

Moser–Trudinger inequalities for singular Liouville systems

Luca Battaglia¹

Received: 15 April 2015 / Accepted: 27 October 2015 / Published online: 23 November 2015 © Springer-Verlag Berlin Heidelberg 2015

Abstract In this paper we prove a Moser–Trudinger inequality for the Euler–Lagrange functional of general singular Liouville systems on a compact surface. We characterize the values of the parameters which yield coercivity for the functional, hence the existence of energyminimizing solutions for the system, and we give necessary conditions for boundedness from below. We also provide a sharp inequality under assuming the coefficients of the system to be non-positive outside the diagonal. The proofs use a concentration-compactness alternative, Pohožaev-type identities and blow-up analysis.

Keywords Liouville systems · Moser–Trudinger inequality · Coercivity · Minimizing solutions

Mathematics Subject Classification 35A23 · 35J47 · 35J50 · 58E35

1 Introduction

An essential tool in the study of the embeddings of Sobolev spaces is the Moser–Trudinger inequality, which gives compact embedding in any L^p space for finite $p \ge 1$ and also exponential integrability.

If we consider a 2-dimensional compact Riemannian manifold (Σ, g) , due to well-known works from Moser [18] and Fontana [13] we get

$$\log \int_{\Sigma} e^{u} \mathrm{d}V_{g} - \int_{\Sigma} u \mathrm{d}V_{g} \le \frac{1}{16\pi} \int_{\Sigma} |\nabla u|^{2} \mathrm{d}V_{g} + C \qquad \forall u \in H^{1}(\Sigma),$$
(1)

☑ Luca Battaglia lbatta@sissa.it

The author has been supported by the PRIN project Variational and perturbative aspects of nonlinear differential problems.

¹ Université Catholique de Louvain, Institut de Recherche en Mathématique et Physique, Chemin du Cyclotron 2, 1348 Louvain-la-Neuve, Belgium

where $\nabla = \nabla_g$ is the gradient given by the metric g and $C = C_{\Sigma,g}$ is a constant depending only on Σ and g.

This inequality has fundamental importance in the study of the Liouville equations of the kind

$$-\Delta u = \rho \left(\frac{he^u}{\int_{\Sigma} he^u \mathrm{d}V_g} - 1 \right),\tag{2}$$

where $\Delta = \Delta_g$ is the Laplace-Beltrami operator, ρ a positive real parameter, *h* a positive smooth function and Σ is supposed, without loss of generality, to have area equal to $|\Sigma| = 1$. In fact, the solutions of (2) are critical points of the functional

$$I_{\rho}(u) = \frac{1}{2} \int_{\Sigma} |\nabla u|^2 \mathrm{d}V_g - \rho \left(\log \int_{\Sigma} h e^u \mathrm{d}V_g - \int_{\Sigma} u \mathrm{d}V_g \right);$$

using the inequality (1) we can control the last term by the Dirichlet energy, thus showing that I_{ρ} is bounded from below on $H^{1}(\Sigma)$ if and only if ρ is smaller or equal to 8π .

Equations like (2) have great importance in different contexts like the Gaussian curvature prescription problem (see for instance [6,7]) and abelian Chern–Simons models in theoretical physics ([21,24]).

An extension of the inequality (1), which takes into consideration power-type weights, was given by Chen [8] and Trojanov [22]. For a given $p \in \Sigma$ and $\alpha \in (-1, 0]$, they showed that

$$(1+\alpha)\left(\log\int_{\Sigma} d(\cdot, p)^{2\alpha} e^{u} \mathrm{d}V_{g} - \int_{\Sigma} u \mathrm{d}V_{g}\right) \leq \frac{1}{16\pi} \int_{\Sigma} |\nabla u|^{2} \mathrm{d}V_{g} + C \qquad \forall u \in H^{1}(\Sigma).$$
(3)

This inequality allows to treat singularities in the Eq. (2), that is to consider equations like

$$-\Delta u = \rho \left(\frac{he^u}{\int_{\Sigma} he^u \mathrm{d}V_g} - 1\right) - 4\pi \sum_{m=1}^M \alpha_m (\delta_{p_m} - 1), \tag{4}$$

where we take arbitrary $p_1, \ldots, p_M \in \Sigma$ and $\alpha_m > -1$ for any $m \in \{1, \ldots, M\}$.

This is a natural extension of (2), which allows to consider the same problems in a more general context. For instance, it arises in the Gaussian curvature prescription problem on surfaces with conical singularities and in Chern–Simons vortices theory.

Defining G_p as the Green function of $-\Delta$ on Σ centered at a point p, through the change of variables

$$u \mapsto u + 4\pi \sum_{m=1}^{M} \alpha_m G_{p_m} \tag{5}$$

Equation (4) becomes

$$-\Delta u = \rho \left(\frac{\tilde{h}e^u}{\int_{\Sigma} \tilde{h}e^u \mathrm{d}V_g} - 1 \right)$$

with $\tilde{h} = he^{-4\pi \sum_{m=1}^{M} \alpha_m G_{p_m}}$.

Since G_p has the same behavior as $\frac{1}{2\pi} \log \frac{1}{d(\cdot,p)}$ around p, then \tilde{h} behaves like $d(\cdot, p_m)^{2\alpha_m}$ around each singular point p_m , hence the energy functional

$$I_{\rho}(u) = \frac{1}{2} \int_{\Sigma} |\nabla u|^2 \mathrm{d}V_g - \rho \left(\log \int_{\Sigma} \tilde{h} e^u \mathrm{d}V_g - \int_{\Sigma} u \mathrm{d}V_g \right)$$

can be estimated, as in the regular case, using (3).

The purpose of this paper is to extend inequality (3) to singular Liouville systems of the type

$$-\Delta u_i = \sum_{j=1}^{N} a_{ij} \rho_j \left(\frac{h_j e^{u_j}}{\int_{\Sigma} h_j e^{u_j} dV_g} - 1 \right) - 4\pi \sum_{m=1}^{M} \alpha_{im} (\delta_{p_m} - 1), \quad i = 1, \dots, N,$$

where $A = (a_{ij})$ is a $N \times N$ symmetric positive definite matrix and ρ_i , h_i , α_{im} are as before.

Applying, similarly to (5), the change of variables

$$u_i \mapsto u_i + 4\pi \sum_{m=1}^M \alpha_{im} G_{p_m},$$

the system becomes

$$-\Delta u_i = \sum_{j=1}^N a_{ij} \rho_j \left(\frac{\widetilde{h}_j e^{u_j}}{\int_{\Sigma} \widetilde{h}_j e^{u_j} \mathrm{d}V_g} - 1 \right), \quad i = 1, \dots, N,$$
(6)

with \tilde{h}_i having the same behavior around the singular points.

The system has a variational formulation with the energy functional

$$J_{\rho}(u) := \frac{1}{2} \sum_{i,j=1}^{N} a^{ij} \int_{\Sigma} \nabla u_i \cdot \nabla u_j \mathrm{d}V_g - \sum_{i=1}^{N} \rho_i \left(\log \int_{\Sigma} \widetilde{h}_i e^{u_i} \mathrm{d}V_g - \int_{\Sigma} u_i \mathrm{d}V_g \right), \quad (7)$$

with a^{ij} indicating the entries of the inverse matrix A^{-1} of A.

A recent paper by the author and Malchiodi ([2]) gives an answer for the particular case of the SU(3) Toda system, that is N = 2 and A is the Cartan matrix

$$\begin{pmatrix} 2 & -1 \\ -1 & 2 \end{pmatrix}.$$

This is a particularly interesting case, due to its application in the description of holomorphic curves in \mathbb{CP}^N in geometry ([3,5,9]) and in the non-abelian Chern–Simons theory in physics ([12,21,24]).

The authors prove a sharp inequality, that is they show that the functional J_{ρ} is bounded from below if and only if both the parameters ρ_i are less or equal than $4\pi \min\{1, 1 + \min_m \alpha_{im}\}$, thus extending the result in the regular case from [15].

Concerning general regular Liouville systems, Wang [23] gave necessary and sufficient conditions for the boundedness from below of J_{ρ} , following previous results in [10,11] for the problem on Euclidean domains with Dirichlet boundary conditions. Analogous results were given in [20] for the standard unit sphere (S^2 , g_0) and in [19] for a similar problem.

In these papers, the authors introduce, for any $\mathcal{I} \subset \{1, ..., N\}$, the following function of the parameter ρ :

$$\Lambda_{\mathcal{I}}(\rho) = 8\pi \sum_{i \in \mathcal{I}} \rho_i - \sum_{i,j \in \mathcal{I}} a_{ij} \rho_i \rho_j.$$

What they prove is boundedness from below for J_{ρ} for any $\rho \in \mathbb{R}^{N}_{+}$ which satisfies $\Lambda_{\mathcal{I}}(\rho) > 0$ for all the subsets \mathcal{I} of $\{1, \ldots, N\}$, whereas $\inf_{H^{1}(\Sigma)^{N}} J_{\rho} = -\infty$ whenever $\Lambda_{\mathcal{I}}(\rho) < 0$ for some $\mathcal{I} \subset \{1, \ldots, N\}$.

The first main result of this paper is an extension of the results from [10, 11, 23] to the case of singularities.

Similarly to Liouville equation, we will have to multiply some quantities by $1 + \alpha_{im}$. Precisely, we have:

Theorem 1.1 Let J_{ρ} be the functional defined by (7) and set, for $\rho \in \mathbb{R}_{>0}^{N}$, $x \in \Sigma$ and $i \in \mathcal{I} \subset \{1, \ldots, N\}$:

$$\alpha_{i}(x) = \begin{cases} \alpha_{im} \text{ if } x = p_{m} \\ 0 \text{ otherwise} \end{cases} \quad \Lambda_{\mathcal{I},x}(\rho) \coloneqq 8\pi \sum_{i \in \mathcal{I}} (1 + \alpha_{i}(x))\rho_{i} - \sum_{i,j \in \mathcal{I}} a_{ij}\rho_{i}\rho_{j} \\ \Lambda(\rho) \coloneqq \min_{\mathcal{I} \subset \{1,\dots,N\}, x \in \Sigma} \Lambda_{\mathcal{I},x}(\rho). \end{cases}$$
(8)

Then, J_{ρ} is bounded from below if $\Lambda(\rho) > 0$, whereas J_{ρ} is unbounded from below if $\Lambda(\rho) < 0$.

Notice that, in the definition of Λ , the minimum makes sense because it is taken in a finite set, since $\alpha_i(x) = 0$ for all points of Σ but a finite number, and for all the former points $\Lambda_{\mathcal{I},x}$ is defined in the same way.

As a consequence of this result, we obtain information about the existence of solutions for the system (6).

Corollary 1.2 The functional J_{ρ} is coercive in $\overline{H}^{1}(\Sigma)$ if and only if $\Lambda(\rho) > 0$. Therefore, if this occurs, then J_{ρ} admits a minimizer u which solves (6).

Theorem 1.1 leaves an open question about what happens when $\Lambda(\rho) = 0$. In this case, as we will see in the following Sections, one encounters blow-up phenomena which are not yet fully known for general systems.

Anyway, we can say something more if we assume in addition $a_{ij} \leq 0$ for any $i, j \in \{1, ..., N\}$ with $i \neq j$. First of all, we notice that in this case

$$\Lambda(\rho) = \min_{i \in \{1, \dots, N\}} \left(8\pi (1 + \widetilde{\alpha}_i) \rho_i - a_{ii} \rho_i^2 \right), \quad \text{where}$$
$$\widetilde{\alpha}_i := \min_{m \in \{1, \dots, M\}, x \in \Sigma} \alpha_i(x) = \min \left\{ 0, \min_{m \in \{1, \dots, M\}} \alpha_{im} \right\}; \quad (9)$$

hence the sufficient condition in Theorem 1.1 is equivalent to assuming $\rho_i < \frac{8\pi(1+\tilde{\alpha}_i)}{a_{ii}}$ for any *i*.

With this assumption, studying what happens when $\Lambda_{\mathcal{I}}(\rho) = 0$ is reduced to a singlecomponent local blow-up, which can be treated by using an inequality from [1]. Therefore, we get the following sharp result:

Theorem 1.3 Let J_{ρ} be defined by (7), $\tilde{\alpha}_i$ as in (9) and $\Lambda(\rho)$ as in Theorem 1.1, and suppose $a_{ij} \leq 0$ for any $i, j \in \{1, ..., N\}$ with $i \neq j$.

Then, J_{ρ} is bounded from below on $H^{1}(\Sigma)^{N}$ if and only if $\Lambda(\rho) \geq 0$, namely if and only if $\rho_{i} \leq \frac{8\pi(1+\widetilde{\alpha}_{i})}{a_{ii}}$ for any $i \in \{1, \ldots, N\}$.

We remark that the assuming A to be positive definite is necessary. If it is not, then J_{ρ} is unbounded from below for any ρ .

In fact, suppose there exists $v \in \mathbb{R}^N$ such that $\sum_{i,j=1}^N a^{ij} v_i v_j \le -\theta |v|^2$ for some $\theta > 0$. Then, we consider the family of functions $u^{\lambda}(x) := \lambda v \cdot x$; by Jensen's inequality we get

$$J_{\rho}\left(u^{\lambda}\right) \leq \frac{1}{2} \sum_{i,j=1}^{N} a^{ij} \int_{\Sigma} \nabla u_{i}^{\lambda} \cdot \nabla u_{j}^{\lambda} \mathrm{d}V_{g} - \sum_{i=1}^{N} \rho_{i} \int_{\Sigma} \log \tilde{h}_{i} \mathrm{d}V_{g}$$
$$\leq -\frac{\theta}{2} \lambda^{2} |v|^{2} + C$$
$$\underset{n \to +\infty}{\longrightarrow} -\infty.$$

We also notice that, with respect to the scalar case, in Theorem 1.1 and Corollary 1.2 the positive coefficients α_{im} 's may affect the definition of $\Lambda(\rho)$, hence the conditions for coercivity and boundedness from below of J_{ρ} .

On the other hand, under the assumptions of Theorem 1.3, coercivity and boundedness from below only depend on the negative α_{im} 's, just like for the scalar equation.

The plan of this paper is the following: in Sect. 2 we will introduce some notations and some preliminary results which will be used throughout the rest of the paper. In Sect. 3 we will show a sort of Concentration-compactness theorem, showing the possible non-compactness phenomena for solutions of the system (6). Finally, in Sects. 4 and 5 we will give the proof of the two main theorems.

2 Notations and preliminaries

In this section, we will give some useful notation and some known preliminary results which will be needed to prove the two main theorems.

Given two points $x, y \in \Sigma$, we will indicate the metric distance on Σ between them as d(x, y). We will indicate the open metric ball centered in *p* having radius *r* as

$$B_r(x) := \{ y \in \Sigma : d(x, y) < r \}.$$

For any subset of a topological space $A \subset X$ we indicate its closure as \overline{A} and its interior part as \mathring{A} .

Given a function $u \in L^1(\Sigma)$, the symbol \overline{u} will indicate the average of u on Σ . Since we assume $|\Sigma| = 1$, we can write:

$$\overline{u} = \int_{\Sigma} u \mathrm{d} V_g = \oint_{\Sigma} u \mathrm{d} V_g.$$

We will indicate the subset of $H^1(\Sigma)$ which contains the functions with zero average as

$$\overline{H}^{1}(\Sigma) := \left\{ u \in H^{1}(\Sigma) : \overline{u} = 0 \right\}.$$

Since the functional J_{ρ} defined by (7) is invariant by addition of constants, it will not be restrictive to study it on $\overline{H}^{1}(\Sigma)^{N}$ rather than on $H^{1}(\Sigma)^{N}$.

We will indicate with the letter *C* large constants which can vary among different lines and formulas. To underline the dependence of *C* on some parameter α , we indicate with C_{α} and so on.

We will denote as $o_{\alpha}(1)$ quantities which tend to 0 as α tends to 0 or to $+\infty$ and we will similarly indicate bounded quantities as $O_{\alpha}(1)$, omitting in both cases the subscript(s) when it is evident from the context.

First of all, we need a result from Brezis and Merle [4]. It is a classical estimate about exponential integrability of solutions of some elliptic PDEs.

Lemma 2.1 ([4], Theorem 1) Take r > 0, $\Omega := B_r(0) \subset \mathbb{R}^2$, $f \in L^1(\Omega)$ with $||f||_{L^1(\Omega)} < 4\pi$ and u solving

$$\begin{cases} -\Delta u = f \text{ in } \Omega\\ u = 0 \quad \text{on } \partial \Omega \end{cases}$$

Then, for any $q \in \left[1, \frac{4\pi}{\|f\|_{L^1(\Omega)}}\right)$ there exists a constant $C = C_{q,r}$ such that $\int_{\Omega} e^{q|u(x)|} dx \le C$.

A crucial role in the proof of both Theorems 1.1 and 1.3 will be played by the concentration values of the sequences of solutions of (6).

For a sequence $u^n = \{u_1^n, \ldots, u_N^n\}_{n \in \mathbb{N}}$ of solutions of (6) with $\rho = \rho^n = \{\rho_1^n, \ldots, \rho_N^n\}$, we define (up to subsequences), for $i \in \{1, \ldots, N\}$, the concentration value of its i^{th} component around a point $x \in \Sigma$ as

$$\sigma_i(x) := \lim_{r \to 0} \lim_{n \to +\infty} \rho_i^n \frac{\int_{B_r(x)} \tilde{h}_i e^{u_i^n} \mathrm{d}V_g}{\int_{\Sigma} \tilde{h}_i e^{u_i^n} \mathrm{d}V_g}.$$
 (10)

In a recent paper ([16], see also [14] for the regular case) it was proved, by a Pohožaev identity, that the concentration values satisfy the following algebraic relation, which involves the same quantities as in Theorem 1.1:

Proposition 2.2 ([14], Lemma 2.2; [16], Proposition 3.1) Let $\{u^n\}_{n\in\mathbb{N}}$ be a sequence of solutions of (6), $\alpha_i(x)$ and $\Lambda_{\mathcal{I},x}$ as in (8) and $\sigma(x) = (\sigma_1(x), \ldots, \sigma_N(x))$ as in (10). Then,

$$\Lambda_{\{1,...,N\},x}(\sigma(x)) = 8\pi \sum_{i=1}^{N} (1 + \alpha_i(x))\sigma_i(x) - \sum_{i,j=1}^{N} a_{ij}\sigma_i(x)\sigma_j(x) = 0.$$

To study the concentration phenomena of solutions of (6) we will use the following simple but useful calculus Lemma:

Lemma 2.3 ([15], Lemma 4.4) Let $\{a^n\}_{n\in\mathbb{N}}$ and $\{b^n\}_{n\in\mathbb{N}}$ two sequences of real numbers satisfying

$$a^n \xrightarrow[n \to +\infty]{} +\infty \qquad \lim_{n \to +\infty} \frac{b^n}{a^n} \le 0.$$

Then, there exists a smooth function $F : [0, +\infty) \to \mathbb{R}$ which satisfies, up to subsequences,

$$0 < F'(t) < 1 \quad \forall t > 0 \qquad F'(t) \xrightarrow[t \to +\infty]{} 0 \qquad F\left(a^n\right) - b^n \xrightarrow[n \to +\infty]{} +\infty.$$

Finally, as anticipated in the introduction, we will need a singular Moser–Trudinger inequality for Euclidean domains by Adimurthi and Sandeep [1], and its straightforward corollary.

Theorem 2.4 ([1], Theorem 2.1) For any r > 0, $\alpha \in (-1, 0]$ there exists a constant $C = C_{\alpha,r}$ such that if $\Omega := B_r(0) \subset \mathbb{R}^2$ and $u \in H_0^1(\Omega)$, then

$$\int_{\Omega} |\nabla u(x)|^2 \mathrm{d}x \le 1 \Rightarrow \int_{\Omega} |x|^{2\alpha} e^{4\pi (1+\alpha)u(x)^2} \mathrm{d}x \le C$$

🖉 Springer

Corollary 2.5 For any r > 0, $\alpha \in (-1, 0]$ there exists a constant $C = C_{\alpha,r}$ such that if $\Omega := B_r(0) \subset \mathbb{R}^2$ and $u \in H_0^1(\Omega)$, then

$$(1+\alpha)\log\int_{\Omega}|x|^{2\alpha}e^{u(x)}\mathrm{d}x \leq \frac{1}{16\pi}\int_{\Omega}|\nabla u(x)|^{2}\mathrm{d}x + C$$

Proof By the elementary inequality $u \le \theta u^2 + \frac{1}{4\theta}$ with $\theta = \frac{4\pi (1+\alpha)}{\int_{\Omega} |\nabla u(y)|^2 dy}$ we get

$$(1+\alpha)\log\int_{\Omega}|x|^{2\alpha}e^{u(x)}dx \le (1+\alpha)\log\int_{\Omega}|x|^{2\alpha}e^{\theta u(x)^{2}+\frac{1}{4\theta}}dx$$
$$=\frac{1}{16\pi}\int_{\Omega}|\nabla u(y)|^{2}dy + (1+\alpha)\log$$
$$\times\int_{\Omega}|x|^{2\alpha}e^{4\pi(1+\alpha)\left(\frac{u(x)}{\sqrt{\int_{\Omega}|\nabla u(y)|^{2}dy}}\right)^{2}}dx$$
$$\le \frac{1}{16\pi}\int_{\Omega}|\nabla u(y)|^{2}dy + C.$$

3 A Concentration-compactness theorem

The aim of this section is to prove a result which describes the concentration phenomena for the solutions of (6), extending what was done for the two-dimensional Toda system in [2, 17].

We actually have to normalize such solutions to bypass the issue of the invariance by translation by constants and to have the parameter ρ multiplying only the constant term.

In fact, for any solution u of (6) the functions

$$v_i := u_i - \log \int_{\Sigma} \widetilde{h}_i e^{u_i} \mathrm{d}V_g + \log \rho_i \tag{11}$$

solve

$$\begin{cases} -\Delta v_i = \sum_{j=1}^N a_{ij} \left(\tilde{h}_j e^{v_j} - \rho_j \right) \\ \int_{\Sigma} \tilde{h}_i e^{v_i} dV_g = \rho_i \end{cases} \quad i = 1, \dots, N.$$
(12)

Moreover, we can rewrite in a shorter way (10) as

$$\sigma_i(x) = \lim_{r \to 0} \lim_{n \to +\infty} \int_{B_r(x)} \tilde{h}_i^n e^{v_i^n} \mathrm{d}V_g.$$

For such functions, we get the following concentration-compactness alternative:

Theorem 3.1 Let $\{u^n\}_{n\in\mathbb{N}}$ be a sequence of solutions of (6) with $\rho^n \xrightarrow[n \to +\infty]{} \rho \in \mathbb{R}^N_+$ and $\widetilde{h}^n_i = V^n_i \widetilde{h}_i$ with $V^n_i \xrightarrow[n \to +\infty]{} 1$ in $C^1(\Sigma)^N$, $\{v^n\}_{n\in\mathbb{N}}$ be defined as in (11) and S_i be defined, for $i \in \{1, ..., N\}$, by

$$S_{i} := \left\{ x \in \Sigma : \exists x^{n} \underset{n \to +\infty}{\longrightarrow} x \text{ such that } v_{i}^{n} \left(x^{n} \right) \underset{n \to +\infty}{\longrightarrow} +\infty \right\}.$$
(13)

Then, up to subsequences, one of the following occurs:

- If $S_i = \emptyset$ for any $i \in \{1, ..., N\}$, then $v^n \xrightarrow[n \to +\infty]{} v$ in $W^{2,q}(\Sigma)^N$ for some q > 1 and some v which solves (12).
- If $S_i \neq \emptyset$ for some *i*, then it is a finite set for all such *i*'s. If this occurs, then there is a subset $\mathcal{I} \subset \{1, \ldots, N\}$ such that $v_j^n \xrightarrow{} -\infty$ in $L^{\infty}_{\text{loc}} \left(\Sigma \setminus \bigcup_{j'=1}^N S_{j'} \right)$ for any $j \in \mathcal{I}$ and $v_j^n \xrightarrow[n \to +\infty]{} v_j$ in $W_{\text{loc}}^{2,q} \left(\Sigma \setminus \bigcup_{j'=1}^N S_{j'} \right)$ for some q > 1 and some suitable v_j , for any $i \in \{1, \ldots, N\} \setminus \mathcal{I}.$

Since \tilde{h}_i is smooth outside the points p_m 's, the estimates in $W^{2,q}(\Sigma)$ are actually in $C^{2,\alpha}_{\text{loc}}\left(\Sigma \setminus \bigcup_{m=1}^{M} p_m\right)$ and the estimates in $W^{2,q}_{\text{loc}}\left(\Sigma \setminus \bigcup_{j'=1}^{N} S_{j'}\right)$ are actually in $C_{\text{loc}}^{2,\alpha}\left(\Sigma \setminus \left(\bigcup_{j'=1}^{N} \mathcal{S}_{j'} \cup \bigcup_{m=1}^{M} p_m\right)\right)$. Anyway, estimates in $W^{2,q}$ will suffice in most of the paper.

To prove Theorem 3.1 we need two preliminary lemmas.

The first is a Harnack-type alternative for sequences of solutions of PDEs. It is inspired by [4,17].

Lemma 3.2 Let $\Omega \subset \Sigma$ be a connected open subset, $\{f^n\}_{n \in \mathbb{N}}$ a bounded sequence in $L^{q}_{loc}(\Omega) \cap L^{1}(\Omega)$ for some q > 1 and $\{w^{n}\}_{n \in \mathbb{N}}$ bounded from above and solving $-\Delta w^{n} = f^{n}$ in Ω .

Then, up to subsequences, one of the following alternatives holds:

- w^n is uniformly bounded in $L^{\infty}_{loc}(\Omega)$. $w^n \xrightarrow[n \to +\infty]{} -\infty$ in $L^{\infty}_{loc}(\Omega)$.

Proof Take a compact set $\mathcal{K} \Subset \Omega$ and cover it with balls of radius $\frac{r}{2}$, with r smaller than the injectivity radius of Σ . By compactness, we can write $\mathcal{K} \subset \bigcup_{h=1}^{H} B_{\frac{r}{2}}(x_h)$. If the second alternative does not occur, then up to relabeling we get $\sup_{B_r(x_1)} w^n \ge -C$.

Then, we consider the solution z^n of

$$\begin{cases} -\Delta z^n = f^n \text{ in } B_r(x_1) \\ z^n = 0 \qquad \text{on } \partial B_r(x_1) \end{cases},$$

which is bounded in $L^{\infty}(B_r(x_1))$ by elliptic estimates. This means that, for a large constant C, the function $C - w^n + z^n$ is positive, harmonic and bounded from below on $B_r(x_1)$, and moreover its infimum is bounded from above; therefore, applying the Harnack inequality (which is allowed since r is small enough) we get that $C - w^n + z^n$ is uniformly bounded in $L^{\infty}\left(B_{\frac{r}{2}}(x_1)\right)$, hence w^n is.

At this point, by connectedness, we can relabel the index h in such a way that $B_{\frac{r}{2}}(x_h) \cap$ $B_{\frac{r}{2}}(x_{h+1}) \neq \emptyset$ for any $h \in \{1, \ldots, H-1\}$ and we repeat the argument for $B_{\frac{r}{2}}(x_2)$: since it has nonempty intersection with $B_{\frac{r}{2}}(x_1)$, we have $\sup_{B_r(x_2)} w^n \ge -C$, hence we get boundedness in $L^{\infty}\left(B_{\frac{r}{2}}(x_2)\right)$. In the same way, we obtain the same result in all the balls $B_{\frac{r}{2}}(x_h)$, whose union contains \mathcal{K} , therefore w^n must be uniformly bounded on \mathcal{K} and we get the conclusion.

The second Lemma basically says that if all the concentration values in a point are under a certain threshold, and in particular if all of them equal zero, then compactness occurs around that point.

On the other hand, if a point belongs to some set S_i , then at least a fixed amount of mass has to accumulate around it; hence, being the total mass uniformly bounded from above, this can occur only for a finite number of points, so we deduce the finiteness of the S_i 's.

Precisely, we have the following, inspired again by [17], Lemma 4.4:

Lemma 3.3 Let $\{v^n\}_{n \in \mathbb{N}}$ and S_i be as in (13) and σ_i as in (10), and suppose $\sigma_i(x) < \sigma_i^0$ for any $i \in \{1, ..., N\}$, where

$$\sigma_i^0 := \frac{4\pi \min\left\{1, 1 + \min_{j \in \{1, \dots, N\}, m \in \{1, \dots, M\}} \alpha_{jm}\right\}}{\sum_{j=1}^N a_{ij}^+}.$$

Then, $x \notin S_i$ for any $i \in \{1, \ldots, N\}$.

Proof First of all we notice that σ_i^0 is well-defined for any *i* because $a_{ii} > 0$, hence $\sum_{j=1}^{N} a_{ij}^{+} > 0.$ Under the hypotheses of the Lemma, for large *n* and small *r* we have

$$\int_{B_r(x)} \tilde{h}_i^n e^{v_i^n} \mathrm{d}V_g < \sigma_i^0.$$
⁽¹⁴⁾

Let us consider w_i^n and z_i^n defined by

$$\begin{cases} -\Delta w_i^n = -\sum_{j=1}^N a_{ij} \rho_j^n \text{ in } B_r(x) \\ w_i^n = 0 & \text{ on } \partial B_r(x) \end{cases}, \quad \begin{cases} -\Delta z_i^n = \sum_{j=1}^N a_{ij}^+ \tilde{h}_j^n e^{v_j^n} \text{ in } B_r(x) \\ z_i^n = 0 & \text{ on } \partial B_r(x) \end{cases}.$$
(15)

Is it evident that the w_i^n 's are uniformly bounded in $L^{\infty}(B_r(x))$.

As for the z_i^n 's, we can suppose to be working on a Euclidean disc, up to applying a perturbation to \tilde{h}_i^n which is smaller as r is smaller, hence for r small enough we still have the strict estimate (14).

Therefore, we get

$$\|-\Delta z_i^n\|_{L^1(B_r(x))} = \sum_{j=1}^N a_{ij}^+ \int_{B_r(x)} \tilde{h}_j^n e^{v_j^n} dV_g < \sum_{j=1}^N a_{ij}^+ \sigma_j^0 \le 4\pi \min\{1, 1+\alpha_i(x)\}$$

and we can apply Lemma 2.1 to obtain $\int_{B_r(x)} e^{q|z_i^n|} dV_g \leq C$ for some $q > \frac{1}{\min\{1, 1+\alpha_i(x)\}}$. If $\alpha_i(x) \ge 0$, then taking $q \in \left(1, \frac{4\pi}{\|-\Delta z_i^n\|_{L^1(B_r(x))}}\right)$ we have $\int_{B_r(x)} \left(\widetilde{h}_i^n e^{z_i^n}\right)^q \mathrm{d} V_g \leq C_r \int_{B_r(x)} e^{q|z_i^n|} \mathrm{d} V_g \leq C.$

On the other hand, if $\alpha_i(x) < 0$, we choose

$$q \in \left(1, \frac{4\pi}{\|-\Delta z_i^n\|_{L^1(B_r(x))} - 4\pi\alpha_i(x)}\right) \qquad q' \in \left(\frac{4\pi}{4\pi - q\|-\Delta z_i^n\|_{L^1(B_r(x))}}, \frac{1}{-\alpha_i(x)q}\right)$$

and, applying Hölder's inequality,

$$\begin{split} \int_{B_r(x)} \left(\widetilde{h}_i^n e^{z_i^n} \right)^q \mathrm{d}V_g &\leq C_r \int_{B_r(x)} d(\cdot, x)^{2q\alpha_i(x)} e^{qz_i^n} \mathrm{d}V_g \\ &\leq C \left(\int_{B_r(x)} d(\cdot, x)^{2qq'\alpha_i(x)} \mathrm{d}V_g \right)^{\frac{1}{q'}} \left(\int_{B_r(x)} e^{q\frac{q'}{q'-1}|z_i^n|} \mathrm{d}V_g \right)^{1-\frac{1}{q'}} \\ &\leq C, \end{split}$$

Springer

because $qq'\alpha_i(x) > -1$ and $q\frac{q'}{q'-1}\alpha_i(x) < \frac{4\pi}{\|-\Delta z_i^n\|_{L^1(B_r(x))}}$. Hence $\widetilde{h}_i^n e^{z_i^n}$ is uniformly bounded in $L^q(B_r(x))$ for some q > 1.

Now, let us consider $v_i^n - z_i^n - w_i^n$: it is a subharmonic sequence by construction, so for any $y \in B_{\frac{r}{2}}(x)$ we get

$$\begin{split} v_i^n(\mathbf{y}) - z_i^n(\mathbf{y}) - w_i^n(\mathbf{y}) &\leq \int_{B_{\frac{r}{2}}(\mathbf{y})} \left(v_i^n - z_i^n - w_i^n \right) \mathrm{d}V_g \\ &\leq C \int_{B_{\frac{r}{2}}(\mathbf{y})} (v_i^n - z_i^n - w_i^n)^+ \mathrm{d}V_g \\ &\leq C \int_{B_r(x)} \left((v_i^n - z_i^n)^+ + (w_i^n)^- \right) \mathrm{d}V_g \\ &\leq C \left(1 + \int_{B_r(x)} \left(v_i^n - z_i^n \right)^+ \mathrm{d}V_g \right). \end{split}$$

Moreover, since the maximum principle yields $z_i^n \ge 0$, taking $\theta = \begin{cases} 1 & \text{if } \alpha_i(x) \le 0 \\ \in \left(0, \frac{1}{1+\alpha_i(x)}\right) & \text{if } \alpha_i(x) > 0 \end{cases}$, we get

$$\begin{split} \int_{B_r(x)} \left(v_i^n - z_i^n \right)^+ \mathrm{d}V_g &\leq \int_{B_r(x)} (v_i^n)^+ \mathrm{d}V_g \\ &\leq \frac{1}{e\theta} \int_{B_r(x)} e^{\theta v_i^n} \mathrm{d}V_g \\ &\leq C \left\| \left(\widetilde{h}_i^n \right)^{-\theta} \right\|_{L^{\frac{1}{1-\theta}}(B_r(x))} \left(\int_{B_r(x)} \widetilde{h}_i^n e^{v_i^n} \mathrm{d}V_g \right)^{\theta} \\ &\leq C. \end{split}$$

Therefore, we showed that $v_i^n - z_i^n - w_i^n$ is bounded from above in $B_{\frac{r}{2}}(x)$, that is $e^{v_i^n - z_i^n - w_i^n}$ is uniformly bounded in $L^{\infty}\left(B_{\frac{r}{2}}(x)\right)$. Since the same holds for $e^{w_i^n}$ and $\tilde{h}_i^n e^{z_i^n}$ is uniformly bounded in $L^q\left(B_{\frac{r}{2}}(x)\right)$ for some q > 1, we deduce that also

$$\widetilde{h}_i^n e^{v_i^n} = \widetilde{h}_i^n e^{z_i^n} e^{v_i^n - z_i^n - w_i^n} e^{w_i^n}$$

is bounded in the same $L^q\left(B_{\frac{r}{2}}(x)\right)$.

Thus, we have an estimate on $\|-\Delta z_i^n\|_{L^q(B_{\frac{r}{2}}(x))}$ for any $i \in \{1, \ldots, N\}$, hence by standard elliptic estimates we deduce that z_i^n is uniformly bounded in $L^\infty(B_{\frac{r}{2}}(x))$. Therefore, we also deduce that

$$v_i^n = \left(v_i^n - z_i^n - w_i^n\right) + z_i^n + w_i^n$$

is bounded from above on $B_{\frac{r}{2}}(x)$, which is equivalent to saying $x \notin \bigcup_{i=1}^{N} S_i$.

From this proof, we notice that, under the assumptions of Theorem 1.3, the same result holds for any single index $i \in \{1, ..., N\}$. In other words, the upper bound on one σ_i implies that $x \notin S_i$.

Corollary 3.4 Suppose $a_{ij} \le 0$ for any $i \ne j$. Then, for any given $i \in \{1, ..., N\}$ the following conditions are equivalent:

- $x \in S_i$.
- $\sigma_i(x) \neq 0$. • $\sigma_i(x) \geq \sigma'_i = \frac{4\pi \min\{1, 1+\min_m \alpha_{im}\}}{a_{ii}}$.

Proof The third statement trivially implies the second and the second implies the first, since if v_i^n is bounded from above in $B_r(x)$ then $\tilde{h}_i^n e^{v_i^n}$ is bounded in $L^q(B_r(x))$. Finally, if $\sigma_i(x) < \sigma'_i$ then the sequence $\tilde{h}_i^n e^{z_i^n}$ defined by (15) is bounded in L^q for q > 1, so one can argue as in Lemma 3.3 to get boundedness from above of v_i^n around x, that is $x \notin S_i$.

We can now prove the main theorem of this Section.

Proof of Theorem 3.1 If $S_i = \emptyset$ for any *i*, then $e^{v_i^n}$ is bounded in $L^{\infty}(\Sigma)$, so $-\Delta v_i^n$ is bounded in $L^q(\Sigma)$ for any

$$q \in \left[1, \frac{1}{-\min_{j \in \{1, \dots, N\}, m \in \{1, \dots, M\}} \alpha_{jm}}\right)$$

Therefore, we can apply Lemma 3.2 to v_i^n on Σ , where we must have the first alternative for every *i*, since otherwise the dominated convergence would give $\int_{\Sigma} \tilde{h}_i^n e^{v_i^n} dV_g \xrightarrow[n \to +\infty]{} 0$

which is absurd; standard elliptic estimates allow to conclude compactness in $W^{2,q}(\Sigma)$.

Suppose now $S_i \neq \emptyset$ for some *i*; from Lemma 3.3 we deduce

$$|\mathcal{S}_i|\sigma_i^0 \leq \sum_{x \in \mathcal{S}_i} \max_j \sigma_j(x) \leq \sum_{j=1}^N \sum_{x \in \mathcal{S}_i} \sigma_j(x) \leq \sum_{j=1}^N \rho_j,$$

hence S_i is finite.

For any $j \in \{1, ..., N\}$, we can apply Lemma 3.2 on $\Sigma \setminus \bigcup_{j'=1}^{N} S_{j'}$ with $f^n = \sum_{j'=1}^{N} a_{jj'} \left(\tilde{h}_{j'}^{n} e^{v_{j'}^{n}} - \rho_{j'}^{n} \right)$, since the last function is bounded in $L_{\text{loc}}^{q} \left(\Sigma \setminus \bigcup_{j'=1}^{N} S_{j'} \right)$. Therefore, either v_{j}^{n} goes to $-\infty$ or it is bounded in L_{loc}^{∞} , and in the last case we get

Therefore, either v_j^n goes to $-\infty$ or it is bounded in L_{loc}^∞ , and in the last case we get compactness in $W_{loc}^{2,q}$ by applying again standard elliptic regularity.

4 Proof of Theorem 1.1

Here we will prove the theorem which gives sufficient and necessary conditions for the functional J_{ρ} to be bounded from below.

In other words, setting

$$E := \left\{ \rho \in \mathbb{R}^N_+ : J_\rho \text{ is bounded from below on } H^1(\Sigma)^N \right\},$$
(16)

we will prove that $\{\Lambda > 0\} \subset E \subset \{\Lambda \ge 0\}$.

As a first thing, we notice that the set E is not empty and it verifies a simple monotonicity condition.

Lemma 4.1 The set E defined by (16) is nonempty.

Moreover, for any $\rho \in E$ *then* $\rho' \in E$ *provided* $\rho'_i \leq \rho_i$ *for any* $i \in \{1, ..., N\}$ *.*

Proof Let $\theta > 0$ be the biggest eigenvalue of the matrix (a_{ij}) . Then,

$$J_{\rho}(u) \geq \sum_{i=1}^{N} \left(\frac{1}{2\theta} \int_{\Sigma} |\nabla u_i|^2 \mathrm{d} V_g - \rho_i \left(\log \int_{\Sigma} \widetilde{h}_i e^{u_i} \mathrm{d} V_g - \overline{u_i} \right) \right).$$

Therefore, from scalar Moser–Trudinger inequality (3), we deduce that J_{ρ} is bounded from below if $\rho_i \leq \frac{8\pi(1+\widetilde{\alpha}_i)}{\theta}$, hence $E \neq \emptyset$.

Suppose now $\rho \in E$ and $\rho'_i \leq \rho_i$ for any *i*. Then, through Jensen's inequality, we get

$$J_{\rho'}(u) = J_{\rho}(u) + \sum_{i=1}^{N} (\rho_i - \rho'_i) \log \int_{\Sigma} e^{u_i - \overline{u_i} + \log \tilde{h}_i} dV_g$$

$$\geq -C + \sum_{i=1}^{N} (\rho_i - \rho'_i) \int_{\Sigma} \log \tilde{h}_i dV_g$$

$$\geq -C$$

for any $u \in H^1(\Sigma)^N$, hence the claim.

It is interesting to observe that a similar monotonicity condition is also satisfied by the set $\{\Lambda > 0\}$ (although one can easily see that it is not true if we replace Λ with $\Lambda_{\mathcal{I},x}$).

Lemma 4.2 Let $\rho, \rho' \in \mathbb{R}^N_+$ be such that $\Lambda(\rho) > 0$ and $\rho'_i \leq \rho_i$ for any $i \in \{1, \ldots, N\}$. Then, $\Lambda(\rho') > 0$.

Proof Suppose by contradiction $\Lambda(\rho') \leq 0$, that is $\Lambda_{\mathcal{I},x}(\rho') \leq 0$ for some \mathcal{I}, x . This cannot occur for $\mathcal{I} = \{i\}$ because it would mean $\rho'_i \geq \frac{8\pi(1+\alpha_i(x))}{a_{ii}}$, so the same inequality would for ρ_i , hence $\Lambda(\rho) \leq \Lambda_{\mathcal{I},x}(\rho) \leq 0$.

Therefore, there must be some \mathcal{I} , x such that $\Lambda_{\mathcal{I},x}(\rho') \leq 0$ and $\Lambda_{\mathcal{I}\setminus\{i\},x}(\rho') > 0$ for any $i \in \mathcal{I}$; this implies

$$0 < \Lambda_{\mathcal{I} \setminus \{i\}, x}(\rho') - \Lambda_{\mathcal{I}, x}(\rho')$$

= $2 \sum_{j \in \mathcal{I}} a_{ij} \rho'_i \rho'_j - a_{ii} {\rho'_i}^2 - 8\pi (1 + \alpha_i(x)) \rho'_i$
< $\rho'_i \left(2 \sum_{j \in \mathcal{I}} a_{ij} \rho'_j - 8\pi (1 + \alpha_i(x)) \right).$ (17)

It will be not restrictive to suppose, from now on, $\rho'_1 \leq \rho_1$ and $\rho'_i = \rho_i$ for any $i \geq 2$, since the general case can be treated by exchanging the indices and iterating.

Assuming this, we must have $1 \in \mathcal{I}$, therefore we obtain:

$$\begin{aligned} 0 &< \Lambda_{\mathcal{I},x}(\rho) - \Lambda_{\mathcal{I},x}(\rho') \\ &= 8\pi (1+\alpha_1(x))(\rho_1 - \rho_1') - a_{11} \left({\rho_1'}^2 - \rho_1^2\right) - 2\sum_{j \in \mathcal{I} \setminus \{1\}} a_{1j}(\rho_1' - \rho_1)\rho_j \\ &= (\rho_1 - \rho_1') \left(8\pi (1+\alpha_1(x)) - a_{11}(\rho_1' + \rho_1) - 2\sum_{j \in \mathcal{I} \setminus \{1\}} a_{1j}\rho_j \right) \\ &< (\rho_1 - \rho_1') \left(8\pi (1+\alpha_1(x)) - 2\sum_{j \in \mathcal{I}} a_{1j}\rho_j' \right), \end{aligned}$$

which is negative by (17). We found a contradiction.

We will now show that if the parameter ρ lies in the interior of E then not only the functional is bounded from below but it is coercive in the space of zero-average functions. In particular, this fact allows to deduce the "if" part in Corollary 1.2 from Theorem 1.1.

On the other hand, if ρ belongs to the boundary of *E*, then the scenario is quite different. Lemma 4.3 Suppose $\rho \in \mathring{E}$. Then, there exists a constant $C = C_{\rho}$ such that

$$J_{\rho}(u) \geq \frac{1}{C} \sum_{i=1}^{N} \int_{\Sigma} |\nabla u_i|^2 \mathrm{d} V_g - C.$$

Moreover, J_{ρ} admits a minimizer which solves (6).

Proof Choosing $\delta \in \left(0, \frac{d(\rho, \partial E)}{\sqrt{N}|\rho|}\right)$ one has $(1 + \delta)\rho \in E$, so

$$\begin{split} J_{\rho}(u) &= \frac{\delta}{2(1+\delta)} \sum_{i,j=1}^{N} a^{ij} \int_{\Sigma} \nabla u_i \cdot \nabla u_j \mathrm{d}V_g + \frac{1}{1+\delta} J_{(1+\delta)\rho}(u) \\ &\geq \frac{\delta}{2\theta(1+\delta)} \sum_{i=1}^{N} \int_{\Sigma} |\nabla u_i|^2 \mathrm{d}V_g - C, \end{split}$$

hence we get the former claim.

To get the latter, we notice that, due to invariance by translation, any minimizer can be supposed to be in $\overline{H}^1(\Sigma)^N$; therefore, we can restrict J_ρ to this subspace. Here, the above inequality implies coercivity, and it is immediate to see that J_ρ is also lower semi-continuous, hence the existence of minimizers follows from direct methods of calculus of variations. \Box

Lemma 4.4 Suppose $\rho \in \partial E$. Then, there exists a sequence $\{u^n\}_{n \in \mathbb{N}} \subset H^1(\Sigma)^N$ such that

$$\sum_{i=1}^{N} \int_{\Sigma} \left| \nabla u_{i}^{n} \right|^{2} \mathrm{d}V_{g} \underset{n \to +\infty}{\longrightarrow} +\infty \qquad \lim_{n \to +\infty} \frac{J_{\rho}\left(u^{n}\right)}{\sum_{i=1}^{N} \int_{\Sigma} \left| \nabla u_{i}^{n} \right|^{2} \mathrm{d}V_{g}} \leq 0$$

Proof We first notice that $(1 - \delta)\rho \in E$ for any $\delta \in (0, 1)$. In fact, otherwise, from Lemma 4.1 we would get $\rho' \notin E$ as soon as $\rho'_i \ge (1 - \delta)\rho_i$ for some *i*, hence $\rho \notin \partial E$.

Now, suppose by contradiction that for any sequence u^n one gets

$$\sum_{i=1}^{N} \int_{\Sigma} \left| \nabla u_{i}^{n} \right|^{2} \mathrm{d} V_{g} \underset{n \to +\infty}{\longrightarrow} +\infty \qquad \Rightarrow \qquad \frac{J_{\rho} \left(u^{n} \right)}{\sum_{i=1}^{N} \int_{\Sigma} \left| \nabla u_{i}^{n} \right|^{2} \mathrm{d} V_{g}} \geq \varepsilon > 0.$$

Therefore, we would have

$$J_{\rho}(u) \geq \frac{\varepsilon}{2} \sum_{i=1}^{N} \int_{\Sigma} |\nabla u_i|^2 \mathrm{d} V_g - C;$$

hence, indicating as θ' the smallest eigenvalue of the matrix A, for small δ we would get

$$J_{\rho}(u) = (1+\delta)J_{(1+\delta)\rho}(u) - \frac{\delta}{2}\sum_{i,j=1}^{N} a^{ij} \int_{\Sigma} \nabla u_i \cdot \nabla u_j dV_g$$
$$\geq \left((1+\delta)\frac{\varepsilon}{2} - \frac{\delta}{2\theta'}\right)\sum_{i=1}^{N} \int_{\Sigma} |\nabla u_i|^2 - C$$
$$\geq -C.$$

So we obtain $(1 + \delta)\rho \in E$; being also $(1 - \delta)\rho \in E$ (by Lemma 4.1), we get a contradiction with $\rho \in \partial E$.

To see what happens when $\rho \in \partial E$, we build an auxiliary functional using Lemma 2.3. Lemma 4.5 Fix $\rho' \in \partial E$ and define:

$$a_{\rho'}^{n} := \frac{1}{2} \sum_{i,j=1}^{N} a^{ij} \int_{\Sigma} \nabla u_{i}^{n} \cdot \nabla u_{j}^{n} \mathrm{d}V_{g} \qquad b_{\rho'}^{n} := J_{\rho'} \left(u^{n} \right)$$
$$J_{\rho',\rho}^{\prime}(u) = J_{\rho}(u) - F_{\rho'} \left(\frac{1}{2} \sum_{i,j=1}^{N} a^{ij} \int_{\Sigma} \nabla u_{i} \cdot \nabla u_{j} \mathrm{d}V_{g} \right),$$

where u^n is given by Lemma 4.4 and $F_{\rho'}$ by Lemma 2.3.

If $\rho \in \mathring{E}$, then $J'_{\rho',\rho}$ is bounded from below on $H^1(\Sigma)^N$ and its infimum is achieved by a solution of

$$-\Delta\left(u_i-\sum_{i,j=1}^N a^{ij}fu_j\right)=\sum_{j=1}^N a_{ij}\rho_j\left(\frac{\widetilde{h}_je^{u_j}}{\int_{\Sigma}\widetilde{h}_je^{u_j}\mathrm{d}V_g}-1\right),\qquad i=1,\ldots,N,$$

with $f = (F_{\rho'})' \left(\frac{1}{2} \sum_{i,j=1}^{N} a^{ij} \int_{\Sigma} \nabla u_i \cdot \nabla u_j \mathrm{d}V_g\right).$

On the other hand, $J'_{\rho',\rho'}$ is unbounded from below.

Proof For $\rho \in \mathring{E}$, we can argue as in Lemma 4.3, since the continuity follows from the regularity of *F* and the coercivity from the behavior of *F'* at the infinity.

For $\rho = \rho'$, if we take u^n as in Lemma 4.4 we get

$$J_{\rho',\rho'}'(u^n) = b_{\rho'}^n - F_{\rho'}\left(a_{\rho'}^n\right) \underset{n \to +\infty}{\longrightarrow} -\infty.$$

Now we can prove the first half of Theorem 1.1, that is J_{ρ} is bounded from below if $\Lambda(\rho) > 0$.

Proof of $\{\Lambda > 0\} \subset E$ Suppose by contradiction there is some $\rho' \in \partial E$ with $\Lambda(\rho) > 0$ and take a sequence $\rho^n \in E$ with $\rho^n \xrightarrow[n \to +\infty]{} \rho'$.

Then, by Lemma 4.5, the auxiliary functional $J_{\rho',\rho''}$ admits a minimizer u^n , so the functions v_i^n defined as in (11) solve

$$\begin{cases} -\Delta v_i^n = \sum_{j,j'=1}^N a_{ij} b^{jj',n} \left(\tilde{h}_j e^{v_j^n} - \rho_j^n \right) \\ \int_{\Sigma} \tilde{h}_i^n e^{v_i^n} dV_g = \rho_i^n \end{cases} \quad i = 1, \dots, N \end{cases}$$

where $b^{ij,n}$ is the inverse matrix of $b_{ij}^n := \delta_{ij} - a^{ij} f^n$, hence $b^{ij,n} \xrightarrow[n \to +\infty]{} \delta_{ij}$.

We can then apply Theorem 3.1. The first alternative is excluded, since otherwise we would get, for any $u \in H^1(\Sigma)^N$,

$$J_{\rho',\rho'}'(u) = \lim_{n \to +\infty} J_{\rho',\rho^n}'(u) \ge \lim_{n \to +\infty} J_{\rho',\rho^n}'(v^n) = J_{\rho',\rho'}'(v) > -\infty,$$

thus contradicting Lemma 4.5.

Therefore, blow up must occur; this means, by Lemma 3.3, that $\sigma_i(p) \neq 0$ for some $i \in \{1, ..., N\}$ and some $p \in \Sigma$.

By Proposition 2.2 follows $\Lambda(\sigma) \leq 0$. On the other hand, since $\sigma_i \leq \rho'_i$ for any *i*, Lemma 4.2 yields $\Lambda(\rho') \leq 0$, which contradicts our assumptions.

To prove the unboundedness from below of J_{ρ} in the case $\Lambda(\rho) < 0$ we will use suitable test functions, whose properties are described by the following:

Lemma 4.6 *Define, for* $x \in \Sigma$ *and* $\lambda > 0$, $\varphi = \varphi^{\lambda, x}$ *as*

$$\varphi_i := -2(1 + \alpha_i(x)) \log \max\{1, \lambda d(\cdot, x)\}.$$

Then, as $\lambda \to +\infty$, one has

$$\begin{split} &\int_{\Sigma} \nabla \varphi_i \cdot \nabla \varphi_j \mathrm{d}V_g = 8\pi (1 + \alpha_i(x))(1 + \alpha_j(x)) \log \lambda + O(1) \\ &\overline{\varphi_i} = -2(1 + \alpha_i(x)) \log \lambda + O(1) \\ &\int_{\Sigma} \widetilde{h}_i e^{\sum_{j=1}^N \theta_j \varphi_j} \mathrm{d}V_g \geq C \lambda^{-2(1 + \alpha_i(x))} \quad if \quad \sum_{i=1}^N \theta_i(1 + \alpha_j(x)) > 1 + \alpha_i(x) \end{split}$$

Proof It holds

$$\nabla \varphi_i = \begin{cases} 0 & \text{if } d(\cdot, x) < \frac{1}{\lambda} \\ -2(1 + \alpha_i(x)) \frac{\nabla d(\cdot, x)}{d(\cdot, x)} & \text{if } d(\cdot, x) > \frac{1}{\lambda} \end{cases}$$

Therefore, being $|\nabla d(\cdot, x)| = 1$ almost everywhere on Σ :

~

$$\int_{\Sigma} \nabla \varphi_i \cdot \nabla \varphi_j dV_g$$

= 4(1 + \alpha_i(x))(1 + \alpha_j(x)) \int_{\sum \lambda_1 \frac{1}{\lambda}} \frac{dV_g}{d(\cdot x)^2}
= 8\pi (1 + \alpha_i(x))(1 + \alpha_j(x)) \log \lambda + O(1).

For the average of φ_i , we get

$$\int_{\Sigma} \varphi_i dV_g = -2(1 + \alpha_i(x)) \int_{\Sigma \setminus B_{\frac{1}{\lambda}}(x)} (\log \lambda + \log d(\cdot, x)) dV_g + O(1)$$
$$= -2(1 + \alpha_i(x)) \log \lambda + O(1).$$

For the last estimate, choose r > 0 such that $\overline{B_{\delta}(x)}$ does not contain any of the points p_m for m = 1, ..., M, except possibly x.

Then, outside such a ball, $e^{\sum_{j=1}^{N} \theta_j \varphi_j} \leq C \lambda^{-2 \sum_{j=1}^{N} \theta_j (1+\alpha_j(x))}$. Therefore, under the assumptions of the Lemma,

$$\int_{\Sigma \setminus B_{\delta}(x)} \widetilde{h}_i e^{\sum_{i=1}^N \theta_j \varphi_j} \mathrm{d} V_g = o\left(\lambda^{-2(1+\alpha_i(x))}\right),$$

hence

$$\begin{split} \int_{\Sigma} \widetilde{h}_{i} e^{\sum_{i=1}^{N} \theta_{j} \varphi_{j}} \mathrm{d}V_{g} &\geq \int_{B_{\delta}(x)} \widetilde{h}_{i} e^{\sum_{i=1}^{N} \theta_{j} \varphi_{j}} \mathrm{d}V_{g} \\ &\geq C \left(\int_{B_{\frac{1}{\lambda}}(x)} d(\cdot, x)^{2\alpha_{i}(x)} \mathrm{d}V_{g} + \frac{1}{\lambda^{2 \sum_{j=1}^{N} \theta_{j}(1+\alpha_{j}(x))}} \right) \end{split}$$

1183

`

$$\int_{A_{\frac{1}{\lambda},\delta}(x)} d(\cdot, x)^{2\alpha_i(x)-2\sum_{i=1}^N \theta_j(1+\alpha_j(x))} \mathrm{d}V_g \Biggr)$$

$$\geq C\lambda^{-2(1+\alpha_i(x))},$$

which concludes the proof.

Proof of $E \subset \{\Lambda \ge 0\}$ Take ρ, \mathcal{I}, x such that $\Lambda_{\mathcal{I},x}(\rho) < 0$ and $\Lambda_{\mathcal{I}\setminus\{i\},x}(\rho) \ge 0$ for any $i \in \mathcal{I}$, and consider the family of functions $\{u^{\lambda}\}_{\lambda>0}$ defined by

$$u_i^{\lambda} := \sum_{j \in \mathcal{I}} \frac{a_{ij} \rho_j}{4\pi (1 + \alpha_i(x))} \varphi_j^{\lambda, x}.$$

By Jensen's inequality we get

$$\begin{split} J_{\rho}\left(u^{\lambda}\right) &\leq \frac{1}{2} \sum_{i,j=1}^{N} a^{ij} \int_{\Sigma} \nabla u_{i}^{\lambda} \cdot \nabla u_{j}^{\lambda} \mathrm{d}V_{g} + \sum_{i \in \mathcal{I}} \rho_{i} \left(\overline{u_{i}^{\lambda}} - \log \int_{\Sigma} \widetilde{h}_{i} e^{u_{i}^{\lambda}} \mathrm{d}V_{g}\right) + C \\ &= \frac{1}{2} \sum_{i,j \in \mathcal{I}} \frac{a_{ij} \rho_{i} \rho_{j}}{16\pi^{2}(1+\alpha_{i}(x))(1+\alpha_{j}(x))} \int_{\Sigma} \nabla \varphi_{i} \cdot \nabla \varphi_{j} \mathrm{d}V_{g} \\ &+ \sum_{i,j \in \mathcal{I}} \frac{a_{ij} \rho_{i} \rho_{j}}{4\pi(1+\alpha_{j}(x))} \overline{\varphi_{j}} - \sum_{i \in \mathcal{I}} \rho_{i} \log \int_{\Sigma} \widetilde{h}_{i} e^{\sum_{j \in \mathcal{I}} \frac{a_{ij} \rho_{j}}{4\pi(1+\alpha_{j}(x))} \varphi_{j}} \mathrm{d}V_{g} + C. \end{split}$$

At this point, we would like to apply Lemma 4.6 to estimate $J_{\rho}(u^{\lambda})$. To be able to do this, we have to verify that

$$\frac{1}{4\pi} \sum_{j \in \mathcal{I}} a_{ij} \rho_j > 1 + \alpha_i(x) \qquad \forall i \in \mathcal{I}.$$

If $\mathcal{I} = \{i\}$, then $\rho_i > \frac{8\pi(1+\alpha_i(x))}{a_{ii}}$, so it follows immediately. For the other cases, it follows from (17).

So we can apply Lemma 4.6 and we get from the previous estimates:

$$J_{\rho}\left(u^{\lambda}\right) \leq \left(\frac{1}{4\pi}\sum_{i,j\in\mathcal{I}}a_{ij}\rho_{i}\rho_{j} - \frac{1}{2\pi}\sum_{i,j\in\mathcal{I}}a_{ij}\rho_{i}\rho_{j} + 2\sum_{i\in\mathcal{I}}\rho_{i}(1+\alpha_{i}(x))\right)\log\lambda + C$$
$$= -\frac{\Lambda_{\mathcal{I},x}(\rho)}{4\pi}\log\lambda + C\underset{n \to +\infty}{\longrightarrow} -\infty.$$

Proof of Corollary 1.2 The coercivity in the case $\Lambda < 0$, hence the existence of minimizing solutions for (6) follows from Theorem 1.1 and Lemma 4.3.

If instead $\Lambda(\rho) \ge 0$, then one can find out the lack of coercivity by arguing as before with the sequence u^{λ} , which verifies

$$\sum_{i=1}^{N} \int_{\Sigma} \left| \nabla u_{i}^{\lambda} \right|^{2} \mathrm{d} V_{g} \underset{\lambda \to +\infty}{\longrightarrow} +\infty \qquad J_{\rho} \left(u^{\lambda} \right) \leq -\frac{\Lambda_{\mathcal{I},x}(\rho)}{4\pi} \log \lambda + C \leq C.$$

Deringer

5 Proof of Theorem 1.3

Here we will finally prove a sharp inequality in the case when the matrix a_{ij} has non-positive entries outside its main diagonal.

As already pointed out in the introduction, the function $\Lambda(\rho)$ can be written in a much shorter form under these assumptions, so the condition $\Lambda(\rho) \ge 0$ is equivalent to $\rho_i \le \frac{8\pi(1+\tilde{\alpha}_i)}{\alpha_i}$ for any $i \in \{1, ..., N\}$.

Moreover, thanks to Lemma 4.1, in order to prove Theorem 1.3 for all such ρ 's it will suffice to consider

$$\rho^0 := \left(\frac{8\pi(1+\widetilde{\alpha}_1)}{a_{11}}, \dots, \frac{8\pi(1+\widetilde{\alpha}_N)}{a_{NN}}\right).$$
(18)

By what we proved in the previous Section, for any sequence $\rho^n \nearrow_{n \to +\infty} \rho^0$ one has

$$\inf_{H^1(\Sigma)^N} J_{\rho^n} = J_{\rho^n}(u^n) \ge -C_{\rho^n},$$

so Theorem 1.3 will follow by showing that, for a given sequence $\{\rho^n\}_{n\in\mathbb{N}}$, the constant $C_n = C_{\rho^n}$ can be chosen independently of n.

As a first thing, we provide a Lemma which shows the possible blow-up scenarios for such a sequence u^n .

Here, the assumption on a_{ij} is crucial since it reduces largely the possible cases.

Lemma 5.1 Let ρ^0 be as in (18), $\{\rho^n\}_{n \in \mathbb{N}}$ such that $\rho^n \nearrow \rho^0$, u^n a minimizer of J_{ρ^n} and v^n as in (11). Then, up to subsequences, there exists a set $\mathcal{I} \subset \{1, \ldots, N\}$ such that:

- If $i \in \mathcal{I}$, then $S_i = \{x_i\}$ for some $x_i \in \Sigma$ which satisfy $\widetilde{\alpha}_i = \alpha_i(x_i)$ and $\sigma_i(x_i) = \rho_i^0$, and $v_i^n \xrightarrow[n \to +\infty]{} -\infty$ in $L^{\infty}_{\text{loc}} \left(\Sigma \setminus \bigcup_{j \in \mathcal{I}} \{x_j\}\right)$.
- If $i \notin \mathcal{I}$, then $S_i = \emptyset$ and $v_i^n \xrightarrow[n \to +\infty]{} v_i$ in $W_{\text{loc}}^{2,q} \left(\Sigma \setminus \bigcup_{j \in \mathcal{I}} \{x_j\} \right)$ for some q > 1 and some suitable v_i .

Moreover, if $a_{ij} < 0$ *then* $x_i \neq x_j$.

Proof From Theorem 3.1 we get a $\mathcal{I} \subset \{1, ..., N\}$ such that $S_i \neq \emptyset$ for $i \in \mathcal{I}$. If $S_i \neq \emptyset$, then by Corollary 3.4 one gets

$$0 < \sigma_i(x) \le \rho_i^0 \le \frac{8\pi (1 + \alpha_i(x))}{a_{ii}}$$

for all $x \in S_i$, hence

$$0 = \Lambda_{\{1,\dots,N\},x}(\sigma(x))$$

$$\geq \sum_{j=1}^{N} \left(8\pi (1 + \alpha_j(x))\sigma_j(x) - a_{jj}\sigma_j(x)^2 \right)$$

$$\geq 8\pi (1 + \alpha_i(x))\sigma_i(x) - a_{ii}\sigma_i(x)^2$$

$$\geq 0.$$
(19)

Therefore, all these inequalities must actually be equalities.

From the last, we have $\sigma_i(x) = \rho_i^0 = \frac{8\pi(1+\alpha_i(x))}{a_{ii}}$, hence $\alpha_i(x) = \tilde{\alpha}_i$. On the other hand, since $\sum_{x \in S_i} \sigma_i(x) \le \rho_i^0$, it must be $\sigma_i(x) = 0$ for all but one $x_i \in S_i$, so Corollary 3.4 yields $S_i = \{x_i\}$.

Let us now show that $v_i^n \xrightarrow[n \to +\infty]{} -\infty$ in L_{loc}^∞ .

Otherwise, Theorem 3.1 would imply $v_i^n \xrightarrow[n \to +\infty]{} v_i$ almost everywhere, therefore by Fatou's Lemma we would get the following contradiction:

$$\sigma_i(x_i) < \int_{\Sigma} \widetilde{h}_i e^{v_i} \mathrm{d}V_g + \sigma_i(x_i) \le \int_{\Sigma} \widetilde{h}_i^n e^{v_i^n} \mathrm{d}V_g = \rho_i^n \le \rho_i = \sigma_i(x_i).$$

Since also inequality (19) has to be an equality, we get $a_{ij}\sigma_i(x_i)\sigma_j(x_i)$ for any $i, j \in \mathcal{I}$, so whenever $a_{ij} < 0$ there must be $\sigma_j(x_i) = 0$, so $x_i \neq x_j$.

Finally, if $S_i = \emptyset$, the convergence in $W_{loc}^{2,q}$ follows from what we just proved and Theorem 3.1.

We basically showed that if a component of the sequence v^n blows up, then all its mass concentrates at a single point which has the lowest singularity coefficient.

The next Lemma gives some more important information about the convergence or the blow-up of the components of v^n .

Lemma 5.2 Let v_i^n , v_i , ρ^0 , \mathcal{I} and x_i as in Lemma 5.1. *Then*,

• If $i \in \mathcal{I}$, then the sequence $v_i^n - \overline{v_i^n}$ converges to some G_i in $W_{\text{loc}}^{2,q}\left(\Sigma \setminus \bigcup_{j \in \mathcal{I}} \{x_j\}\right)$ for some q > 1 and weakly in $W^{1,q'}(\Sigma)$ for any $q' \in (1, 2)$, and G_i solves:

$$\begin{cases} -\Delta G_i = \sum_{j \in \mathcal{I}} a_{ij} \rho_j^0 \left(\delta_{x_j} - 1 \right) + \sum_{j \notin \mathcal{I}} a_{ij} \left(\widetilde{h}_j e^{v_j} - \rho_j^0 \right) \\ \overline{G_i} = 0 \end{cases}$$

• If $i \notin \mathcal{I}$, then $v_i^n \xrightarrow[n \to +\infty]{} v_i$ in the same space, and v_i solves:

$$\begin{cases} -\Delta v_i = \sum_{j \in \mathcal{I}} a_{ij} \rho_j^0 \left(\delta_{x_j} - 1 \right) + \sum_{j \notin \mathcal{I}} a_{ij} \left(\widetilde{h}_j e^{v_j} - \rho_j^0 \right) \\ \int_{\Sigma} \widetilde{h}_i e^{v_i} dV_g = \rho_i^0 \end{cases}$$
(20)

Proof From Lemma 5.1 follows that, for $i \in \mathcal{I}$, $\tilde{h}_i^n e^{v_i^n} \xrightarrow[n \to \infty]{} \rho_i^0 \delta_{x_i}$ in the sense of measures; in fact, for any $\phi \in C(\Sigma)$

$$\begin{split} \left| \int_{\Sigma} \tilde{h}_{i}^{n} e^{v_{i}^{n}} \phi \mathrm{d}V_{g} - \rho_{i}^{0} \phi(x_{i}) \right| &\leq \int_{\Sigma} \tilde{h}_{i}^{n} e^{v_{i}^{n}} |\phi - \phi(x_{i})| \mathrm{d}V_{g} + \left|\rho_{i}^{n} - \rho_{i}^{0}\right| |\phi(x_{i})| \\ &\leq \varepsilon \int_{B_{\delta}(x_{i})} \tilde{h}_{i}^{n} e^{v_{i}^{n}} \mathrm{d}V_{g} + 2 \|\phi\|_{L^{\infty}(\Sigma)} \int_{\Sigma \setminus B_{\delta}(x_{i})} \tilde{h}_{i}^{n} e^{v_{i}^{n}} \mathrm{d}V_{g} \\ &+ \left|\rho_{i}^{n} - \rho_{i}^{0}\right| \|\phi\|_{L^{\infty}(\Sigma)} \\ &\leq \varepsilon \rho_{i}^{n} + 2 \|\phi\|_{L^{\infty}(\Sigma)} o(1) + o(1) \|\phi\|_{L^{\infty}(\Sigma)}, \end{split}$$

which is, choosing properly ε , arbitrarily small. Therefore, v_i solves (20).

On the other hand, if $q' \in (1, 2)$, then $\frac{q'}{q'-1} > 2$, so any function $\phi \in W^{1, \frac{q'}{q'-1}}(\Sigma)$ is actually continuous, hence

$$\left| \int_{\Sigma} \nabla \left(v_i^n - \overline{v_i^n} - G_i \right) \cdot \nabla \phi \mathrm{d} V_g \right|$$
$$= \left| \int_{\Sigma} \left(-\Delta v_i^n + \Delta G_i \right) \phi \mathrm{d} V_h \right|$$

🖄 Springer

$$\leq \sum_{j \in \mathcal{I}} a_{ij} \left| \int_{\Sigma} \widetilde{h}_j e^{v_j^n} \phi dV_g - \rho_j^0 \phi(p) \right| \\ + \sum_{j \notin \mathcal{I}} a_{ij} \left| \int_{\Sigma} \widetilde{h}_j \left(e^{v_j^n} - e^{v_j} \right) \phi dV_g \right| \underset{n \to +\infty}{\longrightarrow} 0.$$

Therefore, we get weak convergence in $W^{1,q'}(\Sigma)$ for any $q' \in (1,2)$; standard elliptic estimates yield convergence in $W^{2,q}_{\text{loc}}(\Sigma \setminus \bigcup_{j \in \mathcal{I}} \{x_j\})$.

In the same way we prove the same convergence of v_i^n to v_i .

From these information about the blow-up profile of v^n we deduce an important fact which will be used to prove the main Theorem:

Corollary 5.3 Let v^n and x_i be as in Lemmas 5.1 and 5.2 and w^n be defined by $w_i^n = \sum_{i=1}^N a^{ij} v_i^n$ for $i \in \{1, ..., N\}$.

Then, $w_i^n - \overline{w_i^n}$ is uniformly bounded in $W_{\text{loc}}^{2,q}(\Sigma \setminus \{x_i\})$ for some q > 1 if $i \in \mathcal{I}$, whereas if $i \notin \mathcal{I}$ it is bounded in $W^{2,q}(\Sigma)$.

Proof Since $-\Delta w_i^n = \tilde{h}_i^n e^{v_i^n} - \rho_i^n$, the claim follows from the boundedness of $e^{v_i^n}$ in $L_{\text{loc}}^{\infty}(\Sigma \setminus \{x_i\})$ and from standard elliptic estimates.

The last Lemma we need is a localized scalar Moser–Trudinger inequality for the blowingup sequence.

Lemma 5.4 Let w_i^n be as in Corollary 5.3 and x_i as in the previous Lemmas. Then, for any $i \in \mathcal{I}$ and any small r > 0 one has

$$\frac{a_{ii}}{2} \int_{B_r(x_i)} \left| \nabla w_i^n \right|^2 \mathrm{d}V_g - \rho_i^n \left(\log \int_{B_r(x_i)} \widetilde{h}_i e^{a_{ii} w_i^n} \mathrm{d}V_g - a_{ii} \overline{w_i^n} \right) \ge -C_r.$$

Proof Since Σ is locally conformally flat, we can choose r small enough so that we can apply Corollary 2.5 up to modifying \tilde{h}_i^n . We also take r so small that $\overline{B_r(x_i)}$ contains neither any x_j for $x_j \neq x_i$ nor any p_m for m = 1, ..., M (except possibly x_i).

Let z^n be the solution of

$$\begin{cases} -\Delta z_i^n = \widetilde{h}_i^n e^{v_i^n} - \rho_i^n \text{ in } B_r(x_i) \\ z_i^n = 0 \qquad \text{ on } \partial B_r(x_i) \end{cases}$$

Then, $w_i^n - \overline{w_i^n} - z_i^n$ is harmonic and it has the same value as $w_i^n - \overline{w_i^n}$ on $\partial B_r(x_i)$, so from standard estimates

$$\left\|w_i^n - \overline{w_i^n} - z_i^n\right\|_{C^1(B_r(x_i))} \le C \left\|w_i^n - \overline{w_i^n}\right\|_{C^1(\partial B_r(x_i))} \le C.$$

From Lemma 5.2 we get

$$\begin{split} \left| \int_{B_r(x_i)} \left| \nabla w_i^n \right|^2 \mathrm{d}V_g - \int_{B_r(x_i)} \left| \nabla z_i^n \right|^2 \mathrm{d}V_g \right| &= \left| \int_{B_r(x_i)} \left| \nabla \left(w_i^n - z_i^n \right) \right|^2 \mathrm{d}V_g \\ &+ 2 \int_{B_r(x_i)} \nabla w_i^n \cdot \nabla \left(w_i^n - z_i^n \right) \mathrm{d}V_g \right| \\ &\leq \int_{B_r(x_i)} \left| \nabla \left(w_i^n - z_i^n \right) \right|^2 \mathrm{d}V_g \end{split}$$

Deringer

$$+2 \left\|\nabla w_i^n\right\|_{L^1(\Sigma)} \left\|\nabla \left(w_i^n - z_i^n\right)\right\|_{L^\infty(B_r(x_i))} \le C_r.$$

Moreover,

$$\begin{split} \int_{B_r(x_i)} \widetilde{h}_i e^{a_{ii} \left(w_i^n - \overline{w_i^n} \right)} \mathrm{d}V_g &\leq e^{a_{ii} \left\| w_i^n - \overline{w_i^n} - z_i^n \right\|_{L^{\infty}(B_r(x_i))}} \int_{B_r(x_i)} \widetilde{h}_i e^{a_{ii} z_i^n} \mathrm{d}V_g \\ &\leq C_r \int_{B_r(x_i)} d(\cdot, x_i)^{2\widetilde{\alpha}_i} e^{a_{ii} z_i^n} \mathrm{d}V_g. \end{split}$$

Therefore, since $\tilde{\alpha}_i \leq 0$ and $a_{ii}\rho_i^n \leq 8\pi(1+\tilde{\alpha}_i)$, we can apply Corollary 2.5 to get the claim:

$$\frac{a_{ii}}{2} \int_{B_r(x_i)} |\nabla w_i^n|^2 \, \mathrm{d}V_g - \rho_i^n \log \int_{B_r(x_i)} \widetilde{h}_i e^{a_{ii} \left(w_i^n - \overline{w_i^n}\right)} \, \mathrm{d}V_g$$

$$\geq \frac{1}{2a_{ii}} \int_{B_r(x_i)} |\nabla \left(a_{ii} z_i^n\right)|^2 \, \mathrm{d}V_g$$

$$-\rho_i^n \log \int_{B_r(x_i)} d(\cdot, x_i)^{2\widetilde{\alpha}_i} e^{a_{ii} z_i^n} \, \mathrm{d}V_g - C_r$$

$$\geq -C_r$$

Proof of Theorem 1.3 As noticed before, it suffices to prove the boundedness from below of $J_{\rho^n}(u^n)$ for a sequence $\rho^n \nearrow_{n \to +\infty} \rho^0$ and a sequence of minimizers u^n for J_{ρ^n} . Moreover, due to invariance by addition of constants, one can consider v^n in place of u^n .

Let us start by estimating the term involving the gradients.

From Corollary 5.3 we deduce that the integral of $|\nabla w_i^n|^2$ outside a neighborhood of x_i is uniformly bounded for any $i \in \mathcal{I}$, and the integral on the whole Σ is bounded if $i \notin \mathcal{I}$.

For the same reason, the integral of $a_{ij} \nabla w_i^n \cdot \nabla w_i^n$ on the whole surface is uniformly bounded. In fact, if $a_{ij} \neq 0$, then $x_i \neq x_j$, then

$$\begin{split} \left| \int_{\Sigma} \nabla w_i^n \cdot \nabla w_j^n \mathrm{d} V_g \right| &\leq \int_{\Sigma \setminus B_r(x_j)} \left| \nabla w_i^n \cdot \nabla w_j^n \right| \mathrm{d} V_g + \int_{\Sigma \setminus B_r(x_i)} \left| \nabla w_i^n \cdot \nabla w_j^n \right| \mathrm{d} V_g \\ &\leq \left\| \nabla w_i^n \right\|_{L^{q'}(\Sigma)} \left\| \nabla w_j^n \right\|_{L^{q''}(\Sigma \setminus B_r\{x_i\})} \\ &+ \left\| \nabla w_i^n \right\|_{L^{q''}(\Sigma \setminus B_r\{x_i\})} \left\| \nabla w_j^n \right\|_{L^{q'}(\Sigma)} \\ &\leq C_r, \end{split}$$

with q as in Corollary 5.3, $q' = \begin{cases} \frac{2q}{3q-2} < 2 \text{ if } q < 2\\ 1 & \text{if } q \ge 2 \end{cases}$ and $q'' = \begin{cases} \frac{2q}{2-q} & \text{if } q < 2\\ \infty & \text{if } q \ge 2 \end{cases}$.

Therefore, we can write

$$\sum_{i,j=1}^{N} a^{ij} \int_{\Sigma} \nabla v_i^n \cdot \nabla v_j^n \mathrm{d}V_g = \sum_{i,j=1}^{N} a_{ij} \int_{\Sigma} \nabla w_i^n \cdot \nabla w_j^n \mathrm{d}V_g$$
$$\geq \sum_{i \in \mathcal{I}} a_{ii} \int_{B_r(x_i)} \left| \nabla w_i^n \right|^2 \mathrm{d}V_g - C_r.$$

Springer

To deal with the other term in the functional, we use the boundedness of w_i^n away from x_i : choosing r as in Lemma 5.4, we get

$$\begin{split} \int_{\Sigma} \widetilde{h}_{i}^{n} e^{v_{i}^{n} - \overline{v_{i}^{n}}} \mathrm{d}V_{g} &\leq 2 \int_{B_{r}(x_{i})} \widetilde{h}_{i}^{n} e^{v_{i}^{n} - \overline{v_{i}^{n}}} \mathrm{d}V_{g} \\ &= 2 \int_{B_{r}(x_{i})} \widetilde{h}_{i} e^{\sum_{j=1}^{N} a_{ij} \left(w_{j}^{n} - \overline{w_{j}^{n}}\right)} \mathrm{d}V_{g} \\ &\leq C_{r} \int_{B_{r}(x_{i})} \widetilde{h}_{i} e^{a_{ij} \left(w_{i}^{n} - \overline{w_{i}^{n}}\right)} \mathrm{d}V_{g}. \end{split}$$

Therefore, using Lemma 5.4 we obtain

$$J_{\rho^{n}}(v^{n}) = \frac{1}{2} \sum_{i,j=1}^{N} a^{ij} \int_{\Sigma} \nabla v_{i}^{n} \cdot \nabla v_{j}^{n} dV_{g} - \sum_{i=1}^{N} \rho_{i}^{n} \left(\log \int_{\Sigma} \widetilde{h}_{i}^{n} e^{v_{i}^{n}} dV_{g} - \overline{v_{i}^{n}} \right)$$

$$\geq \sum_{i \in \mathcal{I}} \left(\frac{a_{ii}}{2} \int_{B_{r}(x_{i})} \left| \nabla w_{i}^{n} \right|^{2} dV_{g} - \rho_{i}^{n} \left(\log \int_{B_{r}(x_{i})} \widetilde{h}_{i} e^{a_{ii} w_{i}^{n}} dV_{g} - a_{ii} \overline{w_{i}^{n}} \right) \right) - C_{r}$$

$$\geq -C_{r}$$

Since the choice of r does not depend on n, the proof is complete.

Acknowledgments The author would like to express his gratitude Professor Andrea Malchiodi for the support and for the discussions about this topic.

References

- Adimurthi, Sandeep, K.: A singular Moser–Trudinger embedding and its applications. Nonlinear Differ. Equ. Appl. (NoDEA) 13(5–6), 585–603 (2007)
- Battaglia, L., Malchiodi, A.: A Moser–Trudinger inequality for the singular Toda system. Bull. Inst. Math. Acad. Sin. (N.S.) 9(1), 1–23 (2014)
- Bolton, J., Woodward, L.M.: Some geometrical aspects of the 2-dimensional Toda equations. In: Geometry, topology and physics (Campinas, 1996), pp. 69–81. de Gruyter, Berlin (1997)
- 4. Brezis, H., Merle, F.: Uniform estimates and blow-up behavior for solutions of $-\Delta u = V(x)e^{u}$ in two dimensions. Commun. Partial Differ. Equ. **16**(8–9), 1223–1253 (1991)
- 5. Calabi, E.: Isometric imbedding of complex manifolds. Ann. Math. 2(58), 1-23 (1953)
- 6. Chang, S.-Y.A., Yang, P.C.: Prescribing Gaussian curvature on S². Acta Math. **159**(3–4), 215–259 (1987)
- Chang, S.-Y.A., Yang, P.C.: Conformal deformation of metrics on S². J. Differ. Geom. 27(2), 259–296 (1988)
- Chen, W.X.: A Trüdinger inequality on surfaces with conical singularities. Proc. Am. Math. Soc. 108(3), 821–832 (1990)
- Chern, S.S., Wolfson, J.G.: Harmonic maps of the two-sphere into a complex Grassmann manifold. II. Ann. Math. 125(2), 301–335 (1987)
- Chipot, M., Shafrir, I., Wolansky, G.: On the solutions of Liouville systems. J. Differ. Equ. 140(1), 59–105 (1997)
- Chipot, M., Shafrir, I., Wolansky, G.: Erratum: "On the solutions of Liouville systems" [J. Differential Equations 140 (1997), no. 1, 59–105; MR1473855 (98j:35053)]. J. Differ. Equ. 178(2), 630 (2002)
- 12. Dunne, G.: Self-dual Chern–Simons Theories. Lecture Notes in Physics. New series m: Monographs. Springer, (1995)
- Fontana, L.: Sharp borderline Sobolev inequalities on compact Riemannian manifolds. Comment. Math. Helv. 68(3), 415–454 (1993)
- Jost, J., Lin, C., Wang, G.: Analytic aspects of the Toda system. II. Bubbling behavior and existence of solutions. Commun. Pure Appl. Math. 59(4), 526–558 (2006)
- Jost, J., Wang, G.: Analytic aspects of the Toda system—I: a Moser–Trudinger inequality. Commun. Pure Appl. Math. 54(11), 1289–1319 (2001)

П

- Lin, C.-S., Wei, J-c, Zhang, L.: Classification of blowup limits for SU(3) singular Toda systems. Anal. PDE 8(4), 807–837 (2015)
- Lucia, M., Nolasco, M.: SU(3) Chern–Simons vortex theory and Toda systems. J. Differ. Equ. 184(2), 443–474 (2002)
- Moser, J.: A sharp form of an inequality by N. Trudinger. Indiana Univ. Math. J. 20, 1077–1092 (1970/1971)
- Shafrir, I., Wolansky, G.: The logarithmic HLS inequality for systems on compact manifolds. J. Funct. Anal. 227(1), 200–226 (2005)
- Shafrir, I., Wolansky, G.: Moser–Trudinger and logarithmic HLS inequalities for systems. J. Eur. Math. Soc. (JEMS) 7(4), 413–448 (2005)
- Tarantello, G.: Selfdual gauge field vortices: an analytical approach. In: Progress in Nonlinear Differential Equations and their Applications, vol. 72. Birkhäuser Boston Inc., Boston, MA (2008)
- Troyanov, M.: Prescribing curvature on compact surfaces with conical singularities. Trans. Am. Math. Soc. 324(2), 793–821 (1991)
- Wang, G.: Moser–Trudinger inequalities and Liouville systems. C. R. Acad. Sci. Paris Sér. I Math. 328(10), 895–900 (1999)
- Yang, Y.: Solitons in Field Theory and Nonlinear Analysis. Springer Monographs in Mathematics. Springer, New York (2001)