Z. angew. Math. Phys. 56 (2005) 592–608
0044-2275/05/040592-17
DOI 10.1007/s00033-004-2061-z
(c) 2005 Birkhäuser Verlag, Basel

Zeitschrift für angewandte Mathematik und Physik ZAMP

On the existence of periodic solutions of Rayleigh equations

Zaihong Wang

Abstract. In this paper, we study the existence of periodic solutions of Rayleigh equation

$$x'' + f(x') + g(x) = p(t)$$

where f, g are continuous functions and p is a continuous and 2π -periodic function. We prove that the given equation has at least one 2π -periodic solution provided that f(x) is sublinear and the time map of equation x'' + g(x) = 0 satisfies some nonresonant conditions. We also prove that this equation has at least one 2π -periodic solution provided that g(x) satisfies $\lim_{|x|\to+\infty} sgn(x)g(x) = +\infty$ and f(x) satisfies $sgn(x)(f(x) - p(t)) \ge c$, for $t \in \mathbf{R}$, $|x| \ge d$ with c, d being positive constants.

Mathematics Subject Classification (2000). 34C15, 34C25.

Keywords. Rayleigh equation, periodic solution, continuation theorem, Lyapunov function.

1. Introduction

We are concerned with the existence of 2π -periodic solutions of Rayleigh equation

$$x'' + f(x') + g(x) = p(t),$$
(1.1)

where $f, g: \mathbf{R} \to \mathbf{R}$ are continuous, $p: \mathbf{R} \to \mathbf{R}$ is continuous and 2π -periodic.

Arising from nonlinear oscillations, Eq.(1.1) has been studied by many authors (see [1-4, 9, 14] and the references therein). In [14], using the method of upper and lower solutions, Habets and Torres studied the existence and multiplicity of 2π -periodic solutions of Eq.(1.1) by assuming that g = g(t, x, x') is bounded (or bounded from below) and other conditions hold.

When $f(x) \equiv 0$, Eq.(1.1) is a conservative system. It is well known that time map plays a crucial role in dealing with the existence and multiplicity of periodic solutions of equation x'' + g(x) = p(t). Assume that g(x) satisfies the following

Research supported by the National Natural Science Foundation of China, No.10001025 and No.10471099, Natural Science Foundation of Beijing, No. 1022003 and by a postdoctoral Grant of University of Torino, Italy.

condition,

(H₁)
$$\lim_{|x| \to +\infty} sgn(x)g(x) = +\infty.$$

Let $G(x) = \int_0^x g(u) du$. Define a function $\tau(c)$ as follows,

$$\tau(c) = \sqrt{2} \left| \int_0^c \frac{du}{\sqrt{G(c) - G(u)}} \right|$$

which is usually called time map and was first introduced in [13]. From condition (H_1) we know that $\tau(c)$ is continuous for |c| large enough. Now, we introduce notations,

$$\tau_{+} = \liminf_{c \to +\infty} \tau(c), \qquad \tau^{+} = \limsup_{c \to +\infty} \tau(c),$$

$$\tau_{-} = \liminf_{c \to -\infty} \tau(c), \qquad \tau^{-} = \limsup_{c \to -\infty} \tau(c).$$

By the asymptotic property of $\tau(c)$, Ding and Zanolin [5] proved that equation x'' + g(x) = p(t) has at least one 2π -periodic solution provided that condition (H_1) holds and there is an integer n > 0 such that

(H₂)
$$\frac{2\pi}{n+1} < \tau_{-} + \tau_{+} \le \tau^{-} + \tau^{+} < \frac{2\pi}{n}.$$

Wang [8] generalized this result to Liénard equation x'' + f(x)x' + g(x) = p(t). He proved that this Liénard equation has at least one 2π -periodic solution provided that conditions (H_1) , (H_2) hold and F(x) is bounded, where $F(x) = \int_0^x f(u) du$.

One aim of this paper is to prove the existence of 2π -periodic solutions of Eq.(1.1) provided that conditons (H_1) , (H_2) hold and f(x) is sublinear. Assume that f(x) satisfies

(H₃)
$$\lim_{|x|\to+\infty} f(x)/x = 0.$$

Since f(x) maybe unbounded under condition (H_3) , the method in [8] is invalid under present situation. By developing a different estimate method, we overcome the difficulty in estimating time owing to the possible unboundedness of f(x). We obtain

Theorem 1. Assume that conditions (H_1) , (H_2) and (H_3) hold. Then Eq.(1.1) has at least one 2π -periodic solution.

If the condition (H_2) is replaced by a condition as follows,

then we also have

Theorem 2. Assume that conditions (H_1) , (H_3) and (H_4) hold. Then Eq.(1.1) has at least one 2π -periodic solution.

Z. Wang

Another aim of this paper is to prove the existence of periodic solutions of Eq.(1.1) provided that f(x) satisfies sign condition. It was proved in [11] that Liénard equation

$$x'' + f(x)x' + g(x) = p(t)$$

has at least one 2π -periodic solution provided that the following conditions hold, (i) f(x) is continuous and $\lim_{|x|\to+\infty} sgn(x)F(x) = +\infty$ (or $-\infty$),

(ii) g(x) is locally Lipschitz continuous and $sgn(x)g(x) \ge 0$, $|x| \ge c_0$, where c_0 is a positive constant and $\int_0^{2\pi} p(t)dt = 0$.

Obviously, the same conclusion still holds if the condition (ii) is replaced by the condition as follows,

(ii)' g(x) is locally Lipschitz continuous and $sgn(x)(g(x) - \bar{p}) \ge 0$, $|x| \ge c_0$, where c_0 is a positive constant and $\bar{p} = (1/2\pi) \int_0^{2\pi} p(t) dt$. For Rayleigh equation, we have a similar result. Assume that f(x) satisfies the

For Rayleigh equation, we have a similar result. Assume that f(x) satisfies the following sign condition,

$$(H_5) \qquad \qquad sgn(x)(f(x) - p(t)) \ge c, \quad \forall t \in \mathbf{R}, \quad |x| \ge d$$

with c, d being arbitrary positive constants. We prove

Theorem 3. Assume that f, g are locally Lipschitz continuous and conditions (H_1) , (H_5) hold. Then Eq.(1.1) has at least one 2π -periodic solution.

Throughout this paper, we always use \mathbf{R} to denote the whole real number set.

2. Periodic solutions via continuation theorem

In this section, we deal with the existence of periodic solutions of (1.1) under conditions (H_1) , (H_2) and (H_3) or (H_1) , (H_3) and (H_4) . Consider an equivalent system of Eq.(1.1),

$$x' = y, \quad y' = -g(x) - f(y) + p(t).$$
 (2.1)

In order to use the Continuation Theorem [6, Theorem 2], we embed (2.1) into a system family with one parameter $\lambda \in [0, 1]$,

$$x' = y, \quad y' = -g(x) - \lambda f(y) + \lambda p(t).$$
 (2.2)

Let $(x_{\lambda}(t), y_{\lambda}(t)) = (x(t, x_0, y_0, \lambda), y(t, x_0, y_0, \lambda))$ be a solution of (2.2) satisfying the initial value condition $(x_{\lambda}(0), y_{\lambda}(0)) = (x_0, y_0)$. In what follows, for simplicity, we always use (x(t), y(t)) to denote $(x_{\lambda}(t), y_{\lambda}(t))$. We have the following lemma.

Lemma 1. Assume that conditions (H_1) , (H_3) hold. Then every solution (x(t), y(t)) of (2.2) exists on the whole t-axis.

Proof. Define a function

$$V(x,y) = \frac{1}{2}y^2 + G(x).$$

Set

$$v(t) = V(x(t), y(t)) = \frac{1}{2}y(t)^2 + G(x(t)).$$

Then we have

$$v'(t) = -\lambda [f(y(t)) - p(t)]y(t).$$
(2.3)

Since $\lim_{|x|\to\infty} f(x)/x = 0$, there exists a constant a > 0 such that

$$|f(x)x| \le x^2, \quad |x| \ge a$$

Furthermore, there exists a positive constant b such that

$$|f(x)x| \le x^2 + b, \quad x \in \mathbf{R}.$$
(2.4)

It follows from (2.3) and (2.4) that

$$|v'(t)| \le y(t)^2 + b + \frac{1}{2}y(t)^2 + \frac{1}{2}p(t)^2 \le \frac{3}{2}y(t)^2 + M,$$
(2.5)

with $M = b + M_p^2/2$, $M_p = max\{|p(t)| : t \in \mathbf{R}\}$. From (H_1) we know that there exists a positive constant M_0 such that

$$G(x) + M_0 \ge 0, \quad x \in \mathbf{R},$$

which, together with (2.5), implies that

$$|v'(t)| \le \frac{3}{2}y(t)^2 + 3G(x(t)) + \bar{M},$$

with $\overline{M} = M + 3M_0$. Thus, we have that $|v'(t)| \leq 3v(t) + \overline{M}$. Hence,

$$v'(t) \le 3v(t) + \bar{M}$$

Multiplying both sides of this inequality by e^{-3t} and integrating over any bounded interval $[0, T_0)$ $(T_0 > 0)$ we have that

$$v(t) \le v(0)e^{3T_0} + \frac{1}{3}\bar{M}(e^{3T_0} - 1), \quad t \in [0, T_0).$$

Therefore, there is no blow-up for solution (x(t), y(t)) on any bounded interval. Furthermore, every solution (x(t), y(t)) of (2.2) exists on the whole *t*-axis.

On the basis of Lemma 1, we have

Lemma 2. Under conditions (H_1) , (H_3) . Then there is a nondecreasing function $\sigma : \mathbf{R}^+ \to \mathbf{R}^+$, with $\sigma(s) \ge s$, for all s > 0, such that for any $\lambda \in [0, 1]$, r > 0 and each solution (x(t), y(t)) of (2.2), the following conclusion holds, (i) If $(x_0^2 + y_0^2)^{1/2} \le r$, then $(x(t)^2 + y(t)^2)^{1/2} \le \sigma(r)$, for $t \in [0, 2\pi]$. (ii) If $(x_0^2 + y_0^2)^{1/2} \ge \sigma(r)$, then $(x(t)^2 + y(t)^2)^{1/2} \ge r$, for $t \in [0, 2\pi]$.

This lemma can be proved by standard methods [7, 12].

Lemma 3. Assume that (H_1) , (H_3) hold. Then there exists a constant $R_0 > 0$ such that if (x(t), y(t)) is a 2π -periodic solution of (2.2) with $x_0^2 + y_0^2 \ge R_0^2$ and

 $x_1 = x(t_1)$ is a local maximum of x(t) and $x_2 = x(t_2)$ is a local minimum of x(t), then

$$x_1 > 0, \quad x_2 < 0.$$

Proof. We only give the proof of $x_1 > 0$. The other case can be treated similarly. Assume by contradiction that $x_1 \leq 0$. Let A > 0 be a constant satisfying

$$|f(0)| + |p(t)| + 1 \le A, \quad t \in \mathbf{R}.$$

It follows from (H_1) that there exists a constant a > 0 such that

$$sgn(x)g(x) \ge A, \quad |x| \ge a.$$
 (2.6)

Define $R_0 = \sigma(a)$, where σ is defined in Lemma 2. Since $x_1 = x(t_1)$ is a local maximum of x(t), we know that $y(t_1) = x'(t_1) = 0$. Therefore,

$$y'(t_1) = -g(x(t_1)) - \lambda f(0) - \lambda p(t_1).$$

From Lemma 2 we have that if $x_0^2 + y_0^2 \ge R_0^2$, then

$$x(t)^2 + y(t)^2 \ge a^2, \quad t \in [0, 2\pi].$$

Therefore,

$$x(t_1)^2 \ge a^2,$$

which implies that $x(t_1) \ge a$ or $x(t_1) \le -a$. By the hypothesis $x_1 \le 0$ we have that $x(t_1) \le -a$, which, together with (2.6), implies that $y'(t_1) > 0$. Then we have $x''(t_1) = y'(t_1) > 0$. From the continuity of x''(t) we know that there exists an interval (α, β) such that $t_1 \in (\alpha, \beta)$ and x''(t) > 0, for $t \in (\alpha, \beta)$. Since $x'(t_1) = 0$, we get that

$$x'(t) < 0, t \in (\alpha, t_1); \quad x'(t) > 0, t \in (t_1, \beta).$$

Thus we obtain that

$$x(t_1) < x(t), t \in (\alpha, t_1); \quad x(t_1) < x(t), t \in (t_1, \beta).$$

This contradicts with the fact that x_1 is a local maximum of x(t).

It follows from Lemma 2 that if $x_0^2 + y_0^2$ is large enough, then we can introduce the polar coordinates. Set $x = r \cos \theta$, $y = r \sin \theta$. Under this transformation, (2.2) becomes

$$\begin{cases} \frac{dr}{dt} = -r\sin\theta\cos\theta - g(r\cos\theta)\sin\theta - \lambda f(r\sin\theta)\sin\theta + \lambda p(t)\sin\theta\\ \frac{d\theta}{dt} = -\sin^2\theta - \frac{g(r\cos\theta)\cos\theta}{r} - \frac{\lambda f(r\sin\theta)\cos\theta}{r} + \frac{\lambda p(t)\cos\theta}{r}. \end{cases}$$

Denote by $(r(t), \theta(t)) = (r(t, r_0, \theta_0, \lambda), \theta(t, r_0, \theta_0, \lambda))$ the solution of above system through the initial point (r_0, θ_0) .

Lemma 4. Assume that (H_1) , (H_3) hold and $\delta(0 < \delta < 1)$ is a given constant. Then there exists $R_{\delta} > 0$ such that for $r_0 \ge R_{\delta}$ and $(r(t) \cos \theta(t), r(t) \sin \theta(t)) \in \{(x, y) : |y| \ge \delta |x|\}$, the following inequality holds,

$$\frac{d\theta(t)}{dt} < 0, \quad t \in [0, 2\pi].$$

Proof. It follows from (H_1) that there exists a constant $a_0 > 0$ such that

 $(g(x) - \lambda p(t))x \ge 0, \quad |x| \ge a_0, t \in \mathbf{R}, \lambda \in [0, 1].$

Therefore,

$$\frac{(g(r\cos\theta) - \lambda p(t))\cos\theta}{r} \ge 0, \quad |r\cos\theta| \ge a_0, t \in \mathbf{R}, \lambda \in [0, 1].$$
(2.7)

On the other hand, since g(x) is bounded in interval $[-a_0, a_0]$, we have

$$-\frac{\epsilon_{\delta}}{4} \le \frac{(g(r\cos\theta) - \lambda p(t))\cos\theta}{r} \le \frac{\epsilon_{\delta}}{4}, \quad |r\cos\theta| \le a_0, t \in \mathbf{R}, \lambda \in [0, 1]$$
(2.8)

for r > 0 large enough, where $\epsilon_{\delta} = \sin^2(\arctan \delta) > 0$. From (H_3) we have that

$$-\frac{\epsilon_{\delta}}{4} \le \frac{f(r\sin\theta)\cos\theta}{r} \le \frac{\epsilon_{\delta}}{4}, \quad \theta \in \mathbf{R}$$
(2.9)

for r > 0 large enough. If $(r(t)\cos\theta(t), r(t)\sin\theta(t)) \in \{(x, y) : |y| \ge \delta |x|\}$ and $|r(t)\cos\theta(t)| \ge a_0, t \in [0, 2\pi]$, then it follows from Lemma 2 and (2.7), (2.9) that

$$\frac{d\theta(t)}{dt} \le -\epsilon_{\delta} + \frac{\epsilon_{\delta}}{4} < 0$$

with r_0 large enough. If $(r(t)\cos\theta(t), r(t)\sin\theta(t)) \in \{(x, y) : |y| \ge \delta |x|\}$ and $|r(t)\cos\theta(t)| \le a_0, t \in [0, 2\pi]$, then it follows from Lemma 2 and (2.8), (2.9) that

$$\frac{d\theta(t)}{dt} \le -\epsilon_{\delta} + \frac{\epsilon_{\delta}}{2} < 0$$

with r_0 large enough. Thus, we have reached the conclusion.

Let (x(t), y(t)) be a 2π -periodic solution of (2.2) with $r_0 = \sqrt{x_0^2 + y_0^2} \ge \max\{R_0, R_\delta\}$, which has polar coordinates expression $(r(t), \theta(t))$. Then we can define the rotation number as follows,

$$n(x,y) = \frac{\theta(0) - \theta(2\pi)}{2\pi} = \frac{1}{2\pi} \int_0^{2\pi} \frac{x'(t)y(t) - x(t)y'(t)}{x^2(t) + y^2(t)} dt.$$

Assume that (H_1) , (H_3) hold. From Lemma 3 that there exists some $t_0 \in [0, 2\pi]$ such that $x(t_0) = 0$. According to Lemma 4, if $(x(t), y(t)) \in \{(x, y) : |y| \ge \delta |x|\}$ and $t \in [0, 2\pi]$, then $\theta(t)$ decreases strictly. Therefore, there exists $t_1 \in [t_0, t_0 + 2\pi]$ such that $y(t_1) = \delta x(t_1)$, $x(t_1) > 0$ and $y(t) \ge \delta x(t)$, $t \in [t_0, t_1]$. Since (x(t), y(t))is 2π -periodic, the solution (x(t), y(t)) must leave the region $\{(x, y) : |y| \le \delta |x|\}$

after it enters this region. Hence, there exists some $t_2 \in [t_0, t_0 + 2\pi]$ such that $y(t_2) = -\delta x(t_2)$ and $(x(t), y(t)) \in \{(x, y) : |y| \le \delta |x|, x > 0\}, t \in [t_1, t_2]$. Thus, we have that

$$\theta(t_1) - \theta(t_2) = 2 \arctan \delta > 0.$$

Every time when the solution (x(t), y(t)) goes through the region $\{(x, y) : |y| \le \delta |x|\}$, we have the same result. Recalling that $\theta(t)$ decreases strictly when $(x(t), y(t)) \in \{(x, y) : |y| \ge \delta |x|\}$, we obtain that $\theta(0) - \theta(2\pi) > 0$, which implies

$$n(x,y) \in \mathbf{N}$$

for $r_0 = \sqrt{x_0^2 + y_0^2}$ large enough.

Lemma 5. Assume that conditions (H_1) , (H_2) and (H_3) hold. Let $k_0 > 0$ be a fixed integer. Suppose that there exists a sequence of 2π -periodic solutions $\{(x_j(t), y_j(t))\}_{j=1}^{\infty}$ of (2.2), with rotation numbers $n(x_j, y_j) = k_0, j = 1, 2, \cdots$, such that

$$\lim_{i \to +\infty} (x_j^2(t) + y_j^2(t)) = +\infty_j$$

then

$$k_0(\tau_+ + \tau_-) \le 2\pi.$$

Proof. For simplicity, we assume that

 $sgn(x)g(x) > 0, \quad x \in \mathbf{R}, x \neq 0.$

Let (x(t), y(t)) be any one of $(x_j(t), y_j(t))$ with j large enough. Then there exist constants $t_1^1 < t_2^1 < t_3^1 = t_1^2 < t_2^2 < t_3^2 = \cdots = t_1^{k_0} < t_2^{k_0} < t_3^{k_0} = t_1^1 + 2\pi$ such that

$$x(t_1^i) = 0; \quad x(t_2^i) = 0; \quad x(t_3^i) = 0; \quad i = 1, 2, \cdots, k_0$$

and

$$x(t) \ge 0, t \in [t_1^i, t_2^i]; \quad x(t) \le 0, t \in [t_2^i, t_3^i], \quad i = 1, 2, \cdots, k_0.$$

For simplicity, let (α, β) $(\alpha < \beta)$ denote any couple of (t_1^i, t_2^i) $(i = 1, 2, \dots, k_0)$. Set $x_* = x(t_*) = \max\{x(t) : \alpha \le x \le \beta\}$. In what follows, we shall estimate $t_* - \alpha$ and $\beta - t_*$, respectively. At first, we estimate the former one. It can be inferred from the first equation of (2.2) that $y(t) \ge 0$, for $t \in [\alpha, t_*]$. From condition (H_3) we know that for any sufficiently small $\varepsilon > 0$, there exists $a_{\varepsilon} > 0$ such that

$$|f(x)| \le \varepsilon |x|, \quad |x| \ge a_{\varepsilon},$$

which implies that

$$f(x)x| \le \varepsilon x^2, \quad |x| \ge a_{\varepsilon}.$$

Thus, there exists a constant $b_{\varepsilon} > 0$ such that

$$|f(x)x| \le \varepsilon x^2 + b_{\varepsilon}, \quad x \in \mathbf{R}.$$
(2.10)

Periodic solutions of Rayleigh equations

Multiplying both sides of $y' = -g(x) - \lambda f(y) + \lambda p(t)$ by y(t) and applying x'(t) =y(t), we have

$$y(t)y'(t) = -g(x(t))x'(t) - \lambda f(y(t))y(t) + \lambda p(t)y(t).$$
(2.11)

Integrating both sides (2.11) over interval $[t, t_*]$ with $\alpha \leq t \leq t_*$ yields

$$\int_{t}^{t_{*}} y(\tau)y'(\tau)d\tau = -\int_{t}^{t_{*}} g(x(\tau))x'(\tau)d\tau - \lambda \int_{t}^{t_{*}} f(y(\tau))y(\tau)dt + \lambda \int_{t}^{t_{*}} p(\tau)y(\tau)d\tau.$$

Since $y(t_{*}) = x'(t_{*}) = 0$, we have that, for $\alpha \le t \le t_{*}$,

$$y^{2}(t) = 2(G(x(t_{*})) - G(x(t))) + 2\lambda \int_{t}^{t_{*}} f(y(\tau))y(\tau)d\tau - 2\lambda \int_{t}^{t_{*}} p(\tau)y(\tau)d\tau.$$
(2.12)

Combining (2.10) and (2.12) we get that

$$y^{2}(t) \leq 2(G(x(t_{*})) - G(x(t))) + 2\varepsilon \int_{t}^{t_{*}} y^{2}(\tau)d\tau + 2M_{p} \int_{t}^{t_{*}} y(\tau)d\tau + 4b_{\varepsilon}\pi.$$
(2.13)
with $M_{\varepsilon} = \max\{|p(t)|: t \in \mathbf{R}\}$ Write

with $M_p = \max\{|p(t)| : t \in \mathbf{R}\}$. Write

$$\Phi(t) = \int_t^{t_*} y^2(\tau) d\tau.$$

Then

$$\Phi'(t) = -y^2(t).$$

Therefore, we have that

$$-\Phi'(t) - 2\varepsilon \Phi(t) \le 2(G(x(t_*)) - G(x(t))) + 2M_p(x(t_*) - x(t)) + 4b_{\varepsilon}\pi.$$

Multiplying both sides of above inequality by $e^{2\varepsilon t}$ and integrating over interval $[t, t_*]$ yields

$$-\int_{t}^{t_{*}} [\Phi(\tau)e^{2\varepsilon\tau}]' d\tau \leq \int_{t}^{t_{*}} [2(G(x(t_{*})) - G(x(\tau))) + 2M_{p}(x(t_{*}) - x(\tau)) + 4b_{\varepsilon}\pi]e^{2\varepsilon\tau} d\tau.$$

Since $\Phi(t_*) = 0$, we have

$$\Phi(t)e^{2\varepsilon t} \le \int_{t}^{t_{*}} [2(G(x(t_{*})) - G(x(\tau))) + 2M_{p}(x(t_{*}) - x(\tau)) + 4b_{\varepsilon}\pi]e^{2\varepsilon\tau}d\tau.$$

On the other hand, from $x'(t) = y(t) \ge 0$ we know that $x(\tau)$ is increasing on the interval $[t, t_*]$. Consequently,

$$\Phi(t)e^{2\varepsilon t} \le e^{2\varepsilon t_*} \int_t^{t_*} [2(G(x(t_*)) - G(x(\tau))) + 2M_p(x(t_*) - x(\tau)) + 4b_\varepsilon \pi]d\tau.$$

Furthermore, for $t \in [\alpha, t_*]$,

$$\Phi(t) \le 2\pi e^{4\pi\varepsilon} [2(G(x(t_*)) - G(x(t))) + 2M_p(x(t_*) - x(t)) + 4b_\varepsilon \pi].$$
(2.14)

It follows from (2.13) and (2.14) that, for $t \in [\alpha, t_*]$,

$$y^{2}(t) \leq 2(G(x(t_{*})) - G(x(t))) + 4\pi\varepsilon e^{4\pi\varepsilon} [2(G(x(t_{*})) - G(x(t))) + 2M_{p}(x(t_{*}) - x(t)) + 4b_{\varepsilon}\pi]$$
(2.15)
+2M_{p}(x(t_{*}) - x(t)) + 4b_{\varepsilon}\pi.

 $+2M_p(x(t_*) - x(t)) + 4\theta_{\varepsilon}\pi.$ Set $\eta = 4\pi\varepsilon e^{4\pi\varepsilon}$. Obviously, $\eta \to 0$ as $\varepsilon \to 0$. By (2.15) we have that

$$y^{2}(t) \leq (1+\eta)[2(G(x(t_{*})) - G(x(t))) + 2M_{p}(x(t_{*}) - x(t)) + 4b_{\varepsilon}\pi].$$

Recalling $x'(t) = y(t)$, we get that, for $t \in [\alpha, t_{*}]$,

$$x'(t) \le \sqrt{1+\eta}\sqrt{2(G(x(t_*)) - G(x(t))) + 2M_p(x(t_*) - x(t)) + 4b_\varepsilon\pi}.$$
 (2.16)

According to (2.16), we have

$$\frac{x'(t)}{\sqrt{1+\eta}\sqrt{2(G(x(t_*))-G(x(t)))+2M_p(x(t_*)-x(t))+4b_{\varepsilon}\pi}} \le 1.$$

Integrating both sides of this inequality over $[\alpha, t_*]$, we obtain that

$$\frac{1}{\sqrt{1+\eta}} \int_0^{x_*} \frac{dx}{\sqrt{2(G(x_*) - G(x)) + 2M_p(x_* - x) + 4b_\varepsilon \pi}} \le t_* - \alpha, \qquad (2.17)$$

where $x_* = x(t_*)$. Take a constant L > 0 and write

$$\int_0^{x_*} \frac{dx}{\sqrt{2(G(x_*) - G(x)) + 2M_p(x_* - x) + 4b_\varepsilon \pi}} = I_1 + I_2$$

with

$$I_1 = \int_0^{x_* - L} \frac{dx}{\sqrt{2(G(x_*) - G(x)) + 2M_p(x_* - x) + 4b_\varepsilon \pi}},$$

$$I_2 = \int_{x_* - L}^{x_*} \frac{dx}{\sqrt{2(G(x_*) - G(x)) + 2M_p(x_* - x) + 4b_\varepsilon \pi}}.$$

If $x \in [0, x_* - L]$, then

 $G(x_*) - G(x) + M_p(x_* - x) \ge G(x_*) - G(x_* - L) + LM_p = [g(\xi) + M_p]L$ (2.18) with $\xi \in [x_* - L, x_*]$. From (2.18) and (H_1) we know that

$$\lim_{x_* \to +\infty} [G(x_*) - G(x) + M_p(x_* - x)] = +\infty.$$

Therefore, I_1 can be expressed in the form

$$I_1 = \int_0^{x_* - L} \frac{dx}{\sqrt{2(G(x_*) - G(x)) + 2M_p(x_* - x)[1 + o(1)]}}$$

for $x_* \to \infty$. From (H_2) we know that $\tau(e)$ is bounded. Thus we obtain

$$I_1 = \int_0^{x_* - L} \frac{dx}{\sqrt{2(G(x_*) - G(x)) + 2M_p(x_* - x)}} + o(1).$$
(2.19)

600

On the other hand, if $x \in [x_* - L, x_*]$, then

$$G(x_*) - G(x) + M_p(x_* - x) \ge [\mu(x_*) + M_p](x_* - x)$$

with $\mu(x_*) = \min\{g(x) : x \in [x_* - L, x_*]\}$. Consequently,

$$I_2 \le \int_{x_*-L}^{x_*} \frac{dx}{\sqrt{2(G(x_*) - G(x)) + 2M_p(x_* - x)}} \le \frac{\sqrt{2L}}{\sqrt{\mu(x_*) + M_p}} = o(1)$$

for $x_* \to \infty$. Furthermore,

$$I_2 = \int_{x_*-L}^{x_*} \frac{dx}{\sqrt{2(G(x_*) - G(x)) + 2M_p(x_* - x)}} + o(1).$$
(2.20)

It follows from (2.17), (2.19) and (2.20) that

$$\frac{1}{\sqrt{1+\eta}} \int_0^{x_*} \frac{dx}{\sqrt{2(G(x_*) - G(x)) + 2M_p(x_* - x)}} + o(1) \le t_* - \alpha$$

for $x_* \to \infty.$ Applying a Lemma in [5, 13] we have that

$$\frac{1}{\sqrt{1+\eta}} \int_0^{x_*} \frac{dx}{\sqrt{2(G(x_*) - G(x))}} + o(1) \le t_* - \alpha,$$

which implies that, for $x_* \to \infty$,

$$\frac{1}{2}\frac{1}{\sqrt{1+\eta}}\tau_{+} + o(1) \le t_{*} - \alpha.$$
(2.21)

Similarly, we have that

$$\frac{1}{2}\frac{1}{\sqrt{1+\eta}}\tau_{+} + o(1) \le \beta - t_{*}.$$
(2.22)

Combining (2.21) and (2.22), we get

$$\frac{1}{\sqrt{1+\eta}}\tau_+ + o(1) \le \beta - \alpha,$$

for $x_* \to \infty$. Sine (α, β) denote any couple of (t_1^i, t_2^i) , $i = 1, 2, \dots, k_0$, we obtain that, for $x_* \to \infty$,

$$\frac{k_0}{\sqrt{1+\eta}}\tau_+ + o(1) \le \sum_{i=1}^{i=k_0} (t_2^i - t_1^i).$$
(2.23)

Using the same methods, we can derive that

$$\frac{k_0}{\sqrt{1+\eta}}\tau_- + o(1) \le \sum_{i=1}^{i=k_0} (t_3^i - t_2^i).$$
(2.24)

From (2.23) and (2.24) we have that

$$\frac{k_0}{\sqrt{1+\eta}}(\tau_+ + \tau_-) + o(1) \le \sum_{i=1}^{i=k_0} (t_2^i - t_1^i) + \sum_{i=1}^{i=k_0} (t_3^i - t_2^i) = 2\pi.$$

Since $\eta \to 0$ as $\varepsilon \to 0$, we get that

$$k_0(\tau_+ + \tau_-) \le 2\pi.$$

Lemma 6. Under the same conditions of Lemma 5. The following conclusion holds,

$$k_0(\tau^+ + \tau^-) \ge 2\pi.$$

Proof. We use the same notations as in Lemma 5. From (2.10) and (2.12) we know that

$$y^{2}(t) \geq 2(G(x(t_{*})) - G(x(t))) - 2\varepsilon \int_{t}^{t_{*}} y^{2}(\tau)d\tau - 2M_{p} \int_{t}^{t_{*}} y(\tau)d\tau - 4\pi b_{\varepsilon},$$

which, together with (2.14), yields

$$\begin{split} y^2(t) &\geq 2(G(x(t_*)) - G(x(t))) \\ &-4\pi\varepsilon e^{4\pi\varepsilon} [2(G(x(t_*)) - G(x(t))) + 2M_p(x(t_*) - x(t)) + 4\pi b_{\varepsilon}] \\ &-2M_p(x(t_*) - x(t)) - 4b_{\varepsilon}\pi. \end{split}$$

Therefore, we have that, for $t \in [\alpha, \beta]$,

 $y^2(t) \ge 2(1-\eta)(G(x(t_*)) - G(x(t))) - 2(1+\eta)M_p(x(t_*) - x(t)) - 4(1+\eta)\pi b_{\varepsilon},$ where $\eta = 4\pi\varepsilon e^{4\pi\varepsilon}$. Let $L_0 > 0$ be a constant. If $x(t) \in [0, x(t_*) - L_0]$, then we have

$$\begin{split} & [(1-\eta)G(x(t_*)) - (1+\eta)M_px(t_*)] - [(1-\eta)G(x(t)) - (1+\eta)M_px(t)] \\ & \geq [(1-\eta)G(x(t_*)) - (1+\eta)M_px(t_*)] \\ & -[(1-\eta)G(x(t_*) - L_0) - (1+\eta)M_p(x(t_*) - L_0)] \\ & = (1-\eta)g(\xi_*)L_0 - (1+\eta)M_pL_0, \quad \xi_* \in [x(t_*) - L_0, x(t_*)]. \end{split}$$

Hence, if $x(t_*)$ is large enough and $x(t) \in [0, x(t_*) - L_0]$, then $y^2(t) \ge 2(1-\eta)(G(x(t_*)) - G(x(t))) - 2(1+\eta)M_p(x(t_*) - x(t)) - 4(1+\eta)\pi b_{\varepsilon} > 0.$ Let $\bar{t}_* \in [\alpha, t_*]$ such that $x(\bar{t}_*) = x(t_*) - L_0$. If $t \in [\alpha, \bar{t}_*]$, then

 $x'(t) \ge \sqrt{2(1-\eta)(G(x(t_*)) - G(x(t))) - 2(1+\eta)M_p(x(t_*) - x(t)) - 4(1+\eta)\pi b_{\varepsilon}}.$ Consequently,

$$\frac{x'(t)}{\sqrt{2(1-\eta)(G(x(t_*))-G(x(t)))-2(1+\eta)M_p(x(t_*)-x(t))-4(1+\eta)\pi b_{\varepsilon}}} \ge 1.$$

602

Integrating both sides of this inequality over $[\alpha, \bar{t}_*]$ results in

$$\int_{0}^{x_{*}-L_{0}} \frac{dx}{\sqrt{2(1-\eta)(G(x_{*})-G(x))-2(1+\eta)M_{p}(x_{*}-x)-4(1+\eta)\pi b_{\varepsilon}}} \ge \bar{t}_{*} - \alpha$$

with $x_* = x(t_*)$. Applying the same methods as in Lemma 5, we have that

$$\int_{0}^{x_{*}-L_{0}} \frac{dx}{\sqrt{2(1-\eta)(G(x_{*})-G(x))-2(1+\eta)M_{p}(x_{*}-x)-4(1+\eta)\pi b_{\varepsilon}}} = \int_{0}^{x_{*}-L_{0}} \frac{dx}{\sqrt{2(1-\eta)(G(x_{*})-G(x))}} + o(1)$$
(2.25)

for $x_* \to \infty$. On the other hand, it is easy to check that

$$\int_{x_*-L_0}^{x_*} \frac{dx}{\sqrt{2(1-\eta)(G(x_*)-G(x))}} = \frac{1}{\sqrt{1-\eta}} \int_{x_*-L_0}^{x_*} \frac{dx}{\sqrt{2(G(x_*)-G(x))}} = o(1).$$
(2.26)

From (2.25) and (2.26) we know that, for $x_* \to \infty$,

$$\frac{1}{\sqrt{1-\eta}} \int_0^{x_*} \frac{dx}{\sqrt{2(G(x_*) - G(x))}} + o(1) \ge \bar{t}_* - \alpha.$$
(2.27)

Next, we estimate $t_* - \overline{t}_*$. Since $x'(t) = y(t) \ge 0$, $t \in [\overline{t}_*, t_*]$, we know that $x(t) \in [x(t_*) - L_0, x(t_*)]$, for $t \in [\overline{t}_*, t_*]$. From x'(t) = y(t) we have that

$$\int_{\bar{t}_*}^{t_*} y(t)dt = \int_{\bar{t}_*}^{t_*} x'(t)dt = x(t_*) - x(\bar{t}_*) = L_0.$$

According to condition (H_3) , for any sufficiently small $\varepsilon > 0$, there exists a c_{ε} such that

$$|f(x)| \le \varepsilon |x| + c_{\varepsilon}, \quad x \in \mathbf{R}.$$
(2.28)

By
$$y'(t) = -g(x(t)) - \lambda f(y(t)) + \lambda p(t)$$
 and (2.28) we get that, for $t \in [\bar{t}_*, t_*]$,
 $y'(t) \leq -g(x(t)) + \varepsilon y(t) + c_\varepsilon + M_p.$ (2.29)

Integrating both sides of (2.29) over interval $[t, t_*]$, with $t \in [\bar{t}_*, t_*]$, we obtain

$$\int_{t}^{t_*} y'(\tau) d\tau \leq -\int_{t}^{t_*} g(x(t)) dt + \varepsilon \int_{t}^{t_*} y(t) dt + 2\pi (c_\varepsilon + M_p) dt$$

Therefore, if $t \in [\bar{t}_*, t_*]$, then

$$y(t) \ge \int_{t}^{t_*} g(x(t))dt - \varepsilon L_0 - 2\pi (c_{\varepsilon} + M_p).$$
(2.30)

Define $\nu(x_*) = \min\{g(x) : x \in [x(t_*) - L_0, x(t_*)]\}$. By condition (H_1) we know that $\nu(x_*) \to +\infty$, as $x(t_*) \to +\infty$. From (2.30) we derive that, for $t \in [\bar{t}_*, t_*]$,

$$y(t) \ge \nu(x_*)(t_* - t) - \varepsilon L_0 - 2\pi (c_\varepsilon + M_p).$$

$$(2.31)$$

Integrating both sides of (2.31) over $[\bar{t}_*, t_*]$ yields

$$\int_{\bar{t}_*}^{t_*} y(t)dt \ge \frac{1}{2}\nu(x_*)(t_* - \bar{t}_*)^2 - [\varepsilon L_0 + 2\pi(c_\varepsilon + M_p)](t_* - \bar{t}_*).$$

Hence, we get that

$$L_0 \ge \frac{1}{2}\nu(x_*)(t_* - \bar{t}_*)^2 - [\varepsilon L_0 + 2\pi(c_\varepsilon + M_p)](t_* - \bar{t}_*),$$

s that

which implies that

$$t_* - \bar{t}_* = o(1), \tag{2.32}$$

for $x_* \to \infty$. Combining (2.27) and (2.32) we obtain

$$\frac{1}{\sqrt{1-\eta}} \int_0^{x_*} \frac{dx}{\sqrt{2(G(x_*) - G(x))}} + o(1) \ge t_* - \alpha.$$

Furthermore,

$$\frac{\tau^+}{2\sqrt{1-\eta}} + o(1) \ge t_* - \alpha.$$
(2.33)

Similarly, we have

$$\frac{\tau^+}{2\sqrt{1-\eta}} + o(1) \ge \beta - t_*.$$
(2.34)

It can be inferred from (2.33) and (2.34) that, for $x_* \to \infty$,

$$\frac{\tau^+}{\sqrt{1-\eta}} + o(1) \ge \beta - \alpha.$$

Since (α, β) denotes any one of (t_1^i, t_2^i) , we reach that

$$\frac{k_0 \tau^+}{\sqrt{1-\eta}} + o(1) \ge \sum_{i=1}^{i=k_0} (t_2^i - t_1^i)$$

Similarly, we have that

$$\frac{k_0 \tau^-}{\sqrt{1-\eta}} + o(1) \ge \sum_{i=1}^{i=k_0} (t_3^i - t_2^i).$$

Therefore, we obtain that

$$\frac{k_0(\tau^+ + \tau^-)}{\sqrt{1 - \eta}} + o(1) \ge \sum_{i=1}^{i=k_0} (t_2^i - t_1^i) + \sum_{i=1}^{i=k_0} (t_3^i - t_2^i) = 2\pi.$$

Recalling that $\eta \to 0$, as $\varepsilon \to 0$, we know that

$$k_0(\tau^+ + \tau^-) \ge 2\pi$$

Lemma 7. Assume that conditions (H_1) , (H_2) and (H_3) hold. Let k > 0 be an arbitrary integer. Then there exists a constant $R_k > 0$ such that for any

604

 2π -periodic solution (x(t), y(t)) of (2.2), with rotation number n(x, y) = k, the following conclusion holds,

$$x(t)^2 + y(t)^2 \le R_k^2, \quad t \in \mathbf{R}.$$

Proof. Assume by contradiction that there exist an integer $k_0 > 0$ and a sequence of 2π -periodic solutions $(x_j(t), y_j(t))$ of (2.2), with the rotation number $n(x_j, y_j) = k_0$ $(j = 1, 2, \dots)$, such that

$$\lim_{j \to +\infty} (x_j^2(t) + y_j^2(t)) = +\infty$$

uniformly for $t \in \mathbf{R}$. From Lemma 5 and Lemma 6 we know that

$$\tau_+ + \tau_- \le \frac{2\pi}{k_0} \le \tau^+ + \tau^-.$$

This contradicts with condition (H_2) .

Proof of Theorem 1. In order to use the Continuation Theorem [6] to prove the existence of 2π -periodic solution of (2.1), we shall check that all conditions of the Continuation Theorem are satisfied. From [8] we know that

(i) There exists B > 0 such that every 2π -periodic solution (x(t), y(t)) of system x' = y, y' = -g(x) satisfies $|x(t)| + |y(t)| \le B, t \in [0, 2\pi]$.

(ii) Define h(x, y) = (y, -g(x)). Then the Brouwer degree d(h, B(0, r), 0) = 1, with r large enough, $B(0, r) = \{(x, y) : x^2 + y^2 \le r^2\}$.

From Lemma 2 we have that

(iii) For any $r_1 > 0$, there exists $r_2 > 0$ such that, for each 2π -periodic solution of (2.2), we have

$$\min_{[0,2\pi]} (x^2(t) + y^2(t)) \le r_1^2 \Longrightarrow \max_{[0,2\pi]} (x^2(t) + y^2(t)) \le r_2^2.$$

From Lemma 7 we know that

(iv) For any integer k > 0, there exists $R_k > 0$ such that, for each 2π -periodic solution of (2.2), we have

$$n(x,y) = k \Longrightarrow \min_{[0,2\pi]} (x^2(t) + y^2(t)) \le R_k^2.$$

Thus, all conditions of the Continuation Theorem are satisfied. Therefore, (2.1) has at least one 2π -periodic solution.

The proof of Theorem 2 can be handled similarly. Indeed, under conditions of Theorem 2, the conditions (i), (ii) and (iii) in proof of Theorem 1 are still satisfied. From Lemma 5 and condition (H_4) we know that (iv) still holds. Therefore, all conditions of the Continuation Theorem are satisfied. Hence, (2.1) has at least one 2π -periodic solution.

Z. Wang

3. Periodic solutions via Lyapunov function

From Massera's theorem [10] we know that if every Cauchy problem of the system

$$\begin{aligned} x' &= h_1(x, y, t), \qquad y' &= h_2(x, y, t) \\ x(0) &= x_0, \qquad y(0) &= y_0 \end{aligned}$$

exists uniquely and is positively bounded, then this system has at least one 2π -periodic solution, where $h_i \in C(\mathbf{R}, \mathbf{R}, \mathbf{R})$ and $h_i(x, y, t + 2\pi) = h_i(x, y, t)$, for $x, y, t \in \mathbf{R}, i = 1, 2$. In this section, by means of Lyapunov function, we will show that all solutions of (2.1) are positively bounded under conditions (H_1) and (H_5) and hence (2.1) possesses at least one 2π -periodic solution. Let us recall that condition (H_5) refers to

$$sgn(x)(f(x) - p(t)) \ge c, \quad \forall t \in \mathbf{R}, \quad |x| \ge d$$

with c, d being positive constants.

Proof of Theorem 3. We follow an argument in [11]. Since f, g are locally Lipschitz continuous, every solution (x(t), y(t)) of (2.1) satisfying the initial value condition $(x(0), y(0)) = (x_0, y_0)$ exists uniquely. Define a potential function V as in Lemma 1,

$$V(x,y) = \frac{1}{2}y^2 + G(x).$$

Set

$$v(t) = V(x(t), y(t)) = \frac{1}{2}y(t)^2 + G(x(t)).$$

Then

$$v'(t) = -y(t)(f(y(t)) - p(t)).$$
(3.1)

Write

$$m_1 = \max\{|f(y)| : -d \le y \le d\}, \quad m_2 = \max\{|p(t)| : t \in \mathbf{R}\}.$$

If $|y| \leq d$, then

$$|(1-y)(f(y) - p(t))| \le m_3, \quad |(1+y)(f(y) - p(t))| \le m_3, \forall t \in \mathbf{R}$$
(3.2)

with $m_3 = (1 + d)(m_1 + m_2)$. Take a constant k > 0 sufficiently large such that

$$|g(x)| \ge m_3, \quad |x| \ge k \quad \text{and} \quad 2d/k \le c.$$
(3.3)

Define a Lyapunov function W(x, y) as follows,

$$W(x,y) = \begin{cases} V(x,y), & |x| < +\infty, \quad y \ge d, \\ V(x,y) - y + d, & x \le -k, \quad |y| \le d, \\ V(x,y) + 2d, & x \le -k, \quad y \le -d, \\ V(x,y) + y - d, & x \ge k, \quad |y| \le d, \\ V(x,y) - 2d, & x \ge k, & y \le -d, \\ V(x,y) - \frac{2d}{k}x, & |x| \le k, & y \le -d. \end{cases}$$

606

It is easy to check that W(x, y) is continuous and locally Lipschitz with respect to $(x, y) \in \{(x, y) : |x| \ge k, |y| \ge d\}$. Moreover, W(x, y) tends to infinity uniformly for $x \in \mathbf{R}$ as $|y| \to +\infty$. Set $\Gamma(x, y) = V(x, y) + 2dx/k + y + 2d$. Then $\Gamma(x, y)$ is continuous and $W(x, y) \le \Gamma(|x|, |y|)$. By using (3.1)-(3.3) and the expression of W(x, y), we have that the derivative W'(x(t), y(t)) of W(x(t), y(t)) with respect to t satisfies

$$W'(x(t), y(t)) \le 0.$$

On the other hand, let l > 0 be a constant. Then there exists a constant $\mathcal{L} > 0$ such that, for $|y| \leq l$,

$$|y - f(y) + p(t)| \le \mathcal{L}, \quad \forall t \in \mathbf{R}.$$

Take a constant r > 0 such that

 $|g(x)| \ge \mathcal{L}, \quad |x| \ge r.$

Define another Lyapunov function

$$U(x,y) = \begin{cases} x+y, & x \ge r, \quad |y| \le l, \\ -x-y, & x \le -r, \quad |y| \le l. \end{cases}$$

Obviously, U(x, y) satisfies the following conclusions.

(1)U(x,y) tends to infinity uniformly for $|y| \le l$ as |x| tends to infinity.

 $(2)U(x,y) \le |x| + l$, for $|y| \le l$.

(3) if $x \ge r$, $|y| \le l$, then $U'(x(t), y(t)) = y(t) - g(x(t)) - f(y(t)) + p(t) \le 0$ and if $x \le -r$, $|y| \le l$, then $U'(x(t), y(t)) = -y(t) + g(x(t)) + f(y(t)) - p(t) \le 0$.

Therefore, all conditions of Theorem 8.9 in [11] are satisfied. Furthermore, all solutions of (2.1) are positively bounded. It follows from Corollary 15.1 in [11] that Eq.(2.1) has at least one 2π -periodic solution.

Acknowledgement

This work was done during the author's stay at the Department of Mathematics, Torino University. The author thanks Prof. Anna Capietto for many valuable discussions and the Department of Mathematics for it's hospitality. The author also thanks the referee for his (or her) valuable comments and suggestions.

References

- S. Radhakrishnan, Exact solutions of Rayleigh's equation and sufficient conditions for inviscid instability of parallel, bounded shear flows, Z. Angew. Math. Phys. 45 (1994), 615–637.
- [2] S. Hal L, On the small oscillations of the periodic Rayleigh equation, Quart. Appl. Math. 44 (1986), 223–247.
- [3] S. Russell A., Period bounds for generalized Rayleigh equation, Internat. J. Non-Linear Mech. 6 (1977), 271–277.

- [4] P. Omari, G. Villari, On a continuation lemma for the study of a certain planar system with applications to Liénard and Rayleigh equations, *Results Math.* 14 (1998), 156–173.
- [5] T. Ding, F. Zanolin, Time maps for the solvability of periodically perturbed nonlinear Duffing equations, Nonlinear Anal. 17 (1991), 635–653.
- [6] A. Capietto, J. Mawhin, F. Zanolin, A continuation approach to superlinear boundary value problems, J. Differential Equations 88 (1990), 347–395.
- [7] T. Ding, F. Zanolin, Periodic solutions of Duffing's equations with superquadratic potential, J. Differential Equations 97 (1992), 328–378.
- [8] Z. Wang, Periodic solutions of the second order forced Liénard equation via time maps, Nonlinear Anal. 48 (2002), 445–460.
- [9] Z. Wang, Periodic solutions of the second order differential equations with asymmetric nonlinearities depending on the derivatives, *Discrete Contin. Dynam. Systems*, 9 (2003), 751– 770.
- [10] J. L. Massera, The existence of periodic solutions of systems of differential equations, Duke Math. J. 17 (1950), 457–475.
- [11] T. Yoshizawa, Stability Theory and the Existence of Periodic Solutions and Almost Periodic Solutions. Springer-Verlag New York, 1975.
- [12] M. A. Krasnosel'skii, The Operator of Translation along the Trajectories of Differential Equations. American Mathematical Society, Providence, RI, 1968.
- [13] Z. Opial, Sur les périodes des solutions de l'équation différentiale x'' + g(x) = 0, Ann. Pol. Math. 10 (1961), 49–72.
- [14] P. Habets and P. J. Torres, Some multiplicity results for periodic solutions of a Rayleigh differential equation, Dyn. Contin. Discrete Impuls. Syst. Ser. A Math. Anal. 8 (2001), 335–351.

Zaihong Wang Department of Mathematics Capital Normal University Beijing 100037 People's Republic of China e-mail: zhwang@mail.cnu.edu.cn and Dipartimento di Matematica Università di Torino Via Carlo Alberto 10 10123 Torino Italia

(Received: July 1, 2002; revised: February 19, 2003)



To access this journal online: http://www.birkhauser.ch