

ELECTRONEGATIVITIES AND HARDNESSES OF SEVERAL ATOMS AND IONS CALCULATED WITH THE $X\alpha$ METHOD HAVING SELF-CONSISTENT PARAMETER α *

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The $X\alpha$ method with self-consistent parameter α_{SCF} has been used to calculate electronegativities and hardnesses of several atoms and ions. Slater's transition state method has been applied, too. The $X\alpha$ results are in agreement with the experimental data and other theoretical values.

1. Introduction

Recently, several chemical concepts have been rigorously established within the density functional theory [1-4]. The electronegativity is defined as the negative of the chemical potential of the density functional theory. The absolute hardness has been recently defined by Parr and Pearson [5] as the derivative of the chemical potential with respect to the number of electrons.

Atomic electronegativities have been calculated using several approximate energy functionals, especially the $X\alpha$ method with the transition state concept [6] has been widely applied in the calculations [7-11]. Hardnesses of elements have also been determined with the $X\alpha$ method [10, 11].

Here, electronegativities and hardnesses of several atoms and ions determined with the $X\alpha$ method having a self-consistent parameter α are presented. (Recently, the ionization energies, electron affinities and chemical potentials of several elements have been calculated with this method [12].) Electronegativities and hardnesses of several singly, doubly and triply ionized positive ions have been calculated. The results are in good agreement with the experimental data.

2. Ab initio self-consistent parameter α in the $X\alpha$ method

The way of calculating the self-consistent parameter α in the $X\alpha$ method has been discussed in several earlier papers [13, 14]. Here, only an outline is presented. The starting point of the method is the free-electron gas model. The exchange potential acting on an electron with spin up is given by

*Dedicated to Academician D. Berényi on his 60th birthday

$$V_{X\uparrow}(\vec{r}) = -8F(\zeta) \left(\frac{3}{4\pi} \rho_{\uparrow}(\vec{r}) \right)^{1/3}, \quad (1)$$

where

$$F(\zeta) = \frac{1}{2} + \frac{1-\zeta^2}{4\zeta} \ln \left| \frac{1+\zeta}{1-\zeta} \right| \quad (2)$$

and

$$\zeta = \frac{p}{p_F} \quad (3)$$

is the reduced momentum of the electron considered, p_F is the Fermi momentum and $\rho_{\uparrow}(\vec{r})$ is the total electron density of the electrons having spin up. The usual way of obtaining the $X\alpha$ exchange potential is to introduce an averaging procedure in the Fermi sphere. If the average is done for the whole Fermi sphere the value proposed by Slater [15], $\alpha = 1$, is given. On the other hand, an average near the Fermi surface for a layer thickness of which goes to zero leads to $\alpha = 2/3$, the value suggested by Gáspár [16] and later Kohn and Sham [17]. A physically more plausible case between these extreme ones can be defined by an average for a layer with $\nu_{i\uparrow} = u_i^* u_i$ electrons in the unit volume near the Fermi surface. It leads to the exchange potential

$$V_{X\alpha\uparrow}^i(\vec{r}) = -6\alpha_{\uparrow}^i(\vec{r}) \left| \frac{3}{4\pi} \rho_{\uparrow}(\vec{r}) \right|^{1/3}, \quad (4)$$

where

$$\alpha_{\uparrow}^i(\vec{r}) = \frac{\rho_{\uparrow}}{\nu_{i\uparrow}} \left[1 - \frac{1}{2}\zeta_i^3 - \frac{1}{2}\zeta_i + \frac{1}{4}(\zeta_i^2 - 1)^2 \ln \left| \frac{1+\zeta_i}{1-\zeta_i} \right| \right] \quad (5)$$

and

$$\zeta_i = \left(1 - \frac{\nu_{i\uparrow}}{\rho_{\uparrow}} \right)^{1/3}. \quad (6)$$

The expression (6) gives different functions $\alpha_{\uparrow}^i(\vec{r})$ for the different spin orbitals u_i . Recently, the r -dependence of the function $\alpha_{\uparrow}^i(\vec{r})$ has been studied [18]. It is possible to introduce an average over the shells:

$$\alpha_{\uparrow}(\vec{r}) = \frac{\sum_{i\uparrow} n_i \alpha_{\uparrow}^i(\vec{r})}{\sum_{i\uparrow} n_i}, \quad (7)$$

where n_i is the occupation number of the shell i . It has been shown that $\alpha_{\uparrow}(\vec{r})$ has only a slight r -dependence, it can be considered a constant value [14]. This constant value can be determined self-consistently and is generally referred to α_{SCF} . This way of calculating the exchange potential contains no fitting or adjustment, this method can be regarded an ab initio one.

3. Calculation of electronegativity and hardness in the $X\alpha$ method

The electronegativity χ has been defined [1, 19] as the negative of the chemical potential μ

$$\chi = -\mu = -\frac{\partial E}{\partial N}. \quad (8)$$

It has been shown [1, 7, 12] that the electronegativity can be approximated by the Mulliken formula [20]

$$\chi_M = \frac{1}{2}(I + A), \quad (9)$$

where I and A are the first ionization energy and the electron affinity, respectively. In the $X\alpha$ method, the electronegativity can be given by the highest occupied orbital energy in open shell systems [7, 12, 21] and with the transition-state method [6] in closed shell systems [7, 12].

The absolute hardness of an atom or a molecule has been defined by Parr and Pearson [5] as

$$\eta = \frac{1}{2} \left(\frac{\partial^2 E}{\partial^2 N} \right)_v = \frac{1}{2} \left(\frac{\partial \mu}{\partial N} \right)_v, \quad (10)$$

where v is the external potential. It has been shown that the hardness can be approximated [5, 10, 11] as

$$\eta = \frac{1}{2}(I - A). \quad (11)$$

Table I
Hardnesses of several atoms (in eV)

| Atom | η ($X\alpha_{SCF}$ method) | $\frac{1}{4} < r^{-1} >_i$ | $\frac{1}{4} < r^{-1} >_i$ (α_{HF}) | Other estimates for η | | |
|------|-------------------------------------|----------------------------|---|----------------------------|-------|------|
| | | | | A[10] | B[23] | C[5] |
| Li | 2.08 | 2.33 | 2.42 | 4.03 | 3.06 | 2.38 |
| B | 2.82 | 4.17 | 4.14 | 4.33 | 4.39 | 6.01 |
| C | 5.57 | 5.37 | 5.32 | 5.50 | 5.49 | 5.00 |
| O | 7.39 | 7.65 | 7.61 | 6.14 | 6.42 | 6.08 |
| F | 8.25 | 8.77 | 8.73 | 7.37 | 7.52 | 7.01 |
| Na | 2.37 | 2.13 | 2.16 | 3.58 | 2.91 | 2.30 |
| Al | 2.96 | 2.62 | 2.64 | 2.83 | 2.94 | 2.77 |
| Si | 3.58 | 3.30 | 3.31 | 3.55 | 3.61 | 3.38 |
| P | 5.20 | 3.92 | 3.93 | 5.93 | 5.42 | 4.86 |
| S | 4.68 | 4.52 | 4.53 | 4.08 | 4.28 | 4.12 |
| Ci | 5.20 | 5.10 | 5.11 | 4.79 | 4.91 | 4.70 |
| K | 1.96 | 1.71 | 1.74 | 2.84 | 2.35 | 1.92 |
| Br | 4.60 | 4.27 | 4.28 | 4.28 | 4.40 | 4.24 |

Electronegativities of several atoms calculated with the $X\alpha$ method having self-consistent parameter α have already been presented [12]. Table I shows the hardnesses of several atoms determined with parameters α_{SCF} . The values in column

A have been obtained from a spin polarized $X\alpha$ calculation [10] with parameter α_{HF} [22]. The column B presents the results of a local-spin-density calculation [23]. The experimental data obtained from experimental ionization energies and electron affinities using the expression (11) [5] can be seen in the column C. The results of the $X\alpha$ method with parameter α_{SCF} are in agreement with the other estimates for the hardness. Gazquez and Ortiz [11] have shown that the hardness of an open shell atom can be approximated as the one-quarter of the expectation value r^{-1} of the highest occupied atomic orbital:

$$\eta = \frac{1}{4} \langle r^{-1} \rangle_i. \quad (12)$$

Table I contains these expectation values determined with parameters α_{SCF} and α_{HF} [11]. They are really close to the other estimates for the hardness.

Parr and Pearson [5] have presented experimental hardnesses for several ions. Calculated values on hardness have not been published. Table II shows the results of the $X\alpha$ calculations with parameter α_{SCF} . We have obtained the Hartree-Fock results from [25] using Koopmans' theorem. Electronegativities of several positive ions have been calculated in the density functional theory [26]. Table II contains

Table II
Electronegativities, electron affinities and hardnesses of
a few ions (in eV)

| Ion | χ | A | η | χ | A | η | χ | A | η |
|------------------|--|-------|--------|--------------------------------|------|--------|------------------------------|-------|--------|
| | Calculated with ($X\alpha_{SCF}$ method) | | | Experimental values [5, 24] | | | Hartree-Fock results [25] | | |
| Al ⁺ | 11.85 | 5.14 | 6.71 | 12.4 | 6.0 | 6.4 | 11.75 | 5.7 | 6.05 |
| Al ²⁺ | 23.18 | 18.56 | 4.62 | 23.6 | 18.8 | 4.8 | 22.95 | 17.8 | 5.15 |
| Al ³⁺ | 75.26 | 27.80 | 47.46 | 74.2 | 28.4 | 45.8 | 75.05 | 28.1 | 46.95 |
| Ca ⁺ | 8.45 | 5.81 | 2.71 | 9.0 | 6.1 | 2.9 | 8.3 | 5.3 | 3.0 |
| Ca ²⁺ | 30.96 | 11.23 | 19.73 | 31.55 | 11.9 | 19.65 | 31.15 | 11.3 | 19.85 |
| Ca ³⁺ | 59.34 | 50.69 | 8.65 | 59.25 | 51.2 | 8.05 | 59.5 | 51.0 | 8.5 |
| Fe ⁺ | 11.71 | 8.00 | 3.72 | 12.05 | 7.9 | 4.15 | 13.04 | 7.0 | 6.04 |
| Fe ²⁺ | 24.64 | 15.43 | 9.21 | 23.45 | 16.2 | 7.25 | 27.44 | 19.07 | 8.36 |
| Fe ³⁺ | 43.37 | 33.85 | 9.52 | 42.75 | 30.7 | 12.05 | 46.1 | 35.8 | 10.3 |
| Br ⁺ | 16.78 | 11.69 | 5.09 | 16.7 | 11.8 | 4.9 | 17.6 | 12.4 | 5.2 |
| Br ²⁺ | 27.49 | 21.87 | 5.62 | 28.75 | 21.6 | 7.15 | 28.95 | 22.8 | 6.15 |
| Br ³⁺ | 39.15 | 33.11 | 6.04 | 41.6 | 35.9 | 5.7 | 40.75 | 35.1 | 5.65 |

electronegativities and electron affinities of a few ions calculated with the $X\alpha$ method using the parameter α_{SCF} and Slater's transition state method [6, 12]. There is good agreement between the calculated and experimental data.

In conclusion, it can be stated that this ab initio way of calculating the exchange potential within the $X\alpha$ method is useful to determine the electronegativities and hardnesses of atoms and molecules. It is a parameter-free method, there is no fitting or adjustment in it.

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