous to point (i).  $\Box$ 

*Proof of Proposition [6.2.](#page-30-2)* We start by Taylor expanding *H* in homogeneous components. The majorant analiticity implies that for a homogeneous component of degree *d* one has

$$
|H^{(d)}|_{r, \eta, w(p,s,a)} \leq |H|_{r, \eta, w(p,s,a)}.
$$

Now let us consider the polinomial map (homogeneous of degree  $d-1$ )  $X_{H^{(d)}}:$  h<sub>p</sub>,*s*,*a* → h*p*,*s*,*a*; as is habitual we identify the polynomial map with the corresponding symmetric multilinear operator  $M^{(d-1)}: h_{p,s,a}^{d-1} \to h_{p,s,a}$ . Since we are in a Hilbert space, one has that

$$
|\underline{M}|_{p,s,a}^{\text{op}} := \sup_{\substack{u_1,\ldots u_{d-1} \in h_{p,s,a} \\ |u_i|_{p,s,a} \le 1}} |\underline{M}^{(d-1)}(u_1,\ldots,u_{d-1})|_{p,s,a} = \sup_{\substack{|u|_{p,s,a} \le 1 \\ |u|_{p,s,a} \le 1}} |\underline{M}^{(d-1)}(u,\ldots,u)|_{p,s,a}
$$

for all  $\eta \geq 0$ . Now let us compute the tame norm on a homogeneous component, i.e.

$$
\sup_{|u|_{p_0,s,a}\leq r-\rho} \frac{|\underline{M}^{(d-1)}(u^{d-1})|_{p,s,a}}{|u|_{p,s,a}} = \sup_{|u|_{p_0,s,a}\leq r-\rho} \frac{|N_p^{(d-1)}(u^{d-1})|_{p_0,s,a}}{|u|_{p,s,a}}
$$

where

$$
\underline{N_p}^{(d-1,j)}(u^{d-1}) = \langle j \rangle^{p-p_0} \sum_{j_1, \dots, j_{d-1}} |M_{j_1, \dots, j_{d-1}}^{(d-1,j)}|u_{j_1} \dots u_{j_{d-1}}
$$

$$
\mathbb{L}
$$

now setting  $\pi = \sum_i j_i - j$  we have

$$
\frac{N_p}{d-1}(u_1, \ldots, u_{d-1})
$$
\n
$$
\leq (d-1)\langle j \rangle^{p-p_0} \sum_{\substack{j_1, \ldots, j_{d-1}: \\ |j_1| \geq |j_i|}} |M_{j_1, \ldots, j_{d-1}}^{(d-1, j)}| u_{j_1} \ldots u_{j_{d-1}}
$$
\n
$$
\leq (d-1) \sum_{\substack{j_1, \ldots, j_{d-1}: \\ |j_1| \geq |j_i|}} \left( \sum_i \langle j_i \rangle + |\pi| \right)^{p-p_0} |M_{j_1, \ldots, j_{d-1}}^{(d-1, j)}| u_{j_1} \ldots u_{j_{d-1}}
$$
\n
$$
\leq (d-1)2^{p-p_0} C(\eta, p) \sum_{\substack{j_1, \ldots, j_{d-1}: \\ j_1, \ldots, j_{d-1}}} e^{\eta |\pi|} |M_{j_1, \ldots, j_{d-1}}^{(d-1, j)}| u_{j_1} \ldots u_{j_{d-1}}
$$
\n
$$
+ (d-1)2^{p-p_0} (d-1)^{p-p_0} \sum_{j_1, \ldots, j_{d-1}} |M_{j_1, \ldots, j_{d-1}}^{(d-1, j)}| \langle j_1 \rangle^{p-p_0} u_{j_1} \ldots u_{j_{d-1}}
$$

which means that for any  $|u|_{p_0,s,a} \leq r - \rho$  one has

$$
\begin{split} & \frac{|N_p^{(d-1)}(u^{d-1})|_{p_0,s,a}}{\leq (d-1)2^{p-p_0}C(\eta,p)|H^{(d)}|_{r-\rho,\eta,w(p_0,s,a)}|u|_{p_0,s,a}} \\ &+2^{p-p_0}(d-1)^{p-p_0+1}|\underline{M}|_{p_0,s,a}^{op}(r-\rho)^{d-2}|u|_{p,s,a}} \\ &\leq (d-1)2^{p-p_0}(C(\eta,p)+(d-1)^{p-p_0})(1-\frac{\rho}{r})^{d-2}|H|_{r,\eta,w(p_0,s,a)}|u|_{p,s,a}. \end{split}
$$

We conclude that

$$
\sup_{|u|_{p_0,s,a}\leq r} \frac{|X_H|_{p,s,a}}{|u|_{p,s,a}} \leq 2^{p-p_0} |H|_{r,\eta,w(p_0,s,a)} \sum_{d\geq 2} (d-1) \big(C(\eta,p) + (d-1)^{p-p_0}\big)(1-\frac{\rho}{r})^{d-2}
$$

and the thesis follows since the right hand side is convergent.

<span id="page-53-0"></span>**Appendix C. Small Divisor Estimates**

Let us start with two techincal lemmata.

**Lemma C.1.** *For p*,  $\beta > 0$  *and*  $x_0 \ge 0$  *we have that* 

$$
\max_{x \ge x_0} x^p e^{-\beta x} = \begin{cases} (p/\beta)^p e^{-p} & \text{if } x_0 \le p/\beta, \\ x_0^p e^{-\beta x_0} & \text{if } x_0 > p/\beta. \end{cases}
$$

**Lemma C.2.** *Let*  $0 < a < 1$  *and*  $x_1 \ge x_2 \ge \cdots \ge x_N \ge 2$ *. Then* 

$$
\frac{\sum_{1 \leq \ell \leq N} x_{\ell}}{\prod_{1 \leq \ell \leq N} x_{\ell}^a} \leq x_1^{1-a} + \frac{2}{ax_1^a}.
$$

*Proof.* By induction over *N*. It is obviously true for  $N = 1$ . Assume that it hols for *N* and prove it for  $N + 1$ .  $\Box$ 

*Proof of Lemma* [6.1.](#page-28-0) The fact that this [\(6.5\)](#page-28-5) holds true when  $\pi = 0$  is proven in [\[Bou96b](#page-63-23)] and [\[CLSY\]](#page-63-24). The bound [\(6.5\)](#page-28-5) is equivalent to proving

$$
\sum_{l\geq 1} \widehat{n}_l^{\theta} - 2\widehat{n}_1^{\theta} + \theta|\pi| - (2 - 2^{\theta}) \sum_{l\geq 3} \widehat{n}_l^{\theta} \geq 0. \tag{C.1}
$$

i.e.

$$
\sum_{l\geq 2} \widehat{n}_l^{\theta} - \widehat{n}_1^{\theta} + \theta |\pi| - (2 - 2^{\theta}) \sum_{l\geq 3} \widehat{n}_l^{\theta} \geq 0.
$$
 (C.2)

Inequality  $(C.2)$  then follows from

$$
f(|\pi|) := \sum_{l\geq 2} \widehat{n}_l^{\theta} - \left(|\pi| + \sum_{l\geq 2} \widehat{n}_l\right)^{\theta} + \theta|\pi| - (2 - 2^{\theta})\sum_{l\geq 3} \widehat{n}_l^{\theta} \geq 0, \qquad (C.3)
$$

which we are now going to prove. We shall show that the function  $f(x)$  is increasing in  $x \ge 0$ ; then the result follows by showing  $f(0) \ge 0$ , which is what was proven by Yuan and Bourgain.

We now verify that  $f'(x) \geq 0$ . By direct computation we see that

$$
f'(x) = -\theta \left( x + \sum_{l \ge 2} \widehat{n}_l \right)^{\theta - 1} + \theta,
$$

so it suffices to prove that

$$
1 \leq \left(x + \sum_{l \geq 2} \widehat{n}_l\right)^{1-\theta},\tag{C.4}
$$

which is indeed true, since  $\sum_{i \ge 2} \hat{n}_i \ge \hat{n}_2 \ge 1$  holds, by mass conservation.

*Proof of Lemma [7.1.](#page-32-1)* In this subsection we will take

$$
\boldsymbol{\alpha}, \boldsymbol{\beta} \in \mathbb{N}^{\mathbb{Z}} \quad \text{with} \quad 1 \leq |\boldsymbol{\alpha}| = |\boldsymbol{\beta}| < \infty. \tag{C.5}
$$

Given  $u \in \mathbb{Z}^{\mathbb{Z}}$ , with  $|u| < \infty$ , consider the set

$$
\{j \neq 0, \quad \text{repeated} \quad |u_j| \text{ times} \},
$$

where  $D < \infty$  is its cardinality. Define the vector  $m = m(u)$  as the reordering of the elements of the set above such that  $|m_1| \ge |m_2| \ge \cdots \ge |m_D| \ge 1$ . Given  $\alpha \ne \tilde{\beta} \in \mathbb{N}^{\mathbb{Z}}$ , with  $|\alpha| = |\beta| < \infty$  we consider  $m = m(\alpha - \beta)$  and  $\hat{n} = \hat{n}(\alpha + \beta)$ . If we denote by *D* the cardinality of *m* and *N* the one of  $\hat{n}$  we have

$$
D + \alpha_0 + \beta_0 \le N \tag{C.6}
$$

and

$$
(|m_1|, \ldots, |m_D|, \underbrace{1, \ldots, 1}_{N-D \text{ times}}) \leq (\widehat{n}_1, \ldots, \widehat{n}_N). \tag{C.7}
$$

Set

$$
\sigma_l = \text{sign}(\pmb{\alpha}_{m_l} - \pmb{\beta}_{m_l}).
$$

For every function  $g$  defined on  $\mathbb Z$  we have that

$$
\sum_{i\in\mathbb{Z}} g(i)|\boldsymbol{\alpha}_i - \boldsymbol{\beta}_i| = g(0)|\boldsymbol{\alpha}_0 - \boldsymbol{\beta}_0| + \sum_{l\geq 1} g(m_l),
$$
  

$$
\sum_{i\in\mathbb{Z}} g(i)(\boldsymbol{\alpha}_i - \boldsymbol{\beta}_i) = g(0)(\boldsymbol{\alpha}_0 - \boldsymbol{\beta}_0) + \sum_{l\geq 1} \sigma_l g(m_l).
$$
 (C.8)

**Lemma C.3.** *Assume that g defined on* Z *is non negative, even and not decreasing on*  $N$ *. Then, if* α  $\neq$  β*,* 

$$
\sum_{i \in \mathbb{Z}} g(i) |\boldsymbol{\alpha}_i - \boldsymbol{\beta}_i| \le 2g(m_1) + \sum_{l \ge 3} g(\widehat{n}_l). \tag{C.9}
$$

*Proof.* By [\(C.8\)](#page-49-1)

$$
\sum_{i \in \mathbb{Z}} g(i) |\boldsymbol{\alpha}_i - \boldsymbol{\beta}_i| = g(0) |\boldsymbol{\alpha}_0 - \boldsymbol{\beta}_0| + \sum_{l \ge 1} g(m_l)
$$
  

$$
\le g(1) (\boldsymbol{\alpha}_0 + \boldsymbol{\beta}_0) + 2g(m_1) + \sum_{l \ge 3} g(m_l)
$$

and  $(C.9)$  follows by  $(C.6)$  and  $(C.7)$ .

We denote as before the momentum by  $\pi$  so by [\(C.8\)](#page-49-1)

$$
\pi = \sum_{i \in \mathbb{Z}} (\alpha_i - \beta_i) i = \sum_l \sigma_l m_l \tag{C.10}
$$

and

$$
\sum_{i} (\alpha_i - \beta_i) i^2 = \sum_{l} \sigma_l m_l^2.
$$
 (C.11)

Analogously

$$
\sum_{i} |\boldsymbol{\alpha}_{i} - \boldsymbol{\beta}_{i}| = D + |\boldsymbol{\alpha}_{0} - \boldsymbol{\beta}_{0}| \stackrel{\text{(C.6)}}{\leq} N. \tag{C.12}
$$

Finally note that

$$
\sigma_l \sigma_{l'} = -1 \qquad \Longrightarrow \qquad m_l \neq m_{l'}.\tag{C.13}
$$

Note that

$$
\alpha \neq \beta \implies N \geq 3 \text{ or } \pi \neq 0,
$$
 (C.14)

indeed, by mass conservation,  $|\alpha| = |\beta| = 1$  therefore if  $N = 2$  we get  $\alpha - \beta = e_{j_1} - e_{j_2}$ so if  $\pi = 0$  we have  $\alpha = \beta$ . Note also that

$$
\alpha \neq \beta \qquad \Longrightarrow \qquad D \geq 1, \tag{C.15}
$$

indeed, if *D* = 0 then  $\alpha_l - \beta_l = 0$  for every  $|l| \ge 1$  and, by mass conservation  $\alpha_0 = \beta_0$ , contradicting  $\alpha \neq \beta$ .

**Lemma C.4.** *Given*  $\alpha \neq \beta \in \mathbb{N}^{\mathbb{Z}}$ , *with*  $1 < |\alpha| = |\beta| < \infty$  *and satisfying* [\(7.1\)](#page-32-6)*, we have*[32](#page-56-0)

<span id="page-56-1"></span>
$$
|m_1| \le 20|\pi| + 31 \sum_{l \ge 3} \widehat{n}_l^2. \tag{C.16}
$$

*Proof.* In the case  $D = 1$  by  $(C.10) |\pi| = |m_1|$  $(C.10) |\pi| = |m_1|$  and  $(C.16)$  follows. Let us now consider the case  $D = 2$ , i.e.

$$
\boldsymbol{\alpha}-\boldsymbol{\beta}=\sigma_1e_{m_1}+\sigma_2e_{m_2}+(\boldsymbol{\alpha}_0-\boldsymbol{\beta}_0)e_0.
$$

Let us start with the case  $\sigma_1 \sigma_2 = 1$ . By mass conservation  $|\sigma_1 + \sigma_2| = |\beta_0 - \alpha_0| = 2$ . By [\(C.12\)](#page-51-0)  $N \geq 4$ . Then conditions [\(7.1\)](#page-32-6) and (C.12) imply that

$$
m_1^2 + m_2^2 \le 20 + 10|\boldsymbol{\alpha}_0 - \boldsymbol{\beta}_0| = 40.
$$

Then

$$
|m_1| \le \sqrt{40} \le \frac{\sqrt{40}}{2} \sum_{\ell=3}^N \widehat{n}_{\ell}^2
$$

since *N*  $\geq$  4 and  $\hat{n}_{\ell} \geq 1$ . When  $\sigma_1 \sigma_2 = -1$  v<br>and by mass conservation  $\alpha_0 - \beta_0 = 0$ . Then since  $N \ge 4$  and  $\hat{n}_{\ell} \ge 1$ . When  $\sigma_1 \sigma_2 = -1$  we have  $m_1 \ne m_2$ ,  $|\pi| = |m_1 - m_2| \ge 1$ 

$$
(|m_1| + |m_2|)(|m_1| - |m_2|) = m_1^2 - m_2^2 \le 20.
$$

If  $|m_1| > |m_2|$  then

$$
|m_1| \le 20 \le 20|\pi|.
$$
 (C.17)

Otherwise  $m_1 = -m_2$  and, therefore,  $|\pi| = 2|m_1|$ , completing the proof in the case  $D = 2$ .

Let us now consider the case  $D \ge 3$ . By [\(7.1\)](#page-32-6), [\(C.11\)](#page-50-3) and [\(C.12\)](#page-51-0)

$$
m_1^2 + \sigma_1 \sigma_2 m_2^2 \le 10N + \sum_{l=3}^D m_l^2 \le 10N + \sum_{l=3}^N \widehat{n}_l^2
$$
  
= 20 +  $\sum_{l=3}^N (10 + \widehat{n}_l^2) \le 20 + 11 \sum_{l=3}^N \widehat{n}_l^2 \stackrel{N \ge 3}{\le} 31 \sum_{l=3}^N \widehat{n}_l^2.$ 

If  $\sigma_1 \sigma_2 = 1$  then

$$
|m_1|, |m_2| \leq \sqrt{31 \sum_{l \geq 3} \widehat{n}_l^2}.
$$

If  $\sigma_1 \sigma_2 = -1$ 

$$
(|m_1| + |m_2|)(|m_1| - |m_2|) = m_1^2 - m_2^2 \le 31 \sum_{l \ge 3} \widehat{n}_l^2.
$$

<span id="page-56-0"></span> $\overline{32}$  Note that by [\(C.14\)](#page-51-1) the r.h.s. of [\(C.16\)](#page-56-1) is at least 20.

Now, if  $|m_1| \neq |m_2|$  then

$$
|m_1| + |m_2| \le 31 \sum_{l \ge 3} \widehat{n}_l^2.
$$

Conversely, if  $|m_1| = |m_2|$ , by [\(C.13\)](#page-51-2),  $m_1 \neq m_2$ , hence  $m_1 = -m_2$ . By substituting this relation into  $(C.10)$ , we have

$$
2|m_1| \leq |\pi| + \sum_{l \geq 3} |m_l| \leq |\pi| + \sum_{l \geq 3} \widehat{n}_l^2,
$$

concluding the proof.

*Conclusion of the proof of Lemma* [7.1.](#page-32-1) As above, given  $\alpha, \beta \in \mathbb{N}^{\mathbb{Z}}$ , with  $1 < |\alpha| =$  $|\beta| < \infty$  we consider  $m = m(\alpha - \beta)$  and  $\hat{n} = \hat{n}(\alpha + \beta)$ . Note that  $N := |\alpha + \hat{\beta}| \ge 2$ . We have <sup>33</sup>

<span id="page-57-1"></span>
$$
\sum_{i} |\alpha_{i} - \beta_{i}| \langle i \rangle^{\theta/2} \stackrel{(C.9)}{\leq} 2|m_{1}|^{\frac{\theta}{2}} + \sum_{l \geq 3} \hat{n}_{l}^{\frac{\theta}{2}}
$$
\n
$$
\stackrel{(C.16)}{\leq} 2 \left(20|\pi| + 31 \sum_{l \geq 3} \hat{n}_{l}^{2}\right)^{\frac{\theta}{2}} + \sum_{l \geq 3} \hat{n}_{l}^{\frac{\theta}{2}}
$$
\n
$$
\leq 2(20|\pi|)^{\frac{\theta}{2}} + 2(31)^{\frac{\theta}{2}} \sum_{l \geq 3} \hat{n}_{l}^{\theta} + \sum_{l \geq 3} \hat{n}_{l}^{\frac{\theta}{2}}
$$
\n
$$
\leq \frac{13}{1-\theta} \left(1-\theta\left||\pi\right| + (2-2^{\theta}) \left(\sum_{l \geq 3} \hat{n}_{l}^{\theta}\right)\right), \tag{C.18}
$$

using that  $1 - \theta \le 2 - 2^{\theta}$  for  $0 \le \theta \le 1$ . Then by Lemma [6.1](#page-28-0) and [\(C.18\)](#page-57-1) we get

$$
\sum_{i} |\boldsymbol{\alpha}_{i} - \boldsymbol{\beta}_{i}| \langle i \rangle^{\theta/2} \leq \frac{13}{1-\theta} \Biggl( (1-\theta)|\pi| + \sum_{i} \langle i \rangle^{\theta} (\boldsymbol{\alpha}_{i} + \boldsymbol{\beta}_{i}) + \theta |\pi| - 2\widehat{n}_{1}^{\theta} \Biggr) \leq \frac{13}{1-\theta} \Biggl[ \sum_{i} \langle i \rangle^{\theta} (\boldsymbol{\alpha}_{i} + \boldsymbol{\beta}_{i}) + |\pi| - 2 \langle j \rangle^{\theta} \Biggr],
$$

proving  $(7.2)$ .

Let us now prove [\(7.3\)](#page-33-2) passing to the logarithm. We have

<span id="page-57-2"></span>
$$
\sum_{i} \ln(1 + |\boldsymbol{\alpha}_{i} - \boldsymbol{\beta}_{i}| \langle i \rangle)
$$
\n
$$
= \sum_{|i| \le 1} \ln(1 + |\boldsymbol{\alpha}_{i} - \boldsymbol{\beta}_{i}|) + \sum_{|i| \ge 2} \ln(1 + |\boldsymbol{\alpha}_{i} - \boldsymbol{\beta}_{i}| |i|)
$$
\n
$$
\le 3 \ln(1 + N) + \sum_{|i| \ge 2} \ln(1 + |\boldsymbol{\alpha}_{i} - \boldsymbol{\beta}_{i}| |i|)
$$
\n
$$
\le 3 \ln 2 + 3 \ln N + \frac{3}{2} \sum_{|i| \ge 2} |\boldsymbol{\alpha}_{i} - \boldsymbol{\beta}_{i}| \ln |i|,
$$
\n
$$
(C.19)
$$

<span id="page-57-0"></span><sup>33</sup> Using that for *x*,  $y \ge 0$  and  $0 \le c \le 1$  we get  $(x + y)^c \le x^c + y^c$ .

using that  $1 + cx \le \frac{3}{2}x^c$  for  $c \ge 1$ ,  $x \ge 2$ . If  $\alpha_i - \beta_i = 0$  for every  $|i| \ge 2$  then [\(7.3\)](#page-33-2) follows. Assume now that  $\alpha_i - \beta_i \neq 0$  for some  $|i| \geq 2$ . By [\(C.14\)](#page-51-1) we have

<span id="page-58-2"></span>
$$
N \ge 3 \quad \text{or} \quad \pi \ne 0. \tag{C.20}
$$

We claim that, when  $N \geq 3$ ,

<span id="page-58-0"></span>
$$
\ln\left(\sum_{l=3}^{N}\widehat{n}_{l}^{2}\right) \leq \ln N + \sum_{l=3}^{N}\ln\widehat{n}_{l}^{2}.
$$
 (C.21)

Let Let  $S := \{3 \le l \le N, \text{ s.t. } \hat{n}_l \ge 2\}$ . If  $S = \emptyset$  we have the equality in [\(C.21\)](#page-58-0). Otherwise  $\sum_{l \in S} \hat{n}_l^2 \ge 4$  and  $\hat{n}_l^3$  $l \in S$   $\hat{n}_l^2 \geq 4$  and  $\hat{d}^3$ 

$$
\ln\left(\sum_{l=3}^N \widehat{n}_l^2\right) \le \ln\left(N + \sum_{l \in \mathcal{S}} \widehat{n}_l^2\right) \le \ln N + \sum_{l \in \mathcal{S}} \ln \widehat{n}_l^2,
$$

proving  $(C.21)$ . We claim that

<span id="page-58-3"></span>
$$
\ln\left(20|\pi| + 31\sum_{l=3}^{N}\widehat{n}_{l}^{2}\right) \le \ln(1+|\pi|) + \ln N + \sum_{l=3}^{N}\ln\widehat{n}_{l}^{2} + \ln 20 + \ln 31. \quad (C.22)
$$

Indeed consider first the case  $\pi = 0$ , then  $N \ge 3$  by [\(C.20\)](#page-58-2) and [\(C.22\)](#page-58-3) follows by [\(C.21\)](#page-58-0). Consider now the case  $|\pi| \ge 1$ . If  $N < 3$  [\(C.22\)](#page-58-3) follows (there is no sum). If  $N > 3$  we have  $35$ 

$$
\ln\left(20|\pi| + 31\sum_{l=3}^{N}\widehat{n}_{l}^{2}\right) \le \ln\left(20|\pi|\right) + \ln\left(31\sum_{l=3}^{N}\widehat{n}_{l}^{2}\right)
$$

$$
\le \ln(|\pi|) + \ln\left(\sum_{l=3}^{N}\widehat{n}_{l}^{2}\right) + \ln 20 + \ln 31.
$$

Recalling  $(C.21)$  this complete the proof of  $(C.22)$ .

Let us continue the proof of [\(7.3\)](#page-33-2). Set  $g(i) := 0$  if  $|i| \le 1$  and  $g(i) := \ln |i|$  if  $|i| \ge 2$ and apply  $(C.9)$  to  $(C.19)$ ; we get

$$
\sum_{|i|\geq 2} |\alpha_i - \beta_i| \ln |i| \leq 2 \ln |m_1| + \sum_{l\geq 3} \ln |\widehat{n}_l|
$$
  
\n(C.16)  
\n
$$
\leq 2 \ln \left( 20|\pi| + 31 \sum_{l\geq 3} \widehat{n}_l^2 \right) + \sum_{l\geq 3} \ln \widehat{n}_l
$$
  
\n(C.22)  
\n
$$
\leq 2 \ln(1 + |\pi|) + 2 \ln N + 5 \sum_{l=3}^N \ln \widehat{n}_l + 16.
$$

<span id="page-58-4"></span><span id="page-58-1"></span>

<sup>&</sup>lt;sup>34</sup> Use that  $ln(x + y) \le ln x + ln y$  if  $x, y \ge 2$ .<br><sup>35</sup> Recall footnote [34.](#page-58-1)

Inserting in  $(C.19)$  we obtain

$$
\sum_{i} \ln(1 + |\alpha_i - \beta_i| \langle i \rangle) \le 3 \ln(1 + |\pi|) + 6 \ln N + \frac{15}{2} \sum_{l=3}^{N} \ln \widehat{n}_l + 27
$$

concluding the proof of [\(7.3\)](#page-33-2).

*Proof of Lemma [7.2.](#page-34-2)* First of all we note that

$$
\sum_{i} f_i(|\ell_i|) = \sum_{i \text{ s.t. } \ell_i \neq 0} f_i(|\ell_i|)
$$

since  $f_i(0) = 0$ . We have that<sup>[36](#page-59-0)</sup>

$$
f_i(x) \leq -\frac{\sigma}{C_*} \langle i \rangle^{\frac{\theta}{2}} x + 2 \ln(x) + (2+q) \ln \langle i \rangle + 1, \quad \forall x \geq 1.
$$

We have that

$$
\max_{x \ge 1} \left( -\frac{\sigma}{C_*} \langle i \rangle^{\frac{\theta}{2}} x + 2 \ln(x) \right) = \begin{cases} -\frac{\sigma}{C_*} \langle i \rangle^{\frac{\theta}{2}} & \text{if } \langle i \rangle \ge i_0, \\ -2 + 2 \ln \frac{2C_*}{\sigma} - \theta \ln \langle i \rangle & \text{if } \langle i \rangle < i_0, \end{cases}
$$

where

$$
i_0 := \left(\frac{2C_*}{\sigma}\right)^{\frac{2}{\theta}},
$$

since the maximum is achieved for  $x = 1$  if  $\langle i \rangle \ge i_0$  and  $x = \frac{2C_*}{\sigma \langle i \rangle^{\theta/2}}$  if  $\langle i \rangle < i_0$ . Note that  $i_0 \geq e$ . Then we get

$$
\sum_{i} f_i(|\ell_i|) = \sum_{i \text{ s.t. } \ell_i \neq 0} f_i(|\ell_i|) \le
$$
\n
$$
\sum_{\langle i \rangle \ge i_0 \text{ s.t. } \ell_i \neq 0} \left( (2+q) \ln \langle i \rangle + 1 - \frac{\sigma}{C_*} \langle i \rangle^{\frac{\theta}{2}} \right) + \sum_{\langle i \rangle < i_0} \left( 2 \ln \frac{2C_*}{\sigma} + \left( 2 + q - \theta \right) \ln \langle i \rangle \right).
$$

We immediately have that

$$
\sum_{\langle i \rangle < i_0} \left( 2 \ln \frac{2C_*}{\sigma} + \left( 2 + q - \theta \right) \ln \langle i \rangle \right) \le 3i_0 \left( 2 \ln \frac{2C_*}{\sigma} + (2 + q) \ln i_0 \right)
$$
\n
$$
= 3 \left( 2 + \frac{2}{\theta} (2 + q) \right) \left( \frac{2C_*}{\sigma} \right)^{\frac{2}{\theta}} \ln \frac{2C_*}{\sigma}.
$$

Moreover, in the case  $\langle i \rangle \ge i_0 \ge e$ ,

$$
(2+q)\ln\langle i\rangle + 1 - \frac{\sigma}{C_*}\langle i\rangle^{\frac{\theta}{2}} \le (2+q+1)\ln\langle i\rangle - \frac{\sigma}{C_*}\langle i\rangle^{\frac{\theta}{2}}
$$

$$
= \frac{2}{\theta}(2+q+1)\Big(\ln\langle i\rangle^{\frac{\theta}{2}} - 2\mathfrak{C}\langle i\rangle^{\frac{\theta}{2}}\Big)
$$

<span id="page-59-0"></span> $\frac{36}{36}$  Using that  $\ln(1 + y) \le 1 + \ln y$  for every  $y \ge 1$ .

$$
\Box
$$

where

$$
\mathfrak{C} := \frac{\sigma \theta}{4C_*(2+q+1)} < 1.
$$

Therefore

$$
S_* := \sum_{\langle i \rangle \ge i_0 \text{ s.t. } \ell_i \neq 0} \left( (2+q) \ln \langle i \rangle + 1 - \frac{\sigma}{C_*} \langle i \rangle^{\frac{\theta}{2}} \right)
$$

satisfies

$$
S_{*} \leq \sum_{\langle i \rangle \geq i_0 \text{ s.t. } \ell_i \neq 0} \frac{2}{\theta} (2+q+1) \Big( \ln \langle i \rangle^{\frac{\theta}{2}} - 2 \mathfrak{C} \langle i \rangle^{\frac{\theta}{2}} \Big).
$$

We have that  $37$ 

$$
\ln \langle i \rangle^{\frac{\theta}{2}} - 2 \mathfrak{C} \langle i \rangle^{\frac{\theta}{2}} \leq - \mathfrak{C} \langle i \rangle^{\frac{\theta}{2}}, \quad \text{when} \quad \langle i \rangle \geq i_* := \left( \frac{2}{\mathfrak{C}} \ln \frac{1}{\mathfrak{C}} \right)^{\frac{2}{\theta}}.
$$

Note that

$$
i_{\sharp} \geq \max\{i_0, i_*\}.
$$

Therefore

$$
S_{*} \leq \frac{2}{\theta}(2+q+1)\left(\sum_{\langle i\rangle < i_{\sharp}} \ln \langle i\rangle^{\frac{\theta}{2}} - \sum_{\langle i\rangle \geq i_{\sharp} \text{ s.t. } \ell_{i} \neq 0} \left(\mathfrak{C}\langle i\rangle^{\frac{\theta}{2}}\right)\right)
$$

$$
\leq (2+q+1)\left(3i_{\sharp}\ln i_{\sharp} - \frac{2\mathfrak{C}}{\theta}M_{\ell}^{\frac{\theta}{2}}\right)
$$

where

$$
M_{\ell} := \max\{|i| \ge i_{\sharp}, \text{ s.t. } \ell_i \neq 0\}
$$

and  $M_{\ell} := 0$  if  $|\ell_i| = 0$  for every  $|i| \ge i_{\sharp}$ . In conclusion we get

$$
\sum_{i} f_{i}(|\ell_{i}|) \leq 3\left(2 + \frac{2}{\theta}(2 + q)\right) \left(\frac{2C_{*}}{\sigma}\right)^{\frac{2}{\theta}} \ln \frac{2C_{*}}{\sigma} + (2 + q + 1)\left(3i_{\sharp}\ln i_{\sharp} - \frac{2\mathfrak{C}}{\theta}M_{\ell}^{\frac{\theta}{2}}\right)
$$
  

$$
\leq 6(q + 3)i_{\sharp}\ln i_{\sharp} - \frac{\sigma}{2C_{*}}M_{\ell}^{\frac{\theta}{2}}
$$
  

$$
\leq 7(q + 3)i_{\sharp}\ln i_{\sharp} - \frac{\sigma}{2C_{*}}(\widehat{n}_{1}(\ell))^{\frac{\theta}{2}},
$$

noting that  $\widehat{n}_1(\ell) = M_\ell$  if  $M_\ell \neq 0$ , otherwise  $\widehat{n}_1(\ell) < i_\sharp$ , and, therefore,

$$
\frac{\sigma}{2C_*} \left(\widehat{n}_1(\ell)\right)^{\frac{\theta}{2}} < \frac{\sigma}{2C_*} i_{\sharp}^{\frac{\theta}{2}} \leq (q+3)i_{\sharp} \ln i_{\sharp}.
$$



<span id="page-60-0"></span><sup>&</sup>lt;sup>37</sup> Using that, for every fixed  $0 < \mathfrak{C} \le 1$ , we have  $\mathfrak{C}x \ge \ln x$  for every  $x \ge \frac{2}{\mathfrak{C}} \ln \frac{1}{\mathfrak{C}}$ .

*Proof of Lemma [4.1.](#page-16-0)* For  $\ell \in \mathbb{Z}^{\mathbb{Z}}$  with  $0 < |\ell| < \infty$  we define

$$
\mathcal{R}_\ell:=\left\{\omega\in\Omega_q\,:\,|\omega\cdot\ell|\leq\frac{\gamma}{1+|\ell_0|^{\mu_1}}\prod_{n\neq 0}\frac{1}{(1+|\ell_n|^{\mu_1}|n|^{\mu_2+q})}\right\}
$$

 $-$  if  $\ell$  is such that  $\ell_n = 0 \,\forall n \neq 0$  then

$$
\mu(\mathcal{R}_{\ell}) = \frac{\gamma}{1 + |\ell_0|^{\mu_1}}.
$$

− Otherwise: let  $s = s(\ell) > 0$  be the smallest positive index *i* such that  $|\ell_i| + |\ell_{-i}| \neq 0$ and  $S = S(\ell)$  be the biggest. Then we have<sup>38</sup>

<span id="page-61-3"></span>
$$
\mu(\mathcal{R}_{\ell}) \leq \frac{\gamma s^{q}}{(1+|\ell_0|^{\mu_1})} \prod_{n \neq 0} \frac{1}{(1+|\ell_n|^{\mu_1}|n|^{\mu_2+q})}.
$$

Let us write

$$
\frac{1}{1+|\ell_0|^{\mu_1}} \prod_{n\neq 0} \frac{1}{(1+|\ell_n|^{\mu_1}|n|^{\mu_2+q})}
$$
\n
$$
= \frac{1}{1+|\ell_0|^{\mu_1}} \prod_{n>0} \frac{1}{(1+|\ell_n|^{\mu_1}|n|^{\mu_2+q})} \frac{1}{(1+|\ell_{-n}|^{\mu_1}|n|^{\mu_2+q})}
$$
\n
$$
= \frac{1}{1+|\ell_0|^{\mu_1}} \prod_{s(\ell)\leq n\leq S(\ell)} \frac{1}{(1+|\ell_n|^{\mu_1}|n|^{\mu_2+q})} \frac{1}{(1+|\ell_{-n}|^{\mu_1}|n|^{\mu_2+q})}.
$$

Now

$$
\mu(\Omega_q \setminus D_{\gamma,q}) \le \sum_{\ell} \mu(\mathcal{R}_{\ell}) = \sum_{\ell_0} \frac{\gamma}{1 + |\ell_0|^{\mu_1}}
$$
\n(C.23)

$$
+\sum_{s>0} \sum_{\ell:s(\ell)=S(\ell)=s} \frac{1}{1+|\ell_0|^{\mu_1}} \frac{\gamma s^q}{|\ell_s|(1+|\ell_s|^{\mu_1}|s|^{\mu_2+q})} \frac{1}{(1+|\ell_{-s}|^{\mu_1}|s|^{\mu_2+q})} \qquad (C.24)
$$

$$
+\sum_{0
$$

Let us estimate [\(C.24\)](#page-61-2)

<span id="page-61-2"></span><span id="page-61-0"></span>
$$
\sum_{s>0} \sum_{\ell_0 \in \mathbb{Z}} \frac{1}{1+|\ell_0|^{\mu_1}} \sum_{\substack{\ell_s, \ell_{-s} \in \mathbb{Z} \\ |\ell_s| + |\ell_{-s}| > 0}} \frac{\gamma s^q}{(1+|\ell_s|^{\mu_1}|s|^{\mu_2+q})} \frac{1}{(1+|\ell_{-s}|^{\mu_1}|s|^{\mu_2+q})}
$$
  

$$
\leq c(\mu_1) \gamma \sum_{s>0} s^q \sum_{\substack{\ell_s, \ell_{-s} \in \mathbb{Z} \\ |\ell_s| + |\ell_{-s}| > 0}} \frac{1}{(1+|\ell_s|^{\mu_1}|s|^{\mu_2+q})} \frac{1}{(1+|\ell_{-s}|^{\mu_1}|s|^{\mu_2+q})}.
$$

<span id="page-61-1"></span>38 Assume, e.g. that  $\ell_s \neq 0$ , then  $|\partial_{\xi_s} \omega \cdot \ell| \geq s^{-q}$ .

Now since

$$
\sum_{h=1}^{\infty} \frac{1}{(1+h^{\mu_1}|n|^{\mu_2+q})} \le \sum_{h=1}^{\infty} \frac{1}{h^{\mu_1}|n|^{\mu_2+q}} \le \frac{c(\mu_1)}{|n|^{\mu_2+q}}
$$

we have

$$
\sum_{h\in\mathbb{Z}}\frac{1}{(1+|h|^{\mu_1}|n|^{\mu_2+q})}\leq 1+\frac{2c(\mu_1)}{|n|^{\mu_2+q}}.
$$

Then we have

$$
\sum_{\substack{\ell_s,\ell_{-s}\in\mathbb{Z}\\|\ell_s|+|\ell_{-s}|>0}}\frac{1}{(1+|\ell_s|^{\mu_1}|s|^{\mu_2+q})}\frac{1}{(1+|\ell_{-s}|^{\mu_1}|s|^{\mu_2+q})}\le\frac{c_1(\mu_1)}{|s|^{\mu_2+q}}
$$

and consequently [\(C.24\)](#page-61-2) is bounded by

$$
c_2(\mu_1)\gamma\sum_{s>0}|s|^b\leq c_3(\mu_1,\mu_2)\gamma.
$$

Regarding the third line in [\(C.23\)](#page-61-3), we note that for all *n* we have

$$
\sum_{\ell_n,\ell_{-n}\in\mathbb{Z}}\frac{1}{(1+|\ell_n|^{\mu_1}|n|^{\mu_2+q})}\frac{1}{(1+|\ell_{-n}|^{\mu_1}|n|^{\mu_2+q})}\leq \left(1+2\frac{c(\mu_1)}{|n|^{\mu_2+q}}\right)^2.
$$

Hence

$$
\sum_{\ell:s(\ell)=s, \atop S(\ell)=S} \frac{1}{1+|\ell_0|^{\mu_1}} \prod_{s\leq n\leq S} \frac{1}{(1+|\ell_n|^{\mu_1}|n|^{\mu_2+q})} \frac{1}{(1+|\ell_{-n}|^{\mu_1}|n|^{\mu_2+q})}
$$
\n
$$
= \sum_{\ell_0\in\mathbb{Z}} \frac{1}{1+|\ell_0|^{\mu_1}} \times \sum_{\substack{\ell_s,\ell_{-s}\in\mathbb{Z} \\ |\ell_s|+|\ell_{-s}|>0}} \frac{1}{(1+|\ell_s|^{\mu_1}|s|^{\mu_2+q})} \frac{1}{(1+|\ell_{-s}|^{\mu_1}|s|^{\mu_2+q})}
$$
\n
$$
\times \sum_{\substack{\ell_s,\ell_{-s}\in\mathbb{Z} \\ |\ell_s|+|\ell_{-s}|>0}} \frac{1}{(1+|\ell_s|^{\mu_1}|S|^{\mu_2+q})} \frac{1}{(1+|\ell_{-s}|^{\mu_1}|S|^{\mu_2+q})}
$$
\n
$$
\times \prod_{s\n
$$
\leq \frac{c_4(\mu_1)}{s^{\mu_2+q}S^{\mu_2+q}} \prod_{s\n
$$
\leq \frac{c_4(\mu_1)}{s^{\mu_2+q}S^{\mu_2+q}} \exp\left(\sum_{n\geq 1} \ln\left(1+2\frac{c(\mu_1)}{|n|^{\mu_2+q}}\right)^2\right)
$$
\n
$$
\leq \frac{c_5(\mu_1, \mu_2)}{s^{\mu_2+q}S^{\mu_2+q}}.
$$
$$
$$

Then, multiplying by  $\gamma s^q$  and taking the  $\sum_{0 < s < S}$ , we have that also [\(C.25\)](#page-61-0) is bounded by some constant  $C_{\text{meas}}(\mu_1, \mu_2)\gamma$ .

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