

# Chapter 3

## Hafnium-Based Gate Dielectric Materials

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**Abstract** In this chapter, we focus on hafnium-based gate dielectrics.  $\text{HfO}_2$  is regarded as the most promising material for the high- $k$  gate dielectrics owing to its large dielectric constant and large band-gap energy. In the first part of this chapter, these characteristics are addressed in a comparison with  $\text{SiO}_2$  and other high- $k$  materials. Thermal stability is a major issue for the application of high- $k$  materials to MOSFETs in LSIs. Although  $\text{HfO}_2$  satisfies this requirement, severer process conditions may cause problems even with this material. Suppression of these issues is also addressed in the second part of this chapter. In order to enhance the characteristics of  $\text{HfO}_2$ , transformation of the monoclinic phase to the tetragonal and the cubic phases with larger dielectric constant has been pursued recently. We mention the issue in the last part of this chapter.

### 3.1 Introductory Remarks and Brief Outline of the Chapter

Since the proposal of the materials in the early 2000s [1–3], hafnium-based gate dielectrics have been regarded as the most promising materials for the application to large scale integrations (LSIs). In 2007, Intel announced the start of manufacturing of 45 nm-node LSIs with high- $k$ /metal gate stack, with the remark that the world's first high- $k$  material in commercial production was hafnium-based [4].

There are several characteristics that the high- $k$  materials should meet for the application to metal oxide semiconductor field effect transistors (MOSFETs) in LSIs. First, we clarify the properties so that readers can understand their

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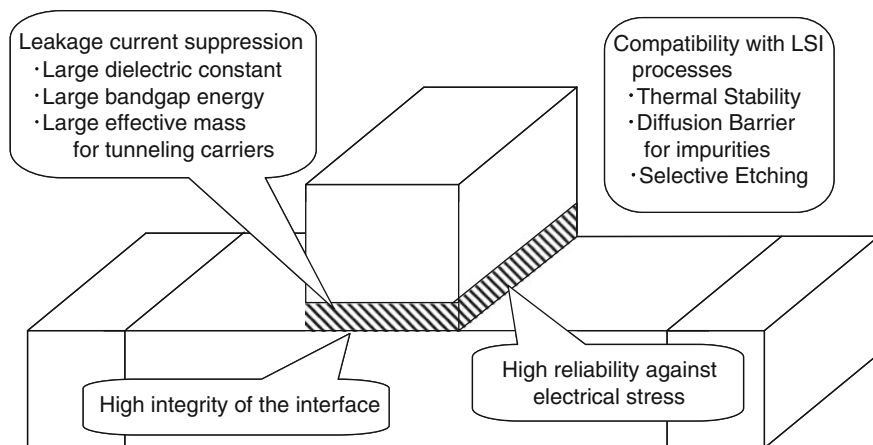
importance in the selection process of high-k materials. Next, we present concrete data on  $\text{HfO}_2$ , and identify the reasons why this material is considered the most promising candidate. This is followed by data on its alloy with  $\text{SiO}_2$  or  $\text{Al}_2\text{O}_3$ . Along with pure  $\text{HfO}_2$ , these two alloys have been investigated intensively by many companies, institutes and universities world-wide in order to enhance thermal stability and compatibility with the LSI processes. Since  $\text{HfO}_2$  takes several phases and they influence the high-k properties, this topic is also addressed in this chapter. Taking the higher-k  $\text{HfO}_2$  phase is regarded as one of the measures to extend the usability of this material through several LSI generations.

## 3.2 Properties Required for Gate Dielectrics

Listed in Fig. 3.1 are the properties required for high-k materials as alternative gate dielectrics of advanced ultra large scale integrations (ULSIs). They are large dielectric constant, large band-gap energy, large effective mass for tunneling carriers, good integrity of high-k/substrate interface, high reliability against electrical stresses, and compatibility with ULSI processes. Obviously, good uniformity of properties throughout the wafer is essential.

The probabilities of the quantum mechanical tunneling of electrons and holes are governed by the integral of the barrier height that carriers encounter in their attempt to tunnel through the insulators [5]. The simplified equation of the direct tunneling  $J_g$  through a gate dielectric is presented in Ref. [6] as follows:

$$J_g \propto \exp\left(-\frac{4\pi\sqrt{2qm^*\phi_b}}{h} \frac{k}{3.9} t_{EOT}\right), \quad (3.1)$$



**Fig. 3.1** Properties required for high-k materials applied as alternative gate dielectrics of advanced MOS transistors

where  $t_{EOT}$  and  $\phi_b$ ,  $k$  and 3.9 correspond to the electrical thickness of the gate dielectric, the energy barrier height formed at the gate dielectric/substrate interface, dielectric constant of the gate insulator and that of  $\text{SiO}_2$ , respectively.  $m^*$ ,  $q$ , and  $h$  correspond to the effective tunneling mass of the carrier, elementary charge and the Planck constant. The electrical thickness is called the ‘equivalent oxide thickness (EOT)’, which is the  $\text{SiO}_2$  thickness equivalent to the high- $k$  gate dielectric inducing the same amount of carrier density in the channel region with the same gate voltage.

According to the International Technology Roadmap for Semiconductors (ITRS), the EOT decreases as the device scaling proceeds in order to maintain a high sheet carrier concentration with lower voltages [7]. However, this leads to a drastic increase in power consumption due to the severe gate leakage current from the smaller  $t_{EOT}$  in (3.1) [7]. Therefore, material with a larger dielectric constant is used. The physical thickness of high- $k$  gate dielectric  $t_{gate}$  with the dielectric constant  $k$  is calculated using the following equation:

$$t_{gate} = \frac{k}{3.9} t_{EOT} \quad (3.2)$$

3.9 is the dielectric constant of  $\text{SiO}_2$ . Equation (3.2) indicates the thicker physical thickness is obtained with larger dielectric constant and this leads to a decrease in  $J_g$ .

In addition to the physical thickness, the energy barrier formed between the high- $k$  material and the substrate should be large. Since silicon has a conduction band edge and a valence band edge at about 4.0 and 5.1 eV from the vacuum level, the conduction band edge of high- $k$  material should be much nearer the vacuum level, whereas the valence band edge should be further from the vacuum level, in order to diminish the tunneling probability of electron transport from the conduction band of silicon in n-channel MOSFETs (NMOS) and hole transport from the valence band of silicon in p-channel MOSFETs (PMOS). A schematic view of the band alignment concept for NMOS is shown in Fig. 3.2.

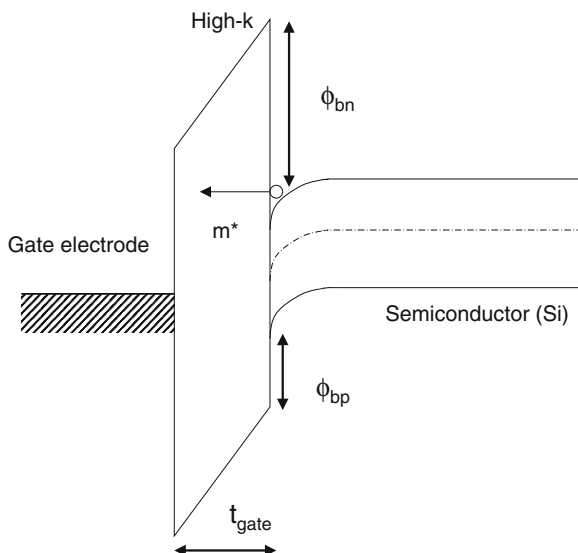
The tunneling probability is also governed by the effective mass of the carriers: a heavier mass leads to lower probability. Therefore, the high- $k$  materials through which carriers tunnel with heavier mass are desirable for the gate insulator. The effective mass of electrons in  $\text{HfO}_2$  was extracted with the fit between the experimental tunneling leakage current data and theoretical calculations. Those studies have provided figures such as  $0.22 m_0$  [6],  $0.15 m_0$  [8], or  $0.1 m_0$  [9], where  $m_0$  represents the free electron mass. It should be noted that the tunneling mass of electrons through  $\text{SiO}_2$  is as large as  $0.4 m_0$  [6].

Using the function inside the brackets in (3.1), we can define the figure of merit (F.O.M.) of the high- $k$  dielectrics as follows [6]:

$$F.O.M. \equiv k \sqrt{m^* \phi_b} \quad (3.3)$$

A large value of F.O.M. of a high- $k$  material indicates its high potential to diminish the amount of  $J_g$  with the fixed  $t_{EOT}$  value.

**Fig. 3.2** Schematic view of the concept of the band alignment among gate electrode, high-k material, and semiconductor substrate.  $\phi_{bn}$ ,  $\phi_{bp}$  and  $m^*$  represent energy barrier for electrons, energy barrier for holes and effective mass for tunneling carriers, respectively



Since the integrity of high-k/semiconductor substrate interface and film reliability issues against electrical stresses are discussed in other chapters, the thermal stability issue is addressed in this chapter. The LSI processes consist of several types of high-temperature annealing such as oxidation annealing, nitrogen annealing, and chemical vapor deposition. When the high-k dielectrics are applied in the conventional LSI processes, a few high-temperature processes including the source/drain impurity activation annealing are performed after the deposition process. Currently, several methods such as spike rapid thermal annealing, flash lamp annealing, and laser thermal annealing are proposed in order to limit the thermal budget of the impurity activation process. The concept of these processes is that thermal energy is provided to incorporate impurity atoms in the substrate crystalline network, limiting the diffusion of the atoms in the device structures. Since the typical temperature does not go below 900 °C, high-k materials on semiconductor substrate should maintain their above-mentioned properties through the high-temperature annealing. It is particularly important to avoid a thermal reaction between high-k materials and the substrate, since such a reaction ruins the integrity of the high-k properties, in many cases.

### 3.3 Physical and Electrical Properties of Hafnium Oxide

Hafnium is a transition metal with a steel grey color. The atomic number of this element is 72. With its very high neutron-capture cross-section, it is regarded as a good material for the control rods of nuclear reactors [3]. Like other elements in group IV (titanium and zirconium), it is a very reactive material in air and forms

several phases of oxides. Among several oxide phases of hafnium,  $\text{HfO}_2$ , called ‘hafnia’, is the only stable compound showing the insulating behavior.

### 3.3.1 Dielectric Constant and Microscopic Polarization

In LSI processes,  $\text{HfO}_2$  is deposited either by physical vapor deposition (PVD) such as sputtering [2], chemical vapor deposition (CVD) [1], oxidation of Hf metals [10], or by atomic layer deposition (ALD) [11–14]. In many cases,  $\text{HfO}_2$  forms monoclinic phase just after the deposition process. Even in the case in which it has amorphous phase after deposition, low-temperature annealing transforms the material to polycrystalline monoclinic  $\text{HfO}_2$  [10].

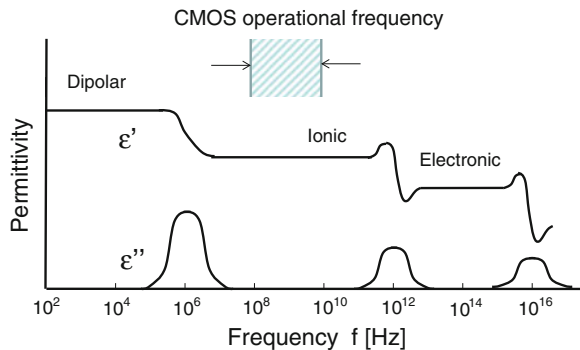
In the monoclinic crystal structure, the coordination number of Hf is 7, whereas that of oxygen is 3 or 4 [15]. The static dielectric constant of monoclinic  $\text{HfO}_2$  ranges from 17 [1, 2] to 22 [16]. Throughout this chapter, ‘dielectric constant’ means relative permittivity  $k$ . The absolute permittivity  $\epsilon$  is the vacuum permittivity  $\epsilon_0$  multiplied by this number. The dielectric constant of materials comes from a few components as shown in Fig. 3.3.

For high- $k$  gate dielectrics, the ionic and the electronic polarization components are important, because they follow a high clock frequency (about GHz order) used in advanced ULSIs. The electronic polarization component of the dielectric constant can be extracted using the following equation:

$$k_e = n^2, \quad (3.4)$$

where  $n$  represents reflectivity. This component is called the ‘the dynamic dielectric constant’. This is also called the ‘optical dielectric constant’ because it is the dielectric constant in the optical frequency region. Knowing that the reflectivity of  $\text{HfO}_2$  is about 2.1 [17], the electronic polarization component of this material is estimated to be around 4.4. Since this value is small compared to the static dielectric constant value [1, 2, 14], one can assume that the reason for the high static dielectric constant of this material resides in a large ionic polarization.

**Fig. 3.3** The frequency dependence of permittivity.  $\epsilon'$  represents the real part, while  $\epsilon''$  represents the imaginary part of the permittivity. Since advanced ULSI operation clock frequency is in the GHz domain, ionic and electronic polarization contributes to the permittivity of high- $k$  materials at the frequencies



The Clausius–Mossotti theory correlates the static dielectric constant  $k$  and the structural parameters as follows:

$$\frac{k - 1}{k + 2} = \frac{4\pi\alpha}{3V_m}, \quad (3.5)$$

where  $V_m$  and  $\alpha$  are the molar volume in  $\text{\AA}^3$  and the polarizability of each molecule, respectively. This equation can be rewritten as follows:

$$k = \frac{1 + \frac{8\pi\alpha}{3V_m}}{1 - \frac{4\pi\alpha}{3V_m}} \quad (3.6)$$

This clearly states that the increase in  $\alpha$  and the decrease in  $V_m$  lead to the increase in the dielectric constant of the material. In the case of  $\text{HfO}_2$ , the polarizability of each molecule and molar volume is about 7.3 and  $35 \text{\AA}^3$ , and its combination results in large dielectric constant among many metal oxides [18]. The polarizability of each molecule and molar volume of  $\text{SiO}_2$  is 4.8 and  $38 \text{\AA}^3$ , for comparison [18]. Although there is little difference between the molecular volumes of the two materials, the polarizability is much larger for  $\text{HfO}_2$  than for  $\text{SiO}_2$ . This large polarization can be ascribed to the large difference between the electro-negativity of Hf and O: Hf electronegativity is 1.3 (Pauling), whereas that of oxygen is 3.44 [19]. It should be noted that the electronegativity of Si is as large as 1.9, which is larger than that of Hf and nearer to the value of oxygen (Table 3.1).

### 3.3.2 Bandgap and Band Alignment with Silicon

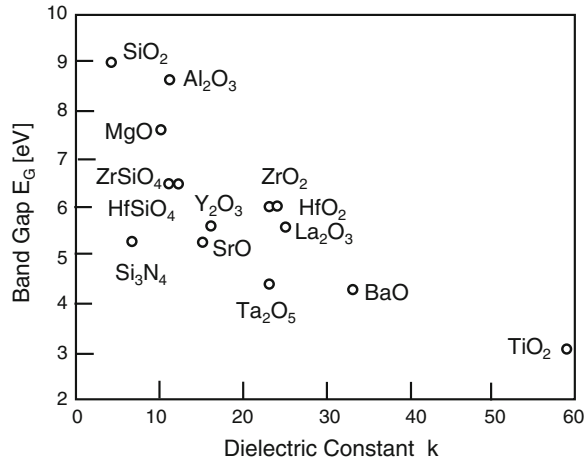
As mentioned in Sect. 3.2, band-gap energy of a high- $k$  material is the basic parameter influencing the F.O.M. of the material. Figure 3.4 shows the relationship of band-gap energy and dielectric constant of metal oxides.

The general trend of these two parameters indicates that as the dielectric constant of material increases, the bandgap energy decreases. Considering that the transition probability of electrons from the stable states to the excited states increases owing to the decrease in the bandgap, it can be understood that the

**Table 3.1** Pauling's electronegativity for typical cation and anion elements [19]

<i>Cation</i>							
Al	Si	Ti	Y	Zr	La	Hf	Ta
1.6	1.9	1.5	1.2	1.3	1.1	1.3	1.5
<i>Anion</i>							
N				O			
3				3.4			

**Fig. 3.4** Correlation of static dielectric constant with band gap of metal oxides. (Reprinted with permission from [26]. Copyright 2002, American Institute of Physics)



dynamic dielectric constants increase for high-k materials with narrower bandgap energies [20].

Bandgap energy of materials can be measured with a few analytical techniques such as XPS (X-ray photoelectron spectroscopy) [21], ellipsometry [22], and REELS (reflection electron energy loss spectroscopy) [23]. The measured bandgap of the monoclinic HfO<sub>2</sub> is 5.7 eV [22] or 5.8 eV [23, 24]. Although the bandgap is smaller than that of SiO<sub>2</sub> (about 9 eV) [23, 25], the value is still large among high-k materials. Theoretical calculation of this material elucidates that its conduction band originates from 5d band of Hf, whereas the valence band originates from 2p states of oxygen [26].

XPS measurement is also a very useful technique for the extraction of the band alignment of high-k materials to the underlying substrate. It is well known that work functions of pure metals are not realized when they contact semiconductors and form metal/semiconductor systems. This is certainly the case for a semiconductor with a relatively small bandgap such as Si (about 1.1 eV) [27]. For example, aluminum, which has been used as the interconnection metal in LSIs for a long time, has bulk work function of about 4.0 eV; however, its effective work function in Al/Si system is about 4.7 eV [27]. This discrepancy is explained using the CNL (Charge Neutrality Level) of semiconductor and the formation of dipole at the metal/semiconductor interface [28]. A similar theory is applicable to the insulator/semiconductor systems using the CNL of insulator as an alternative to Fermi level of metal in metal/semiconductor systems [26].

Normally, valence band spectra from XPS measurements are used to extract the band alignment of the high-k/semiconductor systems [29]. This measurement reveals binding energies of the tops of the valence bands for high-k materials and the underlying semiconductor substrate. The important point is that those two binding energy values come from one sample. By calculating the energy difference of the two, one can extract valence band offset between them. In addition, if the band-gap energies of two materials are known by means of the other methods

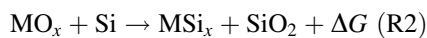
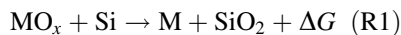
mentioned above, the conduction band offset can be readily obtained. See Fig. 3.2 for the reference of this procedure. Another common method of extracting the band offset is IPES (internal photo-emission spectroscopy) measurement: When electrons are provided with sufficient energy to overcome the band offset, they are detected as the current [30]. The extracted conduction band offset of HfO<sub>2</sub> was 1.2 eV [31] or 1.9 eV [32]. This relatively large value along with the large dielectric constant promises a dramatic decrease in leakage current by 4 or 5 orders of magnitude through HfO<sub>2</sub>, compared to that through SiO<sub>2</sub> with the same EOT [1, 2].

### 3.3.3 Compatibility with LSI Processes

Another important factor considered in high-k material selection is the thermal stability in the ULSI fabrication processes. ‘Thermal stability’ has several aspects as follows:

1. The melting point of the material must be higher than the temperature in ULSI fabrication process
2. The metal oxide should not be reduced by the neighboring materials including the substrate semiconductor during ULSI fabrication processes
3. Reaction with neighboring materials other than reduction should not occur in ULSI fabrication process.

Hubbard et al. [33] used these criteria to select possible high-k insulators among many metal oxides. Regarding the second and the third criteria, they supposed the following reactions and tried to find metal oxides (MO<sub>x</sub>) whose reaction Gibbs free energies,  $\Delta G$ , become positive. With positive  $\Delta G$  value, these reactions become endothermic, indicating that the metal oxide is stable in contact with the substrate silicon. Through this procedure, they selected several metal oxide candidates, including HfO<sub>2</sub>.



With a melting point of 2,758 °C and a boiling point of 5,400 °C, HfO<sub>2</sub> is expected to be a very stable material. Moreover, Hubbard et al. calculated the Gibbs free energy of R1 at 1,000 K and concluded that R1 reactions are not expected for HfO<sub>2</sub> at least at this temperature. They could not calculate Gibbs free energy of R2 reactions for HfO<sub>2</sub>, because of the lack of thermodynamic data for HfSi<sub>2</sub>. However, given that R1 and R2 reactions are both endothermic for ZrO<sub>2</sub> at 1,000 K, one can expect the stability of HfO<sub>2</sub> against R2 reactions, considering the similarity of the two materials. This high thermal stability along with the drastic



reduction of the gate leakage current makes  $\text{HfO}_2$ , as well as  $\text{ZrO}_2$ , a promising material for the high-k gate dielectrics in advanced ULSIs.

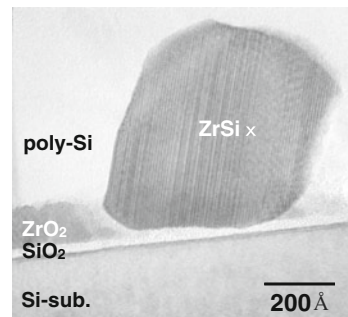
However, in ULSI processes, severer process condition may cause problems even for these materials: zirconium silicide formation was reported when high-temperature annealing (around 950–1,000 °C) was performed under high vacuum or in an oxygen-deficient ambient [34–36] as shown in Fig. 3.5 [37]. Although silicide formation temperature is higher for  $\text{HfO}_2$  than for  $\text{ZrO}_2$  [38], a similar reaction was detected [39, 40]. A different reaction model that includes reduction of the interfacial  $\text{SiO}_2$  layer was proposed as follows (R3) [41] and this may be the probable reason for the increase of the leakage current of  $\text{HfO}_2/\text{SiO}_2$  gate stack after 950 °C annealing in  $\text{N}_2$  ambient [42]:



Another issue concerning the compatibility of  $\text{HfO}_2$  with ULSI process is the severe diffusion of oxygen through the film [43, 44]. High-temperature annealing in oxygen-containing atmosphere is commonly practiced to improve the leakage characteristics of the gate stack. Moreover, the residual atmosphere in the production furnaces contains oxygen. The oxygen diffusion during high-temperature annealing in oxygen-containing ambient leads to EOT increase due to the interfacial layer growth and ruins the leakage suppression characteristics. The experimental result shows that the activation energy of this phenomenon is low, 0.6 eV [45]. This fast interfacial layer growth is attributed to the fast oxygen diffusion in  $\text{HfO}_2$  and the enhanced oxidation at the interface owing to the catalytic function of the  $\text{HfO}_2$  layer [40, 46, 47].

The diffusion of common impurities such as arsenic and boron would pose a problem when poly-Si gate electrode is used with the high-k material [48]. When poly-silicon is used as the gate electrode, it must be highly doped with phosphorous and arsenic for NMOSFETs and with boron for PMOSFETs to reduce the sheet resistance of the electrode and adjust the threshold voltage. Thin  $\text{HfO}_2$  cannot work as the diffusion barrier for arsenic [49] or boron [50, 51] and this diffusion leads to unintentional decrease in the threshold voltage or even to severe leakage current between source and drain (punch-through phenomenon).

**Fig. 3.5**  $\text{ZrO}_2/\text{SiO}_2$  gate stack reacts with Si substrate when annealed in a  $\text{N}_2$  ambient at 1,000 °C, resulting in the  $\text{ZrSi}_x$  formation [37]



Therefore, suppression of the above-mentioned issues is the incentive for incorporation of other elements such as nitrogen, aluminum and silicon into  $\text{HfO}_2$ . The change in the material's properties as a result of their incorporation is also addressed in the following sections.

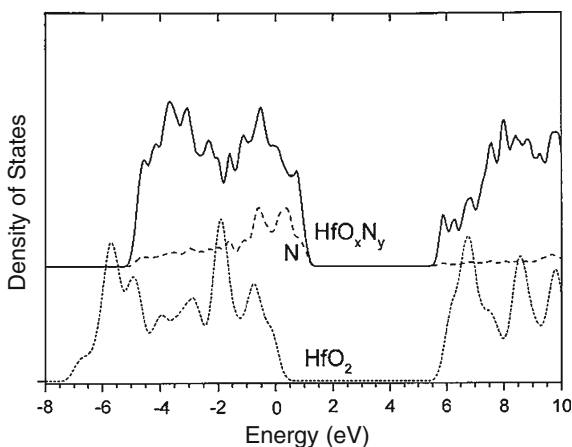
### 3.4 Hafnium–Nitrogen-Based Gate Dielectrics

Nitrogen has been introduced into  $\text{HfO}_2$  in order to enhance the thermal stability of the material.  $\text{HfON}$  is deposited by  $\text{N}_2/\text{O}_2$  reactive sputtering using Hf target [52].  $\text{HfON}$  can also be formed by  $\text{N}_2$  reactive sputtering from Hf target followed by rapid thermal oxidation [42]. CVD deposition of this material has also been reported [53]. Kang et al. [42] claimed that nitrogen incorporation at the gate dielectric/Si interface region and  $\text{SiN}$  formation may lead to the suppression of  $\text{SiO}$  formation in the reaction R3 mentioned above, resulting in the suppression of the leakage current increase. The suppression of diffusion of oxygen [42, 43] as well as impurities such as boron [48, 53, 54] through the film has also been reported to be a beneficial influence of the nitrogen incorporation.

As in the case of  $\text{SiON}$  [55], nitrogen incorporation reduces the bandgap of  $\text{HfO}_2$ . This is ascribed to the replacement of the valence band edge from  $\text{O}2p$  orbital to  $\text{N}2p$  orbital [56]. Although it leads to the decrease of its valence-band barrier offset to Si substrate by about 1.1 eV, a sufficiently large energy barrier (about 2.2 eV) is retained to suppress the hole injection from the substrate in PMOSFETs [56] (Fig. 3.6).

Pure nitrides of fourth-column transition metals (Ti, Zr and Hf) share similar characteristics. All form  $\text{MN}$ -type nitrides with NaCl structure, and with more nitrogen in the structures, they form  $\text{M}_3\text{N}_4$ -type nitrides. Hafnium pure nitrides

**Fig. 3.6** Calculated densities of states for  $\text{HfO}_2$  and  $\text{HfON}$ . Band-gap modification due to the nitrogen incorporation into  $\text{HfO}_2$  is shown. (Reprinted with permission from [56]. Copyright 2004, American Institute of Physics)

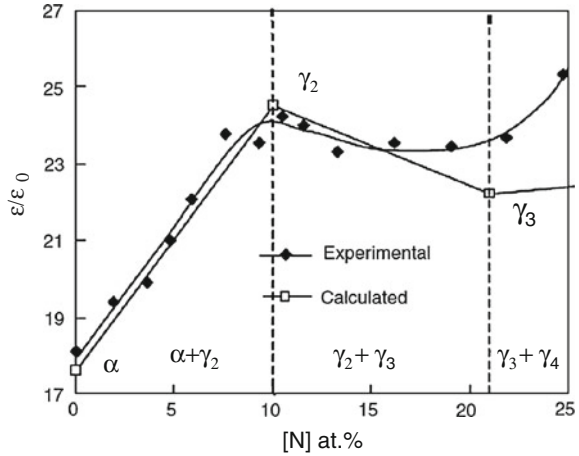


such as HfN [57, 58] and Hf<sub>3</sub>N<sub>4</sub> [59, 60] have been reported. Among them, Hf<sub>3</sub>N<sub>4</sub> shows insulating behavior. It is reported that Hf<sub>3</sub>N<sub>4</sub> may solidify in different crystalline structures according to temperature and pressure around the material [61, 62]. First-principle calculation predicted the band-gap of Hf<sub>3</sub>N<sub>4</sub> to be as large as 1.8 eV [62] and electrical measurement of CVD Hf<sub>3</sub>N<sub>4</sub> revealed the dielectric constant to be about 30 [60]. Since HfON can be regarded as an alloy of HfO<sub>2</sub> and Hf<sub>3</sub>N<sub>4</sub> [63], intermediate characteristics are anticipated for HfON. Having thoroughly investigated the physical properties of HfON with various nitrogen content, Ino et al. concluded that the above-mentioned Clausius–Mossotti equation (3.6) describes the dielectric constants of the alloys well [64]. They calculated the static polarizabilities  $\alpha_{total}$  using the additive rule of the static polarizability of each constituent element [18] as follows:

$$\alpha_{total} = C_{Hf}\alpha_{Hf^{4+}} + C_O\alpha_{O^{2-}} + C_N\alpha_{N^{3-}}, \tag{3.7}$$

where  $C_{Hf}$ ,  $C_O$ , and  $C_N$  represent content of hafnium, oxygen, and nitrogen in the alloy, respectively. Given the polarizability of the oxygen ion [18], they estimated the polarizability of hafnium ion by using the material data of HfO<sub>2</sub> [65]. They also estimated that of the nitrogen ion by using data for some metal nitrides [64]. Figure 3.7 shows dielectric constants calculated using (3.7) and estimated molar volumes for  $\alpha$  (m-HfO<sub>2</sub>),  $\gamma_2$  (Hf<sub>7</sub>O<sub>11</sub>N<sub>2</sub>),  $\gamma_3$  (Hf<sub>7</sub>O<sub>8</sub>N<sub>4</sub>), and  $\gamma_4$  (Hf<sub>2</sub>ON<sub>2</sub>). The calculated values fit experimental data very well, especially for N content of less than 18 %. Ino et al. also extracted the band alignment of the alloys using REELS and valence spectra of XPS and concluded that this material is one of the better candidates for high-k gate dielectrics.

**Fig. 3.7** Dielectric constant of HfON versus N content. Calculation results based on the Clausius-Mossotti theory are compared to the experimental values [64]. In this figure,  $\alpha$ ,  $\gamma_2$ ,  $\gamma_3$ , and  $\gamma_4$  represent m-HfO<sub>2</sub>, Hf<sub>7</sub>O<sub>11</sub>N<sub>2</sub>, Hf<sub>7</sub>O<sub>8</sub>N<sub>4</sub> and Hf<sub>2</sub>ON<sub>2</sub>, respectively



### 3.5 Hafnium–Silicon-Based Gate Dielectrics

HfO<sub>2</sub>–SiO<sub>2</sub> alloy is formed using various techniques such as chemical vapor deposition (CVD) [66, 67], plasma-enhanced CVD [25], and atomic layer deposition (ALD) [68, 69]. For silicon deposition by the CVD method, organic precursors such as SiH[N(C<sub>2</sub>H<sub>5</sub>)<sub>2</sub>]<sub>3</sub>, tetraethoxy-silane (TEOS:Si(OC<sub>2</sub>H<sub>5</sub>)<sub>4</sub>), and tetrakis-dimethylamino-silane (TDMAS:Si[N(CH<sub>3</sub>)<sub>2</sub>]<sub>4</sub>) were normally used. By changing the ratio of precursor flow, Hf content and Si content were varied in the material. Physical vapor deposition methods are also used such as pulsed-laser deposition (PLD) [70], sputtering from composite target [71], from HfO<sub>2</sub> target with SiH<sub>4</sub> gas [72], reactive sputtering from Hf and Si targets [73], and electron beam evaporation [74].

Since Hf is an element in group IV, tetravalent coordination of Hf in the silicate is expected for material with low Hf concentration, in which the SiO<sub>2</sub> tetrahedron is the main component and a few Hf cations are expected to replace some Si in the network [75]. On the other hand, as mentioned above, the coordination number of Hf is 7 and that of oxygen is 3 or 4 in the monoclinic HfO<sub>2</sub> crystal structure. Therefore, modification of the coordination number is expected when Hf content increases in the pseudo-alloy materials [76]. Some microscopic inhomogeneity originating from the coordination variability has been observed as the broadening of Si2p and Hf4f peaks in the XPS spectra from the alloy [77].

Unfavorable Hf silicide formation observed in the case of HfO<sub>2</sub> is suppressed even at high temperature by the inclusion of SiO<sub>2</sub> [74]. In addition, the oxygen diffusion in HfSiO is retarded, leading to the suppression of interfacial layer growth in an oxygen-containing atmosphere at high temperature [78]. SiO<sub>2</sub> incorporation is also effective for suppressing the impurity penetration phenomena, which are problematic when poly-Si gate electrode is used. Boron diffusion through the HfSiO film with the Hf/Hf+Si ratio of 0.52 (atomic concentration of Hf is about 10–12 %) is suppressed up to 950 °C [79]. However, considering that the process temperature for the activation of impurities in the source/drain of MOSFET is as high as 1,000 °C, the capability of Hf silicate as a diffusion barrier is insufficient. Moreover, it is reported that phosphorous and arsenic diffused through the HfSiO with the same Hf concentration during 1,000 °C rapid thermal annealing (RTA) [80]. Such impurity diffusion is ascribed to the enhanced diffusion at the grain boundaries formed in the resultant ‘phase separation phenomena’ in HfSiO at high temperatures [79, 80] as described below.

When the pseudo-alloy is annealed at high temperature, it is observed that the material loses its uniformity and HfO<sub>2</sub> precipitates in the remaining silica-rich matrix [16, 81]. Transmission electron microscope (TEM) photographs clearly show the crystallization of HfO<sub>2</sub>, which results in the inhomogeneity in the materials after 800 °C–1,000 °C high-temperature annealing [82]. Even in the alloys with relatively low content such as Hf/Hf+Si = 0.1, small HfO<sub>2</sub> crystals emerge after 1,000 °C annealing [83]. Figure 3.8 shows the planar TEM photographs of the HfSiO alloy after high-temperature annealing. Elements with larger

mass in the film have larger capability for the diffraction of incident electrons. Therefore, the dark area corresponds to those containing Hf in the film. Actually, the enlargement of the portion suggests that this area consists of orthorhombic  $\text{HfO}_2$  crystal as shown in the inset of the figure [83].

These phase separation phenomena are ascribed either to the spinodal decomposition or to the nucleation and growth [84]. Cahn treated alloys as a super-cooled liquid solution and reported that the system's total Gibbs free energy change becomes negative owing to the compositional fluctuation for a certain initial compositional region. This is where the following equation is satisfied, and the decomposition takes place spontaneously in this region [85]:

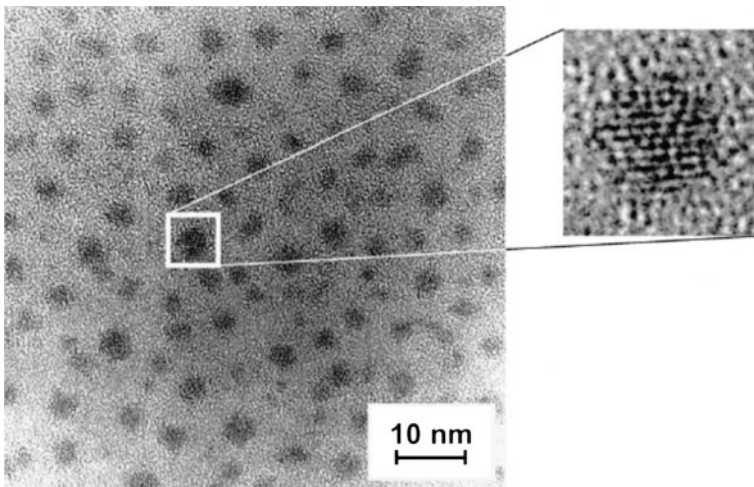
$$\frac{\partial^2 G}{\partial C^2} < 0, \quad (3.8)$$

where  $C$  represents the initial uniform composition of the alloy. This compositional region and its boundary are referred to as the spinodal region and the spinodal boundary.

Figure 3.9 shows the phase diagram of  $\text{HfO}_2$ - $\text{SiO}_2$  system, in which the spinodal boundary is indicated as the inner dotted line [16]. The alloys whose compositions are located inside the spinodal boundary are unstable and tend to suffer from phase separation.

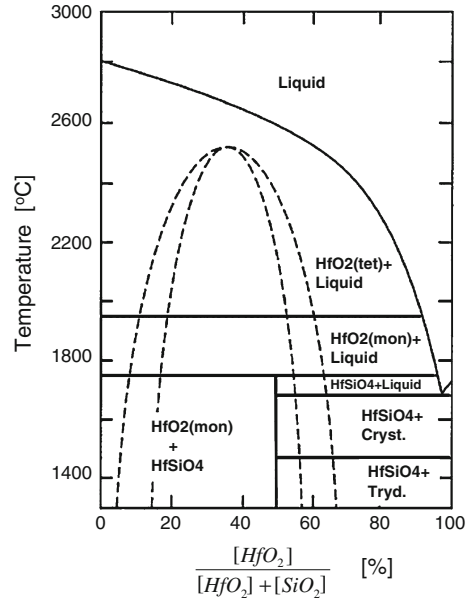
The compositions of the resultant two different phases approach those indicated by the miscibility gap curve (the outer dotted line in Fig. 3.9) when the system reaches equilibrium [16].

Even alloys whose compositions are outside the spinodal boundary may also suffer from phase separation by the nucleation and growth mechanism: once nuclei



**Fig. 3.8**  $\text{HfO}_2$  precipitation after 1,000 °C annealing of the  $\text{HfSiO}$  alloy with Hf relative content ( $\text{Hf}/(\text{Hf}+\text{Si})$ ) of 10 % [83]

**Fig. 3.9** Phase diagram of  $\text{HfO}_2\text{-SiO}_2$  compositional systems. Inner dotted line and outer dotted line represent the spinodal boundary and the miscibility gap, respectively. (Reprinted with permission from [16]. Copyright 2002, Materials Research Society)

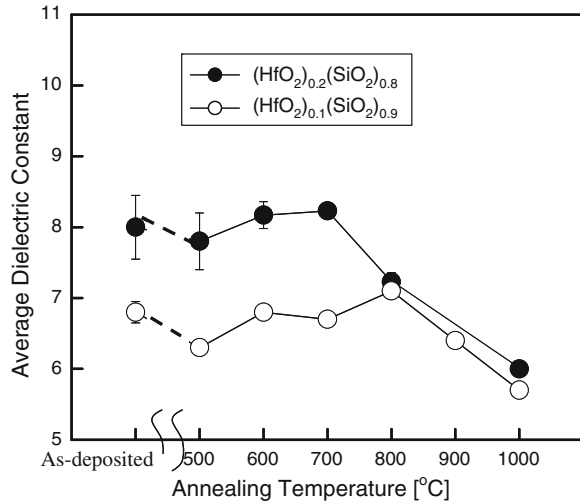


of a certain phase with a diameter larger than the critical size are formed, the growth of the nuclei is favorable for reducing the free energy of the system [86]. The phase separation observed for an alloy with low  $\text{HfO}_2$  contents ( $\text{Hf}/\text{Hf}+\text{Si} = 0.1$ ) as shown in Fig. 3.8 [83] may be caused by the nucleation and growth mechanism. This assumption is supported by the fact that the  $\text{HfO}_2$  crystals precipitated in the alloy are round, which is typical for this mechanism [86]. On the other hand,  $\text{HfO}_2$  crystal forms rather dendritic shapes in samples with larger  $\text{HfO}_2$  content when the phase separation takes place. This phenomenon is shown in Fig. 3.11a for the material with  $\text{Hf}/\text{Hf}+\text{Si} = 26\%$  and in Ref. [87] for that with  $\text{Hf}/\text{Hf}+\text{Si} = 0.4$ . The kinetics of these phenomena is considered to be limited by the diffusion of the constituent material, i.e. Hf, Si or both in the alloy system [84].

The phase separation has several adverse effects in addition to the above-mentioned enhancement of the impurity diffusion through the alloy. One of the influences is the degradation of the dielectric constant: Ino et al. measured the dielectric constant of  $\text{HfSiO}$  alloy with  $\text{HfO}_2$  content of 0.1 and 0.2 and concluded that the resultant non-uniformity in the film after the high-temperature annealing leads to a severe decrease in the average dielectric constant as shown in Fig. 3.10 [88]. Ono et al. [83] modeled the film structure due to the phase separation and calculated the change of the film's dielectric constant. They obtained a reasonable match to the experimental results.

It has also been reported that additional scattering of carriers moving through the channel occurs owing to the non-uniformity of the electric field in the gate insulator and this degrades the electron mobility of transistors with the alloy [89].

**Fig. 3.10** Average dielectric constant of the film just after deposition and after the phase separation at high-temperature annealing [88]



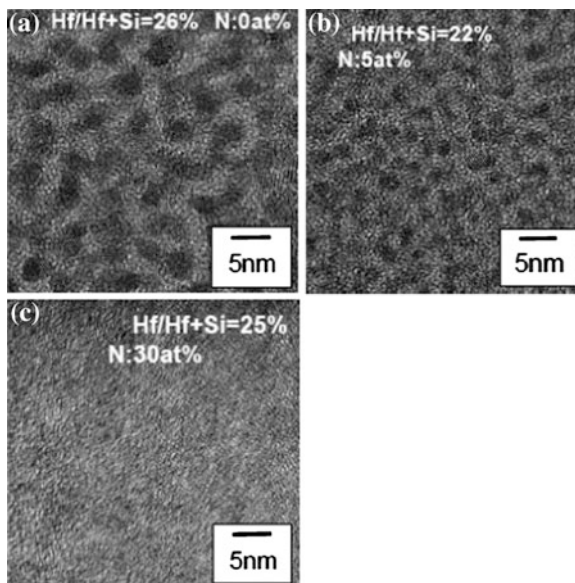
Nitrogen is incorporated in order to suppress the phase separation in HfSiO alloy [90, 91]. Nitrogen is incorporated by the introduction of nitrogen-containing gas in reactive sputtering [90]. Oxidation of HfSiN was also tried in order to obtain the HfSiON structure [92]. There is another method in which HfSiO is firstly deposited and nitrogen is introduced by subsequent NH<sub>3</sub> annealing [93, 94] or nitrogen plasma treatment [94, 95]. Sekine et al. compared the applicability of the two techniques to short-channel MOSFETs and concluded that the plasma treatment is favorable for keeping high carrier mobility in the transistor. As in the case of SiON, nitrogen at the interface between the gate dielectric and the substrate silicon induces charges that scatter carriers in the channel region. Since the plasma treatment is effective for obtaining larger concentration at the surface of HfSiO while maintaining low concentration in the vicinity of the interface, this method is effective for keeping larger carrier mobility of MOSFETs [94].

The effect of the nitrogen incorporation on the phase separation phenomenon is investigated by changing the amount of nitrogen in the material. Figures 3.11 show the TEM planar views of the Hf-silicate films (a) without nitrogen, (b) with a nitrogen atomic concentration [N] of 5 at.%, and (c) 30 at.% after the 1,000 °C annealing [73]. The relative concentration of Hf to Hf+Si (Hf/Hf+Si) was about 25 % for all samples. Microcrystals with a diameter of several nanometers formed in the film without nitrogen during the annealing as shown in Fig. 3.11a. The diameter decreased in the film with the nitrogen incorporation of 5 at.% and uniform amorphous film was obtained for the sample with the 30 at.% nitrogen concentration.

The mechanism of this phase separation suppression due to the nitrogen incorporation has not been investigated in detail yet. However, considering that the phase separation speed is limited by the diffusion of the constituent material, i.e. Hf, Si or both, in the alloy system [84], it is reasonable to think that nitrogen incorporation in the amorphous network retards such diffusion in the materials [96].



**Fig. 3.11** TEM photographs of HfSiON with different nitrogen content after 1,000 °C annealing in nitrogen ambient. Hf ratio (Hf/Hf+Si) is around 22 %. Nitrogen atomic concentration is **a** 0 at.%, **b** 5 at.% and **c** 30 at.%. Annealing was performed in N<sub>2</sub> ambient at 1,000 °C for 30 s. [73]



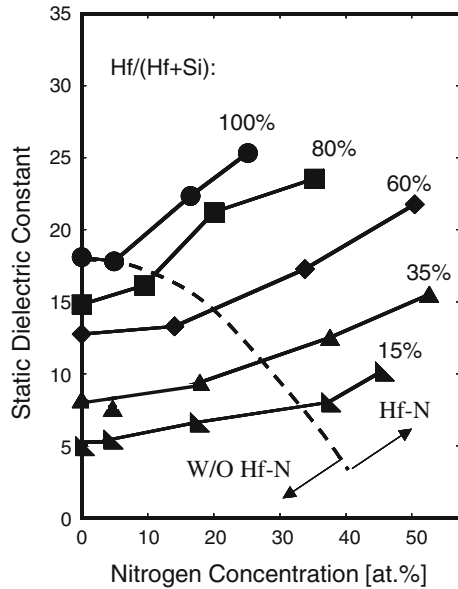
The nitrogen concentration criterion for keeping the thermal stability of HfSiO is definitely dependent on the hafnium content in the film. The lower limit of nitrogen concentration that can sustain the amorphous structure, avoiding the phase separation and crystallization, was investigated extensively for the case of 1,065 °C spike annealing in nitrogen ambient [97]. The lower limit increases gradually when Hf content increases in the film, however, for films with Hf/Hf+Si larger than 80 %, the lower limit increases dramatically.

Another beneficial effect of the nitrogen incorporation is the suppression of the diffusion of the impurities such as boron, phosphorous, and arsenic through the film [73, 98, 99]. As mentioned above, Quevedo-Lopez et al. [99] ascribed high diffusivity of the impurities in HfSiO to the formation of the grain boundary as a result of the phase separation and high diffusion coefficient at the grain boundary. They claim the reason for the lower diffusion coefficient of impurities through HfSiON resides in the film uniformity as well as the fact that the film retains its amorphous structure. Koyama et al. [100] also investigated the correlation between the diffusion coefficient and the film structure, concluding that not only film uniformity but also the film microscopic structure changes the coefficient: larger nitrogen concentration and lower Hf concentration lead to the lower diffusion coefficient of boron at 1,000 °C even among amorphous films.

Nitrogen incorporation in the HfSiO alloy increases the film dielectric constant as shown in Fig. 3.12 [97]. Hf/Hf+Si ratio is taken as a parameter in this figure. The measurement is performed for films without any high-temperature annealing, meaning that there is no influence of the phase separation in this case. As expected, increase in Hf content leads to an increase in the dielectric constant from the SiO<sub>2</sub> value to the HfO<sub>2</sub> value. For any relative Hf concentration, larger nitrogen



**Fig. 3.12** Dielectric constant of HfSiON as a function of the nitrogen incorporation, evaluated from CV measurement for thick (about 100 nm) films with various concentrations. Hf/Hf+Si is taken as a parameter. Nitrogen atomic concentration [N] region where clear Hf–N bonds are observed in XPS is also shown [97]

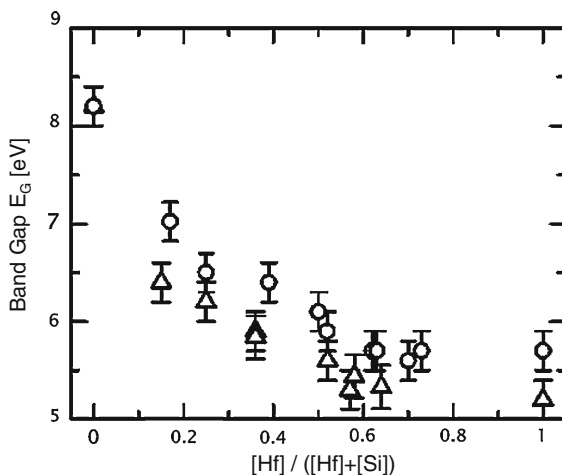


incorporation generally increases the dielectric constant. Considering that  $\text{Si}_3\text{N}_4$  has a larger  $k$ -value (about 7.8) than that of  $\text{SiO}_2$ , nitrogen incorporation is expected to enhance the dielectric constant; however, Koike et al. claimed that the dielectric constant increases drastically with Hf–N bond formation for a certain [N] level depending on Hf content in the film [97]. The dotted line in Fig. 3.12 indicates the critical [N] concentration. It is considered that HfSiON in the high dielectric constant region can be regarded as a pseudo-quaternary alloy consisting of four insulating components:  $\text{SiO}_2$ ,  $\text{HfO}_2$ ,  $\text{Si}_3\text{N}_4$ , and  $\text{Hf}_3\text{N}_4$  [17]. Since  $\text{Hf}_3\text{N}_4$  is an insulating material with a large dielectric constant as mentioned in Sect. 3.4 [60, 62], it is plausible to think that this component contributes to the abrupt increase in the dielectric constant of the alloy.

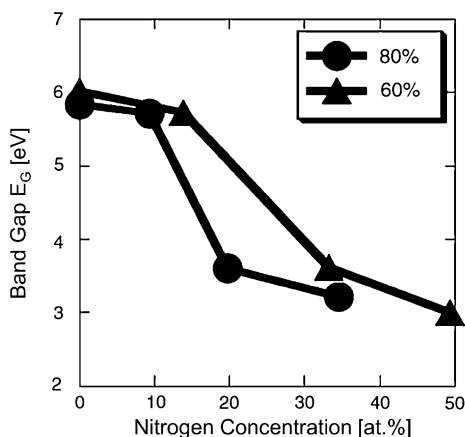
Band-gap energies of the HfSiO with various concentrations have been investigated with XPS, Ellipsometry, and REELS techniques [23, 25]. Kato et al., used XPS and ellipsometry to extract band-gap energy of HfSiO alloy and concluded that Hf incorporation into  $\text{SiO}_2$  network decreases the band-gap energy of the materials from that of  $\text{SiO}_2$  to the  $\text{HfO}_2$  value as in Fig. 3.13. An interesting point is that the value shows steep decrease at low Hf content and saturates at  $\text{HfO}_2$  value already at Hf/Hf+Si of about 0.6. This may be due to the difference in the microscopic structure between  $\text{SiO}_2$  and  $\text{HfO}_2$ . Coordination of Hf atoms in  $\text{SiO}_2$  crystal network may be different from that of Hf atoms in  $\text{HfO}_2$  ionic crystal structure, leading to the different conduction band energies [76].

Nitrogen incorporation further decreases the band-gap [23, 101]. Figure 3.14 shows band-gap energy of HfSiON with Hf/Hf+Si value of 0.6 and 0.8. Although nitrogen incorporation gradually decreases the band-gap energy at low nitrogen concentration [N], it leads to abrupt decreases at critical [N]s. It is ascribed to the

**Fig. 3.13** Band-gap energy estimated from the XPS analyses as a function of the composition  $x(=Hf/Hf+Si)$ . Circles and triangles represent data for HfSiO and ZrSiO, respectively. (Reprinted with permission from [25]. Copyright 2002, American Institute of Physics)



**Fig. 3.14** Band-gap energy of HfSiON with Hf/Hf+Si of 0.6 and 0.8. Nitrogen incorporation gradually decreases the band-gap energy at low nitrogen concentration [N]; however, it leads to abrupt decreases at critical [N]s, depending on the Hf contents in the alloy. It is ascribed to the Hf–N bond formation [23]



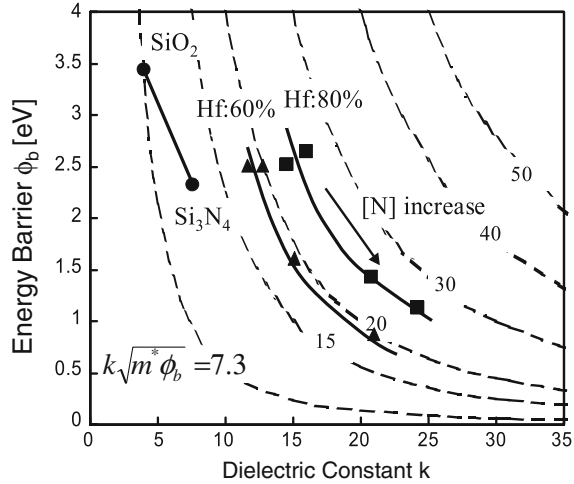
Hf–N bond formation [23]. As mentioned in Sect. 3.4, the nitrogen incorporation moves the valence band edge nearer to the vacuum level by the replacement of the edge from O<sub>2p</sub> band to N<sub>2p</sub> band [56], leading to the valence band offset decrease.

Contrary to the prediction by the theoretical calculation [56], the experimental result shows that nitrogen incorporation in HfO<sub>2</sub> as well as HfSiO network also leads to the decrease in the conduction band offset [23]. It has been reported that a large amount of N in HfO<sub>2</sub> network could induce a large amount of oxygen vacancy, V<sub>o</sub>, and lead to V<sub>o</sub>–V<sub>o</sub> interactions, resulting in the modification of the conduction band structure [102].

As mentioned in Sect. 3.2, the potential of high- $k$  material as a gate leakage suppressor is measured conveniently with the figure of merit (F.O.M.),

$$F.O.M. \equiv k\sqrt{m^*\phi_b}, \quad (3.3)$$

**Fig. 3.15** Contour plots of the F.O.M. for electrons tunneling from Si conduction band. Values for HfSiON films with various Hf and N concentrations are plotted. The effective mass of carriers  $m^*$  was set to be 1 for all concentrations for the purpose of relative comparison [103]



where  $k$ ,  $m^*$  and  $\phi_b$  are the film dielectric constant, effective mass of the carrier and the barrier height, respectively. The F.O.M. for HfSiON films with various Hf and N concentrations are plotted in Fig. 3.15 [103]. Since the effective masses both for electron and hole are not clear yet for HfSiON,  $m^*$  was set to be 1 for all concentrations for the purpose of relative comparison in this figure. Figure 3.15 shows that the F.O.M. becomes larger as Hf concentration increases, indicating the larger potential as a suppressor of gate leakage current. This figure also shows that, contrary to the F.O.M. enhancement due to the Hf relative concentration increase,  $[\text{N}]$  hardly changes the F.O.M. value. This is because the dielectric constant increase leads to the band-gap and barrier height decrease in HfSiON [23]. A similar figure concerning the hole tunneling is shown in Ref. [103]. Considering that larger  $[\text{N}]$  leads to higher thermal stability, HfSiON with high Hf and N concentrations seem to be more desirable. It should be mentioned, however that the discussion above is valid only when modulation of the Hf and N concentration does not change the effective mass of the carriers much. As mentioned in Sect. 3.2, Hf incorporation decreases the effective mass, while nitrogen incorporation into  $\text{SiO}_2$  network increases it [104, 105]. The effective mass should be evaluated for HfSiON before concluding the F.O.M. dependence on the concentration of each constituent material.

### 3.6 Hafnium–Aluminum-Based Gate Dielectrics

$\text{HfO}_2\text{--Al}_2\text{O}_3$  alloy is formed using chemical vapor deposition (CVD) [106], atomic layer deposition (ALD) [107], pulsed laser deposition (PLD) [108], jet vapor deposition (JVD) [109], and sputtering from composite target [110]. Normally, trimethyl-aluminum (TMA:  $\text{Al}(\text{CH}_3)_3$ ) is used for the  $\text{Al}_2\text{O}_3$  deposition in the CVD/

ALD process, whereas organic materials such as tetrakis-dimethylamino-hafnium (TDMAH:  $\text{Hf}[\text{N}(\text{CH}_3)_2]_4$ ) or tetrakis-diethylamino-hafnium TDEAH ( $\text{Hf}[\text{N}(\text{C}_2\text{H}_5)_2]_4$ ) are used for  $\text{HfO}_2$  deposition. Therefore, a certain amount of carbon is expected to be present in the film. The intermittent high-temperature annealing between ALD stages [111] or intermittent  $\text{NH}_3/\text{Ar}$  plasma treatment [112] is effective for reducing the carbon content in the material and this leads to the decrease in the leakage current through the film [112].

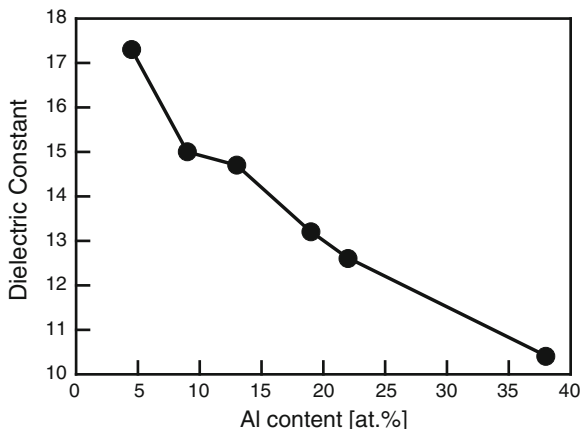
As in the case of Hf silicate, Hf silicide formation that leads to the increase in leakage current through the film is suppressed even at high temperature with  $\text{Al}_2\text{O}_3$  incorporation [108, 113], except for the case in which oxygen deficiency occurs in the deposition process [108]. Oxygen diffusion in  $\text{HfO}_2$  is also retarded by the presence of  $\text{Al}_2\text{O}_3$  component, leading to the suppression of interfacial layer growth in an oxygen-containing atmosphere at high temperature [113]. This phenomenon can be attributed to the very small diffusivity of oxygen in  $\text{Al}_2\text{O}_3$  compared to that in  $\text{HfO}_2$  [114].

However, impurity penetration phenomena such as boron diffusion, which is problematic for poly/high-k gate stack, are hardly suppressed in HfAlO system [115] because of the limited impurity diffusion barrier characteristics of  $\text{Al}_2\text{O}_3$  [116]. Moreover, phase separation between  $\text{HfO}_2$  and  $\text{Al}_2\text{O}_3$  accompanied by  $\text{HfO}_2$  crystallization occurs as in the case of HfSiO, when  $\text{HfO}_2$  content is large relative to  $\text{Al}_2\text{O}_3$  [111] and this partial crystallization may cause the inhomogeneity and the enlargement of the leakage current through HfAlO after the high-temperature annealing [117].

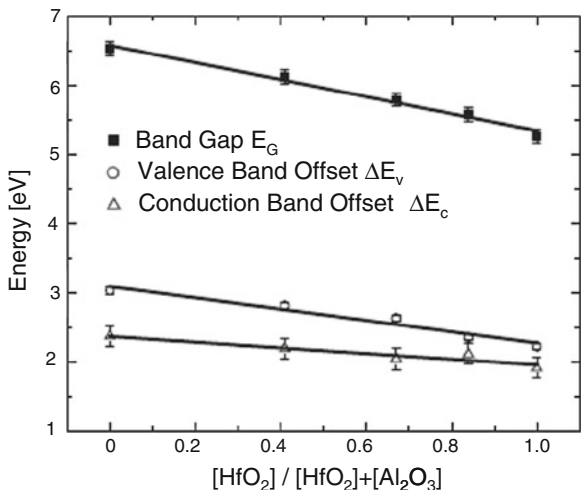
As expected from the dielectric constant values for  $\text{HfO}_2$  and  $\text{Al}_2\text{O}_3$ , the dielectric constants of HfAlO are intermediate between the two. They have been investigated thoroughly by the conventional C–V measurement with a number of samples with various Hf content and Al content as shown in Fig. 3.16 [118]. It should be noted that the measurements were performed with HfAlO layer deposited by metal organic chemical vapor deposition (MOCVD) in which Hf and Al precursors were introduced in the chamber at the same time with various flow ratios. The deposition process was performed at 500 °C, followed by post-metallization annealing at 400 °C. Therefore, homogeneous alloy structure was expected without the phase separation in this case. The film dielectric constant decreases from the value of  $\text{HfO}_2$  to that of  $\text{Al}_2\text{O}_3$  almost linearly as Al atomic concentration increases. Note that Al atomic concentration in pure  $\text{Al}_2\text{O}_3$  is 40 at.%.

The band-gap energy also takes intermediate values between those of  $\text{HfO}_2$  and  $\text{Al}_2\text{O}_3$  [32, 107]. Ohta et al. [119] used XPS measurement in order to extract the band-gap energy and the band alignment to Si using HfAlO with the cation ratio  $\text{Hf}/(\text{Hf}+\text{Al})$  of 0.3. This film was deposited by ALD using  $\text{HfCl}_4$ ,  $\text{Al}(\text{CH}_3)_3$ , and  $\text{H}_2\text{O}$  at 300 °C. Although the film was subjected to 1,050 °C RTA, it was confirmed that this film was not crystallized through the annealing [119]. The extracted band-gap was as large as 6.5 eV, which is an intermediate value between 5.8 eV ( $\text{HfO}_2$ ) and 8.8 eV ( $\text{Al}_2\text{O}_3$ ). From the valence band spectra of XPS, it was confirmed that the valence band and conduction band offsets of this HfAlO are as

**Fig. 3.16** Dielectric constant modification by the increase of Al content in HfAlO alloy [133]



**Fig. 3.17** Dependence of band-gap  $E_G$ , the valence band offset  $\Delta E_v$ , and the conduction band offset  $\Delta E_c$  for  $(\text{HfO}_2)_x(\text{Al}_2\text{O}_3)_{1-x}$ . The  $E_G$  and  $\Delta E_v$  data were obtained by XPS measurements. The  $\Delta E_c$  data are calculated by the  $\Delta E_c = E_G - \Delta E_v - E_{G, \text{Si}}$ . The solid lines are linear least square fits of the data points. (Reprinted with permission from [32]. Copyright 2002, American Institute of Physics)



large as 1.78 and 3.6 eV, respectively [119]. Similar result is reported from another group as shown in Fig. 3.17 [32]. These values are large enough to suppress the leakage current through the high-k film.

Nitrogen is added to the material in order to suppress the thermal stability issues such as impurity penetration as well as phase separation and crystallization. Nitrogen incorporation is realized by co-sputtering of  $\text{HfO}_2$  and  $\text{Al}_2\text{O}_3$  in  $\text{Ar}/\text{O}_2/\text{N}_2$  ambient [115], or  $\text{NH}_3$  annealing during or after the  $\text{HfAlO}$  deposition process [111]. Nitrogen incorporation is effective for suppressing the boron diffusion [115]. It is also effective for enhancing the crystallization temperature [111]. This practice even enhances the dielectric constant of the material [120] as in the case of  $\text{HfSiO}$ , which is beneficial for the realization of smaller EOT gate stack.

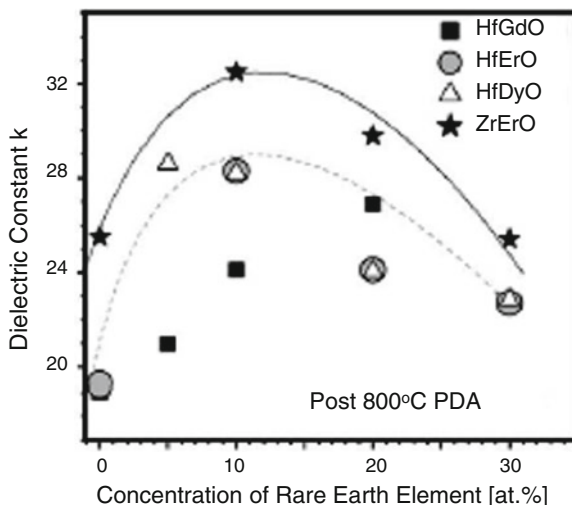
### 3.7 Doped Hafnium-Based Gate Dielectrics

For the further increase in figure of merit (F.O.M.) of the gate insulator, high- $k$  material with a dielectric constant of more than 25 has been searched for, recently. As La has very large ionic polarizability among cation candidates [18], La-based high- $k$  material exemplified by  $\text{La}_2\text{O}_3$  has been investigated intensively. Since this topic is discussed in detail in Chap. 9, research for obtaining a large dielectric constant of more than 25 using Hf-based material is discussed in this section.

A reasonable way of enlarging the dielectric constant is to mix  $\text{HfO}_2$  with other  $\text{MO}_x$  that has larger dielectric constant. Since  $\text{TiO}_2$  has a very large dielectric constant of more than 50 [26], this material may be a good candidate for the mixed oxide.  $\text{HfTiAlO}$  was deposited by co-sputtering for that purpose and it realized a high dielectric constant of 36 [121]. The drawbacks of  $\text{TiO}_2$  are its small band-gap energy (about 3 eV) and almost negligible barrier height from Si conduction band edge [122]. However, this  $\text{HfTiAlO}$  showed a barrier height of 1.3 eV from the Si conduction band edge, which is not much smaller than that of  $\text{HfO}_2$ . Mixing  $\text{TiO}_2$  with  $\text{HfO}_2$  may change the coordination of Ti atoms inside the alloy, resulting in the realization of a larger barrier height.

Since theoretical study predicted that the cubic  $\text{HfO}_2$  and the tetragonal  $\text{HfO}_2$  could have large dielectric constants of 29 and 70 [123] respectively, much work has been done to transform the material to such higher dielectric constant phases. It is well known that the monoclinic phase is the most stable one below 1,750 °C [10, 124]. Therefore, the research interest resides in how to stabilize the cubic and tetragonal phases at temperatures as low as 1,000 °C. Small amount of yttrium doping into  $\text{HfO}_2$  network (about 5 at.%) enhanced the dielectric constant value to 27, as a result of the formation of cubic  $\text{HfO}_2$  phase [125]. The process temperature for the transformation was as low as 800 °C in this case. The transformation was confirmed by the X-ray diffraction (XRD) spectrum, which indicated the molar volume reduction of the material from that of the monoclinic phase [125]. The doping species is not limited to yttrium but other rare earth materials such as La [126], Gd [127], Er [127], Dy [127, 128], Sc [128] or Ce [129] show a similar effect. Figure 3.18 shows the dielectric constant variation of Gd-, Er-, and Dy-doped  $\text{HfO}_2$  as a function of the doping concentration [127]. In the case of Gd, the larger doping with 20 at.% shows the largest dielectric constant, 27, whereas in the case of Er and Dy, doping level of 10 at.% gives the largest value, 28, as a result of the tetragonal  $\text{HfO}_2$  formation. It has also been reported that Dy doping transforms  $\text{HfO}_2$  to its cubic phase and enhances the dielectric constant to 32 [128]. It should be noted that these enhancements cannot be ascribed to the rare earth oxide formation in the materials, because those oxides have smaller dielectric constant than monoclinic  $\text{HfO}_2$ . Actually, larger doping of the rare earth elements resulted in the decrease of the dielectric constant of the materials, as shown in Fig. 3.18. It has also been reported that even a small amount of Si doping (about Si/Hf+Si ratio of 0.05) could lead to the tetragonal  $\text{HfO}_2$  formation, which enhances the dielectric constant to 27 [130].

**Fig. 3.18** Dielectric constant variation of Gd-, Er-, and Dy-doped  $\text{HfO}_2$  as a function of the doping concentration. Data for Er-doped  $\text{ZrO}_2$  is also shown. (Reprinted with permission from [127]. Copyright 2007, American Institute of Physics)



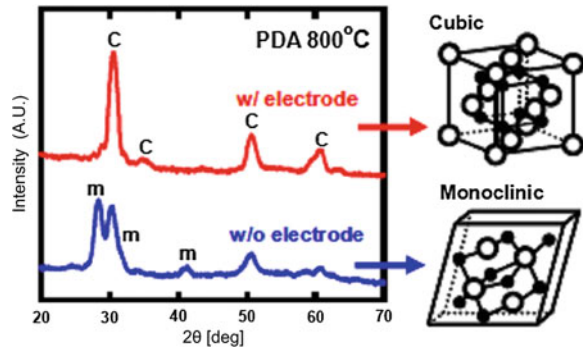
The doping was realized by using the rare earth oxide or  $\text{SiO}_2$  target simultaneously when  $\text{HfO}_2$  is sputter-deposited. Doping is also performed by adding Dy or Sc containing precursors in  $\text{HfO}_2$  MOCVD. Although the doping effect on the  $\text{HfO}_2$  phase transformation is not clear yet, some authors claim that the larger ionic radii of rare earth elements raise internal compressive stress in the  $\text{HfO}_2$  network, leading to the realization of the smaller molar volume (tetragonal or cubic) phases [127]. Theoretical study using first-principles calculations of doped  $\text{HfO}_2$  indicated that the doping modifies the energy of the monoclinic and the tetragonal phases and concluded that Si and Ge are favorable elements for the modification. However, it could not elucidate any clear dependence of the energy modification on the ionic radii of the dopants [131].

There is another method in which the stress is induced not by the doping but by the annealing with a cap layer. Migita et al. indicated that the low-temperature post-deposition annealing (PDA) of  $\text{HfO}_2$  with the cap gate electrode (TiN or TaN) enhances the dielectric constant up to about 50 [132]. Figure 3.19 shows the XRD spectra and indicates that the crystal phase after PDA with the cap electrode is that of cubic  $\text{HfO}_2$ , whereas the monoclinic phase is obtained when PDA is performed without the cap. The advantage of this method is that there is no need to use elements other than Hf and oxygen. This could circumvent the problems associated with the use of other elements in the LSI process such as cross-contamination.

These larger dielectric constant phases have band-gap energy similar to that of the monoclinic  $\text{HfO}_2$  [128, 132]. Even with the doping of rare earth elements, the band-gap energy does not decrease [24, 128]. Therefore, the gate stacks with these materials show drastic reduction of leakage current by about three orders of magnitude, compared with monoclinic  $\text{HfO}_2$  with similar EOT [127, 128].

A practical issue concerning the usage of tetragonal and cubic phase  $\text{HfO}_2$  materials as gate insulators is that subjecting them to the higher annealing

**Fig. 3.19** XRD spectra of  $\text{HfO}_2$  films after PDA at  $800^\circ\text{C}$  with and without electrode. ‘C’ represent peaks from cubic  $\text{HfO}_2$ , whereas ‘m’ represents those of monoclinic phase. The cubic  $\text{HfO}_2$  is successfully formed by PDA only with the electrode cap. (©2008 IEEE, Reprinted, with permission, from [132])



temperature typical of LSI processes diminishes their dielectric constant. For all the cases mentioned above, low-temperature annealing (typically at  $800^\circ\text{C}$ ) was performed for the realization of the large dielectric constant phases. However, it was reported that annealing at higher temperatures such as  $1,000^\circ\text{C}$  drastically reduced the dielectric constant of tetragonal  $\text{HfO}_2$  from 50 to 25 [132]. Yttrium-doped  $\text{HfO}_2$  also showed the decrease after  $1,000^\circ\text{C}$  annealing [125]. This degradation phenomenon seems reasonable because the stable phase of  $\text{HfO}_2$  is monoclinic below  $1,750^\circ\text{C}$  [124]. This tetragonal or cubic  $\text{HfO}_2$  is thought to be unstable at  $1,000^\circ\text{C}$ . Therefore, higher-temperature annealing and longer annealing yield the stable monoclinic phase, leading to the reduction of the dielectric constant. Some measures for keeping the high dielectric constant phase even at  $1,000^\circ\text{C}$  annealing should be developed for the application of this measure to advanced LSIs. The subject of phase modifiers in promoting the higher-k phases of the high-k oxides has been analyzed in detail in Chap. 10.

### 3.8 Summary

$\text{HfO}_2$  is regarded as one of the most promising materials for the high-k gate dielectrics. This is mainly because this material possesses a large dielectric constant originating from a large polarizability of the  $\text{Hf-O}$  molecule, and large band-gap energy. These two characteristics lead to the dramatic reduction of the leakage current through the thin film of this material, compared to the conventional  $\text{SiO}_2$  gate dielectrics. Thermal stability is also required for the high-k materials in order to maintain the integrity of the MOSFET through the LSI processes.  $\text{HfO}_2$  satisfies this requirement, however severer process condition may cause some problems even with this material: anomalous silicide formation in the film may enhance the leakage current, enhanced interfacial layer growth may increase the equivalent oxide thickness, and dopant-impurity penetration through the film may cause threshold voltage fluctuation. Suppression of those issues is the incentive for the incorporation of other components such as  $\text{SiO}_2$  and  $\text{Al}_2\text{O}_3$  into  $\text{HfO}_2$ .



The pseudo-alloy shows good characteristics to overcome most of those issues. Nitrogen incorporation is very effective for keeping homogeneous structure avoiding the phase separation of the pseudo-alloy. As  $\text{HfO}_2$  can have several phases, transformation of the monoclinic phase to the tetragonal and the cubic phases with larger dielectric constant has been pursued recently in order to further enhance the figure of merit of  $\text{HfO}_2$  as the gate dielectric.

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133. This figure was formed using data in Ref. [78]