Published for SISSA by 2 Springer

RECEIVED: April 5, 2023 ACCEPTED: September 11, 2023 PUBLISHED: November 2, 2023

# Adjacency for scattering amplitudes from the Gröbner fan

# L. Bossinger,<sup>a</sup> J.M. Drummond<sup>b</sup> and R. Glew<sup>b</sup>

<sup>a</sup> Instituto de Matemáticas, Unidad Oaxaca, Universidad Nacional Autónoma de México, León 2, 68000 Oaxaca, Mexico

<sup>b</sup>School of Physics & Astronomy, University of Southampton, Highfield, Southampton, SO17 1BJ, U.K.

*E-mail:* lara@im.unam.mx, j.m.drummond@soton.ac.uk, r.glew@herts.ac.uk

ABSTRACT: Scattering amplitudes in planar  $\mathcal{N} = 4$  super Yang-Mills theory exhibit singularities which reflect various aspects of the cluster algebras associated to the Grassmannians  $\operatorname{Gr}(4, n)$  and their tropical counterparts. Here we investigate the potential origins of such structures and examine the extent to which they can be recovered from the Gröbner structure of the underlying Plücker ideals, focussing on the Grassmannians corresponding to finite cluster algebras.

Starting from the Plücker ideal, we describe how the polynomial cluster variables are encoded in non-prime initial ideals associated to certain maximal cones of the positive tropical fan. Following [1] we show that extending the Plücker ideal by such variables leads to a Gröbner fan with a single maximal Gröbner cone spanned by the positive tropical rays. The associated initial ideal encodes the compatibility relations among the full set of cluster variables. Thus we find that the Gröbner structure naturally encodes both the symbol alphabet and the cluster adjacency relations exhibited by scattering amplitudes without invoking the cluster algebra at all.

As a potential application of these ideas we then examine the kinematic ideal associated to non-dual conformal massless scattering written in terms of spinor helicity variables. For five-particle scattering we find that the ideal can be identified with the Plücker ideal for Gr(3, 6) and the corresponding tropical fan contains a number of non-prime ideals which encode all additional letters of the two-loop pentagon function alphabet present in various calculations of massless five-point finite remainders.

**KEYWORDS:** Differential and Algebraic Geometry, Scattering Amplitudes

ARXIV EPRINT: 2212.08931

OPEN ACCESS,  $\bigcirc$  The Authors. Article funded by SCOAP<sup>3</sup>.



# Contents

1	Introduction	1	
2	The Gröbner fan         2.1       Polynomial cluster variables and forbidden pairs         2.2       GP(L_1)	<b>3</b> 6	
	2.2 $GF(I_{2,5})$	7	
3	$\operatorname{Trop}^+(I_{k,n})$ and the Speyer-Williams fan	8	
	3.1  Gr(2,5)	10	
4	Alphabets and adjacency from the Gröbner fan	11	
	4.1 $Gr(3,6)$	11	
	4.2 $Gr(3,7) \cong Gr(4,7)$	17	
	4.3 $Gr(3,8)$	19	
5	Missing binomials and cluster variables	21	
	5.1 The g-vectors as multiweights	21	
	5.2 Cluster variables as missing binomials	22	
6	Massless scattering ideals	24	
	6.1 The five-point two-loop symbol alphabet	25	
	6.2 Non-planar alphabet from a Gröbner fan	26	
7	Conclusion	29	
$\mathbf{A}$	Details of $\operatorname{Trop}^+(I_{3,7})$	30	

### 1 Introduction

The connection between the branch cut structure of perturbative *n*-point scattering amplitudes in planar  $\mathcal{N} = 4$  super Yang-Mills theory and the cluster algebras associated to the Grassmannians  $\operatorname{Gr}(4, n)$  is by now well established [2-4].<sup>1</sup> Such amplitudes can be expressed (at least for small enough *n* and loop order) as generalised polylogarithmic functions, a class of iterated integrals built upon a specified symbol alphabet A. The set of cluster variables of the Grassmannian cluster algebras  $\operatorname{Gr}(4, 6)$  and  $\operatorname{Gr}(4, 7)$  explain the 9 and 42 letter alphabets for the hexagon and heptagon amplitudes respectively, the knowledge of which has since facilitated calculations in the context of the analytic bootstrap up to high loop orders [8–17]. Furthermore, the link between the cluster algebras and the singularities of the amplitudes was deepened with the discovery of *cluster adjacency* [18, 19]

<sup>&</sup>lt;sup>1</sup>The notion of cluster algebras was originally developed in [5-7].

which states that consecutive singularities can only appear in the symbol if there exists a cluster in which both letters are contained. This notion of cluster adjacency is closely related to the Steinmann relations, utilised by the bootstrap programme, and their extended counterparts [13, 20].

Recently, efforts have been made to explore the connection between amplitudes and cluster algebras at eight points and beyond where the respective Grassmannian cluster algebra Gr(4, 8) is no longer of finite type. Two complications arise in this instance. Firstly, some truncation of the infinite set of cluster coordinates must be introduced in order to obtain a finite symbol alphabet and secondly, we must find a way of extracting algebraic square root letters known to appear in the symbol alphabet at eight points, for example from the four-mass one-loop box integral present in N<sup>2</sup>MHV amplitudes. This has led to the study of the eight-point case and related problems from many closely related approaches including tropical geometry, plabic graphs, Schubert problems and more [21–29]. Notably [22], by assigning expressions to certain rays appearing in the tropicalisation of Gr(4, 8), which can also be obtained via infinite sequences of mutations in the cluster algebra, one obtains the set of 18 multiplicatively independent square root letters appearing in the two-loop NMHV amplitude found in [30].

Since the tropicalisation of a polynomial ideal is a subset of the Gröbner fan of the ideal by the fundamental theorem of tropical geometry ([31], section 3.2), the connection between the cluster algebra and the tropical Grassmannian leads us to ask whether the salient features of the cluster algebra can be obtained by considering the Gröbner structure of the underlying ideal generated by the Plücker relations. While connections between tropical Grassmannians [32] and cluster algebras had been discussed already in e.g. [33, 34], it is only recently that the relationship between cluster algebras and Gröbner theory has been investigated in detail [1, 35]. We shall see, using methods developed in these references, that both the alphabet and the adjacency relations can indeed be obtained from the Gröbner structure of the Plücker ideal or an appropriate extension thereof.

In the case of  $\operatorname{Gr}(2, n)$ , all cluster variables are Plücker variables. For the finite cases  $\operatorname{Gr}(3, n)$  for n = 6, 7, 8, we will see that cluster variables which are polynomial in Plücker coordinates are associated to certain non-prime initial ideals associated to maximal cones of the positive tropical fan  $\operatorname{Trop}^+(I_{k,n})$  (a subfan of the Gröbner fan). Furthermore, by extending the Plücker ideal by this additional set of cluster variables, a particular maximal cone of the Gröbner fan will be singled out by the rays of  $\operatorname{Trop}^+(I_{k,n}^{ext})$ . The associated monomial initial ideal provides the set of forbidden pairs of cluster variables i.e. variables which do not appear together in a cluster. For the cases of  $\operatorname{Gr}(2,6) \cong \operatorname{Gr}(4,6)$  and  $\operatorname{Gr}(3,7) \cong \operatorname{Gr}(4,7)$  this provides the information required for the amplitude bootstrap programme in the form of the symbol alphabet  $\mathbb{A}$  and adjacency rules for the symbol of the amplitude.

Motivated by the fact that interesting physical information can be obtained from the Gröbner structure of the kinematic ideal for amplitudes in planar  $\mathcal{N} = 4$  super Yang-Mills theory, we then consider the kinematics of non-planar and non-dual conformal massless scattering described in spinor helicity variables. We consider an ideal generated by the polynomial relations satisfied by the spinor brackets  $\langle ij \rangle$  and [ij] and investigate its Gröbner

structure. In particular we consider the ideal  $I_{5pt}$  corresponding to the case of scattering five massless particles. In fact we observe that, upon an appropriate identification of variables, this ideal can be identified with the Plücker ideal for Gr(3,6). By searching for non-prime initial ideals, this time in the full tropical space  $Trop(I_{5pt})$ , not just the positive part, we *almost* recover the entire non-planar alphabet found in recent two-loop massless five-point amplitudes calculations in gauge theory [36–41] and gravity [42, 43]. Interestingly, the one letter we do not obtain also seems to drop out of appropriately defined finite remainders [39, 40, 42, 43].

#### 2 The Gröbner fan

The Grassmannian  $\operatorname{Gr}(k, n)$  is the space of k-planes in n dimensions. A point in the Grassmannian can be specified by k n-component vectors which can be organised into a  $k \times n$  matrix. These matrices are defined up to row operations which leave the plane invariant, the resulting space of matrices modulo  $\operatorname{GL}(k)$  transformations is k(n-k) dimensional.

Alternatively, the Grassmannian can be described through the set of  $\binom{n}{k}$  maximal minors  $p_{i_1...i_k}$  or Plücker coordinates. On the set of Plücker coordinates row operations act as an overall scaling and, modulo the overall scaling, the vector of Plücker coordinates may be thought of as a point in the projective space  $\mathbb{P}^{\binom{n}{k}-1}$ . An arbitrary point in  $\mathbb{P}^{\binom{n}{k}-1}$  is not necessarily realisable as a matrix, since the set of  $k \times k$  minors of any  $k \times n$  matrix are not independent, and instead obey homogenous quadratic relations known as the Plücker relations which take the form

$$p_{i_1\dots i_r[i_{r+1}\dots i_k}p_{j_1\dots j_{r+1}]j_{r+2}\dots j_k} = 0.$$
(2.1)

We call the ideal generated by the Plücker relations inside the ring of polynomials in the Plücker coordinates the Plücker ideal  $I_{k,n}$ . The Grassmannian can then be thought of as the projective variety inside  $\mathbb{P}^{\binom{n}{k}-1}$  whose points vanish on the Plücker ideal, i.e. the vanishing set  $V(I_{k,n})$ . As an example consider the case of  $\operatorname{Gr}(2,n)$ , whose Plücker ideal is generated by the relations

$$I_{2,n} = \langle p_{ij} p_{kl} - p_{ik} p_{jl} + p_{il} p_{jk} : 1 \le i < j < k < l \le n \rangle.$$
(2.2)

Here we will actually consider the Grassmannian modulo the action of the n rescalings which can be applied to the n columns of the  $k \times n$  matrix, so that each column actually represents an element of  $\mathbb{P}^{k-1}$ . Since the overall scaling is taken into account already in the action of  $\operatorname{GL}(k)$ , the dimension of the space is reduced to k(n-k)-(n-1) = (k-1)(n-k-1). As is common in the literature we will also refer to this space simply as the Grassmannian and henceforth when we discuss the Grassmannian we will always mean it in this reduced sense. The same scalings are present as an invariance of the Plücker relations (2.1) which are not only homogeneous in the total degree of all variables but homogeneous in the presence of each label  $i_k$  on the Plücker coordinates.

Using the general viewpoint of the Grassmannian as an ideal generated by polynomial relations we can introduce a fan structure on  $\mathbb{R}^{\binom{n}{k}}$  known as the Gröbner fan. To begin to

understand the structure of the Gröbner fan we must first introduce the notion of monomial orderings, initial ideals, and Gröbner bases. Our presentation follows that of [1].

Let f be a polynomial in n variables  $(x_1, \ldots, x_n)$  with coefficients in an algebraically closed field K for which we use the notation

$$f = \sum_{\vec{\alpha}} c_{\vec{\alpha}} \mathbf{x}^{\vec{\alpha}},\tag{2.3}$$

where we have introduced  $\vec{\alpha} = (\alpha_1, \ldots, \alpha_n) \in \mathbb{Z}_{\geq 0}^n$ , and  $\mathbf{x}^{\vec{\alpha}}$  is understood as the monomial  $x_1^{\alpha_1} \ldots x_n^{\alpha_n}$ . Given some weight vector  $\vec{w} \in \mathbb{R}^n$  we can define the *initial form* of f with respect to  $\vec{w}$  as

$$\operatorname{in}_{\vec{w}}(f) = \sum_{\vec{\alpha}: \ \vec{\alpha} \cdot \vec{w} = m} c_{\vec{\alpha}} \mathbf{x}^{\vec{\alpha}}, \tag{2.4}$$

where  $m = \min\{\vec{\alpha} \cdot \vec{w} : c_{\vec{\alpha}} \neq 0\}$ . Furthermore, given an ideal  $I \subset \mathbb{K}[x_1, \ldots, x_n]$ , we can define its *initial ideal* with respect to  $\vec{w}$  as the ideal generated by the initial forms of all functions  $f \in I$  written as

$$\operatorname{in}_{\vec{w}}(I) = \langle \operatorname{in}_{\vec{w}}(f) : f \in I \rangle.$$
(2.5)

If, for some finite set of generators  $\mathcal{G} = \{g_1, \ldots, g_r\} \in I$ , we have  $\operatorname{in}_{\vec{w}}(I) = \langle \operatorname{in}_{\vec{w}}(g) : g \in \mathcal{G} \rangle$ we call  $\mathcal{G}$  a *Gröbner basis* for I with respect to  $\vec{w}$ .

The next definition we need is that of a *monomial order*. A monomial order < on the set of monomials  $\mathbf{x}^{\vec{\alpha}} \in \mathbb{K}[x_1, \dots, x_n]$  is a total order which satisfies

$$\begin{array}{ll} i) \ 1 \leq \mathbf{x}^{\vec{\alpha}}, \\ ii) \ \mathrm{if} \ \mathbf{x}^{\vec{\alpha}} < \mathbf{x}^{\vec{\beta}} \implies \ \mathbf{x}^{\vec{\alpha} + \vec{\gamma}} < \mathbf{x}^{\vec{\beta} + \vec{\gamma}}. \end{array}$$

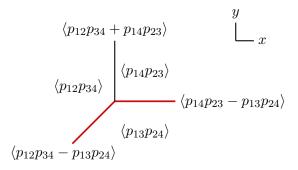
This allows us to define the leading monomial of the polynomial f as  $\operatorname{in}_{<}(f) = c_{\vec{\beta}} \mathbf{x}^{\vec{\beta}}$ , where  $\mathbf{x}^{\vec{\beta}}$  is the leading monomial with respect to < appearing in f with non-zero coefficient i.e.  $\mathbf{x}^{\vec{\beta}} = \max_{<} \{\mathbf{x}^{\vec{\alpha}} : c_{\vec{\alpha}} \neq 0\}$ . Similarly, we can define the initial ideal of I with respect to < as

$$\operatorname{in}_{<}(I) = \langle \operatorname{in}_{<}(f) : f \in I \rangle.$$
(2.6)

Note, we may always choose some weight vector  $\vec{w} \in \mathbb{N}^n$  such that  $\operatorname{in}_{\vec{w}}(I) = \operatorname{in}_{<}(I)$  ([44], Theorem 3.2.1). The converse is not generally true however.

By varying the weight vector  $\vec{w}$  we may study all possible initial ideals of I. This leads us to the notion of the Gröbner fan GF(I) on  $\mathbb{R}^n$  as follows: two weight vectors  $\vec{w}_1$  and  $\vec{w}_2$  lie in the relative interior of the same cone C if and only if  $\operatorname{in}_{\vec{w}_1}(I) = \operatorname{in}_{\vec{w}_2}(I)$  i.e. they generate the same initial ideal [45]. Note, each full-dimensional (maximal) Gröbner cone is associated to a monomial initial ideal specified by some monomial order <, consisting of all weight vectors  $\vec{w} \in \mathbb{R}^n$  such that  $\operatorname{in}_{\vec{w}}(I) = \operatorname{in}_{<}(I)$ . A weight vector will lie on the boundary of a maximal cone when the associated initial ideal is no longer monomial. This collection of maximal Gröbner cones and their intersections is the Gröbner fan GF(I).

Note that the Gröbner fan may have a linear subspace (or lineality space) consisting of elements  $\vec{l}$  such that  $\ln_{\vec{l}}(I) = I$ . This is the case, for example, if the ideal is homogeneous. We are always free to shift a weight vector by any element of the lineality space of the



**Figure 1.** The Gröbner fan structure of  $GF(I_{2,4})$  with each region labelled by its initial ideal. Each point in the (x, y) plane corresponds to a 4-dimensional linear subspace of  $\mathbb{R}^6$  consisting of all weight vectors lineality equivelent to (x, y, 0, 0, 0, 0). The tropical fan corresponds to the three rays, whilst the positive tropical fan corresponds to the two red rays.

Gröbner fan GF(I) without altering the initial ideal. Therefore it makes sense to consider the Gröbner fan modulo the action of this linear subspace. In the case of the Grassmannians Gr(k, n) we have an *n*-dimensional lineality space corresponding to the *n* column rescalings mentioned previously. This is precisely because the Plücker relations are homogeneous with respect to the appearance of all labels  $i_r$  on the Plücker variables  $p_{i_1...p_k}$ .

We will also be interested in certain subfans of the Gröbner fan, the first being the tropical fan Trop(I) defined as the subfan

$$\operatorname{Trop}(I) = \{ \vec{w} \in \mathbb{R}^n : \operatorname{in}_{\vec{w}}(I) \text{ contains no monomial} \}.$$

$$(2.7)$$

We may restrict further and define the totally positive tropical fan  $\operatorname{Trop}^+(I)$  given by

$$\operatorname{Trop}^{+}(I) = \{ \vec{w} \in \operatorname{Trop}(I) : \operatorname{in}_{\vec{w}}(I) \text{ is totally positive} \},$$
(2.8)

where we an ideal  $I \subset \mathbb{R}[x_1, \ldots, x_n]$  is called *totally positive* if it does not contain any non-zero polynomial with all positive coefficients (or equivalently, it has an initial ideal whose vanishing set intersects the positive orthant [46]).

The above discussion is most easily demonstrated with an example, the simplest case being Gr(2, 4), whose Plücker ideal is generated by a single polynomial

$$I_{2,4} = \langle p_{12}p_{34} - p_{13}p_{24} + p_{14}p_{23} \rangle \subset \mathbb{R}[p_{12}, p_{13}, p_{14}, p_{23}, p_{24}, p_{34}].$$
(2.9)

Let  $\vec{w} = (w_{12}, w_{13}, w_{14}, w_{23}, w_{24}, w_{34}) \in \mathbb{R}^6$  and  $f = p_{12}p_{34} - p_{13}p_{24} + p_{14}p_{23}$  which, being the sole generator of  $I_{2,4}$ , constitutes a Gröbner basis for any choice of weight vector  $\vec{w}$ . A generic weight vector  $\vec{w} \in \mathbb{R}^6$  can always be bought to the form  $(x, y, 0, \dots, 0)$  with some suitable choice of lineality shift in the four-dimensional linear subspace. The resulting Gröbner fan  $GF(I_{2,4})$  is depicted in the (x, y) plane in figure 1. The Gröbner fan  $GF(I_{2,4})$ consists of three maximal cones labelled by the monomial initial ideals  $\langle p_{12}p_{34}\rangle$ ,  $\langle p_{13}p_{24}\rangle$  and  $\langle p_{14}p_{23}\rangle$ . The maximal cones intersect to give the three rays of the tropical fan Trop $(I_{2,4})$ given by  $e_{12} = (1,0), e_{13} = (0,1)$  and  $e_{14} = (-1,-1)$ , whose corresponding binomial initial ideals are  $\langle p_{14}p_{23} - p_{13}p_{24}\rangle$ ,  $\langle p_{12}p_{34} + p_{14}p_{23}\rangle$  and  $\langle p_{12}p_{34} - p_{13}p_{24}\rangle$ . The positive part of the tropical fan Trop<sup>+</sup>( $I_{2,4}$ ) consists of the rays  $e_{12} = (1,0)$  and  $e_{14} = (-1,-1)$  highlighted in red whose generators contain terms of opposite sign, the ray  $e_{13} = (0,1)$  is not contained in Trop<sup>+</sup>( $I_{2,4}$ ) since the corresponding initial ideal is generated by a polynomial with all positive coefficients.

#### 2.1 Polynomial cluster variables and forbidden pairs

Moving forward, the final definition needed is that of a prime (or alternatively non-prime) ideal. An ideal I is non-prime if there exists two polynomials  $f \notin I$  and  $g \notin I$  such that their product  $f \cdot g \in I$ . In this case we call f and g non-prime factors of I. Note, a non-prime ideal can always be decomposed into the intersection of finitely many prime components.

With all the necessary material reviewed let us remind ourselves of our goal: to extract (at least in the finite cases) the cluster variables and adjacency relations of the Grassmannian cluster algebras  $\operatorname{Gr}(k, n)$  from the Gröbner fan  $GF(I_{k,n})$ , which in the cases of  $\operatorname{Gr}(4, 6)$ and  $\operatorname{Gr}(4, 7)$  provide vital information for the amplitude bootstrap in the form of the symbol alphabet and adjacency rules. The ideas we will make use of have been presented in [1] for the case  $\operatorname{Gr}(2, n)$  and  $\operatorname{Gr}(3, 6)$  and more generally for any cluster algebra of geometric finite type in [35].

**Cluster variables.** The Plücker ideal is defined on the  $\binom{n}{k}$  Plücker coordinates  $p_{i_1,\ldots,i_k}$ . In the case of the Grassmannians  $\operatorname{Gr}(2,n)$  these make up the full set of cluster variables. However, for  $\operatorname{Gr}(3,6)$ , and more importantly  $\operatorname{Gr}(3,7)$  relevant for heptagon amplitudes, cluster variables quadratic in the Plücker coordinates start to appear. As we shall explain these missing cluster variables appear as non-prime factors of initial ideals inside the maximal cones of  $\operatorname{Trop}^+(I_{k,n})$ .

Forbidden pairs. Generally, the rays of the positive tropical fan  $\operatorname{Trop}^+(I_{k,n})$  will span multiple maximal Gröbner cones. That is to say taking a suitably generic<sup>2</sup> set of weight vectors  $\vec{w}$  lying in the span of the rays of  $\operatorname{Trop}^+(I_{k,n})$  we will generate multiple monomial initial ideals  $\operatorname{in}_{\vec{w}}(I_{k,n})$ . However, upon extending the ideal  $I_{k,n}$  by the missing cluster variables to the ideal  $I_{k,n}^{\text{ext}}$ , the rays of  $\operatorname{Trop}^+(I_{k,n}^{\text{ext}})$  span a single maximal Gröbner cone in  $GF(I_{k,n}^{\text{ext}})$ . The minimal generating set of the initial ideal of this maximal Gröbner cone provides us with a list of monomials which are exactly the forbidden pairs of cluster variables. Note, in the case of  $\operatorname{Gr}(2, n)$  all cluster variables are present already for  $I_{2,n}$  and no extension procedure is needed. The rays of  $\operatorname{Trop}^+(I_{2,n})$  already span a unique maximal Gröbner cone.

In the remainder of this section we review these ideas in more detail for the case of Gr(2,5). Later in section 4 we return to the case of Gr(3,6), discussed in [1], and further apply this discussion to  $Gr(3,7) \cong Gr(4,7)$  relevant for heptagon amplitudes. We also discuss the remaining finite case Gr(3,8) and the outlook beyond the finite type Grassmannians. Finally, in section 6 we will apply some of the ideas to a non-Grassmannian case relevant for scattering in more general massless theories.

 $<sup>^{2}</sup>$ Weight vectors not lying in the intersection of maximal cones of the Gröbner fan.

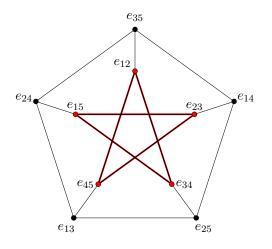


Figure 2. The 10 vertices and 15 edges of  $\text{Trop}(I_{2,5})$  space. The positive part is highlighted red.

# **2.2** $GF(I_{2,5})$

We conclude this section with the example of Gr(2, 5). In the space of Plücker coordinates  $(p_{12}, \ldots, p_{45})$  the five Plücker relations are given by

$$p_{ij}p_{kl} - p_{ik}p_{jl} + p_{il}p_{jk} = 0, \quad 1 \le i < j < k < l \le 5.$$

$$(2.10)$$

The Gröbner fan  $GF(I_{2,5})$  is simplicial, containing 132 maximal cones and twenty rays. Arranging the coordinates in lexicographic order

$$\{w_{12}, w_{13}, w_{14}, w_{15}, w_{23}, w_{24}, w_{25}, w_{34}, w_{35}, w_{45}\},$$
(2.11)

the rays are defined as

$$e_{12} = (1, 0, 0, 0, 0, 0, 0, 0, 0),$$
  

$$\vdots$$
  

$$e_{45} = (0, 0, 0, 0, 0, 0, 0, 0, 0, 1),$$
(2.12)

along with ten more given by  $-e_{ij}$ .

The Gröbner fan contains twelve maximal cones with five rays given by  $e_{ij}$  from the list (2.12). A further 60 maximal cones have four of the  $e_{ij}$  and one of the  $-e_{ij}$  rays and 60 more maximal cones have three  $e_{ij}$  and two  $-e_{ij}$  rays. The ten  $e_{ij}$  vectors defined in (2.12) make up the rays of the tropical fan Trop $(I_{2,5})$ . They are connected in a Petersen graph topology shown in figure 2. The positive tropical fan Trop<sup>+</sup> $(I_{2,5})$  has five rays given by

$$\{e_{12}, e_{23}, e_{34}, e_{45}, e_{15}\}$$

$$(2.13)$$

which are highlighted in red.

There is exactly one maximal Gröbner cone spanned by the five rays (2.13) of  $\operatorname{Trop}^+(I_{k,n})$ . This cone has a Gröbner basis whose initial monomials are all the crossing chords of the pentagon, i.e.

$$\operatorname{in}_{\vec{w}}(I_{2,5}) = \langle p_{13}p_{24}, p_{13}p_{25}, p_{14}p_{25}, p_{14}p_{35}, p_{24}p_{35} \rangle$$
(2.14)

for any  $\vec{w}$  given by a strictly positive linear combination of the five rays of  $\operatorname{Trop}^+(I_{2,5})$ .

The case of  $\operatorname{Gr}(2, n)$  was studied in detail in [1], where a similar construction was shown to hold for all n. The rays of the positive tropical fan  $\operatorname{Trop}^+(I_{2,n})$  span a single maximal Gröbner cone inside  $GF(I_{2,n})$ , whose initial ideal is generated products of Plücker variables labelled by crossing chords of the n-gon. Note that this includes the case  $\operatorname{Gr}(2, 6) \cong \operatorname{Gr}(4, 6)$ relevant for hexagon amplitudes in planar  $\mathcal{N} = 4$  super Yang-Mills theory. As was also shown in [1] for the case of  $\operatorname{Gr}(3, 6)$ , and as we shall explain in section 4, it is not the case for k > 2 that the rays of  $\operatorname{Trop}^+(I_{k,n})$  span a single maximal Gröbner cone inside  $GF(I_{k,n})$ . Instead, to identify a single maximal Gröbner cone, whose rays are given by the positive tropical part, the ideal must first be extended by the missing cluster variables which are polynomials in Plücker variables.

# 3 Trop<sup>+</sup> $(I_{k,n})$ and the Speyer-Williams fan

In practice, it is only possible to compute the entire Gröbner fan for the most simple of cases. Therefore, it is highly desirable to have an efficient route to calculating the positive tropical part  $\text{Trop}^+(I_{k,n})$  directly without having it embedded in the Gröbner fan, or even in  $\text{Trop}(I_{k,n})$ . In this section we review the methods of [33] on this direct construction. For further details of the structure of the resulting fans see [21, 47]. The cases of interest will be Gr(3,n) for n = 6, 7, 8 which, along with Gr(2,n), make up the finite type Grassmannian cluster algebras.

Generally, the Gr(k, n) initial cluster has the form of a  $(k-1) \times (n-k-1)$  array of active (mutable) nodes, in addition to k frozen nodes, each labelled by an  $\mathcal{A}$ -coordinate. An example of the initial cluster for the case of Gr(2, 5) is given in figure 3, with frozen nodes indicated as the boxed vertices. We may also assign to each active node an  $\mathcal{X}$ -coordinate, given by the product of incoming  $\mathcal{A}$ -coordinates over the product of the outgoing ones. Following the array of active nodes, these also organise themselves into a  $(k-1) \times (n-k-1)$ array X with elements  $x_{rs}$ .

Using the  $x_{rs}$  we can define the  $k \times n$  web matrix  $W^{(k,n)} = (\mathbf{I}_k | M)$ . Where M is given by the  $k \times (n-k)$  matrix elements

$$m_{ij} = (-1)^{i+k} \sum_{\vec{\lambda} \in Y_{ij}} \prod_{r=1}^{k-i} \prod_{s=1}^{\lambda_r} x_{rs}, \qquad (3.1)$$

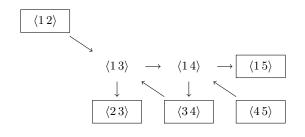
with the summation range  $Y_{ij}$  given by  $0 \le \lambda_{k-i} \le \ldots \le \lambda_1 \le j-1$ . The web matrix thus allows us to evaluate all  $\mathcal{A}$ -coordinates as subtraction free polynomials in the  $\mathcal{X}$ -coordinates by identifying the Plücker coordinates with the maximal minors of the web matrix i.e.

$$p_{i_1\dots i_k} = \det(W_{i_1,\dots,i_k}^{(k,n)})(x_{rs})$$
(3.2)

where  $W_{i_1,\ldots,i_k}^{(k,n)}$  is understood as the matrix formed by taking columns  $i_1 \ldots i_k$  of  $W^{(k,n)}$ .

Let us illustrate this with the example of Gr(2,5). The initial cluster is depicted in figure 3 from which we can read off the  $\mathcal{X}$ -coordinates, they are given by

$$x_{11} = \frac{\langle 12 \rangle \langle 34 \rangle}{\langle 14 \rangle \langle 23 \rangle}, \qquad \qquad x_{12} = \frac{\langle 13 \rangle \langle 45 \rangle}{\langle 34 \rangle \langle 15 \rangle}. \tag{3.3}$$



**Figure 3**. The initial cluster of Gr(2,5).

With the  $\mathcal{X}$ -coordinates to hand we can write down the web matrix,

$$W^{(2,5)} = \begin{bmatrix} 1 & 0 & -1 & -1 & -x_{11} & -1 & -x_{11} & -x_{11} & x_{12} \\ 0 & 1 & 1 & 1 & 1 \end{bmatrix}.$$
 (3.4)

By identifying the Plücker coordinates  $p_{ij}$  with the maximal minor formed by columns iand j of the web matrix as

$$p_{ij} = \det(W_{ij}^{(2,5)})(x_{11}, x_{12}), \tag{3.5}$$

we immediately arrive at an expression for all  $\mathcal{A}$ -coordinates as subtraction free polynomials in the  $\mathcal{X}$ -coordinates. As an example we have

$$p_{25} = x_{11} + x_{11}x_{12},$$

We can now consider tropicalising the expressions for the  $\mathcal{A}$ -coordinates, which amounts to replacing  $(+, \times)$  by their tropical counterparts (min, +). The tropical version of the above expression for  $p_{25}$  is given by

$$\tilde{p}_{25} = \min(\tilde{x}_{11}, \tilde{x}_{11} + \tilde{x}_{12}), \qquad (3.6)$$

where  $\tilde{p}$  and  $\tilde{x}$  are used to emphasise that we are dealing with tropical expressions. This tropical expression defines a piecewise linear map on  $\mathbb{R}^2$  with coordinates  $\{\tilde{x}_{11}, \tilde{x}_{12}\}$ , where regions of linearity are separated by tropical hypersurfaces, and as such provide a fan structure on the space of  $\{\tilde{x}_{11}, \tilde{x}_{12}\}$ . For example the tropical hypersurface of  $\tilde{p}_{25}$  is given by

$$\tilde{x}_{12} = 0.$$
 (3.7)

By tropicalising different subsets of the  $\mathcal{A}$ -coordinates we can define different tropical fans given by the common refinement of all fans in the subset of tropical expressions. In practice, we calculate the refinement of the tropical fan for a subset,  $\mathcal{S}$ , of  $\mathcal{A}$ -coordinates via the Minkowski sum of their Newton polytopes. The tropical expressions for the frozen coordinates do not contain any tropical hypersurface and hence do not contribute to the structure of the fan.

Our focus will be on two fans in particular: the Speyer-Williams fan [33], which is obtained by tropicalising the set of all Plücker coordinates; and the cluster fan, where we choose to tropicalise the entire set of  $\mathcal{A}$ -coordinates. Note, for the case of  $\operatorname{Gr}(2, n)$  the Plücker coordinates do make up the entire set of  $\mathcal{A}$ -coordinates and hence the Speyer-Williams and cluster fans coincide. However, when considering Gr(3,6)  $\mathcal{A}$ -coordinates quadratic in the Plücker variables begin to appear and hence the structure of the two fans begins to differ.

While we have used the initial cluster to explain the definition of the  $x_{rs}$  we would like to emphasise that we are merely using them as a convenient choice of coordinates. We will not be using any of the mutation structure of the cluster algebra and indeed one of the main points of this investigation is to recover information about the Grassmannian cluster algebras (including the set of  $\mathcal{A}$ -coordinates) just by carefully studying the Gröbner structure of the Plücker ideals.

#### **3.1** Gr(2, 5)

The Gr(2,5) web matrix, written in (3.4), allows us to write the 10 Plücker variables  $p_{ij}$  in terms of the two  $\mathcal{X}$ -coordinates  $(x_{11}, x_{12})$  as

$$p_{1i} = p_{23} = 1, \qquad p_{24} = 1 + x_{11}, \qquad p_{25} = 1 + x_{11} + x_{11}x_{12}, p_{34} = x_{11}, \qquad p_{35} = x_{11} + x_{11}x_{12}, \qquad p_{45} = x_{11}x_{12}.$$
(3.8)

Taking these expressions we can compute the corresponding tropical fan via the Minkowski sum operation of e.g. gfan [48] and obtain the Speyer-Williams fan. The resulting fan depicted in figure 4. It has five regions of linearity whose boundaries are given by the five rays

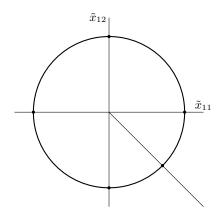
$$\{(1,0), (0,1), (-1,0), (0,-1), (1,-1)\}.$$
(3.9)

Note, the tropical fan as described is parameterised in the space of the  $(\tilde{x}_{11}, \tilde{x}_{12})$  variables. However, as discussed in [21, 47], we can map the five rays above to those of the positive tropical Grassmannian [32] presented in section 2.2, by taking the scalar product of the unit vectors  $e_{ij}$  with the vector of tropicalised Plücker coordinates evaluated at particular values of the  $\tilde{x}$  variables, i.e.

$$\mathbf{ev}: (\tilde{x}_{11}, \tilde{x}_{12}) \mapsto \sum_{1 \le i < j \le 5} \tilde{p}_{ij}(\tilde{x}_{11}, \tilde{x}_{12}) e_{ij}.$$
(3.10)

The five rays given above up to lineality map to  $\{e_{12}, e_{45}, e_{23}, e_{15}, e_{34}\}$ . In particular, the regions between the rays of figure 4 map to the red edges between the corresponding rays in figure 2.

The positive tropical Grassmannian  $\operatorname{Trop}^+(I_{2,5})$  is the image under the map (3.10) of the Speyer-Williams fan. In particular, note that in this case the fan structures map precisely which in fact is the case for all  $\operatorname{Gr}(2, n)$  Grassmannians. It is important to note that the fan structures are not the same in general, as is already the case for  $\operatorname{Gr}(3, 6)$ . While the support of the Speyer-Williams and positive tropical fans coincide, their fan structures do not, and  $\operatorname{Trop}^+(I_{k,n})$  can be viewed as a refinement of the Speyer-Williams fan.



**Figure 4**. The tropical fan  $\operatorname{Trop}^+(I_{2,5})$  via the Speyer-Williams construction.

#### 4 Alphabets and adjacency from the Gröbner fan

In the case of Gr(2, n), the Speyer-Williams fan and the cluster fan coincide since the full set of  $\mathcal{A}$ -coordinates are given solely by Plückers. For the remaining finite-type Grassmannians this is no longer the case. In this section we wish to begin with the Speyer-Williams fan, by tropicalising only the Plücker coordinates, and see where the additional information of the missing  $\mathcal{A}$ -coordinates and adjacency conditions is hidden inside the structure of  $\operatorname{Trop}^+(I_{k,n})$  and  $GF(I_{k,n})$ . This question is inspired by the ideas appearing in [1]. Before we begin let us remind ourselves where we will find this additional information:

**Cluster variables.** The missing  $\mathcal{A}$ -coordinates appear as non-prime factors in the initial ideals of maximal cones of Trop<sup>+</sup>( $I_{k,n}$ ).

**Forbidden pairs.** Upon extending the ideal by the missing  $\mathcal{A}$ -coordinates the rays of  $\operatorname{Trop}^+(I_{k,n})$  span a single maximal Gröbner cone. The initial ideal of this maximal Gröbner cone, of the extended ideal, provides us with a list of monomials which are exactly the forbidden pairs of  $\mathcal{A}$ -coordinates.

#### **4.1** Gr(3,6)

The case of Gr(3, 6) was covered in detail in [1] and we review here the relevant parts of the discussion. The Plücker Ideal  $I_{3,6}$  is generated by three and four-terms relations of the form

$$p_{123}p_{145} + p_{125}p_{134} - p_{124}p_{135} = 0, \dots$$

$$(4.1)$$

$$p_{123}p_{456} - p_{156}p_{234} + p_{146}p_{235} - p_{145}p_{236} = 0, \dots$$

$$(4.2)$$

in the ring of polynomials in the 20 Plücker coordinates  $p_{ijk}$ .

The tropical fan Trop( $I_{3,6}$ ) was studied in [32, 49, 50]. It is simplicial with an *f*-vector given by

$$(1, 65, 550, 1395, 1035),$$
 (4.3)

The 65 rays are given by the permutation copies of the following three basic types [32, 51],

$$e_{123},$$

$$e_{123} + e_{124} + e_{134} + e_{234},$$

$$e_{123} + e_{124} + e_{134} + e_{234} + e_{125} + e_{126}.$$
(4.4)

The positive part of the tropical fan  $\operatorname{Trop}^+(I_{3,6})$  is spanned by 16 rays given by the cyclic copies of the rays listed above in (4.4) (six of the first type, six of the second and four of the third). By restricting to the subfan spanned by these rays we obtain the positive tropical fan  $\operatorname{Trop}^+(I_{3,6})$  with the *f*-vector

$$(1, 16, 68, 104, 52). \tag{4.5}$$

To obtain the corresponding Speyer-Williams fan we first evaluate all Plücker coordinates as minors of the web matrix  $W^{(3,6)}$ , as defined in (3.1). Then we tropicalise the resulting polynomials as explained in section 3. Note, the *frozen* Plücker coordinates  $p_{ii+1i+2}$  are themselves monomial and do not contribute to the structure of the fan. The fan obtained is spanned by 16 rays given by [33]

in  $\tilde{x}$  space with the ordering  $(\tilde{x}_{11}, \tilde{x}_{21}, \tilde{x}_{12}, \tilde{x}_{22})$ . The maximal cones of the fan are fourdimensional regions within which all minors are linear, which can be intersected with the unit sphere to produce 3-dimensional facets of a polyhedral complex. The fan has 48 maximal facets given by 46 terahedra and 2 bipyramids. They themselves have 2dimensional boundaries corresponding to some minor being between two regions of linearity. There are 98 of these 2-dimensional boundaries, which themselves are bounded by 66 edges, which are further bounded by 16 points. The 16 points correspond to the intersection of the rays in (4.6) with the unit sphere. This information can be summarised by the *f*-vector given by  $f_{3,6} = (1, 16, 66, 98, 48)$ . Sometimes we would also like to keep information on the number of vertices of each facet, for this we use the notation

$$f_{3,6} = (1_0, 16_1, 66_2, 98_3, 46_4 + 2_5),$$

where we understand the right most element as 46 tetrahedrons (4-vertex objects) and two, non-simplicial, bipyramids (5-vertex objects). The two bipyramids will play a role in the following discussion so let us discuss their structure in more detail. They are each spanned by 5 rays from (4.6) as follows,

$$b_{1} = \operatorname{span}\{(-1, 0, 0, 1), (0, 0, 1, 0), (-1, 0, 0, 0), (0, 1, 0, 0), (0, 1, 1, -1)\} \\ := \operatorname{span}\{b_{11}, b_{12}, b_{13}, b_{14}, b_{15}\}, \\ b_{2} = \operatorname{span}\{(1, -1, -1, 0), (1, 0, -1, 0), (1, -1, 0, 0), (0, 0, 0, -1), (1, 0, 0, -1)\} \\ := \operatorname{span}\{b_{21}, b_{22}, b_{23}, b_{24}, b_{25}\}.$$

$$(4.7)$$

As discussed before in the example of Gr(2,5), we can map the Speyer-Williams fan to the positive tropical Grassmannian discussed in section 2 via the map

$$\mathbf{ev}: (\tilde{x}_{11}, \tilde{x}_{12}, \tilde{x}_{21}, \tilde{x}_{22}) \mapsto \sum_{1 \le i < j < k \le 6} \tilde{p}_{ijk}(\tilde{x}_{11}, \tilde{x}_{12}, \tilde{x}_{21}, \tilde{x}_{22}) e_{ijk},$$
(4.8)

with the  $e_{ijk}$  comprising the unit vectors in  $\mathbb{R}^{20}$ ,

$$e_{123} = (1, 0, \dots, 0),$$
  
 $\vdots$   
 $e_{456} = (0, \dots, 0, 1).$  (4.9)

For example, modulo the lineality shift of  $I_{3,6}$ , the five rays of the bipyramid  $b_1$  map to (as discussed in [21])

$$b_{11} \mapsto e_{123} + e_{124} + e_{125} + e_{126} + e_{134} + e_{234},$$
  

$$b_{12} \mapsto e_{123} + e_{124} + e_{134} + e_{234},$$
  

$$b_{13} \mapsto e_{125} + e_{126} + e_{156} + e_{256},$$
  

$$b_{14} \mapsto e_{345} + e_{346} + e_{356} + e_{456},$$
  

$$b_{15} \mapsto e_{123} + e_{124} + e_{134} + e_{234} + e_{345} + e_{346}.$$
  
(4.10)

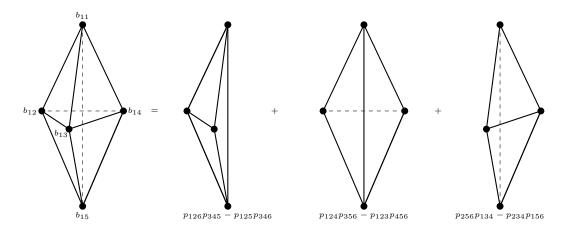
In particular note that image of the centre of the bipyramid can be represented in two equivalent ways (as always modulo lineality),

$$\mathbf{ev}(b_{11}) + \mathbf{ev}(b_{15}) = \mathbf{ev}(b_{12}) + \mathbf{ev}(b_{13}) + \mathbf{ev}(b_{14}).$$
(4.11)

The sixteen rays given in (4.6) map to the sixteen rays of  $\operatorname{Trop}^+(I_{3,6})$  but the fan structures differ slightly. In fact,  $\operatorname{Trop}^+(I_{3,6})$  which, as introduced in section 2, inherits its fan structure from the Gröbner fan  $GF(I_{3,6})$ , has two additional edges and six additional triangles in comparison to the Speyer-Williams fan. In other words, as a fan  $\operatorname{Trop}^+(I_{3,6})$ is a refinement of (the image under **ev** of) the Speyer-Williams fan. The additional edges and triangles actually slice each of the two bipyramids into three tetrahedra, each with its own initial ideal, as illustrated in figure 5.

Related to the splitting of the bipyramids (and in contrast with the Gr(2, n) case), the 16 rays of  $Trop^+(I_{3,6})$  do not single out an individual maximal Gröbner cone of  $GF(I_{3,6})$ . In fact, calculating initial ideals inside the span of the 16 positive rays of  $Trop^+(I_{3,6})$  we find nine maximal Gröbner cones. The splitting of the bipyramids can be understood as each bipyramid intersecting three of the nine maximal Gröbner cones spanned by the rays of  $Trop^+(I_{3,6})$ .

As shown in [1], we can resolve the span of the positive tropical rays into a single maximal Gröbner cone by extending the ideal. To decide how to extend the ideal we search the maximal cones of  $\text{Trop}^+(I_{3,6})$  for initial ideals which are not prime and whose factors will provide for us the missing  $\mathcal{A}$ -coordinates. In fact the bipyramids themselves



**Figure 5.** The image of the bipyramid  $b_1$  inside  $\operatorname{Trop}^+(I_{3,6})$ , on the left hand we have the full bipyramid with its 5 rays, on the right the bipyramid is split into three tetrahedra by the structure of the Gröbner fan. Each tetrahedron is labelled by the quadratic non-prime factor found in the initial ideal. Note all three expressions are equivalent modulo the Plücker relations.

are the source of the non-prime initial ideals. The three non-prime initial ideals generated inside  $b_1$  can be written as the intersection of two prime ideals as

$$\begin{split} &\text{in}_{b_1 \setminus \{b_{12}\}}(I_{3,6}) = \langle \text{in}_{b_1 \setminus \{b_{12}\}}(I_{3,6}) \cup M_1 \rangle \cap \langle \text{in}_{b_1 \setminus \{b_{12}\}}(I_{3,6}) \cup \{p_{256}p_{134} - p_{234}p_{156}\} \rangle, \\ &\text{in}_{b_1 \setminus \{b_{13}\}}(I_{3,6}) = \langle \text{in}_{b_1 \setminus \{b_{13}\}}(I_{3,6}) \cup M_1 \rangle \cap \langle \text{in}_{b_1 \setminus \{b_{13}\}}(I_{3,6}) \cup \{p_{124}p_{356} - p_{123}p_{456}\} \rangle, \\ &\text{in}_{b_1 \setminus \{b_{14}\}}(I_{3,6}) = \langle \text{in}_{b_1 \setminus \{b_{14}\}}(I_{3,6}) \cup M_1 \rangle \cap \langle \text{in}_{b_1 \setminus \{b_{14}\}}(I_{3,6}) \cup \{p_{126}p_{345} - p_{125}p_{346}\} \rangle. \end{split}$$

Similarly for  $b_2$  we have the cyclic copy of the above given by

$$\begin{split} &\operatorname{in}_{b_{2}\setminus\{b_{22}\}}(I_{3,6}) = \langle \operatorname{in}_{b_{2}\setminus\{b_{22}\}}(I_{3,6}) \cup M_{2} \rangle \cap \langle \operatorname{in}_{b_{2}\setminus\{b_{22}\}}(I_{3,6}) \cup \{p_{145}p_{236} - p_{123}p_{456}\} \rangle, \\ &\operatorname{in}_{b_{2}\setminus\{b_{23}\}}(I_{3,6}) = \langle \operatorname{in}_{b_{2}\setminus\{b_{23}\}}(I_{3,6}) \cup M_{2} \rangle \cap \langle \operatorname{in}_{b_{2}\setminus\{b_{23}\}}(I_{3,6}) \cup \{p_{136}p_{245} - p_{126}p_{345}\} \rangle, \\ &\operatorname{in}_{b_{2}\setminus\{b_{24}\}}(I_{3,6}) = \langle \operatorname{in}_{b_{2}\setminus\{b_{24}\}}(I_{3,6}) \cup M_{2} \rangle \cap \langle \operatorname{in}_{b_{2}\setminus\{b_{24}\}}(I_{3,6}) \cup \{p_{156}p_{234} - p_{146}p_{235}\} \rangle. \end{split}$$

In the above equations we understand  $in_{b_1 \setminus \{b_{11}\}}(I_{3,6})$  for instance, as the initial ideal of  $I_{3,6}$  associated to the cone spanned by the rays  $b_1 \setminus \{b_{11}\}$  and we have defined the sets of monomials

$$M_1 = \{ p_{235}, p_{236}, p_{245}, p_{246}, p_{135}, p_{136}, p_{145}, p_{146} \}, M_2 = \{ p_{124}, p_{125}, p_{134}, p_{135}, p_{246}, p_{256}, p_{346}, p_{356} \}.$$

Most importantly notice the three quadratic non-prime factors appearing in each cone modulo the Plücker ideal are equivalent to either

$$p_{12[34]56} \text{ or } p_{23[45]61},$$
 (4.12)

where we have defined  $p_{ij[kl]mn} = p_{ijl}p_{kmn} - p_{ijk}p_{lmn}$ . These are exactly the two missing  $\mathcal{A}$ -coordinates, which along with the 14 monomials contained in  $M_1 \cup M_2$ , make up the full set of active  $\mathcal{A}$ -coordinates.

The appearance of the quadratic  $\mathcal{A}$ -coordinates (4.12) suggests extending our original ideal by including them as new variables. Extending the Plücker ideal as

$$I'_{3,6} = I_{3,6} \cap \langle q_1 - p_{12[34]56} \rangle \subset \mathbb{R}[p_{123}, \dots, p_{456}, q_1], \tag{4.13}$$

defines a new ideal and hence a new Gröbner fan  $GF(I'_{3,6})$ . The Speyer-Williams fan can also be refined by evaluating  $q_1 = p_{12[34]56}$ ,<sup>3</sup> in terms of minors of the web matrix and including it in the set of tropical polynomials defining the fan. In this way we obtain a refined Speyer-Williams fan with *f*-vector

$$f'_{3,6} = (1_0, 16_1, 66_2, 99_3, 48_4 + 1_5), \tag{4.14}$$

where we now have only a single five-vertex bipyramid  $(b_2)$ . The bipyramid  $b_1$  has been broken into two tetrahedra with an additional triangle separating them. Indeed, let us consider the natural extension of the map (4.8)

$$\mathbf{ev}': (\tilde{x}_{11}, \tilde{x}_{12}, \tilde{x}_{21}, \tilde{x}_{22}) \mapsto \sum_{1 \le i < j < k \le 6} \tilde{p}_{ijk}(\tilde{x}_{11}, \tilde{x}_{12}, \tilde{x}_{21}, \tilde{x}_{22}) e_{ijk} + \tilde{q}_1(\tilde{x}_{11}, \tilde{x}_{12}, \tilde{x}_{21}, \tilde{x}_{22}) e_{q_1}$$
(4.15)

which takes a point in the refined Speyer-Williams fan into  $\operatorname{Trop}^+(I'_{3,6})$ , where we introduced  $e_{q_1}$  as a unit vector for the  $q_1$  direction. Under this map the images of the five rays of  $b_1$  are now [47]

$$b_{11} \mapsto e_{123} + e_{124} + e_{125} + e_{126} + e_{134} + e_{234} + e_{q_1},$$
  

$$b_{12} \mapsto e_{123} + e_{124} + e_{134} + e_{234} + e_{q_1},$$
  

$$b_{13} \mapsto e_{125} + e_{126} + e_{156} + e_{256} + e_{q_1},$$
  

$$b_{14} \mapsto e_{345} + e_{346} + e_{356} + e_{456} + e_{q_1},$$
  

$$b_{15} \mapsto e_{123} + e_{124} + e_{134} + e_{234} + e_{345} + e_{346} + e_{q_1},$$
  

$$(4.16)$$

and the relation (4.11) no longer holds if we replace ev with ev'.

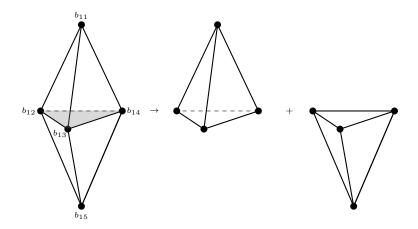
The transition from  $f_{3,6}$  to  $f'_{3,6}$  can be seen as adding a triangle to the equator of the bipyramid  $b_1$  as shown in figure 6 which now splits into two tetrahedra whose images in  $\operatorname{Trop}^+(I'_{3,6})$  now correspond to prime initial ideals. In other words, at least as far as these five rays are concerned, the fan structures of the extended Speyer-Williams fan and the positive tropical fan of  $I'_{3,6}$  match. Also we find that the sixteen rays of  $\operatorname{Trop}^+(I'_{3,6})$  now span only 3 maximal Gröbner cones inside  $GF(I'_{3,6})$ .

Extending the ideal further as

$$I_{3,6}^{\text{ext}} = I_{3,6}' \cap \langle q_2 - p_{23[45]61} \rangle \subset \mathbb{R}[p_{123}, \dots, p_{456}, q_1, q_2], \tag{4.17}$$

defines yet again a new Gröbner fan  $GF(I_{3,6}^{\text{ext}})$ . We can further refine the Speyer-Williams fan by tropically evaluating  $q_2 = p_{23[45]61}$ . This fan was referred to as the cluster fan in [47]

<sup>&</sup>lt;sup>3</sup>Similar to the notation in (2.1)  $p_{12[34]56} = p_{124}p_{356} - p_{123}p_{456}$ 



**Figure 6**. The bipyramid  $b_1$  maps to two tetrahedra inside  $\operatorname{Trop}^+(I'_{3,6})$ , both of which are associated to prime ideals.

because it coincides with the g-vector fan of the Gr(3,6) cluster algebra. The cluster fan has f-vector

$$f_{3.6}^{\text{ext}} = (1_0, 16_1, 66_2, 100_3, 50_4) \tag{4.18}$$

and now its fan structure coincides with that of  $\operatorname{Trop}^+(I_{3,6}^{\text{ext}})$  which is the image of the cluster fan under the natural extension of (4.15),

$$\mathbf{ev}^{\text{ext}} : (\tilde{x}_{11}, \tilde{x}_{12}, \tilde{x}_{21}, \tilde{x}_{22}) \mapsto \sum_{1 \le i < j < k \le 6} \tilde{p}_{ijk}(\tilde{x}_{11}, \tilde{x}_{12}, \tilde{x}_{21}, \tilde{x}_{22}) e_{ijk}$$

$$+ \tilde{q}_1(\tilde{x}_{11}, \tilde{x}_{12}, \tilde{x}_{21}, \tilde{x}_{22}) e_{q_1} + \tilde{q}_2(\tilde{x}_{11}, \tilde{x}_{12}, \tilde{x}_{21}, \tilde{x}_{22}) e_{q_2}.$$

$$(4.19)$$

Note that  $\operatorname{Trop}^+(I_{3,6}^{\text{ext}})$  contains no non-prime maximal cones! The effect of the further extension can be viewed as adding a triangle to the equator of bipyramid  $b_2$  in an exact copy of figure 6, while the rays of  $b_2$  are extended in the new  $e_{q_2}$  direction in an analogous manner to (4.16).

Furthermore, the rays of  $\operatorname{Trop}^+(I_{3,6}^{\text{ext}})$  now span a single maximal Gröbner cone whose initial ideal is generated by

 $\{q_1q_2, p_{124}q_2, p_{125}q_2, p_{134}q_2, p_{135}q_1, p_{135}q_2, p_{124}p_{135}, p_{136}q_1, p_{124}p_{136}, p_{125}p_{136}, p_{145}q_1, p_{145}q_{15}, p_{145}q_{15}$ 

 $p_{146}q_1, p_{125}p_{146}, p_{135}p_{146}, p_{235}q_1, p_{124}p_{235}, p_{134}p_{235}, p_{146}p_{235}, p_{236}q_1, p_{124}p_{236}, p_{125}p_{236}, p_{125}p$ 

 $p_{134}p_{236}, p_{135}p_{236}, p_{145}p_{236}, p_{245}q_1, p_{134}p_{245}, p_{135}p_{245}, p_{136}p_{245}, p_{246}q_1, p_{246}q_2, p_{125}p_{246}, p_{145}p_{246}, p_{145}p_{24}, p_{145}p_{24}, p_{145}p_{24}, p_{145}p_{24}, p_{145}p_{24}, p_{145}p_{24}, p_{145}p_{24}, p_{145}p_{24}, p_{145}p_{24}, p$ 

 $p_{134}p_{246}, p_{135}p_{246}, p_{136}p_{246}, p_{145}p_{246}, p_{235}p_{246}, p_{256}q_2, p_{134}p_{256}, p_{135}p_{256}, p_{136}p_{256}, p_{145}p_{256}, p_{145}p_{25}, p_{145}p_{1$ 

 $p_{146}p_{256}, p_{346}q_2, p_{125}p_{346}, p_{135}p_{346}, p_{145}p_{346}, p_{235}p_{346}, p_{245}p_{346}, p_{356}q_2, p_{124}p_{356}, p_{145}p_{356}, p_{145}p_{156}, p_{156}p_{156}, p_{156}p_{156}, p_{156}p_{156}, p_{156}p_{156}, p_{1$ 

 $p_{146}p_{356}, p_{245}p_{356}, p_{246}p_{356}\},$ 

which are exactly the 54 forbidden pairs of  $\mathcal{A}$ -coordinates (i.e. those pairs which never sit together in a cluster).

In summary, we have seen explicitly that the non-prime initial ideals provided the missing quadratic  $\mathcal{A}$  coordinates. These new variables also tell us how to extend the

Plücker ideal  $I_{3,6}$  to a new ideal  $I_{3,6}^{\text{ext}}$  so as to find a single maximal Gröbner cone inside the span of the rays of the corresponding positive tropical fan  $\text{Trop}^+(I_{3,6}^{\text{ext}})$ . The initial ideal associated to this maximal Gröbner cone provides the adjacency information of the cluster algebra. Thus these two key pieces of information ( $\mathcal{A}$ -coordinates and adjacency) can be obtained from the Gröbner structure of the relevant ideals without direct reference to the cluster algebra at all.

## **4.2** $Gr(3,7) \cong Gr(4,7)$

We now go beyond the results of [1] and consider  $\operatorname{Gr}(3,7) \cong \operatorname{Gr}(4,7)$ . We again wish to compare the structure of the positive tropical fan  $\operatorname{Trop}^+(I_{3,7})$  to the Speyer-Williams fan obtained from tropicalising the 35 Plücker coordinates evaluated as minors of the web matrix.

Let us begin by detailing the structure of the full tropical fan  $\text{Trop}(I_{3,7})$  which was calculated in [49, 52]. The tropical fan  $\text{Trop}(I_{3,7})$  is simplicial with an *f*-vector given by

$$(1, 721, 16800, 124180, 386155, 522585, 252000).$$

$$(4.20)$$

where the 721 rays are given by the permutation copies of [53]

 $b_{1,1234567} = e_{123},$   $b_{2,1234567} = e_{123} + e_{124} + e_{134} + e_{234},$   $b_{3,1234567} = e_{123} + e_{124} + e_{125} + e_{126} + e_{127},$   $b_{4,1234567} = e_{123} + e_{124} + e_{125} + e_{126} + e_{127} + e_{134} + e_{234},$   $b_{5,1234567} = e_{123} + e_{124} + e_{125} + e_{126} + e_{127} + e_{134} + e_{156} + e_{234} + e_{256},$   $b_{6,1234567} = b_{3,1234567} + b_{3,3456712} + b_{3,6712345}.$ (4.21)

The positive part of the tropical fan  $\operatorname{Trop}^+(I_{3,7})$  is spanned by 49 rays given by the cyclic copies of

$$\{b_{1,1234567}, b_{2,1234567}, b_{3,1234567}, b_{4,1234567}, b_{4,1562347}, b_{5,1234675}, b_{6,1234567}\}.$$
(4.22)

By restricting to the subfan spanned by these rays we obtain the positive tropical fan  $\operatorname{Trop}^+(I_{3,7})$  with the *f*-vector

$$(1, 49, 490, 1964, 3633, 3192, 1064).$$
 (4.23)

Now we turn our attention to the Speyer-Williams fan [33] whose structure is summarised by the f-vector

$$f_{3,7} = (1, 42, 392, 1463, 2583, 2163, 595_6 + 63_7 + 28_8 + 7_9), \tag{4.24}$$

where we have included the information on the number of vertices for the maximal cones only (in the final element above). This fan is spanned by 42 rays given by the first six elements of (4.22) along with their cyclic copies and has 6-dimensional maximal cones. Note that  $\text{Trop}^+(I_{3,7})$  has seven more rays than the Speyer-Williams fan. These additional

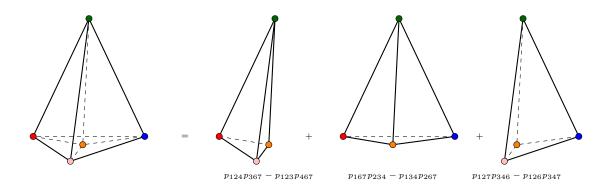


Figure 7. On the left a 6-vertex simplex of the Speyer-Williams fan (note the green node at the top really stands for a combination of three vertices). On the right the refinement of the Speyer-Williams cone by the positive tropical fan  $\operatorname{Trop}^+(I_{3,7})$ . Note the appearance of a new orange vertex.

rays correspond to  $b_{6,1234567}$  and cyclic copies. As discussed in [21, 53] these extra rays actually appear in the middle of triangular faces formed by three rays of  $b_3$ -type.

As discussed in previous cases, we can map the Speyer-Williams fan to the positive tropical Grassmannian  $\text{Trop}^+(I_{3,7})$  via the evaluation map,

$$\mathbf{ev}: (\tilde{x}_{11}, \dots, \tilde{x}_{23}) \mapsto \sum \tilde{p}_{ijk}(\tilde{x}_{11}, \dots, \tilde{x}_{23})e_{ijk}.$$

$$(4.25)$$

Once again we find that the fan structure of  $\operatorname{Trop}^+(I_{3,7})$  is a refinement of the (image of the) Speyer-Williams fan. Let us explore this refinement in more detail.

Let us look at each of the maximal cones in turn starting with the six-vertex cones of the Speyer-Williams fan. There 595 simplicial maximal cones of the Speyer-Williams fan: 567 are un-refined by the positive tropical fan, and 28 are refined by the positive tropical fan coming in 2 dihedral classes. As an example, a representative from the first dihedral class has the rays given by

$$\{b_{1,5671234}, b_{4,1234567}, b_{4,6712345}, b_{3,6712345}, b_{3,3456712}, b_{3,1234567}\}$$
(4.26)

and is refined by the tropical fan with three simplices given by

$$\begin{pmatrix} b_{1,5671234} & b_{4,1234567} & b_{4,6712345} & b_{3,6712345} & b_{3,1234567} & b_{6,1234567} \\ b_{1,5671234} & b_{4,1234567} & b_{4,6712345} & b_{3,6712345} & b_{3,3456712} & b_{6,1234567} \\ b_{1,5671234} & b_{4,1234567} & b_{4,6712345} & b_{3,3456712} & b_{3,1234567} & b_{6,1234567} \\ \end{pmatrix}.$$
(4.27)

Note, in particular the appearance of the extra positive ray  $b_{6,1234567}$ , which is a ray of  $\text{Trop}^+(I_{3,7})$  but not of the Speyer-Williams fan. We give a sketch showing the appearance of the extra ray in the middle of a face as shown in figure 7. Below each part of the refinement in figure 7 we note the missing binomial contained as a factor in the non-prime initial ideal associated to this cone of  $\text{Trop}^+(I_{3,7})$ .

The second dihedral set of 6-vertex cones which are refined by the positive tropical fan have the rays (in the Speyer-Williams fan) given by

$$\{b_{1,5671234}, b_{1,3456712}, b_{4,6712345}, b_{3,6712345}, b_{3,3456712}, b_{3,1234567}\}$$
(4.28)

which are split into three simplices (in the positive tropical fan) with the rays

$$\begin{pmatrix} b_{1,5671234} & b_{1,3456712} & b_{4,6712345} & b_{3,6712345} & b_{3,1234567} & b_{6,1234567} \\ b_{1,5671234} & b_{1,3456712} & b_{4,6712345} & b_{3,6712345} & b_{3,3456712} & b_{6,1234567} \\ b_{1,5671234} & b_{1,3456712} & b_{4,6712345} & b_{3,3456712} & b_{3,1234567} & b_{6,1234567} \end{pmatrix}.$$

$$(4.29)$$

Again we note the appearance of a spurious positive ray. The picture for this cone looks identical to figure 7, with the same missing binomials appearing, but with the coloured nodes referring to (4.28) and (4.29) instead of (4.26) and (4.27).

The remaining 6-vertex cones coincide in both fans, i.e. are unrefined. The remaining maximal cones have 7, 8 or 9 vertices. We explore their refinements in appendix A, with particular focus on the 'missing binomials' appearing as factors of non-prime initial ideals.

In summary, we find all 14 missing quadratic  $\mathcal{A}$ -coordinates (with each one arising as a factor of non-prime initial ideals in many cones).

$$q_{51} = p_{12[34]56} \text{ and cyclic,}$$

$$q_{61} = p_{61[23]45} \text{ and cyclic,}$$
(4.30)

appearing as non-prime factors of the  $7_9$  maximal cones. By extending the ideal by all 14 variables as

$$I_{3,7}^{\text{ext}} = I_{3,7} \cap \langle q_{51} - p_{12[34]56}, \dots, q_{67} - p_{56[12]34} \rangle \subset \mathbb{R}[p_{123}, \dots, p_{567}, q_{51}, \dots, q_{67}], \quad (4.31)$$

the extended Speyer-Williams fan (again the cluster fan) and the positive tropical fan  $\operatorname{Trop}^+(I_{3,7}^{\text{ext}})$  become identical simplicial fans with the *f*-vector

$$f_{3.7}^{\text{ext}} = (1, 42, 399, 1547, 2856, 2499, 833).$$

Moreover, the rays of  $\operatorname{Trop}^+(I_{3,7}^{\text{ext}})$  span a single maximal Gröbner cone in  $GF(I_{3,7}^{\text{ext}})$  whose initial ideal gives us precisely the list of <u>462 forbidden neighbours</u> for 7-point scattering amplitudes as given in [18].

#### **4.3** Gr(3,8)

In the space of variables  $(\tilde{x}_{11}, \tilde{x}_{12}, \tilde{x}_{13}, \tilde{x}_{14}, \tilde{x}_{21}, \tilde{x}_{23}, \tilde{x}_{24})$  the Speyer-Williams fan is spanned by 120 rays with 8-dimensional maximal cones. The structure of the fan is summarised by the f-vector

 $f_{3,8} = (1, 120, 2072, 14088, 48544, 93104, 100852, 57768, 13612),$ 

with the 13612 maximal cones given by

$$9672_8 + 1696_9 + 1092_{10} + 480_{11} + 416_{12} + 104_{13} + 88_{14} + 32_{15} + 24_{16} + 8_{17}.$$

For the case of Gr(3, 8) there two new features which we wish to emphasise. First, there are 8 more  $\mathcal{A}$ -coordinates than there are rays of the Speyer-Williams fan, which must appear

when we begin to extend the ideal. Second, the  $\mathcal{A}$ -coordinates not only contain expressions quadratic in the Plückers but also contain cubic expressions given by

$$\{p_{12[34]5[67]89}, p_{12[35]8[67]45}, p_{12[34]8[67]35}\}$$
 and cylic, (4.32)

where we have made the definition  $p_{ij[kl]m[nr]st} = p_{ijl}p_{km[nr]st} - p_{ijk}p_{lm[nr]st}$ .

Up to this point we have been searching the entire positive part of the tropical fan for the appearance of non-prime factors. However, for Gr(3,8) this calculation becomes cumbersome and instead we satisfy ourselves with a local approach. Concretely, we begin by searching for non-prime factors appearing in the 17-vertex cone spanned by the rays

1	1	2	3	4	5	6	7	8	9	10	11	12	13	14	15	16	17	)			
	$\begin{array}{c} 0 \\ 0 \end{array}$	$\begin{array}{c} 0 \\ 0 \end{array}$	${\begin{array}{c}1\\1\\-1\\0\\-1\\0\\1\end{array}}$	$_0^{-1}$	$\begin{array}{c} 0 \\ 0 \end{array}$	$_{-1}^{0}$	$\begin{array}{c} 0 \\ -1 \\ 0 \end{array}$	$     \begin{array}{c}       0 \\       -1 \\       -1     \end{array} $	$\begin{array}{c} -1 \\ 0 \\ 0 \end{array}$	$\begin{array}{c} -1 \\ 0 \\ 0 \end{array}$	$\begin{array}{c} 0 \\ 0 \end{array}$	$\begin{array}{c} -1 \\ 0 \\ 0 \end{array}$	$-1 \\ 0 \\ -1$	0	$\begin{array}{c} 0 \\ 0 \end{array}$	$\begin{array}{c} 0 \\ 0 \end{array}$	$\begin{array}{c} 0 \\ 0 \end{array}$	),		(4.33)	)

where we find the following 6 non-prime factors<sup>4</sup>

$$\{p_{81[23]45}, p_{81[23]46}, p_{56[78]14}, p_{56[78]23}, p_{23[45]71}, p_{23[46]71}\}.$$
(4.34)

Including these factors into the Speyer-Williams calculation we arrive at a fan with 121 rays and maximal cones given by

$$11454_8 + 1696_9 + 971_{10} + 412_{11} + 328_{12} + 89_{13} + 69_{14} + 28_{15} + 17_{16} + 7_{17}$$

where the new ' $18^{\text{th}}$ ' ray is given by

$$(1, 1, -2, 0, 0, -1, 0, 1).$$
 (4.35)

An interesting question to ask is which of the maximal cones contain this new ray, they are given by

$$270_8 + 70_9 + 46_{10} + 18_{11} + 10_{12} + 4_{13} + 2_{14}. \tag{4.36}$$

By searching inside the  $2_{14}$  cones we are able to find the final cyclic type of quadratic non-prime factor  $p_{56[71]23}$ .

Alternatively, we can consider the subset of cones of the extended Speyer-Williams fan contained in the span of the 18 rays described above, we find 24 such cones interestingly all of which contain the new ray; 18 of these cones are simplicial; 4 contain 9 vertices; and 2 contain 10 vertices and are spanned by the rays

$$(1, 2, 3, 4, 5, 6, 10, 15, 17, 18)$$
 and  $(6, 8, 9, 11, 12, 13, 14, 16, 17, 18)$ . (4.37)

Searching in these two cones we find three additional non-prime factors

 $\{p_{23[46]71}p_{578} - p_{23[45]71}p_{678}, p_{81[23]46}p_{578} - p_{81[23]45}p_{678}, p_{56[78]23}p_{146} - p_{123}p_{456}p_{678}\}, (4.38)$ 

which are the three classes of cubic  $\mathcal{A}$ -coordinates appearing in the Gr(3,8) cluster algebra.

<sup>&</sup>lt;sup>4</sup>Note, we do not include non-prime factors related to those in the list by rotation.

To summarise we are able to find one representative from each cyclic class of  $\mathcal{A}$ coordinate appearing as non-prime factors of initial ideals of  $\operatorname{Trop}^+(I_{3,8})$  and its various
extensions. Therefore, cyclic symmetry would suggest, by following the same procedure
outlined above considering each of the 7 remaining 17-vertex cones in turn we would recover
the entire set of  $\mathcal{A}$ -coordinates in a similar manner. After extending the ideal by the full
set of  $\mathcal{A}$ -coordinates the results of [35], applying to any finite cluster algebra of geometric
type, tell us that we would again see the rays of  $\operatorname{Trop}^+(I_{3,8}^{ext})$  spanning a single maximal
Gröbner cone whose initial ideal is generated by the forbidden pairs of the cluster algebra.

The procedure outline above for locally resolving the Speyer-Williams fan by inclusion of non-prime factors may prove useful when considering how square root letters appear for Gr(4,8) from the structure of the Gröbner fan. Such square root letters were accounted for in [22] by considering coordinates associated to so called *limit rays* arising from limits of affine sequences of mutations inside the Gr(4,8) cluster algebra. Interestingly, these limit rays also show up in the Speyer-Williams fan, where in particular they appear in non-simplicial maximal cones with large numbers of vertices. Presumably, by attempting to locally resolve these special cones, as we have done for the case of Gr(3,8) above, we will see that instead an infinite number of non-prime factors is now needed corresponding to the infinite mutation sequences in the cluster picture. It would be interesting to study this in further detail.

#### 5 Missing binomials and cluster variables

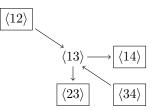
In this section we explain the techniques used above from the point of view of cluster algebras.

#### 5.1 The g-vectors as multiweights

We have already seen Fomin-Zelevinsky's g-vectors appear as rays in the Speyer-Williams fan. However, the g-vectors of Plücker coordinates or more general cluster variables carry even more information, they can detect whether a given cone in the Speyer-Williams fan corresponds to a prime cone in the positive tropicalization. To be more precise we need to understand how g-vectors can be used as *multiweights* to compute initial ideals. For this purpose it is important to work with *extended* g-vectors that can be computed recursively (just like normal g-vectors) with the only difference that also frozen nodes are taken into account. For example, for the quiver in figure 3, extended g-vectors lie in  $\mathbb{Z}^2$  (with entries corresponding to the active nodes  $\langle 13 \rangle$  and  $\langle 14 \rangle$ ). Given a quiver Q with n active and m frozen nodes we associate a matrix  $B := B_Q$  with entries for i an active node and jany node

$$b_{ij} := \#\{\text{arrows } i \to j\} - \#\{\text{arrows } j \to i\}.$$

This matrix defines a partial order on  $\mathbb{Z}^{n+m}$  (called the *dominance order*) where  $p >_B q$  if and only if there is an element  $r \in \mathbb{Z}_{\geq 0}^{n+m}$  such that p = q + Br. **Example 5.1.** Consider Gr(2,4) with Plücker ideal  $I = (p_{12}p_{34} - p_{13}p_{24} + p_{14}p_{23})$ . For the cluster algebra consider the quiver



with associated matrix  $B = (0, 1, -1, -1, 1)^T$  and cluster variables  $s = \{p_{13}, p_{12}, p_{14}, p_{23}, p_{34}\}$ . The extended g-vectors of Plücker coordinates are

$$\begin{split} g_s(p_{13}) &= (1,0,0,0,0), \qquad g_s(p_{24}) = (-1,0,1,1,0), \qquad g_s(p_{12}) = (0,1,0,0,0), \\ g_s(p_{14}) &= (0,0,1,0,0), \qquad g_s(p_{23}) = (0,0,0,1,0), \qquad g_s(p_{34}) = (0,0,0,0,1). \end{split}$$

In particular, the monomials in the Plücker relation have multiweights

$${p_{12}}^{(0,1,0,0,1)}_{p_{12}} - {p_{13}}^{(0,0,1,1,0)}_{p_{13}} + {p_{14}}^{(0,0,1,1,0)}_{p_{14}}_{p_{23}}.$$

 $As (0,0,1,1,0) = (0,1,0,0,1) + (0,1,-1,-1,1) we deduce that (0,0,1,1,0) >_B (0,1,0,0,1).$ Hence, the initial form is  $-p_{13}p_{24} + p_{14}p_{23}$  which corresponds to a prime cone in Trop<sup>+</sup>(I<sub>2,4</sub>).

We are interested in the Gröbner fan structure of the positive part of the tropicalization. However, it is computationally expensive to compute initial ideals which is necessary to determine this fan structure. The Speyer-Williams fan (a coarsening of the g-vector fan) on the other hand is easy and fast to compute. We can use the g-vectors to compare the fan structures as follows.<sup>5</sup>

**Corollary 5.1.** [55, 56] Let s be a seed in the cluster algebra of a finite type Grassmannian with n mutable and m frozen nodes. Then the initial ideal of the Plücker ideal with respect to the g-vectors of Plücker coordinates obtained from s is prime if and only if every point in  $\mathbb{Z}^n \times \mathbb{Z}_{>0}^m$  can be obtained as a positive combination of g-vectors of Plücker coordinates.

In fact, it suffices to consider the usual g-vectors in  $\mathbb{Z}^n$  to verify the Corollary as the extended g-vectors of frozen Plücker coordinates always span  $\mathbb{Z}_{\geq 0}^m$ . Whenever the Corollary fails to apply to a seed s, the initial ideal obtained from extended g-vectors is not prime and we expect to find missing binomials.

#### 5.2 Cluster variables as missing binomials

The phenomenon observed in section 4 can be explained from the point of view of cluster algebras and tropical geometry: for every cluster algebra of finite type A there exists a presentation

$$A \cong \mathbb{K}[x_1, \dots, x_N]/I \tag{5.1}$$

<sup>&</sup>lt;sup>5</sup>This result relies on the theory of valuations and its connection to tropicalization [54]. Further, it uses the fact that the g-vector fan is complete for finite type cluster algebras.

where  $x_1, \ldots, x_N$  correspond to all cluster variables and I is the ideal obtained by saturating<sup>6</sup> the ideal generated by all exchange relations (see e.g. [57], section 6.8). In the cases of the Grassmannian Gr(3, n) with  $n \in \{6, 7, 8\}$  the Plücker ideal  $I_{3,n}$  is obtained from  $I = I^{\text{ext}_{3,n}}$  by eliminating those cluster variables that are not Plücker coordinates.<sup>7</sup>

Let us assume additionally that A is positively graded, as is true for the cases of interest to us. As mentioned above [35] implies that there exists a unique maximal cone C in the Gröbner fan of I whose initial ideal encodes the adjacencies of all cluster variables. In other words, the monomial initial ideal is the Stanley-Reisner ideal of the cluster complex (cluster fan) that is minimally generated by degree two monomials of cluster variables that do not occur in any seed together (these are called *non-compatible* by Fomin-Zelevinsky). The results predict that the procedure of adding (the right) missing binomials as variables to the presentation converges.

The presentation of A from (5.1) is of such a form that is for every seed s there exists a maximal cone  $\tau_s \in \text{Trop}^+(I)$  such that the initial ideal with respect to the multiweights determined by g-vectors of s coincides with the initial ideal of  $\tau_s$  [56]. Moreover, in I we have a trinomial corresponding to every exchange relation. It is not hard to see that for every such trinomial

$$f = xx' - M_1 - M_2$$

and every seed s we have either

Suppose that one of the monomials  $M_1, M_2$  consists of a single variable, say  $M_1 = y$ . Then  $y = xx' - M_2$  holds in A. All exchange relations involving y yield 4-term relations in I. For example,

$$yy' = M'_1 + M'_2$$
 gives  $y'(xx' - M_2) = M'_1 + M'_2$ .

What we observe in the case of the Grassmannians is that there exist weight vectors  $w \in \operatorname{Trop}^+(I)$  such that

$$in_w(y'(xx' - M_2) - M_1' - M_2') = y'(xx' - M_2).$$

In this case  $xx' - M_2$  is a missing binomial contributing to the fact that the initial ideal is not prime. If on the other hand y is a variable in the chosen presentation of the cluster algebra then  $xx' - y - M_2 \in I$  and so  $xx' - M_2 \in in_w(I)$  is not 'missing'. This may happen precisely for those seeds (and their g-vectors) for which the Corollary 5.1 does not apply.

**Example 5.2.** For Gr(3,6) take the seed s with mutable cluster variables  $p_{124}, p_{246}, p_{256}$  and  $p_{346}$  (see figure 8). According to the computations in section 4.1 and the results in [1]

<sup>&</sup>lt;sup>6</sup>The saturation of an ideal is an ideal that contains all elements f for which there exist a monomial  $x_1^{m_1} \cdots x_N^{m_N}$  such that  $f \cdot x_1^{m_1} \cdots x_N^{m_N} \in I$ . <sup>7</sup>The elimination ideal is obtained from I by intersecting with the smaller polynomial ring generated

<sup>&</sup>lt;sup>7</sup>The *elimination ideal* is obtained from I by intersecting with the smaller polynomial ring generated only by Plücker variables.

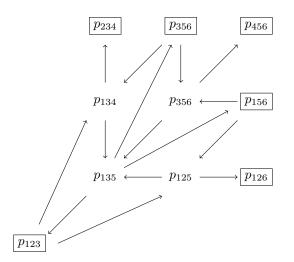


Figure 8. The Gr(3,6) quiver of the seed used in Example 5.2.

this seed is associated with one of the bipyramids. Hence, Corollary 5.1 should fail in this case. The (truncated) g-vectors of mutable Plücker coordinates with respect to s are

$g_s(p_{124}) = (0, 1, 0, 0),$	$g_s(p_{125}) = (0, 0, 0, -1),$	$g_s(p_{134}) = (-1, 0, 0, 0),$
$g_s(p_{135}) = (-1, -1, 1, -1),$	$g_s(p_{136}) = (-1, -1, 1, 0),$	$g_s(p_{145}) = (-1, 0, 1, -1),$
$g_s(p_{146}) = (-1, 0, 1, 0),$	$g_s(p_{235}) = (0, -1, 1, -1),$	$g_s(p_{236}) = (0, -1, 1, 0),$
$g_s(p_{245}) = (0, 0, 1, -1),$	$g_s(p_{246}) = (0, 0, 1, 0),$	$g_s(p_{256}) = (1, 0, 0, 0),$
$g_s(p_{346}) = (0, 0, 0, 1),$	$g_s(p_{356}) = (0, -1, 0, 0).$	

Notice that all g-vectors lie in the half space  $\{(a_1, a_2, a_3, a_4) \in \mathbb{Z}^4 : a_3 \geq 0\}$ . So, Corollary 5.1 does not apply.

#### 6 Massless scattering ideals

So far our discussion has been focused on extracting  $\mathcal{A}$ -coordinates and adjacency rules from the Gröbner fan of the Plücker ideal for the finite type Grassmannians  $\operatorname{Gr}(3,n)$ for n = 6,7,8. In fact, in [35], it was shown that the data of all cluster variables and their adjacencies in finite type cluster algebras of geometric type can be recovered from polyhedral data of a single maximal Gröbner cone of an appropriate ideal. However, what we are most interested in is how much physical information can be extracted from the Gröbner fan. For the cases of  $\operatorname{Gr}(4,6)$  and  $\operatorname{Gr}(4,7)$  the answer is the entire symbol alphabet and adjacency rules relevant for constructing the hexagon and heptagon amplitudes of planar  $\mathcal{N} = 4$  SYM. This motivates the question of whether the Gröbner fan provides a useful tool for the study of other kinematic ideals beyond the dual-conformal invariant case?

In this section we hope to provide a positive answer to this question by considering the example of general five-point massless scattering relevant for e.g. QCD processes. At two loops, all functions relevant for planar five-particle scattering were computed in [36], leading to the 26 letter alphabet  $\mathbb{A}_p$ . This was then extended to 31 letters relevant for the non-planar case  $\mathbb{A}_{np}$  in [37] where it was used to bootstrap individual two-loop Feynman integrals.

The goal of this section is to demonstrate how an analogous exploration of the Gröbner fan associated to a suitably defined five-point kinematic ideal  $I_{5\text{pt}}$  can generate (almost) the entire non-planar alphabet relevant for constructing (at least at two loops) five-point massless amplitudes. Note that we do in fact miss one symbol letter, known as  $W_{31}$  in [37]. However, the failure to recover  $W_{31}$  is consistent with the various calculations made for five-point processes, where it has been observed to be absent from (the suitably defined finite part of) the two-loop  $\mathcal{N} = 4$  SYM [39, 40] and  $\mathcal{N} = 8$  SUGRA [42, 43] amplitudes at two-loops. Similar two-loop observations have been made for  $q\bar{q} \to \gamma\gamma\gamma$  processes [58] and gluon amplitudes [59, 60] in QCD.

It is important to note that the computation we present here is only an analogy to the Grassmannian cases discussed in the previous sections for two reasons. Firstly, in the case of the Grassmannian we imposed positivity conditions by considering only nonprime factors appearing in the positive part of the tropical fan  $\operatorname{Trop}^+(I_{k,n})$ . However, for the five-point ideal, we do not impose any such positivity conditions and consider nonprime factors appearing in the full tropical fan  $\text{Trop}(I_{5\text{pt}})$ . In part this is because we are interested in non-planar theories, but also it is not totally clear, even in the planar case, which positivity conditions to impose. We suspect that there may well be multiple relevant positive regions. Secondly, in the case of the Grassmannian, having obtained the non-prime factors, we subsequently used them in order to extend the ideal, perhaps repeating this procedure multiple times as detailed for Gr(3, 8). This had the effect of eventually resolving the positive tropical fan of the fully extended ideal into a collection of simplices all with prime initial ideals. Furthermore, this singled out a single maximal Gröbner fan whose initial ideal contained the forbidden pairs of  $\mathcal{A}$ -coordinates, providing us with physical adjacency conditions on the symbol alphabet in the cases of Gr(4,6) and Gr(4,7). For the five-point case we perform no such extension however, again because we do not (yet) have a canonical notion of the relevant positivity conditions. Having obtained the non-prime factors appearing in the tropical fan we terminate the procedure since we already find the full non-planar alphabet. Note therefore that we do not extract any adjacency rules. That being said, it is encouraging that the same idea of symbol letters appearing as non-prime factors of the kinematic ideal follows through to the case of five-point massless scattering.

#### 6.1 The five-point two-loop symbol alphabet

The kinematics of five-point massless scattering is described on the five external momenta  $p_i^{\mu}$  subject to the massless on-shell condition  $p_i^2 = 0$  and momentum conservation  $\sum_i p_i^{\mu} = 0$ . Out of the momenta we can construct ten scalar products  $s_{ij} = 2p_i \cdot p_j$ , five of which are independent upon imposing momentum conservation. Following the choice of [38] they are given by<sup>8</sup>

$$v_i = s_{ii+1} = 2p_i \cdot p_{i+1}. \tag{6.1}$$

<sup>&</sup>lt;sup>8</sup>The five remaining non-adjacent scalar products may be written as  $s_{13} = s_{45} - s_{12} - s_{23}$  and its cyclic rotations.

Note that there are four independent dimensionless ratios that can be formed from these variables. It will also prove useful to introduce the following Gram determinant

$$\Delta = \det(2p_i \cdot p_j) = (\operatorname{tr}_5)^2, \tag{6.2}$$

where we have introduced the notation  $\text{tr}_5 = \text{tr}(\gamma_5 \not p_4 \not p_5 \not p_1 \not p_2)$ . Note, when written in terms of ' $\beta$ -variables' [61]  $\sqrt{\Delta} = \text{tr}_5$  can be expressed as a purely rational function.

The planar (two-loop) five point alphabet  $\mathbb{A}_p = \{W_1, \ldots, W_{20}\} \cup \{W_{26}, \ldots, W_{31}\}$  was originally obtained in [36] and consists of 26 letters given by

$$W_{i} = v_{i}, \qquad W_{5+i} = v_{i+2} + v_{i+3}, \qquad W_{10+i} = v_{i} - v_{i+3},$$
$$W_{15+i} = v_{i+3} - v_{i} - v_{i+1}, \qquad W_{25+i} = \frac{a_{i} - \sqrt{\Delta}}{a_{i} + \sqrt{\Delta}}, \qquad W_{31} = \sqrt{\Delta}, \qquad (6.3)$$

where the i indices run from 1 to 5 and we have introduced the notation

$$a_i = v_i v_{i+1} - v_{i+1} v_{i+2} + v_{i+2} v_{i+3} - v_i v_{i+4} - v_{i+3} v_{i+4}.$$

By closing the planar alphabet under permutations the authors of [37] generalised the alphabet to the non-planar case  $\mathbb{A}_{np} = \mathbb{A}_p \cup \{W_{21}, \dots, W_{26}\}$ , where we introduce the five additional non-planar letters given by

$$W_{20+i} = v_{i+2} + v_{i+3} - v_i - v_{i+1}.$$
(6.4)

As for the planar  $\mathcal{N} = 4$  SYM case the non planar alphabet  $\mathbb{A}_{np}$  provides the starting point for the construction of integrable polylogarithmic symbols relevant for the bootstrap of five-point massless non-planar amplitudes/integrals [37].

#### 6.2 Non-planar alphabet from a Gröbner fan

Inspired by the appearance of symbol letters ( $\mathcal{A}$ -coordinates) as non-prime factors of the Plücker ideal we wish to apply similar ideas to a suitably defined five-point ideal  $I_{5\text{pt}}$  in order to generate the non-planar alphabet  $\mathbb{A}_{np}$ .

To define the kinematic space for general n-point massless scattering, instead of using momentum twistor variables, it is instructive to consider spinor-helicity variables. We introduce these in the usual way by defining each null momentum as a bispinor,

$$p_i^{\alpha\dot{\alpha}} = (\sigma_\mu)^{\alpha\dot{\alpha}} p_i^\mu = \lambda_i^\alpha \tilde{\lambda}_i^{\dot{\alpha}} \,. \tag{6.5}$$

Then we consider the Lorentz invariant brackets<sup>9</sup>

$$\langle ij \rangle = \lambda_i^{\alpha} \lambda_{j\alpha}, \qquad [ij] = \tilde{\lambda}_i^{\dot{\alpha}} \tilde{\lambda}_{i\dot{\alpha}}.$$
 (6.6)

<sup>&</sup>lt;sup>9</sup>Note that we are interested in the generic case without dual conformal symmetry. One could use momentum twistors also in this case, allowing for the appearance of the infinity bitwistor, i.e. for both four-brackets  $\langle ijkl \rangle$  and two-brackets  $\langle ij \rangle = \langle ijI \rangle$ , however these variables are most natural in the case of planar amplitudes with ordered external legs. Here we will not impose any particular ordering.

The spinor brackets defined above are constrained by two sets of relations, the Plücker (or Schouten) identities,

$$\langle ij\rangle\langle kl\rangle - \langle ik\rangle\langle jl\rangle + \langle il\rangle\langle jk\rangle = 0; \quad [ij][kl] - [ik][jl] + [il][jk] = 0, \tag{6.7}$$

for  $1 \le i < j < k < l \le n$  and momentum conservation,

$$\sum_{j=1}^{n} \langle ij \rangle [jk] = 0, \tag{6.8}$$

for  $i, k \in \{1, ..., n\}$ . Note that the Plücker relations are three-term relations, while the momentum conservation relations can be three-term (for the off-diagonal case  $i \neq k$ ) or four-term (for the diagonal case i = k).

As we wish to recast the kinematic space in the language of polynomial ideals, we will introduce new variables  $a_{ij}$  and  $\tilde{a}_{ij}$  which obey the corresponding polynomial relations,

$$a_{ij}a_{kl} - a_{ik}a_{jl} + a_{il}a_{jk} = 0, \quad \tilde{a}_{ij}\tilde{a}_{kl} - \tilde{a}_{ik}\tilde{a}_{jl} + \tilde{a}_{il}\tilde{a}_{jk} = 0, \quad \sum_{j=1}^{n} a_{ij}\tilde{a}_{jk} = 0.$$
(6.9)

We take these relations to define an ideal  $I_{npt}$  on  $2 \times {n \choose 2}$  variables  $a_{ij}$  and  $\tilde{a}_{ij}$ .

We will study the ideal,  $I_{5pt}$ , defined on the ten  $a_{ij}$  and ten  $\tilde{a}_{ij}$  variables with i < jorganised as  $(a_{12}, \ldots, a_{45}, \tilde{a}_{12}, \ldots, \tilde{a}_{45})$  and subject to (6.9) for n = 5. In particular, we focus our attention on the Gröbner fan  $GF(I_{5pt})$  and tropical fan  $Trop(I_{5pt})$ . Note the ideal has a six-dimensional linear subspace (or lineality space), corresponding to the five little group rescalings of the spinor variables,

$$\lambda_i \to \alpha_i \lambda_i, \qquad \tilde{\lambda}_i \to \alpha_i^{-1} \tilde{\lambda}_i; \qquad a_{ij} \to \alpha_i \alpha_j a_{ij}, \qquad \tilde{a}_{ij} \to \alpha_i^{-1} \alpha_j^{-1} \tilde{a}_{ij}. \tag{6.10}$$

and the single overall dimension rescaling,

$$\lambda_i \to \beta \lambda_i , \qquad \tilde{\lambda}_i \to \beta \tilde{\lambda}_i ; \qquad a_{ij} \to \beta^2 a_{ij} , \qquad \tilde{a}_{ij} \to \beta^2 \tilde{a}_{ij} .$$
 (6.11)

In fact we can identify the ideal  $I_{5pt}$  with the Grassmannian Plücker ideal  $I_{3,6}$  via the following identification of variables,

$$a_{ij} = p_{ij6}, \qquad \tilde{a}_{ij} = (-1)^{j-i-1} p_{klm}, \qquad (6.12)$$

where in the second equality  $\{k, l, m\}$  is the ordered complement of  $\{i, j\}$  in the set  $\{1, 2, 3, 4, 5\}$ . Thus the  $a_{ij}$  correspond to all the Plücker variables involving the label 6, while the  $\tilde{a}_{ij}$  give all those which do not. The Plücker relations for the  $a_{ij}$  and the  $\tilde{a}_{ij}$  together with the off-diagonal three-term momentum conservation relations, give the three-term Plücker relations for Gr(3, 6) while the five diagonal four-term momentum conservation relations are the minimal set of four-term Plücker relations needed to complete a generating set of the Plücker ideal for Gr(3, 6).

Thus to study the tropical fan  $\text{Trop}(I_{5\text{pt}})$  we can simply refer to results already described for Gr(3, 6). In particular we know the fan has the f-vector (4.3) and, in particular,

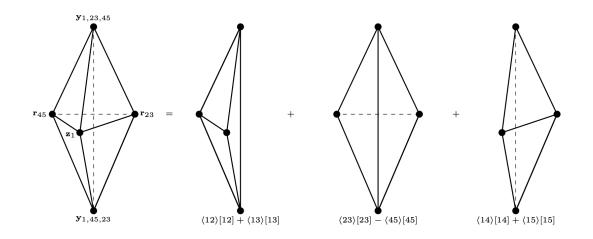


Figure 9. A bipyramid  $b_1$  of  $\text{Trop}(I_{5\text{pt}})$ , on the left hand side we have the full bipyramid labelled by its 5 rays, on the right hand side the bipyramid is split into three tetrahedra. Each tetrahedron is labelled by the quadratic non-prime factor found in the initial ideal, where all three expressions are equivalent modulo  $I_{5\text{pt}}$ .

has 65 rays. In the variables  $a_{ij}$  and  $\tilde{a}_{ij}$  used to define  $I_{5pt}$  these rays take the following form: 20 unit vectors  $\mathbf{e}_{ij}$  and  $\tilde{\mathbf{e}}_{ij}$  where  $\mathbf{e}_{ij}$  is the unit vector in the  $a_{ij}$  direction and respectively  $\tilde{\mathbf{e}}_{ij}$  is the unit vector in the  $\tilde{a}_{ij}$  direction; 5 vectors of the form  $\mathbf{z}_i = \sum_{j \neq i} \mathbf{e}_{ij}$ , lineality equivalent to  $\tilde{\mathbf{z}}_i = \sum_{j \neq i} \tilde{\mathbf{e}}_{ij}$ ; 10 permutations of  $\mathbf{r}_{45} = \mathbf{v}_{123} + \tilde{\mathbf{e}}_{45}$ , lineality equivalent to  $\tilde{\mathbf{r}}_{45} = \tilde{\mathbf{v}}_{123} + \mathbf{e}_{45}$ , where we have defined  $\mathbf{v}_{ijk} = \mathbf{e}_{ij} + \mathbf{e}_{ik} + \mathbf{e}_{jk}$  and similarly for  $\tilde{\mathbf{v}}_{ijk}$ ; finally the 30 permutations of  $\mathbf{y}_{1,23,45} = \mathbf{v}_{123} + \tilde{\mathbf{v}}_{145} = \tilde{\mathbf{y}}_{1,23,45}$ .

Amongst the 1035 maximal cones only 45 are non-prime: 30 given by the permutation copies of  $\{\mathbf{y}_{1,23,45}, \mathbf{r}_{23}, \mathbf{r}_{45}, \mathbf{y}_{1,45,23}\}$ ; and an additional 15 given by the permutation copies of  $\{\mathbf{y}_{1,23,45}, \mathbf{r}_{23}, \mathbf{r}_{45}, \mathbf{y}_{1,45,23}\}$ . As already described in our discussion of Gr(3,6), these tetrahedra fit together into bipyramids. In fact we have 15 such bipyramids across the full tropical fan. As an example consider figure 9 where we have a single bipyramid: the left tetrahedron with the rays  $\{\mathbf{y}_{1,23,45}, \mathbf{z}_1, \mathbf{r}_{45}, \mathbf{y}_{1,45,23}\}$ , produces the non-prime factor  $a_{12}\tilde{a}_{12} +$  $a_{13}\tilde{a}_{13} = \langle 12 \rangle [12] + \langle 13 \rangle [13]$ ; transposing 2  $\leftrightarrow$  4 and 3  $\leftrightarrow$  5 we find the right tetrahedron with the rays  $\{\mathbf{y}_{1,23,45}, \mathbf{z}_1, \mathbf{r}_{23}, \mathbf{y}_{1,45,23}\}$ , and non-prime factor  $a_{14}\tilde{a}_{14} + a_{15}\tilde{a}_{15} = \langle 14 \rangle [14] +$  $\langle 15 \rangle [15]$ ; finally the middle tetrahedron has the rays given by  $\{\mathbf{y}_{1,23,45}, \mathbf{r}_{23}, \mathbf{r}_{45}, \mathbf{y}_{1,45,23}\}$ , which produces the non-prime factor  $a_{23}\tilde{a}_{23} - a_{45}\tilde{a}_{45} = \langle 23 \rangle [23] - \langle 45 \rangle [45]$ . The three non-prime factors appearing in the bipyramid are equivalant modulo the ideal  $I_{5pt}$ .

To generate the full set of non-prime factors modulo the ideal  $I_{5pt}$  we need only take the permutation copies of  $\langle 23 \rangle [23] - \langle 45 \rangle [45]$  which produces 15 quadratic expressions given by

$$\begin{array}{ll} \langle 23\rangle [23] - \langle 45\rangle [45], & \langle 24\rangle [24] - \langle 35\rangle [35], & \langle 25\rangle [25] - \langle 34\rangle [34], \\ \langle 13\rangle [13] - \langle 45\rangle [45], & \langle 14\rangle [14] - \langle 35\rangle [35], & \langle 15\rangle [15] - \langle 34\rangle [34], \\ \langle 12\rangle [12] - \langle 45\rangle [45], & \langle 14\rangle [14] - \langle 25\rangle [25], & \langle 15\rangle [15] - \langle 24\rangle [24], \\ \langle 12\rangle [12] - \langle 35\rangle [35], & \langle 13\rangle [13] - \langle 25\rangle [25], & \langle 15\rangle [15] - \langle 23\rangle [23], \\ \langle 12\rangle [12] - \langle 34\rangle [34], & \langle 13\rangle [13] - \langle 24\rangle [24], & \langle 14\rangle [14] - \langle 23\rangle [23]. \end{array}$$

$$\begin{array}{l} (6.13) \end{array}$$

Along with the  $\{a_{ij} = \langle ij \rangle, \tilde{a}_{ij} = [ij]\}$  this provides us 35 expressions from which to form homogenous combinations.

To see that we are in fact recovering the same content as the symbol alphabet  $\mathbb{A}_{np}$  we re-write the entire non-planar alphabet in terms of spinor helicity variables given by the cyclic copies of<sup>10</sup>

$$W_{1} = \langle 12 \rangle [12], \qquad \qquad W_{6} = \langle 34 \rangle [34] + \langle 45 \rangle [45], W_{11} = \langle 34 \rangle [34] + \langle 35 \rangle [35], \qquad \qquad W_{16} = \langle 13 \rangle [13], W_{21} = \langle 13 \rangle [13] + \langle 34 \rangle [34], \qquad \qquad W_{26} = \frac{\langle 45 \rangle [51] \langle 12 \rangle [24]}{[45] \langle 51 \rangle [12] \langle 24 \rangle}, W_{31} = [45] \langle 51 \rangle [12] \langle 24 \rangle - \langle 45 \rangle [51] \langle 12 \rangle [24]. \qquad (6.14)$$

With this representation it is clear that letters  $\{W_i\}_{i=1}^5 \cup \{W_i\}_{i=16}^{20} \cup \{W_i\}_{i=26}^{30}$  are given by multiplicative combinations of the  $\{a_{ij}, \tilde{a}_{ij}\}$  variables. Furthermore, the remaining 15 letters  $\{W_i\}_{i=6}^{10} \cup \{W_i\}_{i=11}^{15} \cup \{W_i\}_{i=21}^{25}$ , themselves related by the  $S_5$  permutation symmetry, are exactly the 15 non-prime factors appearing in the Gröbner fan of the spinor-helicity ideal! To see this explicitly note we have

$$W_{6} = \langle 34 \rangle [34] + \langle 45 \rangle [45] = \langle 12 \rangle [12] - \langle 35 \rangle [35],$$
  

$$W_{11} = \langle 34 \rangle [34] + \langle 35 \rangle [35] = \langle 12 \rangle [12] - \langle 45 \rangle [45],$$
  

$$W_{21} = \langle 13 \rangle [13] + \langle 34 \rangle [34] = \langle 14 \rangle [14] - \langle 25 \rangle [25],$$
  
(6.15)

all of which appear in (6.13). It follows then that taking homogenous combinations of letters  $\{W_i\}_{i=1}^{30}$  is equivalent to taking homogenous combinations of  $\{\langle ij \rangle, [ij]\}$  and the permutations of the non-prime factors  $\langle 23 \rangle [23] - \langle 45 \rangle [45]$ . Note, as already emphasised, we do not recover the letter  $W_{31}$ . However, this is consistent with  $W_{31}$  not appearing in the expressions for suitably defined amplitudes.

## 7 Conclusion

We have presented a prescription for exploring the *alphabet* associated to the kinematic space of scattering amplitudes via the Gröbner fan. The construction was demonstrated for the cases Gr(3, n) with n = 6, 7, 8 to extract the  $\mathcal{A}$ -coordinates from a combination of the Speyer-Williams and Gröbner fans. Let us recap the various steps taken in the procedure. Our starting point is to identify the maximal cones of the positive tropical fan,  $\operatorname{Trop}^+(k, n)$ , whose associated initial ideals are non-prime. An efficient way of identifying such maximal cones is by comparison to a closely related object, the Speyer-Williams fan, which has the same support as  $\operatorname{Trop}^+(I_{k,n})$  but, in general, differing fan structure. In particular, the Speyer-Williams fan contains non-simplicial maximal cones which are further resolved by  $\operatorname{Trop}^+(I_{k,n})$ . These maximal cones are of particular interest, as their associated initial ideals are non-prime and allow us to identify certain *non-prime factors* which are exactly the non-plücker  $\mathcal{A}$ -coordinates of the cluster algebra. By searching for non-prime factors

<sup>&</sup>lt;sup>10</sup>Let us emphasise the fact that in spinor-helicity variables the rationality of letters  $\{W_{26}, \ldots, W_{31}\}$  becomes manifest.

in this way we are able to recover the full set of  $\mathcal{A}$ -coordinates. Having obtained the full set of  $\mathcal{A}$ -coordinates, following the results of [1], we are able to extract the forbidden pairs by computing the initial ideal of a Gröbner cone singled out by the positive tropical part of the extended ideal.<sup>11</sup> In the case of Gr(3,7) this allows us to extract the alphabet and adjacency rules relevant for constructing seven point amplitudes in planar  $\mathcal{N} = 4$  SYM.

An interesting example for future study is the case of Gr(4, 8) whose corresponding cluster algebra is no longer of finite type. The infinities of this cluster algebra arise due to alternate mutation on the nodes of affine  $A_2$  subalgebras. The series of mutations obtained in this way gives rise to a series of g-vectors which asymptote to certain special directions referred to as *limit rays*. It is to these limit rays the 18 multiplicatively independent square root letters appearing in the two-loop NMHV amplitude [30] are associated, see [22–24] for further details. It would be very interesting to see such infinities arising in the context of the Gröbner fan. We should expect that extending  $I_{4,8}$  by some *finite* set of non-prime factors is not enough to render all cones of the extended ideal prime.

We have also provided hints that the Gröbner fan may be a useful tool when considering more general scattering processes. We were able to find the entire non-planar alphabet relevant for five-point non dual conformal scattering processes by considering the Gröbner fan of a suitably defined kinematic ideal. In fact, this ideal turns out to be none other than the Grassmannian Plücker ideal for Gr(3, 6). Note, as emphasised in the main text no positivity criteria was imposed in this calculation. It would be interesting to study this case further to determine what is the correct notion of positivity for reducing to the planar alphabet for five-point scattering (clearly this will not be the same as reducing to the positive region in Gr(3, 6)) and whether this sheds light on any adjacency properties.

Finally, the six-point alphabet was recently studied at one-loop by [62], it would be interesting to apply the techniques developed here to see how much of the alphabet we can obtain.

# A Details of $\operatorname{Trop}^+(I_{3,7})$

Let us look at the 7-vertex cones in the Speyer-Williams fan for Gr(3,7). There are 63 cones with 7 vertices which come in five dihedral orbits each of length  $\{14, 14, 7, 14, 14\}$ . They have representatives with rays given by

$$\begin{pmatrix} b_{4,2345671} & b_{3,2345671} & b_{2,7123456} & b_{5,2345716} & b_{4,2371456} & b_{4,4567123} & b_{3,4567123} \\ b_{5,7123564} & b_{1,1234567} & b_{4,7123456} & b_{3,2345671} & b_{4,5671234} & b_{2,5671234} & b_{3,5671234} \\ b_{5,7123564} & b_{1,1234567} & b_{4,7123456} & b_{5,5671342} & b_{2,5671234} & b_{4,5634712} & b_{1,3456712} \\ b_{5,7123564} & b_{4,7123456} & b_{4,2345671} & b_{2,5671234} & b_{4,7156234} & b_{1,3456712} & b_{3,7123456} \\ b_{1,1234567} & b_{3,2345671} & b_{4,4567123} & b_{3,4567123} & b_{4,6745123} & b_{3,6712345} & b_{2,4567123} \end{pmatrix}$$

$$(A.1)$$

Starting from top to bottom the cones are refined as

$$\begin{pmatrix} b_{2,7123456} & b_{3,4567123} & b_{3,2345671} & b_{4,2345671} & b_{4,2371456} & b_{4,4567123} \\ b_{2,7123456} & b_{3,4567123} & b_{4,2345671} & b_{4,2371456} & b_{4,4567123} & b_{5,2345716} \end{pmatrix},$$
(A.2)

<sup>&</sup>lt;sup>11</sup>Note, having extended the Plücker ideal by the full set of  $\mathcal{A}$ -coordinates ensures that the Speyer-Williams fan and the positive tropical fan (of the extended ideal) are identical.

containing the missing binomials  $p_{145}p_{237} - p_{123}p_{457}$  and  $p_{137}p_{245} - p_{127}p_{345}$  respectively for the cone defined by the first row. We then have

$$\begin{pmatrix} b_{1,1234567} & b_{2,5671234} & b_{3,2345671} & b_{4,7123456} & b_{4,5671234} & b_{5,7123564} \\ b_{1,1234567} & b_{3,5671234} & b_{3,2345671} & b_{4,7123456} & b_{4,5671234} & b_{5,7123564} \\ b_{1,1234567} & b_{2,5671234} & b_{3,5671234} & b_{3,2345671} & b_{4,5671234} & b_{5,7123564} \end{pmatrix},$$
(A.3)

for the second row containing the missing binomials  $\{p_{167}p_{345}-p_{157}p_{346}, p_{167}p_{245}-p_{157}p_{246}\}, \{p_{147}p_{356}-p_{137}p_{456}, p_{147}p_{256}-p_{127}p_{456}\}$  and  $\{p_{156}p_{347}-p_{134}p_{567}, p_{156}p_{247}-p_{124}p_{567}\}$  respectively. Note that  $p_{167}p_{245}-p_{157}p_{246}$  is in a different cyclic class to the missing binomials found up to this point. Thus it is already clear at this stage that we will find all 14 quadratic  $\mathcal{A}$ -coordinates appearing as factors of non-prime initial ideals. For the third row we have

$$\begin{pmatrix} b_{1,3456712} & b_{1,1234567} & b_{2,5671234} & b_{4,7123456} & b_{5,7123564} & b_{5,5671342} \\ b_{1,3456712} & b_{1,1234567} & b_{4,7123456} & b_{4,5634712} & b_{5,7123564} & b_{5,5671342} \\ b_{1,3456712} & b_{1,1234567} & b_{2,5671234} & b_{4,5634712} & b_{5,7123564} & b_{5,5671342} \end{pmatrix},$$
(A.4)

with the missing binomials  $p_{167}p_{245} - p_{157}p_{246}$ ,  $p_{147}p_{256} - p_{127}p_{456}$  and  $p_{156}p_{247} - p_{124}p_{567}$ . For the fourth row we have

$$\begin{pmatrix} b_{1,3456712} & b_{2,5671234} & b_{4,2345671} & b_{4,7156234} & b_{4,7123456} & b_{5,7123564} \\ b_{1,3456712} & b_{2,5671234} & b_{3,7123456} & b_{4,2345671} & b_{4,7156234} & b_{4,7123456} \end{pmatrix},$$
(A.5)

containing the missing binomials  $p_{156}p_{237} - p_{123}p_{567}$  and  $p_{167}p_{235} - p_{157}p_{236}$ .

Before moving onto the final row of (A.1) let us observe here that none of the above refinements has utilised an extra positive ray of  $b_6$  type. The remaining case does however. For convenience we repeat the seven vertices here,

 $\{b_{1,1234567}, b_{4,4567123}, b_{4,6745123}, b_{3,2345671}, b_{3,4567123}, b_{3,6712345}, b_{2,4567123}\}.$  (A.6)

This 7-cone is refined as follows,

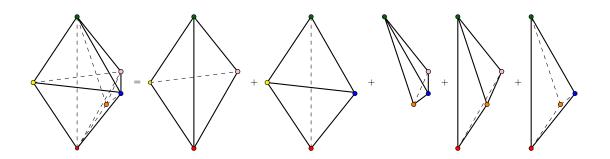
$$\begin{pmatrix} b_{1,1234567} & b_{4,4567123} & b_{4,6745123} & b_{2,4567123} & b_{3,2345671} & b_{3,4567123} \\ b_{1,1234567} & b_{4,4567123} & b_{4,6745123} & b_{2,4567123} & b_{3,2345671} & b_{3,6712345} \\ b_{1,1234567} & b_{4,4567123} & b_{4,6745123} & b_{6,4567123} & b_{3,2345671} & b_{3,4567123} \\ b_{1,1234567} & b_{4,4567123} & b_{4,6745123} & b_{6,4567123} & b_{3,2345671} & b_{3,6712345} \\ b_{1,1234567} & b_{4,4567123} & b_{4,6745123} & b_{6,4567123} & b_{3,2345671} & b_{3,6712345} \\ b_{1,1234567} & b_{4,4567123} & b_{4,6745123} & b_{6,4567123} & b_{3,6712345} & b_{3,4567123} \end{pmatrix},$$
(A.7)

with each part providing the following missing binomials

$$\begin{pmatrix} p_{237}p_{456} - p_{236}p_{457}, p_{137}p_{456} - p_{136}p_{457}, p_{127}p_{456} - p_{126}p_{457} \\ p_{235}p_{467} - p_{234}p_{567}, p_{135}p_{467} - p_{134}p_{567}, p_{125}p_{467} - p_{124}p_{567} \\ p_{237}p_{456} - p_{236}p_{457}, p_{167}p_{345} - p_{145}p_{367}, p_{167}p_{245} - p_{145}p_{267} \\ p_{235}p_{467} - p_{234}p_{567}, p_{167}p_{345} - p_{145}p_{367}, p_{167}p_{245} - p_{145}p_{267} \\ p_{267}p_{345} - p_{245}p_{367}, p_{167}p_{345} - p_{145}p_{367}, p_{167}p_{245} - p_{145}p_{267} \end{pmatrix}$$

$$(A.8)$$

We illustrate this refinement below in figure 10.



**Figure 10**. A 7-vertex Speyer-Williams cone is refined into 5 simplices in the tropical fan, three of which use the spurious orange vertex and two of which do not.

Next we consider the 8-vertex cones of the Speyer-Williams fan. There are 28 of these coming in two dihedral classes. The first has a representative given by

 $\{b_{4,7123456}, b_{4,2345671}, b_{3,2345671}, b_{2,7123456}, b_{5,2345716}, b_{4,2371456}, b_{3,4567123}, b_{3,7123456}\}, \quad (A.9)$ 

which has the refinement

$$\begin{pmatrix} b_{2,7123456} & b_{3,4567123} & b_{3,2345671} & b_{4,2371456} & b_{4,7123456} & b_{4,2345671} \\ b_{2,7123456} & b_{3,4567123} & b_{4,2345671} & b_{4,2371456} & b_{4,7123456} & b_{5,2345716} \\ b_{3,4567123} & b_{3,7123456} & b_{4,2345671} & b_{4,2371456} & b_{4,7123456} & b_{5,2345716} \\ b_{2,7123456} & b_{3,4567123} & b_{3,7123456} & b_{4,2371456} & b_{4,7123456} & b_{5,2345716} \\ b_{3,4567123} & b_{3,2345671} & b_{4,2345671} & b_{4,2371456} & b_{4,7123456} & b_{6,2345671} \\ b_{3,4567123} & b_{3,7123456} & b_{4,2345671} & b_{4,2371456} & b_{4,7123456} & b_{6,2345671} \\ b_{3,4567123} & b_{3,7123456} & b_{4,2345671} & b_{4,2371456} & b_{4,7123456} & b_{6,2345671} \\ b_{3,4567123} & b_{3,7123456} & b_{4,2345671} & b_{4,2371456} & b_{4,7123456} & b_{6,2345671} \\ b_{3,4567123} & b_{3,7123456} & b_{4,2345671} & b_{4,2371456} & b_{4,7123456} & b_{6,2345671} \\ b_{3,4567123} & b_{3,7123456} & b_{4,2345671} & b_{4,2371456} & b_{4,7123456} & b_{6,2345671} \\ b_{3,4567123} & b_{3,7123456} & b_{4,2345671} & b_{4,2371456} & b_{4,7123456} & b_{6,2345671} \\ b_{3,4567123} & b_{3,7123456} & b_{4,2345671} & b_{4,2371456} & b_{4,7123456} & b_{6,2345671} \\ b_{3,4567123} & b_{3,7123456} & b_{4,2345671} & b_{4,2371456} & b_{4,7123456} & b_{6,2345671} \\ b_{3,4567123} & b_{3,7123456} & b_{4,2345671} & b_{4,2371456} & b_{4,7123456} & b_{6,2345671} \\ b_{3,4567123} & b_{3,7123456} & b_{4,2345671} & b_{4,2371456} & b_{4,7123456} & b_{6,2345671} \\ b_{3,4567123} & b_{3,7123456} & b_{3,2345671} & b_{4,2371456} & b_{4,7123456} & b_{6,2345671} \\ b_{3,4567123} & b_{3,7123456} & b_{3,2345671} & b_{3,2345671} & b_{4,2371456} & b_{4,7123456} & b_{6,2345671} \\ b_{3,4567123} & b_{3,7123456} & b_{3,2345671} & b_{4,2371456} & b_{4,7123456} & b_{6,2345671} \\ b_{3,4567123} & b_{3,7123456} & b_{3,2345671} & b_{3,2345671} & b_{4,2371456} & b_{4,7123456} & b_{4,2345671} \\ b_{3,4567123} & b_{3,7123456} & b_{3,2345671} & b_{3,2345671} & b_{4,2371456} & b_{4,7123456} & b_{4,7123456} & b_{3,7123456} & b_{3,7156} & b_{3,7123456} & b_{3,7156} & b_{3,7156} & b_{$$

The missing binomials associated to each piece of the refinement are

$$\begin{pmatrix} p_{156}p_{237} - p_{123}p_{567}, p_{146}p_{237} - p_{123}p_{467}, p_{145}p_{237} - p_{123}p_{457} \\ p_{137}p_{245} - p_{127}p_{345}, p_{156}p_{237} - p_{123}p_{567}, p_{146}p_{237} - p_{123}p_{467} \\ p_{137}p_{245} - p_{127}p_{345}, p_{167}p_{235} - p_{157}p_{236}, p_{167}p_{234} - p_{147}p_{236} \\ p_{137}p_{256} - p_{127}p_{356}, p_{137}p_{246} - p_{127}p_{346}, p_{137}p_{245} - p_{127}p_{345} \\ p_{145}p_{237} - p_{123}p_{457}, p_{167}p_{235} - p_{157}p_{236}, p_{167}p_{234} - p_{147}p_{236} \\ p_{167}p_{235} - p_{157}p_{236}, p_{167}p_{234} - p_{147}p_{236}, p_{167}p_{234} - p_{147}p_{235} \\ p_{137}p_{245} - p_{127}p_{345}, p_{167}p_{235} - p_{157}p_{236}, p_{167}p_{234} - p_{147}p_{236} \end{pmatrix}$$
(A.11)

The second class has a representative with rays

 $\{b_{5,7123564}, b_{1,1234567}, b_{4,7123456}, b_{5,5671342}, b_{4,5671234}, b_{2,5671234}, b_{4,5634712}, b_{3,5671234}\}, (A.12)$ 

which has the refinement

 $\begin{pmatrix} b_{1,1234567} & b_{2,5671234} & b_{4,7123456} & b_{4,5671234} & b_{5,7123564} & b_{5,5671342} \\ b_{1,1234567} & b_{4,7123456} & b_{4,5634712} & b_{4,5671234} & b_{5,7123564} & b_{5,5671342} \\ b_{1,1234567} & b_{3,5671234} & b_{4,7123456} & b_{4,5634712} & b_{4,5671234} & b_{5,7123564} \\ b_{1,1234567} & b_{2,5671234} & b_{4,5634712} & b_{4,5671234} & b_{5,7123564} & b_{5,5671342} \\ b_{1,1234567} & b_{2,5671234} & b_{4,5634712} & b_{4,5671234} & b_{5,7123564} & b_{5,7123564} \end{pmatrix}$ (A.13)

and the missing binomials

$$\begin{pmatrix} p_{167}p_{345} - p_{157}p_{346}, p_{167}p_{245} - p_{157}p_{246} \\ p_{167}p_{345} - p_{157}p_{346}, p_{147}p_{256} - p_{127}p_{456} \\ p_{147}p_{356} - p_{137}p_{456}, p_{147}p_{256} - p_{127}p_{456} \\ p_{167}p_{345} - p_{157}p_{346}, p_{156}p_{247} - p_{124}p_{567} \\ p_{156}p_{347} - p_{134}p_{567}, p_{156}p_{247} - p_{124}p_{567} \end{pmatrix}$$
(A.14)

Finally, we arrive at the 9-vertex cones of which there are 7 in a single dihedral class. A representative has the rays

 $\{b_{4,2345671}, b_{3,2345671}, b_{5,2345716}, b_{4,2371456}, b_{4,4567123}, b_{3,4567123}, b_{5,4567231}, b_{4,4523671}, b_{2,2345671}\}$  (A.15)

and a refinement given by

 $\begin{pmatrix} b_{3,4567123} & b_{3,2345671} & b_{4,2345671} & b_{4,2371456} & b_{4,4567123} & b_{4,4523671} \\ b_{3,2345671} & b_{4,2345671} & b_{4,2371456} & b_{4,4567123} & b_{4,4523671} & b_{5,4567231} \\ b_{2,2345671} & b_{3,2345671} & b_{4,2371456} & b_{4,4567123} & b_{4,4523671} & b_{5,4567231} \\ b_{3,4567123} & b_{4,2345671} & b_{4,2371456} & b_{4,4567123} & b_{4,4523671} & b_{5,2345716} \\ b_{4,2345671} & b_{4,2371456} & b_{4,4567123} & b_{4,4523671} & b_{5,2345716} & b_{5,4567231} \\ b_{2,2345671} & b_{4,2345671} & b_{4,2371456} & b_{4,4523671} & b_{5,2345716} & b_{5,4567231} \\ b_{2,2345671} & b_{3,4567123} & b_{4,2371456} & b_{4,4567123} & b_{4,4523671} & b_{5,2345716} \\ b_{2,2345671} & b_{3,4567123} & b_{4,2345671} & b_{4,4567123} & b_{4,4523671} & b_{5,2345716} \\ b_{2,2345671} & b_{4,2345671} & b_{4,2345671} & b_{4,4567123} & b_{4,4523671} & b_{5,2345716} \\ b_{2,2345671} & b_{4,2345671} & b_{4,4567123} & b_{4,4523671} & b_{5,2345716} & b_{5,4567231} \\ b_{2,2345671} & b_{4,2345671} & b_{4,4567123} & b_{4,4523671} & b_{5,2345716} & b_{5,4567231} \\ b_{2,2345671} & b_{4,2345671} & b_{4,4567123} & b_{4,4523671} & b_{5,2345716} & b_{5,4567231} \\ b_{2,2345671} & b_{4,2345671} & b_{4,4567123} & b_{4,4523671} & b_{5,2345716} & b_{5,4567231} \\ b_{2,2345671} & b_{4,2345671} & b_{4,4567123} & b_{4,4523671} & b_{5,2345716} & b_{5,4567231} \\ b_{2,2345671} & b_{4,2345671} & b_{4,4567123} & b_{4,4523671} & b_{5,2345716} & b_{5,4567231} \\ b_{2,2345671} & b_{4,2345671} & b_{4,4567123} & b_{4,4523671} & b_{5,2345716} & b_{5,4567231} \\ b_{2,2345671} & b_{4,2345671} & b_{4,4567123} & b_{4,4523671} & b_{5,2345716} & b_{5,4567231} \\ b_{2,2345671} & b_{4,2345671} & b_{4,4567123} & b_{4,4523671} & b_{5,4567231} \\ b_{2,2345671} & b_{4,2345671} & b_{4,4567123} & b_{4,4523671} & b_{5,4567231} \\ b_{2,2345671} & b_{4,2345671} & b_{4,4567123} & b_{4,4523671} & b_{5,4567231} \\ b_{2,2345671} & b_{4,2345671} & b_{4,4567123} & b_{4,4523671} & b_{5,4567231} \\ b_{2,2345671} & b_{4,2345671} & b_{4,4567123} & b_{4,4523671} & b_{5,4567231} \\ b_{2,23456$ 

with the missing binomials

 $\begin{pmatrix} p_{237}p_{456} - p_{236}p_{457}, p_{145}p_{237} - p_{123}p_{457}, p_{145}p_{236} - p_{123}p_{456} \\ p_{235}p_{467} - p_{234}p_{567}, p_{145}p_{237} - p_{123}p_{457}, p_{145}p_{236} - p_{123}p_{456} \\ p_{235}p_{467} - p_{234}p_{567}, p_{157}p_{234} - p_{147}p_{235}, p_{156}p_{234} - p_{146}p_{235} \\ p_{235}p_{467} - p_{234}p_{567}, p_{137}p_{245} - p_{127}p_{345}, p_{145}p_{236} - p_{123}p_{456} \\ p_{235}p_{467} - p_{234}p_{567}, p_{137}p_{245} - p_{127}p_{345}, p_{145}p_{236} - p_{123}p_{456} \\ p_{235}p_{467} - p_{234}p_{567}, p_{137}p_{245} - p_{127}p_{345}, p_{156}p_{234} - p_{146}p_{235} \\ p_{267}p_{345} - p_{224}p_{567}, p_{137}p_{245} - p_{127}p_{345}, p_{156}p_{234} - p_{146}p_{235} \\ p_{235}p_{467} - p_{234}p_{567}, p_{137}p_{245} - p_{127}p_{345}, p_{136}p_{245} - p_{126}p_{345} \\ p_{235}p_{467} - p_{234}p_{567}, p_{137}p_{245} - p_{127}p_{345}, p_{136}p_{245} - p_{126}p_{345} \\ p_{235}p_{467} - p_{234}p_{567}, p_{137}p_{245} - p_{127}p_{345}, p_{136}p_{245} - p_{126}p_{345} \\ p_{235}p_{467} - p_{234}p_{567}, p_{137}p_{245} - p_{127}p_{345}, p_{136}p_{245} - p_{126}p_{345} \\ p_{235}p_{467} - p_{234}p_{567}, p_{137}p_{245} - p_{127}p_{345}, p_{136}p_{245} - p_{126}p_{345} \\ p_{235}p_{467} - p_{234}p_{567}, p_{137}p_{245} - p_{127}p_{345}, p_{136}p_{245} - p_{126}p_{345} \\ \end{pmatrix}$ 

#### Acknowledgments

We are grateful to the Isaac Newton Institute and the organizers of the hybrid workshop "Interdisciplinary applications of cluster algebras" held in 2021 where the discussions that lead to this article started. Furthermore L.B. acknowledges support of the PAPIIT project IA100122, Dirección General de Asuntos del Personal Académico, Universidad Nacional Autónoma de México 2022.

**Open Access.** This article is distributed under the terms of the Creative Commons Attribution License (CC-BY 4.0), which permits any use, distribution and reproduction in any medium, provided the original author(s) and source are credited.

#### References

- L. Bossinger et al., Families of Gröbner degenerations, Grassmannians and universal cluster algebras, SIGMA 17 (2021) 059 [arXiv:2007.14972].
- [2] J. Golden, A.B. Goncharov, M. Spradlin, C. Vergu and A. Volovich, Motivic Amplitudes and Cluster Coordinates, JHEP 01 (2014) 091 [arXiv:1305.1617] [INSPIRE].
- [3] J. Golden, M.F. Paulos, M. Spradlin and A. Volovich, Cluster Polylogarithms for Scattering Amplitudes, J. Phys. A 47 (2014) 474005 [arXiv:1401.6446] [INSPIRE].
- [4] A.B. Goncharov, M. Spradlin, C. Vergu and A. Volovich, Classical Polylogarithms for Amplitudes and Wilson Loops, Phys. Rev. Lett. 105 (2010) 151605 [arXiv:1006.5703]
   [INSPIRE].
- [5] S. Fomin and A. Zelevinsky, Cluster algebras i: Foundations, J. Am. Math. Soc. 15 (2001) 497 [math/0104151].
- [6] S. Fomin and A. Zelevinsky, Cluster algebras II. Finite type classification, Invent. Math. 154 (2003) 63 [math/0208229].
- S. Fomin and A. Zelevinsky, Cluster algebras IV. Coefficients, Compos. Math. 143 (2007) 112 [math/0602259].
- [8] L.J. Dixon, J.M. Drummond and J.M. Henn, Bootstrapping the three-loop hexagon, JHEP 11 (2011) 023 [arXiv:1108.4461] [INSPIRE].
- [9] L.J. Dixon, J.M. Drummond, M. von Hippel and J. Pennington, *Hexagon functions and the three-loop remainder function*, *JHEP* **12** (2013) 049 [arXiv:1308.2276] [INSPIRE].
- [10] L.J. Dixon, J.M. Drummond, C. Duhr and J. Pennington, The four-loop remainder function and multi-Regge behavior at NNLLA in planar N = 4 super-Yang-Mills theory, JHEP 06 (2014) 116 [arXiv:1402.3300] [INSPIRE].
- [11] L.J. Dixon and M. von Hippel, Bootstrapping an NMHV amplitude through three loops, JHEP 10 (2014) 065 [arXiv:1408.1505] [INSPIRE].
- [12] L.J. Dixon, M. von Hippel and A.J. McLeod, The four-loop six-gluon NMHV ratio function, JHEP 01 (2016) 053 [arXiv:1509.08127] [INSPIRE].
- S. Caron-Huot, L.J. Dixon, A. McLeod and M. von Hippel, Bootstrapping a Five-Loop Amplitude Using Steinmann Relations, Phys. Rev. Lett. 117 (2016) 241601
   [arXiv:1609.00669] [INSPIRE].
- S. Caron-Huot, L.J. Dixon, F. Dulat, M. von Hippel, A.J. McLeod and G. Papathanasiou, Six-Gluon amplitudes in planar N = 4 super-Yang-Mills theory at six and seven loops, JHEP 08 (2019) 016 [arXiv:1903.10890] [INSPIRE].
- [15] J.M. Drummond, G. Papathanasiou and M. Spradlin, A Symbol of Uniqueness: The Cluster Bootstrap for the 3-Loop MHV Heptagon, JHEP 03 (2015) 072 [arXiv:1412.3763] [INSPIRE].
- [16] L.J. Dixon, J. Drummond, T. Harrington, A.J. McLeod, G. Papathanasiou and M. Spradlin, Heptagons from the Steinmann Cluster Bootstrap, JHEP 02 (2017) 137 [arXiv:1612.08976] [INSPIRE].
- [17] J. Drummond, J. Foster, O. Gürdoğan and G. Papathanasiou, Cluster adjacency and the four-loop NMHV heptagon, JHEP 03 (2019) 087 [arXiv:1812.04640] [INSPIRE].

- [18] J. Drummond, J. Foster and O. Gürdoğan, Cluster Adjacency Properties of Scattering Amplitudes in N = 4 Supersymmetric Yang-Mills Theory, Phys. Rev. Lett. 120 (2018) 161601 [arXiv:1710.10953] [INSPIRE].
- [19] J. Drummond, J. Foster and O. Gürdoğan, Cluster adjacency beyond MHV, JHEP 03 (2019) 086 [arXiv:1810.08149] [INSPIRE].
- [20] S. Caron-Huot, L.J. Dixon, F. Dulat, M. Von Hippel, A.J. McLeod and G. Papathanasiou, The Cosmic Galois Group and Extended Steinmann Relations for Planar  $\mathcal{N} = 4$  SYM Amplitudes, JHEP **09** (2019) 061 [arXiv:1906.07116] [INSPIRE].
- [21] J. Drummond, J. Foster, O. Gürdogan and C. Kalousios, Tropical Grassmannians, cluster algebras and scattering amplitudes, JHEP 04 (2020) 146 [arXiv:1907.01053] [INSPIRE].
- [22] J. Drummond, J. Foster, O. Gürdogan and C. Kalousios, Algebraic singularities of scattering amplitudes from tropical geometry, JHEP 04 (2021) 002 [arXiv:1912.08217] [INSPIRE].
- [23] N. Arkani-Hamed, T. Lam and M. Spradlin, Non-perturbative geometries for planar  $\mathcal{N} = 4$ SYM amplitudes, JHEP 03 (2021) 065 [arXiv:1912.08222] [INSPIRE].
- [24] N. Henke and G. Papathanasiou, How tropical are seven- and eight-particle amplitudes?, JHEP 08 (2020) 005 [arXiv:1912.08254] [INSPIRE].
- [25] A. Herderschee, Algebraic branch points at all loop orders from positive kinematics and wall crossing, JHEP 07 (2021) 049 [arXiv:2102.03611] [INSPIRE].
- [26] N. Henke and G. Papathanasiou, Singularities of eight- and nine-particle amplitudes from cluster algebras and tropical geometry, JHEP 10 (2021) 007 [arXiv:2106.01392] [INSPIRE].
- [27] Q. Yang, Schubert problems, positivity and symbol letters, JHEP 08 (2022) 168
   [arXiv:2203.16112] [INSPIRE].
- [28] S. He, J. Liu, Y. Tang and Q. Yang, The symbology of Feynman integrals from twistor geometries, arXiv:2207.13482 [INSPIRE].
- [29] S. He, Z. Li and C. Zhang, A nice two-loop next-to-next-to-MHV amplitude in  $\mathcal{N} = 4$  super-Yang-Mills, JHEP 12 (2022) 158 [arXiv:2209.10856] [INSPIRE].
- [30] S. He, Z. Li and C. Zhang, Two-loop octagons, algebraic letters and Q equations, Phys. Rev. D 101 (2020) 061701 [arXiv:1911.01290] [INSPIRE].
- [31] D. Maclagan and B. Sturmfels, Introduction to Tropical Geometry, Graduate Studies in Mathematics. Vol. 161, American Mathematical Society, Providence, U.S.A. (2015), https://doi.org/10.1090/gsm/161.
- [32] D. Speyer and B. Sturmfels, The tropical Grassmannian, Adv. Geom. 4 (2004) 389 [math/0304218].
- [33] D. Speyer and L. Williams, The tropical totally positive grassmannian, J. Algebr. Comb. 22 (2005) 189 [math/0312297].
- [34] S.B. Brodsky, C. Ceballos and J.-P. Labbé, Cluster algebras of type D<sub>4</sub>, tropical planes, and the positive tropical Grassmannian, Beitr. Algebra Geom. 58 (2017) 25 [arXiv:1511.02699].
- [35] N. Ilten, A. Nájera Chávez and H. Treffinger, Deformation Theory for Finite Cluster Complexes, arXiv:2111.02566.
- [36] T. Gehrmann, J.M. Henn and N.A. Lo Presti, Analytic form of the two-loop planar five-gluon all-plus-helicity amplitude in QCD, Phys. Rev. Lett. 116 (2016) 062001 [arXiv:1511.05409] [Erratum ibid. 116 (2016) 189903] [INSPIRE].

- [37] D. Chicherin, J. Henn and V. Mitev, *Bootstrapping pentagon functions*, *JHEP* **05** (2018) 164 [arXiv:1712.09610] [INSPIRE].
- [38] T. Gehrmann, J.M. Henn and N.A. Lo Presti, Pentagon functions for massless planar scattering amplitudes, JHEP 10 (2018) 103 [arXiv:1807.09812] [INSPIRE].
- [39] S. Abreu, L.J. Dixon, E. Herrmann, B. Page and M. Zeng, The two-loop five-point amplitude in  $\mathcal{N} = 4$  super-Yang-Mills theory, Phys. Rev. Lett. **122** (2019) 121603 [arXiv:1812.08941] [INSPIRE].
- [40] D. Chicherin, T. Gehrmann, J.M. Henn, P. Wasser, Y. Zhang and S. Zoia, Analytic result for a two-loop five-particle amplitude, Phys. Rev. Lett. 122 (2019) 121602 [arXiv:1812.11057]
   [INSPIRE].
- [41] S. Badger et al., Analytic form of the full two-loop five-gluon all-plus helicity amplitude, Phys. Rev. Lett. 123 (2019) 071601 [arXiv:1905.03733] [INSPIRE].
- [42] D. Chicherin, T. Gehrmann, J.M. Henn, P. Wasser, Y. Zhang and S. Zoia, *The two-loop five-particle amplitude in N = 8 supergravity*, *JHEP* 03 (2019) 115 [arXiv:1901.05932] [INSPIRE].
- [43] S. Abreu, L.J. Dixon, E. Herrmann, B. Page and M. Zeng, The two-loop five-point amplitude in  $\mathcal{N} = 8$  supergravity, JHEP 03 (2019) 123 [arXiv:1901.08563] [INSPIRE].
- [44] J. Herzog and T. Hibi, Monomial Ideals, Graduate Texts in Mathematics. Vol. 260, Springer-Verlag London, London, U.K. (2011), https://doi.org/10.1007/978-0-85729-106-6.
- [45] T. Mora and L. Robbiano, The Gröbner fan of an ideal, J. Symb. Comput. 6 (1988) 183.
- [46] M. Einsiedler and S. Tuncel, When does a polynomial ideal contain a positive polynomial?, J. Pure Appl. Algebra 164 (2001) 149.
- [47] J. Drummond, J. Foster, O. Gürdoğan and C. Kalousios, Tropical fans, scattering equations and amplitudes, JHEP 11 (2021) 071 [arXiv:2002.04624] [INSPIRE].
- [48] A.N. Jensen, *Gfan, a software system for Gröbner fans and tropical varieties*, available at http://home.imf.au.dk/jensen/software/gfan/gfan.html.
- [49] S. Herrmann, A. Jensen, M. Joswig and B. Sturmfels, How to Draw Tropical Planes, Electron. J. Combin. 16 (2009) R6 [arXiv:0808.2383].
- [50] A.N. Jensen, *Planes in five dimensional tropical projective space*, https://users-math.au.dk/jensen/Research/G3\_7/grassmann3\_6.html.
- [51] F. Cachazo, N. Early, A. Guevara and S. Mizera, Scattering Equations: From Projective Spaces to Tropical Grassmannians, JHEP 06 (2019) 039 [arXiv:1903.08904] [INSPIRE].
- [52] A.N. Jensen, *Planes in six dimensional tropical projective space*, https://users-math.au.dk/jensen/Research/G3\_7/grassmann3\_7.html.
- [53] F. Cachazo and J.M. Rojas, Notes on Biadjoint Amplitudes, Trop G(3,7) and X(3,7)Scattering Equations, JHEP 04 (2020) 176 [arXiv:1906.05979] [INSPIRE].
- [54] K. Kaveh and C. Manon, Khovanskii bases, higher rank valuations, and tropical geometry, SIAM J. Appl. Algebra Geom. 3 (2019) 292.
- [55] L. Bossinger, Full-rank valuations and toric initial ideals, Int. Math. Res. Not. 2021 (2020) 7715.

- [56] L. Bossinger, Tropical totally positive cluster varieties, arXiv:2208.01723.
- [57] S. Fomin, L. Williams and A. Zelevinsky, Introduction to Cluster Algebras. Chapter 6, arXiv:2008.09189.
- [58] S. Abreu, B. Page, E. Pascual and V. Sotnikov, Leading-Color Two-Loop QCD Corrections for Three-Photon Production at Hadron Colliders, JHEP 01 (2021) 078 [arXiv:2010.15834] [INSPIRE].
- [59] S. Badger, C. Brønnum-Hansen, H.B. Hartanto and T. Peraro, Analytic helicity amplitudes for two-loop five-gluon scattering: the single-minus case, JHEP 01 (2019) 186 [arXiv:1811.11699] [INSPIRE].
- [60] S. Abreu, J. Dormans, F. Febres Cordero, H. Ita, B. Page and V. Sotnikov, Analytic Form of the Planar Two-Loop Five-Parton Scattering Amplitudes in QCD, JHEP 05 (2019) 084 [arXiv:1904.00945] [INSPIRE].
- [61] Z. Bern, L.J. Dixon and D.A. Kosower, One loop corrections to five gluon amplitudes, Phys. Rev. Lett. 70 (1993) 2677 [hep-ph/9302280] [INSPIRE].
- [62] J.M. Henn, A. Matijašić and J. Miczajka, One-loop hexagon integral to higher orders in the dimensional regulator, JHEP 01 (2023) 096 [arXiv:2210.13505] [INSPIRE].