

Sharp Fourier Extension on Fractional Surfaces

Boning Di1,[2](http://orcid.org/0000-0002-2609-4645) · Dunyan Yan[1](http://orcid.org/0000-0001-7462-3823)

Received: 9 July 2023 / Revised: 9 June 2024 / Accepted: 18 June 2024 / Published online: 27 June 2024 © The Author(s), under exclusive licence to Springer Science+Business Media, LLC, part of Springer Nature 2024

Abstract

We investigate a class of Fourier extension operators on fractional surfaces $(\xi, |\xi|^{\alpha})$ with $\alpha \geq 2$. For the corresponding α -Strichartz inequalities, we characterize the precompactness of extremal sequences by applying the missing mass method and bilinear restriction theory. Our result is valid in any dimension. In particular for dimension two, our result implies the existence of extremals for $\alpha \in [2, \alpha_0)$ with some $\alpha_0 > 5$.

Keywords Sharp Fourier restriction theory · Extremals · Fractional Schrödinger equations · Strichartz inequalities

Mathematics Subject Classification Primary 42B10; Secondary 42B37 · 35B38 · 35Q41

1 Introduction

For $\alpha > 2$ and the corresponding α -order free Schrödinger equation, the classical α-*Strichartz inequality* of [\[26](#page-34-0), Theorem 3.1] states the following estimate

$$
\left\|[D^{\frac{\alpha-2}{q}}][e^{it|\nabla|^{\alpha}}]f\right\|_{L_t^q(\mathbb{R})L_x^r(\mathbb{R}^d)} \leq \mathbf{M}_{d,q,\alpha} \|f\|_{L^2(\mathbb{R}^d)},\tag{1}
$$

Communicated by Fulvio Ricci.

 \boxtimes Boning Di diboning18@mails.ucas.ac.cn

> Dunyan Yan ydunyan@ucas.ac.cn

- ¹ School of Mathematical Sciences, University of Chinese Academy of Sciences, Beijing 100049, People's Republic of China
- ² Academy of Mathematics and Systems Science, Chinese Academy of Sciences, Beijing 100190, People's Republic of China

where $2/q + d/r = d/2$ with $q > 2$ and

$$
\mathbf{M}_{d,q,\alpha} := \sup \left\{ \left\| [D^{\frac{(\alpha-2)}{q}}] [e^{it|\nabla|^{\alpha}}] f \right\|_{L_t^q(\mathbb{R})L_x^r(\mathbb{R}^d)} : \|f\|_{L_x^2(\mathbb{R}^d)} = 1 \right\}
$$

is the sharp constant, as well as

$$
[e^{it|\nabla|^{\alpha}}]f(x) := \mathscr{F}^{-1}e^{it|\xi|^{\alpha}}\mathscr{F}f(x) = \frac{1}{(2\pi)^d} \int_{\mathbb{R}^d} e^{ix\xi + it|\xi|^{\alpha}} \widehat{f}(\xi) d\xi,
$$

$$
[D^s]u(x) := \mathscr{F}^{-1}|\xi|^s \mathscr{F}f(x).
$$

Here $\mathscr F$ denotes the spatial Fourier transform

$$
\mathscr{F}f(\xi) := \widehat{f}(\xi) = \int_{\mathbb{R}^d} e^{-ix\xi} f(x) dx, \quad x\xi := x_1\xi_1 + x_2\xi_2 + \cdots + x_d\xi_d
$$

for $x = (x_1, x_2, \ldots, x_d)$ and $\xi = (\xi_1, \xi_2, \ldots, \xi_d)$ in \mathbb{R}^d . Indeed, this α -Strichartz inequality [\(1\)](#page-0-0) belongs to the wider class of *Fourier extension estimates* since the space-time Fourier support of $[e^{it|\nabla|^{\alpha}}]f$ is on the *fractional surface* $(\xi, |\xi|^{\alpha}) \subset \mathbb{R}^{d+1}$. For convenience, we denote that

$$
[E_{\alpha}]f(t, x) := [D^{\frac{\alpha - 2}{q_0}}][e^{it|\nabla|^{\alpha}}]f(x),
$$

$$
q_0 := \frac{2d + 4}{d}, \quad \mathbf{M}_{d, \alpha} := \mathbf{M}_{d, q_0, \alpha}, \quad \mathbf{S}_d^* := \mathbf{M}_{d, 2}.
$$

Note that S_d^* is the corresponding sharp constant for the classical Schrödinger operator $[e^{it\Delta}]$ and the case $\alpha = 2$ is also known as the *Stein-Tomas estimate for paraboloid*.

The relevant *symmetries* for these α -Strichartz inequalities are the space-time translations and scaling as follows

$$
[g_n]f(x) := (h_n)^{d/2} [e^{it_n|\nabla|^{\alpha}}] f(h_n x + x_n), \quad (h_n, x_n, t_n) \in \mathbb{R}_+ \times \mathbb{R}^d \times \mathbb{R};
$$

and the *associated group G* is defined by

$$
G:=\Big\{[g_n]: (h_n,x_n,t_n)\in\mathbb{R}_+\times\mathbb{R}^d\times\mathbb{R}\Big\}.
$$

Then we say a sequence of functions (f_n) in $L^2(\mathbb{R}^d)$ is *precompact up to symmetries* if there exists a sequence of symmetries $([g_n])$ in *G* such that $([g_n]f_n)$ has convergent subsequence in $L^2(\mathbb{R}^d)$. Meanwhile, a sequence of functions (f_n) in $L^2(\mathbb{R}^d)$ is an *extremal sequence* for $\mathbf{M}_{d,q,q}$ if it satisfies

$$
||f_n||_{L^2(\mathbb{R}^d)} = 1, \quad \lim_{n \to \infty} \left\| [D^{\frac{\alpha-2}{q_0}}] [e^{it|\nabla|^{\alpha}}] f_n \right\|_{L^{q_0}_{t,x}(\mathbb{R}^{d+1})} = \mathbf{M}_{d,q,\alpha}.
$$

Birkhäuser

Furthermore, a function $f_* \in L^2(\mathbb{R}^d)$ is called an *extremal* for $\mathbf{M}_{d,q,\alpha}$ if f_* can make the inequality [\(1\)](#page-0-0) an equality and $|| f_* ||_{L^2} = 1$.

The sharp Fourier restriction theory, equivalently the extremal problems for Strichartz type inequalities, has received much attention recently. Readers are referred to the survey [\[19\]](#page-33-0) and the references therein for some progress on this theory, see also the recent survey [\[35](#page-34-1)]. We sketch briefly some of the works as follows.

- **Paraboloid:** Kunze [\[30\]](#page-34-2) showed the existence of extremals for the Fourier extension operator on one-dimension parabola (ξ, ξ^2) $\subset \mathbb{R}^2$ based on an application of *concentration-compactness principle* from Lions [\[32,](#page-34-3) [33](#page-34-4)]; then Foschi [\[17\]](#page-33-1) proved that the only extremals are Gaussians for the case of one-dimension and two-dimension paraboloids by solving some functional equations and investigating some Cauchy-Schwarz inequalities; meanwhile, Hundertmark and Zharnitsky [\[23\]](#page-33-2) established the same result independently by giving a new representation for Strichartz integral based on some orthogonal projection operators; later, Shao [\[39\]](#page-34-5) showed the existence of extremals for the case of arbitrary dimensional paraboloids by applying the *profile decomposition* consequence of [\[3](#page-33-3)].
- **Cone:** for the case of $(\xi, |\xi|) \subset \mathbb{R}^{d+1}$ with low dimensions $d = \{1, 2, 3\}$, the only extremals are known to be exponentials by the works of Foschi [\[17\]](#page-33-1) and Carneiro [\[8\]](#page-33-4); for the case of higher dimensions $d \geq 4$, the extremals exist due to the work of Ramos [\[37\]](#page-34-6).
- **Sphere:** Christ and Shao [\[12](#page-33-5)] showed the existence of extremals for the Fourier extension operator on the two-dimension sphere \mathbb{S}^2 by following the general concentration compactness framework, as well as establishing some strict comparisons for the sharp constants of sphere and paraboloid; then, for this \mathbb{S}^2 case, Foschi [\[18\]](#page-33-6) proved that the only extremals are constants by investigating the Cauchy-Schwarz type estimates for some quadratic forms based on the geometric feature of \mathbb{S}^2 ; later, Shao [\[40\]](#page-34-7) obtained the existence of extremals for the one-dimension sphere \mathbb{S}^1 by combining the outlines in [\[12\]](#page-33-5) and the profile decomposition ideas in [\[2](#page-33-7), [7\]](#page-33-8); then for arbitrary dimensions, Frank, Lieb and Sabin [\[20\]](#page-33-9) established a characterization for the precompactness of extremal sequences by applying the *missing mass method* from Lieb [\[31](#page-34-8)].
- **Other situations:** There are many related works such as the odd curves [\[6,](#page-33-10) [21,](#page-33-11) [38\]](#page-34-9), hyperboloids [\[9,](#page-33-12) [10\]](#page-33-13), perturbations [\[36](#page-34-10)], and non-endpoint type estimates [\[14,](#page-33-14) [15,](#page-33-15) [22\]](#page-33-16), as well as L^p extremals [\[4](#page-33-17), [11](#page-33-18), [16](#page-33-19), [42](#page-34-11)].

The natural generalization of the paraboloid case is to investigate the sharp Fourier extension on fractional surfaces $(\xi, |\xi|^{\alpha}) \subset \mathbb{R}^{d+1}$, which is corresponding to the fractional Schrödinger equations. For the case $(d, \alpha) = (1, 4)$, Jiang, Pausader and Shao [\[24](#page-33-20)] established a dichotomy result on the existence of extremals by establishing the corresponding linear profile decomposition for one-dimension forth order Schrödinger equations; for the case $(d, \alpha) = (d, 4)$, Jiang, Shao and Stovall [\[25\]](#page-34-12) studied the highdimension forth order Schrödinger equations and established a dichotomy result on the existence of extremals; then for the case $(d, \alpha) = (2, 4)$, Oliveira e Silva and Quilodrán [\[36\]](#page-34-10) resolved this dichotomy and obtained the existence of extremals by applying some comparison principle for convolutions of certain singular measures; later, Brocchi, Oliveira e Silva and Quilodrán [\[6](#page-33-10)] established a dichotomy result for the case of

 $(d, \alpha) = (1, \alpha)$ by following some concentration compactness arguments and further obtained the existence of extremals for all $\alpha \in (1, \alpha_0)$ with some $\alpha_0 > 5$ by applying the aforementioned comparison principle; recently, the authors [\[13](#page-33-21)] established same dichotomy result for the case of $(d, \alpha) = (1, \alpha)$ by establishing the corresponding linear profile decomposition for one-dimension fractional Schrödinger equations, and then further studied the asymmetric as well as non-endpoint Strichartz inequalities. For the case of one-dimension fractional curves, both of the proofs in [\[6\]](#page-33-10) and [\[13](#page-33-21)] are based on some refined Strichartz estimates, which follows from the Hausdorff-Young inequality and Whitney decomposition. However, these techniques cannot deal with the higher-dimension fractional surfaces case.

In this article, we investigate the general (d, α) case. One of our main results is the following existence of extremals consequence Theorem [1.1.](#page-3-0) This result generalizes the aforementioned result of $[36,$ $[36,$ Theorem 1.6] which claims the existence of extremals for the case of $(d, \alpha) = (2, 4)$.

Theorem 1.1 *For dimension* $d = 2$ *, there exists one constant* $\alpha_0 > 5$ *such that for arbitrary* $\alpha \in [2, \alpha_0)$ *the extremal for* $\mathbf{M}_{2,\alpha}$ *exists.*

To prove this result, we need the following precompactness Theorem [1.2,](#page-3-1) which gives one characterization for the precompactness of extremal sequences. With this precompactness theorem in place, as we will show later in Sect. [6,](#page-29-0) our Theorem [1.1](#page-3-0) follows directly from the previous results in [\[36](#page-34-10), Proposition 6.9] and the classical fact that Gaussians are extremals for **S**∗ 2.

Theorem 1.2 *All extremal sequences for* **M***d*,α *are precompact up to symmetries if and only if*

$$
\mathbf{M}_{d,\alpha} > (\alpha - 1)^{\frac{-1}{2d+4}} (\alpha/2)^{\frac{-d}{2d+4}} \mathbf{S}_{d}^{*}.
$$
 (2)

In particular, if the strict inequality [\(2\)](#page-3-2) *holds, then there exists an extremal for* **M***d*,α*.*

It is obvious that the Strichartz norm is invariant under the actions of aforementioned symmetries, $\frac{1}{1}$ hence precompact up to symmetries is the best one can expect. Note that Theorem [1.2](#page-3-1) states some universal property for all extremal sequences instead of identifying the extremals. Similar consequences are also established in previous literature, such as [\[20](#page-33-9), Theorem 1.1] for the sphere and [\[21,](#page-33-11) Theorem 1] for cubic curve. As mentioned above, the one-dimension case of Theorem [1.2](#page-3-1) has been proved in the recent works [\[6](#page-33-10), Theorem 1.3] and [\[13,](#page-33-21) Theorem 1.1] by different methods. For the sharp constant \mathbf{S}_d^* , it is conjectured in [\[17](#page-33-1), [23](#page-33-2)] that the only extremals are Gaussians and then the corresponding constant can be obtained by the residue theorem. It can be seen from the asymptotic Schrödinger Lemma [4.1](#page-18-0) that $(\alpha - 1)^{\frac{-1}{2d+4}}(\alpha/2)^{\frac{-d}{2d+4}}S_d^*$ is a

¹ Indeed, the Strichartz estimates are also invariant under some other transformations such as the rotation symmetries $f(x) \mapsto e^{ix_0} f(x)$. However, we will not use these symmetries here since they do not lead to loss of compactness in $L^2(\mathbb{R}^d)$ and are inessential in our situation. Furthermore, for the special case $\alpha = 2$, there are also frequency-translation symmetries $f(x) \mapsto e^{ix\xi_0} f(x)$ which does not maintain the Strichartz norm for general $α$.

lower bound for $\mathbf{M}_{d,\alpha}$. In addition, the strict inequality [\(2\)](#page-3-2) has been proved for some special cases and further details are discussed in Remark [4.2.](#page-19-0)

Now we give some remarks on the proof of Theorem [1.2.](#page-3-1) Recall that the Banach-Alaoglu theorem implies that all extremal sequences must have weak limits up to subsequences. Hence the precompactness of extremal sequences (as well as the existence of extremals) usually comes from two steps: finding a nonzero weak limit and then upgrading this weak convergence to strong convergence. The first step often relies on some refinement of the original α -Strichartz estimates [\(1\)](#page-0-0), and the second step often follows from some compactness arguments. In this paper, we achieve these two steps by using the bilinear restriction estimates and the missing mass method separately.

The fact that Tao's bilinear restriction estimates [\[43\]](#page-34-13) could deduce some refined Strichartz estimates is first shown by Bégout and Vargas for the classical Schrödinger equations in [\[3](#page-33-3)], which generalizes the previous low-dimensional results in [\[26](#page-34-0), [27,](#page-34-14) [34\]](#page-34-15) to higher dimensions. And then numerous consequences are established by following this idea, see, for instance, $[10, 20, 29]$ $[10, 20, 29]$ $[10, 20, 29]$ $[10, 20, 29]$ $[10, 20, 29]$ $[10, 20, 29]$. In fact, it is also mentioned in $[6,$ $[6,$ Sect. 2] and [\[21,](#page-33-11) Sect. 3.3] that the refined Strichartz-type estimates usually come from the bilinear restriction estimates. However, both of them can simply use the Hausdorff-Young inequality instead, since they both only study the one dimensional fractional curves. This Hausdorff-Young inequality is also used by the authors in the recent work [\[13\]](#page-33-21). Inspired by these previous results, in this paper we will use Tao's bilinear restriction estimates to establish the desired high dimensional refined α-Strichartz estimates Proposition [2.7](#page-14-0) for fractional surfaces and then use this result to establish the non-zero weak limit.

However, there are two potential difficulties we should resolve: one is that the fractional surface $(\xi, |\xi|^{\alpha})$ has zero Gaussian curvature at the origin point, which means that we cannot use Tao's bilinear restriction estimate [\[43,](#page-34-13) Sect. 9] directly; another one is that (except for the origin point) the geometric structure of fractional surface is different from that of paraboloid, which implies that we need to investigate the corresponding quasi-orthogonality of these fractional surfaces in order to apply the bilinear-to-linear arguments in [\[45\]](#page-34-17). As shown later in Sect. [2,](#page-5-0) we settle the first difficulty by applying some annular orthogonality consequence Lemma [2.4](#page-12-0) and then restrict our attention to the annular case, which are inspired by [\[10](#page-33-13), [25,](#page-34-12) [28\]](#page-34-18); furthermore, based on one geometric result Proposition [A.1,](#page-30-0) the second difficulty can be solved by dividing the angle into a large number of regions (see the number $K_{d,\alpha}$ in Sect. [2\)](#page-5-0) and then investigating the quasi-orthogonality of each regions.

To upgrade the weak convergence to strong convergence, we use the missing mass method which is invented by Lieb [\[31\]](#page-34-8) in the content of Hardy-Littlewood-Sobolev inequality; see, for instance, [\[20](#page-33-9), [21](#page-33-11)] for the applications of this method in sharp Fourier restriction theory lately. One crucial tool to apply this method is the Brézis–Lieb type lemma due to [\[5,](#page-33-22) [31](#page-34-8)], and here we use a more general version [\[20](#page-33-9), Lemma 3.1]. There are also various other kinds of generalizations appeared in the literature such as [\[6,](#page-33-10) [9,](#page-33-12) [14\]](#page-33-14). For a sequence of functions, when we decompose each function into different parts, these Brézis-Lieb type lemmas can give some limit-orthogonality properties under some suitable conditions. More specifically, the main required condition is a pointwise convergence assumption which in turn relies on the corresponding local smoothing estimates. Furthermore, in high dimensions, there need some multi-variable

analysis such as the multi-variable Taylor's theorem and decay estimates for multivariable oscillatory integrals.

Let us roughly explain how our strategy works. It seems that this also obeys Lions' concentration-compactness principle [\[32](#page-34-3), [33](#page-34-4)]. Note that the loss of compactness in $L^2(\mathbb{R}^d)$ is the main enemy in our situation. Also recall that there are some symmetries when we investigate the precompactness of extremal sequences for $\mathbf{M}_{d,\alpha}$. These three parameters (h_n, t_n, x_n) in the symmetries represent three possible ways to lose compactness in $L^2(\mathbb{R}^d)$: scaling and space-time translations. Fortunately, by using the terminology "up to symmetries" to eliminate the effect of these three parameters, we do not need to worry about the aforementioned three ways of losing compactness. However, there is obvious another way to lose compactness in $L^2(\mathbb{R}^d)$: the frequency translations $f(x) \mapsto e^{ix\xi_n} f(x)$ with parameters $|\xi_n| \to \infty$. Notice that the Strichartz norm changes when frequency translation occurs. Therefore, to establish the precompactness of extremal sequences (up to symmetries), we should at least understand the effect of frequency translations and exclude this type of possibility of losing compactness. Then we are led to the asymptotic Schrödinger behavior Lemma [4.1](#page-18-0) and this is why the strict inequality [\(2\)](#page-3-2) appears in our Theorem [1.2:](#page-3-1) to exclude the loss of compactness deduced by frequency translations. In this sense, our Theorem [1.2](#page-3-1) states one fact that essentially the aforementioned four ways are the only ways of losing compactness when investigating the extremal sequences of $\mathbf{M}_{d,\alpha}$.

The outline of this paper is as follows. In Sect. [2,](#page-5-0) we use the bilinear restriction theory to deduce the refined α -Strichartz estimates as well as the non-zero weak limit consequence; meanwhile, our arguments in this section rely on an auxiliary geometric result whose proof is postponed to the Appendix A. Then in Sect. [3,](#page-15-0) we present some pointwise convergence results so that we are able to apply the Brézis-Lieb type lemma. Next in Sect. [4,](#page-18-1) we study the effect of frequency parameters including the asymptotic Schrödinger behavior and the corresponding pointwise convergence property. In Sect. [5,](#page-25-0) we apply the missing mass method to establish our precompactness Theorem [1.2.](#page-3-1) Finally in Sect. [6,](#page-29-0) we show our existence of extremals result Theorem [1.1.](#page-3-0)

We end this section with some notations. The familiar notation $x \leq y$ denotes that there exists a finite constant *C* such that $|x| \le C|y|$, similarly for $x \ge y$ and $x \sim y$. If necessary, we may use the notation $x \leq_\alpha y$ to show the dependence of this aforementioned constant $C = C_{\alpha} = C(\alpha)$. Finally the indicator function of a set *E* will be denoted by 1_E and we further define $\hat{f}_E := 1_E \hat{f}$.

2 Bilinear Restriction and Refined Strichartz

In this section, we apply the bilinear restriction theory to establish our desired high dimensional refined α -Strichartz estimates, see Proposition [2.7](#page-14-0) below. To achieve this, we need an auxiliary function and a relevant geometric result which is provided in Appendix A. For a vector $\xi = (\xi_1, \xi_2, \cdots, \xi_d) \in \mathbb{R}^d$, we introduce the notations

```
\|\xi\|_{\max} := \max\{|\xi_1|, |\xi_2|, \ldots, |\xi_d|\}, \quad \|\xi\|_{\min} := \min\{|\xi_1|, |\xi_2|, \ldots, |\xi_d|\}.
```
Let *D* denote the family of all the classical dyadic cubes in \mathbb{R}^d . We further define the *annular dyadic cubes* as follows.

Definition 2.1 For *N* ∈ $2^{\mathbb{Z}}$, we define the *dyadic cube-annular* A_N *on Fourier space* as

$$
\mathcal{A}_N := \left\{ \xi \in \mathbb{R}^d : N \leq \|\xi\|_{\max} < 2N \right\}.
$$

Moreover, for $r \in 2^{\mathbb{Z}_+}$, we decompose \mathcal{A}_N into *annular dyadic cubes* which are

$$
\mathcal{D}_{N,r} := \left\{ \tau \in \mathcal{D} : \tau \subset \mathcal{A}_N, \ell(\tau) = \frac{N}{r} \right\}.
$$

To apply some Whitney-type decomposition and achieve some quasi-orthogonal properties, we need the following preliminary definitions and notations.

Definition 2.2 For dyadic cubes $\tau \subset A_N$ and $\tau' \subset A_N$, we write $\tau \sim \tau' \subset A_N$ if $\ell(\tau) = \ell(\tau')$ and they are not adjacent, their parents are not adjacent, their 2-parents are not adjacent,..., their $(N_{d,\alpha}-1)$ -parents are not adjacent while their $N_{d,\alpha}$ -parents are adjacent. Here we say two regions are *adjacent* if their closures intersect, and the number $N_{d,\alpha} \in \mathbb{Z}_+$ will be determined later in [\(3\)](#page-6-0).

Notice that $\tau \sim \tau'$ will imply $\ell(\tau) = N/r$ with some $r \geq 2^{N_{d,\alpha}}$. To use the bilinear form and the Whitney-type decomposition, we divide the unit sphere \mathbb{S}^{d-1} into $K_{d,\alpha}$ parts which will lead to the following angle decomposition

$$
\mathcal{A}_N = \bigcup_{j=1}^{K_{d,\alpha}} \mathcal{A}_N^j, \quad \theta_N^j := \sup \left\{ \theta(\xi, \eta) : (\xi, \eta) \in \mathcal{A}_N^j \times \mathcal{A}_N^j \right\},
$$

such that $\theta_N^j < \theta_{d,\alpha}$ with the angle range $\theta_{d,\alpha}$ depending on $K_{d,\alpha}$. Here $\theta(\xi, \eta)$ denotes the angle between the vectors ξ and η . By symmetry we may assume $|\xi| \geq |\eta|$, meanwhile we can set $K_{d,\alpha}$ large enough to let the angle region $\theta_{d,\alpha}$ as small as we want. Here the numbers $N_{d,\alpha}$ and $K_{d,\alpha}$ are chosen based on the angle decomposition result Proposition [A.1.](#page-30-0) Indeed, using the notations in Remark [A.2,](#page-32-0) our numbers $N_{d,\alpha}$ and $K_{d,\alpha}$ are chosen such that

$$
\arctan(d^{-1/2}2^{-N_{d,\alpha}}) + \theta_{d,\alpha} \le \bar{\theta}_0. \tag{3}
$$

Since the angle θ_0 is fixed and depends only on (d, α) , we know that this condition [\(3\)](#page-6-0) must can be achieved as long as $N_{d,\alpha}$ large enough and $\theta_{d,\alpha}$ small enough which means $K_{d,\alpha}$ large enough.

The reason we construct this condition (3) is that it will imply the result (11) which is a critical estimate in the proof of Quasi-orthogonality Lemma [2.3.](#page-7-0) Meanwhile it should be pointed out that

$$
\max_{1 \le j \le K_{d,\alpha}} \left\| [e^{it|\nabla|^{\alpha}}] f_{\mathcal{A}_{N}^{j}} \right\|_{L_{t,x}^{q_0}} \le \left\| [e^{it|\nabla|^{\alpha}}] f_{\mathcal{A}_{N}} \right\|_{L_{t,x}^{q_0}} \le \sum_{j=1}^{K_{d,\alpha}} \left\| [e^{it|\nabla|^{\alpha}}] f_{\mathcal{A}_{N}^{j}} \right\|_{L_{t,x}^{q_0}}. (4)
$$

We further introduce the *annular-restricted notation* $\tau \sim \tau' \subset A_N^j$ which means

$$
\tau \sim \tau' \subset \mathcal{A}_N, \quad \tau \cap \mathcal{A}_N^j \neq \emptyset, \quad \tau' \cap \mathcal{A}_N^j \neq \emptyset.
$$

Recall that θ_0 is very small and much smaller than $\pi/8$. Hence by dividing \mathcal{A}_N into *K_{d*, α} parts as above, except for a null set, it can be achieved that for arbitrary (ξ , ξ') \in $A_N^j \times A_N^j$ there exists unique pair $\tau \sim \tau' \subset A_N^j$ satisfying $\xi \in \tau$ and $\xi' \in \tau'$. This fact deduces the Whitney-type decomposition which will be used later in the proof of annular refined estimates Lemma [2.6.](#page-12-1)

In our bilinear setting, the first crucial fact is the following quasi-orthogonality lemma.

Lemma 2.3 *Suppose that* $d \geq 2$ *and* $j_0 \in \{1, 2, ..., K_{d,\alpha}\}$ *. Then the following inequality holds*

$$
\left\|\sum_{\tau\sim\tau'\subset\mathcal{A}_1^{j_0}}[e^{it|\nabla|^{\alpha}}]f_\tau\cdot[e^{it|\nabla|^{\alpha}}]f_{\tau'}\right\|_{L^{\frac{q_0}{2}}_{t,x}}^{\frac{q_0}{2}}\lesssim \sum_{\tau\sim\tau'\subset\mathcal{A}_1^{j_0}}\left\|[e^{it|\nabla|^{\alpha}}]f_\tau\cdot[e^{it|\nabla|^{\alpha}}]f_{\tau'}\right\|_{L^{\frac{q_0}{2}}_{t,x}}^{\frac{q_0}{2}},
$$

for all $f \in L^2(\mathbb{R}^d)$ *satisfying* supp $\widehat{f} \subset A_1^{j_0}$.

Proof of Lemma **[2.3](#page-7-0)** At the beginning, we mention that all the dyadic cubes in this proof have intersections with $A_1^{j_0}$. Given $\tau \in \mathcal{D}_{1,r}$, denote $\xi_0 := c(\tau)$ the center of τ . Then for every $\xi \in \tau$, considering the radial direction and angle differences, some geometry observations give that

$$
\left| |\xi| - |\xi_0| \right| \lesssim 1/r, \quad (|\xi| |\xi_0| - \xi \xi_0)^{1/2} \lesssim 1/r. \tag{5}
$$

Furthermore for $\tau \sim \tau' \in \mathcal{D}_{1,r}$ and $\xi'_0 := c(\tau')$, there holds

$$
\left| |\xi_0| - |\xi'_0| \right| + (|\xi_0| |\xi'_0| - \xi_0 \xi'_0)^{1/2} \sim 1/r. \tag{6}
$$

Denote $\tilde{\tau}$, $\tilde{\tau}'$ the lifts of τ , τ' into the surface $(\xi, |\xi|^{\alpha}) \subset \mathbb{R}^{d+1}$. Based on the quasiorthogonality result [\[37](#page-34-6), Lemma 2.2], our main task is investigating the geometry of the sum-set

$$
\tilde{\tau} + \tilde{\tau}' := \left\{ (\xi + \xi', |\xi|^{\alpha} + |\xi'|^{\alpha}) : (\xi, \xi') \in \tau \times \tau' \right\} \subset \mathbb{R}^{d+1}.
$$

B Birkhäuser

We claim that the collection $\{(\tau, \tau') : \tau \sim \tau' \subset A_1^{j_0}\}\)$ can be finitely decomposed into universal number of subsets $\{\mathcal{I}_m\}_{m=1}^{K_0}$, and then for each subset \mathcal{I}_m there exist a universal number K_1 and a universal constant $\beta > 0$ such that the following property holds: for every $(\tau, \tau') \in \mathcal{T}_m$, we can find K_1 parallelepipeds $\{P_\ell = P_\ell(\tau, \tau')\}_{\ell=1}^{K_1}$ $\ell=1$ satisfying

$$
(\tilde{\tau} + \tilde{\tau}') \subset \left(\bigcup_{\ell=1}^{K_1} P_{\ell}\right), \qquad \left[(1+\beta) \cdot P_{\ell}\right] \cap \left[(1+\beta) \cdot P_{\ell'}\right] = \emptyset \text{ for } \ell \neq \ell'.
$$

Here $(1 + \beta) \cdot P_\ell$ denotes the centered dilation of P_ℓ which means

$$
(1 + \beta) \cdot P_{\ell} := (1 + \beta) [P_{\ell} - c(P_{\ell})] + c(P_{\ell})
$$

with $c(P_\ell)$ denoting the center of P_ℓ . Let us postpone the detailed proof of this claim and use it to prove our final result now. Define $\Psi_{\ell} := 1_{P_{\ell}}$. Since P_{ℓ} is a parallelepiped, for every exponent $q > 1$, the boundedness of Hilbert transform implies the following multiplier boundedness

$$
||F * \widetilde{\Psi}_{\ell}||_{L^{q}(\mathbb{R}^{d+1})} \lesssim ||F||_{L^{q}(\mathbb{R}^{d+1})}
$$
\n(7)

for all $F \in L^q(\mathbb{R}^{d+1})$. On the other hand, note that the space-time Fourier support satisfies

$$
\operatorname{supp}\Big([e^{it|\nabla|^{\alpha}}]f_{\tau}[e^{it|\nabla|^{\alpha}}]f_{\tau'} \Big)^{\bigwedge_{t,x}} \subset (\tilde{\tau} + \tilde{\tau}') .
$$

Hence by the claim and triangle inequality, we only need to show the following estimate holds

$$
\left\| \sum_{(\tau,\tau') \in \mathcal{J}_m} \left([e^{it|\nabla|^{\alpha}}] f_{\tau} [e^{it|\nabla|^{\alpha}}] f_{\tau'} \right) \ast \widehat{\Psi}_{\ell} \right\|_{L_{t,x}^{\frac{q_0}{2}}}^{\frac{q_0}{2}} \right\|_{L_{t,x}^{\frac{q_0}{2}}} \lesssim \sum_{(\tau,\tau') \in \mathcal{J}_m} \left\| [e^{it|\nabla|^{\alpha}}] f_{\tau} [e^{it|\nabla|^{\alpha}}] f_{\tau'} \right\|_{L_{t,x}^{\frac{q_0}{2}}}^{\frac{q_0}{2}} \tag{8}
$$

for each $m \in \{1, 2, \ldots, K_0\}$ and each $\ell \in \{1, 2, \ldots, K_1\}$. This estimate [\(8\)](#page-8-0) comes from a direct application of [\[37](#page-34-6), Lemma 2.2] and the boundedness [\(7\)](#page-8-1). Indeed, since *P*_l is an affine image of the unit cube, for each $(\tau, \tau') \in \mathcal{T}_m$ we can construct a bump function $\varphi = \varphi(\tau, \tau')$ satisfying

$$
\text{supp}\varphi \subset (1+\beta)\cdot P_\ell, \quad \varphi(x) \equiv 1 \text{ for } x \in P_\ell, \quad \|\widehat{\varphi}\|_{L^1(\mathbb{R}^{d+1})} \leq C,
$$

Birkhäuser

$$
\left\| \sum_{(\tau,\tau') \in \mathcal{I}_m} \left([e^{it|\nabla|^{\alpha}}] f_{\tau} [e^{it|\nabla|^{\alpha}}] f_{\tau'} \right) \ast \widehat{\Psi}_{\ell} \right\|_{L_{t,x}^{\frac{q_0}{2}}}^{\frac{q_0}{2}} \leq \sum_{(\tau,\tau') \in \mathcal{I}_m} \left\| \left([e^{it|\nabla|^{\alpha}}] f_{\tau} [e^{it|\nabla|^{\alpha}}] f_{\tau'} \right) \ast \widehat{\Psi}_{\ell} \right\|_{L_{t,x}^{\frac{q_0}{2}}}^{\frac{q_0}{2}}.
$$

This conclusion and the boundedness [\(7\)](#page-8-1) immediately give the desired estimate [\(8\)](#page-8-0).

It remains to prove the aforementioned claim. For the sum-set $\tau + \tau'$, considering the length along radial direction and the angle differences, it is not hard to see that

$$
\left| |\xi + \xi'| - |\xi_0 + \xi'_0| \right| \lesssim 1/r, \quad \left[|\xi + \xi'| |\xi_0 + \xi'_0| - (\xi + \xi')(\xi_0 + \xi'_0) \right]^{1/2} \lesssim 1/r. \tag{9}
$$

Meanwhile some further investigation will imply the following estimate

$$
|\xi|^{\alpha} + |\xi'|^{\alpha} - |\xi + \xi'|^{\alpha} / 2^{\alpha - 1} \sim 1/r^2.
$$
 (10)

Indeed, without loss of generality we may assume $|\xi| \geq |\xi'|$. If we define the auxiliary function

$$
F_{\xi'}(\xi) := |\xi|^{\alpha} + |\xi'|^{\alpha} - |\xi + \xi'|^{\alpha}/2^{\alpha - 1},
$$

then for every fixed $\xi' \in \mathbb{R}^d$ the function $F_{\xi'}(\xi)$ has non-degenerate critical point $\xi = \xi'$. In other words, the Hessian matrix Hess $F_{\xi'}(\xi')$ is positive-definite and

$$
\nabla F_{\xi'}(\xi')=0.
$$

Also notice that $F_{\xi'}(\xi') = 0$ at this point and $F_{\xi'}(\xi)$ is a smooth function. Moreover, since we have

$$
(\xi, \xi') \in \tau \times \tau', \quad \tau \sim \tau' \subset \mathcal{A}_1^{j_0},
$$

the condition (3) and Proposition [A.1](#page-30-0) then imply

$$
\det \text{Hess} F_{\xi'}(\xi) \sim 1. \tag{11}
$$

Hence when $\tau \sim \tau' \in \mathcal{D}_{1,r}$ which leads to $|\xi - \xi'| \sim 1/r$, we can use the multivariable Taylor's theorem to obtain $|F_{\xi'}(\xi)| \sim 1/r^2$. This gives the desired estimate [\(10\)](#page-9-1).

Then we follow some similar arguments in the proof of $[10, \text{Lemma 5.2}]$ $[10, \text{Lemma 5.2}]$. First, we can find a universal constant C_1 such that for every fixed $\tau \sim \tau' \in \mathcal{D}_{1,r}$, the number of corresponding $\rho \sim \rho' \in \mathcal{D}_{1,s}$ satisfies the following quantity bound

$$
\#\big\{(\rho,\rho'):\rho\sim\rho'\in\mathcal{D}_{1,s},\;(\tilde{\tau}+\tilde{\tau}')\cap(\tilde{\rho}+\tilde{\rho}')\neq\emptyset\big\}\leq C_1.
$$

In fact, if $(\tilde{\tau} + \tilde{\tau}') \cap (\tilde{\rho} + \tilde{\rho}') \neq \emptyset$, it is not hard to see that the estimate [\(10\)](#page-9-1) implies *s* ∼ *r*; and for each fixed *s*, the estimates [\(5\)](#page-7-1), [\(6\)](#page-7-2) and [\(9\)](#page-9-2) can imply that for given τ , the number of possible cubes $\rho \in \mathcal{D}_{1,s}$ is universally bounded; then for each fixed ρ , the number of ρ' is obviously uniformly finite. Therefore we obtain the existence of the constant *C*1.

Second, note that [\(10\)](#page-9-1) gives the universal constant $c_2 > 0$ and $c_3 > 0$ with the following relation

$$
\tilde{\tau} + \tilde{\tau}' \subset \Gamma_{\tau,\tau'},
$$

\n
$$
\Gamma_{\tau,\tau'} := \left\{ (\xi,\eta) \in (\tau + \tau') \times \mathbb{R} : |\xi|^{\alpha}/2^{\alpha-1} + c_2/r^2 \leq \eta \leq |\xi|^{\alpha}/2^{\alpha-1} + c_3/r^2 \right\}.
$$

Meanwhile the estimates [\(9\)](#page-9-2) give the rectangle $R_{\tau,\tau'} \subset \mathbb{R}^d$ such that $\tau + \tau' \subset R_{\tau,\tau'}$ with center $c(R_{\tau,\tau'})$ at the point $\gamma_0 := \xi_0 + \xi'_0$ and every side length comparable to 1/*r*, as well as one edge aligned with the vector γ₀. By a centered dilation $R^*_{τ,τ'}$:= $(1 + c_4) \cdot R_{\tau, \tau'}$ with constant $c_4 > 0$ small enough independent of (τ, τ') , we can let the following sets still have bounded overlap

$$
\Sigma_{\tau,\tau'} := \left\{ (\xi,\eta) \in R^*_{\tau,\tau'} \times \mathbb{R} : |\xi|^{\alpha}/2^{\alpha-1} + c_2/2r^2 \leq \eta \leq |\xi|^{\alpha}/2^{\alpha-1} + 2c_3/r^2 \right\}.
$$

Hence we are able to decompose the collection $\{(\tau, \tau') : \tau \sim \tau' \subset A_1^{j_0}\}\$ into universal number of subsets $\{\mathcal{I}_m\}_{m=1}^{K_0}$, such that for each \mathcal{I}_m there is a corresponding set $\{\Sigma_{\tau,\tau'}\}_{m}$ whose elements are pairwise disjoint. Thus, we may fix some $m = m_0$ and investigate one subset \mathcal{T}_{m_0} from now on.

Third, for $\gamma \in \mathbb{R}^d$, we define $T(\gamma)$ to be the tangent plane of the surface $(\xi, |\xi|^{\alpha}/2^{\alpha-1})$ at the point $(\gamma, |\gamma|^{\alpha}/2^{\alpha-1})$ as follows

$$
T(\gamma):=\left\{(\gamma,|\gamma|^{\alpha}/2^{\alpha-1})+v: v\in \mathbb{R}^{d+1},\ v\bot(\alpha|\gamma|^{\alpha-2}\gamma/2^{\alpha-1},-1)\right\}.
$$

Let $(e_1, e_2, \ldots, e_{d+1})$ denote the canonical basis in \mathbb{R}^{d+1} . Without loss of generality, we may assume the center $\gamma_0 = |\gamma_0| e_d$. Consider the point $\gamma = ke_d$ which is very close to γ_0 . In this case, the corresponding hyperplane can be computed as

$$
T(\gamma) = \left\{ (\gamma, |\gamma|^{\alpha}/2^{\alpha-1}) + (v_1, v_2, \dots, v_d, v_d \alpha k^{\alpha-1}/2^{\alpha-1}) : v_i \in \mathbb{R}, i = 1, 2, \dots, d \right\};
$$

and lifting the rectangle $R_{\tau,\tau'}$ to the tangent plane $T(\gamma)$ amounts to choosing

$$
|(v_1,v_2,\ldots,v_{d-1})|\lesssim 1/r,\quad v_d\lesssim 1/r.
$$

Hence we can precisely set $y = (v_1, v_2, \ldots, v_{d-1})$ and assume

$$
|y| \le c_5/r, \quad |v_d| \le c_6/r \tag{12}
$$

with the constants c_5 , c_6 to be determined later. Under these assumptions, in the direction e_{d+1} , we can use the multi-variable Taylor's theorem to estimate the largest displacement between the tangent plane and the surface $(\xi, |\xi|^{\alpha}/2^{\alpha-1})$ as follows

$$
|(y, k + v_d)|^{\alpha}/2^{\alpha} - \left(k^{\alpha}/2^{\alpha-1} + \alpha v_d k^{\alpha-1}/2^{\alpha-1}\right) \lesssim |(y, v_d)|^2 \leq (c_5^2 + c_6^2)/r^2.
$$

Here we have used the fact *k* ∼ 1 which leads to all the second-order partial derivatives and the Hessian of this corresponding function $g(\xi) = |\xi|^\alpha/2^{\alpha-1}$ comparable to 1. Hence we can choose some universal constants c_5 , c_6 small enough such that this displacement is less than $\frac{c_3}{2r^2}$. This conclusion further gives the desired universal constant *K*₁ comparable to $(c_5)^{1-d}$ $(c_6)^{-1}$, such that the rectangle $R_{\tau,\tau'}$ can be decomposed into *K*¹ smaller rectangles

$$
R_{\tau,\tau'}=\bigcup_{\ell=1}^{K_1}R_\ell,
$$

where the smaller rectangles $R_{\ell} = R_{\ell}(\tau, \tau')$ have the same size with disjoint interiors and satisfy the condition [\(12\)](#page-11-0).

Then for each ℓ , let c_{ℓ} be the center of the rectangle P_{ℓ} and let $T(R_{\ell})$ denote the lift of R_ℓ into the tangent plane $T(c_\ell)$. Define the parallelepiped $P_\ell = P_\ell(\tau, \tau') \subset \mathbb{R}^{d+1}$ as the sum-set

$$
P_{\ell} := T(R_{\ell}) + \left\{ s e_{d+1} : \frac{c_2}{r^2} \le s \le \frac{3c_3}{2r^2} \right\}.
$$

Note that distinct P_{ℓ} have disjoint interiors and the following inclusion holds

$$
(\tilde{\tau}+\tilde{\tau}')\subset\Gamma_{\tau,\tau'}\subset\left(\bigcup_{\ell=1}^{K_1}P_{\ell}\right).
$$

Moreover, by the construction of $R^*_{\tau,\tau'}$ and R_ℓ , there exists $\beta > 0$ such that $(1 + \beta) \cdot R_\ell \subset R^*_{\tau, \tau'}$ holds for every $\ell \in \{1, 2, ..., K_1\}$. Taking the aforementioned displacements into consideration, recalling the choice of constants c_5 and c_6 , we can guarantee that

$$
(1+\beta)\cdot P_{\ell}\subset \Sigma_{\tau,\tau'},
$$

possibly choosing a smaller β but still universal. This finishes the proof of our desired claim. \Box

B Birkhäuser

One basic refinement of the original α -Strichartz inequality [\(1\)](#page-0-0) is the following annular orthogonality Lemma [2.4.](#page-12-0) Indeed, a direct application of the Littlewood-Paley theory and the mixed-norm estimates [\(1\)](#page-0-0) will give this refinement. Similar arguments can be found in $[25, \text{Lemma } 2.3]$ $[25, \text{Lemma } 2.3]$. See also $[10, \text{Proposition } 2.1]$ $[10, \text{Proposition } 2.1]$ and $[28, \text{Lemma } 4.1]$ $[28, \text{Lemma } 4.1]$. We omit the proof here for simplicity.

Lemma 2.4 *Let* $d \geq 2$ *. We have the following annular orthogonality estimate*

$$
\| [E_{\alpha}] f \|_{L_{t,x}^{q_0}(\mathbb{R}^{d+1})}^{q_0} \lesssim \sup_{N \in 2^{\mathbb{Z}}} \| [E_{\alpha}] f_N \|_{L_{t,x}^{q_0}(\mathbb{R}^{d+1})}^{q_0-2} \| f \|_{L_x^2(\mathbb{R}^d)}^2
$$

for every $f \in L^2(\mathbb{R}^d)$ *, where* $\widehat{f}_N := \mathbb{1}_{A_N} \widehat{f}$.

Based on this annular orthogonality result, we are able to focus on the annular case and aim to establish some suitable control for the item $\left\| [E_{\alpha}] f_N \right\|_{L_{t,x}^{q_0}}$. At this point, the advantage is that the truncated surface has non-zero Gaussian curvature. Hence we are able to apply Tao's classical bilinear restriction estimates in [\[43](#page-34-13)].

Theorem 2.5 *Suppose that* $\frac{d+3}{d+1} < p < \frac{d+2}{d}$ *and* $j_0 \in \{1, 2, ..., K_{d,\alpha}\}$ *. Then for every* $\tau \sim \tau' \subset A_1^{j_0}$, the following bilinear estimates holds

$$
\left\| [e^{it|\nabla|^{\alpha}}] f_{\tau} \cdot [e^{it|\nabla|^{\alpha}}] f_{\tau'} \right\|_{L^p_{t,x}} \lesssim |\tau|^{1-\frac{q_0}{2p}} \|f_{\tau}\|_{L^2} \|f_{\tau'}\|_{L^2}
$$

for all $f \in L^2(\mathbb{R}^d)$ *satisfying* supp $\widehat{f} \subset A_1$ *. Therefore, by interpolation there exists* $s_0 \in (1, 2)$ *such that*

$$
\left\| [e^{it|\nabla|^{\alpha}}] f_{\tau} \cdot [e^{it|\nabla|^{\alpha}}] f_{\tau'} \right\|_{L_{t,x}^{q_0/2}} \lesssim |\tau|^{1-\frac{2}{s_0}} \|\widehat{f}_{\tau}\|_{L^{s_0}} \|\widehat{f}_{\tau'}\|_{L^{s_0}}.
$$

Proof of Theorem [2.5](#page-12-2) This result follows from a standard parabolic rescaling argument by Tao's bilinear estimates on the paraboloid [\[43\]](#page-34-13). Note that our surface $(\xi, |\xi|^{\alpha})$ restricted on the unit cube-annular A_1 is a elliptic-type compact surface in the sense of [\[45](#page-34-17), Sect. 2]. Hence we can use the bilinear estimates [\[43,](#page-34-13) Sect. 9, third remark] to deduce this desired conclusion. See [\[20](#page-33-9), Theorem A.1] for further details on this parabolic rescaling argument.

As mentioned above, on each annular, we are going to apply this bilinear estimate Theorem [2.5](#page-12-2) to deduce the following annular refined α -Strichartz estimate Lemma [2.6.](#page-12-1) The arguments (from bilinear-type estimate to refined-type Strichartz estimate) are classical, and similar arguments can be seen in the works such as [\[3](#page-33-3), [21](#page-33-11), [44\]](#page-34-19). Here we follow the ideas in these works, and show the details for the convenience of the readers.

Lemma 2.6 *Let* $d \geq 2$ *. There exists* $\gamma \in (0, 1)$ *such that the following estimate*

$$
\left\| [e^{it|\nabla|^{\alpha}}]f \right\|_{L^{q_0}_{t,x}} \lesssim \left[\sup_{r \in 2^{\mathbb{Z}_+}} \sup_{\tau \in \mathcal{D}_{1,r}} |\tau|^{-1/2} \left\| [e^{it|\nabla|^{\alpha}}] f_{\tau} \right\|_{L^{\infty}_{t,x}} \right]^{\gamma} \|f\|^{1-\gamma}_{L^{2}_{x}}
$$

.

holds for all $f \in L^2(\mathbb{R}^d)$ *satisfying* supp $\widehat{f} \subset A_1$ *. Moreover, by rescaling, the following estimate*

$$
N^{\frac{\alpha-2}{q_0}}\left\| [e^{it|\nabla|^{\alpha}}]f \right\|_{L^{q_0}_{t,x}} \lesssim \left[\sup_{r \in 2^{\mathbb{Z}_+}} \sup_{\tau \in \mathcal{D}_{N,r}} |\tau|^{-1/2} \left\| [e^{it|\nabla|^{\alpha}}] f_{\tau} \right\|_{L^{\infty}_{t,x}} \right]^{\gamma} \|f\|_{L^{2}_{x}}^{1-\gamma}
$$

holds for all $f \in L^2(\mathbb{R}^d)$ *satisfying* supp $\widehat{f} \subset A_N$ *.*

Proof of Lemma [2.6](#page-12-1) Due to the relation [\(4\)](#page-7-3), it suffices to fix some $j_0 \in \{1, 2, ..., K_{d, \alpha}\}\$ and prove that

$$
\left\| [e^{it|\nabla|^{\alpha}}] f_{\mathcal{A}_1^{j_0}} \right\|_{L_{t,x}^{q_0}} \lesssim \left[\sup_{r \in 2^{\mathbb{Z}_+}} \sup_{\tau \in \mathcal{D}_{1,r}} |\tau|^{-1/2} \left\| [e^{it|\nabla|^{\alpha}}] f_{\tau} \right\|_{L_{t,x}^{\infty}} \right]^{\gamma} \|f\|_{L_x^{2}}^{1-\gamma}.
$$
 (13)

Our strategy is using the bilinear estimates Theorem [2.5.](#page-12-2) Thus, the Whitney-type decomposition gives

$$
\left\| [e^{it|\nabla|^{\alpha}}] f_{\mathcal{A}_1^{j_0}} \right\|_{L^{q_0}_{t,x}} \le \left\| \sum_{\tau \sim \tau' \subset \mathcal{A}_1^{j_0}} [e^{it|\nabla|^{\alpha}}] f_{\tau} \cdot [e^{it|\nabla|^{\alpha}}] f_{\tau'} \right\|_{L^{q_0/2}_{t,x}}^{1/2}
$$

Then the quasi-orthogonality Lemma [2.3](#page-7-0) and Hölder's inequality imply that

$$
\left\| \sum_{\tau \sim \tau' \subset \mathcal{A}_1^{j_0}} [e^{it|\nabla|^{\alpha}}] f_{\tau} \cdot [e^{it|\nabla|^{\alpha}}] f_{\tau'} \right\|_{L_{t,x}^{q_0/2}}^{1/2}
$$
\n
$$
\lesssim \left(\sum_{\tau \sim \tau' \subset \mathcal{A}_1^{j_0}} \| [e^{it|\nabla|^{\alpha}}] f_{\tau} \cdot [e^{it|\nabla|^{\alpha}}] f_{\tau'} \Big|_{L_{t,x}^{q_0/2}}^{q_0} \right)^{\frac{1}{2q_0}}
$$
\n
$$
\leq \sup_{\tau \sim \tau' \subset \mathcal{A}_1^{j_0}} \| [e^{it|\nabla|^{\alpha}}] f_{\tau} \cdot [e^{it|\nabla|^{\alpha}}] f_{\tau'} \Big|_{L_{t,x}^{q_0/2}}^{q_0-s}
$$
\n(14)

$$
\times \left(\sum_{\tau \sim \tau' \subset \mathcal{A}_1^{j_0}} \left\| [e^{it|\nabla|^{\alpha}}] f_{\tau} \cdot [e^{it|\nabla|^{\alpha}}] f_{\tau'} \right\|_{L_{t,x}^{q_0/2}}^s \right)^{\frac{1}{2q_0}}.
$$
 (15)

Here the inequality is valid for all $0 < s < q_0$, though we will choose $1 < s < q_0$ later. In the remainder of this proof, we will estimate (14) and (15) by different ways which will finally lead to the desired result (13) .

Firstly for the term [\(14\)](#page-13-0), we have that

$$
\| [e^{it|\nabla|^{\alpha}}] f_{\tau} \cdot [e^{it|\nabla|^{\alpha}}] f_{\tau'} \|_{L_{t,x}^{q_0/2}} \leq \left(|\tau|^{-1} \left\| [e^{it|\nabla|^{\alpha}}] f_{\tau} \cdot [e^{it|\nabla|^{\alpha}}] f_{\tau'} \right\|_{L_{t,x}^{\infty}} \right)^{1-\frac{2p}{q_0}} \times \left(|\tau|^{\frac{q_0}{2p}-1} \left\| [e^{it|\nabla|^{\alpha}}] f_{\tau} \cdot [e^{it|\nabla|^{\alpha}}] f_{\tau'} \right\|_{L_{t,x}^p} \right)^{\frac{2p}{q_0}},
$$

where $\frac{d+3}{d+1} < p < \frac{d+2}{d}$. Then the bilinear estimates Theorem [2.5](#page-12-2) gives the following estimate

$$
\| [e^{it|\nabla|^{\alpha}}] f_{\tau} \cdot [e^{it|\nabla|^{\alpha}}] f_{\tau'} \|_{L_{t,x}^{q_0/2}} \lesssim \left(|\tau|^{-1/2} \| [e^{it|\nabla|^{\alpha}}] f_{\tau} \|_{L_{t,x}^{\infty}} \right)^{1 - \frac{2p}{q_0}} \times \left(|\tau'|^{-1/2} \| [e^{it|\nabla|^{\alpha}}] f_{\tau'} \|_{L_{t,x}^{\infty}} \right)^{1 - \frac{2p}{q_0}} \| f \|_{L_x^2}^{\frac{4p}{q_0}},
$$

which immediately means

$$
\sup_{\tau \sim \tau' \subset \mathcal{A}_1^{j_0}} \left\| [e^{it|\nabla|^{\alpha}}] f_{\tau} \cdot [e^{it|\nabla|^{\alpha}}] f_{\tau'} \right\|_{L_{t,x}^{q_0/2}} \n\lesssim \left[\sup_{r \in 2^{\mathbb{Z}_+}} \sup_{\tau \in \mathcal{D}_{1,r}} |\tau|^{-1/2} \| [e^{it|\nabla|^{\alpha}}] f_{\tau} \|_{L_{t,x}^{\infty}} \right]^{2 - \frac{4p}{q_0}} \| f \|_{L_x^2}^{\frac{4p}{q_0}}.
$$
\n(16)

For the term [\(15\)](#page-13-1), a standard application of bilinear estimates Theorem [2.5](#page-12-2) concludes that

$$
\sum_{\tau \sim \tau' \subset \mathcal{A}_1^{j_0}} \left\| [e^{it|\nabla|^{\alpha}}] f_{\tau} \cdot [e^{it|\nabla|^{\alpha}}] f_{\tau'} \right\|_{L_{t,x}^{q_0/2}}^s \lesssim \sum_{\tau \sim \tau' \subset \mathcal{A}_1^{j_0}} \left[|\tau|^{1-2/s_0} \|\widehat{f}_{\tau}\|_{L^{s_0}} \|\widehat{f}_{\tau'}\|_{L^{s_0}} \right]^s
$$

$$
\lesssim \sum_{\tau \sim \tau' \subset \mathcal{A}_1^{j_0}} \left[|\tau|^{1-2/s_0} \|\widehat{f}_{\tau}\|_{L^{s_0}}^2 \right]^s
$$

$$
\lesssim \sum_{\tau \subset \mathcal{A}_1} \left[|\tau|^{1-2/s_0} \|\widehat{f}_{\tau}\|_{L^{s_0}}^2 \right]^s
$$

$$
\lesssim \|f\|_{L^2}^{2s},
$$

where in the last inequality we have used $[3,$ $[3,$ Theorem 1.3] with $s > 1$, see also $[44, p$ $[44, p$. 279]. Thus, this estimate and the previous [\(16\)](#page-14-1) directly imply our desired conclusion $(13).$ $(13).$

Finally, we are able to show our main result in this section, which is the following refined α -Strichartz proposition.

Proposition 2.7 *Let* $d \geq 2$ *. There exists* $\theta \in (0, 1)$ *such that the following estimate*

$$
\| [E_{\alpha}] f \|_{L_{t,x}^{q_0}} \lesssim \left[\sup_{\mathcal{Q} \in \mathcal{D}} |\mathcal{Q}|^{-\frac{1}{2}} \left\| [e^{it|\nabla|^{\alpha}}] f_{\mathcal{Q}} \right\|_{L_{t,x}^{\infty}} \right]^{\theta} \| f \|_{L_x^2}^{1-\theta} \tag{17}
$$

holds for all $f \in L^2(\mathbb{R}^d)$ *.*

Proof of Proposition [2.7](#page-14-0) Combining the annular orthogonality Lemma [2.4](#page-12-0) and annular refined α -Strichartz Lemma [2.6,](#page-12-1) we directly obtain this conclusion by applying Bernstein's inequality.

There are standard arguments to deduce the following non-zero weak limit result Corollary [2.8](#page-15-1) from the aforementioned Proposition [2.7.](#page-14-0) In fact, the $L_{t,x}^{\infty}$ -norm in [\(17\)](#page-15-2) gives space-time translation parameters and the supremum of dyadic cubes gives scaling-frequency parameters. Readers can see [\[20](#page-33-9), Corollary 3.2] for further details and the detailed proof are omitted here for simplicity.

Corollary 2.8 *Let* (*f_n*) *be a bounded sequence in* $L^2(\mathbb{R}^d)$ *with* $d \geq 2$ *. If* $\|[E_\alpha]f_n\|_{L^{q_0}_{t,x}} \nrightarrow$ 0*, then there exists* $(t_n, x_n, \xi_n, h_n) \subset \mathbb{R} \times \mathbb{R}^d \times \mathbb{R}^d \times \mathbb{R}_+$ with $h_n ||\xi_n||_{\min} \geq 1/2$, such *that up to subsequences*

$$
\widehat{[g_n]f_n}(\xi+h_n\xi_n)\rightharpoonup \widehat{V}
$$

in weak topology of $L^2(\mathbb{R}^d)$ *with* $\widehat{V} \neq 0$ *. Moreover, if* $||f_n||_{L^2_x} \leq B$ *and* $\limsup_{n\to\infty}$ $||E_{\alpha} f_n||_{L^{q_0}_{t,x}} \geq A$ *, then we have*

$$
||V||_{L_x^2} \geq C A^{\beta} B^{-\gamma},
$$

where C , β , γ *depend only on the dimension d and* α *.*

3 Local Smoothing and Local Convergence

In this section we show some local smoothing property for the solution $[e^{it|\nabla|^{\alpha}}]f$. Then this property further deduces some local convergence and pointwise convergence results which will provide the desired input for Brézis-Lieb type lemma [\[20,](#page-33-9) Lemma 3.1]. Similar arguments can also be found in [\[20,](#page-33-9) [21\]](#page-33-11). Our local smoothing consequence is the following lemma.

Lemma 3.1 *Let* $\phi \in L^1(\mathbb{R}^d)$ *be a Schwartz function. Then for all* $f \in L^2(\mathbb{R}^d)$ *, there holds*

$$
\int_{\mathbb{R}^{d+1}} \phi(x) \left| [D^{\frac{\alpha-1}{2}}] [e^{it|\nabla|^{\alpha}}] f(x) \right|^2 dx dt \lesssim_{\phi} ||f||_{L^2}.
$$

Proof of Lemma **[3.1](#page-15-3)** We will mimic the proofs in [\[20,](#page-33-9) Lemma 4.4] and [\[21,](#page-33-11) Lemma 4.3], albeit with a small twist.

Let \mathcal{F}_t be the Fourier transform on time space. Then there holds

$$
\mathscr{F}_t e^{ia_0t}(\lambda) = 2\pi \delta(\lambda - a_0).
$$

Here δ denotes the Dirac delta function. Thus, Plancherel theorem implies the following estimate

$$
\int_{\mathbb{R}^{d+1}} \phi(x) \left| [D^{\frac{\alpha-2}{2}}] [e^{it|\xi|^{\alpha}}] f(x) \right|^2 dx dt
$$

\$\sim \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} \widehat{\phi}(\xi'-\xi) |\xi|^{\frac{\alpha-1}{2}} |\xi'|^{\frac{\alpha-1}{2}} \widehat{f}(\xi) \overline{\widehat{f}(\xi')} \delta(|\xi|^{\alpha} - |\xi'|^{\alpha}) d\xi d\xi'.\$

By Schur test, we only need to show

$$
\sup_{\xi}\int_{\mathbb{R}^d}\widehat{\phi}(\xi'-\xi)|\xi|^{\frac{\alpha-1}{2}}|\xi'|^{\frac{\alpha-1}{2}}\delta(|\xi|^{\alpha}-|\xi'|^{\alpha})d\xi'<\infty.
$$

Setting polar coordinate $\xi' = k\theta$ with $\theta \in \mathbb{S}^{d-1}$ and denoting $\xi = |\xi| \theta_{\xi}$, it remains to deduce

$$
\sup_{\xi} |\xi|^{d-1} \int_{\mathbb{S}^{d-1}} \left| \widehat{\phi}\left(|\xi|(\theta_{\xi} - \theta)\right) \right| d\theta < \infty.
$$

The dominated convergence theorem implies that $\int_{\mathbb{S}^{d-1}} |\widehat{\phi}|$ $\left(|\xi| (\theta_{\xi} - \theta) \right) | d\theta$ is a continuous function with respect to ξ , hence it suffices to prove

$$
\int_{\mathbb{S}^{d-1}} \left| \widehat{\phi}\left(|\xi| (\theta_{\xi} - \theta) \right) \right| d\theta = O\left(|\xi|^{1-d} \right) \tag{18}
$$

as $|\xi| \to \infty$. A changing of variables gives

$$
\int_{\mathbb{S}^{d-1}} \left| \widehat{\phi}\left(|\xi|(\theta_{\xi}-\theta)\right) \right| d\theta = \frac{1}{|\xi|^{d-1}} \int_{\mathbb{S}^{d-1}_{\xi}} |\widehat{\phi}(\theta)| d\theta,
$$

where $\mathbb{S}_{\xi}^{d-1} := \{ \theta \in \mathbb{R}^d : |\theta - \theta_{\xi}| = |\xi| \}$ is the sphere centered at θ_{ξ} with radius $|\xi|$. Then the fact that ϕ is a Schwartz function deduces the desired result [\(18\)](#page-16-0) and finishes the proof. \Box

Lemma 3.2 *Let* $\beta < \frac{\alpha - 1}{2}$ *and define the operator* $[\bar{E}_{\beta}] := [D^{\beta}][e^{it|\nabla|^{\alpha}}]$ *. For a bounded sequence of functions* (f_n) *in* $L^2(\mathbb{R}^d)$ *, if we have*

$$
f_n\rightarrow 0
$$

Birkhäuser

weakly in $L^2(\mathbb{R}^d)$ *, then up to subsequences there holds* $[\bar{E}_\beta] f_n(t, x) \to 0$ *almost everywhere in* R*d*+1*.*

Proof of Lemma **[3.2](#page-16-1)** Based on the local smoothing Lemma [3.1,](#page-15-3) the proof is standard. Similar proofs can also be found in [\[20](#page-33-9), Proposition 4.3], [\[21,](#page-33-11) Lemma 4.3] and [\[21,](#page-33-11) Lemma D.1].

We aim to show that $[\bar{E}_{\beta}] f_n \to 0$ strongly in $L^2_{loc}(\mathbb{R}^{d+1})$, which implies the desired result by a Cantor diagonal argument. Let $K \subset \mathbb{R}^{d+1}$ be a compact set and $\varepsilon > 0$. For $\Lambda > 0$, define

$$
[P_{\Lambda}]f(x) := \mathscr{F}^{-1}1_{B_{\Lambda}}\mathscr{F}f(x), \quad [P_{\Lambda}^{\perp}]f(x) := f(x) - [P_{\Lambda}]f(x).
$$

Here B_{Λ} is the ball in \mathbb{R}^d centered at origin with radius Λ . For the $[P_{\Lambda}^{\perp}]$ term, since $\beta < \frac{\alpha - 1}{2}$, we use Lemma [3.1](#page-15-3) to conclude that

$$
\|1_{K}[P_{\Lambda}^{\perp}][\bar{E}_{\beta}]f_{n}\|_{L^{2}_{t,x}} \leq \|1_{K}e^{|x|^{2}}\|_{L^{\infty}_{t,x}}\|e^{-|x|^{2}}[D^{\frac{\alpha-1}{2}}][e^{it|\nabla|^{\alpha}}]\|_{L^{2}_{x}\to L^{2}_{t,x}}
$$

$$
\times \left\|[P_{\Lambda}^{\perp}][D^{\beta-\frac{\alpha-1}{2}}]\right\|_{L^{2}_{x}\to L^{2}_{x}}\|f_{n}\|_{L^{2}_{x}}
$$

$$
\leq C_{K}\Lambda^{\beta-\frac{\alpha-1}{2}},
$$

where C_K is independent of *n* and the notation $||[T]||_{L^p \to L^q}$ denotes the norm of the operator [*T*]. Hence we can choose Λ large enough independent of *n* such that

$$
\left\|1_K[P_\Lambda^\perp][\bar{E}_\beta]f_n\right\|_{L^2_{t,x}}\leq\varepsilon.
$$

Then for this large Λ and the $[P_{\Lambda}]$ term, we claim that for every fixed $t \in \mathbb{R}$ there holds

$$
\mathbb{1}_K(t,\cdot)[P_\Lambda][D^\beta][e^{it|\nabla|^\alpha}]f_n \to 0 \tag{19}
$$

strongly in $L_x^2(\mathbb{R}^d)$ as $n \to \infty$. Indeed, for any $(t, x) \in \mathbb{R} \times \mathbb{R}^d$, since $f_n \to 0$ in $L_x^2(\mathbb{R}^d)$ we have

$$
\mathbb{1}_K(t,\cdot)[P_\Lambda][D^\beta][e^{it|\nabla|^\alpha}]f_n(x)=\mathbb{1}_K(t,x)\int_{\mathbb{R}^d}\mathbb{1}_{\{|\xi|\leq\Lambda\}}|\xi|^{\beta}e^{ix\xi+it|\xi|^\alpha}\widehat{f_n}(\xi)\mathrm{d}\xi\to 0
$$

as $n \to \infty$. Meanwhile, by Cauchy–Schwarz and the boundedness of $||f_n||_{L_x^2}$, there holds

$$
\left|\mathbb{1}_{K}(t,\cdot)[P_{\Lambda}][D^{\beta}][e^{it|\nabla|^{\alpha}}]f_{n}(x)\right|\lesssim \mathbb{1}_{K}(t,x)\Lambda^{\beta+d/2}\|f_{n}\|_{L_{x}^{2}}\lesssim_{\Lambda} \mathbb{1}_{K}(t,x).
$$

Therefore, dominated convergence theorem gives the desired claim [\(19\)](#page-17-0). Notice that there holds

$$
\left\|1_K(t,\cdot)[P_\Lambda][D^\beta][e^{it|\nabla|^\alpha}]f_n(x)\right\|_{L^2_{t,x}} \lesssim_\Lambda \|1_K\|_{L^2_{t,x}}.
$$

Hence, the claim [\(19\)](#page-17-0) and dominated convergence theorem imply that $1_K [P_\Lambda][\bar{E}_\beta] f_n \to$ 0 strongly in $L^2_{t,x}(\mathbb{R}^{d+1})$ as $n \to \infty$. This completes the proof.

4 Approximate Symmetry

As mentioned in the introduction, to establish the main precompactness result Theorem [1.2,](#page-3-1) one should understand the behavior of *approximate symmetries* in the sense of [\[21,](#page-33-11) Remark 2.6]. In other words, we should investigate the effect of frequency parameters ξ_n . In this section, for frequencies $(\xi_n) \subset \mathbb{R}^d$ with $|\xi_n| \to \infty$, then up to subsequences we introduce the following notations

$$
\bar{\xi}_n := \frac{\xi_n}{|\xi_n|}, \quad \xi_0 := \lim_{n \to \infty} \bar{\xi}_n.
$$

For an unit vector $\xi_0 \in \mathbb{S}^{d-1}$, we define the linear transformation A_0 on \mathbb{R}^d as follows

$$
A_0: \xi \mapsto \xi^{\perp} \sqrt{\alpha/2} + \xi^{\shortparallel} \sqrt{\alpha(\alpha - 1)/2}, \quad \xi^{\shortparallel} := (\xi \xi_0) \xi_0, \quad \xi^{\perp} := \xi - \xi^{\shortparallel};
$$

and the associated unitary operator $[\tilde{A}_0]$ on $L^2(\mathbb{R}^d)$ is defined by

$$
[\tilde{A}_0]f(x) := |\det A_0|^{1/2} f(A_0 x). \tag{20}
$$

Note that the absolute value of determinate $|\det A_0| = (\alpha/2)^{d/2}\sqrt{\alpha-1}$. The approximate symmetry mainly investigates the operator

$$
[T_{\alpha}^{n}]f(t,x) := \int_{\mathbb{R}^{d}} |\xi + \xi_{n}|^{\frac{\alpha-2}{q_0}} e^{ix\xi + it|\xi + \xi_{n}|^{\alpha}} \widehat{f}(\xi) d\xi
$$
 (21)

and the following function

$$
F_{\alpha}^{n}(t,x) = [\bar{T}_{\alpha}^{n}]f(t,x) := \int_{\mathbb{R}^{d}} \left| \frac{\xi}{|\xi_{n}|} + \bar{\xi}_{n} \right|^{ \frac{\alpha-2}{q_0}} e^{ix\xi + it\Phi_{n}(\xi)} \widehat{f}(\xi) d\xi, \qquad (22)
$$

where

$$
\Phi_n(\xi) := \frac{1}{|\xi_n|^{\alpha-2}} \left[|\xi + \xi_n|^\alpha - |\xi_n|^\alpha - \alpha |\xi_n|^{\alpha-2} \xi_n \xi \right].
$$

Indeed, it is not hard to see that there holds

$$
\| [T^n_{\alpha}]f \|_{L^{q_0}_{t,x}} = \| F^n_{\alpha} \|_{L^{q_0}_{t,x}}.
$$

Our first result on the aforementioned approximate operator $[T_\alpha^n]$ is the following asymptotic Schrödinger behavior lemma.

.

Lemma 4.1 *Let* $f \in L^2(\mathbb{R}^d)$ *with* $|| f ||_{L^2} = 1$ *and frequencies* (ξ_n) ⊂ \mathbb{R}^d *. Set*

$$
\widehat{f}_n(\xi) := \widehat{f}(\xi - \xi_n), \quad \mathbf{a}_{d,\alpha}^* := (\alpha - 1)^{\frac{-1}{2d+4}} (\alpha/2)^{\frac{-d}{2d+4}}.
$$

Then, up to subsequences, we have the following asymptotic Schrödinger behavior

$$
\lim_{|\xi_n|\to\infty} \|[E_{\alpha}]f_n\|_{L^{q_0}_{t,x}} = \lim_{n\to\infty} \|(T_{\alpha}^n)f\|_{L^{q_0}_{t,x}} = \mathbf{a}^*_{d,\alpha} \left\|[e^{it\Delta}][\tilde{A}_0]f\right\|_{L^{q_0}_{t,x}}
$$

Remark 4.2 Based on the existence of extremals for S_d^* by [\[39\]](#page-34-5), this Lemma [4.1](#page-18-0) immediately gives

$$
\mathbf{M}_{d,\alpha} \geq (\alpha - 1)^{\frac{-1}{2d+4}} (\alpha/2)^{\frac{-d}{2d+4}} \mathbf{S}_d^*.
$$

For the case $d = 1$, after some accurate numerical calculations, it has been shown in [\[6](#page-33-10), Theorem 1.4] that there exists $\alpha_1 \approx 5.485$ such that the strict inequality [\(2\)](#page-3-2) holds for indices $\alpha \in (1, \alpha_1) \setminus \{2\}$. For the case $d = 2$, it has been proved in [\[36](#page-34-10), Proposition 6.9] that there exists $\alpha_0 > 5$ such that the strict inequality [\(2\)](#page-3-2) holds for $\alpha \in (2, \alpha_0)$. These consequences suggest that the aforementioned strict inequality might hold for all general (d, α) with $\alpha > 2$. However, it remains an open question for us.

Proof of Lemma [4.1](#page-18-0) The change of variables can yield that

$$
\begin{aligned} &\|[E_{\alpha}]f_n\|_{L_{t,x}^{q_0}} = \left\|[T_{\alpha}^n]f\right\|_{L_{t,x}^{q_0}},\\ &\|[T_{\alpha}^n]\|_{L_x^2 \to L_{t,x}^{q_0}} = \left\|[T_{\alpha}^n]\right\|_{L_x^2 \to L_{t,x}^{q_0}} = \left\|[E_{\alpha}]\right\|_{L_x^2 \to L_{t,x}^{q_0}} = \mathbf{M}_{d,\alpha}.\end{aligned} \tag{23}
$$

Hence by some dense argument we may assume *f* has compact Fourier support at the beginning. Using our aforementioned definition [\(22\)](#page-18-2), we can obtain that

$$
\begin{split} \left\| [T_{\alpha}^{n}] f \right\|_{L_{t,x}^{q_0}} &= \left\| \int_{\mathbb{R}^d} \left| \xi_n \right|^{\frac{\alpha-2}{q_0}} \left| \frac{\xi}{|\xi_n|} + \bar{\xi}_n \right|^{\frac{\alpha-2}{q_0}} e^{ix\xi + it|\xi + \xi_n|^{\alpha}} \widehat{f}(\xi) \mathrm{d}\xi \right\|_{L_{t,x}^{q_0}} \\ &= \left\| \int_{\mathbb{R}^d} \left| \frac{\xi}{|\xi_n|} + \bar{\xi}_n \right|^{\frac{\alpha-2}{q_0}} e^{ix\xi + it \frac{|\xi + \xi_n|^{\alpha}}{|\xi_n|^{\alpha-2}}} \widehat{f}(\xi) \mathrm{d}\xi \right\|_{L_{t,x}^{q_0}} \\ &= \left\| F_{\alpha}^{n} \right\|_{L_{t,x}^{q_0}} .\end{split}
$$

Then the Taylor's theorem, paired with the assumptions $|\xi_n| \to \infty$ and *f* has compact Fourier support, will directly imply the following pointwise convergence

$$
\lim_{n \to \infty} \Phi_n(\xi) = \frac{\alpha |\xi|^2 + \alpha(\alpha - 2)|\xi \xi_0|^2}{2}.
$$
\n(24)

B Birkhäuser

Therefore it is reasonable to expect the following estimate

$$
\lim_{n\to\infty} \left\| \left[T^n_{\alpha}\right]f \right\|_{L^{q_0}_{t,x}} = \left\| \int_{\mathbb{R}^d} e^{ix\xi + it\left[\alpha|\xi|^2/2 + \alpha(\alpha-2)|\xi\xi_0|^2/2\right]} \widehat{f}(\xi) d\xi \right\|_{L^{q_0}_{t,x}}.
$$
 (25)

Let us postpone the proof of this result and go ahead by using this estimate [\(25\)](#page-20-0). From the definition of A_0 and $[\tilde{A}_0]$, we know that $|A_0\xi|^2 = \alpha |\xi|^2/2 + \alpha(\alpha - 2)|\xi \xi_0|^2/2$. Then a direct computation yields that

$$
\left\| \int_{\mathbb{R}^d} e^{ix\xi + it\left[\alpha|\xi|^2/2 + \alpha(\alpha - 2)|\xi\xi_0|^2/2\right]} \widehat{f}(\xi) d\xi \right\|_{L_{t,x}^{q_0}} = \left\| \int_{\mathbb{R}^d} e^{ix\xi + it|A_0\xi|^2} \widehat{f}(\xi) d\xi \right\|_{L_{t,x}^{q_0}}
$$

\n= $|\det A_0|^{-1/2} \left\| [e^{it\Delta}] [\tilde{A}_0] f(A_0^{-1}x) \right\|_{L_{t,x}^{q_0}}$
\n= $|\det A_0|^{1/q_0 - 1/2} \left\| [e^{it\Delta}] [\tilde{A}_0] f \right\|_{L_{t,x}^{q_0}}$.

Finally, the fact $|\det A_0| = (\alpha/2)^{d/2}\sqrt{\alpha - 1}$ gives the constant $\mathbf{a}^*_{d,\alpha} = (\alpha - 1)^d$ $\frac{-1}{2d+4}$ (α /2) $\frac{-d}{2d+4}$.

Now it remains to prove [\(25\)](#page-20-0). Indeed, we can obtain the following stronger result

$$
F_{\alpha}^{n}(t,x) \rightarrow \int_{\mathbb{R}^{d}} e^{ix\xi + it\left[\alpha|\xi|^{2}/2 + \alpha(\alpha - 2)|\xi\xi_{0}|^{2}/2\right]} \widehat{f}(\xi) d\xi
$$
 (26)

strongly in $L_{t,x}^{q_0}(\mathbb{R}^{d+1})$ as $n \to \infty$. Firstly for any $(t, x) \in \mathbb{R}^{d+1}$, since f is a Schwartz function with compact Fourier support, dominated convergence theorem and the estimate [\(24\)](#page-19-1) imply

$$
\lim_{n \to \infty} F_{\alpha}^{n}(t, x) = \int_{\mathbb{R}^{d}} e^{ix\xi + it\left[\alpha|\xi|^{2}/2 + \alpha(\alpha - 2)|\xi\xi_{0}|^{2}/2\right]} \widehat{f}(\xi) d\xi.
$$
 (27)

On the other hand, for *n* large enough, by a standard application of stationary phase method we can deduce the following two decay estimates

$$
|F_{\alpha}^{n}(t,x)| \le C(1+t^{2}+|x|^{2})^{-\frac{d}{4}}
$$
\n(28)

and

$$
|H_{\alpha}(t,x)| := \left| \int_{\mathbb{R}^d} e^{ix\xi + it\left[\alpha|\xi|^2/2 + \alpha(\alpha - 2)|\xi\xi_0|^2/2\right]} \widehat{f}(\xi) d\xi \right| \le C(1 + t^2 + |x|^2)^{-\frac{d}{4}},\tag{29}
$$

where the constant C is independent of n . Assume for the moment these two decay estimates and let us show the desired convergence result (26) first. From the decay estimates (28) and (29) , after taking a subsequence, we have that for all *n* and some constant C' there holds

$$
\int_{|(t,x)|>R} (|F_n(t,x)|^{q_0} + |H_\alpha(t,x)|^{q_0}) \, \mathrm{d}t \mathrm{d}x \leq \frac{C'}{R}.
$$

By setting $R > 0$ large enough, this term can be arbitrary small uniformly in *n*. Hence it suffices to prove that for any fixed $R > 0$, there holds

$$
\mathbb{1}_{B_R}(t,x)F^n_{\alpha}(t,x)\to \mathbb{1}_{B_R}(t,x)H_{\alpha}(t,x)
$$

strongly in $L_{t,x}^{q_0}(\mathbb{R}^{d+1})$ as $n \to \infty$. By dominated convergence theorem, this follows immediately from the pointwise convergence [\(27\)](#page-20-4) and the following uniform bound (for *n* large enough)

$$
|F_{\alpha}^{n}(t,x)| \lesssim \| \widehat{f} \|_{L^{1}} < \infty.
$$

Thus, it remains to prove the decay estimates [\(28\)](#page-20-2) and [\(29\)](#page-20-3) by means of stationary phase method. In view of the Fourier extension on the surface

$$
\bar{S} := (\xi, \alpha |\xi|^2 / 2 + \alpha (\alpha - 2) |\xi \xi_0|^2 / 2),
$$

then the fact that \bar{S} has non-zero Gaussian curvature and [\[41](#page-34-20), p. 348, Theorem 1] directly give the result [\(29\)](#page-20-3). Indeed, a direct computation shows that the determinant of corresponding Hessian matrix is a constant α^{d} (2 α – 2). To deduce [\(28\)](#page-20-2), since $|\xi_n| \to \infty$, we point out that for *n* large enough there holds

$$
|F_{\alpha}^{n}(t,x)| \leq 3 \left| \int_{\mathbb{R}^{d}} e^{ix\xi + it\Phi_{n}(\xi)} \widehat{f}(\xi) d\xi \right|.
$$

Let $\Phi_n(\bar{x}, \bar{t}, \xi) := \bar{x}\xi + \bar{t}\Phi_n(\xi)$ with $(\bar{x}, \bar{t}) := (x, t)/|(x, t)|$. We investigate the gradient

$$
\nabla \bar{\Phi}_n(\bar{x}, \bar{t}, \xi) = \bar{x} + \frac{\bar{t}}{|\xi_n|^{\alpha - 2}} \left[\alpha |\xi + \xi_n|^{\alpha - 2} (\xi + \xi_n) - \alpha |\xi_n|^{\alpha - 2} \xi_n \right]
$$

and the Hessian

$$
\text{Hess}\bar{\Phi}_n(\bar{x},\bar{t},\xi) = \frac{\bar{t}}{|\xi_n|^{\alpha-2}} \left(\alpha |\xi + \xi_n|^{\alpha-2} E_d + \alpha(\alpha-2) |\xi + \xi_n|^{\alpha-4} L_n^T L_n \right),
$$

where E_d is the identity matrix, $L_n := (\xi + \xi_n)$ is viewed as row vector and L_n^T is the transpose of L_n . Then we can compute the absolute value for the determinant of this Hessian

$$
\left|\det \text{Hess}\,\overline{\Phi}_n(\bar{x},\bar{t},\xi)\right| = \frac{(\alpha\bar{t})^d(\alpha-1)|\xi+\xi_n|^{d(\alpha-2)}}{|\xi_n|^{d(\alpha-2)}}.
$$

Birkhäuser

Since $|\xi_n| \to \infty$ and the ξ -localization, we have the following estimates

$$
\lim_{n\to\infty} |\nabla \bar{\Phi}_n| = |\bar{x} + \alpha(\alpha - 1)\bar{t}\xi|, \quad \lim_{n\to\infty} |\det \text{Hess}\bar{\Phi}_n| = (\alpha\bar{t})^d (\alpha - 1).
$$

Thus, for *n* large enough, the proof is divided into two cases. Firstly if $\bar{t} \geq \varepsilon$ for some small ε to be determined later, then we always have $|\det \text{Hess} \bar{\Phi}_n| \gtrsim \varepsilon^d$. Hence the uniform stationary phase estimates of Alazard, Burq and Zuily [\[1](#page-33-23), Theorem 1] will give the desired estimate [\(28\)](#page-20-2). Secondly if $\bar{t} < \varepsilon$ which means $|\bar{x}| > \sqrt{1 - \varepsilon^2}$, duo to the ξ -localization, we can choose ε small enough such that $|\bar{x}| \ge \alpha(\alpha - 1)t |\xi| + 1/8$ uniformly in ξ . Then we always have the following uniform estimate

$$
\left|\nabla\bar{\Phi}_n\right| \geq |\bar{x}| - \alpha(\alpha - 1)\bar{t}|\xi| \geq 1/8 > 0.
$$

Therefore, the classical integration by parts arguments [\[41](#page-34-20), p. 341, Proposition 4] will deduce an even faster decay than the desired conclusion [\(28\)](#page-20-2) and complete the proof. \Box

As we have done in Sect. [3,](#page-15-0) to use the Brézis-Lieb type lemma, analogously the corresponding local smoothing and local convergence results are required to be established. At this point due to the parameters ξ_n , we need to do some extra estimates by using the Taylor's theorem.

Lemma 4.3 *Suppose that* ϕ *be a Schwartz function and* $|\eta| \ge 100$ *. There exists* $C > 0$ *such that for arbitrary function* $f \in L^2(\mathbb{R}^d)$ *which has Fourier support in* $\{\xi : |\xi|$ < |η|/5}*, we have*

$$
\int_{\mathbb{R}^{d+1}} \phi(x) \Big| [\psi_{\eta}(D)][\bar{T}_{\alpha}^{\eta}] f(t,x) \Big|^2 \mathrm{d}x \mathrm{d}t \leq C \|f\|_{L_x^2}^2,
$$

 \mathbb{R} *where* $[\bar{T}_{\alpha}^{\eta}]$ *is defined as in* [\(22\)](#page-18-2) *with parameter* η *substituting* ξ_n *and*

$$
\psi_{\eta}(\xi):=|\eta|^{\frac{\alpha-2}{q_0}}|\xi|^{1/2}\,|\xi+\eta|^{-\frac{\alpha-2}{q_0}}\,.
$$

Proof of Lemma [4.3](#page-22-0) By Plancherel theorem, as in the proof of Lemma [3.1,](#page-15-3) we investigate

$$
\int_{\mathbb{R}^d} \int_{\mathbb{R}^d} \widehat{\phi}(\xi'-\xi)\psi_{\eta}(\xi)\psi_{\eta}(\xi')|\eta|^{\frac{4-2\alpha}{q_0}} |\xi+\eta|^{\frac{\alpha-2}{q_0}} \left|\xi'+\eta\right|^{\frac{\alpha-2}{q_0}} \widehat{f(\xi')}\widehat{f(\xi)}\delta
$$

$$
\left[\Phi_{\eta}(\xi)-\Phi_{\eta}(\xi')\right] d\xi d\xi'.
$$

Then Schur test implies that it suffices to bound the following term

$$
\sup_{\xi} \int_{\mathbb{R}^d} \widehat{\phi}(\xi' - \xi) |\xi|^{1/2} |\xi'|^{1/2} \delta \left[\Phi_{\eta}(\xi) - \Phi_{\eta}(\xi') \right] d\xi'
$$

independently of η . For every $\xi \in \mathbb{R}^d$ and $\eta \in \mathbb{R}^d$, denote the set

$$
\mathcal{V}_{\xi,\eta} := {\xi' : |\xi' + \eta|^{\alpha} - \alpha |\eta|^{\alpha - 2} \eta \xi' = |\xi + \eta|^{\alpha} - \alpha |\eta|^{\alpha - 2} \eta \xi}.
$$

Then a direct computation shows that the aforementioned term can be bounded by

$$
\sup_{\xi} \sup_{\xi' \in \mathcal{V}_{\xi,\eta}} \frac{\|\widehat{\phi}\|_{L^{\infty}} |\xi|^{1/2} |\xi'|^{1/2} |\eta|^{\alpha-2}}{\alpha \Big| |\xi' + \eta|^{\alpha-2} (\xi' + \eta) - |\eta|^{\alpha-2} \eta \Big|}.
$$

First, we states the following claim whose proof is postponed

$$
\sup_{\xi} \sup_{\xi' \in \mathcal{V}_{\xi,\eta}} \frac{\|\widehat{\phi}\|_{L^{\infty}} |\xi|^{1/2} |\xi'|^{1/2} |\eta|^{\alpha-2}}{\alpha \|\xi' + \eta\|^{\alpha-2} (\xi' + \eta) - |\eta|^{\alpha-2} \eta|} \sim_{\alpha} \sup_{\xi} \sup_{\xi' \in \mathcal{V}_{\xi,\eta}} \frac{\|\widehat{\phi}\|_{L^{\infty}} |\xi|^{1/2}}{|\xi'|^{1/2}}. (30)
$$

Now we turn to prove $|\xi'| \sim_{\alpha} |\xi|$ for all $\xi' \in \mathcal{V}_{\xi,\eta}$ uniformly in η . Then this estimate and the claim [\(30\)](#page-23-0) will immediately give the desired uniform bound $C_{\alpha} ||\phi||_{L^{\infty}}$ and complete the great leads the set \mathcal{V}_{α} can be gaugitten as complete the proof. Indeed, the set $V_{\xi,\eta}$ can be rewritten as

$$
\mathcal{V}_{\xi,\eta} = \left\{ \xi' : |\tilde{\xi}' + \bar{\eta}|^{\alpha} - \alpha |\tilde{\xi}'| \cos \theta(\xi',\eta) = |\tilde{\xi} + \bar{\eta}|^{\alpha} - \alpha |\tilde{\xi}| \cos \theta(\xi,\eta) \right\},\
$$

where

$$
\tilde{\xi}' := \xi'/|\eta|, \quad \bar{\eta} := \eta/|\eta|, \quad \tilde{\xi} := \xi/|\eta|.
$$

Note that $|\xi| \leq 1/5$ and $|\xi'| \leq 1/5$. Without loss of generality, we may assume $\bar{\eta} = (1, 0, \dots, 0)$. Then for $\xi' \in \mathcal{V}_{\xi, \eta}$ we have

$$
(1 + 2\tilde{\xi}_1' + |\tilde{\xi}'|^2)^{\alpha/2} - \alpha \tilde{\xi}_1' = (1 + 2\tilde{\xi}_1 + |\tilde{\xi}|^2)^{\alpha/2} - \alpha \tilde{\xi}_1.
$$
 (31)

Here we have used the notations $\xi = (\xi_1, \xi_2, \dots, \xi_d)$ and $\xi' = (\xi'_1, \xi'_2, \dots, \xi'_d)$. Moreover, we can use Taylor's theorem to deduce the estimates

$$
\left(1+2\tilde{\xi}_1+|\tilde{\xi}|^2\right)^{\alpha/2}-\alpha\tilde{\xi}_1 \le \left(1+2|\tilde{\xi}|+|\tilde{\xi}|^2\right)^{\alpha/2}-\alpha|\tilde{\xi}|
$$

$$
=1+\frac{\alpha(\alpha-1)}{2}\left(1+\theta_0|\tilde{\xi}|\right)^{\alpha-2}|\tilde{\xi}|^2
$$

with $0 < \theta_0 < 1$; and on the other hand there exists $0 < \theta_1 < 1$ such that

$$
\left(1+2\tilde{\xi}'_1+|\tilde{\xi}'|^2\right)^{\alpha/2}-\alpha\tilde{\xi}'_1=1+\frac{\alpha}{2}|\tilde{\xi}'|^2+\frac{\alpha(\alpha-1)}{2}\left[1+\theta_1(2\tilde{\xi}'_1+|\tilde{\xi}'|^2)\right]^{\alpha-2}\left(|\tilde{\xi}'|+2\tilde{\xi}'_1\right)^2.
$$

B Birkhäuser

Hence if there holds $|\xi| \ll |\xi'| \leq 1/5$, then it will deduce a contradiction to the identity [\(31\)](#page-23-1). Similarly $|\xi'| \ll |\xi|$ is not possible. Therefore we obtain the desired estimate $|\xi| \sim_{\alpha} |\xi'|$ and it remains to prove the claim [\(30\)](#page-23-0).

First, we assume $\xi'_1 = x$ and $|(\xi'_2, \dots, \xi'_d)| = y$ for real numbers *x* and *y*. Recall the aforementioned vector $\bar{\eta}$ and consider the vector $(x, y) \in \mathbb{R}^2$ with $|(x, y)| \le 1/5$. To prove the desired estimate (30) , we only need to show that

$$
C_1 \leq \frac{\left| (x+1, y) - (1/|(x+1, y)|^{\alpha-2}, 0) \right|}{\sqrt{x^2 + y^2}} \leq C_2,
$$

which means

$$
C_1^2 \le \frac{(x+1)^2 + [(x+1)^2 + y^2]^{2-\alpha} - 2(x+1)[(x+1)^2 + y^2]^{1-\frac{\alpha}{2}} + y^2}{x^2 + y^2} \le C_2^2.
$$
\n(32)

If $|(x, y)| \ge c_0$ for some $c_0 > 0$, the existence of C_1 and C_2 is obvious. Hence we should investigate the upper and lower bound for the case $|(x, y)| \rightarrow 0$. In this case, by denoting $z = x^2 + y^2 + 2x$ and using Taylor's theorem, we conclude that the middle item in [\(32\)](#page-24-0) equals to

$$
\frac{x^2 + y^2 + (\alpha - 2)x(x^2 + y^2 + 2x) + O(z^2) - 2(x + 1)O(z^2) + O(z^{2\alpha - 4}) + O(z^{\alpha - 2})}{x^2 + y^2}.
$$

The limit does not exist as $(x, y) \rightarrow (0, 0)$, but we can give the bounds for all points (x, y) in the neighborhood of origin. Notice that for $|(x, y)| \le \min\{3/(\alpha - 2), 1/5\}$ there holds

$$
(\alpha - 2)x(x^{2} + y^{2} + 2x) \ge \frac{-x^{2} - y^{2}}{3}, \quad (\alpha - 2)x(x^{2} + y^{2} + 2x)
$$

$$
\le 3(\alpha - 2)(x^{2} + y^{2}).
$$

Hence we obtain the desired lower bound $1/2$ and upper bound 3α for the case $|(x, y)| < c_0$ with c_0 small enough. This further implies the desired conclusion [\(32\)](#page-24-0). \Box

Lemma 4.4 *Let* $(\xi_n) \subset \mathbb{R}^d$ *with* $|\xi_n| \to \infty$ *. If a bounded sequence of functions* (f_n) *in* $L^2(\mathbb{R}^d)$ *satisfies*

$$
\mathrm{supp}\, f_n \subset \{\xi : |\xi| \le |\xi_n|/5\}, \quad f_n \to 0
$$

weakly in $L^2(\mathbb{R}^d)$ *as* $n \to \infty$ *, then up to subsequences we have*

$$
[T^n_\alpha]f_n\to 0
$$

strongly in $L^2_{loc,t,x}(\mathbb{R}^{d+1})$ *and hence* $[T^n_\alpha]f_n(t,x) \to 0$ *almost everywhere in* \mathbb{R}^{d+1} *.*

Proof of Lemma **[4.4](#page-24-1)** Considering the function $\psi_n(\xi)$ defined in Lemma [4.3,](#page-22-0) it is not hard to see that for any given $\varepsilon > 0$, there exists Λ large enough depending only on ε such that for all $|\xi| \ge \Lambda$ with $|\xi| \le |\eta|/5$ there holds

$$
\frac{1}{\psi_{\eta}(\xi)} < \varepsilon.
$$

Hence the desired conclusion follows from a standard argument by imitating the proof of Lemma [3.2,](#page-16-1) which means controlling the $[P_{\Lambda}^{\perp}]$ term by using Lemma [4.3](#page-22-0) and estimating the $[P_\Lambda]$ term by using dominated convergence theorem since we have the pointwise convergence [\(24\)](#page-19-1). Similar proof also can be found in [\[21,](#page-33-11) Lemma 4.3], and the details are omitted here for avoiding too much repetition.

5 Method of Missing Mass

In this section, by adapting the arguments in [\[20,](#page-33-9) Sect. 2] and [\[21](#page-33-11), Sect. 2], we show how the missing mass method can give the desired precompactness result Theorem [1.2.](#page-3-1) First, we introduce some definitions.

Definition 5.1 Let (f_n) ⊂ $L^2(\mathbb{R}^d)$. We write f_n → 0 if for all sequences of symmetries $([g_n])$ ⊂ *G* there holds the weak convergence $[g_n] f_n$ → 0 in $L^2(\mathbb{R}^d)$.

Define the set

$$
\mathcal{P} := \left\{ (f_n) : \|f_n\|_{L^2(\mathbb{R}^d)} = 1, f_n \rightsquigarrow 0 \right\}
$$

and the *sharp* α-*Strichartz constant with respect to P* as follows

$$
\mathbf{M}_{d,\alpha}^* := \sup \left\{ \limsup_{n \to \infty} \left\| [E_{\alpha}] f_n \right\|_{L_{t,x}^{q_0}(\mathbb{R}^{d+1})} : (f_n) \in \mathcal{P} \right\}.
$$

As we consider the precompactness of extremal sequences for $M_{d,\alpha}$, it is obvious that the sequences in P are not precompact up to symmetries and thus are our enemies. On the other hand, our next result Proposition [5.2](#page-25-1) states that all the enemies are in *P*.

Proposition 5.2 *Let* $d \geq 2$ *. All the extremal sequences for* $\mathbf{M}_{d,\alpha}$ *are precompact up to symmetries if and only if*

$$
\mathbf{M}_{d,\alpha} > \mathbf{M}_{d,\alpha}^*.
$$

Proof of Proposition [5.2](#page-25-1) It is clear that $M_{d,\alpha} \geq M^*_{d,\alpha}$. Hence the 'only if' part comes from the definition of $M^*_{d,\alpha}$ and Theorem [5.3](#page-26-0) which claims that the supremum value $\mathbf{M}_{d,\alpha}^*$ can be attained.

Thus, we only need to show the 'if' part. Assume that $\mathbf{M}_{d,\alpha} > \mathbf{M}_{d,\alpha}^*$ and (f_n) is an extremal sequence for $\mathbf{M}_{d,\alpha}$. In this case there exists ([g_n]) $\subset G$ such that, up to subsequences, $[g_n] f_n \to v \neq 0$ weakly in $L^2(\mathbb{R}^d)$ as $n \to \infty$ since $(f_n) \notin \mathcal{P}$. Hence if we write

$$
v_n := [g_n]f_n, \quad r_n := v_n - v,
$$

then for $n \to \infty$ we have $r_n \to 0$ weakly in $L^2(\mathbb{R}^d)$ and further

$$
1 = ||f_n||_{L^2}^2 = ||v_n||_{L^2}^2 = ||v||_{L^2}^2 + ||r_n||_{L^2}^2 + o(1)
$$
\n(33)

.

due to the fact that L^2 is a Hilbert space. Meanwhile the local convergence Lemma [3.2](#page-16-1) implies that $[E_{\alpha}]r_n(t, x) \rightarrow 0$ almost everywhere in \mathbb{R}^{d+1} since $q_0 > 2$. Then as $n \to \infty$, a variant of Brézis-Lieb lemma [\[20,](#page-33-9) Lemma 3.1] gives that

$$
\mathbf{M}_{d,\alpha}^{q_0} = \| [E_{\alpha}] v_n \|_{L_{t,x}^{q_0}}^{q_0} + o(1)
$$

\n
$$
= \| [E_{\alpha}] v \|_{L_{t,x}^{q_0}}^{q_0} + \| [E_{\alpha}] r_n \|_{L_{t,x}^{q_0}}^{q_0} + o(1)
$$

\n
$$
\leq \mathbf{M}_{d,\alpha}^{q_0} \left(\| v \|_{L_x^2}^{q_0} + \| r_n \|_{L_x^2}^{q_0} \right) + o(1).
$$

Combining this estimate with [\(33\)](#page-26-1) and letting $n \to \infty$, we conclude

$$
1 \leq \|v\|_{L^2}^{q_0} + \left(1 - \|v\|_{L^2}^2\right)^{q_0/2}
$$

Therefore the fact $q_0 > 2$ implies either $||v||_{L^2}^2 = 0$ or $1 - ||v||_{L^2}^2 = 0$. Since $v \neq 0$, we obtain that

$$
0 = 1 - ||v||_{L^2}^2 = \lim_{n \to \infty} ||r_n||_{L^2}^2 = \lim_{n \to \infty} ||v_n - v||_{L^2}^2.
$$

This states that $v_n = [g_n] f_n$ convergence strongly to v in $L^2(\mathbb{R}^d)$ and completes the proof. \Box

The Proposition [5.2](#page-25-1) above gives a characterization on the precompactness of extremal sequences for $\mathbf{M}_{d,\alpha}$. Hence our main result Theorem [1.2](#page-3-1) is reduced to the following Theorem [5.3](#page-26-0) which is also used in the proof of Proposition [5.2.](#page-25-1)

Theorem 5.3 *Suppose that*
$$
\mathbf{a}_{d,\alpha}^* := (\alpha - 1)^{\frac{-1}{2d+4}} (\alpha/2)^{\frac{-d}{2d+4}}
$$
 and $d \ge 2$. *There holds*

$$
\mathbf{M}_{d,\alpha}^* = \mathbf{a}_{d,\alpha}^* \mathbf{S}_d^*.
$$

Furthermore, the supremum **M**∗ *^d*,α *is attained. In other words, there is a sequence* $(f_n) \subset L^2(\mathbb{R}^d)$ *with* $||f_n||_{L^2_x} = 1$ *such that* $f_n \rightsquigarrow 0$ *and* $\limsup_{n\to\infty} ||[E_\alpha]f_n||_{L^{q_0}_{t,x}} = 0$ $\mathbf{M}_{d,\alpha}^*$.

Proof of Theorem [5.3](#page-26-0) The proof follows from a more involved version of missing mass method. We first show that $M^*_{d,\alpha} \ge a^*_{d,\alpha} S^*_{d}$. For a sequence (ξ_n) with $|\xi_n| \to \infty$ and $\phi \in L^2(\mathbb{R}^d)$ with $\|\phi\|_{L^2} = 1$, we define $\widehat{f}_n(\xi) := \widehat{\phi}(\xi - \xi_n)$. It is not hard to see that $f_n \rightarrow 0$. Then Lemma [4.1](#page-18-0) implies that

$$
\lim_{n\to\infty} \left\| [E_{\alpha}] f_n \right\|_{L^{q_0}_{t,x}} = \mathbf{a}_{d,\alpha}^* \left\| [e^{it\Delta}] \phi \right\|_{L^{q_0}_{t,x}}.
$$

Hence we conclude

$$
\mathbf{M}_{d,\alpha}^* \geq \mathbf{a}_{d,\alpha}^* \left\| [e^{it\Delta}] \phi \right\|_{L_{t,x}^{q_0}}.
$$

By taking supremum over all such functions ϕ , we obtain $\mathbf{M}_{d,\alpha}^* \ge \mathbf{a}_{d,\alpha}^* \mathbf{S}_{d}^*$.

Furthermore, from [\[39](#page-34-5)] it is known that there exists an extremal for S_d^* . Then taking ϕ in the above argument to be this extremal will give the sequence (f_n) with $|| f_n ||_{L^2} = 1$ and $f_n \rightsquigarrow 0$ such that

$$
\lim_{n\to\infty}\|[E_\alpha]f_n\|_{L^{q_0}_{t,x}}=\mathbf{a}^*_{d,\alpha}\mathbf{S}^*_{d}.
$$

Therefore, it remains to prove the reverse inequality $\mathbf{M}_{d,\alpha}^* \leq \mathbf{a}_{d,\alpha}^* \mathbf{S}_{d}^*$.

We may assume $M^*_{d,\alpha} > 0$ without loss of generality. By the definition of $M^*_{d,\alpha}$, there exists $(f_n) \subset L^2(\mathbb{R}^d)$ satisfying $||f_n||_{L^2} = 1$ and $f_n \rightsquigarrow 0$ such that

$$
\limsup_{n\to\infty} \|[E_{\alpha}]f_n\|_{L_{t,x}^{q_0}} \ge \frac{1}{2}\mathbf{M}_{d,\alpha}^* > 0.
$$

Then Corollary [2.8](#page-15-1) yields that there exist $([g_n]) \subset G$ and ξ_n with $h_n|\xi_n| \geq \frac{1}{2}$ such that up to subsequences there holds

$$
\widehat{[g_n]f_n}(\xi+h_n\xi_n)\to\widehat{V}(\xi),\quad\|V\|_{L^2}\geq\tilde{\gamma}>0.
$$

Here $\tilde{\gamma} := \tilde{\gamma}(d, \alpha)$ depends only on the dimension *d* and α . Note that $f_n \rightsquigarrow 0$ implies $h_n|\xi_n| \to \infty$. Otherwise we have $V = 0$ which is a contradiction. Hence if we set $E_n := \{ \xi \in \mathbb{R}^d : |\xi| \leq h_n |\xi_n|/5 \}$ and write

$$
\widehat{r}_n := \mathbb{1}_{E_n} \left(\widehat{[g_n] f_n} \left(\cdot + h_n \xi_n \right) - \widehat{V} \right), \quad \widehat{q}_n := (\mathbb{1}_{\mathbb{R}^d} - \mathbb{1}_{E_n}) \left(\widehat{[g_n] f_n} \left(\cdot + h_n \xi_n \right) - \widehat{V} \right),
$$

then there holds

$$
\begin{aligned} \|\left[E_{\alpha}\right]f_n\|_{L_{t,x}^{q_0}} &= \|\left[E_{\alpha}\right]\left[g_n\right]f_n\|_{L_{t,x}^{q_0}} = \|\left[T_{\alpha}^n\right](r_n + q_n + V)\|_{L_{t,x}^{q_0}} \\ &= \|\left[\bar{T}_{\alpha}^n\right](r_n + q_n + V)\|_{L_{t,x}^{q_0}}, \end{aligned}
$$

B Birkhäuser

where $[T^n_\alpha]$ and $[T^n_\alpha]$ are defined as

$$
[T_{\alpha}^{n}]f(t,x) := \int_{\mathbb{R}^{d}} |\xi + h_{n}\xi_{n}|^{\frac{\alpha-2}{q_0}} e^{ix\xi + it|\xi + h_{n}\xi_{n}|^{\alpha}} \widehat{f}(\xi) d\xi
$$

and

$$
[\bar{T}_{\alpha}^{n}]f(t,x):=\int_{\mathbb{R}^{d}}\left|\frac{\xi}{h_{n}|\xi_{n}|}+\bar{\xi}_{n}\right|^{\frac{\alpha-2}{q_{0}}}e^{ix\xi+it\Phi_{n}(\xi)}\widehat{f}(\xi)d\xi.
$$

Here $\xi_n := \xi_n / |\xi_n|$ and the function $\Phi_n(\xi)$ is defined by

$$
\Phi_n(\xi) := \frac{1}{|h_n \xi_n|^{\alpha-2}} \left[|\xi + h_n \xi_n|^{\alpha} - |h_n \xi_n|^{\alpha} - \alpha |h_n \xi_n|^{\alpha-2} h_n \xi_n \xi \right].
$$

Notice that $q_n \to 0$ strongly in $L^2(\mathbb{R}^d)$ and further $r_n \to 0$ weakly in $L^2(\mathbb{R}^d)$. Hence the uniform boundedness of the operators $[\bar{T}_{\alpha}^{n}]$ as shown in [\(23\)](#page-19-2) implies that $[\bar{T}_{\alpha}^{n}]q_{n} \to 0$ strongly in $L^{q_0}_{t,x}(\mathbb{R}^{d+1})$. Meanwhile, by the local convergence Lemma [4.4,](#page-24-1) we conclude $[\bar{T}_{\alpha}^{n}]r_{n} \to 0$ almost everywhere in \mathbb{R}^{d+1} . Note that the strong convergence result [\(26\)](#page-20-1) implies

$$
\lim_{n\to\infty}\left\|[\bar{T}_\alpha^n]V(t,x)-\int_{\mathbb{R}^d}e^{ix\xi+it\left[\alpha|\xi|^2/2+\alpha(\alpha-2)|\xi\xi_0|^2/2\right]}\widehat{V}(\xi)\mathrm{d}\xi\right\|_{L_{t,x}^{q_0}}=0.
$$

Thus as $n \to \infty$, applying the variant of Brézis-Lieb lemma [\[20](#page-33-9), Lemma 3.1] with

$$
\pi_n = \int_{\mathbb{R}^d} e^{ix\xi + it\left[\alpha|\xi|^2/2 + \alpha(\alpha - 2)\left|\xi\xi_0\right|^2/2\right]}\widehat{V}(\xi) d\xi, \quad \rho_n = [\overline{T}_\alpha^n]r_n,
$$

we can obtain that

$$
\begin{aligned} &\| [E_{\alpha}] f_n \|_{L_{t,x}^{q_0}}^{q_0} = \| [\bar{T}_{\alpha}^n] V \|_{L_{t,x}^{q_0}}^{q_0} + \| [\bar{T}_{\alpha}^n] r_n \|_{L_{t,x}^{q_0}}^{q_0} + o(1) \\ &= \| [T_{\alpha}^n] V \|_{L_{t,x}^{q_0}}^{q_0} + \| [T_{\alpha}^n] r_n \|_{L_{t,x}^{q_0}}^{q_0} + o(1). \end{aligned} \tag{34}
$$

Due to the asymptotic Schrödinger behavior Lemma [4.1](#page-18-0) and Strichartz inequality [\(1\)](#page-0-0), we obtain the following asymptotic estimate

$$
\lim_{n\to\infty} \left\| [T_\alpha^n] V \right\|_{L_{t,x}^{q_0}} = \mathbf{a}_{d,\alpha}^* \left\| [e^{it\Delta}] [\tilde{A}_0] V \right\|_{L_{t,x}^{q_0}}
$$

$$
\leq \mathbf{a}_{d,\alpha}^* \mathbf{S}_d^* \| V \|_{L_x^2}.
$$

Here $[\tilde{A}_0]$ is an unitary operator on L_x^2 defined in [\(20\)](#page-18-3). On the other hand, a changing of variables deduces $[T_\alpha^n]r_n = [E_\alpha] \omega_n$ with $\widehat{\omega}_n := \widehat{r}_n(\cdot - h_n \xi_n)$. Since $\widehat{[g_n]f_n} \rightsquigarrow 0$

.

and $V(\cdot - h_n \xi_n) \leadsto 0$, as well as the Fourier transform is an automorphism operator on L^2 , we obtain that $\omega_n \rightarrow 0$. This fact implies

$$
\limsup_{n \to \infty} \| [T^n_{\alpha}] r_n \|_{L^{q_0}_{t,x}} = \limsup_{n \to \infty} \| [E_{\alpha}] \omega_n \|_{L^{q_0}_{t,x}} \leq \mathbf{M}^*_{d,\alpha} \left(1 - \| V \|_{L^2}^2 \right)^{1/2}
$$

Inserting these two limit estimates into [\(34\)](#page-28-0) and taking limit $n \to \infty$ can deduce the following inequality

$$
\limsup_{n\to\infty} \|[E_{\alpha}]f_n\|_{L_{t,x}^{q_0}}^{q_0} \leq \left(\mathbf{a}_{d,\alpha}^*\mathbf{S}_{d}^*\right)^{q_0} \|V\|_{L^2}^{q_0} + \left(\mathbf{M}_{d,\alpha}^*\right)^{q_0} \left(1 - \|V\|_{L^2}^2\right)^{q_0/2};
$$

and this inequality can be rewritten as

$$
\left(\mathbf{M}_{d,\alpha}^{*}\right)^{q_{0}}\left[1-\left(1-\|V\|_{L^{2}}^{2}\right)^{\frac{q_{0}}{2}}\right]-\left(\mathbf{a}_{d,\alpha}^{*}\mathbf{S}_{d}^{*}\right)^{q_{0}}\|V\|_{L^{2}}^{q_{0}}\leq\left(\mathbf{M}_{d,\alpha}^{*}\right)^{q_{0}}\\-\limsup_{n\to\infty}\|[E_{\alpha}]f_{n}\|_{L_{t,x}^{q_{0}}}^{q_{0}}.
$$

Since $q_0 > 2$ and $||V||_{L^2} \in [\tilde{\gamma}, 1]$, we conclude

$$
\left[\left(\mathbf{M}_{d,\alpha}^* \right)^{q_0} - \left(\mathbf{a}_{d,\alpha}^* \mathbf{S}_d^* \right)^{q_0} \right] \tilde{\gamma}^{q_0} \leq \left(\mathbf{M}_{d,\alpha}^* \right)^{q_0} - \limsup_{n \to \infty} \left\| [E_\alpha] f_n \right\|_{L_{l,x}^{q_0}}^{q_0}.
$$

Finally, taking supremum over all such sequences (f_n) yields the desired conclusion. \Box

6 Existence of Extremals

Based on the precompactness Theorem [1.2](#page-3-1) established in the previous section, one direct way to show the existence of extremals for $M_{d,\alpha}$ is to compare this sharp constant **_{***,α***} with the constant** $**S**_d[*]$ **. In fact, as mentioned in the introduction, several works** have been done in this direction, and similar strict-inequality phenomena have been appeared in other surfaces such as the sphere, see for example the articles [\[6](#page-33-10), [12,](#page-33-5) [20,](#page-33-9) [36,](#page-34-10) [40\]](#page-34-7).

For our fractional surface situation, with the dimension $d = 2$ and the index α in some region, Oliveira e Silva and Quilodrán [\[36,](#page-34-10) Proposition 6.9] have established the desired strict inequality [\(2\)](#page-3-2) by applying some comparison principle for convolutions of certain singular measures. This result helps us obtain the existence of extremals for $M_{2,\alpha}$ with the index α in this corresponding region. For the convenience of the reader, we recall the result [\[36](#page-34-10), Proposition 6.9] in this article and then show the detailed proof for Theorem [1.1.](#page-3-0)

Here we recall the strict-inequality result of Oliveira e Silva and Quilodrán [\[36,](#page-34-10) Proposition 6.9]. For $d = 2$, they have investigated the following convolution inequality

$$
\|f\sigma_{\alpha}\ast f\sigma_{\alpha}\|_{L^2_{t,x}(\mathbb{R}^3)}\leq \mathcal{Q}_{\alpha}^2 \|f\|_{L^2_x(\mathbb{R}^2)},
$$

where \mathcal{Q}_{α} is the sharp constant and the singular measure σ_{α} is given by

$$
d\sigma_{\alpha}(y,s) := \delta(s - |y|^{\alpha})|y|^{\frac{\alpha-2}{4}}dyds, \quad (y,s) \in \mathbb{R}^2 \times \mathbb{R}.
$$

Then they have established the following proposition.

Proposition 6.1 (Proposition 6.9 of [\[36\]](#page-34-10)) *For the dimension d* = 2*, there exists a constant* $\alpha_0 > 5$ *such that for arbitrary* $\alpha \in (2, \alpha_0)$ *there holds*

$$
\frac{\pi}{\alpha\sqrt{\alpha-1}} < \mathcal{Q}^4_{\alpha} \leq \frac{\pi}{\alpha}.
$$

Finally, we show the desired existence of extremals Theorem [1.1](#page-3-0) and complete this section.

Proof of Theorem [1.1](#page-3-0) Let us first recall the classical result that up to symmetries Gaussians are the only extremals for S_d^* when $d = \{1, 2\}$. Hence, by direct computation, one can obtain the following sharp constants

$$
\mathbf{S}_1^* = 12^{-1/12}, \quad \mathbf{S}_2^* = 2^{-1/2}.
$$

Then applying the precompactness Theorem [1.2,](#page-3-1) we obtain that the extremals for $M_{2,\alpha}$ must exist if we can show the following strict inequality

$$
\mathbf{M}_{2,\alpha}^4 > (2\alpha\sqrt{\alpha - 1})^{-1}.
$$
 (35)

Denoting the space-time Fourier transform by $\mathscr{F}_{t,x}$ with dimension $d = 2$, one can observe that

$$
[D^{\frac{\alpha-2}{4}}][e^{it|\nabla|^{\alpha}}]f(t,x) = (2\pi)^{-2}\mathscr{F}_{t,x}(f\sigma_{\alpha})(-t,-x).
$$

Hence for the sharp constants $M_{2,\alpha}$ and \mathcal{Q}_{α} , by the Plancherel theorem, there holds

$$
(2\pi)\mathbf{M}_{2,\alpha}^4 = \mathcal{Q}_{\alpha}^4.
$$

Then applying the aforementioned Proposition [6.1,](#page-30-0) we obtain that the desired strict inequality [\(35\)](#page-30-1) holds for $\alpha \in (2, \alpha_0)$ with some index $\alpha_0 > 5$. This completes the proof. \Box

Appendix A: One Geometric Consequence

For a parameter η , we define the following function

$$
F_{\eta}(\xi) = |\xi|^{\alpha} + |\eta|^{\alpha} - |\xi + \eta|^{\alpha} / 2^{\alpha - 1}
$$

on the set $E_n := \{ \xi \in \mathbb{R}^d : |\xi| \geq |\eta| \}.$ This appendix is devoted to showing a geometric result related to this function. Roughly speaking, the Proposition [A.1](#page-30-0) below states that if the angle between ξ and η is small, then the Hessian matrix of $F_n(\xi)$ is positive-definite which means the corresponding surface has positive Gaussian curvature. Hence, in this situation, the multi-variable Taylor's theorem will give some nice displacement estimates around the critical points. These estimates are crucial when we establish the quasi-orthogonality Lemma [2.3](#page-7-0) and apply the bilinear restriction theory. We first introduce some notations.

For two positive semi-definite matrices *A* and *B*, we write $A \geq B$ if $A - B \geq 0$ which means $A - B$ is positive semi-definite. Similarly for $A > B$, $A \leq B$ and $A < B$. Meanwhile the vectors in \mathbb{R}^d are viewed as row vectors and the notation ξ^T denotes the transpose of ξ . For two vectors ξ and η in \mathbb{R}^d , we use the notation $\theta(\xi, \eta)$ to denote the angle between ξ and η . In addition, E_d is the $d \times d$ identity matrix. Then our result is as follows.

Proposition A.1 *There exists a number* $\bar{K}_{d,\alpha} \in \mathbb{Z}_+$ *such that the unit cube-annular* A_1 *can be divided into* $\bar{K}_{d,\alpha}$ *parts* $A_1 = \bigcup_{j=1}^{K_{d,\alpha}} R_j$ *, and in each part* R_j *we have the following properties:*

$$
c_1 E_d \leq \text{Hess} F_\eta(\xi) \leq c_2 E_d
$$

holds for all pairs $(\xi, \eta) \in \mathcal{R}_i \times \mathcal{R}_i$ *with* $\xi \in E_n$.

Proof of Proposition [A.1](#page-30-0) Firstly it is not hard to see that, by dividing the unit sphere \mathbb{S}^{d-1} into $\bar{K}_{d,q}$ disjoint parts and then decomposing A_1 accordingly, we can achieve that for every $(\xi, \eta) \in \mathcal{R}_i \times \mathcal{R}_i$ and every $y \in \mathbb{R}^d$ there holds

$$
\left|\cos^2\theta(\xi, y) - \cos^2\theta(\xi + \eta, y)\right| \le \frac{1}{2\alpha - 4}.\tag{36}
$$

Indeed, since $|\theta(\xi + \eta, y) - \theta(\xi, y)| \leq \theta(\xi + \eta, \eta) \leq \theta(\xi, \eta)$, the continuity of trigonometric functions gives the existence of such number $K_{d,\alpha}$.

A direct computation shows

$$
\begin{aligned} \text{Hess} F_{\eta}(\xi) &= \alpha \left[|\xi|^{\alpha - 2} - \frac{|\xi + \eta|^{\alpha - 2}}{2^{\alpha - 1}} \right] E_d \\ &+ \alpha(\alpha - 2) \left[|\xi|^{\alpha - 4} \xi^T \xi - \frac{|\xi + \eta|^{\alpha - 4} (\xi + \eta)^T (\xi + \eta)}{2^{\alpha - 1}} \right]. \end{aligned}
$$

B Birkhäuser

Therefore the inequality $|\xi + \eta|^{\alpha - 2} \leq (|\xi| + |\eta|)^{\alpha - 2} \leq 2^{\alpha - 2} |\xi|^{\alpha - 2}$ implies the following estimate

$$
\text{Hess} F_{\eta}(\xi) \ge \frac{\alpha |\xi|^{\alpha - 2}}{2} \left[E_d + (\alpha - 2) \frac{\xi^T \xi}{|\xi|^2} - (\alpha - 2) \frac{(\xi + \eta)^T (\xi + \eta)}{|\xi + \eta|^2} \right].
$$

We aim to show that for every $y \in \mathbb{R}^d$ there holds

$$
c_3|y|^2 \le y \left[E_d + (\alpha - 2)\frac{\xi^T \xi}{|\xi|^2} - (\alpha - 2)\frac{(\xi + \eta)^T (\xi + \eta)}{|\xi + \eta|^2}\right]y^T,
$$

where *c*₃ is independent of *y*. This fact will give the desired constant *c*₁ since $|\xi| \sim 1$ in our situation. By homogeneity we may assume $y \in \mathbb{S}^{d-1}$. Then we conclude that

$$
y\left[E_d + (\alpha - 2)\frac{\xi^T \xi}{|\xi|^2} - (\alpha - 2)\frac{(\xi + \eta)^T(\xi + \eta)}{|\xi + \eta|^2}\right]y^T
$$

= 1 + (\alpha - 2)\left[\frac{|\xi y^T|^2}{|\xi|^2} - \frac{|(\xi + \eta)y^T|^2}{|\xi + \eta|^2}\right]
\ge 1 - (\alpha - 2)\left|\cos^2\theta(\xi, y) - \cos^2\theta(\xi + \eta, y)\right|
\ge 1/2,

where in the last inequality we have used the condition [\(36\)](#page-31-0). Finally, since $F_n(\xi)$ is smooth and the domain is bounded, the existence of c_2 is obvious. \square

Remark A.2 By the proof, we can define

$$
\bar{\theta}_j := \sup \{ \theta(\xi, \eta) : (\xi, \eta) \in \mathcal{R}_j \times \mathcal{R}_j, |\xi| \ge |\eta| \},
$$

$$
\bar{\theta}_0 := \min \{ \bar{\theta}_j : j = 1, 2, \dots, \bar{K}_{d,\alpha} \}.
$$

Using these notations, Proposition [A.1](#page-30-0) implies that if $\theta(\xi, \eta) \leq \theta_0$ with $|\xi| \geq |\eta|$, then there holds

$$
c_1 E_d \leq \text{Hess} F_\eta(\xi) \leq c_2 E_d.
$$

Acknowledgements The authors would like to thank René Quilodrán for bring their attention to the result [\[36](#page-34-10), Proposition 6.9], and thank Shuanglin Shao for valuable conversations. The first author acknowledges the support from University of Chinese Academy of Sciences Joint Training Program. And this work is completed during the first author's visit to the University of Kansas whose hospitality is also appreciated.

Funding This work was supported by the National Key R&D Program of China [Grant No. 2023YFC3007303], National Natural Science Foundation of China [Grant Nos. 12271501 & 12071052] and China Postdoctoral Science Foundation [Grant No. GZB20230812].

References

- 1. Alazard, T., Burq, N., Zuily, C.: A stationary phase type estimate. Proc. Am. Math. Soc. **145**(7), 2871–2880 (2017). <https://doi.org/10.1090/proc/13199>
- 2. Bahouri, H., Gérard, P.: High frequency approximation of solutions to critical nonlinear wave equations. Am. J. Math. **121**(1), 131–175 (1999). <https://doi.org/10.1353/ajm.1999.0001>
- 3. Bégout, P., Vargas, A.: Mass concentration phenomena for the *L*2-critical nonlinear Schrödinger equation. Trans. Am. Math. Soc. **359**(11), 5257–5282 (2007). [https://doi.org/10.1090/S0002-9947-07-](https://doi.org/10.1090/S0002-9947-07-04250-X) [04250-X](https://doi.org/10.1090/S0002-9947-07-04250-X)
- 4. Biswas, C., Stovall, B.: Existence of extremizers for Fourier restriction to the moment curve. Trans. Am. Math. Soc. **376**(5), 3473–3492 (2023). <https://doi.org/10.1090/tran/8872>
- 5. Brézis, H., Lieb, E.H.: A relation between pointwise convergence of functions and convergence of functionals. Proc. Am. Math. Soc. **88**(3), 486–490 (1983). <https://doi.org/10.2307/2044999>
- 6. Brocchi, G., Oliveira e Silva, D., Quilodrán, R.: Sharp Strichartz inequalities for fractional and higherorder Schrödinger equations. Anal. PDE **13**(2), 477–526 (2020). [https://doi.org/10.2140/apde.2020.](https://doi.org/10.2140/apde.2020.13.477) [13.477](https://doi.org/10.2140/apde.2020.13.477)
- 7. Carles, R., Keraani, S.: On the role of quadratic oscillations in nonlinear Schrödinger equations. II. The *L*2-critical case. Trans. Am. Math. Soc. **359**(1), 33–62 (2007). [https://doi.org/10.1090/S0002-9947-](https://doi.org/10.1090/S0002-9947-06-03955-9) [06-03955-9](https://doi.org/10.1090/S0002-9947-06-03955-9)
- 8. Carneiro, E.: A sharp inequality for the Strichartz norm. Int. Math. Res. Not. IMRN **2009**(16), 3127– 3145 (2009). <https://doi.org/10.1093/imrn/rnp045>
- 9. Carneiro, E., Oliveira e Silva, D., Sousa, M.: Extremizers for Fourier restriction on hyperboloids. Ann. Inst. H. Poincaré C Anal. Non Linéaire **36**(2), 389–415 (2019). [https://doi.org/10.1016/j.anihpc.2018.](https://doi.org/10.1016/j.anihpc.2018.06.001) [06.001](https://doi.org/10.1016/j.anihpc.2018.06.001)
- 10. Carneiro, E., Oliveira e Silva, D., Sousa, M., Stovall, B.: Extremizers for adjoint Fourier restriction on hyperboloids: the higher-dimensional case. Indiana Univ. Math. J. **70**(2), 535–559 (2021). [https://doi.](https://doi.org/10.1512/iumj.2021.70.8323) [org/10.1512/iumj.2021.70.8323](https://doi.org/10.1512/iumj.2021.70.8323)
- 11. Christ, M., Quilodrán, R.: Gaussians rarely extremize adjoint Fourier restriction inequalities for paraboloids. Proc. Am. Math. Soc. **142**(3), 887–896 (2014). [https://doi.org/10.1090/S0002-9939-](https://doi.org/10.1090/S0002-9939-2013-11827-7) [2013-11827-7](https://doi.org/10.1090/S0002-9939-2013-11827-7)
- 12. Christ, M., Shao, S.: Existence of extremals for a Fourier restriction inequality. Anal. PDE **5**(2), 261– 312 (2012). <https://doi.org/10.2140/apde.2012.5.261>
- 13. Di, B., Yan, D.: Extremals for α-Strichartz inequalities. J. Geom. Anal. **33**(4), 136 (2023). [https://doi.](https://doi.org/10.1007/s12220-022-01185-7) [org/10.1007/s12220-022-01185-7](https://doi.org/10.1007/s12220-022-01185-7)
- 14. Fanelli, L., Vega, L., Visciglia, N.: On the existence of maximizers for a family of restriction theorems. Bull. Lond. Math. Soc. **43**(4), 811–817 (2011). <https://doi.org/10.1112/blms/bdr014>
- 15. Fanelli, L., Vega, L., Visciglia, N.: Existence of maximizers for Sobolev–Strichartz inequalities. Adv. Math. **229**(3), 1912–1923 (2012). <https://doi.org/10.1016/j.aim.2011.12.012>
- 16. Flock, T.C., Stovall, B.: On extremizing sequences for adjoint Fourier restriction to the sphere. arXiv Preprint (2022). [arXiv:2204.10361.](http://arxiv.org/abs/2204.10361) <https://doi.org/10.48550/arXiv.2204.10361>
- 17. Foschi, D.: Maximizers for the Strichartz inequality. J. Eur. Math. Soc. (JEMS) **9**(4), 739–774 (2007). <https://doi.org/10.4171/JEMS/95>
- 18. Foschi, D.: Global maximizers for the sphere adjoint Fourier restriction inequality. J. Funct. Anal. **268**(3), 690–702 (2015). <https://doi.org/10.1016/j.jfa.2014.10.015>
- 19. Foschi, D., Oliveira e Silva, D.: Some recent progress on sharp Fourier restriction theory. Anal. Math. **43**(2), 241–265 (2017). <https://doi.org/10.1007/s10476-017-0306-2>
- 20. Frank, R.L., Lieb, E.H., Sabin, J.: Maximizers for the Stein–Tomas inequality. Geom. Funct. Anal. **26**(4), 1095–1134 (2016). <https://doi.org/10.1007/s00208-018-1695-7>
- 21. Frank, R.L., Sabin, J.: Extremizers for the Airy–Strichartz inequality. Math. Ann. **372**(3–4), 1121–1166 (2018). <https://doi.org/10.1007/s00039-016-0380-9>
- 22. Hundertmark, D., Shao, S.: Analyticity of extremizers to the Airy–Strichartz inequality. Bull. Lond. Math. Soc. **44**(2), 336–352 (2012). <https://doi.org/10.1112/blms/bdr098>
- 23. Hundertmark, D., Zharnitsky, V.: On sharp Strichartz inequalities in low dimensions. Int. Math. Res. Not. **1–18**, 2006 (2006). <https://doi.org/10.1155/IMRN/2006/34080>
- 24. Jiang, J.-C., Pausader, B., Shao, S.: The linear profile decomposition for the fourth order Schrödinger equation. J. Differ. Equ. **249**(10), 2521–2547 (2010). <https://doi.org/10.1016/j.jde.2010.06.014>
- 25. Jiang, J.-C., Shao, S., Stovall, B.: Linear profile decompositions for a family of fourth order schrödinger equations. arXiv Preprint (2017). [arXiv:1410.7520.](http://arxiv.org/abs/1410.7520) <https://doi.org/10.48550/arXiv.1410.7520>
- 26. Kenig, C.E., Ponce, G., Vega, L.: Oscillatory integrals and regularity of dispersive equations. Indiana Univ. Math. J. **40**(1), 33–69 (1991). <https://doi.org/10.1512/iumj.1991.40.40003>
- 27. Keraani, S.: On the defect of compactness for the Strichartz estimates of the Schrödinger equations. J. Differ. Equ. **175**(2), 353–392 (2001). <https://doi.org/10.1006/jdeq.2000.3951>
- 28. Killip, R., Stovall, B., Visan, M.: Scattering for the cubic Klein–Gordon equation in two space dimensions. Trans. Am. Math. Soc. **364**(3), 1571–1631 (2012). [https://doi.org/10.1090/S0002-9947-2011-](https://doi.org/10.1090/S0002-9947-2011-05536-4) [05536-4](https://doi.org/10.1090/S0002-9947-2011-05536-4)
- 29. Killip, R., Visan, M.: Nonlinear Schrödinger equations at critical regularity. In: Evolution Equations. Clay Mathematics Proceedings, vol. 17, pp. 325–437. American Mathematical Society, Providence (2013). <https://www.claymath.org/library/proceedings/cmip017c.pdf#page=333>
- 30. Kunze, M.: On the existence of a maximizer for the Strichartz inequality. Commun. Math. Phys. **243**(1), 137–162 (2003). <https://doi.org/10.1007/s00220-003-0959-5>
- 31. Lieb, E.H.: Sharp constants in the Hardy–Littlewood–Sobolev and related inequalities. Ann. Math. (2) **118**(2), 349–374 (1983). <https://doi.org/10.2307/2007032>
- 32. Lions, P.-L.: The concentration-compactness principle in the calculus of variations. The locally compact case. I. Ann. Inst. H. Poincaré Anal. Non Linéaire **1**(2), 109–145 (1984). [https://doi.org/10.1016/](https://doi.org/10.1016/S0294-1449(16)30428-0) [S0294-1449\(16\)30428-0](https://doi.org/10.1016/S0294-1449(16)30428-0)
- 33. Lions, P.-L.: The concentration-compactness principle in the calculus of variations. The locally compact case. II. Ann. Inst. H. Poincaré Anal. Non Linéaire **1**(4), 223–283 (1984). [https://doi.org/10.1016/](https://doi.org/10.1016/S0294-1449(16)30422-X) [S0294-1449\(16\)30422-X](https://doi.org/10.1016/S0294-1449(16)30422-X)
- 34. Merle, F., Vega, L.: Compactness at blow-up time for *L*² solutions of the critical nonlinear Schrödinger equation in 2D. Int. Math. Res. Not. **8**, 399–425 (1998). <https://doi.org/10.1155/S1073792898000270>
- 35. Negro, G., Oliveira e Silva, D., Thiele, C.: When does $e^{-|\tau|}$ maximize Fourier extension for a conic section? In: Harmonic Analysis and Convexity. Advanced Analytic Geometry, vol. 9, pp. 391–426. De Gruyter, Berlin (2023). <https://doi.org/10.1515/9783110775389-009>
- 36. Oliveira e Silva, D., Quilodrán, R.: On extremizers for Strichartz estimates for higher order Schrödinger equations. Trans. Am. Math. Soc. **370**(10), 6871–6907 (2018). <https://doi.org/10.1090/tran/7223>
- 37. Ramos, J.: A refinement of the Strichartz inequality for the wave equation with applications. Adv. Math. **230**(2), 649–698 (2012). <https://doi.org/10.1016/j.aim.2012.02.020>
- 38. Shao, S.: The linear profile decomposition for the airy equation and the existence of maximizers for the airy Strichartz inequality. Anal. PDE **2**(1), 83–117 (2009). <https://doi.org/10.2140/apde.2009.2.83>
- 39. Shao, S.: Maximizers for the Strichartz and the Sobolev–Strichartz inequalities for the Schrödinger equation. Electron. J. Differ. Equ. **2009**(3), 1–13 (2009)
- 40. Shao, S.: On existence of extremizers for the Tomas–Stein inequality for \mathbb{S}^1 . J. Funct. Anal. **270**(10), 3996–4038 (2016). <https://doi.org/10.1016/j.jfa.2016.02.019>
- 41. Stein, E.M.: Harmonic Analysis: Real-Variable Methods, Orthogonality, and Oscillatory Integrals. Princeton Mathematical Series, vol. 43. Princeton University Press, Princeton (1993). (**With the assistance of Murphy, T.S, Monographs in Harmonic Analysis, III**)
- 42. Stovall, B.: Extremizability of Fourier restriction to the paraboloid. Adv. Math. **360**, 1–18 (2020). <https://doi.org/10.1016/j.aim.2019.106898>
- 43. Tao, T.: A sharp bilinear restrictions estimate for paraboloids. Geom. Funct. Anal. **13**(6), 1359–1384 (2003). <https://doi.org/10.1007/s00039-003-0449-0>
- 44. Tao, T.: A pseudoconformal compactification of the nonlinear Schrödinger equation and applications. N. Y. J. Math. **15**, 265–282 (2009)
- 45. Tao, T., Vargas, A., Vega, L.: A bilinear approach to the restriction and Kakeya conjectures. J. Am. Math. Soc. **11**(4), 967–1000 (1998). <https://doi.org/10.1090/S0894-0347-98-00278-1>

Publisher's Note Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.

Springer Nature or its licensor (e.g. a society or other partner) holds exclusive rights to this article under a publishing agreement with the author(s) or other rightsholder(s); author self-archiving of the accepted manuscript version of this article is solely governed by the terms of such publishing agreement and applicable law.