

Surface-Based Electrophysiology Modeling and Assessment of Physiological Simulations in Atria

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Abstract. The objective of this paper is to assess a previously-proposed *surface-based* electrophysiology model with detailed atrial simulations. This model – derived and substantiated by mathematical arguments – is specifically designed to address thin structures such as atria, and to take into account strong anisotropy effects related to fiber directions with possibly rapid variations across the wall thickness. The simulation results are in excellent adequacy with previous studies, and confirm the importance of anisotropy effects and variations thereof. Furthermore, this surface-based model provides dramatic computational benefits over 3D models with preserved accuracy.

1 Introduction

There is a very important medical need for modeling the electrical activity of the heart in general, and in the atria in particular, e.g. with a view to therapy planning assistance in radiofrequency ablation for patients suffering from atrial fibrillation [10]. In addition to generic difficulties inherent to electrophysiology modeling, namely, modeling complexity and computational intensiveness, atria modeling features specific difficulties, in particular due to their very thin walls – mostly apparent as thick surfaces in medical imaging – which requires much refined meshes. Moreover, there is a major challenge in taking into account the anisotropy resulting from the preferred conduction direction along the muscle fibers, which is also known to vary extremely rapidly across the wall thickness [7].

The electrical wave propagating in the cardiac tissue can be represented by a nonlinear reaction-diffusion partial differential equation (PDE), coupled with ordinary differential equations (ODEs) representing cellular activity. In this study, we consider the 2D – namely, surface-based – model proposed and mathematically substantiated in [1], derived from the bidomain model (see for example [11–13]), and defined over the midsurface of the thin region. This surface-based model was specifically designed for thin cardiac structures – the atrial walls, in particular – and takes into account the strong anisotropy variations across the thickness. This

model was already successfully numerically assessed in [1] by comparing the resulting simulations with reference 3D simulations on thin domains of simple geometries, with dramatic benefits in computation times. Our objective here is to further assess this model with physiological simulations of the atrial electrophysiology.

2 Model

2.1 Atrial Mesh

We produced a surface mesh representing the mid-surface of the two atria. Starting from the Zygote¹ heart model – a geometric model based on actual anatomical data – we used the 3-matic² software to obtain a computationally-correct surface mesh, and the Yams [5] meshing software to further process and refine the surface mesh. Figure 1 shows the posterior and anterior views of the mesh, which contains about 26,000 triangles and 13,500 vertices. The anatomy corresponds to ventricular end-systole, namely, when the atrial chamber has its greatest internal volume. This anatomical surface was compared with literature reports of normal human atrial dimension indicators, see for example [2, 8], and found to be within standard ranges, see Table 1.

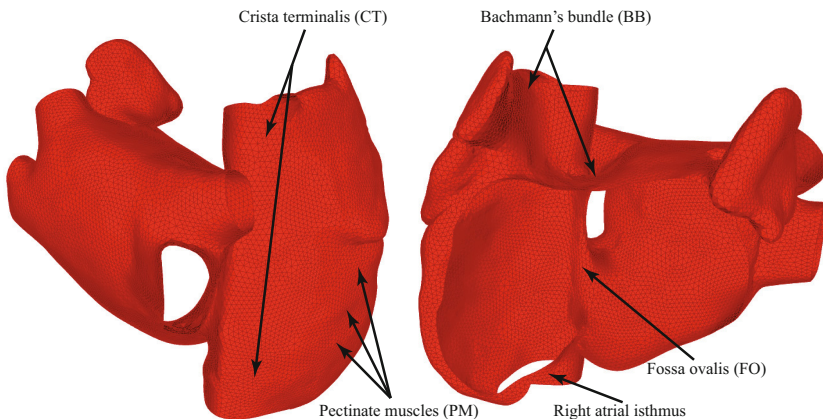


Fig. 1. Atrial mesh posterior (left) and anterior (right) views, with main specific regions

2.2 Fibers

Cardiac tissue has a fiber architecture. As the electrical conductivity is higher along than across the fiber direction, fiber orientation is very important in electric activation propagation. The specificities of the atria are that the walls are very thin, and that fibers orientations may vary extremely rapidly across the thickness.

¹ www.3dscience.com

² www.materialise.com

Table 1. Assessment of left atrial dimension indicators

LA indicators at ventricular end-systole	Normal (from [2, 8])	LA model dimensions
diameter (cm)	2.8–4.0	3.0
major axis (cm)	4.1–6.1	4.65
area (cm ²)	15.0–20.0	15.9
volume (ml)	41–75	47
diameters of pulmonary veins (cm)	1.1–1.5	1.14–1.45

We use [7, 9] to identify and prescribe the fibers directions at the endocardium and epicardium, see Figure 2. This figure also displays the angle θ , defined as half of the angular variation between the endocardium and the epicardium.

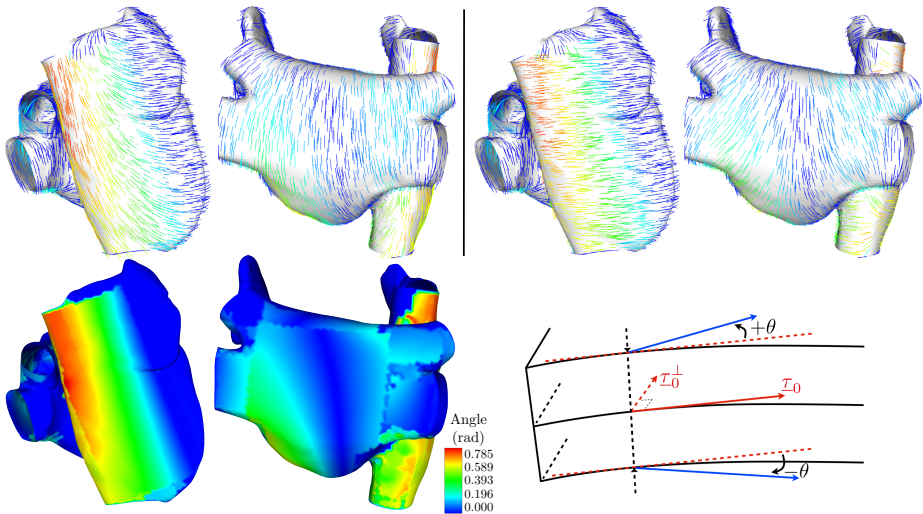


Fig. 2. Fibers directions at the endocardium (top-left) and epicardium (top-right), angle θ (bottom-left), and angular variation visualization (bottom-right)

2.3 Surface-Based Bidomain Model

We denote by \mathcal{S} the midsurface of the wall. The variational bidomain surface model that we propose can be written in terms of the extracellular potential u_e , the transmembrane potential $V_m = u_i - u_e$, with u_i the intracellular potential, as follows [1]. Find (V_m, u_e) with $\int_{\mathcal{S}} u_e dS = 0$, such that for all $t > 0$,

$$\left\{ \begin{aligned} & A_m \int_{\mathcal{S}} \left(C_m \frac{\partial V_m}{\partial t} + I_{ion}(V_m) \right) \phi dS + \int_{\mathcal{S}} \left(\underline{\sigma}_i \cdot (\nabla V_m + \nabla u_e) \right) \cdot \nabla \phi dS \\ & \hspace{15em} = A_m \int_{\mathcal{S}} I_{app} \phi dS, \quad \forall \phi \quad (1) \\ & \int_{\mathcal{S}} \left((\underline{\sigma}_i + \underline{\sigma}_e) \cdot \nabla u_e \right) \cdot \nabla \psi dS + \int_{\mathcal{S}} \left(\underline{\sigma}_i \cdot \nabla V_m \right) \cdot \nabla \psi dS = 0, \quad \forall \psi \end{aligned} \right.$$

also with $\int_S \psi dS = 0$, and where A_m is a positive constant denoting the ratio of membrane area per unit volume, C_m the membrane capacitance per unit surface, $I_{ion}(V_m)$ a reaction term representing the ionic current across the membrane and also depending on local ionic variables satisfying additional ODEs, and I_{app} a given applied stimulus current.

We define the intra- and extra-cellular diffusion tensors $\underline{\underline{\sigma}}_i$ and $\underline{\underline{\sigma}}_e$ by

$$\underline{\underline{\sigma}}_{i,e} = \sigma_{i,e}^t \underline{\underline{I}} + (\sigma_{i,e}^l - \sigma_{i,e}^t) [I_0(\theta) \underline{\underline{\tau}}_0 \otimes \underline{\underline{\tau}}_0 + J_0(\theta) \underline{\underline{\tau}}_0^\perp \otimes \underline{\underline{\tau}}_0^\perp], \tag{2}$$

where $\underline{\underline{I}}$ denotes the identity tensor in the tangential plane – also sometimes called the surface metric tensor – $\underline{\underline{\tau}}_0$ is a unit vector parallel to the local fiber direction on the atria midsurface, and $\underline{\underline{\tau}}_0^\perp$ such that $(\underline{\underline{\tau}}_0, \underline{\underline{\tau}}_0^\perp)$ gives an orthonormal basis of the tangential plane, see Figure 2-bottom-right. The effect of angular variations enters the model through the coefficients $I_0(\theta) = \frac{1}{2} + \frac{1}{4\theta} \sin(2\theta)$ and $J_0(\theta) = 1 - I_0(\theta)$. Note that $J_0(\theta) = 0$ (and $I_0(\theta) = 1$) if and only if $\theta = 0$ – namely, constant direction across the thickness – and then $\underline{\underline{\sigma}}_{i,e} = \sigma_{i,e}^t \underline{\underline{I}} + (\sigma_{i,e}^l - \sigma_{i,e}^t) \underline{\underline{\tau}}_0 \otimes \underline{\underline{\tau}}_0$ as expected for a single fiber direction. By contrast, angular variations make I_0 decrease and J_0 increase in (2), which renders diffusion “more isotropic”. This model derived from a detailed asymptotic analysis thus allows to take into account the rapid variations of the fiber direction.

The current I_{ion} can be described by a physiological or a phenomenological model. In this study, the physiological model proposed by Courtemanche, Ramirez and Nattel in [3] – most widely accepted for atria modeling – is considered.

The values of the membrane parameters are $A_m = 200.0 \text{ cm}^{-1}$ and $C_m = 10^{-3} \text{ mF.cm}^{-2}$. The values of the conductivity parameters vary substantially depending on the specific areas considered and are given in Table 2 (all in S.cm^{-1}). The Bachmann bundle (BB), the *Crista Terminalis* (CT) and the pectinate muscles (PM) are regions of established fast conduction. By contrast, the *Fossa Ovalis* (FO) and the *Isthmus* of the right atrial floor (IRA) are regions of known slow conduction. All these specific regions are depicted in Figure 1.

The two atria are connected via two regions only – in the mesh connectivity, hence also in terms of electrical conduction – namely, the Bachmann bundle and the *Fossa Ovalis*.

Table 2. Conductivity parameters

	regular tissue	PM	CT	BB	IRA	FO	SN
σ_i^t	$2.2 \cdot 10^{-4}$	$3.3 \cdot 10^{-4}$	$6.6 \cdot 10^{-4}$	$1.045 \cdot 10^{-3}$	$8.250 \cdot 10^{-5}$	$2.2 \cdot 10^{-4}$	$2.2 \cdot 10^{-4}$
σ_i^l	$2.2 \cdot 10^{-3}$	$3.894 \cdot 10^{-3}$	$9.570 \cdot 10^{-3}$	$1.639 \cdot 10^{-2}$	$9.735 \cdot 10^{-4}$	$2.002 \cdot 10^{-3}$	$2.002 \cdot 10^{-3}$
σ_e^t	$9.0 \cdot 10^{-4}$	$1.350 \cdot 10^{-3}$	$2.7 \cdot 10^{-3}$	$4.275 \cdot 10^{-3}$	$3.375 \cdot 10^{-4}$	$9.0 \cdot 10^{-4}$	$9.0 \cdot 10^{-4}$
σ_e^l	$2.2 \cdot 10^{-3}$	$3.690 \cdot 10^{-3}$	$8.550 \cdot 10^{-3}$	$1.435 \cdot 10^{-2}$	$9.225 \cdot 10^{-4}$	$2.070 \cdot 10^{-3}$	$2.070 \cdot 10^{-3}$

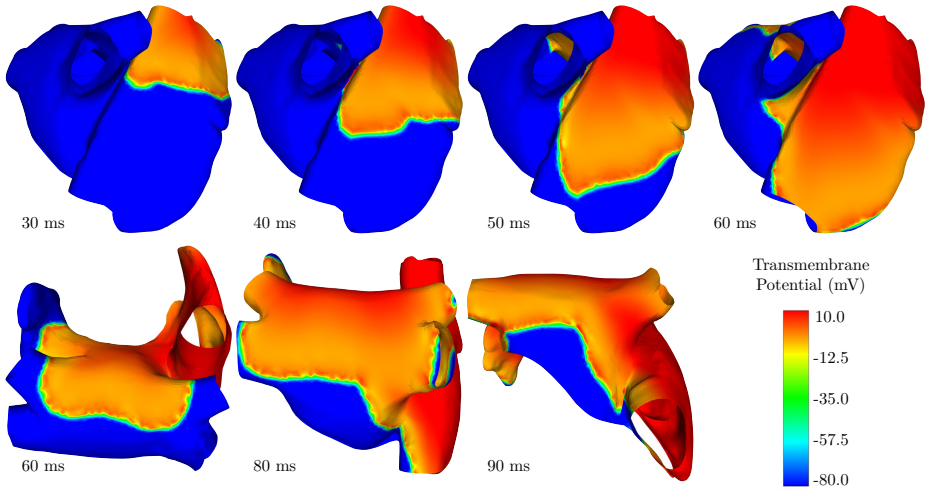


Fig. 3. Simulation of atrial depolarization

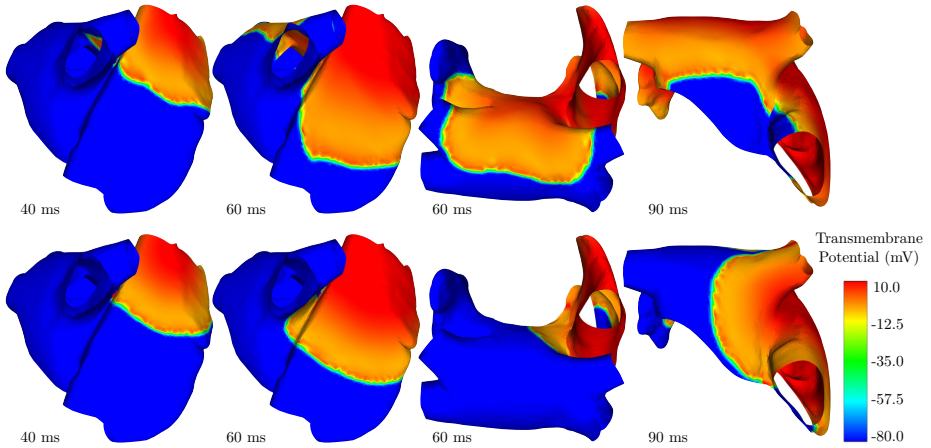


Fig. 4. Simulations of atrial depolarization for homogeneous case i.e $\theta = 0$ (top) and for homogeneous and isotropic case (bottom)

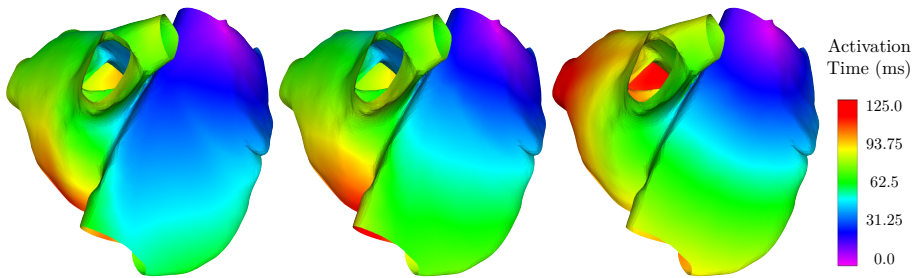


Fig. 5. Activation time maps for the three cases: anisotropic-heterogeneous (left), anisotropic-homogeneous (middle) and isotropic-homogeneous (right)

3 Simulations Results

Figure 3 displays the simulation results obtained with the above-described model. Activation is initiated at the sinus node with a stimulus of 2 ms in duration and sufficient strength to cause the initiation of a propagating wavefront. We compare our results to several 3D modeling studies [4, 6, 9].

As observed in the figure, by 30 ms the wave quickly spreads along the *Crista Terminalis* as a consequence of the high conductivity in this part. Importantly, the depolarizing wave has now traversed the Bachmann bundle and the first activation of the left atrium has occurred at 29.6 ms. This compares very well with the findings of the modeling study [6] giving the left atrium activation at 29.7 ms. At $t = 40$ ms, we clearly see the effect of the anisotropy of the *Crista Terminalis*, and of the pectinate muscles already observed in [9], as the wavefront becomes nearly triangular – as also seen in [6]. The wavefront has encircled the superior *Vena Cava*. The *Fossa Ovalis*, which is the second atrial connection, becomes active (42.8 ms). By 50 ms, because of the rapid conduction in the Bachmann bundle, the wave spreads to the left atrial appendage and has activated a substantial part of the left atrial wall. The *Fossa Ovalis* is very active. The activation of the right atrial appendage is complete. At $t = 60$ ms, in the right atrium the floor is the last part that remains unexcited. The wave has encircled the mouth of the left atrial appendage at 72.8 ms. By $t = 80$ ms, the only part that remains unaffected by the wave in the right atrium is the isthmus, because of the reduced conductivity there. In the left atrium, at $t = 90$ ms we can see three separate wavefronts, also obtained in [4, 6]. The depolarization of the left atrial appendage is complete at 95.8 ms. At 100 ms, only a small part of the left atrium in the shape of a parallelepiped is still inactive. The depolarization of the right and left atrium are complete at 101.8 ms and 109.6 ms, respectively. This is in good adequacy with the timings found in [6], namely, 99.3 ms and 108.2 ms, respectively, while [9] gives 115.0 ms for the left atrium, and [4] 119.0 ms.

4 Discussion

For comparison purposes, we also display in Figure 4-top simulation results obtained when disregarding the fiber direction variations across the thickness, namely, when taking $\theta = 0$. We observe major differences with the above-discussed physiological simulation. In particular, the anisotropy effects characteristic of the *Crista Terminalis* are much less clearly seen, and the triangular shape is hardly observed, indeed. Overall, the activation timings are longer, in particular due to slower propagation in the *Crista Terminalis*, with complete depolarization occurring as late as 116.2 ms and 115.4 ms for the right and left atria, respectively. The last results displayed in Figure 4-bottom correspond to an homogeneous and isotropic tissue throughout. We increased the homogeneous conductivity with $\sigma_i = 8.0 \cdot 10^{-4} \text{ S.cm}^{-1}$ and $\sigma_e = 3.1 \cdot 10^{-3} \text{ S.cm}^{-1}$ to obtain roughly comparable depolarization timings. The resulting activation profile, however, is totally unrealistic – as already noted in [9] – see Figure 5 displaying the activation time maps for the three cases analyzed.

5 Concluding Remarks

We have presented detailed simulation results of a surface-based electrophysiology bidomain model applied with an anatomical model of the atria. The results obtained are in excellent adequacy with previous studies, and confirm the importance of fiber-related anisotropy effects and of their strong variations across the wall thickness. Furthermore, this surface-based model provides dramatic computational benefits over 3D models with preserved accuracy [1]. Further perspectives include applications to radio-frequency ablation planning assistance, for which computational effectiveness – less than 3 minutes for a complete simulation on a standard workstation – will be of great value.

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