#### **ORIGINAL ARTICLE**

# Effect of weld travel speed on solidification cracking behavior. Part 1: weld metal characteristics



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#### Abstract

Solidification cracking is a weld defect common to certain susceptible alloys rendering many of them unweldable. It forms and grows continuously behind a moving weld pool within the two-phase mushy zone and involves a complex interaction between thermal, metallurgical, and mechanical factors. Research has demonstrated the ability to minimize solidification cracking occurrence by using appropriate welding parameters. Despite decade's long efforts to investigate weld solidification cracking, there remains a lack of understanding regarding the particular effect of travel speed. While the use of the fastest welding speed is usually recommended, this rule has not always been confirmed on site. Varying welding speed has many consequences both on stress cells surrounding the weld pool, grain structure, and mushy zone extent. Experimental data and models are compiled to highlight the importance of welding speed on solidification cracking. This review is partitioned into three parts: part I focuses on the effects of welding speed on weld metal characteristics, part II reviews the data of the literature to discuss the importance of selecting properly the metrics, and part III details the different methods to model the effect of welding speed on solidification cracking occurrence.

Keywords Solidification cracking  $\cdot$  Welding  $\cdot$  Welding speed  $\cdot$  Crack initiation  $\cdot$  Crack growth

Abbreviations					
h	Plate thickness				
S	Travel speed				
t	Time				
x	Direction of heat source displacement				
BTR	Brittle temperature range				
CET	Columnar-to-equiaxed transition				
CSZ	Crack-susceptible zone				
G	Temperature gradient				
Н	Heat input				
Ι	Welding current				
Κ	Thermal conductivity				
LHC	Linear heat content				
Q	Welding power				
R	Solidification growth rate				

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$R_{\rm i}$	Radius of weld cross section
SCTR	Solidification cracking temperature range
Т	Temperature
U	Welding voltage
$\alpha$	Thermal diffusivity
ε	Strain
ε	Strain rate
ξ	Traveling coordinates
$\eta$	Welding efficiency

# **1 Introduction**

Solidification cracking is a commonly encountered defect during welding, especially in high-sulfur steels, austenitic steels, and aluminum alloys. Solidification cracks form due to a complex interplay of mechanical, thermal, and metallurgical factors. Their formation is strongly dependent on both material composition and welding parameters. To increase productivity, fabricators aim at reducing manufacturing time by increasing welding speed. This commonly implies using laser and electron-beam welding processes that involve welding speeds  $(10^1-10^2 \text{ mm s}^{-1})$  faster than commonly encountered in arc welding processes  $(10^0-10^1 \text{ mm s}^{-1})$ . However, while the use of the fastest travel speeds in arc welding to avoid solidification cracking has been commonly accepted, it seems not to always apply for the faster speeds encountered during beam welding that can lead to numerous weld bead defects [1, 2] including solidification cracking itself.

Solidification cracking, described by Campbell as "an uniaxial tensile failure in weak materials" [3], appears at the solidification end inside a mushy zone that is subjected to tensile strains. The microstructure forms in the solidification zone, referred to as the mushy zone, located at the rear of the melting zone and bordered by two isothermal surfaces corresponding to liquidus and solidus temperatures. The semi-solid in the mushy zone has little ductility in the terminal stage of solidification, when the liquid fraction is no longer high enough for grains to move around and rearrange in order to accommodate tensile strains. When liquid feeding cannot adequately compensate solidification shrinkage and thermal contraction of the mushy zone, solidification cracking occurs along grain boundaries.

Solidification by epitaxial growth of columnar grains is generally observed in welding from the border of the fusion zone to the center due to high thermal gradients. However, equiaxed dendritic grains can also form in the center area of the fusion zone. A susceptible zone in between coalescence and coherency does exist in which solidification cracking may possibly form (Fig. 1). The coalescence temperature corresponds to the minimum solid fraction so that the mushy zone has a mechanical resistance to shearing, usually induced by dendrite entanglement. The coherency temperature is the temperature at which the first solid bridges form in between grains, leading to a mechanical resistance in tension.

The hot cracking temperature range, also referred to as brittle temperature range (BTR) and solidification cracking

 Table 1
 Low-alloy steel properties used to calculate CSZ length in Eq.

 8 and Fig. 12 [59]

Properties	Values	
Current I (A)	38 + 25·s	
Power Q (W)	$\eta$ ·U·I	
Tension $U(V)$	11	
Welding efficiency $\eta$	0.40	
Lower temperature of CSZ $T_2$ (°C)	1000	
Upper temperature of CSZ $T_1$ (°C)	1500	
Sheet thickness $h$ (mm)	2	
Thermal conductivity K (cal·cm <sup>-1</sup> s <sup>-1</sup> )	0.1	
Thermal diffusivity $\alpha$ (cm <sup>2</sup> s <sup>-1</sup> )	0.2	

temperature range (SCTR), relates to the interval of temperatures over which solidification cracking is likely to occur. It is usually associated to the coalescence-to-coherency range. It is argued that a large solidification range permits a large buildup of strain and a greater likelihood to crack [5]. It is also believed that the faster the welding speed is conducted, the faster the cooling rate and the less time the weld is exposed to thermal strains with the BTR. Moreover, crack growth speed being fixed by welding speed, fast welding speeds tend to reduce cracking susceptibility as the crack tip cannot be continuously maintained in the mushy zone but this valuable effect is countered by the limited time available for backfilling the thermal strain-induced opening of the mushy zone.

Little research work has been reported in the literature on the relationship between travel speed and solidification cracking. Part I of the present work focuses on reviewing the particular effect of welding speed on weld metal characteristics. The observations deal with metallurgical, thermal, and welding conditions, which change when varying travel speed.



#### 2 Grain structure

Grain structure and thus weld texture are affected by the travel speed as the weld pool shape is elongated at fast welding speeds. Grain boundaries are oriented preferentially perpendicular to the isotherms. Therefore, the shapes of the isotherm lines (and in particular the liquidus) that are dependent on the welding process parameters influence the grain structure.

#### 2.1 Ferrous alloys

Ferrous alloys are sensitive to the welding speed in regard to their susceptibility to solidification cracking. Incompletely melted grains in the partially melted zone are nucleation sites from which weld metal may crystallize. The direction of the consequent epitaxial growth is aligned with the maximum temperature gradients that are perpendicular to the isotherms. Faster welding speeds are accompanied by altered weld pool size and shape and modified microstructural textures. With the increase in welding speed, the weld pool changed from an elliptical shape (Fig. 2a) to a teardrop shape (Fig. 2b) because of the deformed isotherms (change in heat flow) and the increased undercooling delaying the solidification process [6]. Both experimental work [7-10]and simulations [6] correlate the weld pool shape to the grain texture and the teardrop shape to the formation of a continuous centerline grain boundary.

At slow welding speeds, an elliptically shaped pool is generated (Fig. 2a). The associated elliptical shape of the isotherms

Fig. 2 Simulated grain structure of URANUS 2202 welds made by GTAW process using a 3D-coupled cellular automaton-finite element model. Welding speed and power are **a** 1 mm s<sup>-1</sup> and 4500 W, **b** 2 mm s<sup>-1</sup> and 7500 W, and **c** 5 mm s<sup>-1</sup> and 15,000 W. Welding power is adjusted to keep the weld width unchanged [6]

promotes curved macroscopic patterns with a competitive growth favoring grains with Miller indices <100> direction nearly parallel to the maximum temperature gradient [8]. The welding joint contains curved grains from the side to the center due to the continuous curvature of the isotherms. Indeed, a columnar grain which survives over any great distance in an elliptical weld pool exhibits considerable curvature due to the progressive change in the favored growth direction [9]. The crack susceptibility is low due to the tortuous path that would have to follow solidification cracking during its growth.

In case of fast welding speeds, the weld pool elongates with the shape of isotherm lines close to straight lines near the weld pool (Fig. 2b). The teardrop-shaped weld pool has an almost invariant direction of maximum thermal gradient at all points on the weld pool edge from the fusion boundary to the weld center [9]. This thermal field geometry results in any grain favorably oriented for growth at the fusion boundary being able to grow at optimum speed (competitive growth) in the direction of maximum gradient, i.e., perpendicular to the trailing edge [8]. Grains grow straight from side to center and meet at mid-way of the weld width without interacting with adjacent grains. As solidification cracks often form preferentially along the weld centerline [10], the presence of straight, long, and continuous centerline grain boundaries are harmful for crack resistance. Indeed, these long continuous grain boundaries along the weld are a preferential location for macrosegregation of sulfur and phosphorous (two elements that lower the solidus temperature, maintain liquid films at the boundaries, and lengthen the BTR), therefore increasing



solidification cracking susceptibility. The microstructure tends to adopt at higher velocity an orientation normal to the weld direction [6]. It must be noticed that the dendrites are assumed to be aligned perfectly in a direction perpendicular to the growth front. However, in reality, dendrites have preferred crystallographic growth directions and cannot bend in the temperature field. As growth competition occurs, the grains oriented closest to the growth direction are selected to minimize undercooling.

# 2.2 Nickel-based alloys

Nickel-based alloys are widely used in various industry facilities due to their high corrosion resistance and high strength and hardness at elevated temperature. However, these alloys are susceptible to solidification cracking [11–15]. EBSD observations reveal that solidification cracks propagate preferentially along two types of high-angle disoriented grain boundaries: solidification grain boundary (SGB) and solidification sub-grain boundary (SSGB) (Fig. 3). Similar to ferrous alloys, the formation of a centerline grain boundary constitutes a potential weldability issue in regard to solidification crack formation because this last region to solidify is enriched in alloying elements and impurities that lower the solidus and thus lengthen the mushy zone.

Welding speed effect on the formation of centerline grain boundary was investigated experimentally and through modeling for autogenous GTA, IN718 superalloy welds on 2-mm-thick sheets [14]. Modeling the formation of a centerline grain boundary during welding required the coupling of models for heat transfer and dendritic growth in multicomponent alloys. Welding power and speed varied from 300 to 1900 W and from 2.5 to 15 mm s<sup>-1</sup>, respectively. The traditional and most common formulation of heat input (*H* in J mm<sup>-1</sup>) is given by:



**Fig. 3** a EBSD and **b** SEM images of solidification crack in autogenous GTA weld of Ni-28W-6Cr alloy. SSGB and SGB referred to solidification sub-grain boundary and solidification grain boundary, respectively [11]

$$H = \frac{Q}{s} \tag{1}$$

where Q is welding power (in W) and s is welding speed (in mm s<sup>-1</sup>). Accordingly, the weld pool shape and microstructure of the welded joint was mapped for variable effective power as a function of welding speed.

A sharp transition from circular to teardrop-shaped weld pools occurred at a threshold heat input of approximately  $150 \text{ J mm}^{-1}$ . However, the teardrop shape was not sufficient to form centerline grain boundaries. Indeed, the formation of centerline grain boundaries required both a teardrop-shaped weld pool (heat input below 150 J mm<sup>-1</sup>) and an effective power above 900 W [14]. These conclusions are summarized in a power-speed weldability map (Fig. 4) and proved that centerline grain boundaries may be avoided in practice through the choice of a suitable combination of welding power and velocity.

Calculations proved that centerline grain boundaries should not form even with a teardrop-shaped weld pool unless growth undercooling exceed a threshold value [14], in contrast to general understanding. The dominant effect of heat power is demonstrated by associating the formation of centerline grain boundary not to the undercooling distance alone ( $\Delta x_{und}$ ) but to the ratio  $\left(\frac{\Delta x_{und}}{\rho}\right)$ , where  $\Delta x_{und}$  is the distance from the liquidus to the solidification front along the weld centerline and  $\rho$  is the curvature of the liquidus isotherm at the trailing edge of the weld pool [14]. For IN718 superalloy, the relationship between dendrite growth velocity  $v_{kin}$  and undercooling  $\Delta T$  is given by [14]:

$$v_{\rm kin} = 1.28 \cdot 10^{-7} \cdot \Delta T^{3.05} \tag{2}$$

#### 2.3 Aluminum alloys

Aluminum alloys have an ambiguous behavior in welding. Indeed, with faster welding velocities, the susceptibility of



Fig. 4 Calculated weldability map showing conditions for formation of centerline grain boundary in GTA welds of IN718 superalloy [14]. These calculations agreed with experimental data

stray grain formation at the weld centerline increases but is simultaneously in competition with the refinement tendency of the microstructure due to very high undercooling.

The solidification condition in the weld is controlled in part by the thermal gradient (G) and solidification growth rate (R) [4, 16]. Thermal gradients (G) are minimum at the weld centerline and maximum on the weld side because of the large heat extraction along the colder base metal, thus inducing the elongation of the weld pool from circular to teardrop shape at fast welding speeds. Assuming that dendrite solidification velocity corresponds to solidification growth rate due to competitive growth, solidification growth rate (R) is zero at the side interface and maximum at the center, accordingly:

$$R = s \cdot \cos(\alpha) \tag{3}$$

where s is welding speed and  $\alpha$  is angle between solidification growth direction and welding direction.

The columnar-to-equiaxed transition (CET) has been associated to specific forms of thermal gradient-solidification rate relationships such as the solidification parameter  $\frac{G}{R}$  [8]. Small  $\frac{G}{P}$  ratios are associated with high constitutional undercooling ahead of the solid-liquid interface and equiaxed grain structures [17]. Hence, equiaxed solidification tends to occur unaided in the centerline region of the weld pool, where solidification rates (R) are the highest and the thermal gradients (G)are the flattest due to the distance of the arc (Fig. 5), and is generally associated to fast welding speeds and important undercooling [9, 16, 18]. Columnar grain structure was always found predominantly at the weld interface where high  $\frac{G}{R}$  ratios exist. The variation in both local solidification rate and thermal gradient in a single weld on moving around the fusion boundary from the side to the weld center causes a progressive change in solidification substructure across an individual weld bead [9].

The ambiguous behavior of aluminum alloy with travel speed is revealed by the existence of the most unsuitable welding velocity at intermediate values. Indeed, autogenous AA6060 GTA welds were investigated from 2 to 6 mm s<sup>-1</sup> with variation of current from 80 to 145 A to maintain a constant

**Fig. 5** Variation in thermal gradient *G*, solidification growth rate *R*, and grain structure along solidification front in GTA weld pool (top-sectional view) [16]

weld width. Large stray grains were observed above 4 mm s<sup>-1</sup> (Fig. 6), causing a deterioration in solidification cracking resistance. Using the Controlled Tensile Weldability (CTW) test, the critical strain rate for solidification crack formation dropped from -0.06 to -0.20% s<sup>-1</sup> when increasing welding speeds from 2 to 6 mm s<sup>-1</sup> [19].

Faster welding velocities were investigated for autogenous GTA full-penetration welds on 3-mm-thick sheets of 1050, 6082 and 5083 aluminum alloys [16]. The welding speed was varied from 2 to 11.5 mm s<sup>-1</sup>, and the current was adjusted between 170 and 200 A to maintain a constant full-penetrated weld bead size. Solidification cracks were prevented for welding speeds faster than 8 mm  $s^{-1}$  (Fig. 7), this threshold velocity corresponding to the minimum velocity for equiaxed grain formation along the weld centerline. The CET and associated grain refining effect have been proven experimentally [16, 19–27] and numerically [25, 28–32] to reduce efficiently the aluminum weld metal susceptibility to solidification cracking because of the better partitioning of the overall straining conditions [28, 33] and the more tortuous paths to follow during crack growth (in comparison with straight columnar grain interfaces). Efficient grain refiners to avoid cracking in aluminum welds are Ti [21, 23, 24, 34, 35], Sc [21, 35, 36], B [23], TiB<sub>2</sub> [16, 19], Zr [23, 35, 37], and even Mg for Al-Si alloys [38]. Moreover, disturbance of the weld pool, e.g., by electromagnetic [9, 39] or arc-oscillated [40] stirring, avoids straight columnar grain formation, refines the grain structure, and subsequently lowers cracking susceptibility.

## 3 Crack-susceptible zone

The region susceptible to solidification cracking within the mushy zone is referred to as the crack-susceptible zone (CZS) (Fig. 8) and is quantified by its length or by a range of temperature, either brittle temperature range (BTR) [7, 41–45] or solidification cracking temperature range (SCTR) [43]. The CSZ extends usually from the coherency ( $T_c$ ) to the solidus ( $T_s$ ) temperature range. The CZS length is calculated







as the ratio of the temperature range from coherency  $T_c$  to solidus  $T_s$  with the temperature gradient G:

$$CSZ = \frac{T_c - T_s}{G} = \frac{BTR}{G}$$
(4)

Therefore, an increase in welding speed should reduce the gradient G and subsequently elongate the CSZ length. Basic premise consists in associating longer distances (i.e., large CSZ) to more difficult feeding of shrinkage induced by longer and narrower interdendritic liquid channels associated with reduced permeability. Consequently, higher solidification cracking susceptibility is expected at faster welding speeds [12]. As fast travel speeds lengthen the mushy zone, it has been hypothesized that the experimentally observed beneficial travel speed effect observed on cracking susceptibility may

possibly be attributed to the dependency of strain on the heat flow condition [46].

The CSZ is measured usually from the lengths of the cracks form by rapidly applied high strains, such as in Varestraint [47] and Trans-Varestraint tests [48, 49]. Care must be taken in weldability testing based upon crack length as cracks form and continue to grow during the entire time of strain application [44, 50], leading to errors in estimating these temperaturebased indexes if strain is too slowly applied [51]. The length of the crack-susceptible zone is usually measured as the length of the cracks generated during a fast-bending Varestraint test. The Varestraint test consists in bending quickly (within 0.1 s) the weld joint during welding to induce high transverse tensile stresses and subsequently form solidification cracks "as long as the crack-susceptible zone." The length of the cracksusceptible zone is more accurately defined at slow welding



Fig. 7 Weld metal grain structure and solidification crack formation for AA6082 GTA welds at different welding speeds [16]



Fig. 8 Illustrated CSZ length measured in cracked mushy zone after Varestraint test

speed, i.e., when the mushy zone is almost stationary during bending. The welding speed must be accounted for the calculation of the BTR length because the crack grows during the time of bending according to [41]:

$$L_{\rm W} = L_{\rm B} + L_{\rm F} = L_{\rm B} + s \cdot \frac{(\varepsilon_u - 2\varepsilon_{\rm min})}{\dot{\varepsilon}} \tag{5}$$

where  $L_W$  is whole crack length,  $L_F$  is front crack length,  $L_B$  is back crack length (i.e., BTR), *s* is welding speed,  $\varepsilon_{\min}$  is strain required for cracking,  $\varepsilon_u$  is given augmented strain, and  $\varepsilon$  is strain rate. Therefore, it is advised to have fast bending speeds. As the mushy zone continues moving during strain application of 0.4 and 17 mm for travel speeds of 4 mm s<sup>-1</sup> (GTAW) and 167 mm s<sup>-1</sup> (LBW), respectively, the BTR measurements are more accurate for slow welding speed processes. If using fast welding speeds, then a calculation must be performed to convert measured crack lengths in BTR and CSZ lengths. Therefore, some works have focused on getting a better control of strain by standardizing the Trans-Varestraint equipment and methodology [52].

Care must be taken when converting the macroscopic bending strain ( $\varepsilon = \frac{t}{2R}$  with *t* as plate thickness and *R* as radius of curvature of mandrel) to locally applied strain in the CSZ because simulations have demonstrated that strains are strongly localized in the trailing edge of molten pool during bending [42, 44, 50, 53, 54]. The local strains in the mushy zone can reach one order of magnitude greater than the average strain applied during bending [42, 55–58].

One issue in investigating the welding speed effect is to isolate it. Welding speed may be varied while keeping the other parameters constant (i.e., constant welding power Q)

but the weld joint size decreases with faster welding speeds. The power Q may also be increased simultaneously with the speed to maintain the weld size constant (i.e.,  $\frac{Q}{s}$  ratio constant). For autogenous GTA welds of low-alloy steels, isogeometries of weld joints are obtained by varying welding speed *s* (mm s<sup>-1</sup>) and welding current *I* (A) accordingly [59]:

$$I = 38 + 25s \tag{6}$$

Varying welding speed and current according to Eq. 6 showed that the CSZ as measured with the Varestraint test increases at fast welding speeds (Fig. 9). This suggests that increasing welding speed deteriorates solidification crack resistance when maintaining weld bead shape constant [59]. Moreover, alloys with high solidification range show longer CSZ and associated higher solidification cracking susceptibility.

The theoretical calculation of the temperature distribution around a moving heat source has been performed to investigate the elongation of isotherms with welding speed. Many complex models have been proposed to simulate isotherms around a weld pool [60–62] but only the point heat source (i.e., Rosenthal model) is presently used to illustrate weld speed effect. In fact, all models will provide similar qualitative trends but the most complex models should lead to more precise quantitative results. The Rosenthal model is applied in 2D or 3D in correlation with heat flow modes away for the weldment.



**Fig. 9** CSZ lengths measured using Varestraint test at constant bead width by increasing current and welding speed according to the relationship I = 38 + 25s [59]



Fig. 10 Coordinates used for Rosenthal calculations

The Rosenthal formula estimates the temperature distributions around a punctual moving heat source. For a given material, the BTR range is defined as the interval between an upper ( $T_1$ ) and lower ( $T_2$ ) temperatures. Applying the 2D Rosenthal model (Fig. 10) to calculate the length of the CSZ (referred to as the distance x) demonstrates that the length of the CSZ related to the pre-defined BTR depends on the  $\frac{Q}{s}$  ratio:

$$x = \xi_1 - \xi_2$$
(7)  
=  $\pi \alpha \frac{Q^2}{s} \left(\frac{1}{2\pi Kh}\right)^2 \left[\frac{1}{(T_2 - T_0)^2} - \frac{1}{(T_1 - T_0)^2}\right]^2$ 

where  $(\xi_1, \xi_2)$  are traveling coordinates  $(\xi = x - s \cdot t)$ , *t* is time,  $\alpha$  is thermal diffusivity, *K* is thermal conductivity, and *h* is plate thickness.

If considering only that the microstructural tendency for cracking is associated with the liquid channel length to be fed, this equation corresponding to full-penetration welding condition shows that at constant Q, increasing *s* reduces the CSZ and therefore reduces the susceptibility to cracking. However, increasing *s* and maintaining the  $\frac{Q}{s}$  ratio constant implies an increase in  $\frac{Q2}{s}$  and subsequently longer CSZ (Fig. 11) and higher susceptibility to cracking. This calculation agrees with experimental evidence (recall Fig. 9). Applying Eq. 7 specifically to low-alloy steels (properties listed in Table 2) leads to the following relationship plotted in Fig. 11 [59]:

$$CSZ = \frac{0.02(1+12s+36s^2)}{s}$$
(8)

Partial-penetration welds are simulated using the 3D Rosenthal model. Similar calculations applied to the 3D Rosenthal model for a point heat source show that the length of the CSZ does not depend on welding speed *s* [60]:



Fig. 11 CSZ theoretical length as a function of welding speed for lowalloy steel welds according to Eq. 8 and Table 1 [59]

$$x = \xi_1 - \xi_2 = \frac{Q}{2\pi K} \left[ \frac{1}{(T_2 - T_0)} - \frac{1}{(T_1 - T_0)} \right]$$
(9)

Equation 9 demonstrates that, for partial penetration welds, increasing *s* at constant *Q* does not change the CSZ and therefore does not affect the susceptibility to cracking. However, increasing *s* while maintaining the  $\frac{Q}{s}$  ratio constant means that the welding power *Q* is increased leading to longer CSZ and higher susceptibility to cracking.

In summary, the Rosenthal modeling correlates the length of CSZ to Q in full-penetration welds (3D Rosenthal model) and  $\frac{Q2}{s}$  in partial penetration welds (2D Rosenthal model). The time  $t_s$  within the CSZ, i.e., to drop temperature from  $T_1$  to  $T_2$ , is given simply by:

$$t_{\rm s} = \frac{\rm CSZ}{\rm s} \tag{10}$$

The time  $t_s$  is thus proportional to  $\frac{Q}{s}$  (in 3D) and  $\frac{Q2}{s2}$  (in 2D). This time corresponds to the time available for backfilling, i.e., liquid feeding of the opening mushy. Consequently, different solidification cracking behaviors are expected when changing welding conditions because of the associated changes in CSZ lengths and available backfilling time.

## 4 Shifting of stress cells

The non-uniform thermal distribution in and around the weldment during welding generates stress and strain gradients that

Travel speed (mm $s^{-1}$ )	Current (A)	Temperature gradient (°C mm-1)			
		In front of Arc	10 mm behind arc	20 mm behind arc	30 mm behind arc
3.3	105	+ 20	-6	-4	-2
13.3	250	+ 30	-7	- 5	-4

Table 2 Temperature gradients in vicinity of partial penetration, GTA weld of Al-Mg alloy [63]

affect solidification crack formation. A tensile cell favors cracking while a compressive cell hinders cracking. The stress-strain patterns formed during welding result from the non-uniform distribution of heat within the metal, with the net effect being localized deformation near the weld [63, 64]. Compression and tensile stress cells exist around the moving weld pool as the magnitude of thermal stresses is quantified as the stress required to suppress thermal expansion or contraction. These cells are formed by thermally induced dilatation gradients and solidification shrinkage (differential dilatation and contraction of the weld) [65]. They are shifted when varying welding speed thus modifying the nature (compression or tensile) of the cell around the mushy zone itself.

Experimental studies on dynamic plastic strains during GTA welding (Fig. 12) have reported the possibility of a compressive cell at the trailing edge of the weld pool [63, 66]. These cells depend on the temperature distribution around the weld pool. In stationary spot GTA welds, welding heat flow consists of concentric circle isotherms. The heated metal surrounding the stationary heat source is restrained by the cooler metal away from the heat source and therefore is in compression. The stress pattern is uniform. However, in case of a moving heat source, the asymmetric temperature distribution along the longitudinal axis generates compressive stresses of higher magnitude in front than behind the weld pool [63]. When plastic flow occurs, there will be a shifting of metal so as to reduce the unbalanced front-to-rear stress pattern [63].

Chihoski attempted to resolve the stress field state around a moving weld pool [66]. The thermal stress pattern behind the weld pool (Fig. 13) was compressively large due to the plastic flow of material from the front to the back. Further behind the

weld pool, the stress changed from compressive to tensile (new tensile cell  $T_2$ ) as both solidification shrinkage and cooling contraction lead to the development of tensile stresses as temperatures and temperature gradients decrease. It demonstrated that changing travel speed brought changes in the extent and magnitude of the stress fields and the location of transition from compressive to tensile. At slow travel speeds, the tensile cells  $T_s$  and  $T_2$  combine and dominate behind the weld pool. The cracking susceptibility is high. In the opposite case, fast travel speeds tend to reduce compressive cell  $C_1$  and shift the compressive cell  $C_2$  toward the mushy zone, thus reducing the solidification crack susceptibility. The compressive cell  $C_2$  can attain very high magnitudes [63]. These observations have been performed on partial penetration, autogenous GTA aluminum welds at constant current (I = 280 A), and varying welding speed (from 2.4 to 8.4 mm s<sup>-1</sup>), therefore at varying  $\frac{Q}{r}$  ratio [66]. Note that these cells are mostly composed of transverse rather than longitudinal stresses. Thermomechanical modeling of welding using the finite element method has later confirmed contractions of the fusion line in the neighborhood of the weld pool [64, 67]. In case of a constant weld bead size, greater plastic flow regions and thus higher rearward plastic deformations are present behind a weld pool for faster welding speeds [63].

Later, Johnson measured strains using a Moiré-Fringe strain analysis technique around autogenous GTA welds on 3-mm-thick aluminum alloy sheets [63]. While Chihoski focused on transverse and longitudinal normal strains [66], Johnson's results indicate that the straining around the arc consists mostly of shear strains that increase with increasing current and welding speed [63].



Fig. 12 Comparison of thermal effects for a stationary and b moving heat source [63]

Fig. 13 In situ measurement of tensile and compressive cells surrounding the weld pool: **a** fast and **b** slow welding speeds [66]



The planes of maximum shear strain along the weldment side are oriented almost perpendicular and parallel to the weld bead whatever the welding parameters are. As welding speed and current increase, there is a marked increase in the temperature gradients ahead of the weld pool [63]. As welding speed increases (3.3 to 13.3 mm  $s^{-1}$ ) at constant penetration (i.e., isogeometry of cross sections), temperature gradients are greater in the transverse direction [63]. Subsequently, the maximum temperature of a point near the weld will be lower, i.e., greater temperature gradients (Table 2) for faster welding speeds [63]. The restraint effect may shift the stress pattern by limiting the expansion (hence increasing the compressive stresses) [68-70]. Moreover, as welding speed increases, the shear strain is increasingly confined to the immediate vicinity of the weld and shifted toward the weld rearwards [63].

This shift has also been observed by Morgan-Warren and Jordan [59] as shown in the graph of the linear heat content (LHC) calculated as:

$$LHC = JTdy \tag{11}$$

where y is the lateral distance from weld centerline. LHC values are mapped as a function of position ( $\xi$ ) along the weld (Fig. 14). The peaks of the curves indicate the point at which overall cooling and hence contraction of the parent metal begins. At fast travel speeds, the tensile cell is shifted behind the end of the mushy zone, and therefore, no cracking forms. Translating these results into strain and stress, it is found that faster welding speeds reduced the overall level of strain-stress and displaced away from the weld pool rear the tensile part induced by contraction [59]. The CSZ is therefore in a compression cell at fast travel speeds.

Modeling the stress fields behind the weld pool in the Sigmajig testing (Fig. 15) aimed at understanding the influence of dynamic stresses, induced by thermal and mechanical loading, on weld metal solidification crack formation [10]. These calculations are based on a mathematical model for arc welding processes that account for conduction and convection modes for heat transfer, buoyancy, electromagnetic forces, and surface tension gradients [60]. Dynamic mechanical factors are influenced by both restraint and welding parameters. A tensile stress surrounding the mushy zone at the weld trailing edge creates a favorable condition for solidification crack formation as observed during full-penetration autogenous GTA welding



Fig. 14 Longitudinal distribution of LHC and associated compression and tensile force cells for GTA low-alloy steel welds of 3 mm sheets at a slow and b fast welding speeds [59]

**Fig. 15** Calculated stress distribution in Sigmajig test specimen for type 316 SS GTA welds after 4 s of welding and a pre-applied stress of 172 MPa [10]





Fig. 16 Relationship between  $\mathbf{a}$ ,  $\mathbf{b}$  mushy zone extent and  $\mathbf{c}$  compressive and tensile stress fields in case of solidification crack  $\mathbf{a}$  formation or  $\mathbf{b}$  avoidance [10]

of type 316 austenitic stainless steel (Fig. 16). A compressive cell behind the weld pool should hinder the formation of solidification cracking if surrounding entirely the mushy zone [10] but may not prevent it, as solidification cracking can form in highly susceptible microstructures under compressive stresses [19, 28] because the compressive stress may not compensate the entire tension generated during solidification shrinkage on liquid films. Numerical simulations have confirmed the compressive and tensile stress field changeovers along the weldment and their variations with travel speed [10].

## 5 Summary

The effect of travel speed on solidification cracking formation is not straightforward. Indeed, increasing travel speed results in opposing effects. It enhances the formation of solidification cracking by decreasing the centerline temperature gradient G, increasing the CSZ length, decreasing the time to feed shrinkage, and generating centerline grain segregation. On the other hand, increasing travel speed hinders solidification crack formation by shifting the compression cell to the mushy zone, reducing the time exposed to strain, and refining weld metal grains. In summary, the travel speed influences solidification cracking formation through both thermal and metallurgical effects. Interestingly, the Rosenthal modeling highlights that the welding speed affects differently the solidification cracking behavior depending on two welding conditions: is the welding speed s increased at constant welding power Q or constant ratio  $\frac{Q}{s}$ ? Is the weld deposited in partial or full-penetration? Part II of the review will focus on classifying the observed effects on solidification cracking behavior to the experimental welding conditions and metrics.

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