# REGULAR ARTICLE

# CFD modeling of the high-temperature HVPE growth of aluminum nitride layers on c-plane sapphire: from theoretical chemistry to process evaluation

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Abstract This study presents numerical modeling based on a relatively limited number of gas-phase and surface reactions to simulate the growth rate of aluminum nitride layers on AlN templates and c-plane sapphire in a broad range of deposition parameters. Modeling results have been used to design particular experiments in order to understand the influence of the process parameters on the crystal quality of AlN layers grown in a high-temperature hydride vapor-phase epitaxy process fed with  $NH_3$ , AlCl<sub>3</sub>, and H<sub>2</sub>. Modeling results allow to access to very interesting local quantities such as the surface site ratio and local supersaturation. The developed universal model starting from local parameters might be easily transferred to other reactor geometry and process conditions. Among the investigated parameters (growth rate, temperature, local supersaturation, gas-phase N/Al ratio, and local surface site N/Al ratio), only the growth rate/supersaturation or growth rate/

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temperature relationships exhibit a clear process window to use in order to succeed in growing epitaxial AlN layers on c-plane sapphire or AlN templates. Gas-phase N/Al ratio and local surface site N/Al ratio seem to play only a secondary role in AlN epitaxial growth.

Keywords AlN · High-temperature HVPE · Numerical modeling - Epitaxial growth

## List of symbols

- (S) Surface site
- (B) Bulk or solid species
- [A] Gas-phase concentration of species A (mol  $m^{-3}$ )
- $a_k$  Pre-exponential factor for reaction k (consistent units)
- $[C(S)]$  Surface site concentration of species C (mol m<sup>-2</sup>)
- $E_{ak}$  Activation energy for reaction k (J mol<sup>-1</sup>)
- $k_k$  Rate constant for reaction  $k (s^{-1}$  or mol<sup>-1</sup> m<sup>3</sup> s or  $mol^{-2}$  m<sup>6</sup> s)
- $M$  Unspecified species  $(-)$
- $M_A$  Molar mass for gaseous specie A (kg mol<sup>-1</sup>)
- $P_A$  Partial pressure of A at the growing AlN surface (Pa)
- $P_A^*$ Equilibrium pressure of  $A$  versus AlN in vacuum (Pa)
- R Ideal gas constant  $(J \text{ mol}^{-1} \text{ K}^{-1})$
- $T$  Temperature  $(K)$
- $\beta_k$  Temperature exponent for reaction  $k$  (–)
- $\gamma_A$  Sticking coefficient for gaseous species A (–)
- $\gamma_0$  Pre-exponential factor for the sticking coefficient  $(-)$
- $\Gamma_{\text{tot}}$  Total surface site concentration (mol m<sup>-2</sup>)
- $v_k$  Production rate of solid surface for reaction  $k \pmod{m^{-2} s^{-1}}$

## 1 Introduction

Aluminum nitride is a promising substrate for AlGaNbased UV LED (intended to be used for water purification and optical information storage) and piezoelectric applications [micro-electro-mechanical systems (MEMS) and surface acoustic waves (SAW) devices]. The UV LED industry requires high-quality single crystals (i.e., deep UV transparency, low density of defect) [\[1](#page-11-0)]. The requirements are more accessible for piezoelectric applications, for which crack-free highly oriented c-axis layers are needed [\[2](#page-11-0), [3\]](#page-11-0). Heteroepitaxy is necessary for the development of these new devices and systems, due to the fact that it is currently the cheapest solution to produce industrial grade AlN single crystals. Among the different available processes for AlN growth (physical vapor transport, PVT [\[4](#page-11-0)], metal–organic chemical vapor deposition, MOCVD [\[5](#page-11-0)], and high-temperature hydride vapor-phase epitaxy (HT– HVPE)  $[6-16]$ , high-temperature HVPE  $(>1,200 \degree C)$ becomes the most prospective technique to produce the required quality for both piezoelectric and semiconductor industry. Despite its considerable interest, this technique is still little studied from a modeling point of view. A better understanding of the phenomena leading to high-quality AlN layers grown on sapphire is the key point to allow high-temperature HVPE becoming a new industrial reference in thick AlN layers processing.

High-temperature CVD processes are commonly used to produce thick layers of AlN on various substrates. The halide process using Al chloride and  $NH<sub>3</sub>$  or  $N<sub>2</sub>$  as Al and N-precursors, respectively, is the most commonly used. The experimental and thermodynamic study of Nickel et al. [\[6](#page-11-0)] very early assessed the possibility to fabricate AlN layers from  $AICI_3$  and  $NH_3$  precursors. Attempts to obtain epitaxial growth on sapphire are more recent one. Kumagai et al. [\[7](#page-11-0), [8](#page-11-0)], Freitas et al. [\[9](#page-11-0)], Nagashima et al. [\[10](#page-11-0)], Eriguchi et al. [[11,](#page-11-0) [12](#page-11-0)], and Tajima et al. [\[13](#page-11-0)] succeeded in growing AlN on (0001) sapphire substrates in a hot-wall reactor, in the  $900-1,550$  °C temperature range. Claudel et al. [\[14–16](#page-11-0)] obtained similar results in a cold-wall reactor.

Kinetic mechanisms of AlN deposition from chlorinated precursors were first initiated by Allendorf's team [\[17](#page-11-0)[–20](#page-12-0)]. They proposed a set of kinetic and thermodynamic data for the B–Cl–N–H system, complemented by the reaction set of  $NH_3$  pyrolysis [\[21](#page-12-0)]. Dollet et al. [\[22](#page-12-0)] extrapolated the gas-phase kinetic data for AlN formation from the work of these previous authors and proposed sticking coefficients for  $AICI_3$  and  $NH_3$  on a growing AlN surface. Their study reported a mechanism where  $AICl<sub>3</sub>$  and  $AICl<sub>2</sub>NH<sub>2</sub>$  are the main aluminum containing species precursors. However, only irreversible gas-phase reactions for the formation of  $AlCl<sub>2</sub>NH<sub>2</sub>$  were used due to a lack of available

thermodynamic data for this species. At the same time, Swihart et al. [[23\]](#page-12-0) proposed a comprehensive kinetic data set for the Al–Cl–H system from ab initio calculations.

Recently, Cai et al. [\[24,](#page-12-0) [25\]](#page-12-0) proposed from the data of [\[22](#page-12-0)] an irreversible gas-phase growth mechanism involving AlCl<sub>3</sub>, AlCl<sub>2</sub>NH<sub>2</sub>, AlCl(NH<sub>2</sub>)<sub>2</sub>, Al(NH<sub>2</sub>)<sub>3</sub>, and AlClNH, as aluminum containing species, neglecting the surface kinetic of  $NH_3$  absorption ( $NH_3$  in large excess), and using the pre-exponential factors of Arrhenius laws as fitting parameters. Segal et al. [\[26](#page-12-0)] similarly conducted a mass transport and quasi-thermodynamic surface kinetic model for AlN growth.

We have also previously developed [\[27](#page-12-0)] and proposed a complete description of heat transfer, mass transport, gas flow dynamics and chemical reactions based on available data  $[21-23]$  and experimental data of AlN growth on AlN templates. This modeling work [[27\]](#page-12-0) included a complete thermodynamic modeling of the species involved in the database reported in Ref. [\[22](#page-12-0)] and proposed forward and backward reactions rates in the Al– Cl–N–H system. The fitting between experimental and numerical results was made via the activation energy of the surface reactions.

The aim of the present study is (1) to extend previous modeling results with the help of new experiments covering a broader range of experimental parameters (temperature from 900 to 1,500  $\degree$ C, N/Al ratio in gas phase from 0.75 to 30, total pressure fixed at 1,333 Pa, partial pressure of precursors of Al- and N-containing reactive species from 0.5 to 120 Pa) and (2) to increase the understanding of the effect of both N/Al ratio and supersaturation in the gas phase on the crystalline quality of AlN layers grown on sapphire.

## 2 Numerical modeling

# 2.1 Model formalism: gas-phase kinetic and equilibrium analysis

All the calculations presented in this paper were performed using the CFD–ACE multiphysics software package (2012 v1[—http://www.esi-group.com\)](http://www.esi-group.com), using the geometry of the vertical and axisymmetric reactor described in Ref. [[27\]](#page-12-0).

In the high-temperature HVPE process used in this study,  $NH_3$  was used as nitrogen source, Al chloride (mainly  $AICI_3$ ) fabricated from Al in situ chlorination as aluminum source, and  $H_2$  as carrier/reactive gas. In consequence, an exhaustive study of the Al–Cl–N–H system has been conducted in order to highlight particular kinetic pathways.

Thermodynamic and kinetic analysis presented in Ref. [\[27](#page-12-0)] leads to the following simplified set of elementary gas-

#### <span id="page-2-0"></span>Table 1 Gas-phase reaction kinetic parameters



Units are given in K s mol m.  $M$  is an unspecified third-body species

phase reactions as relevant to model the AlN growth rate. The rate constant for each reaction is expressed in the form of a classical Arrhenius law (for reaction  $k$ ):

$$
k_k = a_k T^{\beta_k} \exp\left(\frac{-E_{ak}}{RT}\right).
$$
 (1)

Table 1 gives the kinetic parameters for each elementary gas-phase reaction (the order with respect to each reactant is 1).

Reactions in H–Cl system of Table 1 are optional and are only needed to increase the precision of the mass transport modeling (see Table [2](#page-3-0) also). For unknown backward rates, the equilibrium constant is calculated according to Ref. [[28\]](#page-12-0).

#### 2.2 Mass transport modeling and surface reactions

The species transport properties are evaluated by various methods. The diffusion coefficients, mixtures viscosities, and mixtures thermal conductivities are calculated using the Enskog theory [[29\]](#page-12-0). The Lennard–Jones parameters used for each gaseous species are either taken in the data reported in Refs. [[29](#page-12-0), [30](#page-12-0)] or computed according to the empirical method presented in Ref. [[30\]](#page-12-0). The transport equations for gas species do not take into account the multi-component Stephan Maxwell formalism, due to the low concentration of reactive species in  $H_2$  [[29](#page-12-0), [31](#page-12-0)]. Since the Reynolds number in our CVD reactor never exceeds 50, the laminar incompressible Navier–Stokes equation is solved to obtain the convective part of the diffusion-convection problem. The local density of fluid follows an ideal gas law.

To compute the mixture enthalpy, the JANAF formalism [\[32](#page-12-0)] is used. Each thermal property (i.e., entropy, enthalpy, and heat capacity) is calculated by a polynomial relationship with temperature. The used polynomial coefficients are taken from the Sandia National Laboratory Database [\[33\]](#page-12-0).

For surface reactions, the atomic site formalism proposed in Ref. [\[34](#page-12-0)] is adopted. It is considered that each site is an atom chemically bonded to the surface. Two kinds of atomic sites (atoms) are available on growing AlN surfaces: Al(S) and  $N(S)$  sites. The total number of sites is initially evaluated as the atom surface concentration on the polar c-plane of the AlN wurtzite polytype (1.98  $\times$  10<sup>-5</sup> mol m<sup>-2</sup>). The surface reaction rate for gas-phase molecules A approaching complementary atomic sites  $C(S)$  is expressed as [[34\]](#page-12-0):

$$
\nu_k = k_k[A] \times [C(S)] \quad \text{with} \quad k_k = \sqrt{\frac{RT}{2\pi M_A}} \frac{\gamma_A}{[\Gamma_{\text{tot}}]}.
$$
 (2)

The ratio  $C(S)/[\Gamma_{tot}]$  represents the fraction of surface sites C available for the sticking of gas-phase A molecules. The surface production rate is simply expressed as the product of an atom flux (kinetic theory of gases) with a probability that an atom of this flux sticks on the solid surface. We choose to express a simple dependency of the sticking coefficient  $\gamma_A$  with temperature:

$$
\gamma_A = \gamma_0 \exp\left(\frac{-E_{ak}}{RT}\right). \tag{3}
$$

In Eq. (3),  $\gamma_0 = 1$  is taken as simplifying assumption. The activation temperatures for each surface reaction were varied in order to minimize the relative difference between computed and experimental AlN growth rates (results presented in the next section). In order to fit the rate constant  $k_k$  for surface reaction to the Arrhenius law of Eq. (1), the following changes of variables is used as follows:

$$
a_k = \sqrt{\frac{R}{2\pi M_A}} \frac{1}{[\Gamma_{\text{tot}}]}
$$
  
\n
$$
\beta_k = 0.5
$$
  
\n
$$
\frac{E_{ak}}{R} \Big|_{\text{Surface Arehenius law}} = \frac{E_{ak}}{R} \Big|_{\text{Sticking coefficient}}
$$
\n(4)

Surface reactions Reaction	$a_k$ Pre exp. factor	$\beta_k$ Temp. exp.	$E_{ak}/R$ Activ. temp $(K)$
$NH_3 + Al(S) \leftrightarrow N(S) + H_2 + H^{\circ}$	$4.453E + 05$	0.5	8,654
$AICl_3 + N(S) \leftrightarrow AIN(B) + AI(S) + Cl^{\circ} + Cl_2$	$1.590E + 05$	0.5	7.700
$AICl_2NH_2 + Al(S) \leftrightarrow AlN(B) + Al(S) + 2HCl$	$1.720E + 05$	0.5	5,100
$AICl_2NH_2 + N(S) \leftrightarrow ALN(B) + N(S) + 2HCl$	1.720E+05	0.5	5.100

<span id="page-3-0"></span>Table 2 Optimized surface kinetic parameters for AlN growth

Units are given in K s mol m. Data corrected from Ref. [[27](#page-12-0)]

This form is mandatory to enter the surface kinetics into CFD–ACE multiphysics software database. Table 2 gives the kinetic parameters for each elementary surface reaction, according to this change of variables (the order with respect to each reactant is 1). The surface reactions are considered irreversible (no etching taken into account).

It should be noticed that the experimental activation temperature for AlN growth is about 7,460 K  $(62 \text{ kJ mole}^{-1})$  and that the transition between kineticlimited growth and diffusion-limited growth is observed from  $1,200$  to  $1,300$  °C. The modeling presented here covers both diffusion and kinetic regime.

Solving numerically the balance of chemical site on the AlN surface gives access to the surface  $N(S)$ –Al(S) ratio during growth. This balance is simply obtained by assuming that the molar rate of absorption of nitrogen and aluminum atoms is the same (growth of stoichiometric  $\langle AIN \rangle$ ). Then, the site density in Eq. [\(3](#page-2-0)) is calculated in order to satisfy this equality. This balance is performed by the CFD-ACE multiphysics software during calculations.

The validity of the modeling already presented in Ref. [\[27](#page-12-0)] will be confronted to new experimental data. The numerical modeling will give an accurate estimation of growth rates, but will also allow calculating the surface site ratio  $N(S)$ –Al $(S)$  on AlN and estimating the surface supersaturation of gaseous species in front of the growing AlN surface. Indeed, the surface partial pressure of precursors strongly differs from the inlet partial pressure in diffusion-limited regime. These local site ratios and supersaturations will be confronted to the AlN crystal quality to evaluate if there is an apparent correlation.

General trends on the correlation between calculated gasphase supersaturation and morphology of CVD films have already been reported in Ref. [\[35\]](#page-12-0). It is generally accepted that with increasing supersaturation, the sequence obtained for the morphology is epitaxial layer, whiskers, polycrystallites, amorphous layer, and finally homogeneous nucleation.

Several authors have defined supersaturation in the past [\[36–41](#page-12-0)], with different forms within a general concept that represents the difference in thermodynamic potential between the initial and the final equilibrium states. The sum or product of the reactants pressure divided by the sum or product of the species pressure at equilibrium can be used in equations. The sum presents the obvious property of being nondimensional whatever the number of species taken into account in each side of the division. In the frame of our study, the local gas supersaturation near the growing surface has been defined as follows:

$$
\beta = \frac{P_{\text{AlCl}_3} + P_{\text{NH}_3} + 2P_{\text{AlCl}_2\text{NH}_2}}{P_{\text{Al}}^* + 2P_{\text{N}_2}^*}
$$
(5)

where  $P$  is the calculated pressure of species contributing to AlN growth just in front of the growing surface and  $P^*$ the equilibrium vapor pressure of gases  $(N_2 \text{ and } Al)$  in equilibrium with  $\langle AIN \rangle$  in a closed volume, at the same temperature, without hydrogen. As  $\langle AIN \rangle$  becomes theoretically unstable in hydrogen between 1,200 and 1,300  $^{\circ}$ C [\[27](#page-12-0)], the "AlN vapor pressure" in  $H_2$  has no meaning (i.e., the supersaturation is equal to unity in the range 1,200– 1,600  $^{\circ}$ C). This confirms that the growth conditions are far from thermodynamic equilibrium inasmuch as AlN growth is observed in conditions where  $\langle \text{AlN} \rangle$  is theoretically unstable (under  $H_2$  for HT–HVPE). The values  $P^*$  are calculated with Factsage 6.0 software package, with the SGTE database [[42\]](#page-12-0).

All the calculations performed here have been conducted with updated databases for transport, thermodynamic, and kinetic properties. In consequence, the modeling results related to the experimental data of Ref. [\[27](#page-12-0)] may be slightly different from the modeling results presented here.

#### 2.3 Experiments

Experiments were carried out in a vertical tubular quartz cold-wall HTCVD reactor depicted in Ref.  $[27]$  $[27]$ . AlCl<sub>3</sub> is produced in situ by direct chlorination of high-purity aluminum pellets (99.9999 %) under a  $Cl_2$  flux at 640 °C [\[14](#page-11-0)]. The aluminum precursor is then mixed with  $NH<sub>3</sub>$  previously diluted in a large amount of  $H_2$  as carrier and reactive gas. The total pressure is 1,333 Pa. Quarters of two-inchdiameter c-plane  $(\pm 0.1^{\circ})$  sapphire wafers are used as substrates. Prior to the injection of reactive gases, both aluminum pellets and substrate are slowly heated under H2. A 10-min annealing under  $H_2$  at 1,100 °C is used to prepare the sapphire surface (etching of carbon contaminants); then, the ramp-up to the growing temperature is made

<span id="page-4-0"></span>



under  $0.02$ :NH<sub>3</sub>/1:H<sub>2</sub> at 20 °C min<sup>-1</sup>. At the growth temperature, chlorine is injected into the aluminum chloride generator and  $NH<sub>3</sub>$  is readily inserted into  $H<sub>2</sub>$ . At the end of the experiment,  $Cl_2$  flow is first stopped and then  $NH_3$ , and finally, the substrate is slowly cooled down under  $H_2$  at 10 °C min<sup>-1</sup>. The diluting gas flow rate during growth is systematically 1,000 sccm of  $H_2$ . AlN film thickness and surface morphology are measured and observed by scanning electron microscopy (SEM) or field emission gun (FEG)–SEM. The thicknesses are obtained by averaging 5–10 SEM thickness measurements at different locations of a cross section of the grown layer. The substrates are set on a graphite susceptor during experiments. It was observed that a N/Al ratio above 0.75 in the gas phase leads to a strong etching of the susceptor, releasing gaseous chlorine– carbon contaminants and thus limiting the lower value that can be explored for this parameter.

Table [3](#page-4-0) summarizes the experiments performed for this study. The experimental growth rates, experimental conditions, calculated local supersaturations, calculated surface sites ratio, and crystal characteristics are given. The crystal quality is either polycrystalline (X-ray diffraction measurements in Bragg–Brentano configuration presenting many peaks related to different AlN orientations), or epitaxial (only (000n) AlN diffraction peaks). The FWHM (full width at half maximum) of the (0002) diffraction peak of XRD rocking curves is indicated as a criterion for crystalline quality of epitaxial AlN layers (wurtzite phase).

Every growth experiment has been conducted at least two times to ensure the repeatability. The thickness uncertainty is under 10 % of the total layer thickness.

Table [3](#page-4-0) presents three distinct sets of experiments (see comments column of Table [3](#page-4-0)). The first set has been already presented and discussed in our previous study [\[27](#page-12-0)]. It was used to calibrate the numerical modeling of AlN growth. AlN templates were deliberately used to overcome heteroepitaxial issues with growth kinetics. Only one condition leads to AlN epitaxial growth on AlN template. It is obtained for the lowest supersaturation and N/Al ratio. Epitaxial growth cannot be attributed to one parameter or the other [[27\]](#page-12-0).

The second set of experiments has been designed with mass transport modeling in order to vary the N/Al ratio in gas phase at nearly constant surface supersaturation and for a constant layer thickness. Due to the high growth rates obtained, none of the experiments led to epitaxial growth.

Finally, the objective of the third set of experiments also designed by mass transport modeling was to assess the influence of N/Al ratio in gas phase at constant thickness and constant growth rate of the layer. More detailed results can be found in Ref. [\[43\]](#page-12-0). Additionally, in this set of experiments, the influence of a low-temperature protective layer prior the growth of thick high-temperature AlN layers is presented. All



Fig. 1 Parity diagram between experimental and theoretical growth rates

the experiments led to epitaxial layers, with varying crystal qualities due to the presence of etching voids at the sapphire/ AlN interface. These voids are suppressed by growing a lowtemperature protective layer prior to the growth of the hightemperature thick AlN layer. The growth is then divided into two steps: a low-temperature step that leads to the deposition of an AlN protective layer and a high-temperature step that leads to the deposition of the thick AlN layer. The process flowcharts can be found in Ref. [[43](#page-12-0)]. Both the low-temperature protective layer and the high-temperature layer step are described in Table [3](#page-4-0) (two last lines).

The calculated supersaturation can reach values below 1  $[Log<sub>10</sub>(\beta)<sub>0</sub>]$  due to the fact that theoretical equilibrium partial pressure of AlN  $(Al + 1/2N_2)$  on solid  $\langle AlN \rangle$  is greater than the partial pressure of reactants. The supersaturation values can be at least compared between each other.

The following discussion presents a local analysis of the major quantities at the growth front that influence or not the AlN crystal characteristics.

## 3 Results and discussion

#### 3.1 Experimental growth rates versus modeling

First, the overall parity diagram between simulated and experimental growth rates is presented in Fig. 1.

The optimization of surface kinetic parameters has been made by minimizing the relative difference between experimental and simulated growth rate, with a limited number of parameters (only the three activation temperatures for sticking coefficients). The modeling results

<span id="page-7-0"></span>

Fig. 2 Surface morphology at different temperatures for  $N/A = 4.5$  in the gas phase (calibration experiment set on AlN template). The figure also indicates the surface  $N(S)/A(S)$  ratio and the local  $Log<sub>10</sub>$  of the supersaturation



Fig. 3 Surface morphology for different N/Al ratio in the gas phase at constant temperature (calibration experiment set on AlN templates,  $T = 1,400 \degree C$ ). The figure also indicates the surface N(S)/Al(S) ratio and the local Log<sub>10</sub> of the supersaturation

overestimate the lower growth rates and underestimate the higher growth rates. To obtain a better fitting, more calibration parameters could be used. But the model accurately predicts the influence of temperature and N/Al ratio in gas phase on the growth rate in a large range of experimental conditions in both kinetic and transport-limited regimes.

# 3.2 Morphology versus supersaturation and N/Al ratio

As stated in the Numerical Modeling section, the morphology is strongly influenced by the local supersaturation of reactive species. The effect of N/Al ratio both at the surface and in the gas phase is also important.

<span id="page-8-0"></span>

Fig. 4 Surface morphology for different N/Al ratio in the gas phase at constant supersaturation (constant supersaturation set, growth on sapphire,  $T = 1,400 °C$ ). The figure also indicates the surface N(S)/Al(S) ratio, the gas-phase N/Al ratio, and the local Log<sub>10</sub> of the supersaturation



Fig. 5 Surface morphology for different N/Al ratio in the gas phase at constant growth rate and constant thickness (constant thickness and constant growth rate set, growth on sapphire,  $T = 1,500 °C$ ). The figure also indicates the surface N(S)/Al(S) ratio and the local Log<sub>10</sub> of the supersaturation

Typical SEM or FEG–SEM images of the surface morphologies of grown AlN layers with temperature, local supersaturation, gas-phase N/Al, and surface  $N(S)/A(S)$  ratios indicated are given in Figs. [2,](#page-7-0) [3](#page-7-0), 4, and 5.

Figure [2](#page-7-0) shows, with larger grains on the surface morphology, the enhancement of adatoms surface diffusion coefficients with an increase in the temperature. In spite of lower supersaturation at high temperature, the experiment performed at  $1,500$  °C does not lead to epitaxial growth.

<span id="page-9-0"></span>Figure [3](#page-7-0) illustrates the effect of N/Al ratio in the gas phase at nearly constant and low growth rate and slightly varying supersaturation (due to constant temperature  $T = 1,400$  °C). Only the experiment performed with  $N/A = 1.5$  led to epitaxial growth (FWHM  $= 3,204$  arcsec). It is worth noting that both the supersaturation and N/Al ratio in gas phase are in the lowest range for this set of experiments.

The set of experiments with growth at constant supersaturation (Log<sub>10</sub> $\beta \approx 1.15$ ,  $T = 1,400$  °C) on sapphire essentially led to polycrystalline growth with semi-polar planes of AlN parallel to the c-plane of the sapphire substrate AlN {10–11}//sapphire (0001). The typical surface morphologies are illustrated in Fig. [4](#page-8-0). All experiments leading to polycrystalline layers on sapphire globally follow this trend with small variations (presence of non-polar {1–100} or polar (0001) plane also) not easily correlated with any process parameter. The main orientation for polycrystalline growth on sapphire is always  $\langle 10–11 \rangle$ . Preceding studies on graphite as substrate [\[44](#page-12-0), [45\]](#page-12-0) (polycrystalline growth) report a wider variety of orientations and a temperature dependency of the main crystalline orientation.

Figure [5](#page-8-0) illustrates the effect of N/Al ratio in the gas phase at constant low growth rate and constant thickness  $(T = 1,500 \degree C)$ . It is worth noting that the supersaturations used here are low (Log10  $\beta$  < 0.7). All the experiments have led to epitaxial growth with FWHM of the grown layers ranging from 327 to 5,011 arcsec, with AlN (0001)// sapphire (0001) and AlN  $\langle 10–10 \rangle$ //sapphire  $\langle 11–20 \rangle$  crystal relationships. The dependency of FWHM to any process parameter is not obvious. The lowest values of FWHM have been observed for  $N/A = 1.5$  and 15 in the gas phase, essentially due to a better adherence of the AlN layer to the sapphire substrate for  $N/A = 15$ . The minimum for  $N/A1 = 1.5$ , as the layer is locally detached from sapphire, seems to originate from optimal local conditions only. The two last experiments of the set are assessing the following: (1) the quality of a low-temperature (1,200  $^{\circ}$ C) protective layer grown at very low growth rate  $(1 \mu m h^{-1})$ for a low N/Al ratio of 3 since other N/Al ratios led to polycrystalline layers, and (2) the subsequent growth at high temperature (1,500 °C) for a low N/Al ratio of 1.5. The main issue during this set of experiments which led to the synthesis of AlN single-crystal layers was to stabilize the AlN/sapphire interface that undergoes severe etching under hydrogen and high temperatures, degrading the AlN crystal quality [[43,](#page-12-0) [46](#page-12-0)].

The lowest misoriented layer (FWHM of 327 arcsec) was a 4  $\mu$ m AlN layer obtained at 1,500 °C, at a growth rate of 3.22  $\mu$ m h<sup>-1</sup> and N/Al of 1.5. This layer was grown on a thin 170 nm protective layer of AlN deposited at 1,200 °C, at a growth rate of 1  $\mu$ m h<sup>-1</sup> for a N/Al ratio of 3 in the gas phase.



Fig. 6 Dependency of the  $N(S)/A(S)$  ratio versus the N/Al ratio in the gas phase



Fig. 7 N/Al ratio in the gas phase versus experimental growth rate

#### 3.3 Experimental domain of epitaxial growth

We will now summarize the influence of each parameter, namely the growth rate, temperature, surface supersaturation, surface  $N(S)/A I(S)$ , and  $N/A I$  ratio in the gas phase on the crystalline state (epitaxial or polycrystalline) of the deposited AlN layer (the crystal quality of epitaxial layer is another issue not treated here).

Figure 6 presents the plot of  $N(S)/Al(S)$  surface ratio versus the N/Al ratio in the gas phase. These two parameters are completely correlated, and their evolution confirms that the Al containing species is controlling the growth in all of our experiments [The surface is covered with  $N(S)$  waiting for Al containing species]. Even for  $N/Al = 0.75$ , the surface kinetics of  $NH<sub>3</sub>$  is faster than that



of AlCl<sub>3</sub>, leading to a  $N(S)/Al(S)$  close to 1. This trend is shown on Fig. [6.](#page-9-0)

According to Fig. [6](#page-9-0), the rate limiting the AlN growth is the aluminum containing species kinetic (in HVPE conditions and according to our simplified kinetic scheme).

Consequently, N(S)/Al(S) ratio and N/Al ratio are not independent parameters of our study. For sake of clarity, only the N/Al ratio in gas phase will be now presented.

Figure [7](#page-9-0) presents a plot of the experimental growth rate versus the N/Al ratio in the gas phase for epitaxial and polycrystalline layers.

There is no clear correlation between epitaxial growth and N/Al ratio in the gas phase, but the AlN quality is clearly linked to the growth rate. However, at N/Al ratio of 1.5 in the gas phase, the process conditions leading to epitaxial growth seem to be broader in terms of range of growth rates. Clearly, the N/Al ratio parameter is not of prime interest for growing epitaxial AlN layers, but seems to play a role in the fine tuning of crystal quality. The optimum in crystal quality empirically observed at N/Al = 1.5 at 1,500 °C and 3 at 1,200 °C is not correlated to any particular local growing conditions. It could be emphasized that for these ratios, the surface diffusivity of absorbed species (not computed in our simple modeling) is related to the minimization of stacking faults during the AlN lattice growth.

Figure 8 represents the supersaturation versus the experimental growth rate for epitaxial and polycrystalline layers.

Figure 8 clearly allows defining of a process window where growing parameters lead to epitaxial growth of AlN on c-plane sapphire. This window ranges from relatively high supersaturation at low temperature, but very low growth rates to low supersaturation and growth rates as



Fig. 8 Surface supersaturation versus experimental growth rate Fig. 9 Experimental growth rate versus growth temperature

high as 10  $\mu$ m h<sup>-1</sup> (at high temperatures). This indicates that there is perhaps a way of improvement in term of higher growth rates of epitaxial layers by decreasing further the supersaturation (i.e., increasing the temperature). However, these conditions correspond also to the beginning of strong etching phenomena, both for AlN [[47\]](#page-12-0) and sapphire [\[48](#page-12-0)]. It is not sure that this area can be reached in high-temperature HVPE, mainly due to HCl by-product formation and  $H<sub>2</sub>$  etching.

As supersaturation is essentially dependent on temperature due to the exponential relationship between ''AlN vapor pressure'' and temperature, a simple way of plotting all our experiments could be to represent directly the experimental growth rate versus the growth temperature, as represented in Fig. 9. This figure allows defining a similar process window such as Fig. 8, showing that low-temperature epitaxial layers can be grown at very low growth rate and that high growth rate epitaxial layers can be reached only at high temperatures, whatever the N/Al ratio. It can be optimistically emphasized from this figure that the maximal growth rate that can be reached at  $1,500$  °C could be as high as 25  $\mu$ m h<sup>-1</sup> and at 1,600 °C could be around 70  $\mu$ m h<sup>-1</sup>, if no etching reactions occur. Growth rates as high as 57  $\mu$ m h<sup>-1</sup> at 1,450 °C and N/Al of 4 have been reported in another apparatus with similar process conditions [\[49](#page-12-0)].

# 3.4 Limitations of the numerical modeling to understand AlN epitaxial growth

Crystal quality is linked to high diffusion length of adatoms and low incorporation rate (so high temperature, low supersaturation). The numerical modeling proposed here does not take into account the surface diffusion mechanisms, essentially due to a lack of data in the literature.

<span id="page-11-0"></span>The only striking result of this modeling is that contrary to what is classically stated for III–V epitaxy, the IIIcompound containing species  $(AICI<sub>3</sub> and AICI<sub>2</sub>NH<sub>2</sub> here)$ have a relatively low apparent sticking coefficient (at most 0.013 for AlCl<sub>3</sub> and 0.056 for AlCl<sub>2</sub>NH<sub>2</sub> at 1,500 °C in our modeling). As gas-phase aluminum species seem to have a poor affinity with the growing surface (due to steric hindrance or species stability), their surface diffusion length could be high compared to other III-compound or other aluminum precursors.

During III–V epitaxial growth, the III–V ratio is classically kept to a high value to force a kinetic regime limited by a III-compound. This allows increasing the diffusion length of Al adatoms. This condition is clearly not necessary in the case of AlN epitaxial growth in HVPE conditions where the relatively poorly reactive AlCl<sub>3</sub> precursor is the main gas species for aluminum.

### 4 Conclusion

This study presents a numerical modeling of AlN growth by high-temperature-HVPE which comprises a relatively limited number of chemical reactions. This modeling was experimentally validated within a broad range of process parameters (temperature, N/Al ratio in gas phase). Moreover, this study permitted the calculation of various local parameters such as the surface site ratio and the local supersaturation during AlN growth. On the basis of numerical results, sets of experiments have been designed to assess the influence of N/Al ratio in the gas phase at constant supersaturation on the growing AlN surface or at constant AlN growth rate.

In one hand, it appears that the N/Al ratio or  $[N(S)]$  $A(0)$  is not of prime importance to explain the occurrence of epitaxial growth of AlN on sapphire. However, it has to be carefully controlled to enhance the crystal quality when epitaxial growth is obtained. It was found that (1) a value of  $N/A1 = 1.5$  in the gas phase is optimal to regrow AlN on AlN templates at high temperature and (2) a value of N/Al of 3 in the gas phase is necessary to grow a low-temperature protective layer. The simulated local conditions do not explain these ratios.

On the other hand, simple plots representing growth rate–supersaturation versus temperature clearly exhibit a process window where epitaxial growth could occur. As expected, high epitaxial growth rates need low supersaturation or high temperatures. Simple extrapolations indicate that the domain of improvement in epitaxial growth rates on sapphire is not so wide, due to the fast decomposition of sapphire and AlN at temperatures above  $1,500$  °C in hydrogen.

Mass transport simulation and theoretical chemistry were used to design sets of experiments. However, the lack of clear macroscopic relationship (except the well-established supersaturation/quality) between process conditions and crystal quality let us to think that at least another parameter plays a key role. It appears that the chemical stability of the interface AlN/sapphire during the first steps of growth is of prime interest compared to the thickening of the layer which is less sensitive to growth conditions. This primary step is also sensitive to N/Al ratio, but its dependence is still not clear. Use of such modeling for the design of experiments is an interesting tool and has to be pursued. The knowledge of local parameters near the growth surface which are reactor-independent could also be used in the development of new reactor configuration.

#### References

- 1. Kinoshita T, Hironaka K, Obata T, Nagashima T, Dalmau R, Schlesser R, Moody B, Xie JQ, Inoue S, Kumagai Y, Koukitu A, Sitar Z (2012) Appl Phys Express 5:122101
- 2. Ballandras S, Reinhardt A, Laude V, Soufyane A, Camou S, Daniau W, Pastureaud T, Steichen W, Lardat R, Solal M, Ventura P (2004) J Appl Phys 96:7731–7741
- 3. Piazza G, Felmetsger V, Muralt P, Olsson RH, Ruby R (2012) MRS Bull 37:1051–1061
- 4. Slack GA, McNelly TF (1977) J Cryst Growth 42:560–563
- 5. Tian W, Yan WY, Dai JN, Li SL, Tian Y, Hui X, Zhang JB, Fang YY, Wu ZH, Chen CQ (2013) J Phys D Appl Phys 46:065303
- 6. Nickel KG, Riedel R, Petzow G (1989) J Am Ceram Soc 72(10):1804–1810
- 7. Kumagai Y, Yamane T, Miyaji T, Murakami H, Kangawa Y, Koukitu A (2003) Phys Status Solidi (c) 7:2498–2501
- 8. Kumagai Y, Yamane T, Koukitu A (2005) J Cryst Growth 281:62–67
- 9. Freitas JA, Culbertson JC, Mastro MA, Kumagai Y, Koukitu A (2012) J Cryst Growth 350:33–37
- 10. Nagashima T, Harada M, Yanagi H, Kumagai Y, Koukitu A, Takada K (2007) J Cryst Growth 300:42–44
- 11. Eriguchi K, Murakami H, Panyukova U, Kumagai Y, Ohira S, Koukitu A (2007) J Cryst Growth 298:332–335
- 12. Eriguchi K, Hiratsuka T, Marukami H, Kumagai Y, Koukitu A (2008) J Cryst Growth 310:4016–4019
- 13. Tajima J, Murakami H, Kumagai Y, Takada K, Koukitu A (2009) J Cryst Growth 311:2837–2839
- 14. Claudel A, Blanquet E, Chaussende D, Audier M, Pique D, Pons M (2009) J Cryst Growth 311:3371–3379
- 15. Balaji M, Claudel A, Fellmann V, Gelard I, Blanquet E, Boichot R, Pierret A, Attal-Tretout B, Crisci A, Coindeau S, Roussel H, Pique D, Baskar K, Pons M, Alloy J (2012) Compound 526:103–109
- 16. Claudel A, Blanquet E, Chaussende D, Boichot R, Doisneau B, Berthome G, Crisci A, Mank H, Moisson C, Pique D, Pons M (2011) J Cryst Growth 335:17–24
- 17. Allendorf MD, Osterheld TH In: Proceedings of the electrochemical society, vol 96-5
- 18. Mc Daniel AH, Allendorf MD (1998) J Phys Chem A 102:7804–7812
- <span id="page-12-0"></span>19. McDaniel AH, Allendorf MD (1997) In: Proceedings of the CVDXIV/EUROCVD 11, vol 97-25. The Electrochemical Society, Pennington, p 40
- 20. Allendorf MD, Melius CF, Osterheld TH (1996) Mater Res Soc Proc 410:459
- 21. Miller JA, Bowman CT (1989) Prog Energy Combust Sci 15:287
- 22. Dollet A, Casaux Y, Chaix G, Dupuy C (2002) Thin Solid Films 406:1–16
- 23. Swihart MT, Catoire L, Legrand B, Gökalp I, Paillard C (2003) Combust Flame 132:91–101
- 24. Cai D, Zheng LL, Zhang H, Tassev VL, Bliss VF (2005) J Cryst Growth 276:182–193
- 25. Cai D, Zheng LL, Zhang H, Tassev VL, Bliss VF (2006) J Cryst Growth 293:136–145
- 26. Segal AS, Bazarevskiy DS, Bogdanov MV, Yakovlev EV (2009) Phys Status Solidi (c) 6:S329–S332
- 27. Boichot R, Claudel A, Baccar N, Milet A, Blanquet E, Pons M (2010) Surf Coat Technol 205:1294
- 28. Kee RJ, Rupley FM, Meeks E, Miller JA Chemkin-III: a FOR-TRAN chemical kinetics package for the analysis of gas phase chemical and plasma kinetics. Sandia National Laboratories. Livermore, California
- 29. Bird RB, Stewart WE, Lighfoot EN (2002) Transport phenomena, 2nd edn. Wiley, New York
- 30. Svehla RA (1962) NASA technical report R-132. NASA-TR-R-132
- 31. Taylor R, Krishna R (1993) Multicomponent mass transfer. Wiley, New York
- 32. JANAF Thermochemical tables, 3rd edn. The American Chemical Society (Washington) and the American Institute of Physics (New York) for the National Bureau of Standards, USA
- 33. <http://www.sandia.gov/HiTempThermo/>
- 34. Coltrin ME, Kee RJ, Rupley FM (1990) Surface CHEMKIN. Sandia National Laboratories. Livermore, California
- 35. Blocher JM (1974) J Vac Sci Technol 11(4):680–686
- 36. Räback P, Nieminen R, Yakimova R, Tuominen M, Janzén E (1997) J Electrochem Soc 144:1024–1027
- 37. Bloem J, Oei YS, de Moor HHC, Hansen JHL, Giling LJ (1985) J Electron Soc 132(8):1973–1980
- 38. Hwang NM, Yoon DY (1994) J Mater Sci Lett 13(19):1437–1439
- 39. Giling LJ, De Moor HHC, Jacobs WPJH, Saaman AA (1986) J Cryst Growth 78(2):303–321
- 40. Tai C, Shih C-Y (1996) J Cryst Growth 160(1–2):186–189
- 41. Vahlas C, Hwang NM, Gueroudji L, Maury F (1998) Chem Vap Depos 4(3):96–99
- 42. <http://www.sgte.org/>
- 43. Boichot R, Coudurier N, Mercier F, Lay S, Crisci A, Coindeau S, Claudel A, Blanquet E, Pons M (2013) Surf Coat Technol. doi:[10.](http://dx.doi.org/10.1016/j.surfcoat.2013.08.016) [1016/j.surfcoat.2013.08.016](http://dx.doi.org/10.1016/j.surfcoat.2013.08.016)
- 44. Pons M, Boichot R, Coudurier N, Claudel A, Blanquet E, Lay S, Mercier F, Pique D (2013) Surf Coat Technol 230:111–118
- 45. Claudel A, Blanquet E, Chaussende D, Boichot R, Martin R, Mank H, Crisci A, Doisneau B, Chaudouet P, Coindeau S, Pique D, Pons M (2011) J Electrochem Soc 158(3):H328–H332
- 46. Kumagai Y, Enatsu Y, Ishizuki M, Kubota Y, Tajima J, Nagashima T, Murakami H, Takada K, Koukitu A (2010) J Cryst Growth 312:2530–2536
- 47. Claudel A, Blanquet E, Chaussende D, Boichot R, Doisneau B, Berthomé G, Crisci A, Mank H, Moisson C, Pique D, Pons M (2011) J Cryst Growth 335(1):17–24
- 48. Akiyama K, Araki T, Murakami H, Kumagai Y, Koukitu A (2007) Phys Status Solidi (c) 4(7):2297–2300
- 49. Nagashima T, Harada M, Yanagi H, Fukuyama H, Kumagai Y, Koukitu A, Takada K (2007) J Cryst Growth 305:355–359