

## Amplified Phase-Conjugate Reflection of $\lambda = 10.51 \mu\text{m}$ Radiation in Gaseous $\text{SF}_6$

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**Abstract.** Phase-conjugate reflection at the  $P(12)$   $\text{CO}_2$  line ( $\lambda = 10.51 \mu\text{m}$ ) has been obtained in gaseous  $\text{SF}_6$  with a power reflection coefficient of 220% and an energy reflection coefficient of 115%.

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Optical phase conjugation is interesting for its ability to correct optical aberration of radiation passing through inhomogeneous media, to provide optimum focusing of light pulses in plasma studies, etc. [1, 2]. The possibility of optical phase conjugation of  $\text{CO}_2$ -laser radiation in resonant media [3–6] is of a particular interest.

This paper reports on the study of phase-conjugate reflection of  $\text{CO}_2$ -laser radiation at the vibrational molecular transitions in  $\text{SF}_6$  by degenerate four-wave mixing. The reflection has been obtained with the power and the energy of the reflected radiation exceeding those of the probe wave.

### Experimental

The experimental apparatus used is shown in Fig. 1. As a pump source, single-mode radiation ( $\text{TEM}_{00q}$ ) of a pulsed TEA  $\text{CO}_2$ -laser "1" tuned with a grating (100 lines/mm, blaze angle:  $30^\circ$ ) was used. The temporal behavior of output pulse has the shape of a narrow peak, on a wide background, with a duration of  $\tau_{0.5} = 100 \text{ ns}$  and  $\tau_{0.1} = 1.5 \mu\text{s}$ . A germanium 50% beam splitter "2" divided the  $\text{CO}_2$ -laser radiation into 3 beams. Two of them, of approximately equal intensity, were reflected from the mirrors "3" and formed counterpropagating pump waves in the gaseous medium  $\text{SF}_6$ . The third one, of lower intensity, was reflected from another face of the beam splitter "2" to form a probe wave. It was directed into the cell "5" by a ZnSe beam splitter "4". To provide a better overlap with the pump beam, the

lens "6" was used to focus the probe wave. Four-wave mixing of the pump and probe waves in  $\text{SF}_6$  results in a third-order polarization causing a phase-conjugate wave counterpropagating to the probe one. This radiation may be treated as a result of phase-conjugate four-wave-mirror reflection of the signal wave. The reflected wave was selected and recorded by the detector "12" consisting of a photon-drag detector ( $\tau_{\text{const}} = 1 \text{ ns}$ ) and a Ge–Au photoresistance ( $\tau_{\text{const}} \sim 2 \mu\text{s}$ ). To ensure operation in the linear of the photoresistance, a preliminary calibrated attenuators "11" were used. To find the reflection coefficient, the recording system was calibrated with the 100% mirror "7". A precisely backward-travelling probe wave reflect-

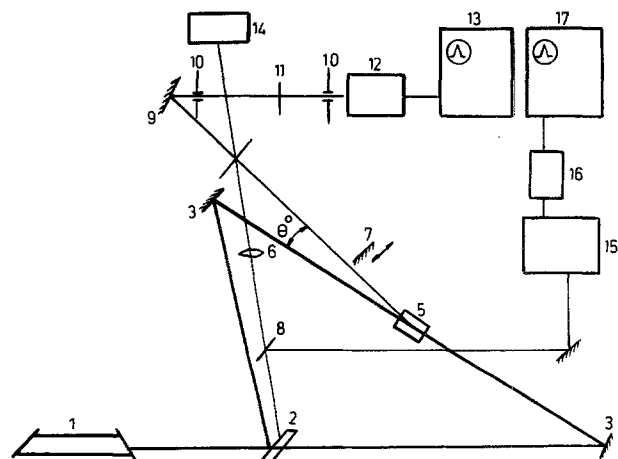


Fig. 1. Experimental arrangement

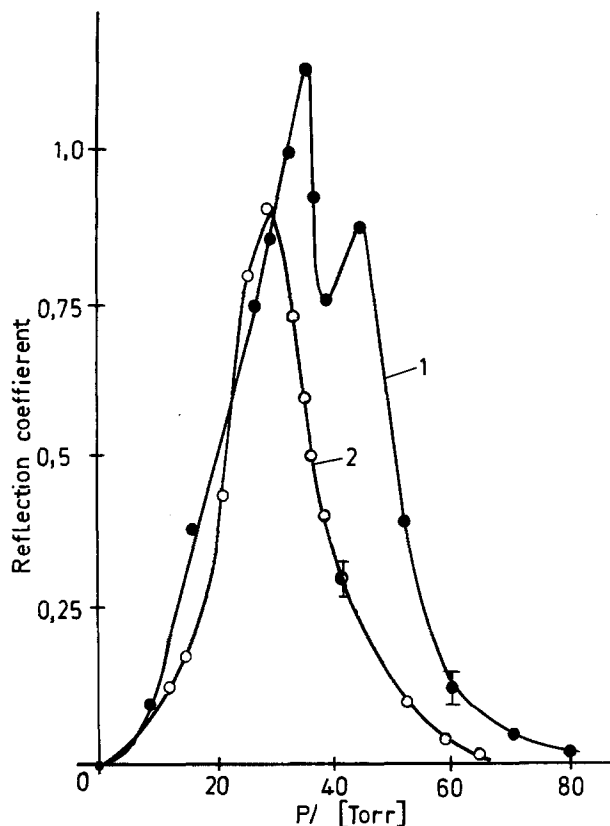


Fig. 2a and b. Reflection-coefficient dependence on  $\text{SF}_6$  pressure [Torr] at  $\lambda=10.51\mu\text{m}$  for the  $P(12)$   $\text{CO}_2$ -laser line: (a)  $I_p=6.25\text{ MW/cm}^2$ ; (b)  $I_p=3\text{ MW/cm}^2$

ed from mirror "7" passed through the ZnSe beam splitter "4" and entered the photodetector "12". The signal obtained corresponded to a 100% reflection. The reflection coefficient for power was measured by a photon-drag detector and that for energy by a photoresistor.

## Results

Experimental studies of the dependence of the reflection coefficient on the pump wavelength for the discussed scheme showed, unlike [4, 5], that the maximum reflection coefficient corresponds to the  $P(12)$

output line ( $\lambda=10.51\mu\text{m}$ ). It should be noted that for each generating line there is own optimum  $\text{SF}_6$  pressure.

Figure 2 illustrates the dependence of energy reflection coefficient  $R_e$  for the  $P(12)$   $\text{CO}_2$ -line on the  $\text{SF}_6$  pressure at two different pump intensities  $I_p$ . For  $I_p=6.25\text{ MW/cm}^2$  (Curve a) the maximum energy reflection coefficient is  $R_e\approx 115\%$  and the power one  $R_p\approx 220\%$ ; they were obtained at a  $\text{SF}_6$  pressure of  $P=34$  Torr. At  $P=38$  Torr a dip is observed, while at lower pump intensities  $I_p=3\text{ MW/cm}^2$  (Curve b) the dip is absent.

The dip appearance is obviously, caused by the competing coherent parametric and noncoherent linear and nonlinear processes in a strong pump field. It is interesting that just in this pressure region ( $P\approx 38$  Torr) at  $I_p=6.25\text{ MW/cm}^2$  the pressure dependence of the transmission coefficient for one of the pump waves has a plateau. This plateau was not observed at  $I_p=3\text{ MW/cm}^2$ . The study of the reflection-coefficient dependence on the pump intensity at the fixed pressure of  $P=34$  Torr showed that the growth is faster than quadratic. A similar peculiarity has been observed in [4] when studying nonlinear optical phase conjugation for the  $P(20)$   $\text{CO}_2$ -laser line. The study of the dependence of the reflected light power on the probe-wave intensity at the fixed pump intensity  $I_p=6.25\text{ MW/cm}^2$  showed a deviation from linearity at  $I_{\text{probe}}/I_p > 5 \times 10^{-2}$ .

In conclusion, using degenerate resonant four-wave mixing in gaseous molecular media provides, at optimum conditions, an effective amplification of the reflected conjugate radiation.

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