

Higher-Dimensional Bianchi Type-VI_h Cosmologies

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Received 1985 February 18; accepted 1985 May 28

Abstract. We present the higher-dimensional perfect fluid generalizations of the extended Bianchi type-VI_h vacuum space-times discussed recently by Demaret & Hanquin (1985). It is shown that the Chodos-Detweiler mechanism of cosmological dimensional-reduction is possible in these cases.

Key words: cosmology, Bianchi type—cosmology, higher dimensional

1. Introduction

The topic of higher-dimensional cosmologies has recently received much attention (Chodos & Detweiler 1980; Freund 1982; Sahdev 1984; Lorenz-Petzold 1984a, b; 1985a-d). It has been suggested that such cosmologies play an important role in the very early universe (Chodos 1984; Chodos & Myers 1984; Appel & Dresden 1984). In the course of time, the influence and effect of extra dimensions $D \geq 1$ ($N=1+d+D$) must have diminished or changed. An easy way to visualize this process is that in the course of time, the usual $d = 3$ dimensions kept on expanding, while the extra dimensions $D \geq 1$ contracted very rapidly to distances of the order of the Planck length. Such $MN=|R \otimes M^d \otimes M^D$ space-time models were first constructed by Chodos & Detweiler (1980) (see also Belinskii & Khalatnikov 1972; Forgacs & Horvath 1979; Barrow 1983; Tosa 1984; Mann & Vincent 1985).

An interesting paper in this respect has been recently published by Demaret & Hanquin (1985). Most of the studies done so far are of the higher-dimensional generalization of the anisotropic (1 + 3)-dimensional Bianchi type-I model (see Lorenz-Petzold 1985c). However, no generalization of the Petrov-Bianchi classification exist in (1 + n)-dimensions ($n = N - 1$) (Petrov 1969; Sahdev 1984). A simple definition was given by Demaret & Hanquin (1985). A N -dimensional spatially homogeneous model is defined as a N -dimensional space-time possessing a ($N - 1$)-dimensional group of isometry acting simply transitively on ($N - 1$)-dimensional space-like hypersurfaces. This definition does not probably recover all possible cases of spatially homogeneous models in (1+3)-dimensions (Ryan & Shepley 1975; Kramer *et al.* 1980). However, the definition given by Demaret & Hanquin is already sufficiently rich as to recover a large set of spatially homogeneous models.

In this paper we discuss the higher-dimensional perfect fluid field equations of a generalization of the (1 + 3)-dimensional Bianchi type-VI_h space-time, which includes also the Bianchi type-V and the Bianchi type-III as special cases. Only the vacuum case was discussed by Demaret & Hanquin (1985).

2. Field equations and solutions

We consider the metrics

$$ds^2 = -dt^2 + R_1^2(dx^1)^2 + \exp(2p_i x^i) R_i^2(dx^i)^2, \quad (1)$$

where $R_j = R_j(t)$ ($j = 1, \dots, N-1$; $n = N-1$) are the cosmic scale functions and $P_i = \text{const}$. The corresponding perfect fluid field equations to be solved are given by

$$H'_1 + H_1(\ln u)' - \sum p_i^2 = \varepsilon(2-\gamma)R_1^2/(N-2), \quad (2a)$$

$$H'_i + H_i(\ln u)' - p_i \sum p_i = \varepsilon(2-\gamma)R_i^2/(N-2), \quad (2b)$$

$$H_1 \sum p_i - (\ln g)' = 0, \quad (2c)$$

$$(\ln u)'^2 + 2(\ln u)'H_1 - \sum H_i^2 - \sum p_i^2 - (\sum p_i)^2 = 2\varepsilon R_1^2, \quad (2d)$$

$$\varepsilon' + \gamma\varepsilon(\ln R_1 u)' = 0, \quad (2e)$$

$$p = (\gamma-1)\varepsilon, \quad 1 \leq \gamma \leq 2, \quad (2f)$$

where $H_j = (\ln R_j)'$ are the Hubble parameters, $u = R_2 \dots R_{N-1}$, $g = R_2^{p_2} \dots R_{N-1}^{p_{N-1}}$, $dt = R_1 d\eta$, $()' = d/d\eta$, $i = 2, \dots, N-1$ and ρ and ε are, respectively, the pressure and energy density of the perfect fluid matter. In cosmology, we are mainly interested in the cases $\varepsilon = 0$ (vacuum), $\gamma = 1$ (dust), $\gamma = 2$ (stiff matter) and $\gamma = N/(N-1)$ (radiation). From Equation (2e) we obtain the conservation equation $\varepsilon = M(R_1 \dots R_{N-1})^{-\gamma}$, $M = \text{const}$.

We first consider the cases $\varepsilon = 0$ and $\gamma = 2$. From the linear combination of Equations (2b) we obtain

$$u'' - p^2 u = 0, \quad p = \sum p_i, \quad (3)$$

with the general solutions

$$(i) \quad u = \sinh p\eta, \quad p \neq 0, \quad (4a)$$

$$(ii) \quad u = \eta, \quad p = 0, \quad (4b)$$

and the special solution

$$(iii) \quad u = \exp p\eta, \quad p \neq 0. \quad (4c)$$

Equations (2a) and (2b) can now be easily integrated (Equation 2d is nothing but the constraint equation on the constants of integration) to give:

$$(i) \quad R_1 = (\sinh p\eta)^{(1/p^2)\sum p_i^2} \left(\tanh \frac{p}{2}\eta \right)^{(1/p)\sum c_i p_i},$$

$$R_i = (\sinh p\eta)^{p_i/p} \left(\tanh \frac{p}{2}\eta \right)^{c_i},$$

$$\sum c_i = 0, \quad p^2 + \sum p_i^2 = p^2 \sum c_i^2 + 2M, \quad (5a)$$

$$(ii) \quad R_1 = \eta^a \exp\left(\frac{1}{4}\sum p_i^2\right),$$

$$R_i = \eta^{c_i},$$

$$\sum c_i = 1, \quad \sum c_i p_i = 0, \quad 2a = \sum c_i^2 - 1 + 2M, \quad (5b)$$

$$(iii) \quad \begin{aligned} R_1 &= \exp\left(\frac{1}{p} \sum p_i^2 \eta\right), \\ R_i &= \exp p_i \eta, \quad M = 0, \end{aligned} \quad (5c)$$

where $c_j, a = \text{const.}$

The stiff matter solutions ($M \neq 0$) are new. By setting $M = 0$ we rediscover the vacuum solutions ($p = 1$, if $p \neq 0$ without loss of generality) discussed by Demaret & Hanquin (1985). The four-dimensional solutions (5a) ($i = 2, 3$) are due to Collins (1971) (in a somewhat different form) and Ruban (1978a) (see also Wainwright, Ince & Marshman 1979), which include the Bianchi type- VI_h vacuum solutions of Ellis & MacCallum (1969) and the Bianchi type-V solution given by Joseph (1966) ($p_i = 1$). The special solution (5b) has no analogue in (1 + 3)-dimensions. Our solution (5c) has been discussed in four-dimensions by Lifshitz & Khalatnikov (1963a, b), Collins (1971, 1977), Evans (1974, 1978), Siklos (1981), Ruban (1978a, b), Siklos (1981), Ruban, Ushakov & Chernin (1981), Belinskii, Khalatnikov & Lifshitz (1982), Wainwright (1983) and Wainwright & Anderson (1984) (see also the recent paper by Rosquist & Jantzen 1985).

We next consider the special power law perfect fluid solutions:

$$R_1 = \exp a\eta, \quad R_i = \exp c_i \eta, \quad (6)$$

where $a, c_i = \text{const.}$ It follows that

$$\begin{aligned} a &= c\gamma/(2 - \gamma), \quad \gamma \neq 2, \quad c = \sum c_i, \\ c_i &= a + [\gamma/a(2 - \gamma)][pp_i - \sum p_i^2], \end{aligned} \quad (7a)$$

and thus a can be re-expressed in terms of p_i :

$$a^2 = \gamma^2 [p^2 - (N - 2) \sum p_i^2] / [(2 - \gamma)(2 - \gamma(N - 1))], \quad (7b)$$

which yields in addition to $\gamma \neq 2$ the condition $\gamma \neq 2(N - 1)$. According to Equation (2c) we have also

$$(i) \quad \begin{aligned} p &\neq 0: \\ a &= \sum c_i p_i / p, \end{aligned} \quad (8a)$$

$$(ii) \quad \begin{aligned} p &= 0: \\ \sum c_i p_i &= 0, \end{aligned} \quad (8b)$$

and the value of M is determined by Equation (2d).

Our solutions (6–8) are the higher-dimensional special perfect fluid generalizations of the (1 + 3)-dimensional solutions first given by Collins (1971). Such solutions are of much interest in cosmology and have been discussed recently in (1 + 3)-dimensions by Ruban, Ushakov & Chernin (1981), Barrow (1982, 1984), Wainwright (1983, 1984) and Wainwright & Anderson (1984) (see also the recent paper by Lorenz-Petzold 1985e). By setting $a = c_2 = c_3, c_4 = b$, we obtain the isotropic solution in (1+4)-dimensions with

$$b = a(2 - 3\gamma)/\gamma, \quad a > 0, \quad (9)$$

leading thus to a contraction of the physical unobservable dimension and an expansion of the three physical dimensions.

We finally consider the space-times defined by $p_i = 1$, which are the higher-dimensional generalization of the (1 + 3)-dimensional Bianchi type-V cosmologies. Exact solutions (besides the ones given by Equation 5) can be obtained in the cases $\gamma = N/(N-1)$ (radiation) and $\gamma = 1$ (dust). Without going into all the details we present only the results (see Lorenz-Petzold 1985f for derivation):

$$\begin{aligned} u &= a \sinh^2 q\eta + b \sinh 2q\eta, & (10a) \\ R_1 &= u^{1/(N-2)} \\ R_i &= u^{1/(N-2)} [1 + (2b/a) \coth q\eta]^{c_i/(N-2)}, \\ a &= 2M/(N-2)(N-1), \quad q = (N-2)/2, \\ \sum c_i &= 0, \quad \sum c_i^2 = (N-1)(N-2), & (10) \end{aligned}$$

where $b, c_i = \text{const.}$ By setting $i = 2, 3$, we rediscover the (1 + 3)-dimensional Bianchi type-V radiation solution first given by Ruban (1977a, b) (Note that these papers are not quoted by Kramer *et al.* 1980).

The dust case can be reduced to the solution of a generalized Friedmann equation:

$$\begin{aligned} R'^2 &= R^2 + (a^2/(N-1)(N-2)) \sum c_i^2 R^{2(1-N)} + (2M/(N-1)(N-2))R^{5-N}, \\ S &= \exp[a \int R^{1-N} d\eta], \end{aligned}$$

where $a = \text{const.}$, $R = R_1$, $R_i = RS^{c_i}$, $\sum c_i = 0$, $c_i = \text{const.}$ The case $N = 4$ is nothing but the (1 + 3)-dimensional Bianchi type-V dust solution first given by Schucking & Heckmann (1958), which can be solved explicitly in terms of elliptic functions. This completes our study of the higher-dimensional Bianchi type-VI_h perfect fluid field equations. According to the relations $\sum c_i = \text{const.}$, it is always possible to construct solutions such that $d = 3$ -dimensions are expanding while the extra dimensions $D \geq 1$ are contracting to an unobservable scale.

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