# On the addition of degrees of freedom to force-balanced linkages.

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Abstract The design of shaking-force balanced linkages can be approached by deriving these linkages from balanced linkage architectures. When desired, a possible step is to add degrees-of-freedom (dof), for instance by substituting a link with a n-dof equivalent linkage for which the balanced design of the other links is not affected. This paper shows how the coupler link of a shaking-force balanced 4R four-bar linkage, applied as a 5R five-bar linkage, can be substituted with an equivalent 2-dof pantograph.

### 1 Introduction

With the increasing speed of manipulators (i.e. mechanisms, robotics), for instance for pick and place tasks, dynamic properties such as shakingforce balance and shaking-moment balance become increasingly important. Acceleration of mass and inertia of moving parts of balanced manipulators do not cause any forces and moments to the base and surrounding, keeping machine vibrations low.

Instead of balancing a manipulator linkage afterwards, it is advantageous to base the design of the linkage on balance properties to minimize complexity, additional mass, and additional inertia (Van der Wijk et al. (2009)). One approach for this is to compose manipulators from balanced linkage sections such as balanced legs (Arakelian and Smith (1999), Wu and Gosselin (2002)). Another approach is to derive manipulators from inherently balanced architectures (Van der Wijk and Herder (2012a)), i.e. linkage architectures that are balanced due to specific kinematic relations. As long as these kinematic relations are maintained, any change to the linkage can be made without affecting the balance properties, for instance by fixing links together and by replacing links with gears.

This paper shows the possibility of adding degrees-of-freedom (dof) to a force-balanced linkage architecture. The substitution of a link with a 2 dof linkage is investigated for which the balanced design of other links is



Figure 1. a) Balanced 4R four-bar linkage with CoM at *S* at link 4; b) Balanced four-bar linkage with coupler *l*<sup>2</sup> replaced with a pantograph.

not affected. As subject of investigation the coupler link of a shaking-force balanced 4R four-bar linkage, applied as a balanced 5R five-bar linkage, is chosen. First an equivalent model of the coupler link is derived with linear momentum equations and subsequently the conditions for an equivalent 2 dof substitute linkage, a pantograph, are obtained.

#### 2 Equivalent model of coupler link

Figure 1a shows a four-bar linkage  $A_0A_1A_2A_3$  of which each link  $l_i$  has a mass  $m_i$  at its link center-of-mass (link CoM)  $S_i$  which are defined with parameters  $p_i$  and  $q_i$  as indicated. From Berkof and Lowen (1969) and by including the mass  $m_4$  of link 4, the balance conditions for which the CoM of the complete linkage is at an invariant point *S* in link 4 are written as

$$
p_1 = -\frac{m_2(l_2 - p_2)}{l_2} \frac{l_1}{m_1} \quad q_1 = \frac{m_2 q_2}{l_2} \frac{l_1}{m_1} \quad c_1 = \frac{1}{m_{tot}} (m_1 l_4 + \frac{m_2(l_2 - p_2)}{l_2} l_4 + m_4 p_4) p_3 = l_3 + \frac{m_2 p_2}{l_2} \frac{l_3}{m_3} \quad q_3 = \frac{m_2 q_2}{l_2} \frac{l_3}{m_3} \quad c_2 = \frac{1}{m_{tot}} (m_4 q_4 - \frac{m_2 q_2}{l_2} l_4)
$$
 (1)

with  $m_{tot} = m_1 + m_2 + m_3 + m_4$ . The linkage of Fig. 1a then is a forcebalanced four-bar linkage when link 4 is stationary with the base and it is a force-balanced five-bar linkage when solely *S* is stationary with the base as being a movable joint.

To substitute a link without affecting the other links, the substitute linkage has to be equivalent. Therefore first an equivalent model of the coupler link is derived from which the substitute linkage can be found.



Figure 2. a-b) Coupler link; c) Equivalent Linear Momentum System

With linear momentum equations an equivalent linear momentum system (ELMS) of the coupler can be created. Figures 2a-b show link *l*2, of which the linear momentum can be written about  $A_2$  w.r.t. frame  $x_1y_1$ (Fig. 2a) and about  $A_1$  w.r.t. frame  $x_2y_2$  (Fig. 2b) respectively as

$$
\frac{\overline{L}_1}{\dot{\theta}} = \begin{bmatrix} m_2(l_2 - p_2) \\ -m_2 q_2 \end{bmatrix} \quad \frac{\overline{L}_2}{\dot{\theta}} = \begin{bmatrix} m_2 p_2 \\ m_2 q_2 \end{bmatrix}
$$
 (2)

Figure 2c shows a model of the coupler with masses  $\mu_1$  in  $A_1$ ,  $\mu_2$  in  $A_2$ , and mass  $\nu_1$  at distance  $l_2$  normal to line  $A_1A_2$  as indicated. Similarly, the linear momentum equations w.r.t. each of the two frames about  $A_1$  and  $A_2$ , respectively, can be written as

$$
\frac{\overline{L}_1}{\theta} = \begin{bmatrix} \mu_1 l_2 \\ -\nu_1 l_2 \end{bmatrix} \quad \frac{\overline{L}_2}{\theta} = \begin{bmatrix} \mu_2 l_2 \\ \nu_1 l_2 \end{bmatrix} \tag{3}
$$

This implies that the model of Fig. 2c is equivalent to Fig. 2a-b for

$$
\mu_1 = \frac{m_2(l_2 - p_2)}{l_2} \quad \mu_2 = \frac{m_2 p_2}{l_2} \quad \nu_1 = \frac{m_2 q_2}{l_2} \tag{4}
$$

Also other equivalent models are possible, e.g. with a mass  $\nu_1$  at  $l_2$  above *A*<sub>1</sub>. When  $S_2$  is on the line  $A_1A_2$ ,  $q_2$  is zero for which  $\nu_1$  becomes zero too.

#### 3 Equivalent pantograph linkage as substitute

To add one dof, *l*<sup>2</sup> can be substituted with the 2-dof linkage. Figure 3a shows an ELMS of 2-dof linkage  $A_1B_2A_2$  defined with  $\mu_1$  in  $A_1$ ,  $\mu_2$  in  $A_2$ , and  $\nu_1$ at two locations normal to and at equal distance from  $B_2$  as, respectively, *l*<sup>21</sup> and *l*<sup>22</sup> as indicated. Also here multiple models can be found for equal ELMSs.



Figure 3. a) Equivalent model of a 2-DoF linkage to replace the coupler; b) Model of the coupler with the CoM of the four masses in *S*2.

From Eqs. 1 follows that  $\mu_1$ ,  $\mu_2$ , and  $\nu_1$  need to remain constant to not affect the balance parameters of the other links. Since  $l_2$  is not constant any longer,  $\mu_1$ ,  $\mu_2$ , and  $\nu_1$  can be written as

$$
\mu_1 = m_2(1 - \kappa_1) \quad \mu_2 = m_2 \kappa_1 \quad \nu_1 = m_2 \kappa_2 \tag{5}
$$

with  $\kappa_1 = p_2/l_2$  and  $\kappa_2 = q_2/l_2$  to be constant for any value of  $l_2$ . This implies that for all lengths  $l_2$ , triangle  $A_1A_2S_2$  has to be similar. In general a real linkage  $A_1B_2A_2$  cannot generate this similarity. A mechanism which is characterized for its properties of similarity is the pantograph linkage (Artobolevskii (1964)). Figure 1b shows how this linkage can be applied to substitute the coupler link.

Figure 4 shows the substitute pantograph linkage in detail, consisting of four links arranged as parallelogram linkage  $B_1B_2B_3B_4$  with each a mass  $m_{2i}$  located at distances  $e_{2i}$  and  $f_{2i}$  from joints  $B_i$  as indicated. The total mass of the linkage then is written as  $m_2 = m_{21} + m_{22} + m_{23} + m_{24}$ . The parallelogram linkage, and in specific joints  $B_1$  and  $B_3$ , are defined with principal lengths  $a_1$  and  $a_2$  from  $B_2$ , respectively, and with angles  $\alpha_1$  and  $\alpha_2$  with the lines  $A_1B_2$  and  $B_2A_2$ , respectively.

The pantograph linkage is equivalent to the coupler link when its CoM is located at  $S_2$  of the similar triangle  $A_1A_2S_2$  at all times. To find the conditions for which this holds, the linear momentum of the pantograph linkage can be written to be equal to the linear momentum of the ELMS of Fig. 3a. These equations can be written for each dof individually as shown in Van der Wijk and Herder (2012a). The linear momentum for  $\dot{\theta}_1$  with



Figure 4. Pantograph linkage  $A_1B_1B_2B_3B_4A_2$  and its parameters.

respect to frame  $x_1y_1$  with  $\dot{\theta}_2 = 0$  and  $A_2B_2$  fixed can be written as

$$
\frac{\overline{L}_{21}}{\dot{\theta}_1} = \begin{bmatrix} u_1 l_{21} \cos \alpha_1 + v_1 l_{21} \sin \alpha_1 \\ u_1 l_{21} \sin \alpha_1 - v_1 l_{21} \cos \alpha_1 \end{bmatrix} =
$$
\n
$$
\begin{bmatrix} (m_{21} e_{21} + m_{23} e_{23}) \cos \alpha_1 - (m_{21} f_{21} + m_{23} f_{23}) \sin \alpha_1 + m_{24} a_1 \\ (m_{21} e_{21} + m_{23} e_{23}) \sin \alpha_1 + (m_{21} f_{21} + m_{23} f_{23}) \cos \alpha_1 \end{bmatrix}
$$
\n(6)

and the linear momentum for  $\dot{\theta}_2$  with respect to frame  $x_2y_2$  with  $\dot{\theta}_1 = 0$ and  $\mathcal{A}_1\mathcal{B}_2$  fixed writes

$$
\frac{\overline{L}_{22}}{\dot{\theta}_2} = \begin{bmatrix} u_2 l_{22} \cos \alpha_2 + v_1 l_{22} \sin \alpha_2 \\ -u_2 l_{22} \sin \alpha_2 + v_1 l_{22} \cos \alpha_2 \end{bmatrix} = (7)
$$
\n
$$
\begin{bmatrix} (m_{22} e_{22} + m_{24} e_{24}) \cos \alpha_2 - (m_{22} f_{22} + m_{24} f_{24}) \sin \alpha_2 + m_{23} a_2 \\ -(m_{22} e_{22} + m_{24} e_{24}) \sin \alpha_2 - (m_{22} f_{22} + m_{24} f_{24}) \cos \alpha_2 \end{bmatrix}
$$
\n(7)

These equations lead to the resulting four conditions for equivalence

$$
X_{11}\cos\alpha_1 + X_{12}\sin\alpha_1 = m_{24}a_1 \quad X_{21}\cos\alpha_2 + X_{22}\sin\alpha_2 = m_{23}a_2
$$
  
\n
$$
X_{11}\sin\alpha_1 - X_{12}\cos\alpha_1 = 0 \quad X_{21}\sin\alpha_2 - X_{22}\cos\alpha_2 = 0
$$
 (8)

with

$$
X_{11} = u_1 l_{21} - m_{21} e_{21} - m_{23} e_{23} \t X_{12} = v_1 l_{21} + m_{21} f_{21} + m_{23} f_{23} X_{21} = u_2 l_{22} - m_{22} e_{22} - m_{24} e_{24} \t X_{22} = v_2 l_{22} + m_{22} f_{22} + m_{24} f_{24}
$$
 (9)

When the mass of each link and the link-CoMs are known, the locations of joints  $B_1$  and  $B_3$  are found with

$$
\tan \alpha_1 = \frac{X_{12}}{X_{11}} \qquad a_1 = \frac{1}{m_{24}} (X_{11} \cos \alpha_1 + X_{12} \sin \alpha_1)
$$
  

$$
\tan \alpha_2 = \frac{X_{22}}{X_{21}} \qquad a_2 = \frac{1}{m_{23}} (X_{21} \cos \alpha_2 + X_{22} \sin \alpha_2)
$$

or by substituting

$$
\cos \alpha_1 = \frac{b_1}{a_1} \quad \sin \alpha_1 = \frac{c_1}{a_1} \quad a_1^2 = b_1^2 + c_1^2
$$
  
\n
$$
\cos \alpha_2 = \frac{b_2}{a_2} \quad \sin \alpha_2 = \frac{c_2}{a_2} \quad a_2^2 = b_2^2 + c_2^2
$$

in the four conditions of Eqs. 8, which results in

$$
X_{11}b_1 + X_{12}c_1 = m_{24}(b_1^2 + c_1^2) \t X_{21}b_2 + X_{22}c_2 = m_{23}(b_2^2 + c_2^2) \t (10)X_{11}c_1 - X_{12}b_1 = 0 \t X_{21}c_2 - X_{22}b_2 = 0
$$

algebraic solutions are obtained for  $a_1$ ,  $b_1$ , and  $c_1$  being

$$
a_1 = \frac{1}{m_{24}} \sqrt{X_{11}^2 + X_{12}^2}, \quad b_1 = \frac{X_{11}}{m_{24}}, \quad c_1 = \frac{X_{12}}{m_{24}} \tag{11}
$$

and for  $a_2$ ,  $b_2$ , and  $c_2$  being

$$
a_2 = \frac{1}{m_{23}}\sqrt{X_{21}^2 + X_{22}^2}, \quad b_2 = \frac{X_{21}}{m_{23}}, \quad c_2 = \frac{X_{22}}{m_{23}} \tag{12}
$$

A pantograph linkage according the conditions of Eqs. 8 can replace link *l*<sup>2</sup> of Fig. 1a as shown in Fig. 1b without affecting the other links for perfect force balance of the complete linkage for all motion.

#### 4 Discussion

In addition to substituting the coupler link, also any of the other three links can be substituted with equivalent linkages to add dofs. Figure 5a shows the result when each of the four links is substituted with an equivalent pantograph with which the mechanism gains four dofs. The procedure to derive the conditions for equivalence is similar to the procedure for the coupler link with some differences due to their specific position within the linkage. Unfortunately this article leaves too few space to discuss them here in detail.

Both branches of the resulting equivalent pantographs can be used, the choice does not affect the design parameters. From Fig. 5a it is observed



Figure 5. a) Resulting linkage when all four links are substituted with a pantograph gaining four dofs; b) With  $\mu_i$  and  $\nu_1$  the mass of the coupler can be distributed equivalently to the other links. For balancing the linkage then each of the three other links can be separately investigated. For links 1 and 3 the CoM of its link mass and the equivalent masses is in joints  $A_0$  and *A*3, respectively. The linkage CoM *S* at link 4 is the CoM of its equivalent system.

that the coupler substitute pantograph is in the other branch as compared to the other three pantographs.

With the approach in this article it is also possible to substitute links with equivalent 3-dof or higher-dof linkages. These linkages will be pantographic linkages as investigated in Van der Wijk and Herder (2012a), consisting of multiple parallelograms.

The equivalent masses  $\mu_i$  and  $\nu_1$  share another feature which is illustrated in Fig. 5b and which can be obtained from Eqs. 1. With  $\mu_i$  and  $\nu_1$ modeled at link 2 as indicated,  $S_2$  is characterized as the CoM of the three equivalent masses. With  $\mu_1$  and  $\nu_1$  modeled at link 1 as indicated,  $A_0$  is the CoM of the two equivalent masses and mass  $m_1$ . With  $\mu_2$  and  $\nu_1$  modeled at link 3 as indicated,  $A_3$  is the CoM of the two equivalent masses and mass *m*3. With the three equivalent masses modeled at link 4 as indicated and with  $m_1$  in  $A_0$  and  $m_3$  in  $A_3$ , S is the CoM of the complete model.  $\mu_i$  and  $\nu_1$  therefore can be interpreted as the distribution of the mass of link 2 onto the other three links. When linkages become complex, this feature is useful for finding the balance conditions of linkages with arbitrary mass distributions without the need of loop equations, which was shown in Van der Wijk and Herder (2012b). There can be various equivalent models. For instance the initial coupler link can also be modeled as in Fig. 3b for which  $S_2$  is the CoM of the four equivalent masses.

## 5 Conclusion

For the purpose of adding degrees of freedom (dof) to shaking-force balanced linkages, the coupler link of a 4R four-bar linkage was substituted with a 2 dof equivalent pantograph linkage. An equivalent model of the coupler link was derived with linear momentum equations with which the conditions for the pantograph linkage were determined. It was discussed that the other links can be substituted in a similar way and that also higher-dof equivalent linkages can be used as substitutes. In addition, it was shown how equivalent masses can be used to distribute the coupler mass to the other three links.

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