Mathematische Annalen

CrossMark

On the boundary behavior of Kähler–Einstein metrics on log canonical pairs

Henri Guenancia[1] · **Damin Wu**[2]

Abstract In this paper, we study the boundary behavior of the negatively curved Kähler–Einstein metric attached to a log canonical pair (X, D) such that $K_X + D$ is ample. In the case where X is smooth and D has simple normal crossings support (but possibly negative coefficients), we provide a very precise estimate on the potential of the KE metric near the boundary D. In the more general singular case (D being assumed effective though), we show that the KE metric has mixed cone and cusp singularities near D on the snc locus of the pair. As a corollary, we derive the behavior in codimension one of the KE metric of a stable variety.

1 Introduction

This paper studies negatively curved Kähler–Einstein metrics on quasi-projective manifolds. This is of course a very broad topic which has witnessed a lot of developments since the foundational works of Aubin, Yau [1,34] in the compact case. Quickly after the resolution of Calabi's conjecture by Yau, many works have revolved around the (non-compact) complete case; let us mention Yau [33], Cheng–Yau [8], Mok–Yau

✉ Henri Guenancia
    henri.guenancia@stonybrook.edu
    http://www.math.sunysb.edu/~guenancia

Damin Wu
    damin.wu@uconn.edu

1 Department of Mathematics, Stony Brook University, Stony Brook, NY 11794-3651, USA

2 Department of Mathematics, University of Connecticut, 196 Auditorium Road, Storrs, CT 06269-3009, USA

[24], Kobayashi [20] and Tian–Yau [26] in the negative scalar curvature case, and [27] in the Ricci-flat case to cite only a few of them.

More recently, a lot of attention has been drawn to conical Kähler–Einstein metrics, which are non-complete metrics living on the complement of a (smooth) divisor in a compact manifold, having a very precise behavior near the divisor, cf. [4,5,9,10,12, 18,19,23,30].

These classes of examples (in the negatively curved case) can be recast in a unified framework. Namely, when we are seeking negatively curved Kähler–Einstein metrics on the complement $X \setminus D$ of a smooth divisor D (or merely with simple normal crossings) in a compact Kähler manifold X, then one has at some point to do an assumption on the positivity of the adjoint canonical bundle $K_X + D$. More precisely, the existence of a negatively curved KE metric with cuspidal singularities along D such as in [20,26] is equivalent to the ampleness of $K_X + D$. In the same vein, the existence of a negatively curved conical KE metric with cone angle $2\pi\beta$ along D is equivalent to the ampleness of $K_X + (1 - \beta)D$.

So what if now, we look at the problem from another angle? That is, what if instead of looking for Kähler–Einstein metrics on $X \setminus D$ having a prescribed behavior along D, we just start by assuming that the line bundle $K_X + aD$ is ample for some real number a, and see what kind of Kähler–Einstein metrics one can construct? Well, if $a \in (0, 1]$, we end up with conical/cuspidal metrics because of what we explained above. In this paper, we will leave aside that case where $a > 1$, and only study the situation where $a \in (-\infty, 1]$. More generally, take $D = \sum D_i$ be a simple normal crossings (snc) divisor, choose real numbers $a_i \in (-\infty, 1]$, and assume that $K_X + \sum a_i D_i$ is ample. Is it possible to construct "reasonable" Kähler–Einstein metrics with negative scalar curvature on $X \setminus D$ that are naturally related to the data of the a_i's?

This question has been studied from various points of view [3,13,28,29] and it seems that the framework of pluripotential theory could be the best fit as it yields a unified approach and treatment of the problems at stake. Indeed, it has been proved in [3] that given a pair (X, D) where $D = \sum a_i D_i$ is a divisor with simple normal crossings support and coefficients $a_i \in (-\infty, 1]$ such that $K_X + D$ is ample, there exists a unique "weak" Kähler–Einstein metric ω_{KE}, smooth on $X \setminus D$ and satisfying Ric $\omega_{KE} = -\omega_{KE} + \sum a_i [D_i]$ in the sense of currents. Moreover, the singularities of ω_{KE} near D are relatively mild as this current has finite energy, cf. [16]. What more do we know about ω_{KE}?

Well, first, if the pair is klt, i.e. $a_i < 1$ for all i (as the a_i's may be negative, this condition is sometimes called "sub-klt" in the literature), then if follows from Kołodziej's estimate [21] that the metric has *bounded* potentials. But as soon as some coefficient a_i equals 1, the potentials have to be unbounded. This can be seen using the Monge–Ampère formulation of the Kähler–Einstein problem which takes the form

$$(\omega + dd^c\varphi)^n = \frac{e^\varphi \omega^n}{\prod_i |s_i|^{2a_i}}$$

where $\omega \in c_1(K_X + D)$ is a Kähler form, and s_i is a defining section for D_i, whose associated line bundle we endow with a suitable hermitian metric (to get the condition on the Ricci curvature). Then, as $|s|^{-2}$ is not integrable, φ has to go to $-\infty$ near

$D_{lc} := \sum_{a_i=1} D_i$ to guarantee the integrability of the rhs. So one cannot expect bounded potentials. If now the divisors has only coefficients equal to 1, then we know from [20,26] that $\varphi = -\sum_i \log(\log |s_i|^2)^2 + O(1)$, and that ω_{KE} has Poincaré singularities along D. We have an analogous expansion (i.e., loglog near D_{lc} + bounded term) if the coefficients are orbifold [26], or more generally of $a_i \in [0, 1]$ [12,13]. More generally, if the irreducible components of D associated to coefficients $a_i < 1$ do not meet any irreducible components of D associated to coefficients $a_i = 1$, the same result holds [29]. The first Theorem of this note aims to prove that the above expansion for the potential always holds regardless of the combinatoric of D:

Theorem A *Let X be a compact Kähler manifold, $D = \sum a_i D_i$ be a divisor with simple normal crossings support having coefficients $a_i \in (-\infty, 1]$ and such that $K_X + D$ is ample. Let $\omega \in c_1(K_X + D)$ a Kähler form, and let $\omega_{KE} = \omega + dd^c \varphi_{KE}$ be the Kähler–Einstein metric of (X, D), i.e., $\mathrm{Ric}\, \omega_{KE} = -\omega_{KE} + [D]$. Then*

$$\varphi_{KE} = -\sum_{a_i=1} \log \log^2 \frac{1}{|s_i|^2} + O(1)$$

We will give two proofs of this result, both based on a approach involving Green's functions but in different contexts. The two proofs share a common core: we start by partially regularizing the Monge–Ampère equation so as to make it of Poincaré-type as in [20,26], and then one will seek for uniform estimates on the potential, independent of the regularizing parameter. The lower bound is obtained using ideas involving approximate cone metrics, and already appearing in [5,13]. Then, using Yau's maximum principle for complete manifolds, we derive an upper bound of the potential involving $\sum_{a_i<0} a_i \log |s_i|^2$. Of course, the right hand side goes to $+\infty$ near the boundary divisor, so this estimate is not sufficient to prove Theorem A. This is where our two proofs take different paths.

The common idea is to estimate the supremum of the potential by its L^1 norm (which is controlled by the previous estimate) using Green's function. The difficulty here is that there is no global positive Green's function for the Laplacian Δ on $X \backslash D_{lc}$ as follows from [7] since a Poincaré-type metric has finite volume. On the other hand, one cannot use the local Green's function of Δ as in [29], since the injectivity radius of the Poincaré metric shrinks to zero as the point tends to D_{lc}. In the first approach, we pull-back the equation to some kind of universal cover to make the Poincaré metric into an euclidian one, so that one can use a standard local Green's function upstairs, derive an upper bound upstairs, and then push it back down to M.

In the second approach, we first construct a global Green's function associated with $\Delta_g - 1$ on any complete Riemannian manifold (M, g). For $M = X \backslash D_{lc}$ endowed with a Poincaré-type metric, we can control the asymptotic behavior of this function with sufficient precision so as to get an upper bound for our potential.

To go beyond Theorem A, it would be natural to expect higher order estimates on the potential of the Kähler–Einstein metric. As for Laplacian type estimates, this has already been done in [12,13] whenever the coefficients of D are non-negative. It would be challenging to extend these results to our more general setting, but if $D = -aH$ for some positive number a, H being a smooth hypersurface, then one of the main

new issues is that we do not really have a global reference metric on $X \backslash H$ that would behave like $|z_1|^{2a} dz_1 \wedge d\bar{z}_1 + \sum_{j>1} dz_j \wedge d\bar{z}_j$ near H, whenever it is locally given by $(z_1 = 0)$.

In the second part, we investigate the case of singular pairs (X, D), i.e., X is now a normal projective variety and D an *effective* Weil divisor on X such that the pair (X, D) has log canonical singularities. If one assumes that $K_X + D$ is ample, then we know from [3] that (X, D) admits a unique Kähler–Einstein metric ω_{KE} (see Sect. 4 for the related definitions), which is smooth on $X_{\text{reg}} \backslash \text{Supp}(D)$. We will study the behavior of ω_{KE} near D, and more precisely at the points where X is smooth and D has simple normal crossings support:

Theorem B *Let (X, D) be a projective log canonical pair such that $K_X + D$ is ample. Then its Kähler–Einstein metric ω_{KE} has mixed cone and cusp singularities along D on the snc locus $(X, D)_{\text{reg}}$ of the pair.*

As a corollary of this theorem, we show that the Kähler–Einstein metric of a stable variety (in the sense of Kollár–Sherpherd–Barron and Alexeev) is cuspidal near the double crossing points, cf. Corollary 4.4.

Let us conclude this introduction by saying that both Theorem A and Theorem B are the crucial analytic inputs in the proof of the polystability of the logarithmic tangent sheaf of a log canonical pair (X, D) such that $K_X + D$ is ample, cf. [15].

2 The smooth case

2.1 The set-up

The setting in this paper is the following one: X is a smooth complex projective variety of dimension n, $D = \sum a_i D_i$ is a $\mathbb{R}$-divisor with simple normal crossing support with coefficients $a_i \in (-\infty, 1]$ such that the adjoint bundle $K_X + D$ is ample (i.e., its Chern class contains a Kähler metric, or equivalently $K_X + D$ is $\mathbb{Q}$-linearly equivalent to a positive $\mathbb{R}$-linear combination of ample $\mathbb{Q}$-line bundles). We stress here that the coefficients of D may be chosen to be negative.

We set $D_{lc} := \sum_{a_i=1} D_i$, and $D_{klt} := D - D_{lc}$, and $M := X \backslash \text{Supp}(D_{lc})$. This notations are borrowed from birational geometry, in the sense that (X, D_{klt}) (resp. (X, D) or (X, D_{lc})) is a –log smooth– Kawamata log terminal (klt) pair (resp. log canonical (lc) pair).

As for endowing $X \backslash \text{Supp}(D_{lc})$ with a natural Kähler–Einstein metric, the viewpoints and definitions vary according to the authors, and we will choose here the following definition which has the advantage to be globally formulated on X, and guarantees the uniqueness of the metric thanks to the formalism developed in [16] and its companion papers:

Definition 2.1 With the previous notations, we say that a closed positive current $\omega \in c_1(K_X + D)$ on X is a Kähler–Einstein metric for (X, D) if it satisfies:

1. The non-pluripolar product ω^n defines an absolutely continuous measure with respect to some smooth volume form dV on X and $\log(\omega^n/dV) \in L^1_{\text{loc}}(X)$,

2. Ric $\omega = -\omega + [D]$,
3. $\int_X \omega^n = c_1(K_X + D)^n$.

This seemingly complicated definition comes from the fact that we know that ω cannot have bounded potentials, hence we have to use the non-pluripolar Monge–Ampère operator [2,16] in order to define ω^n (and thus Ric ω which is defined as $-dd^c \log\langle\omega^n\rangle$ as soon as 1. is satisfied) and have a suitable formulation of the problem in terms of Monge–Ampère equations.

Recall that Bedford–Taylor theory enables to define the Monge–Ampère operator of a closed positive current with bounded potentials and that the non-pluripolar Monge–Ampère operator is defined for any closed positive current T (say in a Kähler class on a compact Kähler manifold) and produces a positive measure denoted by $\langle T^n \rangle$ (or abusively T^n) which does not charge pluripolar sets and whose total mass is less than or equal to the top auto-intersection $\{T\}^n$ of its de Rham cohomology class. If T has bounded potentials, $\langle T^n \rangle$ coincides with the Bedford–Taylor Monge–Ampère T^n and the total mass of T^n is indeed $\{T\}^n$. More generally the equality of those two quantities is an information on the regularity of the current, cf. [2,16] for the definitions and main properties of this construction.

We know from [3] that there exists a unique such current ω; moreover, ω defines a smooth Kähler–Einstein metric on M, and if $\theta_\varepsilon \in c_1(D_{klt})$ is any smooth approximation of $[D_{klt}]$, then ω is the (weak) limit of the twisted Kähler–Einstein metrics ω_ε satisfying Ric $\omega_\varepsilon = -\omega_\varepsilon + \theta_\varepsilon + [D_{lc}]$.

In this part, we assume that the pair (X, D) is log smooth, and we prove Theorem A.

2.2 The lower bound

As we explained in the previous part, it follows from the results of [3] that it is sufficient to obtain uniform estimates for the potential φ_ε solution of

$$(\omega + dd^c \varphi_\varepsilon)^n = \prod_{a_i < 1}(|s_i|^2 + \varepsilon^2)^{-a_i} \frac{e^{\varphi_\varepsilon} \omega^n}{\prod_{a_k=1} |s_k|^2}$$

At that point, it is convenient to work with the complete Poincaré metric $\omega_P := \omega - \sum_{a_k=1} dd^c \log\log^2 |s_k|^2$ on M (up to scaling the hermitian metrics on $\mathcal{O}(D_k)$, it defines indeed a smooth complete Kähler metric with bounded geometry on M); so we set $u_\varepsilon := \varphi_\varepsilon + \sum_{a_k=1} \log\log^2 |s_k|^2$, so that the equation becomes (on M)

$$(\omega_P + dd^c u_\varepsilon)^n = \prod_{a_i < 1}(|s_i|^2 + \varepsilon^2)^{-a_i} e^{u_\varepsilon + F} \omega_P^n \qquad (2.1)$$

where F is known to be a bounded smooth function on M (which is even smooth in the quasi-coordinates, cf. [20,26]).

The first step is to introduce the regularized cone metric [5,6,13]. To sum up the construction therein, there exists a smooth ω-psh (and ω_P-psh) potential ψ_ε which is

uniformly bounded, and such that the metric $\omega_{P,\varepsilon} := \omega_P + dd^c \psi_\varepsilon$ on M is complete, with bounded bisectional curvature and satisfies

$$\prod_{0 < a_i < 1} (|s_i|^2 + \varepsilon^2)^{-a_i} \omega_P^n = e^{G_\varepsilon} \omega_{P,\varepsilon}^n$$

for some smooth function G_ε which is uniformly bounded in ε. Therefore, setting $v_\varepsilon := u_\varepsilon - \psi_\varepsilon$, Eq. (2.1) becomes

$$(\omega_{P,\varepsilon} + dd^c v_\varepsilon)^n = \prod_{a_i < 0} (|s_i|^2 + \varepsilon^2)^{-a_i} e^{v_\varepsilon + F_\varepsilon} \omega_{P,\varepsilon}^n \qquad (2.2)$$

where $F_\varepsilon = F + G_\varepsilon + \psi_\varepsilon$, and by the remarks above, $|F_\varepsilon| \leqslant C$ for some *uniform* $C > 0$.

If we apply Yau's minimum principle [32] on the manifold $(M, \omega_{P,\varepsilon})$, we get that $\inf v_\varepsilon \geqslant -\sup F_\varepsilon + \inf \sum_{a_i < 0} a_i \log(|s_i|^2 + \varepsilon^2)$ and therefore

$$\inf_M u_\varepsilon \geqslant -C \qquad (2.3)$$

for some uniform $C > 0$.

2.3 The upper bound I

Let us get now to the upper bound. We cannot apply the same method here as one sees immediately, so we perform a change of function by setting $w_\varepsilon := v_\varepsilon - \sum_{a_i < 0} a_i \log(|s_i|^2 + \varepsilon^2)$. As $\log(|s_i|^2 + \varepsilon^2)$ is $C\omega$-psh for some uniform C, it is also $C\omega_{P,\varepsilon}$-psh (up to changing C eventually), and therefore

$$
\begin{aligned}
e^{w_\varepsilon + F_\varepsilon} \omega_{P,\varepsilon}^n &= \prod_{a_i < 0} (|s_i|^2 + \varepsilon^2)^{-a_i} e^{v_\varepsilon + F_\varepsilon} \omega_{P,\varepsilon}^n \\
&= \left(\omega_{P,\varepsilon} + \sum_{a_i < 0} dd^c a_i \log(|s_i|^2 + \varepsilon^2) + dd^c w_\varepsilon \right)^n \\
&\leqslant (C\omega_{P,\varepsilon} + dd^c w_\varepsilon)^n
\end{aligned}
$$

hence the maximum principle yields $\sup w_\varepsilon \leqslant -\inf F_\varepsilon + n \log C$, hence

$$u_\varepsilon \leqslant C + \sum_{a_i < 0} a_i \log(|s_i|^2 + \varepsilon^2) \qquad (2.4)$$

for some uniform $C > 0$.

Moreover, we know from [3] that φ_ε converges to φ_{KE}, the potential of the Kähler–Einstein metric of the pair (X, D), which is a quasi-psh function. Hence, by Hartog's Theorem (cf. [17, Theorem 3.2.12]), we know that there exists $C > 0$ such that $\varphi_\varepsilon \leqslant C$.

As a consequence, $u_\varepsilon = \varphi_\varepsilon + \sum_{a_k=1} \log \log^2 \frac{1}{|s_k|^2}$ is locally uniformly bounded above on $X \setminus D_{lc}$. Therefore, if we want to bound u_ε from above, we just need to do it locally around points at the intersection of D_{lc} and $\sum_{a_i>0} D_i$.

2.4 The upper bound II

Now that we have a partial upper bound (2.4) on u_ε, one can derive a true upper bound using Green's functions based on ideas appearing in [29, p. 141]. From now on, one can forget about the cone approximation and just remember the two bounds (2.3)–(2.4) satisfied by our potential u_ε solution of (2.1).

We fix a point $p \in D_{lc}$, and one may assume that p admits a neighborhood $\Omega \simeq \mathbb{D}^n$ where D_{lc} is given by $(z_1 \cdots z_r = 0)$ and $\sum_{a_i<0} D_i$ by $(z_{r+1} \ldots z_s = 0)$ in the holomorphic coordinates $z_1, \ldots, z_n$ induced on Ω under the identification $\Omega \simeq \mathbb{D}^n$. Equation (2.4) can be reformulated as follows

$$u_\varepsilon(z) \leqslant C \left(1 - \sum_{j=r+1}^{s} \log |z_j| \right) \tag{2.5}$$

If we knew that u_ε were quasi-psh, then we could derive a uniform upper bound from the inequality above and the arguments of [29]. But our function u_ε is only ω_P-psh, so that one cannot apply these arguments unless we have a good knowledge of (local or global) Green's functions for the Poincaré metric. In the next section, we will build a global Green's function on $X \setminus D_{lc}$ for $\Delta_{\omega_P} - 1$, study its properties, and use it to get the desired upper bound.

But before that, we will give an alternative solution consisting in using the very particular geometry of the Poincaré metric. Indeed, the Poincaré metric behaves in some way like an euclidian one when we pull it back to some appropriate "cover". The right way to formalize this is to use the quasi-coordinates for the Poincaré metric (cf. [20, 26]): they are maps from an open subset $V \subset \mathbb{C}^n$ to $\mathbb{D}_r := (\mathbb{D}^*)^r \times \mathbb{D}^{n-r}$ having maximal rank everywhere. So they are just locally invertible, but these maps are not injective in general.

To construct such quasi-coordinates on $\mathbb{D}_r$, we start from the universal covering map $\pi : \mathbb{D} \to \mathbb{D}^*$, given by $\pi(w) = e^{\frac{w+1}{w-1}}$. Formally, it sends 1 to 0. The idea is to restrict π to some fixed ball $B(0, R)$ with $1/2 < R < 1$, and compose it (at the source) with a biholomorphism Φ_η of $\mathbb{D}$ sending 0 to η, where η is a real parameter which we will take close to 1. If one wants to write an explicit formula, we set $\Phi_\eta(w) = \frac{w+\eta}{1+\eta w}$, so that the quasi-coordinate maps are given by $\Psi_\eta = (\pi \circ \Phi_\eta)^r \times \mathrm{Id}_{\mathbb{D}^{n-r}} : V = B(0, R)^r \times \mathbb{D}^{n-r} \to \mathbb{D}_r$, i.e. $\Psi_\eta(v_1, \ldots, v_r, v_{r+1}, \ldots, v_n) = (e^{\frac{1+\eta}{1-\eta} \frac{v_1+1}{v_1-1}}, \ldots, e^{\frac{1+\eta}{1-\eta} \frac{v_r+1}{v_r-1}}, v_{r+1}, \ldots, v_n)$.

Once we have said this, it is easy to see that $\mathbb{D}_r$ is covered by the images $\Psi_\eta(V)$ when η goes to 1. Now, an easy computation shows that $\Psi_\eta^* \omega_P$ is a Kähler metric on $V \subset \mathbb{C}^n$ which is uniformly (in η) quasi-isometric to the euclidian flat metric; moreover all the covariants derivatives of this metric are uniformly bounded with respect to η, but we will not need this property.

Let us go back to our situation. We started from an ω_P-psh function u_ε satisfying (2.5). Pulling it back by Ψ_η, we get a smooth function $u_{\varepsilon,\eta} := u_\varepsilon \circ \Psi_\eta$ on V which is $\Psi_\eta^* \omega_P$-psh hence (uniformly) quasi-psh by the observation above. Furthermore, as Ψ_η acts trivially on the component $\mathbb{D}^{n-r}$, we have

$$u_{\varepsilon,\eta}(v) \leqslant C \left(1 - \sum_{j=r+1}^{s} \log |v_j| \right) \tag{2.6}$$

for all $v \in V$ thanks to (2.5).

We are now in position to apply the arguments of [29], so let us set up a precise framework. For $\rho > 0$ large enough ($\rho > 2n$ would be sufficient), we have $V \Subset B(0, \rho)$; let us also pick $1/2 < R' < R$ and set $V' = B(0, R')^r \times \mathbb{D}(0, 1/2)^{n-r}$. As above, the images of V' by Ψ_η when η goes to 1 cover $(\mathbb{D}^*)^r \times \mathbb{D}(0, 1/2)^{n-r}$. We choose a cut-off function χ such that $\mathrm{Supp}(\chi) \Subset V$, and $\chi = 1$ on V' so that $d(\mathrm{Supp}(\nabla\chi), V') > 0$. Finally, we denote by $G : \bar{B}(0, \rho) \times \bar{B}(0, \rho) \to [-\infty, 0]$ the Green's function of $B(0, \rho)$. If $x \in V'$, we denote by G_x the function $G(x, \cdot)$. Then for any $x \in V'$, the function χG_x satisfies $\Delta(\chi G_x) = \delta_x + G_x \Delta\chi + \nabla G_x \cdot \nabla\chi$ (this can be verified locally, first near x, and then away from x). Therefore, if dV is the Lebesgue measure of $\mathbb{C}^n$, we have:

$$\int_{B(0,\rho)} \chi G_x \Delta u_{\varepsilon,\eta}\, dV = u_{\varepsilon,\eta}(x) + \int_{\mathrm{Supp}(\nabla\chi)} u_{\varepsilon,\eta} \left(G_x \Delta\chi + \nabla G_x \cdot \nabla\chi \right) dV$$

Remember that $u_{\varepsilon,\eta}$ is quasi-psh, so that $\Delta u_{\varepsilon,\eta} \geqslant -C$. As a consequence,

$$u_{\varepsilon,\eta}(x) \leqslant C \left(||G_x||_{L^1} + \int_{\mathrm{Supp}(\nabla\chi)} u_{\varepsilon,\eta} \left[|G_x \Delta\chi| + |\nabla G_x \cdot \nabla\chi| \right] dV \right)$$

Of course, $\nabla\chi$ and $\Delta\chi$ are bounded by some constants depending only on R', R and n. As for G_x and ∇G_x, these functions are bounded in terms of (negative powers of) $d(x, \cdot)$, therefore they are uniformly (in x) bounded on $\mathrm{Supp}(\nabla\chi)$ by the above observation that $d(\mathrm{Supp}(\nabla\chi), V') > 0$. Therefore, we have:

$$u_{\varepsilon,\eta}(x) \leqslant C \left(||G_x||_{L^1} + ||u_{\varepsilon,\eta}||_{L^1} \right) \tag{2.7}$$

Applying the Green–Riesz representation formula to the function $y \mapsto |y|^2$, we easily get that $||G_x||_{L^1} = (\rho^2 - |x|^2)/2n \leqslant \rho^2$. Moreover, thanks to Eq. (2.6), we have a uniform control $||u_{\varepsilon,\eta}||_{L^1} \leqslant C$ (remember that u_ε hence $u_{\varepsilon,\eta}$ are uniformly bounded from below already). Putting these two estimates together, we infer from (2.7):

$$u_{\varepsilon,\eta}(x) \leqslant C$$

for some constant C independent of $x \in V'$, η and ε. Pushing this inequality downwards to $(\mathbb{D}^*)^r \times \mathbb{D}(0, 1/2)^{n-r}$, we obtain

$$u_\varepsilon \leqslant C$$

on this latter open set, which ends the proof.

3 A new global Green's function

In this section, we investigate the question of the existence of appropriate global Green's functions on the complete Kähler manifold $(X \backslash D_{lc}, \omega_P)$. Adapting the arguments of [22,25], we will construct on this Riemannian manifold a positive Green's function for the operator $\Delta - 1$ whose behavior is well understood at infinity. This will enable us to give an alternative proof of Theorem A, cf. Sect. 3.3.

3.1 Existence of the Green's function

Let (M^m, g) be a complete Riemannian manifold of real dimension m, and Δ_g be the Laplacian of g. Similar to [25, p. 81], a function G defined on $M \times M \backslash \mathrm{diag}(M \times M)$ is called a *global positive Green's function* for $\Delta_g - 1$ on M if G satisfies the following properties:

(i) For any fixed $x \in M$, $(\Delta_{g(y)} - 1)G(x, y) = 0$ and $G(x, y) > 0$, for all $y \in M$, $y \neq x$;

(ii) $G(x, y) = G(y, x)$;

(iii) As $y \to x$ for fixed x, $G(x, y) = [(m-2)\sigma_{m-1}]^{-1}\mathrm{dist}(x, y)^{2-m}\big(1 + o(1)\big)$.

Here $\mathrm{dist}(x, y)$ denotes the geodesic distance between x and y in M and σ_{m-1} is the volume of the unit $(m-1)$-sphere in $\mathbb{R}^m$.

The following lemma constructs on *any* complete Riemannian manifold a global positive Green's function for $\Delta_g - 1$. This is in sharp contrast to the global positive Green's function for Δ_g (cf. [22]).

Lemma 3.1 *Let (M, g) a complete Riemannian manifold. Then (M, g) admits a global positive Green's function G for $\Delta_g - 1$ on M. Furthermore, for any $x \in M$ and any compact set B containing x, $G(x, y) \leqslant C_B$ for all $y \in M \backslash B$, where the constant $C_B > 0$ depending only on B.*

Proof The existence of G follows from almost the same argument for Theorem A.1 in Schoen–Yau [25, p. 82] (see also Li–Tam [22]), using the monotone increasing sequence of positive Dirichlet Green's functions $\{G_i\}$ on the exhaustion $\{\Omega_i\}$. The only difference is that here the operator $\Delta_g - 1$ allows us to compare G_i with the constant function.

More precisely, let $\{\Omega_i\}$ be an exhaustion of M and G_i be the positive Dirichlet Green's function on Ω_i (cf. [11, p. 157] or Remark 3.2). Fix an arbitrary $x \in M$. We need to show that the monotonic sequence

$$m_i = \sup_{y \in \partial B_r(x)} G_i(x, y)$$

is bounded from above for all $r > 0$ (to see the monotonicity, apply the maximum principle to $G_i - (1+\delta)G_{i+1}$ on $\Omega_i \smallsetminus \{x\}$ for any fixed $\delta > 0$). Suppose the contrary, i.e., there exists an $r > 0$ such that $m_i \to +\infty$. Let

$$v_i(y) = \frac{1}{m_i} G_i(x, y) \quad \text{for all } i \geqslant 1.$$

By the maximum principle

$$v_i \leqslant 1 \quad \text{on } \Omega_i \backslash B_r(x).$$

On the other hand, for any $\varepsilon > 0$, by property (iii) of G we have

$$v_i(y) \leqslant \varepsilon G_1(x, y) + 1 \quad \text{on } \overline{B}_r(x)\backslash\{x\}$$

for all sufficiently large i such that $1/m_i < \varepsilon$. Applying the diagonal process we obtain that a subsequence of v_i converges uniformly on compact subsets of $M\backslash\{x\}$ to a function v in $M\backslash\{x\}$ satisfying

$$(\Delta_g - 1)v = 0 \quad \text{in } M\backslash\{x\}$$

and $0 \leqslant v \leqslant 1$ on $M\backslash B_r(x)$ and $v(y) \leqslant \varepsilon G_1(x, y) + 1$ on $\overline{B}_r(x)\backslash\{x\}$. Letting $\varepsilon \to 0$ yields

$$0 \leqslant v \leqslant 1 \quad \text{in } M\backslash\{x\}.$$

Since $\max_{\partial B_r(x)} v_i = 1$ for all i, the function v attains its maximum value 1 at an interior point of $M\backslash\{x\}$. Applying the maximum principle to $(\Delta_g - 1)v = 0$ at the interior point yields $v \leqslant 0$ on $M\backslash\{x\}$, which contradicts $\max v = 1$. Hence, the sequence $\{m_i\}$ is bounded from above for all $r > 0$. Then as in Schoen–Yau [25, p. 83] we apply the diagonal process to obtain a global positive Green's function G on M.

For the second statement, for a given $x \in M$, assume that $x \in B \subset\subset \Omega_{i_0}$ for some $i_0 \geqslant 1$. By the previous step

$$\sup_{i \geqslant i_0} \sup_{y \in \partial B} G_i(x, y) \leqslant C$$

where $C > 0$ depends only on B. Applying the maximum principle yields

$$G_i(x, y) \leqslant C \quad \text{on } M\backslash B.$$

By the proof of the first statement [25, p. 83], a subsequence of G_i converges uniformly on compact subsets to G; hence,

$$G(x, y) \leqslant C \quad \text{on } M\backslash B.$$

Remark 3.2 The proof of Lemma 3.1 makes use of a classical fact that for a bounded domain Ω with smooth boundary in a Riemannian manifold, there exists a positive Dirichlet Green's function G satisfies the properties (i)–(iii). This fact can be proved as follows: By Duff [11, p. 104, 5.3] one obtains a local fundamental solution $\gamma(P, Q)$ for operator $\Delta - 1$ on a sufficiently small neighborhood $\mathcal{U}$ of diag($\Omega \times \Omega$). That is, given $Q \in \Omega$, $\gamma(P, Q)$ is smooth and satisfies

$$(\Delta - 1)\gamma(P, Q) = 0$$

for any P near Q and $P \neq Q$. Furthermore,

$$\gamma(P, Q) \sim \frac{1}{(m-2)\sigma_{m-1}}\text{dist}(P, Q)^{2-m}, \quad \text{as dist}(P, Q) \to 0.$$

Let

$$\Gamma(P, Q) = \eta(\text{dist}(P, Q)/\epsilon)\gamma(P, Q),$$

where $\eta = \eta(t)$ satisfies that $\eta \in C^\infty(\mathbb{R})$, $0 \leqslant \eta \leqslant 1$, $\eta \equiv 1$ for $0 \leqslant t \leqslant 1/2$ and $\eta \equiv 0$ for $t \geqslant 1$, and $\epsilon > 0$ is a small constant such that the compact support of η is contained in the neighborhood $\mathcal{U}$ of the diagonal. Fix an arbitrary $Q \in \Omega$. It follows that

$$(\Delta - 1)\Gamma(\cdot, Q) = -\delta_Q + F,$$

where $F \equiv 2\nabla\eta \cdot \nabla\gamma(\cdot, Q) + \gamma(\cdot, Q)\Delta\eta \in C_c^\infty(\Omega)$. We can solve

$$(\Delta - 1)w = -F \quad \text{in } \Omega, \quad w = 0 \quad \text{on } \partial\Omega$$

for a smooth function $w(\cdot, Q)$ on Ω. Then

$$G(P, Q) = \Gamma(P, Q) + w(P, Q)$$

is the desired Green's function. That $G(P, Q) = G(Q, P)$ is proven in Duff [11, p. 158].

From now on we let $M = X \backslash D_{lc}$. Then (M, ω) is a complete Kähler manifold of finite volume, where we have a notion of Hölder spaces relatively to quasi-coordinates as introduced by Cheng–Yau [8, p. 515], see also [20,26] for the definitions that we don't recall here as we won't use them in the following.

Corollary 3.3 *Let us endow $M := X \backslash D_{lc}$ with the metric ω. Then M admits a global positive Green's function G which in particular belongs to $L^1(M)$. Furthermore, for any $x \in M$ and any $r > 0$, $G(x, y)$ as a function of $y \in M$ satisfies*

$$\|G(x, y)\|_{C^{k,\alpha}(M \backslash B_r(x))} \leqslant C(k, \alpha) \sup_{y \in \partial B_r} G(x, y)$$

where $C^{k,\alpha}$, $k \geqslant 0$, $0 < \alpha < 1$, *is the Hölder space in the sense of Cheng–Yau,* $C(k,\alpha) > 0$ *is a constant depending only on k and α.*

Proof That G is in $L^1(M)$ follows immediately from Lemma 3.1 and the finiteness of volume. To see the estimate, note that (M, ω) has bounded geometry in terms of the quasi-coordinates. In constructing G we apply the Schauder interior estimates to the Cheng–Yau's Hölder spaces, and then, using a diagonal process we can pass from G_i to G. □

3.2 Properties of the Green's function

The following result is a slight variant of [29, Lemma 2 p. 138].

Lemma 3.4 *Let* $M = X \backslash D_{lc}$ *with metric* ω. *Given* $x \in M$ *and a small ball* $B \equiv B_r(x) \subset M$. *For any* $f, h \in C^{k,\alpha}(M\backslash B)$, $k \geqslant 2$, $0 < \alpha < 1$,

$$\int_{M\backslash B} \mathrm{div}(f\nabla h) = -\int_{\partial B} f\frac{\partial h}{\partial \nu}$$

where the divergence div *and gradient* ∇ *are both with respect to* ω, *and* ∂B *is oriented according to the outer unit normal* ν.

Proof As in [29, p. 138] we use the cutoff function $\chi_m(\rho)$ on M such that $\chi_m \equiv 0$ for $\rho \leqslant m$, $0 \leqslant \chi_m \leqslant 1$ for $m \leqslant \rho \leqslant m+1$, and $\chi_m \equiv 1$ for $\rho \geqslant m+1$. Here

$$\rho = \sum_{i=1}^{k} \log(-\log|s_i|^2) \to +\infty \quad \text{as } x \to D_{lc},$$

and $m \geqslant 1$ such that $B = B_r(x) \subset\subset \{\rho < m\}$. Write

$$\int_{M\backslash B} \mathrm{div}(f\nabla h) = \int_{M\backslash B} \mathrm{div}(\chi_m f\nabla h) + \int_{M\backslash B} \mathrm{div}[(1-\chi_m)f\nabla h].$$

Note that $1 - \chi_m(\rho)$ has compact support $\{\rho \leqslant m+1\}$ in M and $1 - \chi_m \equiv 1$ on $\overline{B}$. Applying the usual Stokes' theorem yields

$$\int_{M\backslash B} \mathrm{div}[(1-\chi_m)f\nabla h] = -\int_{\partial B}(1-\chi_m)f\nabla h = -\int_{\partial B}(1-\chi_m)f\nabla h.$$

Then following [29, p. 138] we obtain

$$\lim_{m\to+\infty}\int_{M\backslash B}\mathrm{div}(\chi_m f\nabla h) = 0$$

by using Lebesgue's dominated convergence theorem. □

Corollary 3.5 *Let us endow $M = X \backslash D_{lc}$ with the metric ω, and let G be the global positive Green's function obtained in Corollary 3.3. For any $x \in M$ and any $f \in C^{k,\alpha}(M)$ with $k \geqslant 2$,*

$$f(x) = -\int_M G(x, y)(\Delta_{\omega(y)} - 1)f(y)\omega^n.$$

Proof Fix an arbitrary $\varepsilon > 0$. By Lemma 3.4,

$$\int_{M \backslash B_\varepsilon(x)} G(x, \cdot)(\Delta - 1)f$$
$$= \int_{M \backslash B_\varepsilon(x)} f(\Delta - 1)G - \int_{\partial B_\varepsilon(x)} G \frac{\partial f}{\partial \nu} + \int_{\partial B_\varepsilon(x)} f \frac{\partial G}{\partial \nu} = -f(x),$$

in view of properties (i) and (iii) of G. Since $G(x, \cdot) \in L^1(M)$, by Lebesgue's dominated convergence theorem

$$\int_M G(x, \cdot)(\Delta - 1)f = \lim_{\varepsilon \to 0^+} \int_{M \backslash B_\varepsilon(x)} G(x, \cdot)(\Delta - 1)f = -f(x).$$

$\square$

3.3 An alternative proof of Theorem A

We borrow the notations of Sect. 2.2: let X be a compact Kähler manifold X endowed with a snc divisor D_{lc}, let us set $M := X \backslash D_{lc}$, and let us consider the Monge–Ampère equation on M:

$$(\omega + dd^c u)^n = \frac{e^{u+F} \omega^n}{\prod_i |s_i|^{2a_i}}$$

where $D_{klt} = \sum a_i(s_i = 0)$ is a divisor with snc support whose coefficients a_i belong to $(-\infty, 1)$, and ω is a metric with Poincaré-type singularities along D_{lc}, the latter divisor being also assumed to have normal crossings with $\text{Supp}(D_{klt})$. Finally, F is a smooth function when read on the quasi-coordinate, i.e. $F \in C^{k,\alpha}(M)$ for all $k \geqslant 2$ and $0 < \alpha < 1$.

We will assume part of the results of Sect. 2.2, namely that u is bounded below. We claim that by Corollary 3.3 and Corollary 3.5, we can derive a true upper bound for u.

Indeed, let us perturb (as in Sect. 2.2) the above Monge–Ampère equation by the following one, for $\varepsilon > 0$,

$$(\omega + dd^c u_\varepsilon)^n = e^{u_\varepsilon + F + f_\varepsilon} \omega^n$$

where $f_\varepsilon = -\sum_i a_i \log(|s_i|^2 + \varepsilon^2)$. We know that the latter equation has a unique solution $u_\varepsilon \in C^{k,\alpha}(M)$ for all $k \geqslant 2$ and $0 < \alpha < 1$, and that u_ε converges to u weakly on X, and smoothly on the compact sets of $M \backslash D_{klt}$.

Combining the arithmetic mean–geometric mean inequality

$$\left(\frac{(\omega + dd^c u_\varepsilon)^n}{\omega^n}\right)^{\frac{1}{n}} \leqslant \frac{n + \Delta_\omega u_\varepsilon}{n}$$

with the basic inequality $e^t \geqslant 1 + t$ yields

$$(\Delta_\omega - 1)u_\varepsilon \geqslant F + f_\varepsilon$$

Multiplying this inequality by $G(x, y)$ gives

$$u_\varepsilon(x) = -\int_M G(x, y)(\Delta_\omega - 1)u_\varepsilon(y)\omega^n(y)$$

$$\leqslant -\int_M G(x, y)(F + f_\varepsilon)(y)\omega^n(y) \leqslant C$$

in view of the fact that

$$\int_{M\backslash B} G(x, y)(-\log|s_i|^2)(y)\omega(y)^n \leqslant C_B \int_{M\backslash B} (-\log|s_i|^2)(y)\omega(y)^n \leqslant C$$

for each i. Here B is a compact set in M containing x, and $C > 0$ is a constant depending only on n, a_i and C_B.

4 The case of singular pairs

The goal of this section is to explain and prove Theorem B about the behavior near the boundary divisor of the Kähler–Einstein metric associated with a log canonical pair (X, D) such that $K_X + D$ is ample.

4.1 Mixed cone and cusp singularities

Let (X, D) be a pair consisting in a complex manifold X and a $\mathbb{R}$-divisor D having simple normal crossing support and coefficients in $[0, 1]$. A Kähler metric ω on $X_0 := X\backslash\text{Supp}(D)$ is said to have mixed cone and cusp (also called Poincaré) singularities along D if ω is locally quasi-isometric to the model

$$\omega_{\text{mod}} := \sum_{j=1}^r \frac{i dz_j \wedge d\bar{z}_j}{|z_j|^{2(1-\beta_j)}} + \sum_{k=r+1}^s \frac{i dz_k \wedge d\bar{z}_k}{|z_k|^2 \log^2 |z_k|^2} + \sum_{l=r+s+1}^n i dz_l \wedge d\bar{z}_l$$

whenever (X, D) is locally isomorphic to $(X_{\text{mod}}, D_{\text{mod}})$, where $X_{\text{mod}} = (\mathbb{D}^*)^r \times (\mathbb{D}^*)^s \times \mathbb{D}^{n-(s+r)}$, $D_{\text{mod}} = (1 - \beta_1)[z_1 = 0] + \cdots + (1 - \beta_r)[z_r = 0] + [z_{r+1} = 0] + \cdots + [z_{r+s} = 0]$; where $\beta_j \in (0, 1)$ and $\mathbb{D}$ (resp. $\mathbb{D}^*$) is the disc (resp. punctured disc) of radius $1/2$ in $\mathbb{C}$.

In [13] (and later in full generality in [12]), it was proved that given a compact Kähler manifold X and a divisor D with simple normal crossing support and coefficients in [0, 1] such that $K_X + D$ is ample, there exists a unique Kähler metric ω on X_0 with mixed cone and cusp singularities along D such that $\mathrm{Ric}\,\omega = -\omega$.

Of course this metric coincides with the Kähler–Einstein metric constructed in [3] in the more general case of singular log canonical pairs. Our goal in this second section is to generalize the result of [12,13] to this singular setting, as we will explain in the next paragraph after recalling the necessary definitions.

4.2 Log canonical pairs

Definition 4.1 A log canonical pair (X, D) consists of a complex normal variety X and an *effective* Weil divisor D such that $K_X + D$ is $\mathbb{Q}$-Cartier, and such that for any log resolution $\pi : X' \to X$ of (X, D), the coefficients a_i defined by the formula $K_{X'} = \pi^*(K_X + D) + \sum a_i E_i$ satisfy $a_i \geqslant -1$ (here E_i is either exceptional or the strict transform of a component of D).

Definition 4.2 Let (X, D) be a log pair. The simple normal crossing (snc) locus of the pair, denoted by $(X, D)_{\mathrm{reg}}$, is the locus of points $x \in X$ such that the pair (X, D) is log smooth at x, i.e., such that there exists a Zariski open set $U \ni x$ satisfying that $U \subset X_{\mathrm{reg}}$ and that the divisor $D_{|U}$ has simple normal crossing support.

The snc locus is a Zariski open set whose complement has codimension at least 2 by normality of X. If now X is projective, (X, D) is log canonical and $K_X + D$ is ample, then the main result of [3] provides a unique *Kähler–Einstein metric* ω_{KE} with negative curvature, which is smooth on $X_{\mathrm{reg}} \backslash \mathrm{Supp}(D)$. What about further regularity? So far, it is really hard to tell anything about the local behavior of this metric near the singular points of X; but if we look at what happens at points of the boundary divisor D where it is smooth (or merely snc), then we have a better understanding of how ω looks like. Indeed, if (X, D) is klt (i.e., the coefficients a_i above satisfy $a_i > -1$), it was proved first partially in [14] and then in full generality in [12] that ω_{KE} has cone singularities along D on the snc locus $(X, D)_{\mathrm{reg}}$. We now aim to generalize this result to the log canonical case:

Theorem 4.3 *Let (X, D) be a projective log canonical pair such that $K_X + D$ is ample. Then its Kähler–Einstein metric ω_{KE} has mixed cone and cusp singularities along D on the snc locus $(X, D)_{\mathrm{reg}}$ of the pair.*

We can deduce from this statement how the Kähler–Einstein metric of a stable variety (i.e., a projective variety X with semi-log canonical singularities such that K_X is ample, cf. [3]) behaves near the double crossing points. Recall that a double crossing point is a point near which the variety is locally analytically isomorphic to $0 \in \{xy = 0\} \subset \mathbb{C}^{n+1}$.

Corollary 4.4 *Let X be a stable variety. Then its Kähler–Einstein metric is locally quasi-isometric to a cusp near the double crossing points.*

Let us explain what it means. If p is such a point and $\nu : X^\nu \to X$ is the normalization morphism, then $\nu^{-1}(p)$ consists of two distinct points q', q'' sitting on the conductor divisor D^ν, and the pair (X^ν, D^ν) is log smooth at q', q'' (actually D^ν is even smooth near those points). The corollary expresses that the pull-back $\nu^* \omega_{KE}$ of the Kähler–Einstein metric has cusp singularities along D^ν near q' and q''.

This generalizes the picture existing for stable curves. Indeed, if C is a stable curve, let C' be its normalization and D' be the reduced divisor on C' whose support consists of the preimage of the nodes. Then $K_{C'} + D'$ is ample, and each connected component of $C' \backslash D' (= C_{reg})$ has a unique hyperbolic metric which has a cusp near each point in the support of D'.

Proof There is not much more left to say. Indeed, with the above notations, the conductor D^ν is a reduced divisor; moreover, (X^ν, D^ν) is log canonical, $K_{X^\nu} + D^\nu$ is ample, and the pair is log smooth at each point above a double crossing point. So if we apply Theorem 4.3 at those points, we get exactly the statement claimed in the corollary. $\qquad\qquad\square$

4.3 Proof of Theorem B

4.3.1 The set-up

In order to keep more usual notations, we assume that the initial log pair is (Y, Δ), and we consider a log resolution $\pi : (X, D) \to (Y, \Delta)$ of the pair. Here, $D = \sum a_i D_i$ is a divisor on X with snc support, consisting of π-exceptional divisors (with arbitrary coefficients in $(-\infty, 1]$) and of the strict transforms of the components of Δ (with coefficients in $[0, 1]$). The Kähler–Einstein ω_{KE} for (X, D), or equivalently the pullback of the KE metric for (Y, Δ) by π can be written as $\omega_{KE} = \theta + dd^c \varphi$ where $\theta \in c_1(\pi^*(K_Y + \Delta))$ is a smooth semipositive and big form and φ is a θ-psh function solving the Monge–Ampère equation

$$MA(\varphi) = \frac{e^\varphi dV}{\prod_i |s_i|^{2a_i}}$$

where s_i are non-zero sections of $\mathcal{O}_X(D_i)$, $|\cdot|_i$ are smooth hermitian metrics on $\mathcal{O}_X(D_i)$, and dV is a smooth volume form on X. Let us also introduce as before the convenient notation $D_{lc} := \sum_{a_k=1} D_k$.

By [3, Theorem 3.5] we know that the solution φ is the limit (as both ε and t tend to 0) of the quasi-psh functions $\varphi_{t,\varepsilon}$ solving

$$(\theta + t\omega_0 + dd^c \varphi_{t,\varepsilon})^n = \frac{e^{\varphi_{t,\varepsilon}} dV}{\prod_{a_j < 1}(|s_i|^2 + \varepsilon^2)^{a_i} \prod_{a_k=1} |s_k|^2} \tag{4.1}$$

where ω_0 is some fixed Kähler form.

We will divide the proof of Theorem B in three steps. In the first two, we will be dealing with the L^∞ estimate on the potential (upper bound then lower bound), and in the last one, we will focus on the Laplacian estimate.

4.3.2 The upper bound

To find the upper bound, we mimic what we did in the case of log smooth pair, as the loss of positivity will not hinder the previous method. We set $\varphi_P := -\sum_{a_k=1} \log \log^2 |s_k|^2$, $u_\varepsilon := \varphi_{t,\varepsilon} - \varphi_P$, and $w_\varepsilon := u_\varepsilon - \sum_{a_i < 0} a_i \log(|s_i|^2 + \varepsilon^2)$. Actually, u_ε and w_ε depend on t, but we choose not to underline this dependence so as to keep the notations lighter. If ω_P denotes a metric with Poincaré singularities along $\sum_{a_k=1} D_k$, e.g. $\omega_P = \omega_0 + dd^c \varphi_P$, then we have

$$(\theta + t\omega_0 + dd^c \varphi_P + dd^c u_\varepsilon)^n = \frac{e^{w_\varepsilon + f_\varepsilon} \omega_P^n}{\prod_{0 < a_i < 1}(|s_i|^2 + \varepsilon^2)^{a_i}} \tag{4.2}$$

for some uniformly bounded function f_ε on $X \backslash D_{lc}$ (this function does not depend on t).

The function w_ε is bounded on the complete manifold $(X \backslash D_{lc}, \omega_P)$ so one can apply Yau's maximum principle [31] to this function. So let (x_m) be a sequence such that $w_\varepsilon(x_m) \to \sup w_\varepsilon$, and $dd^c w_\varepsilon(x_m) \leqslant \frac{1}{m}\omega_P$. As v_ε satisfies $dd^c w_\varepsilon \geqslant dd^c u_\varepsilon - C\omega_0 \geqslant dd^c u_\varepsilon - C'\omega_P$, we have

$$\theta + t\omega_0 + dd^c \varphi_P + dd^c u_\varepsilon \leqslant C\omega_P + dd^c w_\varepsilon$$

and therefore

$$\begin{aligned}
\frac{e^{w_\varepsilon(x_m) + f_\varepsilon(x_m)} \omega_P^n(x_m)}{\prod_{0 < a_i < 1}(|s_i(x_m)|^2 + \varepsilon^2)^{a_i}} &= (\theta + t\omega_0 + dd^c \varphi_P + dd^c u_\varepsilon)^n(x_m) \\
&\leqslant (C\omega_P + dd^c w_\varepsilon)^n(x_m) \\
&\leqslant (C + 1/m)\omega_P^n(x_m)
\end{aligned}$$

so that

$$w_\varepsilon(x_m) \leqslant -f_\varepsilon(x_m) + \sum_{0 < a_i < 1} a_i \log(|s_i(x_m)|^2 + \varepsilon^2) + \log(C + 1/m)$$

hence $\sup w_\varepsilon \leqslant C$, or equivalently

$$u_\varepsilon \leqslant C + \sum_{a_i < 0} a_i \log(|s_i|^2 + \varepsilon)^2$$

But u_ε is $(\theta + t\omega_0 + dd^c \varphi)$-psh, hence also $C\omega_P$-psh, so the arguments of the first part of this article can be applied the same way in this situation, and they yield:

$$u_\varepsilon \leqslant C \tag{4.3}$$

4.3.3 The lower bound

This is where we have to pay for the loss of positivity of $K_X + D$. We know that there exists an effective $\mathbb{R}$-divisor $E = \sum c_\alpha E_\alpha$, π-exceptional, such that $K_X + D - E =$

$\pi^*(K_Y + \Delta) - E$ is ample. Therefore, one can find a Kähler metric ω_0 on X, non-zero sections s_α of $\mathcal{O}_X(E_\alpha)$, and hermitian metrics $|\cdot|_\alpha$ on these bundles such that the function $\chi = \sum c_\alpha \log|s_\alpha|^2$ satisfies:

$$\theta + dd^c\chi = \omega_0 + [E]$$

Recall that in Sect. 2.2, we introduced the potential ψ_ε of the regularized cone metric; it is a uniformly bounded ω_0-psh function on X, such that the metric $\omega_{P,\varepsilon} := (1+t)\omega_0 + dd^c\varphi_P + dd^c\psi_\varepsilon$ on $X\backslash D_{lc}$ is complete, with bounded bisectional curvature and satisfies

$$\prod_{0<a_i<1}(|s_i|^2 + \varepsilon^2)^{-a_i}\omega_P^n = e^{G_\varepsilon}\omega_{P,\varepsilon}^n$$

for some smooth function G_ε which is uniformly bounded in ε (and t, which is why we choose not to emphasize the dependence of G_ε on t). One should emphasize that this metric $\omega_{P,\varepsilon}$ has approximate cone singularities not only along the strict transform of $\Delta - \lceil\Delta\rceil$ but also along some exceptional divisors. Setting $v_\varepsilon := u_\varepsilon - \psi_\varepsilon - \chi$, Eq. (4.1) becomes, on $X\backslash E$:

$$(\omega_{P,\varepsilon} + dd^c v_\varepsilon)^n = e^{v_\varepsilon + F_\varepsilon}\omega_{P,\varepsilon}^n \qquad (4.4)$$

where $F_\varepsilon = G_\varepsilon + \psi_\varepsilon + \chi - \sum_{a_i<0} a_i \log(|s_i|^2 + \varepsilon^2)$, and from the remarks above, $\sup F_\varepsilon \leqslant C$ for some C independent of ε and t. The job would be done if one could apply Yau's maximum principle to v_ε on the complete manifold $(X\backslash D_{lc}, \omega_{P,\varepsilon})$. But v_ε is not smooth along E, so we should be careful. Fortunately, $v_\varepsilon = -\chi + O(1)$ tends to $+\infty$ near E, so one can run the proof of Yau's maximum principle without any change: let us first introduce, for every positive integer m, the function $h_m := v_\varepsilon - \frac{1}{m}\varphi_P$. This function is smooth on $X\backslash(D_{lc} \cup E)$ and tends to $+\infty$ near the boundary. Therefore, it attains its minimum at some point x_m in $X\backslash(D_{lc} \cup E)$. Then, $0 \leqslant dd^c h_m(x_m) = dd^c v_\varepsilon(x_m) - \frac{1}{m}dd^c\varphi_P(x_m)$ so that $dd^c v_\varepsilon(x_m) \geqslant -\frac{C}{m}\omega_0(x_m) \geqslant -\frac{C'}{m}\omega_{P,\varepsilon}(x_m)$ as φ_P and ψ_ε are uniformly quasi-psh. Plugging this inequality into (4.4), we find $\inf v_\varepsilon \geqslant -\sup F_\varepsilon \geqslant -C$, hence

$$u_\varepsilon \geqslant C + \chi \qquad (4.5)$$

4.3.4 The Laplacian estimate

The metric $\omega_{P,\varepsilon}$ on $X\backslash D_{lc}$ is complete and has bounded curvature, but when ε goes to zero, its curvature may blow up (in both directions) due to the conic part. In [12], a new Laplacian estimate has been introduced to deal specifically with that kind of geometries (cf. Section 6.3). More precisely, if we write $\omega := \omega_{P,\varepsilon}$ and $\omega' = \omega_{P,\varepsilon} + dd^c v_\varepsilon$, then we get from (4.4) that $\omega'^n = e^{v_\varepsilon + F_\varepsilon}\omega^n$, and it is shown in [12] that there exists a smooth and uniformly bounded function Ψ_ε on X satisfying on $X\backslash(D_{lc} \cup E)$:

$$\Delta_{\omega'}(\log \mathrm{tr}_\omega\omega' + \Psi_\varepsilon) \geqslant -C\mathrm{tr}_{\omega'}\omega$$

for some constant C *independent of ε*. This constant takes into account a lower bound for the ω-Laplacian of $v_\varepsilon + F_\varepsilon$ (the existence of this bound is also proved in [12]). As $\omega' = \omega + dd^c v_\varepsilon$, we infer:

$$\Delta_{\omega'}(\log \operatorname{tr}_\omega \omega' + \Psi_\varepsilon - (C+1)v_\varepsilon) \geqslant \operatorname{tr}_{\omega'} \omega - n(C+1)$$

The function inside the Laplacian is smooth on $X \setminus (D_{lc} \cup E)$ and tends to $-\infty$ near E. Therefore, one can apply the same maximum principle as in the last subsection: introduce $H := \log \operatorname{tr}_\omega \omega' + \Psi_\varepsilon - (C+1)v_\varepsilon$ and $H_m := H + \frac{1}{m}\varphi_P$. By construction H_m tends to $-\infty$ near D_{lc} and E, so we can choose a point x_m outside of these divisors where H_m attains its maximum. At this point, we have $0 \geqslant dd^c H_m = dd^c H + \frac{1}{m}dd^c\varphi$ so that at this point again, we find $dd^c H \leqslant \frac{C}{m}\omega_0 \leqslant \frac{C_\varepsilon}{m}\omega'$. Using a basic inequality, we find

$$\log \operatorname{tr}_\omega \omega'(x_m) \leqslant (v_\varepsilon + F_\varepsilon)(x_m) + (n-1)\log\left[n(C+1) + nC_\varepsilon/m\right] \qquad (4.6)$$

Therefore, as Ψ_ε is uniformly bounded and v_ε is uniformly bounded below, we have:

$$
\begin{aligned}
\log \operatorname{tr}_\omega \omega' \quad &= \quad H + (C+1)v_\varepsilon - \Psi_\varepsilon \\
&\leqslant \quad \sup_m H(x_m) + (C+1)v_\varepsilon + C \\
&\underset{(4.6)}{\leqslant} \quad \sup_m \left[(v_\varepsilon + F_\varepsilon)(x_m) + (n-1)\log\left[n(C+1) + nC_\varepsilon/m\right]\right. \\
&\qquad\qquad \left. -(C+1)v_\varepsilon)(x_m)\right] + (C+1)v_\varepsilon + C \\
&\leqslant \quad -C\inf v_\varepsilon + \sup F_\varepsilon + (C+1)v_\varepsilon + C \\
&\underset{(4.3)-(4.5)}{\leqslant} \quad C - C\chi
\end{aligned}
$$

So in the end, we have proved that the approximate KE metric $\omega_{\mathrm{P},\varepsilon} + dd^c v_\varepsilon$ satisfies on $X \setminus E$:

$$C^{-1}e^{C\chi}\omega_{\mathrm{P},\varepsilon} \leqslant \omega_{\mathrm{P},\varepsilon} + dd^c v_\varepsilon \leqslant Ce^{-C\chi}\omega_{\mathrm{P},\varepsilon}$$

for some constant $C > 0$ independent of ε and t. As χ is locally bounded on $X \setminus E = \pi^{-1}((Y, \Delta)_{\mathrm{reg}})$, this ends the proof of Theorem B.

References

1. Aubin, T.: Équations du type Monge-Ampère sur les variétés Kählériennes compactes. Bull. Sci. Math. **102**, 63–95 (1978)
2. Boucksom, S., Eyssidieux, P., Guedj, V., Zeriahi, A.: Monge–Ampère equations in big cohomology classes. Acta Math. **205**(2), 199–262 (2010)
3. Berman, R.J., Guenancia, H.: Kähler-Einstein metrics on stable varieties and log canonical pairs. GAFA. arXiv:1304.2087 (2013, to appear)
4. Brendle, S.: Ricci flat Kähler metrics with edge singularities. IMRN. arXiv:1103.5454 (2011, to appear)
5. Campana, F., Guenancia, H., Păun, M.: Metrics with cone singularities along normal crossing divisors and holomorphic tensor fields. Ann. Sci. Éc. Norm. Sup. **46**, 879–916 (2013)

6. Claudon, B.: Γ-reduction for smooth orbifolds. Manuscr. Math. **127**(4), 521–532 (2008)
7. Cheng, S.Y., Yau, S.T.: Differential equations on Riemannian manifolds and their geometric applications. Commun. Pure Appl. Math. **28**(3), 333–354 (1975)
8. Cheng, S.Y., Yau, S.-T.: On the existence of a complete Kähler metric on non-compact complex manifolds and the regularity of Fefferman's equation. Commun. Pure Appl. Math. **33**, 507–544 (1980)
9. Donaldson, S.K.: Kähler metrics with cone singularities along a divisor. In: Essays in Mathematics and its Applications, pp. 49–79. Springer, Heidelberg (2012)
10. Donaldson, S., Sun, S.: Gromov–Hausdorff limits of Kähler manifolds and algebraic geometry. arXiv:1206.2609 (2012)
11. Duff, G.F.D.: Partial Differential Equations, Mathematical Expositions, vol. 9. University of Toronto Press, Toronto (1956)
12. Guenancia, H., Păun, M.: Conic singularities metrics with prescribed Ricci curvature: the case of general cone angles along normal crossing divisors. arXiv:1307.6375 (2013)
13. Guenancia, H.: Kähler–Einstein metrics with mixed Poincaré and cone singularities along a normal crossing divisor. Ann. Inst. Fourier. arXiv:1201.0952 (2012, to appear)
14. Guenancia, H.: Kähler–Einstein metrics with cone singularities on klt pairs. Int. J. Math. **24** (2013)
15. Guenancia, H.: Semistability of the tangent sheaf of singular varieties. arXiv:1502.03711 (2015)
16. Guedj, V., Zeriahi, A.: The weighted Monge–Ampère energy of quasi plurisubharmonic functions. J. Funct. Anal. **250**, 442–482 (2007)
17. Hörmander, L.: Notions of Convexity. Birkhäuser, Basel (1994)
18. Jeffres, T.: Uniqueness of Kähler–Einstein cone metrics. Publ. Mat. **44**(2), 437–448 (2000)
19. Jeffres, T., Mazzeo, R., Rubinstein, Y.: Kähler–Einstein metrics with edge singularities. Ann. Math. arXiv:1105.5216 (with an appendix by C. Li and Y. Rubinstein) (2011, to appear)
20. Kobayashi, R.: Kähler-Einstein metric on an open algebraic manifolds. Osaka J. Math. **21**, 399–418 (1984)
21. Kołodziej, S.: The complex Monge–Ampère operator. Acta Math. **180**(1), 69–117 (1998)
22. Li, P., Tam, L.F.: Symmetric Green's functions on complete manifolds. Am. J. Math. **109**(6), 1129–1154 (1987)
23. Mazzeo, R.: Kähler–Einstein metrics singular along a smooth divisor. Journées Équations aux dérivées partielles (Saint Jean-de-Mont, 1999) (1999)
24. Mok, N., Yau, S.-T.: Completeness of the Kähler–Einstein metric on bounded domains and the characterization of domains of holomorphy by curvature conditions. In: The mathematical heritage of Henri Poincaré, Part 1 (Bloomington, Ind.) Proc. Sympos. Pure Math., vol. 39, pp. 41–59. Amer. Math. Soc., Providence (1980)
25. Schoen, R., Yau, S.-T.: Lectures on differential geometry. Conference Proceedings and Lecture Notes in Geometry and Topology, I. International Press, Cambridge (1994)
26. Tian, G., Yau, S.-T.: Existence of Kähler–Einstein metrics on complete Kähler manifolds and their applications to algebraic geometry. Adv. Ser. Math. Phys. **11**, 574–628 (1987) [Mathematical aspects of string theory (San Diego, Calif., 1986)]
27. Tian, G., Yau, S.-T.: Complete Kähler manifolds with zero Ricci curvature. I. J. Am. Math. Soc. **3**(3), 579–609 (1990)
28. Wu, D.: Kähler–Einstein metrics of negative Ricci curvature on general quasi-projective manifolds. Commun. Anal. Geom. **16**(2), 395–435 (2008)
29. Wu, D.: Good Kähler metrics with prescribed singularities. Asian J. Math. **13**(1), 131–150 (2009)
30. Yao, C.: Existence of weak conical Kähler–Einstein metrics along smooth hypersurfaces. arXiv:1308.4307 (2013)
31. Yau, S.-T.: Harmonic functions on complete Riemannian manifolds. Commun. Pure Appl. Math. **28**, 201–228 (1975)
32. Yau, S.-T.: A general Schwarz lemma for Kähler manifolds. Am. J. Math. **100**, 197–203 (1978)
33. Yau, S.-T.: Métriques de Kähler–Einstein sur les variétés ouvertes. In: (Première Classe de Chern et courbure de Ricci: Preuve de la conjecture de Calabi), Séminaire Palaiseau, vol. 58. Astérisque, pp. 163–167 (1978)
34. Yau, S.-T.: On the Ricci curvature of a compact Kähler manifold and the complex Monge–Ampère equation. I. Commun. Pure Appl. Math. **31**, 339–411 (1978)