# **Phase Diagram Investigation and Proposition of a Thermodynamic Evaluation of the Tin-Selenium System**

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*(Submitted December 1, 1995)* 

**Differential scanning calorimetry (DSC) and x-ray diffraction (XRD) measurements provide phase diagram data in the whole composition range. From the results and literature data, an optimization and calculation study was done to generate the thermodynamic functions of each phase.** 

## **Introduction**

According to the assessment of [86Sha], the phase diagram for the Sn-Se binary system shows two intermediate stoichiometric compounds, SnSe with low ( $\alpha$ SnSe) and high ( $\beta$ SnSe) temperature forms and  $SnSe<sub>2</sub>$ , and a monotectic invariant on the Sn-rich side. Whereas, the general shape of the phase diagram seems to be established, some problems remain concerning the melting points of the compounds and the order of transition in SnSe.

Liquid thermodynamic data available in the literature show a large negative departure from ideality with a characteristic shape indicating the presence of short-range order. A thermodynamic evaluation of the system is presented using the optimization and calculation program of [77Luk] and [92Luk].

# **Experimental Investigation**

## **Experimental** Details

Twenty-eight alloys were carefully weighed from pure Sn (Fluka, 6N, Fluka, L'Isle d'Abeau Chesnes BP 701, F-38297 St-Quentin Fallavier Cedex, France) and pure Se (Aldrich, 5N, Sigma-Aldrich-Chimie, L'Isle d'Abeau Chesnes BP 701, F38297 St-Quentin Fallavier Cedex, France). Nonsoluble impurities were removed from Sn by filtration under purified argon on silica wool. Because the high vapor pressure of Se voids preparation of alloys by direct melting, special attention was paid to the procedure for alloy preparation.

To prevent oxidation, the pure elements were placed in a silica tube (5 mm inside diameter, 6 mm outside diameter, 25 mm length), which was then sealed under vacuum  $(10^{-1}$  Pa). The total weights mixed were approximately 0.200 g.

The alloys were then subjected to the following heat treatment in a furnace that had the capacity of controlling temperatures

up to 1200  $\degree$ C and maintaining temperature uniformity to  $\pm 0.5$ °C: (1) 24 h at 250 °C, (2) slow heating until 500 °C followed by 24 h annealing, (3) slow heating until 900 °C ( $\approx$  20 °C above the temperature of fusion of SnSe) for alloys with  $X_{S_e} < 0.55$  or until 700 °C ( $\approx$  50 °C above the temperature of fusion of SnSe<sub>2</sub>) for alloys with  $X_{\text{Se}}$  < 0.55, (4) to remove any possible remaining condensed vapor, the alloys were then introduced into a Netzsch DTA apparatus (Netzsch-Frères, 32-34 Avenue des Chardons, F-77340 Pontault-Combault, France) equipped with a "low head" to obtain a gradient of temperature, and  $(5)$ 24 h annealing at 500  $^{\circ}$ C.

DSC measurements were subsequently carried out on a DSC 121 (Setaram, 7 rue de l'Oratoire, BP 34, F-69641 Caluive Cedex, France)and a multi-HTC (Setaram) whose main characteristics are listed in Table 1. A heating rate of  $1 °C/min$  was used except for alloys with  $0.55 \leq X_{\text{Se}} \leq 0.667$  where the eutectic and the liquidus temperatures were too close together. In those cases, to get a good differentiation between the thermal arrests, the heating rate was lowered to  $0.1 \degree C/\text{min}$ . Calibration of the DSC apparatus was made by checking the temperatures of fusion of In, Sn, Te, and Sb for the DSC 121 and Ag and Ge for the multi-HTC.

Identification of phases was made with powder diffraction measurements on a PW 1729 (Philips Electronique Industrielle S.A.S., Analyse Rayons X, 22 Avenue Descartes, F-94454) x-ray diffractometer (equipped with goniometer driven

**Table 1 Characteristics of DSC Measurements** 

<b>Characteristics</b>	<b>DSC 121</b>	<b>Multi-HTC</b>		
Temperature range	$-120$ to 830 °C	20 to $1600^{\circ}$ C		
	$5$ to $15 \mu V$	$500 \mu V$		
Heating and cooling rate	$0.01$ to 30 $^{\circ}$ C/min	$0.01$ to $30^{\circ}$ C/min		

X <sub>Se</sub>	Temperature, <sup>o</sup> C					<b>Identification of phases</b>	
at.%	<b>Liquidus</b>	<b>Solidus</b>		<b>Invariants</b>		at room temperature	
	775.6	$\cdots$	231.7	526.7		$\beta$ Sn + $\alpha$ SnSe	
	822.7	$\cdots$	231.7	526.8		$\beta$ Sn + $\alpha$ SnSe	
	831.4	$\sim$ $\sim$ $\sim$	231.6	526.7		$\beta$ Sn + $\alpha$ SnSe	
	$\ldots$	$\sim$ $\sim$ $\sim$	231.5	526.0	832.5	$\beta$ Sn + $\alpha$ SnSe	
	$\cdots$	$\cdots$	231.5	526.0	832.0	$\beta$ Sn + $\alpha$ SnSe	
	$\cdots$	$\cdots$	231.5	526.0	832.1	$\beta$ Sn + $\alpha$ SnSe	
45.0	$\cdots$		231.7	526.5	831.7	$\beta$ Sn + $\alpha$ SnSe	
	865.5	$\ldots$	231.6	526.1	831.8	$\beta$ Sn + $\alpha$ SnSe	
	871.5	524.3	231.7	831.6		$\beta$ Sn + $\alpha$ SnSe	
	$\cdots$	521.1	872.5			$\alpha$ SnSe	
50.5	870.9	521.0	629.1			$\alpha$ SnSe + SnSe <sub>2</sub>	
51.0	865.0	521.1	629.4			$\alpha$ SnSe + SnSe,	
53.0	$\cdots$	$\cdots$	520.4	628.9		$\alpha$ SnSe + SnSe <sub>2</sub>	
	793.6	$\cdots$	520.6	629.5		$\alpha$ SnSe + SnSe,	
	770.5	$\cdots$	519.7	629.9		$\alpha$ SnSe + SnSe <sub>2</sub>	
	744.4	$\cdots$	520.0	628.7		$\alpha$ SnSe + SnSe <sub>2</sub>	
	722.0	$\ddotsc$	520.0	628.7		$\alpha$ SnSe + SnSe <sub>2</sub>	
59.0	690.0	$\cdots$	518.9	629.5		$\alpha$ SnSe + SnSe <sub>2</sub>	
	660.0	$\cdots$	519.6	629.6		$\alpha$ SnSe + SnSe,	
	$\cdots$	$\ddotsc$	520.0	628.9		$\alpha$ SnSe + SnSe <sub>2</sub>	
	640.0	$\cdots$	520.0	629.5		$\alpha$ SnSe + SnSe <sub>2</sub>	
	643.6	$\cdots$	629.4			$\alpha$ SnSe + SnSe,	
	$\cdots$	$\cdots$	645.0			SnSe <sub>2</sub>	
	642.4	$\cdots$	219.9			$Se + SnSe2$	
	635.0	$\cdots$	219.9			$Se + SnSe2$	
80.0	624.2	$\cdots$	220.0			$Se + SnSe$	
	613.8	$\cdots$	219.7			$Se + SnSe2$	
90.0	599.3	$\cdots$	219.9			$Se + SnSe2$	

**Table 2 Experimental Phase Diagram Data Obtained in This Work** 

by a software developed by [95Fra]) using the  $K_{\alpha_1}$  radiation of copper ( $\lambda_{\text{CuK}_{\alpha}} = 1.5405~\text{\AA}$ ).

## **Results**

#### *Liquidus and Invariant Temperatures, Enthalpy of Fusion, Temperature*

Table 2 gives the temperatures of both liquidus and invariants, whereas Table 3 presents the measured temperatures and enthalpies of fusion of the compounds available in the literature. The melting temperature obtained for the intermetallic compound  $\beta$ SnSe is 1145.65 K, which agrees well with the most recently reported value [94Yam]. The present value for the enthalpy of fusion,  $\Delta_{\text{fus}}H = 14 \pm 1$  kJ/g-atom, is 15 to 20% lower than the reported values of [81Bal] and [94Yam], which were obtained form heat content measurements. Both of those investigations used the drop calorimetric method and thus worked with open tubes. Because Se is a volatile element, some evaporation would yield an overestimation of the heat effect.

The melting temperature obtained for  $SnSe<sub>2</sub>$  is 918.15 K with  $\Delta_{\text{fus}}H = 17 \pm 1$  kJ/g-atom, which is much lower than the value given by [82Ale], who gave no information on the method for this determination.

## *Invariants*

The calorimetric study of the invariants led to the invariant equilibria listed in Table 4.

#### *Phase Transition in SnSe*

According to the assessment of [86Sha] and as shown in Table 5, all workers agree with the presence of a phase transformation at  $X_{\text{Se}} = 0.500$ . There is conflicting evidence with regard to the phase transition in SnSe, and doubt remains as to the order of this transition. The works of [81Bal] and [94Yam] reveal the presence of a small latent heat at the transition temperature, which seems to indicate a first-order transition in the sense of Ehrenfest and Landau. On the contrary, the continuous variation of the primary order parameter with temperature obtained by neutron diffraction measurements [86Cha] is more typical of a second-order transition. Still, the sharp decrease of this order parameter in the transition temperature range may be interpreted as a first-order transition with order parameter [67Lan, 81Wey]. These data have to be compared with the structural results of [81Sch], who concluded from their XRD measurements that there is an indication of a possible first-order transition near the critical temperature for the  $\lambda$ -type phase transition.

We found invariant reactions on both sides of the composition,  $X_{\text{Se}} = 0.500$  at different temperatures. Verification was made

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#### **Table 4 Sn-Se Invariant Equilibria**



by checking the variation of the invariant temperatures with different heating rates (Table 6).

The small variations in temperature with different heating rates are within the instrumental error range. Thus, two possibilities arise as to the representation of the phase diagram in the vicinity of  $X_{\text{Se}} = 0.500$ . Figure 1(a) shows first-order transition of the SnSe phase involving metatectic and eutectoid invariants. Figure l(b) shows reaction between two different phases involving peritectic and eutectoid invariants.

As shown in Table 2, a thermal arrest was found with slightly decreasing temperature from 524.3 to 521.0  $^{\circ}$ C when  $X_{Se}$  increases from 0.495 to 0.510. This temperature variation seems difficult to interpret in Fig.  $1(a)$ , but may appear as the Sn-rich border of the two-phase field  $\alpha$ SnSe +  $\beta$ SnSe in Fig. 1(b).

This interpretation agrees well with the results of [63Dem] and [81Bal], who gave a range of temperature for the transition instead of a definite value. We believe we cannot deal with a phase transition, but with two different phases involving peritectic and eutectoid invariants (according to Fig. lb), whose compositions are, respectively,  $0.490 < X_S<sup>pl</sup> < 0.495$ and  $0.510 < X_{\text{Se}}^{23} < 0.530$ .

The composition range of both phases,  $\alpha$ SnSe and  $\beta$ SnSe, are approximately 2 to 3% at the temperature of the invariants. From XRD patterns, we calculated the parameters of the or-

#### **Table 5 Transition Temperature and Transition Order**  in αSnSe to βSnSe



thorhombic cell of  $\alpha$ SnSe for three compositions,  $X_{\text{Se}} = 0.450$ , 0.500, and 0.550. The values observed are identical within the experimental error. Thus, a significant composition range for this phase at room temperature could not be established.

#### *Phase Sn2Sea*

The DSC curves carried out at composition  $X_{\text{Se}} = 0.600$  did not show any more thermal arrest than those related to the invariants, E2, E3, and the liquidus. Phase  $Sn<sub>2</sub>Se<sub>3</sub>$  does not undergo a phase transition. XRD patterns in this composition range, showing only the presence of the two phases,  $\alpha$ SnSe and  $SnSe<sub>2</sub>$ , confirm the results of differential calorimetry. This result agrees well with the conclusion of the thermal, metallographic, and XRD study [06Pel, 66Kar], the Mössbauer works  $[70Bol, 71Bar]$ , the nuclear  $\gamma$ -resonance investigation [71Mel], and the DTA contribution [76Rou], which did not include this compound.  $Sn<sub>2</sub>Se<sub>3</sub>$  may possibly exist in a metastable state in thin films, as reported by [54Pal] from XRD results.

with <b>itemating</b> ituto					
<b>Heating rate,</b>	Invariant temperature, <sup>o</sup> C				
$^{\circ}$ C/min	Alloy $X_{S_0} = 0.490$	Alloy $X_{S_0} = 0.580$			
	526.5	520.0			
	526.0	520.0			
	526.1	519.6			

**Table 6 Comparison of Sn-Se Invariant Temperatures with Heating Rates** 

# **Optimization and Calculation**

## **Phase Diagram Data**

In earlier works, [06Pel] and [09Pel] claimed the existence of SnSe and  $SnSe<sub>2</sub>$  and gave some liquidus points and melting temperatures of these two compounds. The first complete shape of the liquidus was established by [09Bil], whose interpretation points out the existence of compound  $Sn<sub>2</sub>Se<sub>3</sub>$ . Later, using hardness measurements and crystal growth, [38Jen] found a solubility of Se in  $\beta$ Sn of 0.05 at.% at 200 °C. From electrical measurements, [63Alb] gave the composition limits of stability at the Se-rich side for  $\alpha$ SnSe as 50.005 at.% at 300  $^{\circ}$ C and 50.01 at.% at 400  $^{\circ}$ C. A solubility curve of Se in Sn for 300 to 550  $\degree$ C is proposed by [64Vas], who used a diffusion method. The first description of the whole phase diagram including the monotectic invariant was given by [66Kar], who used metallography, XRD, and DTA investigations. More recently, from measurements on enthalpy as a function of composition and DSC, [87Kot] confirmed the general shape of the description given by [66Kar]. The melting temperature of the two compounds are presented in the section, "Experimental Investigation" (Table 3).

## **Enthalpy Data**

The enthalpies of mixing in the liquid state were studied by [87Kot], who reported results for six temperatures from 921 to 1243 K. As explained by these authors, the large negative enthalpy of formation (-45 kJ/g-atom at  $T = 1175$  K) located at  $X_{S_{\rm e}}$  $= 0.5$  indicates the presence of short-range order based on  $\beta$ SnSe.

Heat content measurements on alloys of composition  $X_{Se}$  = 0.500 carried out by different authors [74Bla1, 81Bal, 91Ras, 94Yam] are in good agreement for the low-temperature phase. For the high-temperature modification, the slope of the curve  $H(T) - H(298.15) = f(T)$  given by [91Ras] is less pronounced than those of other authors. Only [77Kul] provided heat content results for  $SnSe<sub>2</sub>$  but with no indication of the experimental method used.

In a careful study, [81Wie] determined the heat capacities of both compounds in the temperature range from 230 to 580 K using DSC. More recently, [88Ras] reported a few  $C_p$  values for  $\alpha$ SnSe in the range from 300 to 800 K, which show a bigger temperature dependence than the measured values of [81Wie]. Because [88Ras] did not give any experimental information, these data were not used in this optimization.

The enthalpies of formation of both compounds,  $\alpha$ SnSe and  $SnSe<sub>2</sub>$ , have been measured by different techniques: direct combination calorimetry [60Haj, 70Haj, 92Boo], Knudsen effusion [64Col], dissociation pressure [67Kar], emf measure-



ments  $[71$ Mell for  $\alpha$ SnSe and direct combination calorimetry [92Boo], equilibrium data [67Rau], dissociation pressure [67Kar,  $77$ Kul] for  $\beta$ SnSe. If the values are in relatively good agreement for  $\alpha$ SnSe, the enthalpy of formation of SnSe, coming from dissociation pressure is 20% higher [67Kar] or 12% lower [77Kul] than the average value of the other authors.

## **Gibbs Energy Data**

Because of the high temperature of fusion of  $\alpha$ SnSe and the high vapor pressure of Se at high temperature, the only available chemical potential data were determined in the solid state by [71Mel] from emf measurements.

## **Previous** Calculations

By applying a regular associated treatment to the melts, [88Kot] and [89Cas] showed that the liquid is better described with the use of one associate ( $\beta$ SnSe) instead of two ( $\beta$ SnSe +  $SnSe<sub>2</sub>$ ). The standard enthalpy and entropy of formation are 45.5 kJ/g-atom and 4.15 J/K.g-atom, respectively.

## Analytical **Description of the Phase** Stabilities

## *Pure Elements*

According to the Scientific Group Thermodata Europe (SGTE), the temperature dependence of the molar Gibbs energy of the pure stable elements, referred to standard state, is given by the following expression:

$$
G_i^0 - H_i^{SER} = A + BT + CT \ln T + DT^2 + ET^{-1} + FT^3
$$
  
+ 
$$
IT^7 + JT^{-9}
$$
 (Eq 1)

The coefficients of this equation, available from the database of [91Din], are given in the appendix.  $H_1^{SER}$  is the enthalpy of the pure element  $i$  at 298.15 K in its stable state.

## *Liquid Phase*

As discussed in the section, "Enthalpy Data," the liquid phase was described using the association model presented by [82Soml] and [82Som2]. Then the Gibbs energy (given for one mole of formula units) is expressed as the sum of a refer-

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ence part, G<sup>ref</sup>, formation of the associate part,  $\Delta_f G$ , ideal part,  $G<sup>id</sup>$ , and an excess part,  $G<sup>ex</sup>$ :

$$
Gliq - HSER = Gref + \DeltafG + Gid,liq + Gex,liq
$$
 (Eq 2)

with:

$$
G^{\text{ref}} = \left\{ X_{\text{Sn}} \cdot \left[ G_{\text{Sn}}^{0,\text{liq}}(T) - H_{\text{Sn}}^{\text{SER}}(298.15) \right] \right\} + X_{\text{Se}}
$$

$$
\left[ G_{\text{Se}}^{0,\text{liq}}(T) - H_{\text{Se}}^{\text{SER}}(298.15) \right] \cdot (y_{\text{Sn}} + 2y_{\text{SnSe}} + y_{\text{Se}}) \left( \text{Eq } 3 \right)
$$

$$
\Delta_{\mathbf{f}} G = y_{\mathbf{Sn}} \cdot G_{\mathbf{Sn}}^{\mathbf{0},\text{liq}} + y_{\mathbf{SnSe}} \cdot G_{\mathbf{SnSe}}^{\mathbf{0},\text{liq}} + y_{\mathbf{Se}} \cdot G_{\mathbf{Se}}^{\mathbf{0},\text{liq}} \tag{Eq 4}
$$

$$
G^{\text{id},\text{liq}} = R \cdot T \cdot z \cdot (y_{\text{Sn}} \cdot \ln y_{\text{Sn}} + y_{\text{SnSe}} \cdot \ln y_{\text{SnSe}} + y_{\text{Se}} \cdot \ln y_{\text{Se}})
$$
\n
$$
(Eq 5)
$$

$$
G^{\text{ex,liq}} = L^{\text{0,liq}}_{\text{Sn,SDSe}} \cdot y_{\text{Sn}} \cdot y_{\text{SnSe}} + L^{\text{0,liq}}_{\text{SnSe,Se}} \cdot y_{\text{SnSe}} \cdot y_{\text{Se}} + L^{\text{1,liq}}_{\text{Sn,SnSe}}
$$

$$
\cdot y_{\text{Sn}} \cdot y_{\text{SnSe}} \cdot (y_{\text{Sn}} - y_{\text{SnSe}}) + L_{\text{SnSe},\text{Se}}^{1,\text{lig}} \cdot y_{\text{SnSe}} \cdot y_{\text{Se}} \cdot (y_{\text{SnSe}} - y_{\text{Se}})
$$

+ 
$$
L_{\text{Sn,SnSe}}^{2, \text{lig}} \cdot y_{\text{Sn}} \cdot y_{\text{SnSe}} \cdot (y_{\text{Sn}} - y_{\text{SnSe}})^2 + L_{\text{SnSe,Se}}^{2, \text{lig}} \cdot y_{\text{SnSe}} \cdot y_{\text{Se}}
$$
  
\n $\cdot (y_{\text{SnSe}} - y_{\text{Se}})^2$  (Eq 6)

where  $T$  is the temperature; Sn, SnSe, and Se are the three species in the liquid;  $y_i$  are their site fractions; and z is the number of lattice sites in one mole of formula units:  $z = y_{Sn} + y_{SnSe} +$  $y_{\text{Se}}$ . The  $L_{\text{L}}^{\text{y}}$  are polynomial interaction terms described with the Redlich-Kister formalism [48Red], which can be interpreted as a modification of the polynomial description used by Margules [1895Mar]. The terms  $G_t^{0,\text{liq}}$  and  $L_t^{y,\text{liq}}$  are expressed as functions of temperature from Eq 1.

#### *Solid Phases*

As discussed in the section, "Experimental Investigation," the width of  $\alpha$ SnSe is too small at room temperature to define significant phase limits. The reported values of [63Alb] (see section, "Phase Diagram Data") confirm the narrow range of homogeneity of this phase, but the experimental data are insufficient to enable modeling this phase as nonstoichiometric. No data were reported for the composition limits of  $\beta$ SnSe and  $SnSe<sub>2</sub>$ . Thus  $\alpha SnSe$ ,  $\beta SnSe$ , and  $SnSe<sub>2</sub>$  were treated as stoichiometric compounds. The Gibbs energy of one mole of formula units is expressed as:

$$
G^{\text{Sn}}_{p}S_{q} - p \cdot H_{\text{Sn}}^{\text{SER}} - q \cdot H_{\text{Se}}^{\text{SER}} = A + BT + CT \ln T + DT^{2}
$$

$$
+ ET^{-1} + FT^{3} \qquad (\text{Eq 7})
$$

where  $p$  and  $q$  denote the stoichiometric numbers.

The coefficients were fitted to the experimental data (384 data were used) by a least squares technique described and programmed (BINGSS) by [77Luk]. Table 7 gives the data sources used in the final data set.

## **Results and Discussion**

One of the problems encountered in the optimization was the definition of the stoichiometry of the two solid phases  $\alpha$ SnSe and  $\beta$ SnSe. According to our experimental results, we made the following approximations.

Since the widths of these phases at the invariant temperatures extend on both sides of the composition  $X_{\text{Se}} = 0.500$ , we chose the stoichiometric composition  $Sn_{0.5}Se_{0.5}$  for both  $\alpha$  and  $\beta$ .

Then came the problem of the phase transformation temperature. With the assumption of the same composition for  $\alpha$  and  $\beta$ , it is not possible, from a thermodynamic point of view, to state two different temperatures for the phase transition. Thus we

<b>Function</b>	A	B	С	D	E	
$G_{\text{SnSe}}^{0,\text{lig}}$	$-134591.86$	278.59972	$-32.99$			
	43 462.43	$-21.60740$				
$L_{\text{SnSe.Se}}^{0,\text{lig}}$	15 432.80	$-19.06391$				
	$-1032.50$					
$L_{\text{SnSe.Se}}^{1, \text{lig}}$	$-895.55$					
$G(\alpha SnSe) - H_{Sn}^{SER} - H_{Se}^{SER}$	$-63$ 171.87	134.4964	$-26.31$	$-1.82$	61 422.50	0.11076
$G(\beta$ SnSe) – $H_{\text{Sn}}^{\text{SER}}$ – $H_{\text{Se}}^{\text{SER}}$ …………	$-62529.58$	133.69002	$-26.31$	$-1.82$	61 422.50	0.11076
$G(Sn_1Se_2) - H_{Sn}^{SER} - 2 \cdot H_{Se}^{SER}$	$-46104.05$	121.16205	$-24.84$	$-1.64$	39 347.00	0.1185733

**Table 8 Optimized Coefficients According to the Analytical Description of the Phases** 

**Note: Functions are** expressed in J/g-atom.



defined a hypothetical transformation temperature as  $T_{trans} =$  $(T_{\rm P1} + T_{\rm E3})/2$   $(T_{\rm trans} = 523.35$  °C).

Heat capacity data are available [81Wie] only for  $\alpha$ . To get homogeneous descriptions of the Gibbs functions for both phases,  $C_p$  coefficients of  $\beta$  were constrained during the optimization to be the same as for  $\alpha$ .

#### The optimized coefficients are given in Table 8.

The associate model defines a fixed  $C_p$  that just depends on the definition of the associate. This value is intrinsic and cannot be adjusted. To fit the results found for the enthalpy of mixing [87Kot] for various temperatures, an extra  $C_p$  term was optimized in the description of the Gibbs energy of formation of the associate. We always got a significantly lower mean square of errors by removing the data obtained for the highest temperature (1243 K). The high vapor pressure of Se and compounds may suggest evaporation problems for this temperature, and the related data were removed from the final data set. Moreover, we did not include the data for  $T = 1118$  K because they show a less negative departure from ideality than the values measured at  $T = 1175$  K; this would introduce a

negative excess  $C_p$ . As shown in Fig. 2, the calculated enthalpies of mixing  $(921 \text{ and } 1175 \text{ K})$  are in good agreement with the experimental data. The convexity shown by the two parts of these curves have to be compared to the endothermic values of the enthalpic terms  $L_{\text{Sn,SnSe}}^{\text{0,liq}}$  and  $L_{\text{Sn,Se,Se}}^{\text{0,liq}}$  (Table 8), which indicate repulsion between the associate and the pure elements. The consequence in the phase diagram is the presence of a miscibility gap in the Sn-rich portion and of a flat liquidus in the Se-rich region. The degree of association, which is 99.7% at 921 K and 97.4% at 1175 K, decreases slowly with temperature. This confirms the presence of strong interactions in this system, as shown by [88Kot] and [89Cas],

In the optimization of the parameters of solid phases, we first assumed a constant  $C_p$  term and obtained values for coefficients  $A$ ,  $B$ , and  $C$  (Eq 7). Then, since heat capacity data were available, the temperature dependence of  $C_p$  could be described by optimizing coefficients  $D$ ,  $E$ , and  $F$  of Eq 7. During this step, all the other parameters were fixed, and only the  $C_p$ data of [81Wie] were used. As shown in Fig. 3 and 4, the calculated heat capacities of both phases are in excellent agreement with the experimental values.

The heat content calculations are in very good agreement with the available data (Fig. 5 and 6). With the assumption made **for**  the description of phases  $\alpha$  and  $\beta$ , the calculated heat of transformation is 640 J/mol in accordance with the estimated values of [81Bali (800 J/g-atom) and [94Yam] (600 J/g-atom). The good fit of the heat content in the liquid phase indicates a good consistency between the heat content data and the temperature dependence of the enthalpies of mixing.

The use of the temperature and enthalpy of fusion given by [82Ale] for  $SnSe<sub>2</sub>$  led to a worse description of the phase diagram and heat content of phases. Thus, these values were not used in the final optimization. The calculated enthalpy of fusion of the two compounds, given in Table 3, are in fair agreement with the used data.

Several runs of optimization were tried with the different enthalpy of formation data given in the literature. (See "Enthalpy Data.") The values of [67Rau], [71Mel], and [92Boo] showed the best consistency with all the other data and were used in the final optimization. The result of the calculation gave a reasonable agreement with the experimental enthalpy of formation as shown in Table 9.







In Fig. 7, the calculated Sn partial Gibbs energies of the twophase fields are plotted. The measured values of [71 Mel] are in



fair agreement with the calculation. Since there is no twophase field,  $\alpha$ SnSe +  $\beta$ SnSe, and because these two phases have the same stoichiometric composition, Fig. 7 does not show any three-phase equilibrium point for the transformation temperature. For this temperature ( $T = 796.5$  K), we have the following equality:  $\mu_{Sn}^{L1} = \mu_{Sn}^{asnSe} = \mu_{Sn}^{SnSe} = \mu_{Sn}^{SnSe}$ . We have an apparent four-phase equilibrium in contradiction with the phase rule. But, this is only a consequence of the coincidence of the phase transformation temperatures according to the above-mentioned approximations. In fact, at  $T = 796.5$  K, we have the two following equilibria:  $\alpha$ SnSe  $\leftrightarrow$  L1 +  $\beta$ SnSe and  $\beta$ SnSe  $\leftrightarrow \alpha$ SnSe + SnSe<sub>2</sub>.

The results of the phase diagram calculation together with the experimental data are given in Fig. 8. The liquidus data of [66Kar] on the Se-rich side, being very different from the values of the other authors, were not used in the final optimization run. Contamination of elements is not an explanation as the degenerate invariant temperature on the Se side is in fair agreement with the other reported values. An explanation may be proposed after looking at the DTA curve given in the publication: the liquidus temperature does not seem to correspond to the end of melting, and thus the liquidus temperature may be underestimated. As shown in Table 3, some divergence occurs in the temperatures of fusion of the two compounds. In order to avoid getting an average value, we chose to use our experimental measurements, which were included with a normal weight to be sure not to force the optimization to fit our data. The same treatment was applied to the monotectic invariant for which some differences can be found in the temperature. Figure 8 shows that the optimized phase diagram fits the data well. The two invariants related to the transformation  $\alpha$ SnSe  $\rightarrow \beta$ SnSe are given by dashed lines for the reasons explained above. We included in Table l0 experimental and calculated invariants obtained in phase diagram investigations. The reported transformation temperatures of  $\alpha$ SnSe $\rightarrow \beta$ SnSe are given in Table 5, The calculated critical temperature of the miscibility gap, 1141 K, is in reasonable agreement with the proposed value of [87Kot] (1114 K).

## **Summary'**

DSC analysis and XRD investigation provide a reaction between phases instead of a phase transition around the composition  $X_{\text{S}_e} = 0.500$ . Nevertheless, the lack of experimental data defining the narrow range of solubility of both phases made it necessary to assume the phases as stoichiometric with the same stoichiometry. The good agreement observed between the calculated results and the literature experimental values validates the assumption. Considering the difficulty in obtaining valuable experimental data on the concentration range of phases  $\alpha$ SnSe and  $\beta$ SnSe and in determining the kind of defects that enable the deviation from the ideal composition, the thermodynamic description presented here seems to represent a fair modeling of the experimentally determined properties of this system.

#### **Acknowledgments**

The authors thank H.L. Lukas for supplying the programs, additional information on how to use them, and helpful discussions.

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#### **Table 9 Calculated and Experimental Enthalpy of Formation**





Fig. 7 Temperature versus Sn chemical potential diagram of the two phases field.



## **Appendix**

#### **Table 10 Experimental and Calculated Invariants in the Sn-Se Phase Diagram**







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