p, V, T and Derived Thermodynamic Data for Toluene, Trichloromethane, Dichloromethane, Acetonitrile, Aniline, and n-Dodecane

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Experimentally determined p, V, T data are reported for toluene, trichloromethane, dichloromethane, acetonitrile, aniline, and n-dodecane at 278, 288, 298, 313, and 323 K, except for dichloromethane, for which the highest temperature was 298 K. At each temperature, measurements were done at pressures up to about 280MPa or (for aniline and n-dodecane) at a lower pressure slightly below the freezing pressure at the temperature of measurement. Values of the isobaric expansivity isothermal compressibility and (for toluene, trichloromethane, dichloromethane, and acetonitrile) internal pressure, derived from the *p,V,T* data, are presented.

KEY WORDS: acetonitrile; aniline; dichloromethane; high pressure; ndodecane; p, V, T data; toluene; trichloromethane.

1. INTRODUCTION

During the past several years a bellows volumometer technique for determining volume ratios (i.e., volumes at elevated pressure relative to the volume at normal atmospheric pressure) has been developed in this laboratory [1]. The purpose of that development was to determine volume ratios with sufficient accuracy (of the order of $\pm 0.1\%$) to provide ancillary data for the measurement and interpretation of tracer and self-diffusion coefficients of liquids under pressure. However, the accuracy of the volumometer exceeds the original expectation and the instrument yields *p,V,T* data which can be used to derive reliable values of isothermal com-

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pressibility, isobaric expansivity, and internal pressure. The accuracy of p, V, T data obtained with the volumometer for water + acetonitrile mixtures at 298 K was demonstrated in an earlier paper [2] by a stringent test, viz., comparison of indirectly determined excess volumes with directly measured values [3]. The agreeement was within the probable uncertainty of the directly measured data. In subsequent work the temperature and pressure ranges of the volumometer were increased, and the water $+$ methanol system has been studied at temperatures from 278 to 323 K and pressures up to about 280 MPa [4, 5].

The objectives of the present investigation were to extend the temperature and/or pressure ranges at which p, V, T data are currently available for a number of pure liquids, all of which are relevant to our experimental program on diffusion in liquids under pressure. In this report we present volume-ratio data for six liquids, viz., toluene, trichloromethane, dichloromethane, acetonitrile, aniline, and n -dodecane, and derived values of compressibility, expansivity, and (except for aniline and n -dodecane) internal pressure for these liquids.

There are features of these derived quantities which are particularly interesting and justify the present grouping of these otherwise diverse substances.

2. EXPERIMENTAL

2.1. Measurements

Materials used were AnalaR or equivalent grade, used without further purification except for storage over molecular sieve, to remove water and other low molecular weight impurities, for several weeks before use. The trichloromethane, as received, contained 1-2 % by volume ethanol (added as stabilizer), which was not removed.

Volume ratios, k , defined by

$$
k = V_p / V_0 \tag{1}
$$

where V_0 and V_n are volumes at 0.1 MPa and pressure p, respectively, were determined using the bellows volumometer, whose construction, calibration, and use have been described previously [1,4]. The measurements were done at 278.15, 288.15, 298.15, 313.15, and 323.15 K (the temperatures at which the volumometer had been calibrated), except for dichloromethane, for which, owing to the low boiling point (approximately 313 K) at 0.1 MPa, the highest temperature was 298.15 K. The pressure range for the measurements was from about 2.5 MPa to

either about 280 MPa or (for aniline and *n*-dodecane) a lower pressure slightly below the freezing pressure at the temperature of measurement.

Temperatures were controlled to within $+0.005$ K and the uncertainty of temperature measurement (with a calibrated platinum resistance thermometer) was $+0.01$ K. Pressures were measured with two Heise bourdontube gauges. Above 25 MPa pressures were measured with a maximum uncertainty of $+0.4$ MPa with a 400-MPa gauge: in the region of 0-25 MPa a 25-MPa gauge was used, with an accuracy of $+0.025$ MPa. The overall accuracy of the volume ratio determinations corresponds to about $+0.05\%$ in density except (as for aniline and *n*-dodecane at 278K) where the experimental pressure range at a given temperature does not exceed about 50 MPa. In such cases the uncertainty is likely to be greater than 0.05 % in density.

2.2. Treatment of Experimental Data

Experimentally determined volume ratios were expressed as the secant bulk modulus, K , defined by

$$
K = p/(1 - k) \tag{2}
$$

and fitted to cubic equations of the form

$$
K = \sum_{i=0}^{3} A_i p^i \tag{3}
$$

We have found in the present and in previous work that Eq. (3) is a very good representation of the pressure dependence of experimental volume ratios at constant temperature. It should not, however, be extrapolated to pressures which substantially exceed the highest experimental pressure. The form of the Tait equation which has been extensively utilized by other workers, viz.,

$$
K = p/C \ln[(B+p)/(B+p_0)] \tag{4}
$$

where p_0 is normal atmospheric pressure, gives an inferior fit of volumeratio data but, unlike Eq. (3), can be reliably extrapolated to pressures beyond the experimental pressure range.

Isothermal compressibilities, κ , were calculated from the coefficients of Eq. (3) at each temperature by using the relationships

$$
\kappa = -(1/V)(\partial V/\partial p)_{T} = -(1/k)(\partial k/\partial p)_{T}
$$

\n
$$
\kappa = -[1/(p - K)][1 - (p/K)(\partial K/\partial p)_{T}]
$$
\n(5)

The probable uncertainty in κ is believed to be $+1-2\%$ for pressures higher than 50 MPa, provided that the experimental pressure range for volume-ratio measurement is at least 150 MPa.

The densities of the liquids at 0.1 MPa were taken from Timmermans $\lceil 6 \rceil$ and from Gibson and Loeffler $\lceil 7 \rceil$ for aniline at temperatures above 298 K. Isobaric expansivities, α , were calculated from densities at 0.1 MPa and volume ratios from Eq. (3) by evaluating the molar volume (V_m) as a function of temperature at 25, 50, 100, 150, 200, 250, and 275 MPa, fitting V_m to a quadratic in T (or to a linear equation in T where volume ratios at a given pressure were measured at fewer than three temperatures), then utilizing the relationship

$$
\alpha = (1/V_{\rm m})(\partial V_{\rm m}/\partial T)_{p} \tag{6}
$$

The probable uncertainty in α is estimated to be $+2-3\%$ for pressures higher than 50 MPa and where the experimental pressure range at a given temperature was at least 150 MPa.

Internal pressures, π , were calculated for toluene, trichloromethane, dichloromethane, and acetonitrile from the equation

$$
\pi = T(\partial p/\partial T)_v - p \tag{7}
$$

For each liquid, at each temperature the values of p corresponding to selected values of V_m were determined by interpolation of cubic spline fits of 10 values of p against V_m . The thermal pressure coefficients $(\partial p/\partial T)_v$ and hence π were evaluated from quadratic fits of these values of p against T. The probable uncertainty in π is of the order of $+3\%$. For aniline and ndodecane, since the experimental pressure range was considerably smaller at low temperatures than at high temperatures, π was not calculated.

3. RESULTS

The coefficients of Eq. (3) are listed in Table I, expansivities and compressibilities in Table II, and internal pressures in Table III. It should be noted that the entries in Table II for aniline and n -dodecane at temperatures of less than 298 K are approximate values, owing to the restricted accessible pressure range for these liquids at low temperatures. Internal pressures for aniline and n-dodecane are not tabulated. Values can be obtained via the thermal pressure coefficient by using Eq. (7) and the relationship

$$
(\partial p/\partial T)_v = \alpha/\kappa \tag{8}
$$

^{*a*} Maximum pressure approximately 280 MPa except where specified.
 ${}^b A_3 = 5164.9 \text{ GPa}^{-2}$.
 ${}^c A_3 = 1329.3 \text{ GPa}^{-2}$.

Thermodynamic Properties of Liquids

			p(MPa)						
T(K)	Property	0.1	25	50	100	150	200	250	275
				(a) Toluene					
278.15	$10^4\alpha$	10.5	9.47	8.54	7.03	5.89	5.08	4.55	4.38
	$10^4\kappa$	8.05	6.68	5.74	4.63	4.09	3.90	3.93	4.02
288.15	$10^4\alpha$	10.7	9.48	8.54	7.12	6.14	5,44	4.95	4.77
	10^4 _K	8.56	7.11	6.11	4.89	4.28	4.01	3.98	4.03
298.15	$10^4\alpha$	10.8	9.50	8.53	7.22	6.38	5.81	5.37	5.17
	$10^4\kappa$	9.21	7.54	6.41	5.09	4.47	4.23	4.26	4.35
313.15	104 α	11.0	9.52	8.52	7.36	6.76	6.39	6.02	5.79
	10^4 K	10.0	8.20	6.93	5.36	4.51	4.05	3.82	3.78
323.15	$10^4\alpha$	11.2	9.53	8.81	7.45	7.02	6.77	6.44	6.20
	10^4 K	11.0	8.74	7.24	5.48	4.58	4.12	3.94	3.92
				(b) Trichloromethane					
278.15	104α	12.3	11.2	10.2	8.66	7.67	7.00	6.48	6.22
	10^4 K	9.08	7.36	6.26	5.09	4.66	4.68	5.06	5.39
288.15	$10^4\alpha$	12.4	11.0	9.98	8.47	7.52	6.88	6.37	6.11
	$10^4\kappa$	9.74	7.81	6.59	5.31	4.84	4.86	5.26	5.60
298.15	104α	12.5	10.9	9.79	8.27	7.36	6.76	6.27	6.01
	$10^4\kappa$	10.2	8.31	7.00	5.44	4.64	4.26	4.14	4.16
313.15	$10^4\alpha$	12.6	10.8	9.49	7.96	7.12	6.58	6.11	5.86
	$10^4\kappa$	11.6	9.19	7.63	5.90	5.12	4.85	4.92	5.06
323.15	104 α	12.7	10.6	9.29	7.74	6.95	6.45	5.99	5.73
	10^4 K	12.5	9.81	8.02	5.97	4.95	4,43	4.22	4.20
				(c) Dichloromethane					
278.15	$10^4\alpha$	13.3	11.8	10.7	9.15	7.66	6.33	5.47	5.34
	$10^4\kappa$	8.68	7.26	6.23	4.88	4.09	3.60	3.32	3.23
288.15	104α	13.5	11.9	10.8	9.18	7.95	6.94	6.16	5.88
	10^4 K	9.44	7.76	6.62	5.26	4.58	4.28	4.24	4.29
298.15	$10^4\alpha$	13.7	12.1	10.8	9.21	8.24	7.57	6.87	6.43
	10^4 K	10.2	8.31	7.03	5.51	4.74	4.39	4.30	4.34
				(d) Acetonitrile					
278.15	104 α	13.2	11.4	10.2	8.61	7.25	5.84	4.38	3.67
	$10^4\kappa$	9.28	7.94	6.92	5.53	4.66	4.10	3.73	3.61
288.15	104 a	13.5	11.7	10.49	8.83	7.60	6.45	5.32	4.77
	$10^4\kappa$	10.6	8.65	7.35	5.85	5.17	4.95	5.04	5.19
298.15	$10^4\alpha$	13.8	12.0	10.7	9.06	798	7.09	6.28	5.90
	$10^4\kappa$	11.1	9.11	7.74	6.07	5.18	4.73	4.55	4.54
313.15	$10^4\alpha$	14.3	12.4	11.0	9.42	8.54	8.06	7.77	7.66
	$10^4\kappa$	12.6	10.1	8.43	6.60	5.78	5.50	5.60	5.76
323.15	104α	14.6	12.7	11.3	9.66	8.93	8.71	8.76	8.81
	$10^4\kappa$	13.7	10.7	8.85	6.68	5.61	5.11	4.95	4.96

Table II. Isobaric Expansivities $(\alpha, \text{ in } K^{-1})$ and
Isothermal Compressibilities $(\kappa, \text{ in } MPa^{-1})$

	Property	p(MPa)							
T(K)		0.1	25	50	100	150	200	250	275
				(e) Aniline					
278.15	$10^4\alpha$	(8.3)	(7.7)	(7.4)					
	$10^4\kappa$	(4.6)	(3.8)	(3.4)					
288.15	104α	(8.4)	(7.9)	(7.5)	(6.8)				
	$10^4\kappa$	(4.4)	(4.0)	(3.7)	(3.1)				
298.15	$10^4\alpha$	8.53	7.99	7.56	6.75	6.52			
	$10^4\kappa$	4.80	4.21	3.84	3.34	2.73			
313.15	$10^4\alpha$	8.66	8.19	7.68	7.00	6.49	6.23	5.59	
	10^4 _K	5.12	4.54	4.09	3.44	2.99	2.64	2.34	
323.15	$10^4\alpha$	8.75	8.32	7.76	7.16	6.46	6.19	5.56	
	$10^4\kappa$	5.16	4.72	4.31	3.58	3.05	2.71	2.56	
				(f) Dodecane					
278.15	104α	(9.6)	(8.6)	(7.9)					
	$10^4\kappa$	(9.3)	(7.2)	(6.1)					
288.15	104α	(9.6)	(8.5)	(7.8)	(6.6)				
	$10^4\kappa$	(10)	(7.6)	(6.4)	(4.6)				
298.15	104α	9.68	8.47	7.66	6.56	5.76			
	10^4 _K	9.64	8.02	6.73	5.07	4.43			
313.15	104α	9.76	8.41	7.49	6.39	6.04	5.68	4.35	
	$10^4\kappa$	11.0	8.63	7.10	5.34	4.32	3.56	2.87	
323.15	$10^4\alpha$	9.81	8.38	7.37	6.28	6.22	5.64	4.35	
	$10^4\kappa$	11.6	9.08	7.43	5.46	4.38	3.71	3.26	

Table II. *(Continued)*

but values of π calculated in that fashion are more uncertain than the values listed for toluene, trichloromethane, dichloromethane, and acetonitrile.

3.1. Comparisons with Literature Data

3.1.1. Toluene

The density of toluene at four temperatures (223, 248, 273, and 298 K) and pressures up to about 200 MPa has been determined by Mopsik [8] and expressed in terms of a Tait-type equation at each temperature. Volume ratios evaluated from the parameters given by Mopsik at several pressures at 298.15 K are compared with those obtained in the present work in Table IV, which includes a comparison of values, interpolated at

			T(K)		
V_m (cm ³ · mol ⁻¹)	278.15	288.15	298.15	313.15	323.15
		(a) Toluene			
95.01 95.97 96.92 97.88 98.84 99.80 100.76 101.72 102.68 103.64 104.59	253 265 283 303 322 338 350 360 365 368 369	277 286 300 314 327 338 347 353 358 360 360	298 307 316 324 332 338 343 347 349 351 351	327 336 340 340 340 338 337 337 336 336 337	355 361 359 352 345 339 333 330 328 327 327
		(b) Trichloromethane			
71.09 71.85 72.61 73.37 74.13 74.89 75.64 76.40 77.16 77.92 78.68 55.93 56.62 57.30 57.99 58.67 59.36 60.04 60.72 61.41 62.09	406 387 383 386 391 396 400 401 399 395 389 273 313 357 394 417 429 432 431 427 424	383 371 369 372 376 381 384 385 384 381 377 (c) Dichloromethane 322 347 373 393 408 416 420 421 420 418	358 $\frac{555}{355}$ 355 358 362 365 367 368 368 366 364 372 383 389 393 398 404 408 411 413 412	315 327 333 337 339 341 342 343 343 344 344	292 310 318 321 322 323 323 324 326 328 331
62.78	422	416	409		
		(d) Acetonitrile			
46.39 46.89 47.40 47.90 48.40 48.91 49.41 49.92 50.42 50.93 51.43	286 303 317 329 338 346 353 360 367 375 383	312 323 333 341 348 353 358 363 367 371 377	336 342 347 352 355 358 361 363 365 367 369	373 371 371 370 370 369 368 366 364 361 358	403 394 387 383 379 376 372 368 363 357 350

Table III. Internal Pressures $(\pi, \text{ in MPa})$

 \bar{z}

278.15 and 288.15 K from Mopsik's parameters, with our data. The two sets of data are clearly in very good agreement at pressures up to at least 100 MPa; the only significant divergence is at pressures approaching the upper limit of Mopsik's experimental pressure range.

The compressibility of toluene at 0.1 MPa has been measured over a wide temperature range, and literature data [9-11] are compared with values from Eq. (5) in Fig. 1. The present data are in close agreement with the literature data, even though our values of κ for pressures below 50 MPa are less accurate than at higher pressures.

3.1.2. Trichloromethane and Dichloromethane

Volume ratios for these liquids at 298 K have been reported by Newitt and Weale $\lceil 12 \rceil$ for pressures up to about 100 MPa and their values at selected pressures are compared with our data in Table IV. The two sets of values for both liquids agree to within about 0.1% , which we believe to be the accuracy of the previous data.

Compressibilities for trichloromethane at 0.1 MPa are compared with literature data $\lceil 12^{-14} \rceil$ in Fig. 2. The agreement is within the probable uncertainty in our values of κ at 0.1 MPa. Compressibilities for both trichloromethane and dichloromethane at 298 K and 100 MPa agree with values from Newitt and Weale [12] to within the expected uncertainty in κ .

It is interesting to note that Newitt and Weale reported identical values for the volume ratios (and hence κ) of trichloromethane and dichloromethane at 298 K. The very close similarity at low pressures has been confirmed in the present work, but at higher pressures (and particularly at low temperatures) those properties of these two liquids diverge significantly.

3.1.3. Acetonitrile

Relative densities (i.e., values of $1/k$) for acetonitrile at 283, 298, and 313 K and pressures up to 250 MPa have been reported by Srinivasan and Kay [15], and their data (expressed as volume ratios) are compared with values from the present work in Table IV. There are discrepancies of up to about 0.8 % between the two sets of data, and the magnitude of the discrepancies is not easily accounted for in view of the established reliability of the bellows volumometer used in the present work. We believe that there are practical problems associated with the piston-type volumometer which was used by Srinivasan and Kay and that, despite the calibration check

 $\bar{1}$

f Volume Patios with Literature Data

 \overline{a}

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' (a) This work.

b The values from Ref. 17 were interpolated from data at other temperatures.

The temperatures are 298.27, 310.95, and 323.35 K. The data from this work were interpolated using the parameters in Table I.

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Fig. 1. Isothermal compressibility for toluene at 0.1 MPa. \bullet . This work; \triangle , Ref. 9; Φ , Ref. 10; Φ , Ref. 11.

Fig. 2. Isothermal compressibility for trichloromethane at 0.1 MPa . \bullet , This work; \triangle , Ref. 9; \oplus , Ref. 10; ∇ , Ref. 12; \square , Ref. 13; **(a**, Ref. 14.

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which was made, using water, by these workers, their volume-ratio data for acetonitrile are significantly in error.

The compressibility of acetonitrile at 0.1 MPa has been directly measured by Grant-Taylor and MacDonald [16]: their values of κ at 298.15 and 318.15 K, 11.54×10^{-4} and 13.00×10^{-4} MPa⁻¹, respectively, are in good agreement with the values $(11.07 \times 10^{-4}$ and 13.18×10^{-4} MPa⁻¹, the latter value obtained by interpolation) from the present work.

3.1.4. Aniline

Volume ratios for aniline at temperatures from 298 to 358 K and pressures up to 100 MPa have been reported by Gibson and Loeffler [7, 17]. Data from those sources are compared with our data in Table IV. The agreement is good: only at 100 MPa and 298.15 K does the discrepancy exceed 0.02 %.

Compressibilities at 0.1 MPa are compared with literature data [5, 18] in Fig. 3. The agreement is satisfactory except possibly at 298 K, where the discrepancy is larger than we would have expected. Values of κ at 100 MPa are compared with values from Gibson and Loeffler in the inset in Fig. 3. The agreement is within the uncertainty in κ , except at 298 K. It appears that, for reasons which are not apparent, our values of κ at 298 K may be $2-3\%$ high.

Fig. 3. Isothermal compressibility for aniline at 0.1 MPa and (inset) 100 MPa. \bullet , This work; \triangle , Ref. 9; Φ , Ref. 7.

3.1.5. n-Dodecane

 $p.V.T$ data for *n*-dodecane have been reported by Cutler and coworkers [18] and by Dymond and co-workers [19], for temperatures of 298.17 K and above. Volume ratios calculated from these data are compared with values from the present work interpolated at 298.27 and 310.95 K and extrapolated to 323.35 K in Table IV. The agreement is satisfactory for pressures up to 50 MPa but at higher pressures there is an increasing divergence between the sets of data. In view of the general concordance between our volume ratio-ratio data (or derived values for quantities such as excess volume) and literature data, for liquids other than ndodecane at pressures up to 250 MPa, we believe that there is some doubt concerning the reliability of the literature data for n-dodecane at elevated pressures. We have previously commented on the necessity to calibrate bellows-type volumometers over the whole pressure range at which they are to be used. The discrepancies which are apparent in the data in Table IV may be due to the fact that the (bellows-type) volumometers which were used to obtain the previous data for n -dodecane were calibrated only at 0.1 MPa.

Compressibilities at 0.1 MPa are compared with literature data $\lceil 18, 20 \rceil$ in Fig. 4. The agreement is within 2.5% with the directly measured values [20]. As for aniline, our values at the lower temperatures

Fig. 4. Isothermal compressibility for n-dodecane at 0,1MPa and (inset) 100 MPa. \bullet , This work; \triangle , Ref. 20; \square , Ref. 18.

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are only approximate, owing to the limited accessible pressure range. The agreement between our values of κ and literature data at 100 MPa, shown by the inset in Fig. 4, is good.

4. DISCUSSION

The volume dependence at constant temperature of the internal pressure of toluene is shown in Fig. 5 and has some striking features. At large (molar) volumes, i.e., low pressures, the V, T dependence of π is essentially that of a typical nonassociated liquid. However, this behavior is not preserved as the volume is decreased, and $(\partial \pi/\partial T)_v$ changes sign at about $100 \text{ cm}^3 \cdot \text{mol}^{-1}$. At elevated pressures (roughly above 75 MPa at 278 K) and 125 MPa at 323 K) π becomes quite strongly temperature dependent and the behavior becomes more like that of an associated liquid. The

Fig. 5. Volume and temperature dependence of the internal pressure for toluene. \Box , 278 K; \blacksquare , 288 K; \odot , 298 K; \bigcirc , 313 K; \triangle , 323 K.

behavior at high pressures is consistent with strong short-range intermolecular interactions becoming effective as the average interparticle separation decreases. No specific chemical (e.g., hydrogen-bonding) interactions can occur in toluene and the molecule has only a small dipole moment, but the quadrupole moment is relatively large. It seems likely therefore, that multipole-multipole interactions are moderately strong in compressed toluene.

For dichloromethane, the volume and temperature dependence of π (Fig. 6) is qualitatively similar to that of toluene. At large average interparticle distances the behavior is characteristic of nonassociated liquids, but when interparticle distances are sufficiently reduced by the application of pressure, short-range intermolecular interactions can evidently become significant. The dipole moment of dichloromethane is substantial (about

Fig. 6. Volume and temperature dependence of the internal pressure for dichloromethane. Symbols as in the legend to Fig. 5.

 5×10^{-30} Cm) so that the variation of π with V and T at high pressures can be ascribed principally to dipole-dipole interactions.

The progressive change in the *V*,*T* dependence of π shown in Figs. 5 and 6, including the increase in $(\partial \pi/\partial T)$, at low volume, is even more pronounced for acetonitrile (Fig. 7), for which there is only a small volume range (at low pressure) where behavior typical of a nonassociated liquid [i.e., $(\partial \pi / \partial T)_v > 0$] is observed. The behavior of acetonitrile, compared with toluene and dichloromethane, is consistent with the larger dipole moment (of the order 13×10^{-30} Cm) for acetonitrile. It is interesting to note that for the extensively hydrogen-bonded liquid methanol, for which internal pressure calculated from the p, V, T data in Ref. 4 are shown in Fig. 8, the volume dependence of π is very like that of acetonitrile (except for some complexities near the highest experimental pressures), but there is virtually no volume range for which nonassociated liquid behavior is obser-

Fig. 7. Volume and temperature dependence of the internal pressure for acetonitrile. Symbols as in the legend to Fig. 5.

Fig. 8. Volume and temperature dependence of the internal pressure for methanol. Symbols as in the legend to Fig. 5.

ved (for pressures up to about 280 MPa). The extreme case of associated liquid behavior is that shown by water (Fig. 9), for which π was evaluated from literature data [21-23]. A feature of the *V*,*T* dependence of π for water is that $(\partial \pi/\partial T)$, decreases with increasing pressure, presumably due to, first, distortion and, eventually, disruption of the hydrogen-bonded network structure as the pressure is increased. It seems, from Fig. 8, that a similar effect occurs in methanol when the molar volume is reduced to less than about 36 cm³ · mol⁻¹. Apparently at higher densities the molecules are forced to pack in such a manner that fewer extended hydrogen-bonded structures can exist.

Following the analyses due to Wiehe and Bagley [24] and Bagley and co-workers [25], it is instructive to compare the product πV_m with the cohesive energy calculated as $E_y + 3/2RT$, where E_y is the energy of vaporization evaluated from $\Delta H_v - RT$. Such a comparison, based on ΔH_v

Fig. 9. Volume and temperature dependence of the internal pressure for water. Symbols as in the legend to Fig. 5.

values from Ref. 6, is presented in Table V (for 298 K and 0.1 MPa). The quantity of interest in Table V is δ , the difference between $E_v + 3/2RT$ and πV_m expressed as a percentage of $E_v + 3/2RT$. For the nonpolar, nonasociated carbon tetrachloride δ is almost zero, and πV_m accounts for virtually the whole magnitude of the cohesive energy. The increase in δ from toluene to dichloromethane, acetonitrile, and methanol, which can be interpreted as an increase in the contribution from dipole-dipole (or higherorder multipole) interactions to the cohesive energy, is consistent with the change in the observed V, T dependence of π for these liquids.

The negative value of δ in Table V for *n*-dodecane is difficult to explain, since the uncertainty in π at 298 K and 0.1 MPa should not be sufficiently large to lead to an error of 15% in πV_m .

The case of trichloromethane is clearly anomalous. The *V,T* dependence of π (Fig. 10) is complex but, nonetheless, typical of a nonassociated rather than an asociated liquid. As for toluene, $(\partial \pi / \partial T)_v$ increases rapidly

Liquid	$\Delta E_{v}^{b} + (3/2) RT$	$\pi V_{\rm m}$	δ (%) ^c
Methanol	38.67	11.87	69
Acetonitrile	34.46	19.76	43
Trichloromethane	34.55	29.85	14
Aniline	57.00	49.83	13
Dichloromethane	29.80	27.09	9
Toluene	39.23	37.61	4
Carbon tetrachloride ^d	33.65	33.38	
n -Dodecane	62.54	71.78	-15

Table V. Comparison of the Product πV_m^a with Cohesive Energy,^a at 298 K and 0.1 MPa

^a In kJ·mol⁻¹.

^b Calculated as $\Delta H_v - RT$, where ΔH_v is the enthalpy of vaporization. ΔH_v values from Ref. 6.

^cThe difference between $\Delta E_v + (3/2)RT$ and πV_m , expressed as a percentage of AE _v + (3/2) RT.

^d Molar volume and internal pressure from Ref. 26.

Fig. 10. Volume and temperature dependence of the internal pressure for trichloromethane. Symbols as in the legend to Fig. 5.

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in magnitude as the volume is decreased (below about $73 \text{ cm}^3 \cdot \text{mol}^{-1}$ for **trichloromethane), but unlike toluene the negative temperature coefficient** of π is maintained as the pressure is increased. In view of the substantial value of δ in Table V, it is surprising that the V,T dependence of π is not **intermediate between that of dichloromethane and that of acetonitrile.**

REFERENCES

- 1. P. J. Back, A. J. Easteal, R. L. Hurle, and L. A. Woolf, J. *Phys. E Sci. Instrum.* 15:360 (1982).
- 2. A. J. Easteal and L. A. Woolf, *J. Chem. Thermodyn.* 14:755 (1982).
- 3. G. G6etze and G. M. Schneider, *J. Chem. Thermodyn.* 12:661 (1980).
- 4. A. J. Easteal and L. A. Woolf, *J. Chem. Thermodyn.* 17:49 (1985).
- 5. A. J. Easteal and L. A. Woolf, *J. Chem. Thermodyn.* 17:69 (1985).
- 6. J. Timmermans, *Physico-chemical Constants of Pure Organic Compounds* (Elsevier, Amsterdam, 1970).
- 7. R. E. Gibson and O. H. Loeltler, *J. Am. Chem. Soc.* 61:2515 (1939).
- 8. F. I. Mopsik, *J. Chem. Phys.* 50:2259 (1969).
- 9. D. Tyrer, *J. Chem. Soc.* 105, 2534 (1914).
- 10. E. B. Fryer, J. C. Hubbard, and D. H. Andrews, *J. Am. Chem. Soc.* 51:759 (1929).
- 11. J. G. Marshall, L. A. K. Staveley, and K. R. Hart, *Trans. Faraday Soc.* 52:23 (1956).
- 12. D. M. Newitt and K. E. Weale, *J. Chem. Soc.* 3092 (1951).
- 13. L. A. K. Staveley, W. I. Tupman, and K. R. Hart, *Trans. Faraday Soc.* 51:323 (1955).
- 14. D. Harrison and E. A. Moelwyn-Hughes, *Proc. Roy. Soc.* A239:230 (1957).
- 15. K. R. Srinivasan and R. L. Kay, *J. Solut. Chem.* 6:357 (1957).
- 16. D. F. Grant-Taylor and D. D. MacDonald, *Can. J. Chem.* 54:2813 (1976).
- 17. R. E. Gibson and O. H. Loeffler, *J. Phys. Chem.* 43:207 (1939).
- 18. W. G. Cutler, R. H. McMickle, W. Webb, and R. W. Schiessler, *J. Chem. Phys.* 29:727 (1958).
- 19. J. H. Dymond, J. Robertson, and J. D. Isdale, *Int. J. Thermophys.* 2:133 (1981).
- 20. M. Diaz Pena and G. Tardajos, *J. Chem. Thermodyn.* 10:19 (1978).
- 21. E. H. Amagat, *Ann. Chim. Phys.* 29:68, 505 (1893).
- 22. G. S. Kell and E. Whalley, *J. Chem. Phys.* 62:3496 (1975).
- 23. T. Grindley and J. E. Lind, *J. Chem. Phys.* 54:3983 (1971).
- 24. I. A. Wiehe and E. B. Bagley, *A. I. ChE. J.* 13:836 (1967).
- 25. E. B. Bagley, T. P. Nelson, J. W. Barlow, and S. A. Chen, *Ind. Eng. Chem. Fund.* 9:93 (1970).
- 26. R. E. Gibson and O. H. Loeffler, J. *Am. Chem. Soc.* 63:898 (1941).