

Tools for CICYs in F-theory

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ABSTRACT: We provide a set of tools for analyzing the geometry of elliptically fibered Calabi-Yau manifolds, starting with a description of the total space rather than with a Weierstrass model or a specified type of fiber/base. Such an approach to the subject of F-theory compactification makes certain geometric properties, which are usually hidden, manifest. Specifically, we review how to isolate genus-one fibrations in such geometries and then describe how to find their sections explicitly. This includes a full parameterization of the Mordell-Weil group where non-trivial. We then describe how to analyze the associated Weierstrass models, Jacobians and resolved geometries. We illustrate our discussion with concrete examples which are complete intersections in products of projective spaces (CICYs). The examples presented include cases exhibiting non-abelian symmetries and higher rank Mordell-Weil group. We also make some comments on non-flat fibrations in this context. In a companion paper [1] to this one, these results will be used to analyze the consequences for string dualities of the ubiquity of multiple fibrations in known constructions of Calabi-Yau manifolds.

KEYWORDS: F-Theory, String Duality

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1 Introduction

In many approaches to compactifications of F-theory, the identification of the fiber and base of the internal manifold is built in from the start. Typically, one begins with a choice of base manifold, and then fibers an elliptic curve, described in terms of an appropriate complete intersection in some toric variety, over that space. Systematic scans over Calabi-Yau (CY) geometries constructed in this manner can be found in refs. [2–12]. Such a construction has many advantages, including the fact that one is guaranteed that the associated manifold is genus-one fibered and therefore is indeed suitable for use in F-theory. Nevertheless, this methodology has the drawback that, instead of simply using the large data sets of Calabi-Yau manifolds that have already been constructed (for example in refs. [13–23]), one is essentially starting all over again in reconstructing those manifolds with the desired internal structure. In addition, as we will discuss, certain properties of the resulting compactifications can be hard to see in such descriptions.

In this paper we present tools for systematically pursuing a different approach to F-theory compactification. We describe how to take any smooth Calabi-Yau manifold and extract the F-theory physics associated to this “resolved space” directly (for some related work see ref. [22]). In particular we describe how to do the following:

1. Isolate genus-one fibrations in one of the conventional data sets of Calabi-Yau manifolds. Here we build on technology first developed in refs. [2, 19, 22, 24].
2. Determine whether or not each of these fibrations have a section. In cases where a section is present we describe how to obtain an explicit form for it in terms of the original description of the manifold. This methodology is closely related to descriptions of holomorphic functions used in recent constructions of “generalized complete intersection CY manifolds” (gCICYs) [25].
3. Obtain an explicit Weierstrass model associated to blowing down all components of the fibers in the original manifold that have generic vanishing intersection number with a chosen zero section. Here we follow the construction of refs. [28, 29].

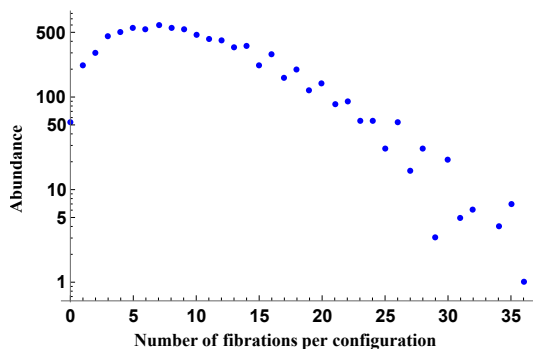


Figure 1. The abundance of CICY threefold configuration matrices, in the standard list, exhibiting a given number of genus-one fibrations which are visible directly in the configuration matrix [32].

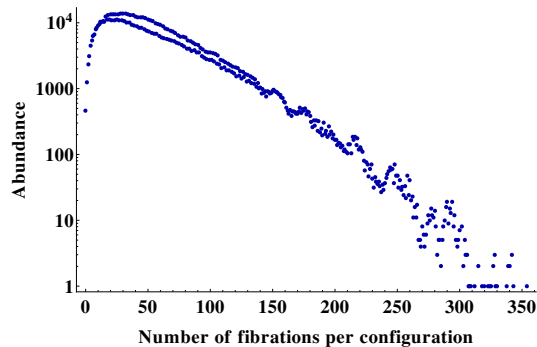


Figure 2. The abundance of CICY fourfold configuration matrices, in the standard list, exhibiting a given number of genus-one fibrations which are visible directly in the configuration matrix [19].

4. Obtain the Jacobian manifold associated to the original compactification, making use of techniques from [2, 30, 31].

Once this data has been obtained, standard techniques can be employed to study the F-theoretical physics of the compactifications in question. We illustrate all of this with concrete examples taken from the data set of complete intersection Calabi-Yau (CICY) manifolds [13–19]. We expect that similar techniques should allow the methodology to be generalized to other data sets such as that due to Kreuzer and Skarke [20–23], or the gCICYs [25–27], in a straightforward manner. Note that some steps in this direction have already been taken with Jacobian forms associated to complete intersection fibers ref. [30].

In addition to utilizing pre-existing data sets of Calabi-Yau manifolds, this approach to analyzing global F-theory compactifications makes evident some features of fibrations that are not as obvious in more standard methodologies. One of the most important of these features, which will be explored extensively in the context of dualities in a companion paper [1] to this one, is that of multiple fibrations in a single CY geometry. It is known [8, 12, 19, 22] that the vast majority of known Calabi-Yau manifolds are genus-one fibered. It is also suspected that essentially all such manifolds can be written in a myriad of different ways as such a fibration — indeed this has been proven in the case of complete intersections in products of projective spaces [19, 32]. This abundance of possible rewritings of the CICY geometries is illustrated in figures 1 and 2. If one simply constructs a Calabi-Yau manifold as a particular fiber type over a given base (e.g. in Weierstrass form), the existence of other descriptions of the manifold as fibrations over different bases can be difficult to see. In the approach we are following here, such collections of descriptions are manifest, and thus the obviously closely related F-theory models that they correspond to can be studied as a set [1, 22].

As a simple illustration of multiple genus one fibrations, consider the following Calabi-Yau threefold, described as a complete intersection of two polynomial equations in $\mathbb{P}^1 \times$

$\mathbb{P}^2 \times \mathbb{P}^1 \times \mathbb{P}^1$.

$$X = \left[\begin{array}{c|cc} \mathbb{P}^1 & 1 & 1 \\ \mathbb{P}^2 & 1 & 2 \\ \mathbb{P}^1 & 1 & 1 \\ \mathbb{P}^1 & 1 & 1 \end{array} \right], \tag{1.1}$$

Here the columns describe the two defining relations of the Calabi-Yau manifold by giving their polynomial degree in the homogeneous coordinates of the ambient projective spaces. This construction is described in detail in the next section. The manifold X can be written in two different ways as a genus-one fibration. Below the fiber is described by the configuration matrix above the dotted line:

$$\left[\begin{array}{c|cc} \mathbb{P}^1 & 1 & 1 \\ \mathbb{P}^2 & 1 & 2 \\ \hline \mathbb{P}^1 & 1 & 1 \\ \mathbb{P}^1 & 1 & 1 \end{array} \right], \quad \left[\begin{array}{c|cc} \mathbb{P}^1 & 1 & 1 \\ \mathbb{P}^1 & 1 & 1 \\ \hline \mathbb{P}^1 & 1 & 1 \\ \mathbb{P}^2 & 1 & 2 \end{array} \right]. \tag{1.2}$$

The base of the fibration is then simply $\mathbb{P}^1 \times \mathbb{P}^1$ in the first case and \mathbb{P}^2 in the second.

For the first of these fibrations we will provide, in section 6, an explicit description of the associated Weierstrass model (over $\mathbb{P}^1 \times \mathbb{P}^1$) as one simple example of our method. The second of these fibrations is genus one but does not have a section. We provide several examples in sections 5–8 with some exhibiting multiple elliptic fibrations giving rise to the same total space.

A second feature that we have observed in applying this approach to the CICY data set, is that a great deal of these manifolds exhibit at least one fibration with a relatively high rank Mordell-Weil group. Indeed, even for the very simple example given above, the first fibration discussed admits two sections. As will be detailed in section 6, these can be described as the global holomorphic sections of the line bundles $\mathcal{O}_X(-1, 1, 1, 1)$ and $\mathcal{O}_X(2, -1, 4, 4)$, respectively. Here we are using the standard notation where the integers represent the coefficients in an expansion of the first Chern class of the line bundle in a basis provided by restricting the ambient Kähler forms to the Calabi-Yau manifold. As a somewhat larger example, in section 8 we provide a case with Mordell-Weil rank 4.

As a final comment, it is clear that certain features — such as the exact nature of the resolved geometry corresponding to an F-theory model, and thus its M-theory limit — are obvious in this construction. In the more conventional approach to building F-theory models this information can be highly non-trivial to obtain and a lot of interesting work has been carried out in this regard (see refs. [37–44] for some recent advances).

The outline of this paper is as follows. In section 2, we describe how to identify genus-one fibrations of CICYs and obtain explicit expressions for their sections, if present. In section 3, we describe how to use this information to obtain the Weierstrass models and the Jacobians associated to the initial Calabi-Yau manifolds. We also give some more information on the resolved geometries and discuss the decomposition of the Picard lattice according to the theorem of Shioda-Tate-Wazir [45]. Sections 5–8 contain explicit examples of these techniques as applied to cases taken from the CICY data set.

2 Elliptic fibrations and sections

Consider a Calabi-Yau n -fold X_n embedded as a complete intersection of K hypersurfaces in a product of projective spaces $\mathcal{A} = \prod_{r=1}^m \mathbb{P}_{\mathbf{x}_r}^{n_r}$, where the subscripts $\mathbf{x}_r = (x_{r,0} : \dots : x_{r,n_r})$ denote the homogeneous coordinates of the corresponding projective space \mathbb{P}^{n_r} . Then, it can be characterized by a so-called configuration matrix,

$$X_n = \left[\begin{array}{c|ccc} \mathbb{P}_{\mathbf{x}_1}^{n_1} & a_1^1 & \dots & a_K^1 \\ \mathbb{P}_{\mathbf{x}_2}^{n_2} & a_1^2 & \dots & a_K^2 \\ \vdots & \vdots & \ddots & \vdots \\ \mathbb{P}_{\mathbf{x}_m}^{n_m} & a_1^m & \dots & a_K^m \end{array} \right]. \quad (2.1)$$

Here the matrix entry a_j^r denotes the degree of the j -th hypersurface equation in the homogeneous coordinates of the r -th projective space factor of the ambient space. The Calabi-Yau condition leads to the degree constraints,

$$n_r + 1 = \sum_{j=1}^K a_j^r, \quad (2.2)$$

for each $r = 1, \dots, m$, while the condition that the Calabi-Yau be an n -fold is given as

$$\sum_{r=1}^m n_r = n + K. \quad (2.3)$$

In this paper, we will be analyzing such Calabi-Yau manifolds realized as a CICY.¹ Further, since we wish to study F-theory vacua, we will restrict our considerations to those CICYs with at least one “obvious” genus-one fibration, as we will describe in the next subsection. To avoid potential confusions, before we proceed, however, we clarify the terminology that we will use throughout the paper. “Genus-one fibration” refers to a fibration of genus-one curves, whether or not it has a section, while “elliptic fibration” implies the existence of a section. Note also that we will oftentimes hide the manifold subscripts that indicate dimensions, for example X_n and X will be used interchangeably unless confusions arise.

2.1 Genus-one fibration structures in CICYs

In general it is rather difficult to take a CICY configuration matrix and enumerate all of the ways in which the associated variety can be written as a genus-one fibration. However, there exists a class of genus-one fibrations which can be readily classified from the configuration matrix alone.

It is possible to perform arbitrary row and column permutations on a configuration matrix without changing the geometry that is described. These operations simply correspond to relabelling the \mathbb{P}^{n_r} ambient factors and the hypersurface equations, respectively. Let us

¹In many computations we have made use of the “CICY Package” [46], based on methods described in [47–51].

suppose that upon appropriate use of such permutations, the configuration matrix (2.1) can be put in the following block form,

$$X_n = \left[\begin{array}{c|cc} \mathcal{A}_F & O & \mathcal{F} \\ \hline \mathcal{A}_B & \mathcal{B} & \mathcal{T} \end{array} \right], \tag{2.4}$$

where \mathcal{A}_F and \mathcal{A}_B are two products of projective spaces such that

$$\mathcal{A}_F \times \mathcal{A}_B = \prod_{r=1}^m \mathbb{P}_{\mathbf{x}_r}^{n_r} \equiv \mathcal{A}, \tag{2.5}$$

while \mathcal{F} , \mathcal{B} , \mathcal{T} , and $O = 0$ are submatrices of the full configuration matrix. Such a configuration describes a fibration of the fiber $F = [\mathcal{A}_F || \mathcal{F}]$ over the base $B = [\mathcal{A}_B || \mathcal{B}]$, where \mathcal{T} describes the variation of the fiber over the base. Since the top-left block is a zero matrix, the Calabi-Yau condition (2.2) guarantees that the fiber F obeys the analogous degree constraints and hence is Calabi-Yau itself too. Therefore, as long as the number of columns of \mathcal{F} and the dimension of \mathcal{A}_F are such that F is of complex dimension 1, the fibers will be Calabi-Yau one-folds: that is genus-one curves as desired. It follows that the base will then be of complex dimension $n - 1$.

Such genus-one fibration structures can easily be found at the level of configuration matrix via permutations of rows and columns, and will hence be referred to as “obvious genus-one fibrations” (OGFs).² Almost all CICY configurations have an OGF, and in many cases multiple such structures [19, 32]. For example, among the 7,890 CICY threefolds, 7,837 can be brought into the OGF form (2.4) and a CICY threefold admits 9.85 OGFs on average, with the number of OGFs per configuration ranging from 0 to 36. Similarly, among the 921,497 CICY fourfolds, all but 477 have an OGF and a CICY fourfold admits 54.6 OGFs on average, with the number of OGFs per configuration ranging from 0 to 354.

As an illustration of OGF, let us consider the $K3$ surface with the following configuration matrix,

$$X_2 = \left[\begin{array}{c|cc} \mathbb{P}_{\mathbf{x}_1}^1 & 1 & 1 \\ \hline \mathbb{P}_{\mathbf{x}_2}^2 & 1 & 2 \\ \mathbb{P}_{\mathbf{x}_3}^1 & 1 & 1 \end{array} \right]. \tag{2.6}$$

The $K3$ surface X_2 admits an OGF structure over the base $B_1 = \mathbb{P}_{\mathbf{x}_3}^1$, where the configuration of the fiber is given as

$$F = \left[\begin{array}{c|cc} \mathbb{P}_{\mathbf{x}_1}^1 & 1 & 1 \\ \hline \mathbb{P}_{\mathbf{x}_2}^2 & 1 & 2 \end{array} \right]. \tag{2.7}$$

In this particular example the matrix \mathcal{B} has 0 columns. For the rest of this section and the next two, we will use this configuration as a simple and explicit example with which to clarify various techniques. The entire analysis for this $K3$ surface is put together in

²Note that in ref. [19] such OGFs were referred to as obvious elliptic fibrations (OEFs). In this paper, however, we need to carefully distinguish between fibrations with and without sections and thus will avoid this earlier nomenclature.

section 5 in a self-contained manner for the reader who prefers a complete worked example to an illustrated general analysis.

One further point that should be made before we proceed is that the OGFs just described need not be flat fibrations. As a simple example consider the following configuration,

$$X_3 = \left[\begin{array}{c|ccc} \mathbb{P}_{\mathbf{x}_1}^1 & 1 & 1 & 0 \\ \mathbb{P}_{\mathbf{x}_2}^1 & 1 & 0 & 1 \\ \mathbb{P}_{\mathbf{x}_3}^2 & 1 & 0 & 2 \\ \mathbb{P}_{\mathbf{x}_4}^2 & 1 & 1 & 1 \end{array} \right]. \quad (2.8)$$

Here we can take the fiber to be described by

$$F = \left[\begin{array}{c|ccc} \mathbb{P}_{\mathbf{x}_1}^1 & 1 & 1 & 0 \\ \mathbb{P}_{\mathbf{x}_2}^1 & 1 & 0 & 1 \\ \mathbb{P}_{\mathbf{x}_3}^2 & 1 & 0 & 2 \end{array} \right], \quad (2.9)$$

with $\mathbb{P}_{\mathbf{x}_4}^2$ being the base. Consider the defining relation given by the second column in eq. (2.8). This takes the following form,

$$P_2 = x_{1,0}l_0(\mathbf{x}_4) + x_{1,1}l_1(\mathbf{x}_4), \quad (2.10)$$

where l_0 and l_1 are linear functions in the homogeneous coordinates of the base $\mathbb{P}_{\mathbf{x}_4}^2$. For general choices of complex structure, at a certain point in the base we have $l_0 = l_1 = 0$. At such point $P_2 = 0$ holds automatically, and this equation does not provide a constraint in $\mathcal{A}_F = \mathbb{P}_{\mathbf{x}_1}^1 \times \mathbb{P}_{\mathbf{x}_2}^1 \times \mathbb{P}_{\mathbf{x}_3}^2$. Thus over that point in the base the fiber is actually two dimensional, not a curve, and thus the fibration is not flat.

More generally, if there is any choice of point on the base, such that the associated choice of complex structure describing the fiber eq. (2.9) is not a complete intersection, then the fibration will not be flat. For most of this paper, we will restrict ourselves to examining flat fibrations for ease and physical motivations. However, in appendix A, we will apply our methods to a non-flat fibration and make some comments about the connection between non-flatness of fibrations and singularities in the associated Weierstrass model in which f and g vanish to orders 4 and 6.

Finally, it should be noted that although the statistics described above focused on OGFs in CICYs, an elliptic fibration structure can also be found via a set of criteria purely in terms of the intersection theory:

Conjecture [52]: Let X be a Calabi-Yau n -fold. Then X is genus-one fibered iff there exists a $(1, 1)$ -class D in $H^2(X, \mathbb{Q})$ such that $(D \cdot C) \geq 0$ for every algebraic curve $C \subset X$, $(D^{\dim(X)}) = 0$ and $D^{\dim(X)-1} \neq 0$.

For $n = 3$ (i.e. a CY 3-fold) this conjecture has been proven subject to the additional constraints that D is effective or $(D \cdot c_2(X)) \neq 0$ [53, 54]. It is straightforward to see in the case of many CICYs, including for example the 3-fold given in (1.1), the criteria

above can be used to verify that the OGF fibrations are in fact *all fibrations* for the given CICY (see [55, 56] for general formulas on intersection numbers and topology of CICYs). In summary, the CICY dataset provide a rich data set of examples of multiply genus-one fibered Calabi-Yau manifolds which we will exploit in ref. [1] with the techniques developed in this paper.

2.2 Putative sections to genus-one fibrations

Given a smooth, genus-one fibered Calabi-Yau n -fold, $X_n \rightarrow B_{n-1}$, the fibration may or may not have a section. Although many of the techniques in this paper apply also to the cases without a section, for a clearer interpretation of the corresponding F-theory vacuum, we will always start our geometric exploration by determining whether or not a section exists (see, e.g., refs. [57–70] for examples of recent work on the physics of genus-one fibrations without a section). We will separate the procedure of finding a section into two steps. In the first step we will find a candidate divisor class in which a section could live by imposing topological constraints. In the second step the section itself will be constructed explicitly as a map from the base to the fiber.

We call a codimension-one subvariety $S \subset X_n$ a “putative section” if it is an element of a divisor class $[S]$ that meets the following two criteria, both of which necessarily hold for a section:

- (a) Let \hat{D}_α^b for $\alpha = 1, \dots, h^{1,1}(B)$ be a basis of divisor classes in the base and D_α^b their pull backs to X . Then, due to Oguiso [53], the following equality between two intersection products, one on X_n and the other on B_{n-1} , should hold,

$$[S] \cdot \prod_{k=1}^{n-1} D_{\alpha_k}^b = \prod_{k=1}^{n-1} \hat{D}_{\alpha_k}^b, \tag{2.11}$$

for each $(n - 1)$ -tuple $(\alpha_1, \dots, \alpha_{n-1})$ with $\alpha_k = 1, \dots, h^{1,1}(B)$. This is necessary if the fiber F_p at a generic point $p \in B$ is to intersect S at a single point. We will oftentimes refer to the criterion (a) as the Oguiso criterion.

- (b) For S to be a section it should be birational to the base B , and thus necessarily obeys the following intersection criterion [72],

$$[S] \cdot [S] \cdot \prod_{k=1}^{n-2} D_{\alpha_k}^b = -[c_1(B)] \cdot [S] \cdot \prod_{k=1}^{n-2} D_{\alpha_k}^b, \tag{2.12}$$

for each $(n - 2)$ -tuple $(\alpha_1, \dots, \alpha_{n-2})$ with $\alpha_k = 1, \dots, h^{1,1}(B)$. Note that this criterion has been well-established for $n = 2$ and 3, and is believed to also hold for CY n -folds with $n > 3$.

When applied to any particular case, the two criteria above lead to a family of possible divisor classes for putative sections. Given that any putative section S is an element of a divisor class $[S]$, there exists a corresponding line bundle $\mathcal{O}_X(S)$. In what follows

we will frequently find it useful to consider these line bundles rather than the associated divisor classes.

As an illustration, let us consider the $K3$ configuration of eq. (2.6). If the first Chern class of $\mathcal{O}_X(S)$ is $c_1(\mathcal{O}_X(S)) = b_1 J_1 + b_2 J_2 + b_3 J_3$, where $J_{r=1,2,3}$ are the harmonic $(1, 1)$ -forms descending from the ambient projective pieces $\mathbb{P}_{\mathbf{x}_r}^{n_r}$, then we will denote the line bundle as $\mathcal{O}_X(S) = \mathcal{O}_X(b_1, b_2, b_3)$. Here we have labeled the basis of forms such that the base $B = \mathbb{P}_{\mathbf{x}_3}^1$ has a unique harmonic $(1, 1)$ -form generator J_3 . The right hand side of criterion (a) is then simple to compute:

$$\prod_{k=1}^1 \hat{D}_1^b = \int_{\mathbb{P}_{\mathbf{x}_3}^1} J_3 = 1. \quad (2.13)$$

Here we have computed the intersection number of the divisor class in the base by integrating the dual $(1, 1)$ form, which we take to be normalized in the usual manner, over that manifold. The left hand side of eq. (2.11) can be computed in terms of a similar integral over wedge products of dual forms — this time over the Calabi-Yau manifold itself. Remembering that our unique divisor in the base is dual to J_3 and how we are parameterizing the first Chern class of $\mathcal{O}_X(S)$ we find

$$[S] \cdot \prod_{k=1}^1 D_1^b = \int_X (b_1 J_1 + b_2 J_2 + b_3 J_3) \wedge J_3. \quad (2.14)$$

We will denote the form dual to the Calabi-Yau manifold itself, inside the ambient space, as μ_X . In general this form can be read directly off from the configuration matrix (2.1) as $\mu_X = \bigwedge_j a_j^r J_r$. Using this, we find that the integral in eq. (2.14) can now be rewritten as follows,

$$\begin{aligned} [S] \cdot \prod_{k=1}^1 D_1^b &= \int_{\mathbb{P}_{\mathbf{x}_1}^1 \times \mathbb{P}_{\mathbf{x}_2}^2 \times \mathbb{P}_{\mathbf{x}_3}^1} [(b_1 J_1 + b_2 J_2 + b_3 J_3) \wedge J_3] \wedge \mu_X \\ &= \int_{\mathbb{P}_{\mathbf{x}_1}^1 \times \mathbb{P}_{\mathbf{x}_2}^2 \times \mathbb{P}_{\mathbf{x}_3}^1} [(b_1 J_1 + b_2 J_2 + b_3 J_3) \wedge J_3] \wedge [(J_1 + J_2 + J_3) \wedge (J_1 + 2J_2 + J_3)] \\ &= 2b_1 + 3b_2. \end{aligned} \quad (2.15)$$

Combining eqs. (2.13) and (2.15) we finally find that condition (2.11) results in the following constraint on the divisor class of any potential section,

$$2b_1 + 3b_2 \stackrel{!}{=} 1. \quad (2.16)$$

The criterion (b) can be written in terms of b_r via a very similar intersection computation and leads to

$$6b_1 b_2 + 2b_2^2 + 4b_1 b_3 + 6b_2 b_3 \stackrel{!}{=} -2. \quad (2.17)$$

These two putative-section criteria, (2.16) and (2.17), can be solved as

$$b_1 = -1 - 3k, \quad b_2 = 1 + 2k, \quad b_3 = 1 + 11k + 14k^2, \quad (2.18)$$

with an integer parameter $k \in \mathbb{Z}$. For some small values of $k = \{1, 0, -1, -2\}$, for instance, we obtain the putative section classes, $\mathcal{O}_X(-4, 3, 26)$, $\mathcal{O}_X(-1, 1, 1)$, $\mathcal{O}_X(2, -1, 4)$, and $\mathcal{O}_X(5, -3, 35)$, respectively.

2.3 Sections to elliptic fibrations

To make further progress towards finding a section, beyond the necessary topological conditions imposed in the previous subsection, we must fix a complex structure for the Calabi-Yau manifold X_n . This is simply achieved by choosing an explicit set of defining equations $P_j(\mathbf{x}_r) = 0$ where $j = 1, \dots, K$, in terms of the homogeneous coordinates, $\mathbf{x}_r = (x_{r,0} : \dots : x_{r,n_r})$, of the ambient space, $\mathcal{A} = \prod_{r=1}^m \mathbb{P}_{\mathbf{x}_r}^{n_r}$.

Given a parameterization of putative section classes, such as eq. (2.18), and a choice of complex structure, the next step is to select a divisor S_0 satisfying criteria (a) and (b) which also has $h^0(X, \mathcal{O}_X(S_0)) = 1$.³ The unique global holomorphic section (GHS),

$$z = z(\mathbf{x}_1, \dots, \mathbf{x}_m) \in \Gamma(X, \mathcal{O}_X(S_0)), \tag{2.19}$$

will then be explicitly found and proven to yield a true section to the fibration, whether holomorphic or rational.⁴

The condition $h^0(X, \mathcal{O}_X(S_0)) = 1$ follows, in the case where the section describes a smooth element of an effective divisor class, from the birationality of the desired section to the base. For an elliptically fibered Calabi-Yau manifold the cohomology of the trivial bundle over the base is always zero, except for the presence of a single global section. This is a simple consequence of the Calabi-Yau condition and the fact that pulling back any further cohomology elements under the projection map would lead to harmonic forms on the total space that are known not to exist [73]. Since these cohomologies are a birational invariant, and the section is birational to the base, we require that these statements should also hold for S_0 . For a smooth section we can then use the Koszul sequence to relate trivial bundle cohomology on S_0 to that of $\mathcal{O}_X(S_0)$ as follows. The short exact Koszul sequence describing the section inside the Calabi-Yau gives rise to

$$0 \rightarrow \mathcal{O}_X(-S_0) \rightarrow \mathcal{O}_X \rightarrow \mathcal{O}_X|_{S_0} \rightarrow 0. \tag{2.20}$$

The associated long exact sequence in cohomology then contains the following piece.

$$0 \rightarrow H^3(X, \mathcal{O}_X(-S_0)) \rightarrow H^3(X, \mathcal{O}_X) \rightarrow 0 \tag{2.21}$$

Using Serre duality on the Calabi-Yau manifold, together with $h^3(X, \mathcal{O}_X) = 1$, we then find that $h^0(X, \mathcal{O}_X(S_0)) \stackrel{!}{=} 1$ as claimed.⁵

³It should be noted that the criterion (a) together with the line-bundle cohomology condition, $h^0(X, \mathcal{O}_X(S)) = 1$, are not sufficient conditions for the existence of a section. In particular, it happens for some divisors that these two conditions are satisfied while the putative-section criterion (b) in eq. (2.12) is not.

⁴Note that for the rest of the paper we will use the acronym GHS for global holomorphic section of a line bundle and reserve the word “section” for the (putative) section to a genus-one fibration.

⁵The condition $h^0(X, \mathcal{O}_X(S_0)) = 1$ will in fact also hold in the case of singular sections. In the case of smooth sections, the higher cohomology groups also vanish: $h^j(X, \mathcal{O}_X(S_0)) = 0 \forall j > 0$. These vanishings are not a necessity for us however, since rational sections to an elliptic fibration may be singular, in which case $\mathcal{O}_X(S_0)$ may have a non-trivial higher cohomology (see eq. (6.13) for an example). We do not add the condition $h^0(X, \mathcal{O}_X(S_0)) = 1$ to our definition of a putative section, since the cohomology condition does not lead to a closed-form constraint unlike the criteria (a) and (b) and thus it is practically employed in a different manner.

It should be noted that putative-section criteria (a) and (b), even when equipped with the cohomology condition, $h^0(X, \mathcal{O}_X(S_0)) = 1$, do not guarantee that the zero locus of eq. (2.19) is a section to the genus-one fibration. Therefore, in order to ensure that the putative section S_0 leads to a true section, it is desired to provide an explicit rational map from B to S_0 , specified by parametrization of the form,

$$\mathbf{x}_F = \mathbf{x}_F(\mathbf{x}_B), \tag{2.22}$$

where \mathbf{x}_F and \mathbf{x}_B collectively denote the ambient homogeneous coordinates of \mathcal{A}_F and \mathcal{A}_B , respectively. Once an explicit expression for the GHS is found in terms of the ambient coordinates \mathbf{x}_r , as implied by eq. (2.19), one can indeed solve for the fiber coordinates in terms of the base coordinates (up to some subtleties which we will discuss further in the Example sections). However, the Oguiso criterion (a), together with the fact that the intersection numbers in the projective-space product \mathcal{A}_F are non-negative, tend to force the line bundle $\mathcal{O}_X(S_0)$ to simultaneously have positive and negative degrees along the \mathcal{A}_F direction. For example, in the $K3$ case above, the resulting constraint was eq. (2.16), which indeed forces one of b_1 and b_2 to be negative and the other positive. Thus, the GHS of $\mathcal{O}_X(S_0)$ cannot be written as a polynomial in \mathbf{x}_r . Instead, it takes a rational form and can be constructed using the techniques developed in the context of generalized CICYs [25].

Let us briefly review how such a rational expression for the GHS of $\mathcal{O}_X(S_0)$ can be obtained, in the presence of both positive and negative degrees. We first label the line bundle $\mathcal{O}_X(S_0)$ in terms of its first Chern class as $\mathcal{O}_X(b_1, \dots, b_m)$ as in section 2.2. We then take the following rational ansatz for its GHS,

$$\frac{N(\mathbf{x}_1, \dots, \mathbf{x}_m)}{D(\mathbf{x}_1, \dots, \mathbf{x}_m)} \in \Gamma(X, \mathcal{O}_X(b_1, \dots, b_m)), \tag{2.23}$$

where N and D are polynomials in $\mathbf{x}_1, \dots, \mathbf{x}_m$ of multi-degree $(b_1, \dots, b_m)_+$ and $(b_1, \dots, b_m)_-$, respectively. Here, the signs in subscript indicate that only the degrees with the specified sign are extracted (and are flipped to be positive for the ‘-’ case). For example, for $\mathcal{O}_X(-1, 1, 1)$, which is one of the putative sections we found for our simple $K3$ case (2.6)–(2.7), the polynomials N and D are of multi-degrees $(-1, 1, 1)_+ = (0, 1, 1)$ and $(-1, 1, 1)_- = (1, 0, 0)$, respectively. For a given denominator D with the right degree, one is not allowed to choose a generic numerator N . This is because the ratio N/D would then behave irregularly at a generic point where D vanishes, while the GHS of the line bundle $\mathcal{O}(-1, 1, 1)$ is known to be associated to a polynomial in coordinates on X itself which can not exhibit any such singularities. Instead N must be chosen such that on the complete intersection X the vanishing of D is completely cancelled by that of N , yielding a regular function. Such a tuning of coefficients in N provides an explicit GHS in terms of \mathbf{x}_r .

As an illustration, let us go back to the $K3$ example and take the putative section $\mathcal{O}_X(-1, 1, 1)$ from eq. (2.18) with $k = 0$, whose cohomology is computed as $h^\bullet(X, \mathcal{O}_X(-1, 1, 1)) = (1, 0, 0)$. To find the (unique) rational expression of the form (2.23), where $D(\mathbf{x}_1)$ is linear in \mathbf{x}_1 and $N(\mathbf{x}_2, \mathbf{x}_3)$ is bi-linear in \mathbf{x}_2 and \mathbf{x}_3 , we need to make use of the first defining equation of $K3$, which we denote by $P_1(\mathbf{x}_1, \mathbf{x}_2, \mathbf{x}_3)$. This defining equation

can be expanded as

$$P_1 = x_{1,0} p_1(\mathbf{x}_2, \mathbf{x}_3) + x_{1,1} p_2(\mathbf{x}_2, \mathbf{x}_3), \tag{2.24}$$

for bi-linear polynomials p_1 and p_2 (note from the first column of eq. (2.6) that P_1 is tri-linear). Without loss of generality, we may choose to use $D(\mathbf{x}_1) = x_{1,0} + 10x_{1,1}$, for example, and proceed to determine the bi-linear polynomial $N(\mathbf{x}_2, \mathbf{x}_3)$, that vanishes whenever D does in X . Substituting the solution $x_{1,0} = -10x_{1,1}$ to $D(\mathbf{x}_1) = 0$ into P_1 we immediately have a perfect candidate for a numerator, $-10p_1 + p_2$ which vanishes on X whenever D does. This particular choice of N can be thought of as a tuning of the coefficients of the six bi-linear monomials in $\mathbb{P}_{\mathbf{x}_2}^2 \times \mathbb{P}_{\mathbf{x}_3}^1$. Thus, the GHS in this case is constructed (uniquely as a function on X) as

$$\frac{N(\mathbf{x}_2, \mathbf{x}_3)}{D(\mathbf{x}_1)} = \frac{-10p_1(\mathbf{x}_2, \mathbf{x}_3) + p_2(\mathbf{x}_2, \mathbf{x}_3)}{x_{1,0} + 10x_{1,1}}. \tag{2.25}$$

In particular, this expression can be proven equivalent when evaluated on X to any other expression obtained from a different choice of the denominator polynomial D [25]. In some cases, the GHS can have sufficiently complex dependence on the defining equations of the CY that it can be difficult to determine the section analytically along the lines above. In this case it is still possible to determine the appropriate regular rational function numerically. This can be done by locating a large number of points on the denominator, D — found by intersecting the CY defining equations, together with the denominator, with an appropriate number of generic multi-linear hypersurface constraints. By requiring that the numerator also vanish (along with the denominator) for this collection of points, the coefficients of a generic numerator can be fully fixed, leading to a complete description of the GHS [25].

It is worth mentioning that there is a related method for the numerator tuning, which shares the same spirit as the previous method, and which is applicable specifically when the base X of the line bundle in question is a fibration itself (as in our case). Again, one starts from the ansatz (2.23), together with a choice of D with the right degree. At a generic point $p \in B$, $D = 0$ can be solved for a discrete set of points on the fiber F_p over p . Then, the numerator polynomial N , when evaluated at each of these fibral points, should vanish. By choosing sufficiently many points on the base B one can then obtain enough constraints on the coefficients in N to uniquely determine the numerator. Despite being essentially the same as the first method described above, this approach benefits from the fact that one does not need to spot the correct combination of defining relations that must be used in order to derive the appropriate numerator. This method therefore lends itself better to automation on a computer. This alternative method will be used in section 5 for the same $K3$ geometry and will be shown to give the same global section expression as in eq. (2.25).

Using any of the above methods, one can obtain the GHS expression and hence, also the explicit section map of the form (2.22). If the section map can explicitly be shown to be rational this way, we will have found a legitimate section. It should be noted, however, that the corresponding divisor is not necessarily smooth. For example, putative sections may be reducible, containing a genuine section as an irreducible component as well as a

vertical divisor therein (see appendix A for explicit examples); these cases can (and will) be ruled out by testing the irreducibility of the divisor. In general, even a genuine section can be a singular divisor and some CY geometries presented in this paper will include such instances.

As has been mentioned above, the sections that are found may be rational and thus wrap fibral \mathbb{P}^1 's over certain points of the base. For the case of smooth sections, therefore, where the Euler number $\chi(S_0)$ of S_0 is well-defined, this number should be closely related to that of the base $\chi(B)$ via the manner in which those \mathbb{P}^1 's are wrapped. In particular, for $X_2 \rightarrow B_1$, the sections are always holomorphic, while for $X_3 \rightarrow B_2$, the sections may wrap \mathbb{P}^1 's over a finite number of base points. For the latter case, the Euler number difference, $\chi(S_0) - \chi(B)$, counts the total number of fibral \mathbb{P}^1 's being wrapped by the rational section, S_0 . All of these statements will be confirmed for explicit examples in later sections.

Finally, let us fix some nomenclature. Since eq. (2.19) is to define the z coordinate of the Weierstrass model in section 3.1, we denote the associated line bundle as

$$L_z := \mathcal{O}_X(S_0) . \tag{2.26}$$

Likewise, we further define the two line bundles,

$$L_x := \mathcal{O}_X(2S_0) \otimes K_B^{-2} , \tag{2.27}$$

$$L_y := \mathcal{O}_X(3S_0) \otimes K_B^{-3} , \tag{2.28}$$

where K_B is the canonical bundle of the base B pulled back to X . The remaining Weierstrass coordinates x and y will then be constructed as GHS's of L_x and L_y , respectively, in manner similar to the method described above. This will be discussed further in section 3.1.

3 Singular fiber analysis

In order to understand the F-theory effective physics associated to an elliptically fibered Calabi-Yau manifold, X it is necessary to obtain a minimal limit of the geometry in which the fibers are irreducible and possibly singular. The structure of the singular fibers encodes information about non-Abelian gauge symmetries, charged matter and more in the effective theory [80, 81, 83].

In general, for an elliptically fibered CY n -fold, $X \rightarrow B$, with section there are three possible routes to a “minimal” form the geometry suitable for an F-theory limit. 1) By Nakayama’s theorem, any such elliptic fibration is birational to a Weierstrass model [29]. 2) To X we can associate the Jacobian, $J(X)$ [31] and finally 3) All reducible components of fibers can be explicitly blown down to form a “minimal model” in the sense of the Minimal Model Program (MMP) (see e.g. [33]). In the case of smooth, elliptically fibered CY 2-folds these three procedures all lead to the same simple/minimal geometry. However, in the case of CY n -folds with $n \geq 3$ these approaches can differ (and “minimal models” in the sense of the MMP are non-unique). For example, while the Weierstrass models of smooth CY 3-folds are birational to X , the topology of $J(X)$ and X may differ still further. Moreover, all three approaches can lead to different singular fibers at codimension 2 and higher in

the base. As a result, our focus will be primarily on Weierstrass models as constructed by Nakayama [29], however, as we will discuss below, the discriminant loci, $\Delta \subset B$, of these different forms are in fact identical and this observation, as well as information from the different approaches will be used to simply extract the gauge symmetries and charged matter from the original geometry.

We summarize below the ways that these three different approaches can be used to analyze the singular fibers of X . The first approach will be to construct a Weierstrass model X_W (built as a hypersurface defined in an ambient \mathbb{P}^2 -bundle defined by the projectivization of three line bundles over B), the second will be to form the Jacobian $J(X)$ of X (or that of its blow down), and the third will be to analyze the smooth (resolved) geometry X itself in order to directly study the singular fibers. A key object in studying the singular fibers of our manifolds is the discriminant locus for the elliptic fibration, and the following is expected of the triple of geometries, X_W , $J(X)$, and X :

- X is birational to its Weierstrass model X_W . Since the singular fibers of X remain singular after blowing down any \mathbb{P}^1 's therein, the discriminant loci of the two elliptic fibrations agree as an algebraic variety in the base B . We then have, in particular,

$$\Delta_{\text{res}} = \Delta_W, \tag{3.1}$$

where Δ_{res} and Δ_W denote the discriminant polynomials in B for X and X_W , respectively.

- The Jacobian⁶ $J(X)$ of X is constructed as a fibration over the same base B , in such a way that the fiber over each point is the Jacobian of the curve defined by the original fiber (that is, the moduli space of degree zero line bundles on that curve). This construction insures that the j -invariant of each fiber of $J(X)$ agrees with that of the original fiber of X . The two discriminant loci in B are identified and it is expected that

$$\Delta_{\text{res}} = \Delta_J, \tag{3.2}$$

where Δ_J denotes the discriminant polynomial in B for $J(X)$. Furthermore, at generic points on a codimension-one locus in the base, the singular fibers have the same Kodaira type [34, 35] in both geometries. As described above, the behaviors of the singular fibers at codimension two (or higher) may differ in general.

- Due to the two points made above, the Weierstrass model X_W and the Jacobian $J(X)$ also share the discriminant loci in B and we have

$$\Delta_W = \Delta_J. \tag{3.3}$$

Furthermore, the singular fibers of X_W and $J(X)$ are of the same Kodaira type generically at a codimension-one locus in the base.

⁶To avoid confusion it should be noted that in the mathematics literature, the “Jacobian” of a genus one fibered manifold $X \rightarrow B$ is sometimes taken to refer to *any* fibration of the form $X \rightarrow B$ with a section. Here we will reserve the terminology of “Jacobian” (or $J(X)$) to refer explicitly to a fibration constructed via the variable changes/procedure outlined in [2, 30, 31].

In what follows, we will sketch how the codimension-one locus in the base is obtained for each of the three approaches in turn, oftentimes returning to our simple $K3$ example whenever illustration is in need.

3.1 Weierstrass models

An elliptic fibration with a section can be associated to a Weierstrass model. Here we will follow a procedure due to Deligne (for elliptic curves) and Nakayama (for elliptic fibrations) [28, 29] which is well known in the physics literature (see refs. [72, 74] for some explicit examples of its application in such contexts).⁷ Schematically, a Weierstrass model for an elliptic fibration is built as a hypersurface constraint (with cubic fiber) within a projectivization of three line bundles $\mathbb{P}(\mathcal{O}_B \oplus \mathcal{L}^2 \oplus \mathcal{L}^3)$. The Calabi-Yau condition on X fixes $\mathcal{L} = K_B^{-1}$ and a change of variables makes it possible to describe the elliptic fiber as a degree 6 hypersurface in a weighted \mathbb{P}_{123} . In this description then, the fiber coordinates z, x, y (of weights $(1, 2, 3)$) are associated to global holomorphic sections of the following line bundles over X :

$$z \sim \mathcal{O}(S_0), \quad x \sim \mathcal{O}(2S_0) \otimes \mathcal{L}^2, \quad y \sim \mathcal{O}(3S_0) \otimes \mathcal{L}^3 \tag{3.4}$$

We first choose a zero section, with respect to which a Weierstrass model will be found, and denote the associated line bundle as L_z , which in particular satisfies $h^0(X, L_z) = 1$.

- Using the technique reviewed in the previous section (see the paragraph including eq. (2.23)), we can obtain the rational expression for the unique GHS of L_z . This will be the z coordinate of the Weierstrass form, expressed in terms of the coordinates of our original ambient space:

$$z = z(\mathbf{x}_1, \dots, \mathbf{x}_m) . \tag{3.5}$$

- Next we take the line bundle L_x as defined in eq. (2.27). The dimension of the space of global sections of this bundle, $h^0(X, L_x)$, is such that it is one larger than the subset of elements of that space that is spanned by a basis that can be written as z^2 multiplied by polynomials in the base coordinates. In essence, that additional element will describe the x coordinate of the Weierstrass model in terms of the original ambient space coordinates.

In practice one may take a generic element of $H^0(X, L_x)$ to describe x . This is simply because such a generic element will indeed not be proportional to z^2 , containing some portion of the remaining basis element in the cohomology group. There are of course a plethora of different generic elements that could then be chosen. This freedom simply corresponds to making different choices of the Weierstrass coordinate x , which are related under coordinate transformations mixing x and z which maintain the Weierstrass form. After choosing a generic element of $H^0(X, L_x)$ we can now employ

⁷We would like to thank T. Pantev for very useful conversations about algorithmically applying this procedure.

the methodology described in the proceeding section to obtain an explicit description of the x coordinate of the Weierstrass form.

$$x = x(\mathbf{x}_1, \dots, \mathbf{x}_m) \tag{3.6}$$

- The final Weierstrass coordinate y is obtained in a very similar manner as a generic element of $H^0(X, L_y)$. In this instance, the dimension of the cohomology group, $h^0(X, L_y)$, is such that it is one larger than the subset of elements of that space that is spanned by a basis that can be written as z^3 or xz multiplied by a polynomial in the base. As before, different choices for the coordinate y will lead to Tate forms for the fibration which are related by coordinate transformations mixing x, y and z which leave the Weierstrass form invariant. Once more, we employ the technology of the previous section to obtain the explicit description of the coordinate at hand in terms of those of the original description of the ambient space of the manifold.

$$y = y(\mathbf{x}_1, \dots, \mathbf{x}_m) \tag{3.7}$$

- Finally, given that we now have explicit expressions for x, y and z in terms of our original ambient space coordinates, we can now, by straightforward calculation, find a relationship between them that is of the Tate form [36] (up to scaling),

$$y^2 + c_1xyz + c_3yz^3 + c_0x^3 + c_2x^2z^2 + c_4xz^4 + c_6z^6 = 0. \tag{3.8}$$

Here, the c_i 's are GHS's of K_B^{-i} , that is functions of the base coordinates of specific degrees. That there is a unique such relation follows from similar arguments to those given in the previous bullet points. The left hand side of the relation eq. (3.8) is associated to an element of $H^0(\mathcal{O}_X(6S_0) \otimes K_B^{-6})$. The dimension of this space is one less than that naively spanned by elements of appropriate degree that can be written as $y^2, xyz, yz^3, x^3, x^2z^2, xz^4$ and z^6 multiplied by elements of the relevant $H^0(X, K_B^{-i})$. Thus there must be one relation between these quantities which vanishes as in eq. (3.8).

Practically to find the relationship in eq. (3.8) we employ a similar technique to that discussed in the paragraph under eq. (2.25). We write out a generic relation of the correct form with undetermined numerical coefficients. In particular, we form a basis M^{l_i} of $H^0(X, K_B^{-i})$ where $l_i = 1, \dots, h^0(X, K_B^{-i})$. Expanding the c_i 's we then have the following,

$$c_i = c_{i l_i} M^{l_i}, \tag{3.9}$$

which is taken to define the $c_{i l_i}$ — our numerical coefficients.

Writing eq. (3.8) in terms of the $c_{i l_i}$ and the ambient coordinates of the original description of the manifold, we then substitute in coordinates of a point on the manifold (solved for using our original description of the space) to obtain a relationship between the $c_{i l_i}$. Repeating this procedure with sufficient numbers of points on the

manifold we obtain a system of linear equations in the numerical coefficients which, by a naive counting of equations, would seem to be over-constrained. Nevertheless, this system can then be solved to uniquely determine the $c_i l_i$, and thus the relationship eq. (3.8). The fact that such parameter values can be found satisfying the equation system is, in itself, a reassuring check of the method.

As an illustration, let us return to the $K3$ geometry with configuration (2.6) with the line bundle $L_z = \mathcal{O}_X(-1, 1, 1)$, which has already been proven to give a section to the elliptic fibration. The GHS z of L_z can be obtained explicitly as eq. (2.25), and similarly, one may easily find the generic GHS's associated to x and y using the same technique. Here we will simply note that, in performing this computation, since $K_B^{-1} = \mathcal{O}(0, 0, 2)$ and $\mathcal{O}_X(S_0) = \mathcal{O}_X(-1, 1, 1)$, the other relevant line bundles are given as $L_x = \mathcal{O}_X(2S_0) \otimes K_B^{-2} = \mathcal{O}_X(-2, 2, 6)$ and $L_y = \mathcal{O}_X(3S_0) \otimes K_B^{-3} = \mathcal{O}_X(-3, 3, 9)$, respectively. Similarly, the relationship eq. (3.8) is associated with an element of the zeroth cohomology of $\mathcal{O}_X(6S_0) \otimes K_B^{-6} = \mathcal{O}_X(-6, 6, 18)$. We omit the explicit expressions for x and y here, and the detailed form of the final Weierstrass form due to their length. More details can be found for this specific example in section 5.2.1.

Once the relation (3.8) is obtained, via an appropriate rescaling of x we can put eq. (3.8) into the standard Tate form,

$$y^2 + a_1xyz + a_3yz^3 = x^3 + a_2x^2z^2 + a_4xz^4 + a_6z^6, \tag{3.10}$$

where $a_i \in \Gamma(B, K_B^{-i})$, and then also into the Weierstrass form,

$$y^2 = x^3 + f_Wxz^4 + g_Wz^6, \tag{3.11}$$

where

$$f_W = -\frac{1}{48}(b_2^2 - 24b_4), \tag{3.12}$$

$$g_W = -\frac{1}{864}(-b_2^3 + 36b_2b_4 - 216b_6), \tag{3.13}$$

with

$$b_2 = a_1^2 + 4a_2, \tag{3.14}$$

$$b_4 = 2a_4 + a_1a_3, \tag{3.15}$$

$$b_6 = a_3^2 + 4a_6. \tag{3.16}$$

In particular, the discriminant polynomial for such a Weierstrass model is given as

$$\Delta_W = 4f_W^3 + 27g_W^2. \tag{3.17}$$

Once the Tate form/Weierstrass form has been obtained one can use the standard techniques in order to analyze the singular fibers.

3.2 Jacobians

In the case where the fiber of a CICY is realized as a complete intersection of codimension one or two, the Jacobian of X_n can be formed by using the results in refs. [2, 30, 31].⁸ This work provides a list of Jacobians for all elliptic fibers realized as a complete intersection of codimension 1 or 2 in any toric variety, thus including products of projective spaces as a special case. More generally, however, the fibers we will encounter can be of higher codimension. Rather than generalizing the results of refs. [2, 30, 31] in these cases, we find it more expedient to blow down the fiber until it reaches a codimension one or two description by utilizing the process of “contraction” [15].

Contraction refers to the procedure of making the configuration matrix smaller by removing a row of 1’s as follows,

$$\left[\begin{array}{c|cccc} \mathbb{P}^a & 1 & \cdots & 1 & 0 \\ \mathcal{A}' & \mathbf{u}_1 & \cdots & \mathbf{u}_{a+1} & C \end{array} \right] \longrightarrow \left[\mathcal{A}' \left\| \sum_{i=1}^{a+1} \mathbf{u}_i C \right. \right]. \quad (3.18)$$

Here the first $a + 1$ polynomials have merged to a single determinantal polynomial. We perform this procedure in such a manner that, while the description of the fiber in eq. (2.4) is changed, the description of the base, $B = [\mathcal{A}_B || \mathcal{B}]$, remains invariant. As has been described in ref. [15], this procedure corresponds to blowing down \mathbb{P}^1 ’s in a manner that may or may not be associated with a geometric transition. In the case where the final manifold is different from the initial one we say the contraction is effective and otherwise it is ineffective. In our case we are clearly blowing down \mathbb{P}^1 ’s in the fiber.

Even though the final manifold after completing this contraction may be different from our starting configuration it will have the same discriminant locus in the same base (and thus, so will its Jacobian). This is simply because the \mathbb{P}^1 ’s in the fiber that are being blown down are singular before and after the process and therefore project to a point on the discriminant in both cases. Since the discriminant is what we are trying to obtain here, we are able to contract to get a codimension one or two fiber and then make use of the aforementioned existing results without any loss of information.

In making use of the results of refs. [2, 30, 31], for a given complex structure for X_n (or its blow down), we consider the defining equations for the fiber as polynomials in the fibral homogeneous coordinates, demoting the base coordinates to parameters. Then the Jacobian of the form,

$$y^2 = x^3 + f_J x z^4 + g_J z^6, \quad (3.19)$$

can be immediately read off, and its discriminant locus is obtained in turn by the zero locus of

$$\Delta_J = 4f_J^3 + 27g_J^2, \quad (3.20)$$

in the normal way. This is a polynomial in the base coordinates that we now promote to variables again.

⁸The results of [30] can be searched from the following website: <http://wwwth.mpp.mpg.de/members/jkeitel/Weierstrass/>.

For an illustration, let us return to our toy example of the $K3$ configuration (2.6). Its fiber configuration (2.7),

$$F = \left[\begin{array}{c|cc} \mathbb{P}_{\mathbf{x}_1}^1 & 1 & 1 \\ \mathbb{P}_{\mathbf{x}_2}^2 & 1 & 2 \end{array} \right], \quad (3.21)$$

is already a complete intersection of codimension two in $\mathcal{A}_F = \mathbb{P}_{\mathbf{x}_1}^1 \times \mathbb{P}_{\mathbf{x}_2}^2$. The defining relations are bi-degree (1, 1) and (1, 2) polynomials and this codimension-two fiber has the PALP ID (4, 0), which, via the result of ref. [30], can straightforwardly be transformed into the Jacobian. The two defining polynomials, due to the base twist in the full configuration (2.6), also depend on \mathbf{x}_3 in our case. However, we demote these variables to parameters so that the various monomial coefficients in the defining equations can be thought of as a polynomial parameterized by the base coordinates \mathbf{x}_3 . The expressions for f_J and g_J in terms of those monomial coefficients are immediately found and thereby one obtains $f_J \in \Gamma(B, K_B^{-4})$ and $g_J \in \Gamma(B, K_B^{-6})$ explicitly.

In every case we have computed the discriminant obtained in this manner has matched that obtained by the other two methods discussed in this section. This is a highly non-trivial check of the above procedure, several explicit examples of which will be provided in later sections.

3.3 Resolved geometries

One of the benefits to the approach being espoused here for constructing F-theory compactifications is that the resolved space associated to the models is known from the start. Bertini's theorem (see refs. [75, 76]) guarantees that a CICY of the form being considered is smooth, presuming of course that a generic enough complex structure is chosen.

In the original description of the manifold, the discriminant locus of the fibration can be explored by directly computing where over the base the fiber becomes singular. Starting with the configuration matrix of the form (2.4),

$$X_n = \left[\begin{array}{c|cc} \mathcal{A}_F & \mathcal{O} & \mathcal{F} \\ \mathcal{A}_B & \mathcal{B} & \mathcal{T} \end{array} \right], \quad (3.22)$$

one can perform such a computation as follows.

- We will denote the coordinates of the base ambient space, \mathcal{A}_B by \mathbf{x}_B , and the coordinates of the fiber ambient space, \mathcal{A}_F by \mathbf{x}_F .
- We write as $P_{\hat{j}}$ those defining relations associated to the last block of columns in eq. (3.22), that is those defining equations associated to the following portion of the configuration matrix.

$$X_n = \left[\begin{array}{c|c} \mathcal{A}_F & \mathcal{F} \\ \mathcal{A}_B & \mathcal{T} \end{array} \right]. \quad (3.23)$$

- We then form the equation system

$$P_1 = \cdots = P_K = 0, \quad \wedge_{\hat{j}} d_F P_{\hat{j}} = 0, \quad (3.24)$$

where the exterior derivative d_F is only taken with respect to the variables \mathbf{x}_F and not with respect to \mathbf{x}_B . The differentiated conditions here describe when the normal form to the fiber is ill defined — that is they describe for what values of the ambient space coordinates the fiber becomes singular. Including the defining relations in the equation system then gives us the points on the Calabi-Yau itself where the fiber becomes singular.

- Finally we want to project the equation system eq. (3.24) to obtain those points on the base above which there are singularities in the fiber. This projection is equivalent algebraically to the process of elimination. We must eliminate the variables \mathbf{x}_F to obtain a necessary and sufficient set of relations on the \mathbf{x}_B such that there is a solution to eq. (3.24) for some value of the fiber coordinates. Thinking of the equations as generators of an ideal we wish to form

$$\langle P_1, \dots, P_K, \wedge_j d_F P_j \rangle \cap \mathbb{C}[\mathbf{x}_B]. \quad (3.25)$$

Here by an abuse of notation we have denoted the polynomial coefficients of an expansion of $\wedge_j d_F P_j$ in a basis of forms of an appropriate degree by the expression itself. Such an elimination can easily be performed with a Gröbner basis computation and results in a set of equations in the variables \mathbf{x}_B describing the discriminant locus in the base of the fibration.

The last step in this procedure, the elimination process, is computationally expensive in large examples. Nevertheless, the full description (3.24) of location of fiber singularities in the total space of the Calabi-Yau manifold is already useful in comparing to computations performed using the preceding methods described in this section. We will return to comparing the discriminants we have found in various ways, which, as expected, have matched in every case we have investigated, in the Example sections to follow. It should also be noted that, given a discriminant locus derived using one of the other methods in this section, one could use the initial description of the manifold to investigate the nature of the associated singular fibers. One would simply choose values of \mathbf{x}_B which lie on the discriminant and substitute these into the P_j to obtain an explicit description of the singular fiber as a variety in \mathcal{A}_F . We will return to this point later.

4 The Mordell-Weil group

4.1 Decomposition of the Picard lattice

For an elliptic Calabi-Yau manifold X_n with a section, the Shioda-Tate-Wazir theorem [45] states that the Picard lattice of X_n is generated by linearly independent basis elements of the following four types:⁹ (1) base divisor classes pulled back to X_n , (2) “fibrals” divisors associated to blow-ups in the fiber (i.e., vertical divisor classes that are not pulled back from the base), (3) the zero section, and (4) a basis of the rational sections generating (the

⁹The statement is for the Néron-Severi lattice, which coincides with the Picard lattice for a Calabi-Yau manifold.

free part) of the Mordell-Weil (MW) group (the additive group of sections to the elliptic fibration — see, for example, ref. [77]). This in particular implies the following dimensional relation for $n \geq 3$:

$$h^{1,1}(X_n) = h^{1,1}(B_{n-1}) + \sum_A \text{rk } G_A + 1 + \text{rk } MW(X_n) . \quad (4.1)$$

Here G_A are the non-abelian Lie groups, each associated with the reducible fiber type over an irreducible component of the discriminant locus in B_{n-1} , and $MW(X_n)$ denotes the Mordell-Weil group of X_n .

It is worth emphasizing how useful eq. (4.1) is for our purposes. Given a specific configuration for X_n and B_{n-1} in the form (2.4), it is a straightforward exercise in algebraic topology to compute $h^{1,1}(X_n)$ and $h^{1,1}(B_{n-1})$. Furthermore, factorization of the discriminant equation is straightforward, from which one can easily read off the enhancement pattern of the fiber singularity, and in particular, $\text{rk } G_i$. On the other hand, determination of $\text{rk } MW(X_n)$ involves a careful analysis of the section structure, which in many cases is a difficult task. Thus, in analyzing the MW group structure, the relation (4.1) can be used as either a consistency check on a direct computation or an indirect method to determine the MW rank, as will be illustrated with examples in later sections.

From the physical perspective, eq. (4.1) also plays an important role in systematic exploration of the F-theory vacua from the plethora of elliptically fibered CICY threefolds and fourfolds [19, 32]. Upon compactifying F-theory over an elliptic Calabi-Yau manifold X_n , one obtains a $(12 - 2n)$ -dimensional effective theory with gauge group of the form,

$$G = \text{U}(1)^{\text{rk } MW(X)} \times \prod_A G_A . \quad (4.2)$$

Since $h^{1,1}(X_n)$ and $h^{1,1}(B_{n-1})$ can be computed in a systematic manner for CICYs, eq. (4.1) makes it easy to classify the F-theory vacua with a fixed total rank, $\text{rk } G$, of the gauge group. Furthermore, a relatively straightforward analysis of the discriminant locus and of the enhancement pattern of the fiber singularities of the manifold can be used to determine the non-abelian part of the gauge group in a systematic manner. Thus, it is possible to systematically explore F-theory vacua with a fixed gauge group in the context of CICY manifolds, which is a topic that we will return to in future work [32].

4.2 Arithmetic of the sections

Although the decomposition of the Picard lattice described in subsection 4.1 reveals the rank of the MW group in a systematic manner, it is a rather indirect procedure in that one still does not have explicit forms for the generating sections. In this subsection, we review what is known about the arithmetic of rational sections which, when combined with the section construction technology described in section 2, allows us to obtain an explicit description of the MW group.

The arithmetic of sections was derived at the level of divisor classes in refs. [67, 71, 72], resulting in the following group law under the section addition, ‘ \oplus ’:

$$\text{Div}(\sigma_1 \oplus \underbrace{\sigma_2 \oplus \cdots \oplus \sigma_2}_{k \text{ times}}) = S_1 + k(S_2 - S_0) - k\pi((S_1 - kS_0) \cdot (S_2 - S_0)) . \quad (4.3)$$

Here, $S_m := \text{Div}(\sigma_m)$, for $m = 0, 1, 2$, denote the divisor classes associated to the sections σ_m and we have chosen to identify the zero section as σ_0 . The projection π of the intersection, $D \cdot D'$, of two divisors D and D' in X_n is defined in refs. [67, 71]. In particular, for the $n = 3$ case, which we will give examples of in sections 6–8, the projection is given by

$$\pi(D \cdot D') := (D \cdot D' \cdot D^{\text{b},\alpha})D_\alpha^{\text{b}}, \tag{4.4}$$

where the index $\alpha = 1, \dots, h^{1,1}(B)$ is raised and lowered by the intersection matrix,

$$\eta_{\alpha\beta} = \hat{D}_\alpha^{\text{b}} \cdot \hat{D}_\beta^{\text{b}} \tag{4.5}$$

of the base two-fold. For the simpler case of $n = 2$, one can also show that the appropriate projection has to be defined as

$$\pi(D \cdot D') := (D \cdot D')D^{\text{b}}, \tag{4.6}$$

where D^{b} is the pull-back of a hyperplane class in the base.

The procedures described in section 2 can be used to find divisor classes corresponding to true sections for a given elliptically fibered Calabi-Yau manifold. If enough sections are found this way, given the rank of the MW group, one can choose a zero section and then use the remaining $\text{rk } MW(X)$ generators and the addition law (4.3) to form a complete basis of the MW group.

As an illustration, we return to the $K3$ example. Let us choose the zero section σ_0 to be the one we have obtained in section 2 with the class $\mathcal{O}_X(S_0) = \mathcal{O}_X(-1, 1, 1)$. In section 5, it will also be shown that $\mathcal{O}_X(S_1) = \mathcal{O}_X(2, -1, 4)$ represents another section, call it σ_1 , and here we will use this fact. The divisor classes of $k\sigma_1$ are then given by substituting $\sigma_1 \rightarrow \sigma_0$ and $\sigma_2 \rightarrow \sigma_1$ in the formula (4.3),

$$\text{Div}(k\sigma_1) = S_0 + k(S_1 - S_0) + k(k - 1) \pi(S_0 \cdot (S_1 - S_0)), \tag{4.7}$$

$$= (-1 + 3k)J_1 + (1 - 2k)J_2 + (1 - 11k + 14k^2)J_3, \tag{4.8}$$

which reproduces all the putative sections in eq. (2.18). Therefore, given that S_0 and S_1 are sections, each of those putative section classes has to also correspond to a true section, and furthermore, this proves that σ_1 fully generates the rank-one MW group.

In the remaining sections of this paper we will demonstrate the details of the above discussions with a series of explicit examples.

5 Example 1: a K3 example

In this section, we provide a complete analysis of the $K3$ geometry, (2.6), that we used to illustrate our general analysis in the preceding sections

$$X_2 = \left[\begin{array}{c|cc} \mathbb{P}_{\mathbf{x}_1}^1 & 1 & 1 \\ \mathbb{P}_{\mathbf{x}_2}^2 & 1 & 2 \\ \mathbb{P}_{\mathbf{x}_3}^1 & 1 & 1 \end{array} \right]. \tag{5.1}$$

This $K3$ surface admits an obvious genus-one fibration structure over the base $B_1 = \mathbb{P}_{\mathbf{x}_3}^1$, where the configuration of the fiber is given by,

$$F = \left[\begin{array}{c} \mathbb{P}_{\mathbf{x}_1}^1 \\ \mathbb{P}_{\mathbf{x}_2}^2 \end{array} \middle| \begin{array}{c} 1 \ 1 \\ 1 \ 2 \end{array} \right]. \quad (5.2)$$

For the purpose of giving explicit examples of results for this configuration, we choose the following generic complex structure:

$$\begin{aligned} P_1(\mathbf{x}_1, \mathbf{x}_2, \mathbf{x}_3) = & 5x_{1,0}x_{2,0}x_{3,0} + 9x_{1,1}x_{2,0}x_{3,0} - 11x_{1,0}x_{2,1}x_{3,0} + 13x_{1,0}x_{2,2}x_{3,0} + 7x_{1,1}x_{2,2}x_{3,0} \\ & - 17x_{1,0}x_{2,0}x_{3,1} - 17x_{1,1}x_{2,0}x_{3,1} + 19x_{1,0}x_{2,1}x_{3,1} - 14x_{1,1}x_{2,1}x_{3,1} + 6x_{1,0}x_{2,2}x_{3,1} \\ & - 12x_{1,1}x_{2,2}x_{3,1}, \end{aligned} \quad (5.3)$$

$$\begin{aligned} P_2(\mathbf{x}_1, \mathbf{x}_2, \mathbf{x}_3) = & -8x_{1,0}x_{2,0}^2x_{3,0} - 5x_{1,1}x_{2,0}^2x_{3,0} - 11x_{1,0}x_{2,0}x_{2,1}x_{3,0} + 5x_{1,1}x_{2,0}x_{2,1}x_{3,0} + 7x_{1,0}x_{2,1}^2x_{3,0} \\ & + 16x_{1,1}x_{2,1}^2x_{3,0} - 13x_{1,0}x_{2,0}x_{2,2}x_{3,0} + x_{1,1}x_{2,0}x_{2,2}x_{3,0} - 20x_{1,0}x_{2,1}x_{2,2}x_{3,0} \\ & - 20x_{1,1}x_{2,1}x_{2,2}x_{3,0} + 15x_{1,0}x_{2,2}^2x_{3,0} - 12x_{1,1}x_{2,2}^2x_{3,0} + 12x_{1,0}x_{2,0}^2x_{3,1} + 6x_{1,1}x_{2,0}^2x_{3,1} \\ & - 8x_{1,0}x_{2,0}x_{2,1}x_{3,1} - 13x_{1,1}x_{2,0}x_{2,1}x_{3,1} - 9x_{1,0}x_{2,1}^2x_{3,1} - 16x_{1,1}x_{2,1}^2x_{3,1} \\ & - 16x_{1,0}x_{2,0}x_{2,2}x_{3,1} + 19x_{1,1}x_{2,0}x_{2,2}x_{3,1} + 9x_{1,0}x_{2,1}x_{2,2}x_{3,1} + 13x_{1,1}x_{2,1}x_{2,2}x_{3,1} \\ & - 13x_{1,0}x_{2,2}^2x_{3,1} + 15x_{1,1}x_{2,2}^2x_{3,1}. \end{aligned} \quad (5.4)$$

5.1 Section analysis

Putative sections. Let us start with the classification of putative sections. For a putative section S labelled by $\mathcal{O}_X(S) = \mathcal{O}_X(b_1, b_2, b_3)$, given that the minimal base-point form integrates to unity,

$$\int_{\mathbb{P}_{\mathbf{x}_3}^1} J_3 = 1, \quad (5.5)$$

the intersection of S with the generic fiber F_p is computed, as in section 2.2, as follows.

$$\int_{\mathbb{P}_{\mathbf{x}_1}^1 \times \mathbb{P}_{\mathbf{x}_2}^2 \times \mathbb{P}_{\mathbf{x}_3}^1} [(b_1 J_1 + b_2 J_2 + b_3 J_3) \wedge J_3] \wedge [(J_1 + J_2 + J_3) \wedge (J_1 + 2J_2 + J_3)] = 2b_1 + 3b_2, \quad (5.6)$$

Here J_1 and J_2 are the Kähler forms of $\mathbb{P}_{\mathbf{x}_1}^1$ and $\mathbb{P}_{\mathbf{x}_2}^2$, respectively. Therefore, the Oguiso criterion (a) demands that

$$2b_1 + 3b_2 \stackrel{!}{=} 1. \quad (5.7)$$

Via a similar intersection computation, the second criterion (2.12) leads to¹⁰

$$6b_1b_2 + 2b_2^2 + 4b_1b_3 + 6b_2b_3 \stackrel{!}{=} -2, \quad (5.8)$$

where eq. (5.7) has been used in simplifying the result. These two putative-section conditions, (5.7) and (5.8), can be solved as

$$b_1 = -1 - 3k, \quad b_2 = 1 + 2k, \quad b_3 = 1 + 11k + 14k^2. \quad (5.9)$$

¹⁰For a smooth divisor S , the left hand side of eq. (5.8) is $-\chi(S)$ and hence, the criterion is equivalent to $\chi(S) = \chi(B)$ unless singularities are involved.

with an integer parameter $k \in \mathbb{Z}$. For instance, with $k = 1, 0, -1, -2$, we obtain the following putative section classes, $\mathcal{O}_X(-4, 3, 26)$, $\mathcal{O}_X(-1, 1, 1)$, $\mathcal{O}_X(2, -1, 4)$, and $\mathcal{O}_X(5, -3, 35)$, respectively.

Now the question arises as to whether the putative section classes in eq. (5.9) indeed correspond to sections. In what follows, we will first show that two of them do by providing explicit expressions for the sections themselves (the rest will also prove to be a section later in section 5.3).

Explicit expressions for the sections. Let us proceed with the methodology described in section 2. Here we will use the alternative approach discussed in the paragraph under eq. (2.25) and show that it gives the same result as the method employed explicitly in that section. We will focus on the line bundle $\mathcal{O}_X(S) = \mathcal{O}_X(-1, 1, 1)$, taking the solution $(b_1, b_2, b_3) = (-1, 1, 1)$ from the family (5.9). The divisor class $[S] = -J_1 + J_2 + J_3$ naturally splits into two effective pieces, S_{zero} and S_{poles} , such that

$$[S_{\text{zero}}] = J_2 + J_3, \quad [S_{\text{pole}}] = J_1, \quad (5.10)$$

that intersect with the generic fiber at 3 and 2 points, respectively (see eq. (5.6)). The GHS of $\mathcal{O}_X(S)$ can then be constructed by appropriately choosing two GHS's,

$$s_{\text{zero}} \in H^0(X, \mathcal{O}_X(S_{\text{zero}})), \quad s_{\text{pole}} \in H^0(X, \mathcal{O}_X(S_{\text{pole}})), \quad (5.11)$$

of $\mathcal{O}_X(S_{\text{zero}}) = \mathcal{O}_X(0, 1, 1)$ and $\mathcal{O}_X(S_{\text{pole}}) = \mathcal{O}_X(1, 0, 0)$ so that along the generic fiber F_p over $p \in B$ the two points of $F_p \cap S_{\text{pole}}$ match with two of the three points of $F_p \cap S_{\text{zero}}$. The unmatched point of $F_p \cap S_{\text{zero}}$ should be the single intersection point of $F_p \cap S$.

To be more concrete, let us illustrate the procedure with explicit expressions, given the complex structure in eqs. (5.3) and (5.4). For a generic random choice of GHS of $\mathcal{O}_X(S_{\text{pole}}) = \mathcal{O}_X(1, 0, 0)$, for example,

$$s_{\text{pole}}(\mathbf{x}_1, \mathbf{x}_2, \mathbf{x}_3) = x_{1,0} + 10x_{1,1}, \quad (5.12)$$

we shall look for an appropriate section of $\mathcal{O}_X(S_{\text{zero}}) = \mathcal{O}_X(0, 1, 1)$,

$$s_{\text{zero}}(\mathbf{x}_1, \mathbf{x}_2, \mathbf{x}_3) = C_1x_{2,0}x_{3,0} + C_2x_{2,0}x_{3,1} + C_3x_{2,1}x_{3,0} + C_4x_{2,1}x_{3,1} + C_5x_{2,2}x_{3,0} + C_6x_{2,2}x_{3,1}, \quad (5.13)$$

for which $F_p \cap S_{\text{pole}} \subset F_p \cap S_{\text{zero}}$. Demoting the base coordinates $\mathbf{x}_3 = (x_{3,0} : x_{3,1})$ to parameters and solving the system,

$$P_1(\mathbf{x}_1, \mathbf{x}_2, \mathbf{x}_3) = 0, \quad (5.14)$$

$$P_2(\mathbf{x}_1, \mathbf{x}_2, \mathbf{x}_3) = 0, \quad (5.15)$$

$$s_{\text{pole}}(\mathbf{x}_1, \mathbf{x}_2, \mathbf{x}_3) = 0, \quad (5.16)$$

for \mathbf{x}_1 and \mathbf{x}_2 , one obtains two solutions $(\mathbf{x}_1^{(a)}, \mathbf{x}_2^{(a)}) = (\mathbf{x}_1^{(a)}(\mathbf{x}_3), \mathbf{x}_2^{(a)}(\mathbf{x}_3))$, for $a = 1, 2$. We then substitute each of these to eq. (5.13) and demand that $s_{\text{zero}}(\mathbf{x}_1^{(a)}, \mathbf{x}_2^{(a)}) \equiv 0$ as a function of \mathbf{x}_3 , for $a = 1, 2$. This turns out to fix the section s_{zero} uniquely (up to scaling) as

$$(C_1, C_2, C_3, C_4, C_5, C_6) = \lambda (41, -153, -110, 204, 123, 72), \quad \text{with } \lambda \in \mathbb{C}. \quad (5.17)$$

As promised in the general discussion of section 2.3, one can immediately confirm that this is equivalent to eq. (2.25), where the numerator is given as

$$\begin{aligned}
 -10p_1 + p_2 &= -10(5x_{2,0}x_{3,0} - 11x_{2,1}x_{3,0} + 13x_{2,2}x_{3,0} - 17x_{2,0}x_{3,1} + 19x_{2,1}x_{3,1} + 6x_{2,2}x_{3,1}) \\
 &\quad + (9x_{2,0}x_{3,0} + 7x_{2,2}x_{3,0} - 17x_{2,0}x_{3,1} - 14x_{2,1}x_{3,1} - 12x_{2,2}x_{3,1}) \quad (5.18) \\
 &= -41x_{2,0}x_{3,0} + 153x_{2,0}x_{3,1} + 110x_{2,1}x_{3,0} - 204x_{2,1}x_{3,1} - 123x_{2,2}x_{3,0} - 72x_{2,2}x_{3,1},
 \end{aligned}$$

which agrees with eq. (5.17) for $\lambda = -1$.

Finally, having specified the two divisors, S_{zero} and S_{pole} , of X as the vanishing loci of s_{zero} and s_{pole} , respectively, we can now explicitly parameterize the section, $S = S_{\text{zero}} \setminus S_{\text{pole}}$, of the elliptic fibration in terms of the base coordinates \mathbf{x}_3 . We obtain the following explicit parametric expression for the section,

$$x_{1,0} = A_{1,0}(\mathbf{x}_3), \quad x_{1,1} = A_{1,1}(\mathbf{x}_3); \quad x_{2,0} = A_{2,0}(\mathbf{x}_3), \quad x_{2,1} = A_{2,1}(\mathbf{x}_3), \quad x_{2,2} = A_{2,2}(\mathbf{x}_3). \quad (5.19)$$

Here, the polynomials $A_{1,i}(\mathbf{x}_3)$ for $i = 0, 1$, as well as $A_{2,i}(\mathbf{x}_3)$ for $i = 0, 1, 2$, are, respectively, quintic and quadratic polynomials given by

$$\begin{aligned}
 A_{1,0}(\mathbf{x}_3) &= 241226x_{3,0}^5 - 2444409x_{3,0}^4x_{3,1} + 6970327x_{3,0}^3x_{3,1}^2 - 4889388x_{3,0}^2x_{3,1}^3 \\
 &\quad - 2858859x_{3,0}x_{3,1}^4 + 992331x_{3,1}^5, \\
 A_{1,1}(\mathbf{x}_3) &= 152844x_{3,0}^5 - 1296506x_{3,0}^4x_{3,1} + 3553577x_{3,0}^3x_{3,1}^2 - 8289055x_{3,0}^2x_{3,1}^3 \\
 &\quad + 11322255x_{3,0}x_{3,1}^4 - 5290227x_{3,1}^5, \\
 A_{2,0}(\mathbf{x}_3) &= 77x_{3,0}^2 - 447x_{3,0}x_{3,1} + 144x_{3,1}^2, \\
 A_{2,1}(\mathbf{x}_3) &= -82x_{3,0}^2 - 12x_{3,0}x_{3,1} + 306x_{3,1}^2, \\
 A_{2,2}(\mathbf{x}_3) &= -99x_{3,0}^2 + 428x_{3,0}x_{3,1} - 561x_{3,1}^2. \quad (5.20)
 \end{aligned}$$

These expressions define a rational map from the base, parameterized by \mathbf{x}_3 , to the fiber, parameterized by \mathbf{x}_1 and \mathbf{x}_2 . The map is well defined over every point on the base. That is, for every choice of \mathbf{x}_3 on the base manifold, we obtain a valid set of homogeneous coordinates \mathbf{x}_1 and \mathbf{x}_2 on the fiber. In particular, for no point on the base do we find that all of the homogeneous coordinates in a fiber ambient projective space factor simultaneously vanish. We thus conclude that the map is a holomorphic section, and confirm that the line bundle $\mathcal{O}_X(S) = \mathcal{O}_X(-1, 1, 1)$ is associated to this holomorphic section to the elliptic fibration.

We can perform the same analysis for the second putative section, $\mathcal{O}_X(2, -1, 4)$. In order to find the associated rational map, we first need to find the rational expression for the GHS of this line bundle. With the negative degree in the $\mathbb{P}_{\mathbf{x}_2}^2$, however, the naive ansatz with the linear denominator in \mathbf{x}_2 does not work.¹¹ It turns out that the GHS can

¹¹With a generic linear denominator, global holomorphicity of the rational expression, when imposed on a numerator polynomial of multi-degree $(2, 0, 4)$, leads to an identically vanishing holomorphic section. For a brief hint as to why this degree shifting is frequently necessary consider a generic polynomial defining relation of multi-degree $(d_1, 1, d_3)$ in $\mathbb{P}_{\mathbf{x}_1}^{n_1} \times \mathbb{P}_{\mathbf{x}_2}^2 \times \mathbb{P}_{\mathbf{x}_3}^{n_3}$ with linear dependence in the \mathbb{P}^2 direction: $x_{2,0}p_0 + x_{2,1}p_1 + x_{2,2}p_2 = 0$. For a regular rational function defined over this coordinate ring, a linear monomial denominator — $x_{2,0}$ for instance — implies an associated vanishing polynomial of the form $x_{2,1}p_1 + x_{2,2}p_2 = 0$, which

be found once we shift both the numerator and the denominator multi-degrees by $(0, 2, 0)$, i.e., with the modified ansatz,

$$z = \frac{N(\mathbf{x}_1, \mathbf{x}_2, \mathbf{x}_3)}{D(\mathbf{x}_2)}, \quad (5.21)$$

where $N \in \Gamma(X, \mathcal{O}_X(2, 2, 4))$ and $D \in \Gamma(X, \mathcal{O}_X(0, 3, 0))$. Let us take $D(\mathbf{x}_2) = x_{2,0}x_{2,1}x_{2,2}$ and tune the 90 monomial coefficients in N . That is, we demand that N vanishes on each of the three divisors $\{x_{2,i} = 0\} \subset X$. Practically, we achieve this by substituting a sufficiently large number of points on these loci into the equation for the divisor. Furthermore, we also demand that N appropriately vanishes to order 2 on all of the points in X with $\mathbf{x}_2 = (1 : 0 : 0)$, $(0 : 1 : 0)$, or $(0 : 0 : 1)$. Then we find that only 9 of the 90 coefficients are undetermined. With such a tuning, $z = N/D$ is globally holomorphic. This may at first sound strange since we know that $h^0(X, \mathcal{O}_X(2, -1, 4)) = 1$. The only way to make sense of this result is that the 9-parameter expression we have obtained for the GHS's of $\mathcal{O}_X(2, -1, 4)$ should only span a one-dimensional vector space of GHS. Indeed, this turns out to be the case and in the coordinate ring of X they all lead to one and the same GHS up to scaling. Having specified $z = N/D$, we can now proceed to find a generic parametrization of its zero locus in terms of the base coordinates \mathbf{x}_3 . This results in a parametric expression of the form,

$$x_{1,i} = A'_{1,i}(\mathbf{x}_3, \mathbf{x}_4) \in \Gamma(B, \mathcal{O}_B(5)), \quad \text{for } i = 0, 1, \quad (5.22)$$

$$x_{2,i} = A'_{2,i}(\mathbf{x}_3, \mathbf{x}_4) \in \Gamma(B, \mathcal{O}_B(16)), \quad \text{for } i = 0, 1, 2, \quad (5.23)$$

where $A'_{1,i}$ coincide with the $A_{1,i}$ from eq. (5.20) and $A'_{2,i}$ are some fixed degree-16 polynomials, which we do not display in this paper (each of the 17 integer coefficients has 17 to 24 digits). The map is holomorphic and we thus confirm that the line bundle $\mathcal{O}_X(S) = \mathcal{O}_X(2, -1, 4)$ corresponds to another holomorphic section to the elliptic fibration.

5.2 Locating the singular fibers

5.2.1 Weierstrass model

Having proven that the unique GHS of the line bundle,

$$L_z := \mathcal{O}_X(S_0) = \mathcal{O}_X(-1, 1, 1), \quad (5.24)$$

is a holomorphic section to the elliptic fibration, we choose to use L_z as the bundle, which the Weierstrass coordinate z is a GHS of. Then the other Weierstrass coordinates, x and y , should respectively be a GHS of the following line bundles,

$$L_x := \mathcal{O}_X(2S_0) \otimes K_B^{-2} = \mathcal{O}_X(-2, 2, 6), \quad (5.25)$$

$$L_y := \mathcal{O}_X(3S_0) \otimes K_B^{-3} = \mathcal{O}_X(-3, 3, 9). \quad (5.26)$$

In order to obtain the explicit Weierstrass model, we first need to construct an expression for the three Weierstrass coordinates, z , x , and y , in terms of the homogeneous

still generically contains \mathbf{x}_2 -dependence. For denominators valued in \mathbb{P}^n with $n > 1$, in general the form of the rational function is more complicated, with a precise form that is in principle determined by the Koszul sequence associated to the line bundle in question.

coordinates, \mathbf{x}_1 , \mathbf{x}_2 and \mathbf{x}_3 . We have already obtained an explicit expression for z in the previous subsection. One can similarly choose a quadratic and a cubic polynomial in \mathbf{x}_1 for the denominators of global sections of L_x and L_y , respectively and tune appropriate degree numerators to construct their GHS's. In this way, we obtain explicit rational expressions,

$$z = z(\mathbf{x}_1, \mathbf{x}_2, \mathbf{x}_3) \in \Gamma(X, L_z), \tag{5.27}$$

$$\tilde{x} = \tilde{x}(\mathbf{x}_1, \mathbf{x}_2, \mathbf{x}_3) \in \Gamma(X, L_x), \tag{5.28}$$

$$\tilde{y} = \tilde{y}(\mathbf{x}_1, \mathbf{x}_2, \mathbf{x}_3) \in \Gamma(X, L_y), \tag{5.29}$$

where z is uniquely fixed up to an overall constant and \tilde{x} and \tilde{y} are expressed as a linear combination of 6 and 11 independent rational expressions, respectively. Note that the tilded variables correspond to general sections of the appropriate line bundles. The specific sections corresponding to Weierstrass coordinates will then simply be denoted as x and y , respectively. One can independently compute the dimensions of these line bundles as

$$h^\bullet(X, L_z) = (1, 0, 0), \quad h^\bullet(X, L_x) = (6, 0, 0), \quad h^\bullet(X, L_y) = (11, 0, 0), \tag{5.30}$$

and hence confirm that a complete basis has been obtained for each space of GHS's.

We are now ready to construct the Weierstrass model via the procedure described in subsection 3.1. Before we start off, however, let us first convince ourselves, via cohomology dimension count, that the procedure will work. In the one-dimensional space, $H^0(X, L_z)$, we find the unique z . Next we consider $H^0(X, L_x)$. It has $h^0(X, L_x) = 6$ independent GHS's, 5 of which are of the form,

$$t^{(4)}(\mathbf{x}_3)z(\mathbf{x}_1, \mathbf{x}_2, \mathbf{x}_3), \tag{5.31}$$

with $t^{(4)} \in \Gamma(B, K_B^{-2}) = \Gamma(\mathbb{P}^1, \mathcal{O}_{\mathbb{P}^1}(4))$. The procedure tells us that there are GHS's in $\Gamma(X, L_x)$, which cannot be written in the form (5.31), and such a GHS can be obtained as a generic linear combination of the 6 basis elements for $\Gamma(X, L_x)$. We denote that choice by x . Similarly, we consider $H^0(X, L_y)$. It has $h^0(X, L_y) = 11$ independent GHS's, 10 of which are of the form,

$$t^{(6)}(\mathbf{x}_3)z(\mathbf{x}_1, \mathbf{x}_2, \mathbf{x}_3)^3, \quad t^{(2)}(\mathbf{x}_3)z(\mathbf{x}_1, \mathbf{x}_2, \mathbf{x}_3)x(\mathbf{x}_1, \mathbf{x}_2, \mathbf{x}_3), \tag{5.32}$$

where $t^{(6)} \in \Gamma(\mathbb{P}^1, \mathcal{O}_{\mathbb{P}^1}(6))$ and $t^{(2)} \in \Gamma(\mathbb{P}^1, \mathcal{O}_{\mathbb{P}^1}(2))$. Again, there exists GHS's in $\Gamma(X, L_y)$, which cannot be written in the form (5.32), and such a GHS can be obtained as a generic linear combination of the 11 basis elements for $\Gamma(X, L_y)$. We denote that choice by y . Finally, let us consider $H^0(X, \mathcal{O}_X(-6, 6, 18))$, which turns out to be 38 dimensional. Making use of z , x , and y , one can construct a total of 39 GHS's as follows:

$$y^2, \quad t^{(2)}yxz, \quad t^{(6)}yz^3, \quad x^3, \quad t^{(4)}x^2z^2, \quad t^{(8)}xz^4, \quad t^{(12)}z^6, \tag{5.33}$$

where $t^{(k)} \in \Gamma(\mathbb{P}_{\mathbf{x}_3}^1, \mathcal{O}_{\mathbb{P}_{\mathbf{x}_3}^1}(k))$. Therefore, there must be a linear relation among these 39 GHS's. This can be found by writing down a generic linear combination of the GHS's with unspecified coefficients. We then substitute a number of points on X into this combination

and constrain the coefficients such that the linear combination is zero on each point. If this procedure is repeated for enough points then the coefficients will be completely specified up to an overall scale. Upon an appropriate choice of that overall scale, one thus obtains the following Tate form,

$$y^2 + a_1xyz + a_3yz^3 = x^3 + a_2x^2z^2 + a_4xz^4 + a_6z^6. \quad (5.34)$$

In the case at hand, we obtain the explicit Tate coefficients below.

$$\begin{aligned} a_1 &= (251 - 435i)x_{3,0}^2 - (479 - 829i)x_{3,0}x_{3,1} + (143 - 247i)x_{3,1}^2, \\ a_2 &= (29600 + 51200i)x_{3,0}^4 + (-113000 - 195000i)x_{3,0}^3x_{3,1} + (150000 + 260000i)x_{3,0}^2x_{3,1}^2 \\ &\quad + (-59100 - 102000i)x_{3,0}x_{3,1}^3 + (8490 + 14700i)x_{3,1}^4, \\ a_3 &= -4250x_{3,0}^6 - 24500x_{3,0}^5x_{3,1} + 3.13 \times 10^6 x_{3,0}^4x_{3,1}^2 - 7.25 \times 10^6 x_{3,0}^3x_{3,1}^3 + 6.28 \times 10^6 x_{3,0}^2x_{3,1}^4 \\ &\quad - 4.33 \times 10^6 x_{3,0}x_{3,1}^5 + 1.23 \times 10^6 x_{3,1}^6, \\ a_4 &= (352000 - 609000i)x_{3,0}^8 + (973000 - 1.69 \times 10^6 i)x_{3,0}^7x_{3,1} \\ &\quad + (-3.49 \times 10^8 + 6.05 \times 10^8 i)x_{3,0}^6x_{3,1}^2 \\ &\quad + (1.46 \times 10^9 - 2.52 \times 10^9 i)x_{3,0}^5x_{3,1}^3 + (-2.44 \times 10^9 + 4.23 \times 10^9 i)x_{3,0}^4x_{3,1}^4 \\ &\quad + (2.11 \times 10^9 - 3.65 \times 10^9 i)x_{3,0}^3x_{3,1}^5 + (-1.30 \times 10^9 + 2.25 \times 10^9 i)x_{3,0}^2x_{3,1}^6 \\ &\quad + (4.26 \times 10^8 - 7.37 \times 10^8 i)x_{3,0}x_{3,1}^7 + (-4.63 \times 10^7 + 8.01 \times 10^7 i)x_{3,1}^8, \\ a_6 &= 3.61 \times 10^6 x_{3,0}^{12} + 9.68 \times 10^7 x_{3,0}^{11}x_{3,1} + 2.91 \times 10^9 x_{3,0}^{10}x_{3,1}^2 \\ &\quad - 6.86 \times 10^9 x_{3,0}^9x_{3,1}^3 - 1.85 \times 10^{12} x_{3,0}^8x_{3,1}^4 \\ &\quad + 8.48 \times 10^{12} x_{3,0}^7x_{3,1}^5 - 1.71 \times 10^{13} x_{3,0}^6x_{3,1}^6 + 2.08 \times 10^{13} x_{3,0}^5x_{3,1}^7 - 1.87 \times 10^{13} x_{3,0}^4x_{3,1}^8 \\ &\quad + 1.16 \times 10^{13} x_{3,0}^3x_{3,1}^9 - 5.20 \times 10^{12} x_{3,0}^2x_{3,1}^{10} + 1.50 \times 10^{12} x_{3,0}x_{3,1}^{11} - 1.62 \times 10^{11} x_{3,1}^{12}. \end{aligned}$$

Note that here we have only reproduced the numerical coefficients in these expressions to three significant figures in order to keep the equations of a manageable size.

Via a reparameterization it is possible to obtain the Weierstrass form,

$$y^2 = x^3 + f_Wxz^4 + g_Wz^6, \quad (5.35)$$

where $f_W \in \Gamma(B, K_B^{-4}) = \Gamma(\mathbb{P}^1, \mathcal{O}_{\mathbb{P}^1}(8))$ and $g_W \in \Gamma(B, K_B^{-6}) = \Gamma(\mathbb{P}^1, \mathcal{O}_{\mathbb{P}^1}(12))$. The discriminant locus is then located at the vanishing of $\Delta_W := 4f_W^3 + 27g_W^2$, which is a homogeneous polynomial of degree 24 in $\mathbf{x}_3 = (x_{3,0} : x_{3,1})$ and which has 24 distinct roots. We confirm that at these 24 points, f_W and g_W do not vanish and Δ_W vanishes to order 1. Therefore, the singular fibers are of type I_1 from the Kodaira's classification.

5.2.2 Jacobian

The configuration (5.2) describing the fiber in this example represents a complete intersection in the toric variety $\mathcal{A}_F = \mathbb{P}_{\mathbf{x}_1}^1 \times \mathbb{P}_{\mathbf{x}_2}^2$ of a bi-linear and a multi-degree (1, 2) polynomials. This codimension-two fiber has the PALP ID (4, 0), which, via the results of ref. [30], can straightforwardly be transformed into the Jacobian as follows. The defining equations can be viewed as a function of the fiber coordinates, \mathbf{x}_1 and \mathbf{x}_2 , with the base coordinates

\mathbf{x}_3 being demoted to parameters. Then, the various monomial coefficients in the defining equations are expressed in terms of the base coordinates \mathbf{x}_3 . The expressions for f and g of the Jacobian in terms of those coefficients are found from ref. [30] with the PALP ID (4,0), and thereby one obtains the Jacobian in the form,

$$y^2 = x^3 + f_J x z^4 + g_J z^6, \tag{5.36}$$

where $f_J \in \Gamma(B, K_B^{-4})$ and $g_J \in \Gamma(B, K_B^{-6})$. The discriminant $\Delta_J := 4f_J^3 + 27g_J^2$ of the Jacobian is a homogeneous polynomial of degree 24. We confirm explicitly that this discriminant agrees with that obtained from the Weierstrass equation derived in the previous subsection:

$$\Delta_W \sim \Delta_J. \tag{5.37}$$

5.2.3 Resolved geometry

Finally, we may find the discriminant locus at the level of the smooth geometry, given by eqs. (5.3) and (5.4). As described in subsection 3.3, the fiber is singular at points in X that obeys eqs. (3.24). In this case these equations are the two fiber-defining relations,

$$P_1(\mathbf{x}_1, \mathbf{x}_2, \mathbf{x}_3) = 0, \quad P_2(\mathbf{x}_1, \mathbf{x}_2, \mathbf{x}_3) = 0, \tag{5.38}$$

together with the condition for the degeneration of the normal form,

$$d_F P_1 \wedge d_F P_2 = 0. \tag{5.39}$$

Eq. (5.39) gives rise to three polynomial equations as the fiber is embedded in a threefold $\mathbb{P}_{\mathbf{x}_1}^1 \times \mathbb{P}_{\mathbf{x}_2}^2$, along which the exterior derivatives, d_F , are taken. Given this system of five polynomials in $\mathbf{x}_1, \mathbf{x}_2$, and \mathbf{x}_3 , we can immediately eliminate the fiber coordinates to locate the discriminant locus in the base, using Gröbner basis techniques. As a consequence, we obtain a single degree-24 polynomial, $\Delta_{\text{res}}(\mathbf{x}_3)$, which can easily be seen to agree with the other two discriminant equations that we have already found:

$$\Delta_{\text{res}} \sim \Delta_W \sim \Delta_J. \tag{5.40}$$

5.3 Arithmetic of the sections

The sections to the elliptic fibration form an additive group. Let us take $\mathcal{O}_X(-1, 1, 1)$ as the line bundle associated to the zero section, σ_0 , and $\mathcal{O}_X(2, -1, 4)$ as that associated to the generator section, σ_g , of the MW group. Note that with this choice, σ_g could potentially only generate a subgroup of the MW group. However, we will see shortly that this is not the case in this example. Denoting their divisor classes by S_0 and S_g , the divisor classes of $k\sigma_g$ are given as [67, 72]

$$\text{Div}(k\sigma_g) = S_0 + k(S_g - S_0) + k(k-1) \pi(S_0 \cdot (S_g - S_0)), \tag{5.41}$$

$$= (-1 + 3k)J_1 + (1 - 2k)J_2 + (1 - 11k + 14k^2)J_3. \tag{5.42}$$

This reproduces all of the putative sections in eq. (5.9). Note that the addition law in eq. (5.41) is guaranteed to give true sections corresponding to multiples of σ_g as long as

the two divisors σ_0 and σ_g are indeed both sections themselves. Since S_0 and S_g have both been proven to be sections, we find that each of the putative sections in eq. (5.9) is a section too, and that these sections exhaust the MW group.

This in particular proves that the MW group is of rank 1. The consistency of this result can be checked by analyzing the decomposition of the rank-three Picard lattice of X . According to the Tate-Shioda-Wazir theorem, the Picard lattice of X is comprised of linearly independent basis elements of the following four types:

- (1) Base divisor classes: $B = \mathbb{P}^1$ has a unique generator for the base divisor classes in this example.
- (2) Fibral divisors: as was shown by an analysis of the discriminant in the previous subsection, it turns out that there are no enhancements of fiber singularities at codimension one in this example and hence, blow-up divisors do not exist.
- (3) The zero section: $\mathcal{O}_X(-1, 1, 1)$ has proven to lead to a (holomorphic) section, σ_0 , which we may take as the zero section.
- (4) Rational sections: the arithmetic (5.41) of putative sections shows that the MW group has rank 1 and is generated by σ_g , associated with the line bundle $\mathcal{O}_X(2, -1, 4)$.

We find then, that the contributions from these four different types of generators indeed give rise to the Picard lattice of rank $1 + 0 + 1 + 1 = 3$, as desired in this example.

6 Example 2: a threefold example

We now consider an elliptic Calabi-Yau threefold example, X_3 with the configuration,

$$X_3 = \left[\begin{array}{c|cc} \mathbb{P}_{\mathbf{x}_1}^1 & 1 & 1 \\ \mathbb{P}_{\mathbf{x}_2}^2 & 1 & 2 \\ \mathbb{P}_{\mathbf{x}_3}^1 & 1 & 1 \\ \mathbb{P}_{\mathbf{x}_4}^1 & 1 & 1 \end{array} \right]. \tag{6.1}$$

This threefold can be obtained by fibering the $K3$ surface in the previous section over $\mathbb{P}_{\mathbf{x}_4}^1$. This corresponds to the CICY #7675 in the list of complete intersection Calabi-Yau threefolds [78] and has the Hodge numbers $(h^{1,1}, h^{2,1}) = (4, 50)$. From the configuration matrix, an obvious genus-one fibration structure with the base $B_2 = \mathbb{P}_{\mathbf{x}_3}^1 \times \mathbb{P}_{\mathbf{x}_4}^1$ can be seen, where the fiber has the same configuration as eq. (5.2),

$$F = \left[\begin{array}{c|cc} \mathbb{P}_{\mathbf{x}_1}^1 & 1 & 1 \\ \mathbb{P}_{\mathbf{x}_2}^2 & 1 & 2 \end{array} \right]. \tag{6.2}$$

We start by fixing a generic complex structure, for which the threefold is smooth. Instead of writing out the long expressions for the two defining equations, we will provide the list of coefficients for a given monomial ordering. Firstly, the monomial GHS's of $\mathcal{O}_{\mathcal{A}}(d_1, \dots, d_4)$

are given the lexicographic ordering applied to the exponent list. For instance, to compare the two monomial GHS's, $m_1 = x_{1,0}x_{2,1}^2x_{3,0}x_{4,0}$ and $m_2 = x_{1,0}x_{2,0}x_{2,2}x_{3,1}x_{4,1}$, of $\mathcal{O}_{\mathcal{A}}(1, 2, 1, 1)$, we first read their exponents as a list of degree vectors,

$$m_1 = \mathbf{x}_1^{(1,0)} \mathbf{x}_2^{(0,2,0)} \mathbf{x}_3^{(1,0)} \mathbf{x}_4^{(1,0)} \sim \{(1, 0), (0, 2, 0), (1, 0), (1, 0)\}, \quad (6.3)$$

$$m_2 = \mathbf{x}_1^{(1,0)} \mathbf{x}_2^{(1,0,1)} \mathbf{x}_3^{(0,1)} \mathbf{x}_4^{(0,1)} \sim \{(1, 0), (1, 0, 1), (0, 1), (0, 1)\}, \quad (6.4)$$

then the exponent order leads to $m_2 > m_1$ and hence, m_2 will come earlier. With respect to this ordering for the basis monomials, our choice of the two defining equations,

$$P_1 \in \Gamma(\mathcal{A}, \mathcal{O}_{\mathcal{A}}(1, 1, 1, 1)), \quad \text{with } h^0(\mathcal{A}, \mathcal{O}_{\mathcal{A}}(1, 1, 1, 1)) = 24, \quad (6.5)$$

$$P_2 \in \Gamma(\mathcal{A}, \mathcal{O}_{\mathcal{A}}(1, 2, 1, 1)), \quad \text{with } h^0(\mathcal{A}, \mathcal{O}_{\mathcal{A}}(1, 2, 1, 1)) = 48, \quad (6.6)$$

are specified by the two lists of the monomial coefficients,

$$c_{P_1} = \{45, 4, 17, 7, 35, 69, 82, 43, 44, 46, 19, 95, 74, 100, 36, 71, 4, 43, 95, 2, 36, 9, 50, 28\}, \quad (6.7)$$

$$c_{P_2} = \{97, 60, 78, 22, 11, 27, 36, 15, 12, 53, 65, 81, 64, 12, 81, 70, 17, 9, 12, 89, 91, 43, 38, 43, 6, 99, 66, 3, 58, 26, 62, 62, 43, 23, 76, 6, 40, 63, 2, 85, 42, 80, 63, 37, 75, 33, 74, 67\}. \quad (6.8)$$

6.1 Section analysis

Putative sections. We begin by writing down the criteria that a putative section S , labelled by $\mathcal{O}_X(S) = \mathcal{O}_X(b_1, b_2, b_3, b_4)$, must satisfy. Firstly, the Oguiso criterion, that S must have a single intersection point with a generic fiber, leads to

$$2b_1 + 3b_2 \stackrel{!}{=} 1. \quad (6.9)$$

The base topology criteria (2.12), one for each base divisor, $D_\alpha^b \sim J_3, J_4$, are

$$6b_1b_2 + 2b_2^2 + 4b_1b_\alpha + 6b_2b_\alpha \stackrel{!}{=} -2, \quad \text{for } \alpha = 3, 4, \quad (6.10)$$

where eq. (6.9) has been used in simplifying the result. These three equations (6.9) and (6.10) then lead to the one-parameter family of divisor classes,

$$b_1 = -1 - 3k, \quad b_2 = 1 + 2k, \quad b_3 = 1 + 11k + 14k^2, \quad b_4 = 1 + 11k + 14k^2, \quad (6.11)$$

with an integer parameter $k \in \mathbb{Z}$. This family gives us, for the values $k = \{1, 0, -1, -2\}$, the putative sections, $\mathcal{O}_X(-4, 3, 26, 26)$, $\mathcal{O}_X(-1, 1, 1, 1)$, $\mathcal{O}_X(2, -1, 4, 4)$ and $\mathcal{O}_X(5, -3, 35, 35)$, respectively. In what follows, in order to prove that the MW group is of rank 1, we will provide an explicit description of the associated section maps for the simplest two putative sections,

$$\mathcal{O}_X(-1, 1, 1, 1), \quad \text{with } h^\bullet(X, \mathcal{O}_X(-1, 1, 1, 1)) = (1, 0, 0, 0), \quad (6.12)$$

$$\mathcal{O}_X(2, -1, 4, 4), \quad \text{with } h^\bullet(X, \mathcal{O}_X(2, -1, 4, 4)) = (1, 6, 0, 0). \quad (6.13)$$

Explicit expressions for the sections. Let us start with the simplest putative section, $\mathcal{O}_X(-1, 1, 1, 1)$. We take the ansatz,

$$z = \frac{N(\mathbf{x}_2, \mathbf{x}_3, \mathbf{x}_4)}{D(\mathbf{x}_1)}, \tag{6.14}$$

where $N \in \Gamma(X, \mathcal{O}_X(0, 1, 1, 1))$ and $D \in \Gamma(X, \mathcal{O}_X(1, 0, 0, 0))$. As a specific example, we may choose $D = x_{1,0} - x_{1,1}$. Then, the numerator polynomial N has to be fixed up to scaling so that z can be globally holomorphic. Organizing the first defining relation, P_1 as

$$P_1 = x_{1,0} p_0(\mathbf{x}_2, \mathbf{x}_3, \mathbf{x}_4) + x_{1,1} p_1(\mathbf{x}_2, \mathbf{x}_3, \mathbf{x}_4), \tag{6.15}$$

where p_0 and p_1 are tri-linear polynomials, we see that on the divisor of vanishing denominator, $D(\mathbf{x}_1) = x_{1,0} - x_{1,1} = 0$, the tri-linear polynomial $p_0 + p_1$ must vanish. Therefore, upon choosing

$$\begin{aligned} N &:= p_0 + p_1 && (6.16) \\ &= 119x_{2,0}x_{3,0}x_{4,0} + 39x_{2,1}x_{3,0}x_{4,0} + 80x_{2,2}x_{3,0}x_{4,0} + 53x_{2,0}x_{3,1}x_{4,0} + 177x_{2,1}x_{3,1}x_{4,0} \\ &\quad + 69x_{2,2}x_{3,1}x_{4,0} + 104x_{2,0}x_{3,0}x_{4,1} + 112x_{2,1}x_{3,0}x_{4,1} + 55x_{2,2}x_{3,0}x_{4,1} + 78x_{2,0}x_{3,1}x_{4,1} \\ &\quad + 45x_{2,1}x_{3,1}x_{4,1} + 123x_{2,2}x_{3,1}x_{4,1}, \end{aligned}$$

we obtain the desired GHS, the zero locus of which can easily be seen to have no singularities via the patch-wise analysis developed in ref. [25] (note that in the above we have utilized the specific choice of complex structure that we have made in this case).

Having specified $z = N/D$, we can now try to explicitly parameterize the zero locus in terms of the base coordinates $\mathbf{x}_3, \mathbf{x}_4$, resulting in the following generic parametrization,

$$\begin{aligned} x_{1,i} &= A_{1,i}(\mathbf{x}_3, \mathbf{x}_4) \in \Gamma(B, \mathcal{O}_B(5, 5)), \quad \text{for } i = 0, 1, \\ x_{2,i} &= A_{2,i}(\mathbf{x}_3, \mathbf{x}_4) \in \Gamma(B, \mathcal{O}_B(2, 2)), \quad \text{for } i = 0, 1, 2, \end{aligned} \tag{6.17}$$

which we specify by their coefficient list for the basis monomials, as given in appendix B.1. Unlike for the $K3$ case, this section map is not holomorphic and is a rational section to the elliptic fibration. In particular, there arise 32 points on the base $\mathbb{P}_{\mathbf{x}_3}^1 \times \mathbb{P}_{\mathbf{x}_4}^1$ where the map is ill-defined for the $\mathbb{P}_{\mathbf{x}_1}^1$ direction and 6 points where it is ill-defined for the $\mathbb{P}_{\mathbf{x}_2}^2$ direction; the latter 6 points turn out to entirely belong to the former set. More specifically, at each of these base points, eq. (6.17) returns a vanishing value for all of the homogeneous coordinates of the \mathbb{P}^n 's mentioned above. Note that the map eq. (6.17) gives a single point on the fiber over all of the base except for these 32 points. The issue here is that we have naively tried to solve the system given by setting $z = 0$ to find a unique value of the fiber ambient coordinates as a function of the variables in the base. At points on the base where a rational section has vertical components such a procedure can give spurious results. Let us elaborate upon this further.

Upon investigating each of the problematic base points, we observe two qualitatively different behaviors: over the 6 base points where the rational map is ill-defined along both

$\mathbb{P}_{\mathbf{x}_1}^1$ and $\mathbb{P}_{\mathbf{x}_2}^2$ directions, the rational section wraps a \mathbb{P}^1 with the following configuration matrix,

$$\left[\begin{array}{c|cc} \mathbb{P}_{\mathbf{x}_1}^1 & 1 & 0 \\ \mathbb{P}_{\mathbf{x}_2}^2 & 2 & 1 \end{array} \right], \tag{6.18}$$

while over the other 26 points it wraps a copy of $\mathbb{P}_{\mathbf{x}_1}^1$. This explicitly shows that the rational section is a blowup of the base, $B = \mathbb{P}_{\mathbf{x}_3}^1 \times \mathbb{P}_{\mathbf{x}_4}^1$ at 32 points, which conforms with their Euler number difference,

$$\chi(S) - \chi(B) = 36 - 4 = 32. \tag{6.19}$$

Another observation related to the above structure is as follows. Upon blowing down the smooth Calabi-Yau threefold X via the contraction of the first row of its configuration matrix, one obtains a hypersurface of multi-degree $(3, 2, 2)$ in $\mathbb{P}_{\mathbf{x}_2}^2 \times \mathbb{P}_{\mathbf{x}_3}^1 \times \mathbb{P}_{\mathbf{x}_4}^1$, which is a deformation of a smooth Calabi-Yau threefold with Euler number -144 , while $\chi(X)$ is computed to be -92 . The difference of these two Euler numbers is $52 = 2 \times 26$. Thus the ambient \mathbb{P}^1 that we are blowing down is associated with 26 \mathbb{P}^1 's in the Calabi-Yau geometry. This conforms with the fact that a degeneration associated with this \mathbb{P}^1 in the section map was associated with exactly 26 \mathbb{P}^1 's above specific points in the base.

Let us now move on to analyzing the second putative section, $\mathcal{O}_X(2, -1, 4, 4)$. In order to find the explicit map, we first need to find the rational expression for the GHS of this line bundle. With the negative degree in the $\mathbb{P}_{\mathbf{x}_2}^2$, however, the naive ansatz with the linear denominator in \mathbf{x}_2 does not work. It turns out that the GHS can be found once we shift both the numerator and the denominator multi-degrees by $(0, 2, 0, 0)$, i.e., with the modified ansatz,

$$z = \frac{N(\mathbf{x}_1, \mathbf{x}_2, \mathbf{x}_3, \mathbf{x}_4)}{D(\mathbf{x}_2)}, \tag{6.20}$$

where $N \in \Gamma(X, \mathcal{O}_X(2, 2, 4, 4))$ and $D \in \Gamma(X, \mathcal{O}_X(0, 3, 0, 0))$. As a choice of denominator, we may take $D(\mathbf{x}_2) = x_{2,0}x_{2,1}x_{2,2}$ and tune the 450 monomial coefficients in N . That is, we demand that N vanishes on each of the three divisors $\{x_{2,i} = 0\} \subset X$, in practice employing the same techniques as we have in previous sections. Then, only 129 coefficients of the 450 remain free. This does not guarantee that the rational form tuned as such is regular yet, since the denominator may vanish to order 2 on the loci in X with $\mathbf{x}_2 = (1 : 0 : 0)$, $(0 : 1 : 0)$, or $(0 : 0 : 1)$. Interestingly, we find that N also vanishes to order 2 on such loci. We therefore conclude that the tuned form is indeed regular and corresponds to GHS's. This may at first sound strange since we know that $h^0(X, \mathcal{O}_X(2, -1, 4, 4)) = 1$. The only way to make sense of this result is that the 129-parameter expression we have obtained for the GHS's of $\mathcal{O}_X(2, -1, 4, 4)$ should only span a one-dimensional vector space, and it indeed turns out that in the coordinate ring of X they all lead to one and the same GHS up to scaling.

Having specified $z = N/D$ for this section, we can now proceed to find a generic parametrization of its zero locus in terms of the base coordinates $\mathbf{x}_3, \mathbf{x}_4$, resulting in the

parametric expression of the form,

$$x_{1,i} = A'_{1,i}(\mathbf{x}_3, \mathbf{x}_4) \in \Gamma(B, \mathcal{O}_B(5, 5)), \quad \text{for } i = 0, 1, \quad (6.21)$$

$$x_{2,i} = A'_{2,i}(\mathbf{x}_3, \mathbf{x}_4) \in \Gamma(B, \mathcal{O}_B(16, 16)), \quad \text{for } i = 0, 1, 2. \quad (6.22)$$

Here $A'_{1,i}$ coincide with the $A_{1,i}$ from eq. (6.17) and $A'_{2,i}$ are some fixed polynomials, which we do not display in this paper (each of the 289 integer coefficients has 27 to 37 digits). Equipped with such an explicit map, one can show that this defines a rational section to the elliptic fibration. The map eq. (6.21) is ill-defined in the same manner as for the previous GHS in this section for the $\mathbb{P}^1_{\mathbf{x}_3}$ direction at 32 points on the base $\mathbb{P}^1_{\mathbf{x}_3} \times \mathbb{P}^1_{\mathbf{x}_4}$ and for the $\mathbb{P}^2_{\mathbf{x}_2}$ direction at 92 points on the base, the former 32 points entirely belonging to the latter set. As before, this corresponds to the GHS giving rise to a rational section wrapping \mathbb{P}^1 's in the fiber over some points on the base.

6.2 Locating the singular fibers

All the steps in this subsection are a straightforward analogue of those in subsection 5.2. We will thus be brief here and will mainly state the results of the analysis.

6.2.1 Weierstrass model

Having proven that the unique GHS of the line bundle,

$$L_z := \mathcal{O}_X(S_0) = \mathcal{O}_X(-1, 1, 1, 1), \quad (6.23)$$

is a rational section to the elliptic fibration, we choose to use L_z as the bundle, which the Weierstrass coordinate z is a GHS of, and take

$$L_x := \mathcal{O}_X(2S_0) \otimes K_B^{-2} = \mathcal{O}_X(-2, 2, 6, 6), \quad (6.24)$$

$$L_y := \mathcal{O}_X(3S_0) \otimes K_B^{-3} = \mathcal{O}_X(-3, 3, 9, 9), \quad (6.25)$$

to which the other Weierstrass coordinates, x and y , are associated.

We first need the explicit expressions for the three Weierstrass coordinates z , x and y in terms of the homogeneous coordinates $\mathbf{x}_{r=1,\dots,4}$. We have already performed this computation for z in the previous subsection. Using similar techniques to those outlined already we can choose a quadratic and a cubic polynomial in \mathbf{x}_1 for the denominator and tune the numerator, to construct the other GHS's, x and y , respectively. In this manner, we obtain explicit rational expressions,

$$z = z(\mathbf{x}_1, \mathbf{x}_2, \mathbf{x}_3, \mathbf{x}_4) \in \Gamma(X, L_z), \quad (6.26)$$

$$x = x(\mathbf{x}_1, \mathbf{x}_2, \mathbf{x}_3, \mathbf{x}_4) \in \Gamma(X, L_x), \quad (6.27)$$

$$y = y(\mathbf{x}_1, \mathbf{x}_2, \mathbf{x}_3, \mathbf{x}_4) \in \Gamma(X, L_y), \quad (6.28)$$

where z is uniquely fixed up to an overall constant, while x and y are expressed as a linear combination of 26 and 59 independent rational expressions, respectively. One can independently compute the cohomology dimensions of these line bundles as

$$h^\bullet(X, L_z) = (1, 0, 0, 0), \quad h^\bullet(X, L_x) = (26, 32, 0, 0), \quad h^\bullet(X, L_y) = (59, 128, 0, 0), \quad (6.29)$$

and hence confirm that a complete basis has been obtained for each space of GHS's.

We are then ready to construct the Weierstrass model via the procedure of subsection 3.1. The cohomology dimension count works out as in previous examples, with a GHS of L_x not being proportional to z^2 and a GHS of L_y not being constructed from terms proportional to z^3 or xz . The Weierstrass relation itself is a section of $L_w := L_y^{\otimes 2}$ whose zeroth cohomology is one less in dimension than the space that can be spanned by monomials in x, y, z and the base coordinates of the correct degree, as expected.

As in the $K3$ case, one can find a unique linear relation among the 335 GHS's of L_w , using the same techniques that have been employed in that case. Upon an appropriate rescaling, one thus obtains the Tate form,

$$y^2 + a_1xyz + a_3yz^3 = x^3 + a_2x^2z^2 + a_4xz^4 + a_6z^6, \tag{6.30}$$

with the explicit Tate coefficients a_i , and via the reparameterization, gets to the Weierstrass form,

$$y^2 = x^3 + f_Wxz^4 + g_Wz^6, \tag{6.31}$$

where $f_W \in \Gamma(B, K_B^{-4}) = \Gamma(\mathbb{P}^1 \times \mathbb{P}^1, \mathcal{O}_{\mathbb{P}^1}(8, 8))$ and $g_W \in \Gamma(B, K_B^{-6}) = \Gamma(\mathbb{P}^1 \times \mathbb{P}^1, \mathcal{O}_{\mathbb{P}^1}(12, 12))$. The discriminant locus is then located at the vanishing of $\Delta_W := 4f_W^3 + 27g_W^2$, which is a homogeneous polynomial of multi-degree $(24, 24)$ in $(\mathbf{x}_3, \mathbf{x}_4)$. The expression for the Tate form and the discriminant are rather large and so we do not reproduce them here. The only explicit information we will require going forward with this example is that, given the discriminant polynomial Δ_W , we confirm that all the codimension-one singularities are of type I_1 .

6.2.2 Jacobian

The configuration (6.2) of the elliptic fiber is the same as the $K3$ case and is a complete intersection in the toric variety $\mathcal{A}_F = \mathbb{P}_{\mathbf{x}_1}^1 \times \mathbb{P}_{\mathbf{x}_2}^2$, with the PALP ID $(4, 0)$. Via the results of ref. [30], it can straightforwardly be transformed into the Jacobian as follows. The defining equations can be viewed as a function of the fiber coordinates, \mathbf{x}_1 and \mathbf{x}_2 , with the base coordinates \mathbf{x}_3 and \mathbf{x}_4 demoted to parameters. Then, the various monomial coefficients in the defining equations are expressed in terms of the base coordinates \mathbf{x}_3 and \mathbf{x}_4 . Using the expressions for f and g of the Jacobian in terms of those coefficients, one obtains the Jacobian in the form,

$$y^2 = x^3 + f_Jxz^4 + g_Jz^6, \tag{6.32}$$

where $f_J \in \Gamma(B, K_B^{-4})$ and $g_J \in \Gamma(B, K_B^{-6})$. The discriminant $\Delta_J := 4f_J^3 + 27g_J^2$ of the Jacobian is a homogeneous polynomial of bi-degree $(24, 24)$ and we confirm that it agrees with the Weierstrass form:

$$\Delta_W \sim \Delta_J. \tag{6.33}$$

6.2.3 Resolved geometry

Finally, we can explore the discriminant locus at the level of the smooth geometry associated with the choice of complex structure in eq. (6.7). Elliptic fibers admit a singularity at

points in X obeying eqs. (3.24), which, much like the $K3$ case, are the defining relations for the fiber,

$$P_1 = 0, \quad P_2 = 0, \tag{6.34}$$

together with the degeneration condition,

$$d_F P_1 \wedge d_F P_2 = 0, \tag{6.35}$$

where the latter gives rise to three polynomial equations as the fiber is embedded in a threefold $\mathbb{P}_{x_1}^1 \times \mathbb{P}_{x_2}^2$. Given this system of five polynomials in \mathbf{x}_r , we could in principle eliminate the fiber coordinates as in the $K3$ case to obtain a single polynomial $\Delta_{\text{res}} \in \Gamma(B, K_B^{-12})$ for the discriminant locus embedded in B . However, the elimination process never finishes in a reasonable amount of time. In order to compare this discriminant locus with the previous ones, we take a numerical analysis based on a set of generic points. That is, first we choose a large number of random points that solve the system (6.34) and (6.35) and substitute them into the previously obtained polynomial, $\Delta_W \sim \Delta_J$ to confirm that it vanishes on all of those points. This shows that the discriminant obtained from the smooth geometry is a subset of that obtained from the Weierstrass and Jacobian forms. Similarly by choosing random points on the base that solve $\Delta_W = 0 (= \Delta_J)$ and substituting them into eqs. (6.34) and (6.35) we can show that the discriminant obtained from the Weierstrass form is a subset of that obtained from the analysis of the resolved geometry. In this manner, we can conclude that

$$\Delta_{\text{res}} \sim \Delta_W \sim \Delta_J. \tag{6.36}$$

6.3 Arithmetic of the sections

We have so far obtained two legitimate sections, associated with $\mathcal{O}_X(-1, 1, 1, 1)$ and $\mathcal{O}_X(2, -1, 4, 4)$, respectively. Let us take $\mathcal{O}_X(-1, 1, 1, 1)$ as the line bundle for a zero section, σ_0 , and $\mathcal{O}_X(2, -1, 4, 4)$ for the generator section, σ_g , of (potentially a subgroup of) the MW group. Denoting their divisor classes by S_0 and S_g , the divisor classes of $k\sigma_g$ are given by applying eq. (4.3) as

$$\begin{aligned} \text{Div}(k\sigma_g) &= S_0 + k(S_g - S_0) + k(k-1) \pi(S_0 \cdot (S_g - S_0)), \tag{6.37} \\ &= (-1 + 3k)J_1 + (1 - 2k)J_2 + (1 - 11k + 14k^2)J_3 + (1 - 11k + 14k^2)J_4, \tag{6.38} \end{aligned}$$

which reproduces all the putative sections in eq. (6.11). Note that the addition law guarantees to give the divisor classes of multiples of σ_g in the MW group. Given that S_0 and S_g have both proven to be a section, we learn that each of the putative sections in eq. (6.11) is a section, too, and that there are no more sections.

As in the previous example, we can check what we have learned about this geometry for consistency by analyzing the splitting of the rank-four Picard lattice of X . The four different types of the divisor classes of X are as follows:

- (1) Base divisor classes: $B = \mathbb{P}^1 \times \mathbb{P}^1$ has two independent generators.
- (2) Fibral divisors: we have seen that there are no enhancement of fiber singularities at codimension one and hence, appropriate vertical divisors do not exist.

- (3) Zero section: $\mathcal{O}_X(1, -1, 1, 1)$ has proven to be a rational section, which we may take as the zero section.
- (4) Rational sections: it has also been shown in the previous paragraph that the MW group has rank 1.

Then, the contributions from these four different types of generators indeed give rise to the correct rank, $2 + 0 + 1 + 1 = 4$, as desired.

7 Example 3: a threefold example with non-abelian symmetry

In the two simple examples we have so far looked at, the discriminant locus is irreducible and the singular fibers are only of type I_1 . For a more non-trivial example with enhanced fiber singularities, in this section, we analyze another threefold case with the configuration,

$$X_3 = \left[\begin{array}{c|cccccc} \mathbb{P}_{\mathbf{x}_1}^1 & 1 & 1 & 0 & 0 & 0 & 0 \\ \mathbb{P}_{\mathbf{x}_2}^1 & 1 & 0 & 1 & 0 & 0 & 0 \\ \mathbb{P}_{\mathbf{x}_3}^2 & 0 & 2 & 0 & 0 & 0 & 1 \\ \mathbb{P}_{\mathbf{x}_4}^3 & 0 & 0 & 1 & 1 & 1 & 1 \\ \mathbb{P}_{\mathbf{x}_5}^1 & 1 & 0 & 0 & 1 & 0 & 0 \\ \mathbb{P}_{\mathbf{x}_6}^1 & 0 & 1 & 0 & 0 & 1 & 0 \end{array} \right]. \tag{7.1}$$

This case corresponds to the CICY #5075 of the CICY threefold list [78] and has the Hodge numbers $(h^{1,1}, h^{2,1}) = (7, 30)$. Note in particular that the complete intersection is not “favorable” in that $h^{1,1}(X) = h^{1,1}(\mathcal{A}) + 1$. However, much of our analysis of the sections and fiber types are insensitive to favorability and we can still apply our techniques to this configuration. There also exists, as with almost all CICY threefolds, a nested Calabi-Yau fibration structure of the form,

$$X_3 \longrightarrow B_2 = \mathbb{P}_{\mathbf{x}_5}^1 \times \mathbb{P}_{\mathbf{x}_6}^1 \longrightarrow B_1 = \mathbb{P}_{\mathbf{x}_6}^1, \tag{7.2}$$

and hence can be explored for the Heterotic/F-theory duality (see ref. [1]). Note that the elliptic fibration has the fiber with the following configuration,

$$F = \left[\begin{array}{c|cccccc} \mathbb{P}_{\mathbf{x}_1}^1 & 1 & 1 & 0 & 0 & 0 & 0 \\ \mathbb{P}_{\mathbf{x}_2}^1 & 1 & 0 & 1 & 0 & 0 & 0 \\ \mathbb{P}_{\mathbf{x}_3}^2 & 0 & 2 & 0 & 0 & 0 & 1 \\ \mathbb{P}_{\mathbf{x}_4}^3 & 0 & 0 & 1 & 1 & 1 & 1 \end{array} \right]. \tag{7.3}$$

We start by fixing a generic complex structure, for which the threefold is smooth. We specify our choice of the six polynomials, P_1 - P_6 , each by a list of monomial coefficients, again,

in the lexicographic order for the monomial exponents that was introduced in section 6:

$$c_{P_1} = \{20, 6, 8, 19, 6, 17, 18, 14\}, \quad (7.4)$$

$$c_{P_2} = \{8, 2, 6, 18, 6, 3, 5, 8, 8, 19, 13, 15, 15, 11, 18, 20, 6, 3, 14, 2, 16, 17, 9, 11\}, \quad (7.5)$$

$$c_{P_3} = \{2, 7, 10, 18, 7, 19, 14, 9\}, \quad (7.6)$$

$$c_{P_4} = \{7, 5, 3, 19, 3, 1, 17, 2\}, \quad (7.7)$$

$$c_{P_5} = \{10, 1, 11, 19, 13, 4, 15, 19\}, \quad (7.8)$$

$$c_{P_6} = \{4, 3, 11, 20, 19, 4, 8, 17, 7, 7, 12, 16\}. \quad (7.9)$$

7.1 Section analysis

Putative sections. A putative section S , which we label as $\mathcal{O}_X(S) = \mathcal{O}_X(b_1, \dots, b_6)$, must have a single topological intersection with a generic fiber, leading to

$$2b_1 + 2b_2 + 3b_3 + 2b_4 \stackrel{!}{=} 1. \quad (7.10)$$

In addition, we have the base topology criteria (2.12), one for each base divisor, $D_\alpha^b \sim J_5, J_6$, which give the following conditions,

$$\begin{aligned} &6b_1b_3 + 6b_2b_3 + 2b_3^2 + 4b_1b_4 + 4b_2b_4 + 8b_3b_4 + 2b_4^2 + 4b_1b_6 + 4b_2b_6 \\ &+ 6b_3b_6 + 4b_4b_6 \stackrel{!}{=} -2, \\ &4b_1b_2 + 8b_1b_3 + 6b_2b_3 + 2b_3^2 + 8b_1b_4 + 4b_2b_4 + 8b_3b_4 + 2b_4^2 + 4b_1b_5 + 4b_2b_5 \\ &+ 6b_3b_5 + 4b_4b_5 \stackrel{!}{=} -2, \end{aligned} \quad (7.11)$$

where eq. (7.10) has been used for a simplification. These three equations (7.10) and (7.11) then lead to the following three-parameter family of divisor classes,

$$\begin{aligned} b_1 &= -1 - 3k_1 - k_2 - k_3, \\ b_2 &= k_2, \\ b_3 &= 1 + 2k_1, \\ b_4 &= k_3, \\ b_5 &= 2 + 16k_1 + 20k_1^2 + 3k_2 + 8k_1k_2 + 2k_2^2 + 4k_3 + 12k_1k_3 + 4k_2k_3 + 3k_3^2, \\ b_6 &= 1 + 11k_1 + 14k_1^2 + k_3 + 4k_1k_3 + k_3^2, \end{aligned} \quad (7.12)$$

with $k_1, k_2, k_3 \in \mathbb{Z}$, which, for some small parameter values, $(k_1, k_2, k_3) = (0, 0, -1), (0, -1, 0), (0, 0, 0)$, for example, give the putative sections, $\mathcal{O}_X(0, 0, 1, -1, 1, 1), \mathcal{O}_X(0, -1, 1, 0, 1, 1), \mathcal{O}_X(-1, 0, 1, 0, 2, 1)$, respectively, all of which turn out to have the line-bundle cohomologies,

$$h^\bullet(X, \mathcal{O}_X(S)) = (1, 0, 0, 0). \quad (7.13)$$

We must now confirm that these putative sections do indeed correspond to genuine sections of the fibration.

Explicit expressions for the sections. We will analyze the simplest putative section associated with $L_z := \mathcal{O}_X(0, 0, 1, -1, 1, 1)$, whose GHS of the form,

$$z = \frac{N(\mathbf{x}_3, \mathbf{x}_5, \mathbf{x}_6)}{D(\mathbf{x}_4)} \in \Gamma(X, L_z), \quad (7.14)$$

is sought for, where $N \in \Gamma(X, \mathcal{O}_X(0, 0, 1, 0, 1, 1))$ and $D \in \Gamma(X, \mathcal{O}_X(0, 0, 0, 1, 0, 0))$. We choose $D = x_{4,0}$ and proceed to determine the numerator polynomial by combining the last three defining equations, P_4 - P_6 . We can write these without loss of generality as,

$$P_4 = \sum_{i=0}^3 x_{4,i} u_i(\mathbf{x}_5), \quad P_5 = \sum_{i=0}^3 x_{4,i} v_i(\mathbf{x}_6), \quad P_6 = \sum_{i=0}^3 x_{4,i} w_i(\mathbf{x}_3), \quad (7.15)$$

where u_i, v_i and w_i are linear in $\mathbf{x}_5, \mathbf{x}_6$ and \mathbf{x}_3 , respectively. Then, on the divisor $\{x_{4,0} = 0\} \subset X$, we have

$$\sum_{i=1}^3 x_{4,i} u_i(\mathbf{x}_5) = 0, \quad \sum_{i=1}^3 x_{4,i} v_i(\mathbf{x}_6) = 0, \quad \sum_{i=1}^3 x_{4,i} w_i(\mathbf{x}_3) = 0. \quad (7.16)$$

Therefore, with the denominator $D = x_{4,0}$, if the numerator polynomial is chosen as

$$N := -u_3 v_2 w_1 + u_2 v_3 w_1 + u_3 v_1 w_2 - u_1 v_3 w_2 - u_2 v_1 w_3 + u_1 v_2 w_3 \quad (7.17)$$

$$\begin{aligned} &= 1154x_{3,0}x_{5,0}x_{6,0} + 534x_{3,1}x_{5,0}x_{6,0} + 568x_{3,2}x_{5,0}x_{6,0} + 1794x_{3,0}x_{5,1}x_{6,0} \\ &+ 1864x_{3,1}x_{5,1}x_{6,0} + 543x_{3,2}x_{5,1}x_{6,0} + 1993x_{3,0}x_{5,0}x_{6,1} + 1319x_{3,1}x_{5,0}x_{6,1} \\ &+ 2395x_{3,2}x_{5,0}x_{6,1} - 2380x_{3,0}x_{5,1}x_{6,1} - 1571x_{3,1}x_{5,1}x_{6,1} - 2887x_{3,2}x_{5,1}x_{6,1}, \end{aligned} \quad (7.18)$$

the rational expression (7.14) for z is the desired GHS, the zero locus of which can easily be seen to have no singularities via a direct calculation. With this expression, we obtain an explicit parametrization of its zero locus of the form,

$$\begin{aligned} x_{1,i} &= A_{1,i}(\mathbf{x}_5, \mathbf{x}_6) \in \Gamma(B, \mathcal{O}_B(2, 3)), & \text{for } i = 0, 1, \\ x_{2,i} &= A_{2,i}(\mathbf{x}_5, \mathbf{x}_6) \in \Gamma(B, \mathcal{O}_B(3, 3)), & \text{for } i = 0, 1, \\ x_{3,i} &= A_{3,i}(\mathbf{x}_5, \mathbf{x}_6) \in \Gamma(B, \mathcal{O}_B(1, 1)), & \text{for } i = 0, 1, 2, \\ x_{4,i} &= A_{4,i}(\mathbf{x}_5, \mathbf{x}_6) \in \Gamma(B, \mathcal{O}_B(4, 4)), & \text{for } i = 0, 1, 2, 3, \end{aligned} \quad (7.19)$$

which we present in appendix B.2 for concreteness. These parameterizations provide a rational section to the elliptic fibration. In particular, the map is ill-defined, in a similar manner to what has been seen in previous sections, for the $\mathbb{P}_{\mathbf{x}_4}^3$ direction at 22 points on the base $\mathbb{P}_{\mathbf{x}_5}^1 \times \mathbb{P}_{\mathbf{x}_6}^1$. Amongst these 22 points, for 15 the map is ill-defined also for the $\mathbb{P}_{\mathbf{x}_3}^2$ direction, for 9 the map is ill-defined also for the $\mathbb{P}_{\mathbf{x}_1}^1$ and the $\mathbb{P}_{\mathbf{x}_3}^2$ directions, and for 1 point the map is ill defined for all four ambient projective-space directions. Note that the Euler number difference,

$$\chi(S) - \chi(B) = 26 - 4 = 22, \quad (7.20)$$

indicates that the rational section wraps a \mathbb{P}^1 in the fiber over these 22 base points.

One can perform similar computations for the other two small putative sections, $\mathcal{O}_X(0, -1, 1, 0, 1, 1)$, $\mathcal{O}_X(-1, 0, 1, 0, 2, 1)$. Interestingly, both of them lead to exactly the same section map as in eqs. (7.19). However, they should be understood as vertical divisors attached to the genuine section and provide the first explicit examples of putative sections failing to be a section, even when the strong cohomology condition in eq. (7.13) holds. For more detailed analysis of this geometry, we refer the readers to appendix A.

7.2 Locating the singular fibers

It is possible to apply all three of the different techniques for analyzing singular fibers that we have discussed in this paper to this case. However, in practice, the procedure to compute the Weierstrass model in larger cases is rather expensive in terms of computational resources. The time cost is also high in analyzing the resolved geometry; eliminating the fiber coordinates from the criterion for the fiber singularity to obtain the hypersurface equation for the discriminant locus in the base B_2 is a Gröbner basis calculation and therefore scales badly with the size of the problem. On the other hand, the Jacobian in this case can be analyzed quickly, upon blowing the geometry down until the codimension of the elliptic fiber becomes less than or equal to 2. Therefore, for this example, we only present our exploration of Jacobian to analyze the singular fibers.

Since the fiber configuration (7.3) has codimension 6, Jacobian of X cannot be obtained by the prescription in refs. [2, 30, 31]. We therefore first blow the fiber geometry down by a chain of contractions along fiber directions as described in section 3.2. Let us consider the following chain of three blow downs along the fiber directions, $\mathbb{P}_{x_1}^1$, $\mathbb{P}_{x_2}^1$ and $\mathbb{P}_{x_4}^3$, in turn:

$$\left[\begin{array}{c|c} \mathbb{P}_{x_1}^1 & 1 \ 1 \ 0 \ 0 \ 0 \ 0 \\ \mathbb{P}_{x_2}^1 & 1 \ 0 \ 1 \ 0 \ 0 \ 0 \\ \mathbb{P}_{x_3}^2 & 0 \ 2 \ 0 \ 0 \ 0 \ 1 \\ \mathbb{P}_{x_4}^3 & 0 \ 0 \ 1 \ 1 \ 1 \ 1 \\ \mathbb{P}_{x_5}^1 & 1 \ 0 \ 0 \ 1 \ 0 \ 0 \\ \mathbb{P}_{x_6}^1 & 0 \ 1 \ 0 \ 0 \ 1 \ 0 \end{array} \right] \rightarrow \left[\begin{array}{c|c} \mathbb{P}_{x_2}^1 & 1 \ 1 \ 0 \ 0 \ 0 \ 0 \\ \mathbb{P}_{x_3}^2 & 2 \ 0 \ 0 \ 0 \ 1 \\ \mathbb{P}_{x_4}^3 & 0 \ 1 \ 1 \ 1 \ 1 \\ \mathbb{P}_{x_5}^1 & 1 \ 0 \ 1 \ 0 \ 0 \\ \mathbb{P}_{x_6}^1 & 1 \ 0 \ 0 \ 1 \ 0 \end{array} \right] \rightarrow \left[\begin{array}{c|c} \mathbb{P}_{x_3}^2 & 2 \ 0 \ 0 \ 1 \\ \mathbb{P}_{x_4}^3 & 1 \ 1 \ 1 \ 1 \\ \mathbb{P}_{x_5}^1 & 1 \ 1 \ 0 \ 0 \\ \mathbb{P}_{x_6}^1 & 1 \ 0 \ 1 \ 0 \end{array} \right] \rightarrow \left[\begin{array}{c|c} \mathbb{P}_{x_3}^2 & 3 \\ \mathbb{P}_{x_5}^1 & 2 \\ \mathbb{P}_{x_6}^1 & 2 \end{array} \right], \quad (7.21)$$

to get the fiber blown down to a cubic hypersurface in $\mathbb{P}_{x_3}^2$, with which we can immediately form the Jacobian of the form,

$$y^2 = x^3 + f_J x z^4 + g_J z^6, \quad (7.22)$$

where $f_J \in \Gamma(B, K_B^{-4})$ and $g_J \in \Gamma(B, K_B^{-6})$. The discriminant $\Delta_J := 4f_J^3 + 27g_J^2$ of the Jacobian is a homogeneous polynomial of bi-degree (24, 24). Upon analyzing the singular fibers of the resolved geometry, without eliminating the fiber coordinates to obtain the discriminant locus in the base, we have numerically compared the two discriminant loci and have shown that they are the same:

$$\Delta_J \sim \Delta_{\text{res}}. \quad (7.23)$$

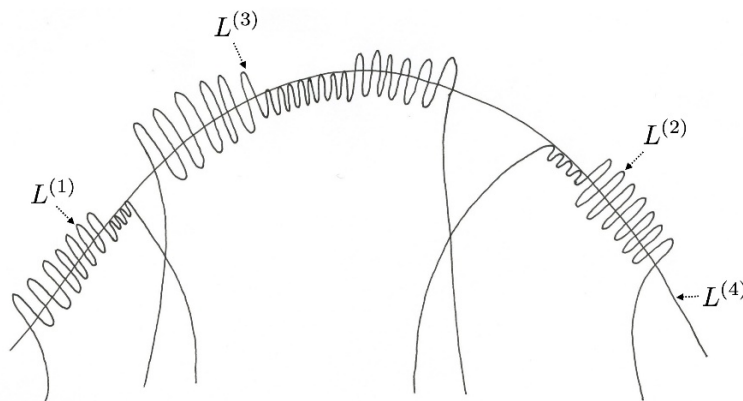


Figure 3. Schematic picture of the discriminant locus, $L \subset B = \mathbb{P}^1_{\mathbf{x}_5} \times \mathbb{P}^1_{\mathbf{x}_6}$, where $\Delta(\mathbf{x}_5, \mathbf{x}_6)$ vanishes. The locus decomposes to four pieces $L^{(I)}$, $I = 1, 2, 3, 4$.

Note that there are five more chains of blow downs that lead to a cubic fibration by performing contractions in different orders. We find that all of these lead to the same Jacobian.

7.3 Discriminant locus analysis

An explicit analysis of the discriminant locus Δ_J in this example reveals that it has the following factorization structure,

$$\Delta_{24,24}(\mathbf{x}_5, \mathbf{x}_6) = (F_{1,0}^{(1)}(\mathbf{x}_5, \mathbf{x}_6))^2 (F_{1,0}^{(2)}(\mathbf{x}_5, \mathbf{x}_6))^2 (F_{1,1}^{(3)}(\mathbf{x}_5, \mathbf{x}_6))^2 F_{18,22}^{(4)}(\mathbf{x}_5, \mathbf{x}_6). \quad (7.24)$$

Here, the subscript pairs denote the bi-degrees. The vanishing locus $L = \{\Delta = 0\} \subset B$ thus decomposes accordingly as

$$L = \bigcup_{I=1}^4 L^{(I)}, \quad (7.25)$$

where $L^{(I)} = \{F^{(I)} = 0\} \subset B$ are the vanishing loci of the individual factors in eq. (7.24).

By further analyzing the vanishing orders of f , g and Δ , we see the following singularity structure and the enhancement pattern (characterized by Kodaira type [34, 35]):

- $L^{(I)}$, for $I = 1, 2, 3$, are of type I_2 and $L^{(4)}$ is of type I_1 ;
- $|L^{(I)} \cap L^{(3)}| = 1$, for $I = 1, 2$, and at each of the two points the singularity enhances to I_4 ;
- $|L^{(I)} \cap L^{(4)}| = 18$, for $I = 1, 2$, and at 14 points the singularity enhances to I_3 and at the remaining 4 points (each with multiplicity 2) to III ;
- $|L^{(3)} \cap L^{(4)}| = 32$ and at 24 points the singularity enhances to I_3 and at the remaining 8 points (each with multiplicity 2) to III .

The decomposition of the discriminant locus is schematically depicted in figure 3.

7.4 Arithmetic of the sections

Since we have presented a complete section analysis in several other cases, in this example we will simply content ourselves with an analysis of the decomposition of the rank-seven Picard lattice of X . The four different types of the divisor classes of X are as follows:

- (1) Base divisor classes: $B = \mathbb{P}^1 \times \mathbb{P}^1$ have two independent generators,

$$\begin{aligned} \mathcal{O}_X(0, 0, 0, 0, 1, 0), \\ \mathcal{O}_X(0, 0, 0, 0, 0, 1). \end{aligned} \tag{7.26}$$

- (2) Fibral divisors: as we have seen, there are three codimension-one loci of singularity enhancement to I_2 and hence, three vertical/blow-up divisors. For completeness, we provide a choice of three independent vertical-divisor classes,

$$\begin{aligned} \mathcal{O}_X(-1, 1, 0, 0, 1, 0), \\ \mathcal{O}_X(1, -1, 0, 0, 1, 0), \\ \mathcal{O}_X(0, -1, 0, 1, 0, 0), \end{aligned} \tag{7.27}$$

each of which is effective with a unique GHS. Note that they do not meet with a generic fiber (when substituted into the left hand side of eq. (7.10) we obtain zero).

- (3) Zero section: $\mathcal{O}_X(0, 0, 1, -1, 1, 1)$ is a rational section, which we may take as the zero section.
- (4) Rational sections: in order for the above six independent divisor classes of three different types, together with the rational sections, to span the rank-seven Picard lattice, there must exist an additional rational section generating a MW group of rank one. Given that the other six divisor classes, all obtained from the favorable sector, generate the rank-six sub-lattice, the Mordell-Weil generator must lie (in part) in the non-favorable sector of the Picard lattice, which we will not discuss in this paper. Nevertheless, the decomposition of the Picard lattice alone is enough to demonstrate that such a group of sections exists.

8 Example 4: a threefold example with higher rank Mordell-Weil group

In this section, we will present an example of an elliptic CICY threefold with $\text{rk } MW(X) = 4$. The configuration of this geometry (CICY #7907 of the CICY threefold list [78]) is given by

$$X_3 = \left[\begin{array}{c|cccccccc} \mathbb{P}_{\mathbf{x}_1}^2 & 1 & 0 & 1 & 0 & 0 & 0 & 1 & 0 \\ \mathbb{P}_{\mathbf{x}_2}^2 & 0 & 1 & 0 & 1 & 0 & 0 & 0 & 1 \\ \mathbb{P}_{\mathbf{x}_3}^3 & 0 & 0 & 0 & 0 & 1 & 1 & 1 & 1 \\ \mathbb{P}_{\mathbf{x}_4}^1 & 0 & 0 & 1 & 1 & 0 & 0 & 0 & 0 \\ \mathbb{P}_{\mathbf{x}_5}^1 & 1 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ \mathbb{P}_{\mathbf{x}_6}^1 & 0 & 0 & 0 & 0 & 1 & 1 & 0 & 0 \\ \mathbb{P}_{\mathbf{x}_7}^1 & 0 & 0 & 0 & 0 & 2 & 0 & 0 & 0 \end{array} \right] \longrightarrow B_2 = \mathbb{P}_{\mathbf{x}_6}^1 \times \mathbb{P}_{\mathbf{x}_6}^1, \tag{8.1}$$

where the elliptic fibers have the following configuration,

$$F = \left[\begin{array}{c|cccccccc} \mathbb{P}_{\mathbf{x}_1}^2 & 1 & 0 & 1 & 0 & 0 & 0 & 1 & 0 \\ \mathbb{P}_{\mathbf{x}_2}^2 & 0 & 1 & 0 & 1 & 0 & 0 & 0 & 1 \\ \mathbb{P}_{\mathbf{x}_3}^3 & 0 & 0 & 0 & 0 & 1 & 1 & 1 & 1 \\ \mathbb{P}_{\mathbf{x}_4}^1 & 0 & 0 & 1 & 1 & 0 & 0 & 0 & 0 \\ \mathbb{P}_{\mathbf{x}_5}^1 & 1 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \end{array} \right] . \quad (8.2)$$

It has the Hodge numbers $(h^{1,1}, h^{2,1}) = (7, 30)$. We start our analysis by first choosing a generic complex structure, described by the following monomial coefficients listed in the lexicographic order for the monomial exponents described in section 6,

$$\begin{aligned} c_{P_1} &= \{18, 14, 14, 13, 14, 13\} , \\ c_{P_2} &= \{18, 16, 19, 13, 9, 15\} , \\ c_{P_3} &= \{2, 20, 7, 16, 2, 15\} , \\ c_{P_4} &= \{4, 15, 3, 2, 5, 19\} , \\ c_{P_5} &= \{6, 7, 12, 20, 4, 15, 15, 2, 16, 20, 2, 1, 4, 15, 19, 14, 10, 10, 6, 9, 16, 19, 7, 16\} , \\ c_{P_6} &= \{11, 4, 10, 7, 11, 8, 17, 9\} , \\ c_{P_7} &= \{10, 7, 5, 7, 1, 13, 20, 7, 3, 13, 18, 3\} , \\ c_{P_8} &= \{3, 4, 6, 6, 20, 4, 17, 3, 4, 20, 20, 20\} . \end{aligned} \quad (8.3)$$

8.1 Section analysis

Putative sections. A putative section S , which we label as $\mathcal{O}_X(S) = \mathcal{O}_X(b_1, \dots, b_6)$, must have a single topological intersection with a generic fiber, leading to

$$3b_1 + 3b_2 + 2b_3 + 2b_4 + 2b_5 \stackrel{!}{=} 1 , \quad (8.4)$$

and the base topology criteria (2.12), one for each base divisor, $D_\alpha^b \sim J_6, J_7$, give the following conditions,

$$\begin{aligned} &2b_1^2 + 8b_1b_2 + 2b_2^2 + 12b_1b_3 + 12b_2b_3 + 4b_3^2 + 4b_1b_4 + 4b_2b_4 + 8b_3b_4 + 4b_1b_5 \\ &+ 4b_2b_5 + 8b_3b_5 + 4b_4b_5 + 6b_1b_\alpha + 6b_2b_\alpha + 4b_3b_\alpha + 4b_4b_\alpha + 4b_5b_\alpha \stackrel{!}{=} -2 , \end{aligned} \quad (8.5)$$

for $\alpha = 6, 7$, where eq. (8.4) has been used for a simplification. These two equations, (8.4) and (8.5), then lead to the following four-parameter ($k_m \in \mathbb{Z}$, for $m = 1, 2, 3, 4$) family of divisor classes.

$$\begin{aligned} b_1 &= -1 - 2k_1 - k_2 , \\ b_2 &= k_2 , \\ b_3 &= k_3 , \\ b_4 &= k_4 , \\ b_5 &= 2 + 3k_1 - k_3 - k_4 , \\ b_6 &= 2 + 10k_1 + 8k_1^2 + 2k_2 + 4k_1k_2 + 2k_2^2 - 4k_3 - 4k_1k_3 + 2k_3^2 - 4k_4 - 6k_1k_4 + 2k_3k_4 + 2k_4^2 , \\ b_7 &= 2 + 10k_1 + 8k_1^2 + 2k_2 + 4k_1k_2 + 2k_2^2 - 4k_3 - 4k_1k_3 + 2k_3^2 - 4k_4 - 6k_1k_4 + 2k_3k_4 + 2k_4^2 . \end{aligned} \quad (8.6)$$

With some small parameter values, this family of putative sections gives rise to the following possibilities:

S	(k_1, k_2, k_3, k_4)	$\mathcal{O}_X(S)$
S_0	$(0, -1, 0, 1)$	$\mathcal{O}_X(0, -1, 0, 1, 1, 0, 0)$
S_1	$(0, -1, 1, 0)$	$\mathcal{O}_X(0, -1, 1, 0, 1, 0, 0)$
S_2	$(0, -1, 1, 1)$	$\mathcal{O}_X(0, -1, 1, 1, 0, 0, 0)$
S_3	$(-1, 0, 0, -1)$	$\mathcal{O}_X(1, 0, 0, -1, 0, 0, 0)$
S_4	$(0, 0, 0, 1)$	$\mathcal{O}_X(-1, 0, 0, 1, 1, 0, 0)$

(8.7)

all of which turn out to have the line-bundle cohomologies,

$$h^\bullet(X, \mathcal{O}_X(S)) = (1, 0, 0, 0) . \quad (8.8)$$

We must next ensure that these putative sections do correspond to true sections of the fibration.

Explicit expressions for the sections. We will fully analyze the putative section associated with the first line bundle in eq. (8.7), $L_z := \mathcal{O}_X(S_0) = \mathcal{O}_X(0, -1, 0, 1, 1, 0, 0)$, whose GHS is of the form,

$$z = \frac{N(\mathbf{x}_4, \mathbf{x}_5)}{D(\mathbf{x}_2)} \in \Gamma(X, L_z) . \quad (8.9)$$

Here $N \in \Gamma(X, \mathcal{O}_X(0, 0, 0, 1, 1, 0, 0))$ and $D \in \Gamma(X, \mathcal{O}_X(0, 1, 0, 0, 0, 0, 0))$. We choose $D = x_{2,0}$ and consider the second and the fourth defining equations,

$$P_2 = x_{2,0}q_0(\mathbf{x}_5) + x_{2,1}q_1(\mathbf{x}_5) + x_{2,2}q_2(\mathbf{x}_5) , \quad (8.10)$$

$$P_4 = x_{2,0}u_0(\mathbf{x}_4) + x_{2,1}u_1(\mathbf{x}_4) + x_{2,2}u_2(\mathbf{x}_4) , \quad (8.11)$$

where q_i and u_i , for $i = 0, 1, 2$, are linear in \mathbf{x}_5 and \mathbf{x}_4 , respectively. Then, on the divisor $\{x_{2,0} = 0\} \subset X$, we have

$$x_{2,1}q_1(\mathbf{x}_5) + x_{2,2}q_2(\mathbf{x}_5) = 0 , \quad (8.12)$$

$$x_{2,1}u_1(\mathbf{x}_4) + x_{2,2}u_2(\mathbf{x}_4) = 0 , \quad (8.13)$$

and hence, with the choice,

$$N := q_1u_2 - q_2u_1 = 68x_{4,0}x_{5,0} + 343x_{4,1}x_{5,0} + 20x_{4,0}x_{5,1} + 217x_{4,1}x_{5,1} , \quad (8.14)$$

the rational expression (8.9) for z is the desired GHS. The zero locus of this GHS can easily be seen to have no singularities via a patch-wise analysis of the system. With this expression, we obtain the explicit parameterization of its zero locus of the divisor of the form,

$$\begin{aligned} x_{1,i} &= A_{1,i}(\mathbf{x}_6, \mathbf{x}_7) \in \Gamma(B, \mathcal{O}_B) , & \text{for } i = 0, 1, 2 , \\ x_{2,i} &= A_{2,i}(\mathbf{x}_6, \mathbf{x}_7) \in \Gamma(B, \mathcal{O}_B(2, 2)) , & \text{for } i = 0, 1, 2 , \\ x_{3,i} &= A_{3,i}(\mathbf{x}_6, \mathbf{x}_7) \in \Gamma(B, \mathcal{O}_B(2, 2)) , & \text{for } i = 0, 1, 2, 3 , \\ x_{4,i} &= A_{4,i}(\mathbf{x}_6, \mathbf{x}_7) \in \Gamma(B, \mathcal{O}_B) , & \text{for } i = 0, 1 , \\ x_{5,i} &= A_{5,i}(\mathbf{x}_6, \mathbf{x}_7) \in \Gamma(B, \mathcal{O}_B) , & \text{for } i = 0, 1 , \end{aligned} \quad (8.15)$$

which we present in appendix B.3 for concreteness. These parameterizations provide a rational section to the elliptic fibration. In particular, the section map is ill-defined, in the same manner as has been seen in preceding sections. This phenomenon occurs for the $\mathbb{P}_{\mathbf{x}_2}^2$ direction at 8 points on the base $\mathbb{P}_{\mathbf{x}_6}^1 \times \mathbb{P}_{\mathbf{x}_7}^1$, amongst which are 4 points where the map is ill-defined also for the $\mathbb{P}_{\mathbf{x}_3}^3$ direction. Note that the Euler number difference,

$$\chi(S) - \chi(B) = 12 - 4 = 8, \tag{8.16}$$

indicates that the rational section indeed wraps a \mathbb{P}^1 in the fiber over these 8 base points.

Similarly, one can analyze the other smooth, rational sections, $S_{m=1,2,3,4}$, using the same procedures. We thereby obtain the explicit rational maps for each of those MW generators and also confirm the way in which fibral \mathbb{P}^1 's are wrapped by each rational section.

8.2 Locating the singular fibers

As in section 7, exploration of the Jacobian is the simplest way in which to obtain the information we need about the singular fibers. With the fiber (8.2) being of codimension 8, to easily read off the Jacobian of X we first need to blow the fiber geometry down. Let us consider the following chain of three contractions along the $\mathbb{P}_{\mathbf{x}_1}^2$, $\mathbb{P}_{\mathbf{x}_3}^3$, and $\mathbb{P}_{\mathbf{x}_4}^1$ directions, in turn:

$$\begin{aligned} & \left[\begin{array}{c} \mathbb{P}_{\mathbf{x}_1}^2 \\ \mathbb{P}_{\mathbf{x}_2}^2 \\ \mathbb{P}_{\mathbf{x}_3}^3 \\ \mathbb{P}_{\mathbf{x}_4}^1 \\ \mathbb{P}_{\mathbf{x}_5}^1 \\ \mathbb{P}_{\mathbf{x}_6}^1 \\ \mathbb{P}_{\mathbf{x}_7}^1 \end{array} \middle\| \begin{array}{c} 1 \ 0 \ 1 \ 0 \ 0 \ 0 \ 1 \ 0 \\ 0 \ 1 \ 0 \ 1 \ 0 \ 0 \ 0 \ 1 \\ 0 \ 0 \ 0 \ 0 \ 1 \ 1 \ 1 \ 1 \\ 0 \ 0 \ 1 \ 1 \ 0 \ 0 \ 0 \ 0 \\ 1 \ 1 \ 0 \ 0 \ 0 \ 0 \ 0 \ 0 \\ 0 \ 0 \ 0 \ 0 \ 1 \ 1 \ 0 \ 0 \\ 0 \ 0 \ 0 \ 0 \ 2 \ 0 \ 0 \ 0 \end{array} \right] \rightarrow \left[\begin{array}{c} \mathbb{P}_{\mathbf{x}_2}^2 \\ \mathbb{P}_{\mathbf{x}_3}^3 \\ \mathbb{P}_{\mathbf{x}_4}^1 \\ \mathbb{P}_{\mathbf{x}_5}^1 \\ \mathbb{P}_{\mathbf{x}_6}^1 \\ \mathbb{P}_{\mathbf{x}_7}^1 \end{array} \middle\| \begin{array}{c} 0 \ 1 \ 1 \ 0 \ 0 \ 1 \\ 1 \ 0 \ 0 \ 1 \ 1 \ 1 \\ 1 \ 0 \ 1 \ 0 \ 0 \ 0 \\ 1 \ 1 \ 0 \ 0 \ 0 \ 0 \\ 0 \ 0 \ 0 \ 1 \ 1 \ 0 \\ 0 \ 0 \ 0 \ 2 \ 0 \ 0 \end{array} \right] \tag{8.17} \\ & \rightarrow \left[\begin{array}{c} \mathbb{P}_{\mathbf{x}_2}^2 \\ \mathbb{P}_{\mathbf{x}_4}^1 \\ \mathbb{P}_{\mathbf{x}_5}^1 \\ \mathbb{P}_{\mathbf{x}_6}^1 \\ \mathbb{P}_{\mathbf{x}_7}^1 \end{array} \middle\| \begin{array}{c} 1 \ 1 \ 1 \\ 1 \ 0 \ 1 \\ 1 \ 1 \ 0 \\ 2 \ 0 \ 0 \\ 2 \ 0 \ 0 \end{array} \right] \rightarrow \left[\begin{array}{c} \mathbb{P}_{\mathbf{x}_2}^2 \\ \mathbb{P}_{\mathbf{x}_5}^1 \\ \mathbb{P}_{\mathbf{x}_6}^1 \\ \mathbb{P}_{\mathbf{x}_7}^1 \end{array} \middle\| \begin{array}{c} 2 \ 1 \\ 1 \ 1 \\ 2 \ 0 \\ 2 \ 0 \end{array} \right]. \end{aligned}$$

After this procedure the fiber has been blown down to a complete intersection of two hypersurfaces in $\mathbb{P}_{\mathbf{x}_2}^2 \times \mathbb{P}_{\mathbf{x}_5}^1$ of degrees (2, 1) and (1, 1), respectively. This blown-down fiber has the PALP ID (4, 0) and, based on the results of ref. [30], we can read the Jacobian of the form,

$$y^2 = x^3 + f_J x z^4 + g_J z^6, \tag{8.18}$$

where $f_J \in \Gamma(B, K_B^{-4})$ and $g_J \in \Gamma(B, K_B^{-6})$. The discriminant $\Delta_J := 4f_J^3 + 27g_J^2$ of the Jacobian is a homogeneous polynomial of degree (24, 24). Upon going through the

analysis of the resolved geometry, without eliminating the fiber coordinates, we are able to numerically compare the the two discriminant loci and confirm that

$$\Delta_J \sim \Delta_{\text{res}} . \tag{8.19}$$

Having obtained the explicit discriminant polynomial Δ_J , one can proceed to analyze its factorization properties and we find that all the codimension-one fibers are of type I_1 .

8.3 Arithmetic of the sections

Let us denote sections corresponding to S_m , $m = 0, 1, \dots, 4$ in (8.7), by $\sigma_m \in MW(X)$. Here we will show that the MW group is of rank 4, and that the lattice is generated by these σ_m . Specifically, based on the group law (4.3) at the level of divisor classes, we will show, upon choosing σ_0 to be the zero section, the following:

- $\sigma_1, \dots, \sigma_4$ span a lattice of rank 4;
- $\sigma_1, \dots, \sigma_4$ generate all the putative section classes in the family (8.6).

These two combined will prove not only that $\sigma_{m=1,2,3,4}$ is a basis of the MW group, but also that all of the putative section classes in eq. (8.6) correspond to a genuine section. Both of the bullet points above can be shown by forming the divisor class of a most general linear combination of σ_m . Upon sequentially applying the addition law (4.3), one sees that

$$\text{Div}\left(\bigoplus_{m=1}^4 l_m \sigma_m\right) = \sum_{r=1}^7 \beta_r J_r , \tag{8.20}$$

where the coefficients β_r are given in terms of the l_m as

$$\begin{aligned} \beta_1 &= l_3 - l_4 , \\ \beta_2 &= -1 + l_3 + l_4 , \\ \beta_3 &= l_1 + l_2 , \\ \beta_4 &= 1 - l_1 - 2l_3 , \\ \beta_5 &= 1 - l_2 - l_3 , \\ \beta_6 &= -2l_1 + 2l_1^2 - 2l_2 + 2l_1l_2 + 2l_2^2 - 2l_3 + 2l_1l_3 + 2l_3^2 - 2l_4 + 2l_4^2 , \\ \beta_7 &= -2l_1 + 2l_1^2 - 2l_2 + 2l_1l_2 + 2l_2^2 - 2l_3 + 2l_1l_3 + 2l_3^2 - 2l_4 + 2l_4^2 . \end{aligned} \tag{8.21}$$

With eqs. (8.20) and (8.21), one can immediately see that

$$\bigoplus_{m=1}^4 l_m \sigma_m = \sigma_0 \quad \Rightarrow \quad \text{Div}\left(\bigoplus_{m=1}^4 l_m \sigma_m\right) = S_0 \quad \Rightarrow \quad l_m = 0, \quad \text{for } m = 1, 2, 3, 4, \tag{8.22}$$

which proves the first bullet point. In addition, the b_r from eq. (8.6) equal β_r for $r = 1, \dots, 7$, with the following choice of l_m ,

$$l_1 = 1 + 2k_1 - k_4 ; \quad l_2 = -1 - 2k_1 + k_3 + k_4 ; \quad l_3 = -k_1 ; \quad l_4 = 1 + k_1 + k_2 , \tag{8.23}$$

which proves the second bullet point. Therefore, we have in particular shown $\text{rk } MW(X) = 4$, which is consistent with the following decomposition of the independent divisor classes:

- (1) Base divisor classes: $B = \mathbb{P}^1 \times \mathbb{P}^1$ has two independent generators,

$$\begin{aligned} \mathcal{O}_X(0, 0, 0, 0, 0, 1, 0), \\ \mathcal{O}_X(0, 0, 0, 0, 0, 0, 1). \end{aligned} \tag{8.24}$$

- (2) Fibral divisors: as we have seen, there do not arise any enhancement of fiber singularities at codimension one and hence, no blow-up divisors either.

- (3) Zero section: σ_0 has been proven to be a rational section, which we may take as the zero section with the divisor class,

$$\mathcal{O}_X(0, -1, 0, 1, 1, 0, 0). \tag{8.25}$$

- (4) Rational sections: $MW(X)$ has been proven to be generated by σ_m , $m = 1, 2, 3, 4$, with the respective divisor classes,

$$\begin{aligned} \mathcal{O}_X(0, -1, 1, 0, 1, 0, 0), \\ \mathcal{O}_X(0, -1, 1, 1, 0, 0, 0), \\ \mathcal{O}_X(1, 0, 0, -1, 0, 0, 0), \\ \mathcal{O}_X(-1, 0, 0, 1, 1, 0, 0). \end{aligned} \tag{8.26}$$

Then, the contributions from these four different types of generators give rise to the correct rank for the Picard lattice, $2 + 0 + 1 + 4 = 7$. Furthermore, the seven divisor classes in eqs. (8.24), (8.25) and (8.26) indeed span the entire Picard lattice and hence form a basis.

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A Reducible putative sections and non-flat fibers

In this section, we provide examples of reducible putative sections that decompose into a genuine section and a vertical divisor. Let us consider the CICY threefold discussed in section 7, with the elliptic fibration given by the CICY and the fiber configuration matrices (7.1) and (7.3), respectively,

$$X_3 = \left[\begin{array}{c|cccccc} \mathbb{P}_{x_1}^1 & 1 & 1 & 0 & 0 & 0 & 0 \\ \mathbb{P}_{x_2}^1 & 1 & 0 & 1 & 0 & 0 & 0 \\ \mathbb{P}_{x_3}^2 & 0 & 2 & 0 & 0 & 0 & 1 \\ \mathbb{P}_{x_4}^3 & 0 & 0 & 1 & 1 & 1 & 1 \\ \mathbb{P}_{x_5}^1 & 1 & 0 & 0 & 1 & 0 & 0 \\ \mathbb{P}_{x_6}^1 & 0 & 1 & 0 & 0 & 1 & 0 \end{array} \right]; \quad F = \left[\begin{array}{c|cccccc} \mathbb{P}_{x_1}^1 & 1 & 1 & 0 & 0 & 0 & 0 \\ \mathbb{P}_{x_2}^1 & 1 & 0 & 1 & 0 & 0 & 0 \\ \mathbb{P}_{x_3}^2 & 0 & 2 & 0 & 0 & 0 & 1 \\ \mathbb{P}_{x_4}^3 & 0 & 0 & 1 & 1 & 1 & 1 \end{array} \right], \tag{A.1}$$

with the base $B = \mathbb{P}_{\mathbf{x}_5}^1 \times \mathbb{P}_{\mathbf{x}_6}^1$. Recall that in section 7 we have identified a family of putative sections (see eq. (7.12)), from which three line bundles with small degrees,

$$L_1 = \mathcal{O}_X(0, 0, 1, -1, 1, 1), \quad L_2 = \mathcal{O}_X(0, -1, 1, 0, 1, 1), \quad L_3 = \mathcal{O}_X(-1, 0, 1, 0, 2, 1), \quad (\text{A.2})$$

were chosen for discussion. Note that all these three line bundles have the cohomology,

$$h^\bullet(X, L_i) = (1, 0, 0, 0), \quad \text{for } i = 1, 2, 3. \quad (\text{A.3})$$

The unique global holomorphic section of L_1 is of the form,

$$s_1 = \frac{N(\mathbf{x}_3, \mathbf{x}_5, \mathbf{x}_6)}{x_{4,0}}, \quad (\text{A.4})$$

with the numerator polynomial N given in eq. (7.18), and its vanishing locus has proven smooth, leading to a genuine section to the elliptic fibration (see eq. (7.19) and appendix B.2). On the other hand, the other two putative sections associated with the line bundles, L_2 and L_3 , are neither smooth nor irreducible, while leading to the same section map as the L_1 case. As claimed in the main text, they should be thought of as the genuine section attached with an additional vertical divisor. Since the two cases are similar in nature, we will only present a full analysis for one of them, L_2 . We start with the following injective mapping for cohomology groups,

$$H^0(X, L_1) \times H^0(X, L) \longrightarrow H^0(X, L_2), \quad (\text{A.5})$$

where $L := \mathcal{O}_X(0, -1, 0, 1, 0, 0)$ satisfies $L_1 \otimes L = L_2$ and $h^0(X, L) = 1$. Given such an injection, the unique GHS of L_2 should be of the form,

$$s_2 = s_1 s, \quad (\text{A.6})$$

where s is the GHS of L that can be written, for instance, as

$$s = \frac{r_1(\mathbf{x}_4)}{x_{2,0}}, \quad (\text{A.7})$$

with $r_1(\mathbf{x}_4)$ read from the expansion of the third defining equation for X ,

$$P_3 = x_{2,0}r_0(\mathbf{x}_4) + x_{2,1}r_1(\mathbf{x}_4). \quad (\text{A.8})$$

The divisor $\{s = 0\} \subset X$ does not intersect with a generic fiber and is in fact vertical. One way to see this is to compute the left hand side of the Oguiso criterion (2.11) and notice that one obtains zero. Via numerical algebraic geometry techniques, one can further check that the singular locus of the putative section associated with L_2 is of dimension 1, which is consistent with the fact that the genuine section and the vertical divisor meets at a curve.

Similarly, the relevant injection for the L_3 case is,

$$H^0(X, L_1) \times H^0(X, L') \longrightarrow H^0(X, L_3), \quad (\text{A.9})$$

where $L' := \mathcal{O}_X(-1, 0, 0, 1, 1, 0)$ satisfies $L_1 \otimes L' = L_3$ and $h^0(X, L') = 1$. The argument for this case follows along exactly analogous lines.

The decomposition structure described above is rather clean. For some configurations, however, one may face a more exotic situation where the elliptic fiber is non-flat and a point on the base pulls back to a vertical divisor. As an illustration, let us consider the CICY threefold (with CICY #5075 [78]), in which an elliptic fibration can be found with the CICY and the fiber configurations,

$$X_3 = \left[\begin{array}{c|cccccc} \mathbb{P}^1_{\mathbf{x}_1} & 1 & 0 & 1 & 0 & 0 & 0 \\ \mathbb{P}^1_{\mathbf{x}_2} & 0 & 1 & 0 & 0 & 1 & 0 \\ \mathbb{P}^2_{\mathbf{x}_3} & 0 & 2 & 0 & 0 & 0 & 1 \\ \mathbb{P}^3_{\mathbf{x}_4} & 0 & 0 & 1 & 1 & 1 & 1 \\ \mathbb{P}^1_{\mathbf{x}_5} & 1 & 1 & 0 & 0 & 0 & 0 \\ \mathbb{P}^1_{\mathbf{x}_6} & 1 & 0 & 0 & 1 & 0 & 0 \end{array} \right]; \quad F = \left[\begin{array}{c|cccccc} \mathbb{P}^1_{\mathbf{x}_1} & 1 & 0 & 1 & 0 & 0 & 0 \\ \mathbb{P}^1_{\mathbf{x}_2} & 0 & 1 & 0 & 0 & 1 & 0 \\ \mathbb{P}^2_{\mathbf{x}_3} & 0 & 2 & 0 & 0 & 0 & 1 \\ \mathbb{P}^3_{\mathbf{x}_4} & 0 & 0 & 1 & 1 & 1 & 1 \end{array} \right], \quad (\text{A.10})$$

over the base $B = \mathbb{P}^1_{\mathbf{x}_5} \times \mathbb{P}^1_{\mathbf{x}_6}$. Note that this is the same CICY geometry as the one analyzed in section 7, with the rows appropriately interchanged (i.e., the total space is the same but we are examining a different fibration structure here). One can easily confirm, via the simple topological checks of section 2.3, that the two line bundles,

$$L_1 = \mathcal{O}_X(1, 0, 1, -1, 0, 1), \quad L_2 = \mathcal{O}_X(0, 0, 1, -1, 1, 2), \quad (\text{A.11})$$

correspond to putative sections. Furthermore, they both have the cohomology,

$$h^\bullet(X, L_i) = (1, 0, 0, 0), \quad \text{for } i = 1, 2. \quad (\text{A.12})$$

Let us expand the first, the third, the fourth, and the sixth defining equations for X as

$$P_1 = \sum_{i=0}^1 x_{1,i} p_i(\mathbf{x}_5, \mathbf{x}_6), \quad (\text{A.13})$$

$$P_3 = \sum_{i=0}^3 x_{4,i} r_i(\mathbf{x}_1), \quad (\text{A.14})$$

$$P_4 = \sum_{i=0}^3 x_{4,i} u_i(\mathbf{x}_6), \quad (\text{A.15})$$

$$P_6 = \sum_{i=0}^3 x_{4,i} w_i(\mathbf{x}_3), \quad (\text{A.16})$$

where r_i , u_i , and w_i , for $i = 0, 1, 2, 3$, are all linear in their respective variables and p_i , for $i = 0, 1$, are bilinear in \mathbf{x}_5 and \mathbf{x}_6 . Then, the GHS of L_1 can be constructed as

$$s_1 = \frac{-r_3 u_2 w_1 + r_2 u_3 w_1 + r_3 u_1 w_2 - r_1 u_3 w_2 - r_2 u_1 w_3 + r_1 u_2 w_3}{x_{4,0}}, \quad (\text{A.17})$$

and its zero locus can easily be proven smooth. Then, exactly the same steps as those used for all the examples in the main text can be applied to show that this putative section is a genuine section to the elliptic fibration (A.10).

For the GHS of L_2 , we consider the injection,

$$H^0(X, L_1) \times H^0(X, L) \longrightarrow H^0(X, L_2), \quad (\text{A.18})$$

where $L := \mathcal{O}_X(-1, 0, 0, 0, 1, 1)$ satisfies $L_1 \otimes L = L_2$ and $h^0(X, L) = 1$. The GHS s_2 of L_2 then factors as

$$s_2 = s_1 s, \quad (\text{A.19})$$

where s is the section of L that can be written, for instance, as

$$s = \frac{p_1(\mathbf{x}_5, \mathbf{x}_6)}{x_{1,0}} \sim -\frac{p_0(\mathbf{x}_5, \mathbf{x}_6)}{x_{1,1}}. \quad (\text{A.20})$$

Thus, the putative section $\{s_2 = 0\} \subset X$ decomposes into the genuine section $\{s_1 = 0\}$ and the vertical divisor $\{s = 0\}$. However, the vertical divisor in this case is the non-flat fiber over the two points in the base with $p_0(\mathbf{x}_5, \mathbf{x}_6) = 0 = p_1(\mathbf{x}_5, \mathbf{x}_6)$.¹²

It is interesting to look in more detail at the structure associated to these two points. In order to do so, we proceed to analyze the Jacobian of the blown-down geometry, obtained via the following chain of contractions:

$$\left[\begin{array}{c|cccccc} \mathbb{P}_{\mathbf{x}_1}^1 & 1 & 0 & 1 & 0 & 0 & 0 \\ \mathbb{P}_{\mathbf{x}_2}^1 & 0 & 1 & 0 & 0 & 1 & 0 \\ \mathbb{P}_{\mathbf{x}_3}^2 & 0 & 2 & 0 & 0 & 0 & 1 \\ \mathbb{P}_{\mathbf{x}_4}^3 & 0 & 0 & 1 & 1 & 1 & 1 \\ \mathbb{P}_{\mathbf{x}_5}^1 & 1 & 1 & 0 & 0 & 0 & 0 \\ \mathbb{P}_{\mathbf{x}_6}^1 & 1 & 0 & 0 & 1 & 0 & 0 \end{array} \right] \rightarrow \left[\begin{array}{c|cccc} \mathbb{P}_{\mathbf{x}_2}^1 & 0 & 1 & 0 & 1 & 0 \\ \mathbb{P}_{\mathbf{x}_3}^2 & 0 & 2 & 0 & 0 & 1 \\ \mathbb{P}_{\mathbf{x}_4}^3 & 1 & 0 & 1 & 1 & 1 \\ \mathbb{P}_{\mathbf{x}_5}^1 & 1 & 1 & 0 & 0 & 0 \\ \mathbb{P}_{\mathbf{x}_6}^1 & 1 & 0 & 1 & 0 & 0 \end{array} \right] \rightarrow \left[\begin{array}{c|cccc} \mathbb{P}_{\mathbf{x}_3}^2 & 2 & 0 & 0 & 1 \\ \mathbb{P}_{\mathbf{x}_4}^3 & 1 & 1 & 1 & 1 \\ \mathbb{P}_{\mathbf{x}_5}^1 & 1 & 1 & 0 & 0 \\ \mathbb{P}_{\mathbf{x}_6}^1 & 0 & 1 & 1 & 0 \end{array} \right] \rightarrow \left[\begin{array}{c|c} \mathbb{P}_{\mathbf{x}_3}^2 & 3 \\ \mathbb{P}_{\mathbf{x}_5}^1 & 2 \\ \mathbb{P}_{\mathbf{x}_6}^1 & 2 \end{array} \right].$$

The Jacobian of the resulting cubic fibration [31] has a Weierstrass form whose f_J , g_J , and Δ_J , at these two base points, vanish to order 4, 6, and 12, respectively. Conversely, one may analyze non-generic codimension-two points as follows. The Δ_J turns out to factorize as

$$\Delta_{24,24}(\mathbf{x}_5, \mathbf{x}_6) = (F_{1,2}^{(1)}(\mathbf{x}_5, \mathbf{x}_6))^2 F_{22,20}^{(2)}(\mathbf{x}_5, \mathbf{x}_6), \quad (\text{A.21})$$

where the subscript pairs denote the bi-degrees. One may then go through all of the codimension-two points obtained by intersecting $F^{(1)}$ and $F^{(2)}$, and analyze the vanishing orders of f_J , g_J , and Δ_J there. It then turns out that at exactly two of those codimension-two points they vanish to order 4, 6, and 12, respectively, and these two points are exactly where the fiber goes non-flat. We take this as a non-trivial piece of evidence that vanishing of f , g , and Δ to those orders is intimately related to the non-flat fibers.

¹²In order to see that this fiber is non-flat at these points note that, when these two terms vanish, the first defining equation in the fiber configuration in eq. (A.10) becomes trivial (see eq. (A.13)).

The above results are somewhat to be expected given results in both the physics [2, 79, 81–84] and mathematics [85, 86] literatures. It is well known that non-flat fibers are dangerous in that they lead to SCFTs when they are blown down to obtain the associated Weierstrass model. So called “(4, 6, 12)” points are also well known to be associated with SCFTs, dual to heterotic small instantons, and so it is not surprising that this is what we obtain. Note that, while normally in the physics literature (4, 6, 12) points are dealt with by blowing up the base, here we see that resolving via a non-flat fiber may also be an option. Indeed it is known that non-flat fibers can always be birationally related to a flat fibration, by going to the Weierstrass model and then blowing up the base [85, 86]. Thus if one can blow up to a non-flat fiber to resolve some given (4, 6, 12) points, then removing them by blowing up the base will also be an option.

B Explicit rational maps

B.1 The rational map in eq. (6.17)

The coefficients for the respective basis monomials in the expression (6.17) are given as

$$\begin{aligned}
 c_{A_{1,0}} &= \{374645428, 2459924454, 6041201902, 6757197031, 3506712251, 580810171, 924301148, 8199431465, \\
 &\quad 25264031186, 36372427880, 22779823774, 4908357591, 1452188771, 13033199001, 40135221388, \\
 &\quad 66573888839, 54108574906, 15515879412, 1248897453, 12094877912, 34197169604, 54193633181, \\
 &\quad 55266938999, 22696202323, 215837430, 6368649137, 15892630074, 22118006553, 22024966369, \\
 &\quad 15692070369, 237121534, -325372751, 6378759403, 1293660759, 3196094763, 3924307629\}, \\
 c_{A_{1,1}} &= \{-734944740, -4326339166, -11130294298, -13907839771, -9009394083, -2347190168, \\
 &\quad -1675319136, -14776875073, -46120017979, -63136826781, -36725168560, -7668395811, \\
 &\quad -2113655527, -18322779267, -72678726227, -110406161553, -65489903673, -10031189305, \\
 &\quad -1349647374, -11165685192, -53929588927, -95298609096, -59057362305, -9086250737, \\
 &\quad -625397263, -1377088015, -21506754436, -41349413418, -24925221695, -5744117643, \\
 &\quad -270489017, 1411511282, -5280035471, -6318376072, -4098094623, -1500696270\}, \\
 c_{A_{2,0}} &= \{1084, 723, -1357, 446, 1061, -1858, 2295, -4617, 1014\}, \\
 c_{A_{2,1}} &= \{1636, 7255, 4564, 128, 11845, 12591, -166, 3943, 6549\}, \\
 c_{A_{2,2}} &= \{-2410, -6655, -6728, -2985, -15122, -8890, -1337, -6671, -3039\},
 \end{aligned}$$

where the coefficients are listed in the lexicographic order of the monomial exponents described in section 6.

B.2 The rational map in eq. (7.19)

The coefficients for the respective basis monomials in the expression (7.19) are given as

$$\begin{aligned}
 c_{A_{1,0}} &= \{2221064631, 11917255459, 28570974597, 16495328929, 8210461686, 7704389866, \\
 &\quad -49864858414, -52760626002, 8036361006, -15034898100, 6923938379, 37655021444\}, \\
 c_{A_{1,1}} &= \{-2813621667, -13737694558, -23649388867, -13410639788, -10885604364, -16063185502, \\
 &\quad 34747606348, 35848540654, -10733842857, 14828715552, 4172982765, -22394056833\},
 \end{aligned}$$

$$\begin{aligned}
c_{A_{2,0}} &= \{32876672958, 151940458372, 197121202830, 109428884752, 127447660413, 193402090199, \\
&\quad -438295120157, -348851016375, 125317972440, -68136505562, 330460624192, 602425176324, \\
&\quad -2417059116, 78061046172, -189976587911, -401928611774\}, \\
c_{A_{2,1}} &= \{27539562618, 155918941832, 429523158738, 249442739852, 64390426983, -104328590424, \\
&\quad -1019425293349, -969130178938, -39468341094, -438573483026, 455036821602, 911597522988, \\
&\quad -134257162533, 161878775784, 112484337279, -154768837497\}, \\
c_{A_{3,0}} &= \{7357, 26839, 10114, -49331\}, \\
c_{A_{3,1}} &= \{537, 15147, -1059, 1817\}, \\
c_{A_{3,2}} &= \{-15452, -30676, -29780, 39679\}, \\
c_{A_{4,0}} &= \{26125015333068, 14093705690638, -253407065151880, -1392929418892058, -969784333605016, \\
&\quad 44111282567346, -403354228064373, 154876615769747, 5112562529762937, 4813145357835159, \\
&\quad -171100860164100, -616912937629296, 868417042593585, -6973432524385874, -8327868925909157, \\
&\quad -436942930362768, 56095959480934, -1955852891330594, 3018949340241019, 5939379450686520, \\
&\quad -233507718098055, -349570577705856, 1272688939129888, 369255734060038, -1443023739731898\}, \\
c_{A_{4,1}} &= \{52147316163654, 182788360653752, 31326182511870, -568882677708652, -448279427660016, \\
&\quad 196582084116597, 18967518423364, -1149983213111555, 512164610746422, 1360945691525348, \\
&\quad 196592267852520, -250589375084200, 1437953867180930, 1652250053613618, -567580289130554, \\
&\quad 32425178124081, 442219027542972, -56965838050663, -1358907015270799, -779145891064596, \\
&\quad 46248190501128, 118523292993366, -333807482278694, -124559271657028, 390617948985922\}, \\
c_{A_{4,2}} &= \{-51721845976638, -211044216465212, -185467081673182, 770647194632056, 613395005398400, \\
&\quad -85575733018929, 267990300915931, 868293730827257, -1937085042649463, -2902140503357580, \\
&\quad 217465884106170, 270811223476316, -3411572449024905, 364374416960336, 3657090237062633, \\
&\quad 361920509039658, -1037318384139832, 2066163139747743, 1800958612708262, -1942857053473884, \\
&\quad -61786399417851, 407958330780186, -684815876214366, -858431195253898, 730146388419428\}, \\
c_{A_{4,3}} &= \{-10832442228972, -816963082358, 131545420796888, 537953669498482, 370184917765880, \\
&\quad -99402081447330, -80480182341872, 20890954975173, -916841939725634, -1003292183336129, \\
&\quad -218570855874462, 341593139545558, 1176891692757061, 661234289887252, 235975445849429, \\
&\quad -46147618216851, 487317156800170, -1349531840686582, -1215954798086102, 875569510180860, \\
&\quad 175304685193347, -456024004562430, 331856671930056, 689389343218620, -468384360246228\},
\end{aligned}$$

where the coefficients are listed in the lexicographic order of the monomial exponents described in section 6.

B.3 The rational map in eq. (8.15)

The coefficients for the respective basis monomials in the expression (8.15) are given as

$$\begin{aligned}
c_{A_{1,0}} &= \{-651897\}, \\
c_{A_{1,1}} &= \{61934\}, \\
c_{A_{1,2}} &= \{548944\}, \\
c_{A_{2,0}} &= \{21105788506648, -11775044280804, 6163099181603, 29703669175731, -14775894485118, \\
&\quad -9206484163372, 9487229871540, -4625233450212, -7147022545227\},
\end{aligned}$$

$$\begin{aligned}
c_{A_2,1} &= \{-4435876022168, 1454959565556, -2028437292847, -6978840453732, 1754250504522, 421432378118, \\
&\quad -2379663273177, 466722799641, 457564831947\}, \\
c_{A_2,2} &= \{-14587835867512, 8657245681028, -3886990601851, -20156267237980, 10899890597746, \\
&\quad 7132962137614, -6361220860757, 3453847947797, 5471027603303\}, \\
c_{A_3,0} &= \{1836343820, -1284971884, 348994373, 2279657494, -1578310010, -1230865198, 611200657, \\
&\quad -492599473, -835170489\}, \\
c_{A_3,1} &= \{-347211138, -689099572, -736001231, -1590588657, -718071560, -1330998724, -950616680, \\
&\quad -206565160, -573918245\}, \\
c_{A_3,2} &= \{798897195, -82629092, 494287724, 1618408765, -182328463, 276360492, 726380168, -72433064, \\
&\quad -4829618\}, \\
c_{A_3,3} &= \{-1500914105, 1290270968, -112711221, -1457099424, 1503524551, 1448568784, -177947468, \\
&\quad 443979836, 821860735\}, \\
c_{A_4,0} &= \{67\}, \\
c_{A_4,1} &= \{4\}, \\
c_{A_5,0} &= \{-92\}, \\
c_{A_5,1} &= \{247\},
\end{aligned}$$

where the coefficients are listed in the lexicographic order of the monomial exponents described in section 6.

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