

Short Note

## Enhancement of *T*-Noninvariant Effects in Neutron-Induced Nuclear Reactions

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A b s t r a c t: The previously developed approach to P-nonconservation in nuclear reactions is applied to the investigation of P- and T-noninvariant effects in neutron-induced reactions.Dynamical and resonance enhancement effects for T-noninvariant quantities are considered.The estimate is given of the expected effect in the compound-nucleus p-resonance.

The possible measurements of simultaneous P- and T-voilation in the transmission of polarized neutron beam through a target with polarized nuclei were considered recently/1,2/. These effects arize from the P- and T-noninvariant combination of the target spin  $\vec{I}$ , neutron spin  $\vec{c}$  and neutron momentum  $\vec{p}$  in the scattering amplitude: $\vec{c} [\vec{p} \times \vec{I}]$ . One of the effects leads to a precession of  $\vec{c}$  around the  $[\vec{p} \times \vec{I}]$  axis. When the target polarization is perpendicular to  $\vec{p}$  the component of  $\vec{c}$  in the  $(\vec{p}, \vec{I})$  plane will be rotating with "velocity"

$$\frac{d \chi}{d z} = \frac{4\pi N}{\kappa} Re(f_x - f_{-x})$$

where z is the sample length, k is the neutron wave vector, N is the density of target nuclei in a sample, while  $f_x$  and  $f_{-x}$  are the zero angle scattering amplitudes for neutrons polarized along and against the  $[\vec{p} \times \vec{I}]$  axis respectively.

The second effect consists in observing the difference between the total neutron cross-sections for  $\overrightarrow{\delta}$  parallel and antiparallel to  $\overrightarrow{p \times 1}$  which is

$$\Delta_{\tau} = \frac{4\pi}{\kappa} \operatorname{Im}(f_{x} - f_{-x})$$

We are going now to find the expressions for (1),(2) demonstrating the possible enhancement mechanisms arising in nucleon-nucleus reactions.Using the formalism of /3/ and the microscopic theory of nuclear reactions /4/ we can follow the procedure of /5/ to obtain the expression for f in the first Born approximation in weak interaction W.Then

$$\frac{d\chi}{dZ} = \frac{\mu_{JTNG}}{\kappa^2} \frac{w(\Gamma_s^n \Gamma_p^n)^{\gamma_2}}{[s] \Gamma_p]} [(E - E_s)(E - E_p) - \frac{\Gamma_s \Gamma_p}{4}]$$

$$\Delta_{\tau} = -\frac{2\pi G}{K^{2}} \frac{w(\Gamma_{s}^{n} \Gamma_{p}^{n})^{1/2}}{[(E - E_{p})\Gamma_{s} + (E - E_{s})\Gamma_{p})]}$$

Here  $G \approx 1$  is the spin factor(omitted further on),  $\mathcal{W}=\mathrm{Im}\int \Psi_{S} \mathcal{W}\Psi_{P} d\mathcal{T}$  is the Pand T-violating part of the weak interaction matrix element between the s- and p-compound resonance wave functions admixed to each other;  $\mathbf{E}_{s,p}$  and  $\Gamma_{s,p}$  are the positions and widths of these resonances;  $[s,p] = (\mathbf{E}-\mathbf{E}_{s,p})^{2} + \Gamma_{s,p}^{2}$ . It is often convenient to use the angle

X observed for a sample thickness equal to the mean free path  $l_{m.f.}$  and the relative value  $\chi = \Delta_T / 26_{tot}$ . Now we are able to relate these values to the corresponding values of P-violating effects (see /5/)  $\Phi$  and P obtained for neutrons with opposite helicities and unpolarized target. Since the letter effects are

essentially defined by the T-invariant part v of the weak interaction matrix element  $v = \operatorname{Re} \int \Psi_{s} \Psi \Psi_{p} d\tau$ we use the procedure of /5/ to get

 $\eta = \mathbf{P} \cdot \frac{w}{\mathbf{v}}$  $\chi = \Phi \cdot \frac{\omega}{\omega}$ Thus we get for the effects caused by the T- and P-violating interaction in the processes going through compound resonances the same two enhancement factors as for P-violating caused by v /5, 6/, namely the dynamical enhancement factor W/D arising from small spacing D between the admixed s- and p-resonances and the resonance enhancement factors D/f or  $(D/f)^2$  arising from the resonance propagators in (3). The quantities  $\chi$  and  $\eta$  (contrary to d)  $(dz \text{ and } \Delta_m)$  would feel resonance enhancement only in the vicinity of p-resonance(see the analogous property of  $\phi$  and P in (5/). Mind that the net enhancement for P was proved both theoretically and experimentally to reach 5-6 orders of magnitude(see e.g./5/ and references therein). In order to estimate the absolute values of  $\mathcal K$ and  $\eta$  one needs to estimate the ratio r= w/v.To do this one can use the resonable hypothesis that the ratio of nuclear matrix elements equals the ratio of the corresponding nucleon-nucleon interaction constants. According to the theories of spontaneous CP-violation in weak interaction /7/ which describe satisfactory experimental data on K-meson decays this ratio is

$$\mathbf{r} \sim \frac{\mathbf{m}_{q}^{2}}{\mathbf{m}_{H}^{2}} \cdot \frac{\mathbf{m}_{N}}{\mathbf{m}_{q}}$$

Where  $m_q, m_N$  and  $m_H$  are the masses of quark, nucleon and the light Higgs' boson responsible for the CP-noninvariant interaction with quarks. Reasonable estimates of  $m_{0}$  and  $m_{H}$  give us  $r \sim 10^{-2} - 10^{-3}$ .

With this estimate we obtain for the p-compound resonance, say in <sup>139</sup>La(where experiment showed /8/ a large value of P≈0,1):

$$\chi$$
 (La)~10<sup>-3</sup>-10<sup>-4</sup>  $\chi$  (La)~10<sup>-3</sup>-10<sup>-4</sup>  
rad

Finding the effects would unambiguously indicate the presence of "milliweak" interaction which violates T-invariance. One can point that resonance behaiour of  $\eta$  and  $\chi$  allows to separate the real effects from possible background caused, say, by internal fields in solid target sample.

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