

ON THE SINGULAR BOCHNER–MARTINELLI INTEGRAL

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The Bochner-Martinelli (B.-M.) kernel inherits, for $n \geq 2$, only some of properties of the Cauchy kernel in \mathbb{C} . For instance it is known that the singular B.-M. operator M_n is not an involution for $n \geq 2$. M. Shapiro and N. Vasilevski found a formula for M_2^2 using methods of quaternionic analysis which are essentially complex-twodimensional. The aim of this article is to present a formula for M_n^2 for any $n \geq 2$. We use now Clifford Analysis but for $n = 2$ our formula coincides, of course, with the above-mentioned one.

1 Theorem on the square of the singular Bochner-Martinelli

1.1 We consider the m -dimensional complex space \mathbb{C}^m of the variable $z = (z_1, \dots, z_m)$. If $\{z, w\} \subset \mathbb{C}^m$ then $\langle z, w \rangle := \sum_{k=1}^m z_k w_k$, $|z| := \sqrt{\langle z, \bar{z} \rangle}$ where $\bar{z} := (\bar{z}_1, \dots, \bar{z}_m)$. Topology in \mathbb{C}^m is determined by the metric $dist(z, w) := |z - w|$. If $z \in \mathbb{C}^m$ then $Re z := (Re z_1, \dots, Re z_m) \in \mathbb{R}^m$, $Re z_j = x_j$, and $Im z = (Im z_1, \dots, Im z_m)$, $Im z_j = y_j$. Hence $\mathbb{C}^m \cong \mathbb{R}^{2m}$. Orientation in \mathbb{C} is defined by the order of coordinates $(x_1, y_1, \dots, x_m, y_m)$ which means that the differential form of volume is of the form: $dv := (-1)^{\frac{m(m-1)}{2}} dx_1 \wedge \dots \wedge dx_m \wedge dy_1 \wedge \dots \wedge dy_m = (-1)^{\frac{m(m-1)}{2}} dx \wedge dy = (-1)^{\frac{m(m-1)}{2}} (\frac{i}{2})^m dz \wedge d\bar{z} = (-1)^{\frac{m(m-1)}{2}} (-\frac{i}{2})^m d\bar{z} \wedge dz$ with $dz := dz_1 \wedge \dots \wedge dz_m$. Standard notation: $\mathbb{B}(z; \varepsilon) := \{\zeta \in \mathbb{C}^m \mid |\zeta - z| < \varepsilon\}$.

1.2 Denote, as usually, by $u(\zeta, z)$ an exterior differential form (of type $(m, m-1)$)

$$u(\zeta, z) := \frac{(m-1)!}{(2\pi i)^m} \sum_{k=1}^m (-1)^{k-1} \frac{\bar{\zeta}_k - \bar{z}_k}{|\zeta - z|^{2m}} d\bar{\zeta}_{[k]} \wedge d\zeta,$$

where $d\bar{\zeta}_{[k]} := d\bar{\zeta}_1 \wedge \dots \wedge d\bar{\zeta}_{k-1} \wedge d\bar{\zeta}_{k+1} \wedge \dots \wedge d\bar{\zeta}_m$. For $m = 1$ the form $u(\zeta, z)$ is nothing more than the Cauchy kernel $\frac{d\zeta}{2\pi i(\zeta - z)}$.

Let $g(\zeta, z)$ be a fundamental solution of the Laplace operator, i.e.

$$g(\zeta, z) := \begin{cases} -\frac{(m-2)!}{(2\pi)^m} \frac{1}{|\zeta-z|^{2m-2}}, & m > 1, \\ \frac{1}{2\pi} \ln|\zeta-z|, & m = 1. \end{cases}$$

Then

$$\begin{aligned} u(\zeta, z) &= \sum_{k=1}^m (-1)^{k-1} \frac{\partial g}{\partial \zeta_k} d\bar{\zeta}_{[k]} \wedge d\zeta = \\ &= (-1)^{m-1} \partial_\zeta g \wedge \sum_{k=1}^m d\bar{\zeta}_{[k]} \wedge d\zeta_{[k]} \end{aligned}$$

where $\partial := \sum_{k=1}^m d\zeta_k \frac{\partial}{\partial \zeta_k}$. $\Delta_{\mathbb{C}^m}$ will stand for the complex Laplace operator

$$\Delta_{\mathbb{C}^m} := \sum_{k=1}^m \frac{\partial^2}{\partial \zeta_k \partial \bar{\zeta}_k} = \frac{1}{4} \sum_{k=1}^m \left(\frac{\partial^2}{\partial x_k^2} + \frac{\partial^2}{\partial y_k^2} \right) = \frac{1}{4} \Delta_{\mathbb{R}^{2m}}.$$

1.3 Let Ω be a bounded domain in \mathbb{C}^m with the piece-wise smooth boundary $\Gamma = \partial\Omega$. For any f holomorphic in Ω and continuous up to the boundary Γ the following Bochner–Martinelli integral representation holds:

$$f(z) = \int_{\partial\Omega} f(\zeta) u(\zeta, z) := F_f(z), \quad z \in \Omega. \tag{1.1}$$

For $m = 1$ one has just the Cauchy integral formula but for $m > 1$ the kernel in (1.1) is not holomorphic (in z and ζ), which causes many difficulties.

In particular, the Bochner–Martinelli–type integral for an integrable function f on Γ ($f \in \mathcal{L}^1(\Gamma)$):

$$F_f(z) := \int_{\Gamma} f(\zeta) u(\zeta, z), \quad z \notin \Gamma,$$

is not holomorphic in Ω unless f is a trace of a holomorphic function. If $z \in \Gamma$ then $F_f(z)$ does not, generally speaking, exist as an improper integral because the integrand has a singularity $\frac{1}{|\zeta-z|^{2m-1}}$. That is why for $z \in \Gamma$ there will be considered the Cauchy principal value of the Bochner–Martinelli integral:

$$F_f(z) := \lim_{\varepsilon \rightarrow 0} \int_{\Gamma \setminus \mathbb{B}(z; \varepsilon)} f(\zeta) u(\zeta, z), \quad z \in \Gamma.$$

We restrict ourselves to the case of a Liapunov surface Γ . It is well-known (see, for instance, [Ky]) for $f \in C^0, \mu$, the Hölder function space, $F_f(z)$ exists everywhere on Γ , and for $f \in L_p$, $p > 1$, $F_f(z)$ exists a.e. on Γ and $F_f \in L_p$.

Let

$$M_m[f](z) := 2 \int_{\Gamma} f(\zeta)u(\zeta, z), \quad z \in \Gamma.$$

For $m = 1$

$$M_1[f](z) = \frac{1}{\pi i} \int_{\Gamma} \frac{f(\zeta)}{\zeta - z} d\tau$$

and it is very well known that M_1 is an involution both on C^0, μ and on L_p :

$$M_1^2 = I, \tag{1.2}$$

I the identity operator. The property plays an extremely important role in many problems.

It is quite natural then to ask the question about M_m^2 for any m . In [Se] it was stated that $M_m^2 = I$ for any m , as well, but then an unremovable error was detected in the proof (see [AiYu], [Ky, §22]). The right correlation for $m = 2$ was given by M. Shapiro and N. Vasilevski, see [VaSh1], [VaSh2]:

$$M_2^2 = I + (N_2Z)^2 \tag{1.3}$$

where the operator N_2 acts by the rule

$$N_2[f](z) := -\frac{1}{2\pi^2} \int_{\Gamma} \frac{(\bar{\zeta}_1 - \bar{z}_1)d\zeta_{[2]} \wedge d\bar{\zeta} + (\bar{\zeta}_2 - \bar{z}_2)d\zeta_{[1]} \wedge d\bar{\zeta}}{|\zeta - z|^4} f(\zeta),$$

and Z is the operator of complex conjugation: $Z(f) := \bar{f}$. It was shown also that

$$M_2N_2 + N_2M_2Z = 0. \tag{1.4}$$

Denote by $w_2(\zeta, z)$ the kernel of the integral $N_2[f]$. It was proved later in [MiSh] that for any $f \in C^1(\Gamma; \mathbb{C})$ and $z \notin \Gamma$ there holds:

$$\int_{\Gamma} w_2(\zeta, z)f(\zeta) = \int_{\Gamma} g(\zeta, z) \cdot \frac{1}{|\text{grad}\rho(\zeta)|} \cdot \frac{\partial(f, \rho)}{\partial(\zeta_1, \zeta_2)} \cdot dS \tag{1.5}$$

where $\Gamma = \{z \in \mathbb{C}^2 | \rho(z) = 0, \rho \in C^1(\mathbb{C}^2; \mathbb{R}), \text{grad } \rho|_{\Gamma} \notin 0\}$. Formula (1.5) explains the commutation relation (1.4).

1.4 Formula (1.3) was obtained by the methods of quaternionic analysis which, by their nature, fits the case of \mathbb{C}^2 . In this paper we treat the general situation of \mathbb{C}^m using Clifford analysis.

1.5 Introduce the following notation for $1 \leq p < q \leq m$:

$$\begin{aligned} N_{p, q}[f](z) &:= 2 \frac{(m-1)!}{(2\pi i)^m} \int_{\Gamma} \frac{1}{|\zeta - z|^{2m}} ((-1)^q(\bar{\zeta}_p - \bar{z}_p)f(\zeta)d\bar{\zeta} \wedge d\zeta_{[q]} \\ &\quad - (-1)^p(\bar{\zeta}_q - \bar{z}_q)f(\zeta)d\bar{\zeta} \wedge d\zeta_{[p]}) \\ M_{p, q} &:= N_{p, q} \circ Z. \end{aligned}$$

Let $w_{p, q}$ be a kernel of this integral, then for $m = 2$ the only possible combination is $p = 1, q = 2$, and hence $w_{1, 2}$ coincides with w_2 and $N_2 = N_{1, 2}$.

1.6 Theorem (*Square of the singular Bochner–Martinelli operator*). *Under the above described conditions the following equality holds on $C^{0, \mu}(\Gamma; \mathbb{C})$, $0 < \mu < 1$; and $L_p(\Gamma; \mathbb{C})$, $p > 1$:*

$$M_m^2 = I + \sum_{1 \leq p < q \leq m} M_{p, q}^2.$$

2 Prof of the Theorem

2.1 Let $Cl_{0, 2m}$ be a complex Clifford algebra with generators e_1, e_2, \dots, e_{2m} . This means that

$$\begin{aligned} e_k^2 &= -1 =: -e_0, & k \in \{1, \dots, 2m\}, \\ e_k e_q + e_q e_k &= 0, & k \neq q, \end{aligned}$$

any element $a \in Cl_{0, 2m}$ is of the form

$$a = \sum_A a_A e_A$$

where $A = (\alpha_1, \dots, \alpha_p)$ with $1 \leq \alpha_1 < \dots < \alpha_p \leq 2m$, $\{a_A\} \subset \mathbb{C}$, $e_A := e_{\alpha_1} \dots e_{\alpha_p}$. All the necessary information the reader can find in [DeSoSo] and in many other sources.

Mention that for the Clifford conjugation of a we use the notation a^* :

$$a^* := \sum_A a_A e_A^*$$

with $e_A^* := e_{\alpha_p}^* \dots e_{\alpha_1}^* := (-e_{\alpha_p}) \dots (-e_{\alpha_1})$, and for the complex conjugation \bar{a} :

$$\bar{a} := \sum_A \bar{a}_A \cdot e_A$$

with $\bar{a}_A := Re a_A - iIm a_A$. Note that sometimes both conjugations are denoted by the same symbol but we prefer to use different ones.

2.2 Let Ω be a domain in \mathbb{R}^{2m} and let $f \in C^1(\Omega; Cl_{0, 2m})$. f is called hyperholomorphic, or monogenic, or regular, in Ω if

$$D[f](t) := \sum_{k=1}^{2m} e_k \cdot \frac{\partial f}{\partial t_k} = 0, \quad \forall t \in \Omega.$$

We shall denote $ker D =: \mathfrak{M}(\Omega)$. For $m = 1$ the definition reduces to the usual holomorphic functions of one complex variable. For an arbitrary m the theory of hyperholomorphic functions inherits many basic structural properties of one-dimensional complex analysis.

Let $t = (t_1, t_2, \dots, t_{2m-1}, t_{2m})$, $\tau = (\tau_1, \tau_2, \dots, \tau_{2m-1}, \tau_{2m})$ be two points in \mathbb{R}^{2m} . The Clifford–Cauchy kernel $\mathcal{K}(\tau - t)$ is defined as

$$\mathcal{K}(\tau - t) := \frac{1}{A_{2m}} \frac{(\tau - t)^*}{|\tau - t|^{2m}} \tag{2.6}$$

where $A_{2m} = |\mathbb{S}^{2m-1}| = \frac{2\pi^m}{\Gamma(m)} = \frac{2\pi^m}{(m-1)!}$,

$$\begin{aligned} \tau - t &= \sum_{k=1}^{2m} (\tau_k - t_k) e_k, \\ (\tau - t)^* &= \sum_{k=1}^{2m} (\tau_k - t_k) e_k^* = - \sum_{k=1}^{2m} (\tau_k - t_k) e_k. \end{aligned}$$

Let $\sigma := \sigma_\tau$ denote a differential form

$$\sigma_\tau := \sum_{j=1}^{2m} (-1)^{j-1} e_j d\tau_{[j]},$$

where $d\tau_{[j]} := d\tau_1 \wedge \dots \wedge d\tau_{j-1} \wedge d\tau_{j+1} \wedge \dots \wedge d\tau_{2m}$, then for any function $f \in \mathfrak{M}(\Omega; Cl_{0, 2m}) \cap C(\bar{\Omega}; Cl_{0, 2m})$ the Clifford–Cauchy integral formula holds:

$$\int_{\Gamma} \mathcal{K}(\tau - t) \sigma_\tau f(\tau) = \begin{cases} f(t), & t \in \Omega, \\ 0, & t \in \mathbb{R}^{2m} \setminus \bar{\Omega}. \end{cases} \tag{2.7}$$

Let $f \in C^0, \mu(\Gamma; Cl_{0, 2m})$ then on Γ

$$S[f](t) := 2 \int_{\Gamma} \mathcal{K}(\tau - t) \sigma_\tau f(\tau) \tag{2.8}$$

exists in the sense of Cauchy’s principal value; what is more on $C^0, \mu(\Gamma; Cl_{0, 2m})$

$$S^2 = I, \tag{2.9}$$

compare with (1.2) and (1.3).

If $f \in L_p(\Gamma; Cl_{0, 2m})$ then the Cauchy principal value of (2.8) exists almost everywhere on Γ and (2.9) holds on L_p .

2.3 For what follows we need some properties of idempotents in Clifford algebra $Cl_{0, 2m}$. Set

$$J_k := \frac{1}{2}(1 + ie_{2k-1}e_{2k})$$

for $k \in \{1, \dots, m\}$. J_k is an idempotent: $J_k^2 = J_k$. Besides $J_j J_k = J_k J_j$. The product $J := J_1 \cdots J_m$ is a primitive idempotent: $J^2 = J$, and $J_j \cdot J = J \cdot J_j$ for any j . Very important for us is: for any j

$$\begin{aligned} e_{2j} J &= ie_{2j-1} J, \\ e_{2j-1} e_{2j} J &= -iJ. \end{aligned}$$

Both J_k and \bar{J}_k are zero divisors:

$$J_k \cdot \bar{J}_k = \bar{J}_k \cdot J_k = 0.$$

2.4 Introduce notation:

$$w_j := e_{2j-1} + ie_{2j}, \quad j \in \{1, \dots, m\},$$

then

$$\bar{w}_j := e_{2j-1} - ie_{2j}.$$

Of course

$$J_k = -\frac{1}{2}e_{2k-1}\bar{w}_k$$

which is equivalent to

$$\begin{aligned} \bar{w}_k &= 2e_{2k-1} \cdot J_k; \\ w_k &= 2e_{2k-1} \cdot \bar{J}_k. \end{aligned}$$

This gives immediately: for any j

$$w_j \cdot J_j = 0, \quad \bar{w}_j \cdot \bar{J}_j = 0,$$

and

$$w_j^2 = 0, \quad (\bar{w}_j)^2 = 0.$$

One can check up directly the following equalities:

$$\begin{aligned} w_j \cdot \bar{w}_k &= -\bar{w}_k \cdot w_j \quad \text{for } j \neq k; \\ \bar{w}_j \cdot \bar{w}_k &= -\bar{w}_k \cdot \bar{w}_j \quad \text{for } j \neq k; \\ \bar{w}_j e_{2p-1} e_{2q-1} &= e_{2p-1} e_{2q-1} \bar{w}_j \quad \text{for } j \neq p, j \neq q, p \neq q; \\ w_j \cdot \bar{w}_j &= -4 \cdot J_j, \\ \bar{w}_j \cdot w_j &= -4 \cdot \bar{J}_j, \\ \bar{w}_j \cdot \bar{w}_k \cdot J &= \begin{cases} 0, & j = k, \\ 4e_{2j-1}e_{2k-1}J, & j \neq k; \end{cases} \\ w_j \cdot \bar{w}_k \cdot J &= \begin{cases} -4J, & j = k, \\ 0, & j \neq k; \end{cases} \\ \bar{w}_j \bar{w}_k \bar{w}_p \bar{w}_q J &= 16e_{2j-1}e_{2k-1}e_{2p-1}e_{2q-1}J \quad \text{for } j, k, p, q \text{ different in paires.} \end{aligned}$$

2.5 On the underlying space, \mathbb{R}^{2m} , of the Clifford algebra $Cl_{0, 2m}$, we introduce now the following complex structure. If $t \in \mathbb{R}^{2m}$ then we write x_j for t_{2j-1} and y_j for t_{2j} ,

setting after this $z_j := x_j + iy_j$, $z := (z_1, \dots, z_m)$. In the same fashion $\tau \in \mathbb{R}^{2m}$ gives us $\zeta_j = \xi_j + i\eta_j$, $\zeta = (\zeta_1, \dots, \zeta_m)$. In these notations, we have:

$$\begin{aligned} (\tau - t) &= \sum_{k=1}^{2m} (\tau_k - t_k) e_k = \\ &= \sum_{k=1}^m ((\xi_k - x_k) e_{2k-1} + (\eta_k - y_k) e_{2k}) = \\ &= \frac{1}{2} \sum_{k=1}^m ((\zeta_k - z_k) e_{2k-1} + (\bar{\zeta}_k - \bar{z}_k) e_{2k-1} + \\ &\quad + i(\bar{\zeta}_k - \bar{z}_k) e_{2k} - i(\zeta_k - z_k) e_{2k}) = \\ &= \frac{1}{2} \sum_{k=1}^m ((\zeta_k - z_k) \cdot \bar{w}_k + (\bar{\zeta}_k - \bar{z}_k) \cdot w_k). \end{aligned}$$

Let

$$\mathcal{X}_j(\zeta - z) := \frac{-1}{2A_{2m}} \frac{\zeta_j - z_j}{|\zeta - z|^{2m}} = -\frac{(m-1)!}{4\pi^m} \frac{\zeta_j - z_j}{|\zeta - z|^{2m}},$$

hence

$$\mathcal{X}(\zeta - z) = \sum_{j=1}^m (\mathcal{X}_j(\zeta - z) \cdot \bar{w}_j + \bar{\mathcal{X}}_j(\zeta - z) \cdot w_j).$$

Introduce the following notations:

$$\begin{aligned} \alpha_j &:= \frac{(-1)^{\lfloor \frac{m}{2} - \frac{1}{4} \rfloor} (-1)^j}{2^m \cdot i^m}, \\ \beta_j &:= \frac{(-1)^{\lfloor \frac{m}{2} - \frac{1}{4} \rfloor} (-1)^{m+j}}{2^m \cdot i^m}, \quad j \in \{1, \dots, m\}. \end{aligned}$$

We have:

$$\begin{aligned} d\tau_{2j-1} &= \frac{1}{2} (d\zeta_j + d\bar{\zeta}_j), \\ d\tau_{2j} &= \frac{i}{2} (d\bar{\zeta}_j - d\zeta_j). \end{aligned}$$

Thus

$$\begin{aligned} d\tau_{[2j-1]} &= \frac{1}{2} (d\zeta_1 + d\bar{\zeta}_1) \wedge \frac{i}{2} (d\bar{\zeta}_1 - d\zeta_1) \wedge \dots \wedge \frac{i}{2} (d\bar{\zeta}_{j-1} - d\zeta_{j-1}) \wedge \\ &\quad \wedge \frac{i}{2} (d\bar{\zeta}_j - d\zeta_j) \wedge \frac{1}{2} (d\bar{\zeta}_{j+1} + d\zeta_{j+1}) \wedge \dots \wedge \frac{i}{2} (d\bar{\zeta}_m - d\zeta_m) \\ &= i^m \cdot \frac{1}{2^{2m-1}} \cdot (d\bar{\zeta}_j - d\zeta_j) \bigwedge_{\substack{k=1 \\ k \neq j}}^m (d\zeta_k + d\bar{\zeta}_k) \wedge (d\bar{\zeta}_k - d\zeta_k) = \\ &= \frac{i^m}{2^m} (d\bar{\zeta}_j - d\zeta_j) \bigwedge_{\substack{k=1 \\ k \neq j}}^m d\zeta_k \wedge d\bar{\zeta}_k = \\ &= \frac{i^m (-1)^{\lfloor \frac{m}{2} - \frac{1}{4} \rfloor + (m-1)}}{2^m} (d\bar{\zeta}_j - d\zeta_j) \wedge d\zeta_{[j]} \wedge d\zeta_{[j]} \end{aligned}$$

and

$$\begin{aligned}
 d\tau_{[2j]} &= i^{m-1} \cdot \frac{1}{2^{2m-1}} \cdot (d\bar{\zeta}_j + d\zeta_j) \bigwedge_{\substack{k=1 \\ k \neq j}}^m (d\zeta_k + d\bar{\zeta}_k) \wedge (d\bar{\zeta}_k - d\zeta_k) = \\
 &= \frac{i^{m-1}}{2^m} (d\bar{\zeta}_j + d\zeta_j) \bigwedge_{\substack{k=1 \\ k \neq j}}^m d\zeta_k \wedge d\bar{\zeta}_k = \\
 &= \frac{(-1)^{\lfloor \frac{m}{2} - \frac{1}{4} \rfloor + (m-1)} \cdot i^{m-1}}{2^m} (d\bar{\zeta}_j + d\zeta_j) \wedge d\bar{\zeta}_{[j]} \wedge d\zeta_{[j]}.
 \end{aligned}$$

All this gives:

$$\begin{aligned}
 \sigma_\zeta &= \frac{(-1)^{\lfloor \frac{m}{2} - \frac{1}{4} \rfloor - 1}}{2^m \cdot i^m} \sum_{j=1}^m (ie_{2j}(d\bar{\zeta}_j + d\zeta_j) + \\
 &\quad + e_{2j-1}(d\bar{\zeta}_j - d\zeta_j)) \wedge d\bar{\zeta}_{[j]} \wedge d\zeta_{[j]} = \\
 &= \frac{(-1)^{\lfloor \frac{m}{2} - \frac{1}{4} \rfloor - 1}}{2^m \cdot i^m} \sum_{j=1}^m (d\bar{\zeta}_j(ie_{2j} + e_{2j-1}) + \\
 &\quad + d\zeta_j(ie_{2j} - e_{2j-1})) \wedge d\bar{\zeta}_{[j]} \wedge d\zeta_{[j]} = \\
 &= \frac{(-1)^{\lfloor \frac{m}{2} - \frac{1}{4} \rfloor - 1}}{2^m \cdot i^m} \sum_{j=1}^m ((-1)^{j-1}(ie_{2j} + e_{2j-1})d\bar{\zeta} \wedge d\zeta_{[j]} + \\
 &\quad + (-1)^{m-j}(ie_{2j} - e_{2j-1})d\bar{\zeta}_{[j]} \wedge d\zeta).
 \end{aligned}$$

Thus finally

$$\sigma_\zeta = \sum_{j=1}^m (\alpha_j w_j d\bar{\zeta} \wedge d\zeta_{[j]} + \beta_j \bar{w}_j d\bar{\zeta}_{[j]} \wedge d\zeta).$$

Consider now

$$\begin{aligned}
 \mathfrak{X}(\zeta - z) \cdot \sigma_\zeta &= \sum_{j, k=1}^m (\mathfrak{X}_j(\zeta - z) \cdot \alpha_k \cdot d\bar{\zeta} \wedge d\zeta_{[k]} \cdot \bar{w}_j \cdot w_k + \\
 &\quad + \bar{\mathfrak{X}}_j(\zeta - z) \cdot \alpha_k \cdot d\bar{\zeta} \wedge d\zeta_{[k]} \cdot w_j \cdot w_k + \\
 &\quad + \mathfrak{X}_j(\zeta - z) \cdot \beta_k \cdot d\bar{\zeta}_{[k]} \wedge d\zeta \cdot \bar{w}_j \cdot \bar{w}_k + \\
 &\quad + \bar{\mathfrak{X}}_j(\zeta - z) \cdot \beta_k \cdot d\bar{\zeta}_{[k]} \wedge d\zeta \cdot w_j \cdot \bar{w}_k).
 \end{aligned}$$

Hence:

$$\begin{aligned}
 S[f](\zeta) &:= 2 \int_{\Gamma} \mathfrak{X}(\zeta - z) \cdot \sigma_\zeta \cdot f(\zeta) \\
 &= 2 \sum_{j, k=1}^m \left(\int_{\Gamma} \mathfrak{X}_j(\zeta - z) \cdot \alpha_k \cdot d\bar{\zeta} \wedge d\zeta_{[k]} \cdot \bar{w}_j w_k \cdot f(\zeta) + \right.
 \end{aligned}$$

$$\begin{aligned}
 & + \int_{\Gamma} \overline{\mathcal{X}}_j(\zeta - z) \cdot \alpha_k \cdot d\overline{\zeta} \wedge d\zeta_{[k]} \cdot w_j w_k \cdot f(\zeta) + \\
 & + \int_{\Gamma} \mathcal{X}_j(\zeta - z) \cdot \beta_k \cdot d\overline{\zeta}_{[k]} \wedge d\zeta \cdot \overline{w}_j \overline{w}_k \cdot f(\zeta) + \\
 & + \int_{\Gamma} \overline{\mathcal{X}}_j(\zeta - z) \cdot \beta_k \cdot d\overline{\zeta}_{[k]} \wedge d\zeta \cdot w_j \overline{w}_k \cdot f(\zeta) \Big).
 \end{aligned}$$

Denote:

$$\begin{aligned}
 S_{jk}[f](\zeta) & := 2 \int_{\Gamma} \beta_k \overline{\mathcal{X}}_j(\zeta - z) f(\zeta) d\overline{\zeta}_{[k]} \wedge d\zeta, \\
 P_{jk}[f](\zeta) & := 2 \int_{\Gamma} \beta_k \mathcal{X}_j(\zeta - z) f(\zeta) d\overline{\zeta}_{[k]} \wedge d\zeta, \\
 Q_{jk}[f](\zeta) & := 2 \int_{\Gamma} \alpha_k \mathcal{X}_j(\tau - z) f(\zeta) d\overline{\zeta} \wedge d\zeta_{[k]}, \\
 R_{jk}[f](\zeta) & := 2 \int_{\Gamma} \alpha_k \overline{\mathcal{X}}_j(\zeta - z) f(\zeta) d\overline{\zeta} \wedge d\zeta_{[k]}.
 \end{aligned}$$

Hence

$$S = \sum_{j,k=1}^m (Q_{jk} \overline{w}_j w_k + R_{jk} w_j w_k + S_{jk} w_j \overline{w}_k + P_{jk} \overline{w}_j \overline{w}_k).$$

2.6 Consider now $S \cdot J$. Taking into account Section 2.4 we have:

$$\begin{aligned}
 S \cdot J & = \sum_{j,k=1}^m (S_{jk} w_j \overline{w}_k \cdot J + P_{jk} \overline{w}_j \overline{w}_k \cdot J) = \\
 & = -4 \sum_{k=1}^m S_{kk} J + 4 \sum_{j \neq k} P_{jk} e_{2j-1} e_{2k-1} J.
 \end{aligned}$$

From here, again applying Section 2.4,

$$\begin{aligned}
 S^2 \cdot J & = 16 \sum_{k=1}^m \sum_{q=1}^m S_{kk} S_{qq} J - \\
 & - 16 \sum_{j \neq k} \sum_{q=1}^m P_{jk} S_{qq} e_{2j-1} e_{2k-1} J + \\
 & + \sum_{k=1}^m \sum_{p \neq q} Q_{kk} P_{pq} \overline{w}_k w_k \overline{w}_p \overline{w}_q J +
 \end{aligned}$$

$$\begin{aligned}
 & + \sum_{k=1}^m \sum_{p \neq q} S_{kk} P_{pq} w_k \bar{w}_k \bar{w}_p \bar{w}_q J + \\
 & + \sum_{j \neq k} \sum_{p \neq q} Q_{jk} P_{pq} \bar{w}_j w_k \bar{w}_p \bar{w}_q J + \\
 & + \sum_{j \neq k} \sum_{p \neq q} R_{jk} P_{pq} w_j w_k \bar{w}_p \bar{w}_q J + \\
 & + \sum_{j \neq k} \sum_{p \neq q} S_{jk} P_{pq} w_j \bar{w}_k \bar{w}_p \bar{w}_q J + \\
 & + \sum_{j \neq k} \sum_{p \neq q} P_{jk} P_{pq} \bar{w}_j \bar{w}_k \bar{w}_p \bar{w}_q J.
 \end{aligned}$$

Compute now all the terms:

- 1)
$$\begin{aligned}
 & \sum_{k=1}^m \sum_{p \neq q} Q_{kk} P_{pq} \bar{w}_k w_k \bar{w}_p \bar{w}_q J = \\
 & = -16 \sum_{p \neq q} (Q_{pp} + Q_{qq}) P_{pq} e_{2p-1} e_{2q-1} J;
 \end{aligned}$$
- 2)
$$\begin{aligned}
 & \sum_{k=1}^m \sum_{p \neq q} S_{kk} P_{pq} w_k \bar{w}_k \bar{w}_p \bar{w}_q J = \\
 & = -16 \sum_{\substack{p \neq q \\ k \neq p \\ k \neq q}} S_{kk} P_{pq} e_{2p-1} e_{2q-1} J;
 \end{aligned}$$
- 3)
$$\begin{aligned}
 & \sum_{j \neq k} \sum_{p \neq q} Q_{jk} P_{pq} \bar{w}_j w_k \bar{w}_p \bar{w}_q J = \\
 & = 16 \sum_{\substack{p \neq q \\ k \neq p \\ k \neq q}} Q_{pk} (P_{qk} - P_{kq}) e_{2p-1} e_{2q-1} J;
 \end{aligned}$$
- 4)
$$\begin{aligned}
 & \sum_{\substack{j \neq k \\ p \neq q}} R_{jk} P_{pq} w_j w_k \bar{w}_p \bar{w}_q J = \\
 & = 16 \sum_{p, q} (R_{qp} - R_{pq}) P_{pq} J;
 \end{aligned}$$
- 5)
$$\begin{aligned}
 & \sum_{j \neq k} \sum_{p \neq q} S_{jk} P_{pq} w_j \bar{w}_k \bar{w}_p \bar{w}_q J = \\
 & = 16 \sum_{\substack{p \neq q \\ k \neq p \\ k \neq q}} S_{pk} (P_{kq} - P_{qk}) e_{2p-1} e_{2q-1} J;
 \end{aligned}$$
- 6)
$$\sum_{j \neq k} \sum_{p \neq q} P_{jk} P_{pq} \bar{w}_j \bar{w}_k \bar{w}_p \bar{w}_q J =$$

$$= 16 \sum_{\substack{j, k, p, q \\ \text{different} \\ \text{in pairs}}} P_{jk} P_{pq} e_{2j-1} e_{k-1} e_{2p-1} e_{2q-1} J.$$

2.7 This means that the equality

$$S^2 \cdot J = J$$

takes the form:

$$\begin{aligned} &M_m^2 J - \sum_{1 \leq p < q \leq m} M_{pq}^2 J - \\ &- 16 \sum_{p \neq q} \sum_{k=1}^m P_{pq} S_{kk} e_{2p-1} e_{2q-1} J - \\ &- 16 \sum_{p \neq q} (Q_{pp} + Q_{qq}) P_{pq} e_{2p-1} e_{2q-1} J - \\ &- 16 \sum_{\substack{p \neq q \\ k \neq p \\ k \neq q}} S_{kk} P_{pq} e_{2p-1} e_{2q-1} J + \\ &+ 16 \sum_{\substack{p \neq q \\ k \neq p \\ k \neq q}} S_{kp} (P_{kq} - P_{qk}) e_{2p-1} e_{2q-1} J + \\ &+ 16 \sum_{\substack{p \neq q \\ k \neq p \\ k \neq q}} Q_{pk} (P_{qk} - P_{kq}) e_{2p-1} e_{2q-1} J + \\ &+ 16 \sum_{\substack{j, k, p, q \\ \text{different} \\ \text{in pairs}}} P_{jk} P_{pq} e_{2j-1} e_{2k-1} e_{2p-1} e_{2q-1} J = J. \end{aligned}$$

2.8 This operator equality can be applied to any function f from $C^0, \mu(\Gamma; Cl_{0, 2m})$ or $L_p(\Gamma; Cl_{0, 2m})$. In particular on the corresponding subsets of \mathbb{C} -valued functions we get:

$$\begin{aligned} M_m^2 &= \sum_{1 \leq p < q \leq m} M_{pq}^2 + I; \\ &((P_{pq} - P_{qp}) M_m - 2(Q_{pp} + Q_{qq})(P_{pq} - P_{qp}) - \\ &- M_m(P_{pq} - P_{qp}) + 2 \sum_{\substack{p \neq q \\ k \neq p \\ k \neq q}} ((S_{kp} - Q_{pk})(P_{kq} - P_{qk}) + \\ &+ (S_{kq} - Q_{qk})(P_{pk} - P_{kq})) = 0 \end{aligned}$$

for any $1 \leq p < q \leq m$. The first equality gives what is written in Theorem in 1.6. Compare also the second equality with (1.4).

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