

**Modification of the Chew-Low Formula for the  $(\frac{3}{2}, \frac{3}{2}) \pi\text{-N}^{\circ}$  Resonance.**

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The theory of Chew and Low for  $\pi\text{-N}^{\circ}$  scattering results in a formula, containing two parameters, which may be used to fit the observed phase shifts. It is most useful in the case of the  $T = \frac{3}{2}, J = \frac{3}{2}$  state which contains the well-known resonance at approximately 200 MeV (lab.). The usual form (1) is

$$(1) \quad \frac{\eta^3}{\omega^*} \text{ctg } \alpha_{33} = \frac{3}{4f^2} \left( 1 - \frac{\omega^*}{\omega_0} \right),$$

where  $\eta = p/m_{\pi}c$  and  $p$  is the pion c.m. momentum,  $(mc^2\omega^*)$  is the  $\pi\text{-N}^{\circ}$  c.m. energy less one nucleon mass,  $\omega_0$  is the resonance value and  $f$  is the coupling constant. For  $\omega^* < \omega_0$  the straight line form of (1) agrees well with experiment, but for  $\omega^* > \omega_0$  there is serious disagreement, as one might expect for a theory derived in a low energy approximation. However, it is possible to obtain good agreement with the data by noting that (1) actually (2) should contain

$\eta^3[V(K)]^2$  in place of  $\eta^3$ . As explained by Wick

$$(2) \quad V(K) = \int \exp [i\mathbf{K} \cdot \mathbf{x}] \varrho(\mathbf{x}) d\mathbf{x},$$

where  $K$  is the wave number

$$\left( K = \eta \left( \frac{\hbar}{m_{\pi}c} \right) \right),$$

and  $\varrho(\mathbf{x})$  is the « extended source distribution » for the nucleon. In formula (1)  $V(K)=1$  and  $\varrho(\mathbf{x})=\delta(\mathbf{x})$ . Actually, in a convergent theory one must introduce a cutoff. Instead of the usual sharp cutoff at  $\omega^* = M/m_{\pi}$  ( $M$ =nucleon mass), I have considered the following:

$$(3) \quad \varrho(\mathbf{x}) = \frac{a^2 \exp [-ar]}{4\pi r}.$$

This is a Yukawa source distribution, and yields a smooth cutoff for  $V(K)$  or  $V(\eta)$ :

$$(4) \quad V(K) = \frac{1}{1 + K^2 a^{-2}}.$$

Introducing this into (1) it is possible to fit the data with  $a^{-1}=0.38$  fermi

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(1) G. PUPPI and A. STANGHELLINI: *Nuovo Cimento*, **5**, 1305 (1957).

(2) G. C. WICK: *Rev. Mod. Phys.*, **27**, 339 (1955).

( $0.27(\hbar/m_\pi c)$  or  $1.8(\hbar/Mc)$ ). The other constants are approximately in agree-

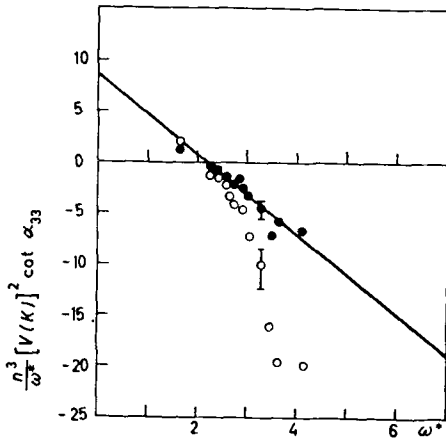


Fig. 1. - A Chew-Low plot for two values of  $V(K)$  is shown above. The points are from the data of references (\*\*). The straight line is for  $f^2 = 0.087$ ,  $\omega_0 = 2.17$ .

(\*) B. PONTECORVO: *Ninth Int. Annual Conf. on High Energy Physics*, p. 108.

(\*) J. H. FOOTE, O. CHAMBERLAIN, E. H. ROGERS, H. M. STEINER, C. WIEGAND and T. YPSILANTIS: *Phys. Rev. Lett.*, **4**, 30 (1960).

(\*) W. J. WILLIS: *Phys. Rev.*, **116**, 753 (1959).

(\*) W. D. WALKER, J. DAVIS and W. D. SHEPARD: *Phys. Rev.*, **118**, 1612 (1960).

ment with those of JOHNSON and CAMAC (?):

$$f^2 = 0.087,$$

$$\omega_0 = 2.17.$$

Since the Chew-Low theory is for a non-relativistic static nucleon problem and relies on the assumption that a complicated integral expression (8) can be approximated as a constant effective range (or equivalent  $\omega_0$ ), the agreement achieved with (4) may be fortuitous. Nevertheless it is interesting to note that the simple application of a more « realistic » nucleon source of approximately nucleonic dimensions brings theory and experiment together, as shown in the figure.

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(\*) G. F. CHEW and F. E. LOW: *Phys. Rev.* **101**, 1570 (1956).