

NLSP gluino search at the Tevatron and early LHC

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ABSTRACT: We investigate the collider phenomenology of gluino-bino co-annihilation scenario both at the Tevatron and 7 TeV LHC. This scenario can be realized, for example, in a class of realistic supersymmetric models with non-universal gaugino masses and $t - b - \tau$ Yukawa unification. The NLSP gluino and LSP bino should be nearly degenerate in mass, so that the typical gluino search channels involving leptons or hard jets are not available. Consequently, the gluino can be lighter than various bounds on its mass from direct searches. We propose a new search for NLSP gluino involving multi-b final states, arising from the three-body decay $\tilde{g} \rightarrow b\bar{b}\tilde{\chi}_1^0$. We identify two realistic models with gluino mass of around 300 GeV for which the three-body decay is dominant, and show that a 4.5σ observation sensitivity can be achieved at the Tevatron with an integrated luminosity of 10 fb^{-1} . For the 7 TeV LHC with 50 pb^{-1} of integrated luminosity, the number of signal events for the two models is $\mathcal{O}(10)$, to be compared with negligible SM background events.

KEYWORDS: Supersymmetry Phenomenology

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1 Introduction

Low-scale supersymmetry (SUSY) is arguably the leading candidate of new physics and will be seriously tested at the CERN Large Hadron Collider (LHC). It provides an elegant solution to the gauge hierarchy problem and with conserved R -parity, the thermal relic abundance of the lightest neutralino (LSP) can often be identified with dark matter, consistent with the current cosmological observations [1]. The recent results reported by the CDMS-II experiment may indicate the presence of dark matter with mass of around $\mathcal{O}(100)$ GeV [2]. To reproduce the required thermal relic abundance, a pure wino or Higgsino dark matter should have mass of a few TeV due to rapid annihilation rate. On the contrary, the annihilation rate of a pure bino with mass of around 100 GeV is too slow, and leads to excessive relic abundance [3]. A variety of constraints from low energy flavor physics and CP violating physics typically favor scenarios with heavy sfermions [4–8]. Consequently, pure bino self annihilation is further suppressed if scalars in the t -channel become heavier. To enhance the annihilation rate of pure bino, there are generally two categories of models available, and both will lead to interesting phenomenology at the LHC. In one scenario the dark matter is a bino-wino or bino-Higgsino mixture [9], so that $\tilde{\chi}_1^0 \tilde{\chi}_1^0$ annihilation is enhanced via the enlarged coupling. The mass degeneracy $M_1 \simeq M_2$ leads to a nearly degenerate chargino-neutralino spectrum [10]. The second scenario is co-annihilation [11–13], and if the scalars are heavy, the required relic abundance can be achieved through co-annihilation with a strongly-interacting particle such as the gluino, namely $\tilde{\chi}_1^0 \tilde{g}, \tilde{g}\tilde{g} \rightarrow f\bar{f}$ [3, 14–18]. This scenario can be realized in a realistic class of $t - b - \tau$ Yukawa unified models [14–16, 19–22]. In the gluino-bino co-annihilation scenarios, the gluino is the next-to-lightest-supersymmetric-particle (NLSP), and the mass splitting between \tilde{g} and $\tilde{\chi}_1^0$ is relatively small, namely [3]

$$\frac{M_{\tilde{g}} - M_{\tilde{\chi}_1^0}}{M_{\tilde{\chi}_1^0}} \lesssim 20\% . \quad (1.1)$$

Note that the Sommerfeld enhancement of bino-gluino co-annihilation cross section is rather mild, such that eq. (1.1) is altered by only 2 – 3% [12, 13, 17, 18].

The relations between gaugino masses will be crucial in understanding the nature of supersymmetry breaking and of the underlying theory at ultra-high energy scale. For instance, supersymmetric grand unification, string dilaton SUSY breaking, and minimal gauge mediation all predict the gaugino mass relations $M_3/g_3^2 = M_2/g_2^2 = M_1/g_1^2$. With NLSP gluino, one must invoke non-universal gaugino masses at M_{GUT} . This not only implies very different physics from the above models but can also give rise to very different phenomenology at hadron colliders. We focus in this paper on the phenomenological implications of bino-gluino co-annihilation scenarios at hadron colliders.

Being a color octet fermion, gluino pair production is the most promising discovery channel for supersymmetry at hadron colliders. In an environment such as this with huge QCD jet backgrounds, isolated charged leptons (μ^\pm, e^\pm) and b -jets usually play an important role in the searches. For the most well studied scenarios where the charginos are lighter than gluinos, the Majorana nature of the gluino will result in same-sign chargino signature (jets plus $\tilde{\chi}^\pm\tilde{\chi}^\pm$). This eventually leads to same-sign dileptons plus jets with very little Standard Model (SM) background. With a NLSP gluino, however, the chargino as well as the leptons are absent in the gluino cascade decay. The conventional search strategy does not work here and the NLSP gluino can evade the current bounds from direct searches at the Tevatron. The NLSP gluino can be relatively light and its production rates can be large both at the Tevatron and especially at the LHC.

On the other hand, as a consequence of t - b - τ Yukawa unification, the third generation squarks, stops or sbottoms, are usually much lighter than those of the first two generations. In the large $\tan\beta$ limit, sbottoms are often the lightest squarks. As a consequence, gluino decays may lead to top-rich or bottom-rich events. In the co-annihilation region with $\Delta M = M_{\tilde{g}} - M_{\tilde{\chi}_1^0} \ll 2m_t$, there is no phase-space for on-shell top quarks. The gluino decay into b -jets, $\tilde{g} \rightarrow b\bar{b}\tilde{\chi}_1^0$, then becomes dominant and this enable us to search for NLSP gluino via multi- b jets, namely

$$p\bar{p}, pp \rightarrow \tilde{g}\tilde{g} \rightarrow b\bar{b}b\bar{b} + \cancel{E}_T. \quad (1.2)$$

Multi- b events have been widely proposed for light scalar searches, for instance in NMSSM with the Higgs decaying into multi- b jets via light scalars ($h \rightarrow aa \rightarrow b\bar{b}b\bar{b}$) [23]. The multi- b jets plus significant missing transverse energy may only appear in Wh, Zh associated productions whose production rates are much smaller than gluino pair production. The reconstruction is also very different.

The paper is organized as follows. In the next section (II), we discuss NLSP gluino decay and the parameter space where multi- b jet production is significant. We very briefly comment models in which NLSP gluino can be realized and discuss its implications. In section III and IV, we study the collider phenomenology of this scenario, with event selection and identification both at Tevatron and the LHC. We summarize our findings in section V.

2 NLSP gluino: decays and benchmark models

As mentioned earlier, the gluino-bino co-annihilation scenario requires the gluino to be NLSP in the sparticle spectrum, and to be nearly degenerate in mass with the bino LSP. The mass difference between the two should be $\lesssim 20\% M_{\tilde{\chi}_1^0}$. In the framework of minimal supergravity, this clearly requires non-universal gaugino masses at M_{GUT} . The leading motivations for grand unification theories (GUTs such as SO(10)) are the explanation of tiny neutrino masses and charge quantization. However, a partial unified model such as $SU(4)_C \times SU(2)_L \times SU(2)_R$ (4-2-2) [24] also provides solutions to both problems. Non-universal asymptotic gaugino masses are naturally accommodated in the supersymmetric 4-2-2 model and have recently been investigated in ref. [14–16]. Other examples with non-universal gaugino masses include a supersymmetric $SU(5) \times SU(3)_{\text{Hypercolor}}$ proposed to explain doublet-triplet splitting problem in SU(5) GUT [25, 26], and GUT models with non-singlet F -term vevs [27].

Since we focus on a spectrum with NLSP gluino, on-shell charginos are kinematically forbidden in gluino decay. The color octet gluino therefore can decay only into colored SM particles such as the gluon octet or a $q\bar{q}$ pair, plus the color singlet LSP $\tilde{\chi}_1^0$,

$$\tilde{g} \rightarrow q\bar{q}\tilde{\chi}_1^0, g\tilde{\chi}_1^0. \quad (2.1)$$

The three-body decay $\tilde{g} \rightarrow q\bar{q}\tilde{\chi}_1^0$ is through off-shell squark interchange, while the two-body decay $g\tilde{\chi}_1^0$ can be realized by the triangle loop involving squarks. The partial widths of these two decay channels are given by [28, 29]

$$\begin{aligned} \Gamma(\tilde{g} \rightarrow g\tilde{\chi}_1^0) &= \frac{(M_{\tilde{g}}^2 - M_{\tilde{\chi}_1^0}^2)^3}{2\pi M_{\tilde{g}}^3} \left[\frac{g_3^2 g_1}{128\pi^2} (M_{\tilde{g}} - M_{\tilde{B}}) \sum_q Q_q \left(\frac{1}{M_{\tilde{q}_L}^2} - \frac{1}{M_{\tilde{q}_R}^2} \right) N_{1B} \right. \\ &\quad \left. + \frac{g_3^2 y_t^2}{32\sqrt{2}\pi^2 \sin\beta} \left(\frac{1}{M_{\tilde{q}_L}^2} + \frac{1}{M_{\tilde{u}_R}^2} \right) N_{1H_u} v \left(1 + \ln \frac{m_t^2}{M_{\tilde{g}}^2} \right) \right]^2, \quad (2.2) \\ \Gamma(\tilde{g} \rightarrow q\bar{q}\tilde{\chi}_1^0) &= \frac{M_{\tilde{g}}^5}{768\pi^3} \left[\left(\frac{g_3 g_1}{6M_{\tilde{q}_L}^2} N_{1B} + \frac{g_3 g_2}{2M_{\tilde{q}_L}^2} N_{1W} \right)^2 + \left(\frac{2g_3 g_1}{3M_{\tilde{u}_R}^2} N_{1B} \right)^2 \right. \\ &\quad \left. + \left(\frac{g_3 g_1}{6M_{\tilde{q}_L}^2} N_{1B} - \frac{g_3 g_1}{2M_{\tilde{q}_L}^2} N_{1W} \right)^2 + \left(\frac{g_3 g_1}{3M_{\tilde{d}_R}^2} N_{1B} \right)^2 \right] f \left(\frac{M_{\tilde{\chi}_1^0}}{M_{\tilde{g}}} \right) \quad (q = u, d), \quad (2.3) \end{aligned}$$

$$\begin{aligned} f(x) &= 1 + 2x - 8x^2 + 18x^3 - 18x^5 + 8x^6 - 2x^7 - x^8 \\ &\quad - 12x^4 \ln x^2 + 12x^3(1+x^2) \ln x^2. \quad (2.4) \end{aligned}$$

Here N_{1B} , N_{1W} and N_{1H_u} are respectively the bino \tilde{B} , wino \tilde{W} and Higgsino \tilde{H}_u components of the LSP neutralino $\tilde{\chi}_1^0$. The three-body decays will be suppressed if the scalar masses are too large, or due to phase space if the mass difference between \tilde{g} and $\tilde{\chi}_1^0$ is too small. The two-body mode may dominate in this case. If either the gluino two-body decay $\tilde{g} \rightarrow g\tilde{\chi}_1^0$, or three-body decay due to the small mass difference into light jets dominates, the final state jets are typically as soft as those from parton showers. In this case the gluino

decay very likely gets buried in the huge QCD background, and consequently the search becomes extremely challenging [17, 18, 30]. Our study will focus on a different region where the gluino three-body decays into multi- b jets dominate.

To illustrate the search strategy, we have selected two benchmark points from a previous study of supersymmetric 4-2-2 models [14–16]. In 4-2-2 models, the matter fields of each family belong in $(4, 2, 1)$ and $(4, 1, 2)$. The third family fermion masses, to a good approximation, arise from a Yukawa coupling to the bi-doublet $(1, 2, 2)$. Thus, $t - b - \tau$ Yukawa unification arises as a natural prediction [14–16, 19–22]

$$Y_t = Y_b = Y_\tau \equiv Y_{\text{Dirac}} . \tag{2.5}$$

In $\text{SO}(10)$, $t - b - \tau$ Yukawa unification typically predicts gluino to be the lightest colored sparticle [31, 32], while in 4-2-2 models one realizes the NLSP gluino scenario through gluino-bino co-annihilation. A large bottom Yukawa Y_b also naturally drives sbottom to be the lightest squark. With $\Delta M \simeq 50 \text{ GeV}$ and $\mathcal{O}(\text{TeV})$ squarks, $\tilde{g} \rightarrow b\bar{b}\tilde{\chi}_1^0$ decay often dominates. Figure 1 shows the dependence of the gluino decay BR in the $\Delta M - M_{\tilde{b}_1}$ plane for the 4-2-2 model with $\mu < 0$. We require consistency of the model with various phenomenological constraints such as $\text{BR}(B_s \rightarrow \mu^+\mu^-)$, $\text{BR}(b \rightarrow s\gamma)$, $\text{BR}(B \rightarrow \tau\nu)$, $\Delta(g - 2)_\mu$, WMAP relic density (in the 5σ range), and all the sparticle mass bounds [14–16]. In figure 1 the points which satisfy $t - b - \tau$ Yukawa unification are a subset of the displayed points and mostly lie in the dense region around 1 TeV sbottom mass, where the mass difference ΔM is between 40 and 60 GeV. For this region the branching fraction of gluino three-body decay $\text{BR}(\tilde{g} \rightarrow b\bar{b}\tilde{\chi}_1^0)$ dominates both the two-body one, $\text{BR}(\tilde{g} \rightarrow g\tilde{\chi}_1^0)$, as well as the three-body decay into light quarks $\text{BR}(\tilde{g} \rightarrow q\bar{q}\tilde{\chi}_1^0)$. For a sufficiently large sbottom mass the two-body decay is dominant, as can be seen from figure 1. Note that we show only those scenarios in figure 1 for which gluino is the NLSP. We have picked two benchmark models, one with $\mu > 0$ (Model A), listed as point 1 in ref. [14, 15]; the second model has $\mu < 0$ (Model B), selected from the large number of points in figure 1. The relevant observables are listed in table 1. Both models can evade the direct search Tevatron gluino bounds due to the dominant decay mode $\tilde{g} \rightarrow b\bar{b}\tilde{\chi}_1^0$.

3 NLSP gluino search at the Tevatron

In both benchmark models, the gluino masses are of order 200-300 GeV, and their pair production rates at the Tevatron are around the pico-barn level. Therefore, a search for a relatively light NLSP gluino appears quite promising. In this section, we illustrate how one could identify NLSP gluinos at the Tevatron for the two benchmark models above.

As shown in table 1 for both models we focus on NLSP gluino with the dominant three-body decay $\tilde{g} \rightarrow b\bar{b}\tilde{\chi}_1^0$. Therefore we wish to identify signal events of gluino pair production with 4 b -jets plus missing transverse energy \cancel{E}_T ,

$$p\bar{p} \rightarrow \tilde{g}\tilde{g} \rightarrow b\bar{b}b\bar{b} + \cancel{E}_T . \tag{3.1}$$

Due to the relatively long lifetime of the B -mesons, their decays on average take place $\mathcal{O}(\text{mm})$ distance away from the primary interacting vertex. With the vertex detector,

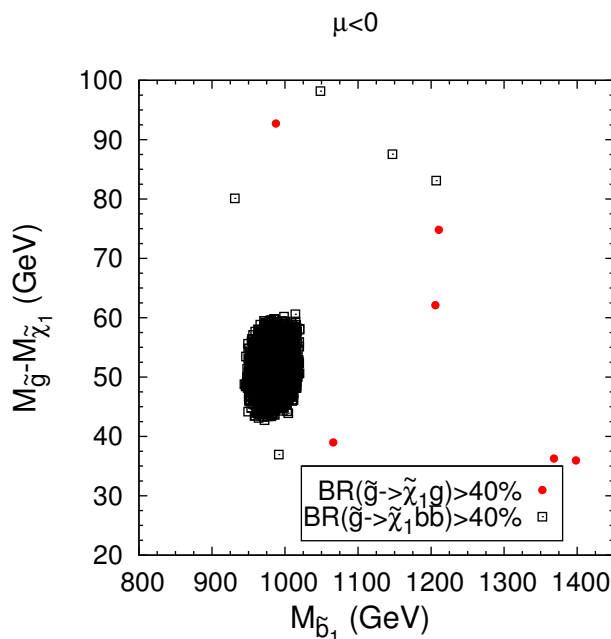


Figure 1. Mass difference $M_{\tilde{g}} - M_{\tilde{\chi}_1^0}$ versus $M_{\tilde{b}_1}$ for 4-2-2 model with $\mu < 0$. The points shown satisfy all the experimental constraints described in the text. The red circular points depict scenarios for which the branching fraction of two-body decay is dominant, i.e. $\text{BR}(\tilde{g} \rightarrow g\tilde{\chi}_1^0) > 40\%$. For the empty square points in black the three-body decay is dominant, $\text{BR}(\tilde{g} \rightarrow b\bar{b}\tilde{\chi}_1^0) > 40\%$. Gluino is the NLSP for all the points shown in the figure.

	Model A ($\mu > 0$)	Model B ($\mu < 0$)
m_0 (GeV)	14110	1513
M_1 (GeV)	499.54	-479.49
M_2 (GeV)	832.03	-845.5
M_3 (GeV)	0.7945	69.53
$\tan\beta$	50.82	47.7
A_0	-34551.2	-1668.84
m_{H_u} (GeV)	6092.74	492.41
m_{H_d} (GeV)	14194.5	1071.75
$M_{\tilde{g}}$ (GeV)	329	261
$M_{\tilde{\chi}_1^0}$ (GeV)	284	207
$M_{\tilde{b}_1}$ (GeV)	5294	950
$\text{BR}(\tilde{g} \rightarrow b\bar{b}\tilde{\chi}_1^0)$	76.3%	50.8%

Table 1. Model parameters at GUT scale (above double line) and low scale (below double line) for two benchmark models. Note that the bino component of $\tilde{\chi}_1^0 \geq 99.9\%$.

tagging jet with decaying B -mesons will significantly reduce the QCD jets background. The b -tagging efficiency at the Tevatron is taken to be $\epsilon_b = 50\%$ [33, 34]. The b -jet production in the SM is either due to gluon splitting or from top quark decay originating

from top pair and/or single production ($t\bar{t}$ pair provides $b\bar{b}$ in the final states, while single top production $tg \rightarrow \bar{b}tj'$ also provides $b\bar{b}$), and so the b -jets always arise as pairs. Hence we only need to tag three b -jets so that we do not have to pay the additional 50% loss in the fourth b -tagging. After multiplying by the b -tagging efficiency, both the signal and the background events are reduced by

$$\epsilon_b^3 = 12.5\% . \quad (3.2)$$

By requiring 3 b -tagged jets, the SM production $4b + X$ becomes the leading irreducible background. Also, there exist reducible backgrounds due to other jets being mis-identified as b -events. About 15% of D -mesons in the jets can be mis-identified as B 's, and the mis- b tagging rate for light jets is 0.4% [33, 34]. Since the c -jet production in the SM is very similar to the b -jet production, the production rate is at comparable level. With the 15% faking rate, we can safely neglect the c -jet production in our study. However, the production rate of light jets is several orders of magnitude higher than that of pure b -jet production [35] and cannot therefore be neglected even with 0.4% mis- b -tagging rate. Therefore, we include $jjb\bar{b} + X$ in our background analyses.

Besides the $b\bar{b}$, the dark matter particles $\tilde{\chi}_1^0$ also appear in the gluino cascade decay. The missing transverse energy \cancel{E}_T is another characteristic feature of the signal. The irreducible SM background for \cancel{E}_T is from Z production, with the branching fraction of Z invisible decay ($Z \rightarrow \nu\bar{\nu}$) as 20%. However, due to the uncertainty of mis-measurement in jet energy or momentum, the events without Z can also induce \cancel{E}_T . The third source is due to leptonic decays of W^\pm bosons, especially $W^\pm \rightarrow \tau^\pm\nu_\tau \rightarrow \ell^\pm\nu_\ell\bar{\nu}_\tau\nu_\tau$ where the leptons from τ three-body decays are below the visible lepton cut ($p_T^\ell > 10$ GeV). The SM backgrounds that we consider in the study are then

$$\begin{aligned} & b\bar{b}b\bar{b}, jjb\bar{b}, \\ & b\bar{b}b\bar{b}Z, jjb\bar{b}Z \quad \text{with } \text{BR}(Z \rightarrow \nu\bar{\nu}) = 20\%, \\ & t\bar{t} \rightarrow b\bar{b}jj\tau^\pm\nu_\tau \quad \text{with } \tau^\pm \text{ leptonic decay} . \end{aligned} \quad (3.3)$$

To simulate the detector effects, we smear the hadronic jet energy by a Gaussian distribution whose width is parameterized as [33, 34]

$$\frac{\Delta E_j}{E_j} = \frac{75\%}{\sqrt{E_j/\text{GeV}}} \oplus 5\% . \quad (3.4)$$

The following basic kinematical cuts on the transverse momentum (p_T), the pseudo-rapidity (η), and the separation in the azimuthal angle-pseudo rapidity plane ($\Delta R = \sqrt{(\Delta\phi)^2 + (\Delta\eta)^2}$) between two jets have been employed for jet-selection [33, 34]

$$p_T^j > 15 \text{ GeV}, |\eta_j| < 1.0, \Delta R_{jj} > 0.4. \quad (3.5)$$

Because of the relatively small mass difference $\Delta M = M_{\tilde{g}} - M_{\tilde{\chi}_1^0} \simeq 50$ GeV, the b -jets in final states could be rather soft. The minimal and maximal p_T distributions of b -jets are plotted in figure 2(a). The plots that we show in this paper are all for benchmark model B, with the relevant features for Model A expected to be very similar. One can see that

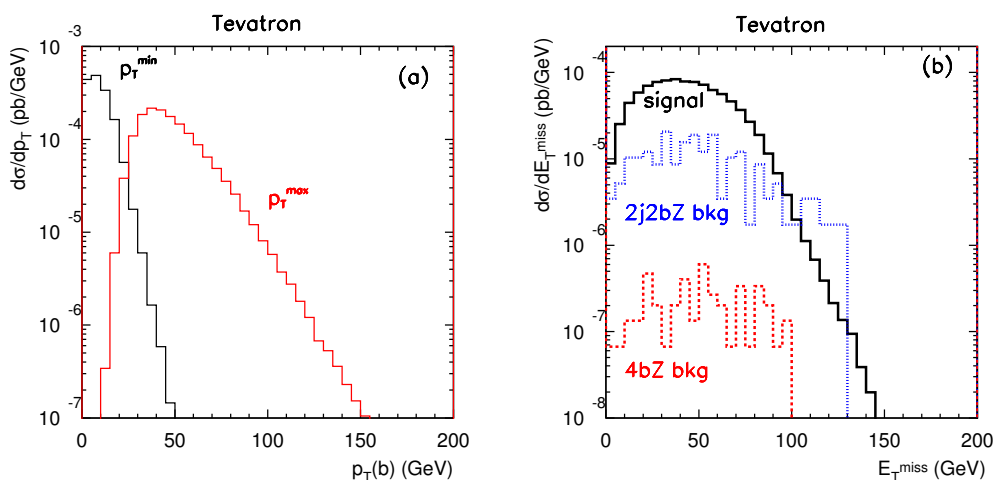


Figure 2. (a) Minimal and maximal p_T distribution of 4 b -jets in the signal $p\bar{p} \rightarrow \tilde{g}\tilde{g} \rightarrow b\bar{b}b\bar{b}\tilde{\chi}_1^0\tilde{\chi}_1^0$ at Tevatron. (b) Missing transverse energy (\cancel{E}_T) distribution in signal events as well as its background $4bZ$ and $jjb\bar{b}Z$ at Tevatron, taking the branching fraction of Z invisible decay ($Z \rightarrow \nu\bar{\nu}$) as 20% and mis- b -tagging rate of light jet as 0.4%.

the b -jet with minimal p_T would be rejected by the basic p_T cut in eq. (3.5). To retain as many signal events as possible, we apply the p_T cut for three b -jets except the softest one. The softest jet will be viewed as unreconstructed calorimeter energy in the detector. It is consistent with the requirement of 3 tagged b -jets above. The missing transverse energy \cancel{E}_T is reconstructed according to the smeared observed particles, namely b -jets in our case. We show the \cancel{E}_T distribution of the signal and backgrounds including basic cuts for η and ΔR in eq. (3.5) and b -jet and mis- b jet tagging efficiency in figure 2(b). We find that by requiring a significant \cancel{E}_T cut in the final states, namely

$$\cancel{E}_T > 30 \text{ GeV}, \tag{3.6}$$

the SM $4b$ and $jjb\bar{b}$ backgrounds at Tevatron can be completely eliminated. Since the signal events do not contain any lepton in the final states, we veto any event with visible leptons of

$$p_T^\ell > 10 \text{ GeV} . \tag{3.7}$$

The leading reducible background of soft leptons type is due to semi-leptonic $t\bar{t}$ events with one W^\pm decaying into τ^\pm , and τ^\pm further decaying into soft leptons. With one of the light jets faking a b -jet, we find that the contribution due to $t\bar{t}$ is below 0.01 fb. The leading irreducible background after \cancel{E}_T cut is then $jjb\bar{b}Z$ with invisible Z decay.

In table 2, we summarize the signal and background cross sections at the Tevatron for the two benchmark models, after implementing the basic cuts and \cancel{E}_T cut. For our numerical analyses, we use the CTEQ6L1 parton distribution function [36]. The SM backgrounds are simulated by the automatic package Madgraph/Madevent [35]. The signal significance is obtained in terms of Gaussian statistics, given by the ratio S/\sqrt{B} of signal and background events with a luminosity of 10 fb^{-1} . For benchmark model A with $M_{\tilde{g}} \sim 300 \text{ GeV}$,

$\sigma(\text{fb})@ \text{TeVatron}$	Model A	Model B	$4b$	$4bZ$	$2j2bZ$	S/\sqrt{B}
basic cuts and 3b tagging	2.3	4.8	2.7×10^3	0.02	1	
$\cancel{E}_T > 30 \text{ GeV}$	1.4	3.3	–	0.019	0.95	4.5(A)/11(B)

Table 2. Production cross section for $p\bar{p} \rightarrow \tilde{g}\tilde{g} \rightarrow b\bar{b}b\bar{b}\tilde{\chi}_1^0\tilde{\chi}_1^0$ in models A and B and for backgrounds $4b, 4bZ, 2j2bZ$, after basic cuts and missing energy cut at the Tevatron with a luminosity of 10 fb^{-1} .

the statistical significance is close to 5σ , while for benchmark model B with a smaller gluino mass, as expected it is 2.5 times larger.

In order to further identify the signal events and extract spectrum information, we propose to reconstruct the events through the invariant mass distribution of di- b jets $M_{b\bar{b}}$ and M_{T2} method. There are 4 b -jets in the final states so we will have three combinations of b -jet pair in event reconstruction. Usually the two hardest b -jets come from different gluinos and the b -jets pair coming from the same \tilde{g} has smaller separation angles. Therefore, following the reconstruction method in ref. [31, 32], we select the two hardest b -jets b_1, b_2 , and let b_3 denote the third jet that minimizes $\Delta\phi(b_2, b_3)$, such that the pair b_2, b_3 come from the same \tilde{g} . The other b -jets pair from the second \tilde{g} consists of two b -jets b_1 and b_4 . For gluino three-body decay, the dijet invariant mass $M_{b\bar{b}}$ must be less than the mass difference between the gluino and the LSP masses [37]

$$M_{b\bar{b}} \leq \Delta M. \tag{3.8}$$

To illustrate this, we display the distribution of $M_{b\bar{b}} = \max[M_{b_1b_4}, M_{b_2b_3}]$ for benchmark model B in figure 3(a); it clearly shows the edge of $M_{b\bar{b}}$ around ΔM . It is important to note that the final states contain two invisible massive particles emanating separately from two parent particles, and therefore it is usually hard to reconstruct the signal events. However we still have sufficient information to fully reconstruct the signal events. Since the signal dominates according to our analyses above, the gluino mass can be estimated from the production rate. Once $M_{\tilde{g}}$ and the mass difference ΔM are known, the LSP mass $M_{\tilde{\chi}_1^0}$ can be easily obtained. By substituting the LSP mass $M_{\tilde{\chi}_1^0}$ into the reconstruction, we can plot the variable M_{T2} which is defined as [38–40]

$$M_{T2}^2(\tilde{g}) = \min_{\mathbf{p}_T^{\chi^{(1)}} + \mathbf{p}_T^{\chi^{(2)}} = \mathbf{p}_T^{\text{miss}}} \{ \max[M_T^{2(1)}, M_T^{2(2)}] \}, \tag{3.9}$$

where the transverse masses $M_T^{(1)}, M_T^{(2)}$ are constructed for two gluino decay chains in terms of the relevant transverse invariant mass and transverse momentum of $b\bar{b}$ system and as function of the trial LSP mass. Figure 3(b) shows the M_{T2} reconstruction of the gluinos in Model B with trial LSP mass as 200 GeV. The consistency with our assumptions confirms the LSP mass $M_{\tilde{\chi}_1^0}$ is the correct one and more importantly, it can also be used as a check for the gluino-bino co-annihilation condition as in eq. (1.1),

$$M_{\tilde{\chi}_1^0} \simeq M_{\tilde{g}} / (1 + 20\%) \simeq 200 \text{ GeV} . \tag{3.10}$$

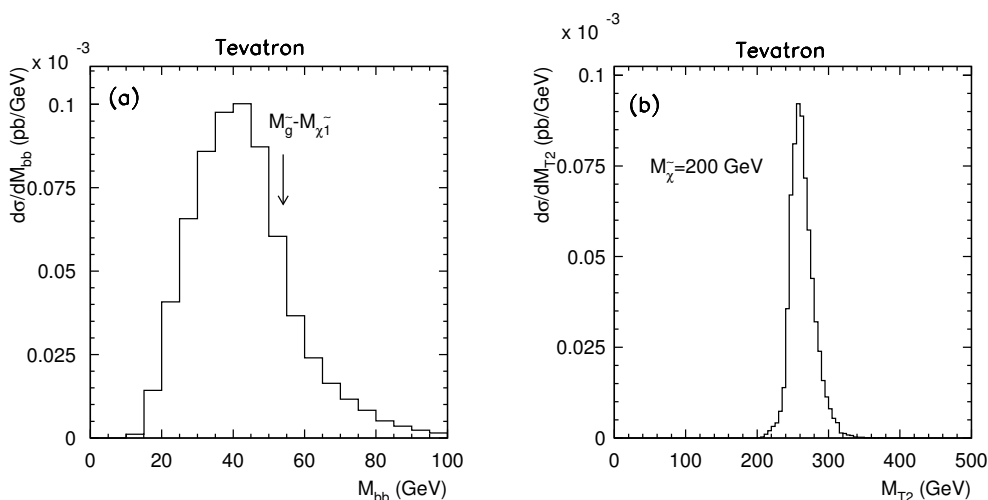


Figure 3. (a) The distribution of di- b jet invariant mass $M_{b\bar{b}}$. (b) M_{T2} reconstruction with 200 GeV trial LSP mass. Both distributions are for Model B.

4 NLSP gluino search at the LHC

In this section we discuss the NLSP gluino search at the early LHC run of 7 TeV c.m. energy. We employ a search strategy very similar to that for the Tevatron. In the LHC environment, the mis- b tagging rate for light jets is about 1/30 for low p_T range (15-50 GeV) [41]. However, since the gluino pair production at the LHC is dominated by $gg \rightarrow \tilde{g}\tilde{g}$, the total cross section is about 30 times larger than at the Tevatron. We can then require to tag all four b -jets and increase the missing transverse energy \cancel{E}_T selection cut:

- 4 b -tagged jets with $p_T^j > 15$ GeV, $|\eta_j| < 2.0$, $\epsilon_b = 50\%$;
- $\Delta R_{jj} > 0.4$;
- veto any event with lepton $p_T^\ell > 10$ GeV ;
- $\cancel{E}_T > 40$ GeV .

The smearing parameterization is given as [41]

$$\frac{\Delta E_j}{E_j} = \frac{50\%}{\sqrt{E_j/\text{GeV}}} \oplus 3\%. \quad (4.1)$$

We show the p_T distribution of b -jet in the signal final states and \cancel{E}_T distribution for signal and backgrounds in figure 4(a) and (b) respectively. One can see that the p_T distribution of b -jets in the final states at the LHC is similar to that at the Tevatron. The missing transverse energy is smaller in comparison with the SM backgrounds. We summarize in table 3 the results of signal and background studies at the LHC for the two benchmark models. After all the selection cuts, the signal events far exceed the SM background. By the first shut-down in winter 2010, the LHC should accumulate about 50 pb^{-1} of data.

$\sigma(\text{fb})@ 7\text{ TeV LHC}$	Model A	Model B	$4b$	$4bZ$	$2j2bZ$
basic cuts and $4b$ tagging	143	271	157×10^3	0.55	4.2
$\cancel{E}_T > 40\text{ GeV}$	59	140	—	0.4	3.3

Table 3. Production cross section for signal $pp \rightarrow \tilde{g}\tilde{g} \rightarrow \bar{b}b\bar{b}b\tilde{\chi}_1^0\tilde{\chi}_1^0$ and backgrounds $4b, 4bZ, 2j2bZ$ after basic cuts and missing energy cut at 7 TeV LHC for the two benchmark models.

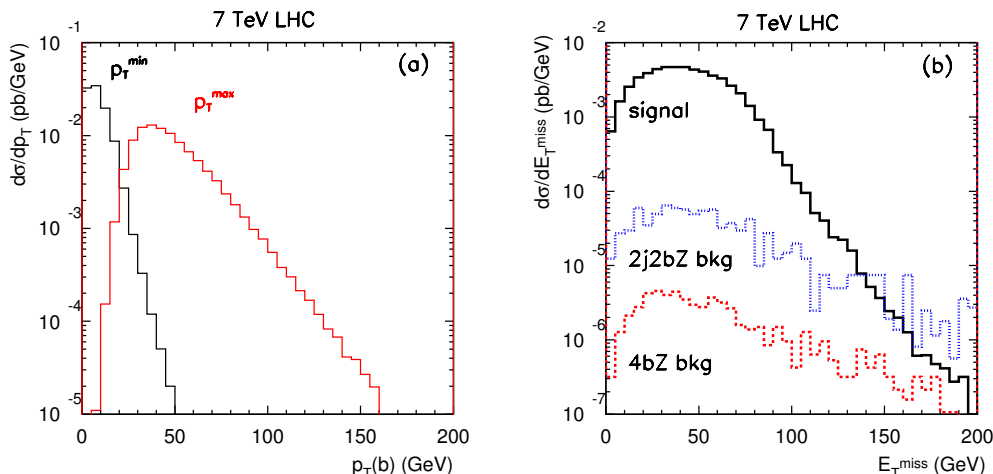


Figure 4. (a) Minimal and maximal p_T distribution of 4 b -jets in $pp \rightarrow \tilde{g}\tilde{g} \rightarrow \bar{b}b\bar{b}b\tilde{\chi}_1^0\tilde{\chi}_1^0$ at 7 TeV LHC. (b) Missing transverse energy (\cancel{E}_T) distribution in signal events and in backgrounds $4bZ$ and $jjbbZ$ at 7 TeV LHC, with the branching fraction of Z invisible decay ($Z \rightarrow \nu\bar{\nu}$) as 20% and mis- b -tagging rate of light jet as $1/30$.

After the cuts, we expect negligible background events and 3 signal events for Model A, while Model B predicts about 7 events. We expect our study to yield important clues about the underlying NLSP gluino scenario during the early stage of LHC operation.

At the LHC, another channel may also become interesting. With a light gluino exchange in the t -channel, the first generation squarks \tilde{u}, \tilde{d} can be produced together with gluinos at significant rates via valence quark-gluon scattering. For Model B with $M_{\tilde{u}} \simeq M_{\tilde{d}} \sim 1.5\text{ TeV}$, the total production cross section for

$$pp \rightarrow gu, gd \rightarrow \tilde{g}\tilde{u}, \tilde{g}\tilde{d} \tag{4.2}$$

is about 120 fb. With a NLSP gluino, $\tilde{q} \rightarrow \tilde{g}q$ decay always dominates and the gluinos will be highly boosted. The signal from heavy squarks decay will consequently be two extremely hard jets with one of them being the collimated gluino. It is then difficult to tag the b -jet in decay products of boosted gluino. However, the b -jets from the other gluino decay can still be tagged. The heavy squark resonance also provides a nice handle to suppress the SM backgrounds. In addition, with the heavy resonance, this channel may enable the search for $\tilde{g} \rightarrow \tilde{g}\tilde{\chi}_1^0$ decay. In either case, the search will require a more careful analyses of the jet substructure and we postpone this for future study.

5 Conclusion

We have explored the collider phenomenology of gluino-bino co-annihilation scenarios for both the Tevatron and 7 TeV LHC. The NLSP gluino is only slightly more massive (~ 50 GeV) than the bino LSP, so that the conventional gluino searches $\tilde{\chi}_1^\pm \tilde{\chi}_1^\pm + \text{jets}$ are not applicable. We propose to search for gluino pairs through multi-b jets final states. By using two benchmark points from a supersymmetric 4-2-2 model, in which NLSP gluino arises naturally, we explicitly show how the search strategy works at hadron colliders. It is shown that with 10 fb^{-1} integrated luminosity, one can reach over 4.5σ at Tevatron. By the end of the first LHC run at 7 TeV with 50 pb^{-1} of accumulated data, the predicted signal events for both benchmark models are $\mathcal{O}(10)$, with negligible SM background events.

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