

## On the Differential Equations of Species in Competition

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### Summary

It is shown that the ordinary differential equation commonly used to describe competing species are compatible with any dynamical behavior provided the number of species is very large.

The goal is to show that the ordinary differential equations used in ecology to describe competing species do not say much in case the number of species is more than three or four. In fact these equations are compatible with any dynamical behavior in a certain reasonable sense.

More precisely consider an ecological system of  $n$  competing species with state space

$$R_+^n = \{x \in R^n \mid x = (x_1, \dots, x_n), x_i \geq 0\}$$

where  $x_i$  represents the population of the  $i$ -th space.

We suppose the dynamics of population growth is given by the system of ordinary differential equations:

$$\frac{dx_i}{dt} = x_i M_i(x), \quad i = 1, \dots, n$$

For background on these matters see references [1], [2], [3], [4], [5], [6].

We will impose the following conditions:

- (1)  $M_i : R_+^n \rightarrow R$  are  $C^\infty$  (i.e., have continuous derivative to all orders).
- (2) For all pairs  $i$  and  $j$  between 1 and  $n$ , if  $x_i > 0$ ,  $\frac{\partial M_i(x)}{\partial x_j} < 0$ . This is a classic condition of competition which can be interpreted simply as "crowding inhibits growth".
- (3) There is a constant  $K > 0$  with the property: For each  $i$ ,  $M_i(x) < 0$  if  $\|x\| > K$  (i.e., the planet is finite).

**Example 1:** Let  $M_i(x) = 1 - \sum_{j=1}^n x_j$  for each  $i=1, \dots, n$ . This choice satisfies the conditions (1)—(3). In this case the boundary  $\partial R_+^n = \{x \in R_+^n \mid \text{some } x_i = 0\}$  is invariant under the flow generated by the differential equations and all other solutions tend as  $t \rightarrow \infty$  to the invariant “attractor”

$$\Delta_1 = \{x \in R_+^n \mid \sum x_i = 1\}.$$

On  $\Delta_1$  the flow is stationary, a very degenerate situation which is remedied by the next example.

**Example 2:** Define  $\Delta_0 = \{x \in R^n \mid \sum x_i = 0\}$ ,  $1_0 = (1, \dots, 1) \in R_+^n$  and let  $\beta: R \rightarrow R$  be a  $C^\infty$  function which is 1 in a neighborhood of 1 and  $\beta(t) = 0$  if  $t \leq \frac{1}{2}$  or if  $t \geq 1 - \frac{1}{2}$ . Now define  $h: R_+^n \rightarrow \Delta_0$  by

$$h(x) = \left(\frac{1}{n}\right) 1_0 - \frac{x}{\sum x_i}$$

and let  $m_i: R_+^n \rightarrow \Delta_0 \subset R^n$  be defined by

$$m_i(x) = \frac{1}{x_i} \beta(\sum x_i) \left(\prod_{j=1}^n x_j\right) h(x).$$

The equations of Example 2 are then  $\frac{dx_i}{dt} = x_i(M_i + \eta m_i) = x_i N_i$ ,  $\eta > 0$ , where the  $M_i$  are as in Example 1. Then clearly  $N_i: R_+^n \rightarrow R$  is  $C^\infty$  and  $N_i$  satisfies Property (2) since for large  $\|x\|$ ,  $N_i = M_i$ . Also if we choose  $\eta$  small enough, which we do, then (3) is also satisfied.

Let us check properties of solutions of  $\frac{dx_i}{dt} = x_i N_i$ . Any solution  $t \rightarrow x(t)$  not lying on the boundary of  $R_+^n$  must satisfy  $\sum x_i(t) \rightarrow 1$  as  $t \rightarrow \infty$ . This follows since  $\sum m_i = 0$ , and from the form of the  $M_i$ . Thus  $\Delta_1$  is the attracting set. On  $\Delta_1$ , the  $M_i = 0$ , and the flow is determined by the  $\eta x_i m_i$ , or up to a scalar factor,  $h$  restricted to  $\Delta_1$ . On  $\Delta_1$ , the differential equation  $\frac{dx}{dt} = h(x)$  has  $\frac{1}{n} 1_0$  as a global linear sink. Thus in Example 2 every solution not on the boundary  $\partial R_+^n$ , tends to  $\frac{1}{n} 1_0 \in R_+^n$  as  $t \rightarrow \infty$ .

**Example 3:** Let  $h_0: \Delta_1 \rightarrow \Delta_0$  be any  $C^\infty$  map and  $h: R_+^n \rightarrow \Delta_0$  any  $C^\infty$  map which agrees with  $h_0$  on  $\Delta_1$ . Example 3 then is Example 2 where  $h$  is replaced by the above. As in Example 2 with  $\eta$  chosen small enough, conditions (1), (2), and (3) will be satisfied. Furthermore as in Example 2,  $\Delta_1$  will be the attractor for all non-boundary solutions, so that after a period of transitions, the dynamics of Example 3 will be the dynamics of  $\frac{dx}{dt} = h_0(x)$  on  $\Delta_1$  (up to a scalar factor, which doesn't affect the qualitative behavior). But this was arbitrary. The system  $\frac{dx}{dt} = h_0(x)$  could be any prescribed system!

We end with a sequence of remarks.

1. Let  $n=2$ . Then in Example 3, since  $\Delta_1$  is one-dimensional, solutions on  $\Delta_1$  tend to equilibria which will be stable in general. Thus the same will be true in general for solutions in Example 3. This is consistent with earlier results, see e. g., [5], [1].
2. If  $n=3$ , in Example 3 one can put on  $\Delta_1$  a system with a stable limit cycle  $\gamma$ ; in fact we can suppose that every solution on  $\Delta_1$  will tend to  $\gamma$  and that perturbations of this system on  $\Delta_1$  will have the same properties. Thus every solution in the interior of  $R_+^n$  will oscillate after a period of transition and this property will be true for perturbed systems. Compare this to the results of May and Leonard [4].
3. Suppose  $n \geq 5$ . Then since  $\Delta_1$  is at least 4-dimensional, one may impose on  $\Delta_1$  many of the complicated dynamical systems that have been found in recent years (see e. g., [7]). In particular the system of Example 3 may not be approximated by a structurally stable, dynamical system, or it may have strange attractors with an infinite number of periodic solutions, etc.
4. If the constant  $\eta$  as in Example 3 has been chosen small enough, then by the theory of invariant manifolds, after small perturbations of the system, there will remain an invariant manifold near  $\Delta_1$ . Then in fact if the system on  $\Delta_1$  is structurally stable so will be the system on  $R_+^n$ .

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