# Mechanical Characterization of Graphene 4

Changhong Cao, Xuezhong Wu, Xiang Xi, Tobin Filleter, and Yu Sun

#### Keywords

Experimental mechanics • Graphene • Mechanical properties • Molecular dynamics simulations • Monolayer

# Introduction

Graphene is a planar monolayer of strongly  $sp^2$ -bonded carbon atoms arranged into a two-dimensional honeycomb lattice with a carbon–carbon bond length of 0.142 nm [\[1](#page-12-0)]. It was initially assumed not to exist in free state and was described as 'academic material' [[2\]](#page-12-0) until 2004 when Novoselov et al. [\[3](#page-12-0)] successfully separated single-layer graphene experimentally. It can be wrapped up into 0D fullerenes, rolled into 1D nanotubes or stacked into 3D graphite (Fig. [4.1\)](#page-1-0) [[4\]](#page-12-0). Five typical methods are typically used to make graphene sheets which include [\[5](#page-12-0)] (i) chemical vapor deposition (CVD) on metals, (ii) micromechanical exfoliation of graphite, (iii) reduction of graphene oxide, (iv) epitaxial growth on large band gap semiconductor  $S_iC_i$ , (v) and the creation of colloidal suspension.

In terms of mechanical properties, graphene is a super strong nanomaterial. A suspended monolayer graphene sheet over circular holes on a Si substrate was measured by AFM (atomic force microscopy) nano-indentation revealing that it exhibits a Young's modulus of  $\sim$ 1TPa and a critical failure stress and strain of

X. Wu • X. Xi

C. Cao  $(\boxtimes) \cdot$  T. Filleter  $\cdot$  Y. Sun

Department of Mechanical and Industrial Engineering, University of Toronto, Toronto, ON, Canada

e-mail: [changhongcao@hotmail.com;](mailto:changhongcao@hotmail.com) [filleter@mie.utoronto.ca](mailto:filleter@mie.utoronto.ca); [sun@mie.utoronto.ca](mailto:sun@mie.utoronto.ca)

College of Mechatronics Engineering and Automation, National University of Defense Technology, Changsha, China

e-mail: [www.67113@gmail.com;](mailto:www.67113@gmail.com) [fordada@126.com](mailto:fordada@126.com)

<span id="page-1-0"></span>

Fig. 4.1 Mother of all graphitic forms. Graphene is a 2D building material for carbon materials of all other dimensionalities. It can be wrapped up into 0D buckyballs, rolled into 1D nanotubes or stacked into 3D graphite (Figure/Caption reproduced (Adapted) with permission from Macmillan Publishers Ltd: [Nature Material] [\[4\]](#page-12-0), Copyright (2007))

130GPa and 25 %, respectively [[6\]](#page-12-0). These extraordinary mechanical properties, in addition to the well-documented beneficial electrical properties of graphene, have attracted a great deal ofinterestin areas such as nano-/microelectromechanical system (NEMS/MEMS) [\[7](#page-12-0), [8](#page-12-0)], nanoelectronics [[9,](#page-12-0) [10\]](#page-12-0), as well as nanocomposites [[11\]](#page-12-0). A detailed fundamental understanding of its mechanical properties is of importance in its application to a number of these fields, in particular nanocomposites and MEMS.

## Mechanical Characterization Techniques

Optical microscopy, electron microscopy, AFM, and Raman spectroscopy have been applied to the mechanical characterization of graphene. Optical microscopy is used to image samples 'macro'scopically, and electron microscopy provides higher imaging resolutions down to nanometer and sub-nanometer scales, facilitating both structural and mechanical characterization of graphene samples. AFM is a standard tool to obtain surface images of graphene and can also be applied in mechanical testing modes to measure mechanical properties such as Young's modulus, friction, and strength. Raman spectroscopy is used to investigate the vibrational properties of graphene and measure graphene thickness accurately, which is essential in the interpretation and analysis of graphene's mechanical behavior.

Optical microscopy is primarily used to detect graphene 'macro'scopically. Substrate design is of great importance in order to enhance the visibility of graphene under optical microscopy [[12](#page-12-0), [13](#page-12-0)]. Based on the Fabry–Perot interference mechanism, various materials have been employed to enhance imaging contrast. For example, a  $SiO<sub>2</sub>$  layer is usually created on the surface of a silicon substrate [\[14\]](#page-13-0). By adjusting the  $SiO<sub>2</sub>$  thickness to 90 or 300 nm, the intensity of the reflected light is at the maximum, which is also the maximum sensitivity of human eye [[3\]](#page-12-0). In addition, 50 nm  $Si<sub>3</sub>N<sub>4</sub>$  and 72 nm  $Ai<sub>2</sub>O<sub>3</sub>$  substrates have also been used to improve the contrast of graphene  $[14, 15]$  $[14, 15]$  $[14, 15]$  $[14, 15]$ . Additionally, fluorescence quenching microscopy (FQM) was used to image graphene, reduced graphene oxide (RGO), and graphene oxide (GO) for sample evaluation and manipulation so that the synthesis process can be improved  $[16]$  $[16]$ .

Scanning electron microscopy (SEM) and transmission electron microscopy (TEM) are widely used in nanomaterials research. Similar to optical microscopy, an underlying substrate is needed for using SEM to image graphene sheets. The magnification of SEM is order of magnitude higher than that of optical microscopy typically achieving a spatial resolution on the order of tens to a few nms.

TEM is used to observe morphological and structural features of graphene and measure the number of graphene layers accurately. For samples in which the edges of graphene films fold back, the observation of these edges by TEM provides an accurate way to count the number of layers at multiple locations on the films, as shown in Fig. [4.2a](#page-3-0) [\[17\]](#page-13-0). In addition, TEM is often assisted with electron diffraction pattern analysis, which enables the observation of the crystal structure of graphene sheets [[18\]](#page-13-0) (Fig. [4.2b\)](#page-3-0). Defects in graphene can also be detected by using this method [\[18](#page-13-0)]. However, it should be noted that the resolution of electron microscopy is limited by the accelerating voltage. High accelerating voltages can damage the monolayer of graphene. Aberration-corrected TEM has been reported to achieve a 1 A resolution at an acceleration voltage of only 80 kV  $[19, 20]$  $[19, 20]$  $[19, 20]$  $[19, 20]$ .

Atomic force microscopy (AFM) can provide a direct way to observe the topography of single to few layer graphene films which are atomically thin. In theory, the thickness of a single layer of graphene is about 0.34 nm. However, due to the surface adsorption of graphene, the actual measured value consistently appears to be 0.8–1.2 nm. The number of layers can be obtained by measuring the thickness of a single-layer sheet, and it can be calculated according to [\[21\]](#page-13-0)

$$
N = \frac{P - X}{0.34} + 1
$$

where  $N$  is the number of layers of graphene,  $P$  is the measured thickness, and  $X$  is the measured thickness of a single-layer graphene. Furthermore, AFM can be used

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Fig. 4.2 TEM images: (a) Edges of graphene films (Figure/Adaption Reprinted (Adapted) by permission from Macmillan Publishers Ltd. on behalf of Cancer Research UK: [Nature] [\[17\]](#page-13-0), Copyright (2009)) (b) Hexagonal pattern of the graphene structure (Figure/Caption Reprinted (adapted) with permission from [\[18\]](#page-13-0), Copyright (2008) American Chemical Society)

to investigate the mechanical properties of graphene such as elastic modulus, tensile strength, bending stiffness, and friction. In this case, several mechanical testing modes of AFM are implemented including friction force microscopy (FFM), for determining the frictional characteristics of graphene  $[22-24]$ , and AFM deflection for measuring elastic modulus and strength [\[6](#page-12-0)].

Raman spectroscopy provides a fast and nondestructive way to gain insight into the electron–phonon interactions in graphene. The frequency of scattered photons is changed due to the interaction of incident light and the material's molecular structure. Thus, the structural features of graphene can be identified, such as the number of graphene layers, the molecular structure, and defects in graphene [[25\]](#page-13-0). Figure [4.3](#page-4-0) shows the spectra of single-layer and double-layer graphene sheets [[26\]](#page-13-0). Raman spectroscopy can also be used to measure the mechanical properties of graphene [[27\]](#page-13-0). Mingyuan Huang et al. [[28\]](#page-13-0) presented Raman spectra of optical phonons in graphene monolayers under tunable uniaxial tensile stress. They showed that all the prominent bands exhibit significant red shifts and the resulting shift rates can be used to calibrate strain in graphene.

# Mechanical Characterization of Graphene

### Young's Modulus

#### Simulation

The elastic properties of graphene can be estimated by using numerical simulations, based on elasticity theory. It is known that graphene is a two-dimensional crystal of carbon atoms bonded by  $sp^2$  hybridized bonds. Interatomic force field models for graphene thus can be built. Based on the continuum elasticity theory, a simple

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Fig. 4.3 Raman spectra of single and double layer grapheme (Figure/Caption Reprinted (adapted) with permission from [[26](#page-13-0)], Copyright (2007) American Chemical Society)



valence force field model was formulated by Keating [[29\]](#page-13-0) for semiconductors and then extended to graphene by Lobo [\[30](#page-13-0)]. Results predicting the Young's modulus of graphene were reported based on different modeling theories. Table 4.1 [\[31–35](#page-13-0)] summarizes these results from different simulation studies. It can be seen that with different modeling approaches, the Young's modulus values are close to 1 TPa.

## Experiment

Reported experimental results of the Young's modulus are  $\sim$  1.0 TPa [[6\]](#page-12-0) which is in good agreement with simulations. A suspended monolayer graphene sheet over circular holes on a Si substrate was measured using AFM nano-indentation (Fig. [4.4\)](#page-5-0) [[36\]](#page-13-0). The thickness of the graphene sample was accurately determined using contact mode AFM imaging [\[37](#page-13-0)]. AFM was also used to accurately resolve the small forces involved in deformation and friction of graphene [[36](#page-13-0), [38](#page-14-0)]. Young's modulus is calculated according to

<span id="page-5-0"></span>

$$
\sigma = E\varepsilon + D\varepsilon^2
$$

where  $\sigma$  is the stress,  $\varepsilon$  is the strain, and D is the third-order elastic modulus.

Elastic properties of graphene with different boundary conditions were also measured. Wong et al. [\[29](#page-13-0)] investigated the mechanical properties of suspended graphene drums via static deflection experiments. Deflection was detected using AFM, which is highly sensitive to topographical changes in the out-of-plane direction and can pick up height variations as small as 1 nm. The AFM scans provided information on both the peak displacement amplitude and the deflection mode shape of the drum structure. It was found that the structures have linear spring constants ranging from 3.24 to 37.4 N  $m^{-1}$  and could be actuated to about 18–34 % of their thickness before exhibiting nonlinear deflections [[29\]](#page-13-0). The result indicates a Young's modulus value of 1TPa which agrees well with the results reported by Lee et al.  $[6]$  $[6]$ .

## Bending Stiffness

## Simulation

Theoretical studies have suggested that the bending stiffness of graphene is critical to attain structural stability for suspended graphene sheets, which in turn affects their mechanical properties significantly [[39–41\]](#page-14-0). Molecular mechanics simulations for graphene bending rigidity were reported through calculations of the strain energy for graphene sheets subjected to a point loading [[42\]](#page-14-0). Based on the firstgeneration Brenner potential, an analytical form was derived for the bending modulus of monolayer graphene under infinitesimal bending curvature [\[43](#page-14-0), [44\]](#page-14-0), namely,

Fig. 4.4 Measurement of graphene elastic modulus:

from [[36](#page-13-0)], Copyright

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$$
D_{\text{bending}} = \frac{\sqrt{3}}{2} \frac{\partial V_{ij}}{\partial \cos \theta_{ijk}} = \frac{1}{2} V(r_0) \frac{\partial b_{ij}}{\partial \theta_{ijk}}
$$

where the interatomic potential takes the form

$$
V_{ij} = V_R(r_{ij}) - \overline{b}V_A(r_{ij})
$$

and  $\overline{b} = (b_{ij} + b_{ji})/2$  is a function of the bond angles.

Simulation results show that the bending rigidity increases from a small value of 0.819 eV with the size of 1.87 nm to an asymptotic value of 2.385 eV for sheets with sizes larger than 12 nm. In addition, rigidity changes from the asymptotic value, 2.385 eV, for a square sheet to a smaller value, 0.360 eV, for a sheet with the shape ratio of  $b/a = 2.15$ , as shown in Fig. 4.5 [[42\]](#page-14-0).



Fig. 4.6 (continued)

#### Experiment

While in-plane mechanical properties such as elastic modulus and strength of monolayer graphene have been deduced from experiments [[6,](#page-12-0) [45\]](#page-14-0), direct measurements of bending stiffness of monolayer graphene have not been reported. By using nano-indentation on suspended multilayer graphene flakes, the bending stiffness for 8–100 layers has been measured and determined to be in the range of  $2 \times 10^{-14}$  N m<sup>-2</sup>  $\times 10^{-11}$  N m<sup>-1</sup> [\[46](#page-14-0)].

# Friction

#### Simulation

Molecular dynamics simulations can also be applied to extract frictional properties of graphene. A probing tip using a short-capped single-walled carbon nanotube is able to capture the frictional characteristics and resolve the graphene lattice through measuring oscillatory lateral forces or normal forces. As shown in Fig. 4.6a [[47\]](#page-14-0), the graphene sheet (atoms are colored in red and blue) has a length l and a diameter d. The initial tip–surface distance is  $h$ , and the sliding distance is  $D$ . The atoms in blue are prescribed with a moving velocity while the atoms in red follow the molecular dynamics. The graphene sheet with lateral dimensions of  $7.6 \times 7.6$  nm was placed horizontally with its boundary atoms (in pink) fixed and its interior atoms (in green) following the molecular dynamics. By averaging the oscillatory lateral force and normal force along the tip moving path, the friction coefficient was extracted [\[47](#page-14-0)].

Based on the simulation, it was found that the friction coefficient decreases with an increase in the initial tip–surface distance and the number of graphene layer [\[47](#page-14-0)]. The calculated results for the friction force  $\overline{F}_x$  and the friction coefficient  $C_x$  at different initial tip–surface distances and the number of graphene layers are shown in Fig. 4.6b [\[47](#page-14-0)].

#### Experiment

The frictional properties of graphene can be measured using AFM under applied loads. Initially, the AFM tip is made to contact the specimen surface at a fixed normal load, and the friction force is measured over a distance for multiple cycles in which friction loops are recorded. Then, the applied load is varied, and the friction force is measured again for multiple applied loads. This process is repeated until the tip becomes detached from the specimen surface due to excessive negative loads and the relationship between normal force and frictional forces is determined [\[23](#page-13-0), [24](#page-13-0), [38\]](#page-14-0).

<sup>≺</sup> 

Fig. 4.6 (A) Schematic of the probing interactions between the capped nanotube and graphene layer (B) Tip sliding along the x-direction, (a) the friction force (c) the friction coefficient; along the y-direction (b) the friction force (d) the friction coefficient (Figure/Caption Reprinted (Adapted) from Carbon [[47](#page-14-0)], Copyright (2011), with permission from Elsevier)

Measurements down to single atomic sheets revealed that friction monotonically increased as the number of layers decreased for suspended graphene [\[23](#page-13-0)]; a similar trend revealed that monolayer epitaxial graphene on SiC exhibited higher friction than bilayer graphene, and the single and bilayer graphene reduced friction on the SiC substrate by a factor of  $\sim$ 10 [[24\]](#page-13-0). In another study, binding the graphene strongly to a mica surface suppressed the friction trend as a function of thickness [\[23](#page-13-0), [24\]](#page-13-0). Li-Yu Lin et al. [[38\]](#page-14-0) investigated the friction and wear characteristics of multilayer graphene films deposited on a  $S_i$  substrate. The graphene films consisted of a few layers of carbon basal plane. The number of graphene layers was determined by AFM and Raman spectroscopy. It was found that graphene films exhibited much lower friction (from 0.36 to 0.62 nN) than bare Si surface (from 1.1 to 4.3 nN) when applied loads varied from 3 to 30 nN.

## Strength and Fracture

Depending on single or polycrystalline nature, the fracture mechanisms of graphene can vary. For single crystalline graphene, the failure mechanism (take double vacancy as an example) by molecular dynamics simulation is illustrated in Fig. [4.7a](#page-10-0) [\[48\]](#page-14-0). Before the propagation of fracture, pentagons and heptagons initiate around the defects. As load increases, fracture disperses in the direction parallel to the loading direction. Two chains of atoms are formed at the region of the vacancy defects and the other two chains on the edges of the sheet during fracture propagation. The critical stress and critical strain relationship depends on the number of vacancy defects [[48,](#page-14-0) [49\]](#page-14-0). As shown in Fig. [4.7b](#page-10-0) [[48](#page-14-0)], a pristine single-layer graphene sheet has the highest ultimate strength and strain. The presence of vacancy defects can reduce these features of the graphene sheet. A single-layer graphene sheet with a single vacancy defect has a critical stress and strain reduction of 6.4 % and 9.7 %, respectively, compared with the pristine sheet. A reduction of 7.3 % and 11.46 % for a double vacancy defect graphene sheet was also estimated.

Simulation performed by Cao et al. [[50\]](#page-14-0) shows that for polycrystalline graphene, fracture initiates from either a grain boundary triple junction [Fig. [4.8a–d](#page-11-0)] or an array of vacancies on a preferential grain boundary [Fig. [4.8e–h\]](#page-11-0) by unzipping atomic bonds along a preferential grain boundary. Crack propagation takes only 4.0 ps from crack initiation to the final failure of the entire sample. Furthermore, polycrystalline graphene exhibits 'flaw tolerance' as reported in the simulation study reported by Zhang et al. [\[51](#page-14-0)]. As illustrated in Fig. [4.9,](#page-11-0) fracture behavior of polycrystalline graphene can become insensitive to a preexisting flaw (hole or notch) below a critical length scale, which means that there is no stress concentration near the flaw site.

Experimental fracture behavior of graphene is relatively understudied. Research in experimentally revealing fracture mechanisms is a promising future field, which can significantly enhance the understanding of failure modes of graphene in various applications.

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Fig. 4.7 (A) Failure process of a SLGS with double vacancy (shortest possible separation distance). (a) after initial relaxation, (b) formation of pentagons and heptagons at adjacent of the vacancies, (c) early stages of fracture propagation, (d) formation of two chains at the location of vacancy defects, and (e) formation of four chains. (B) Stress–strain curves of graphene sheets with double vacancies, single vacancy, and a perfect structure (Figures/Captions Reprinted (Adapted) from Solid State Communications [\[48\]](#page-14-0), Copyright (2011), with permission from Elsevier)

<span id="page-11-0"></span>



Fig. 4.9 Schematic shows the fracture behavior of the polycrystalline graphene can be insensitive to preexisting flaw (Figure/Caption Reprinted (adapted) with permission from [\[51\]](#page-14-0), Copyright (2012) American Chemical Society)

# Summary

The synthesis, characterization, and applications of graphene are rapidly progressing. Mechanically, Young's modulus of mechanical exfoliated monolayer

<span id="page-12-0"></span>graphene was determined to be  $\sim 1.0$ TPa, and a critical failure stress and strain of 130 GPa and 25  $\%$  were measured, respectively; bending stiffness was reported only for 8–100 layers to be in the range between  $2 \times 10^{-14}$  N m<sup>-1</sup> and  $2 \times 10^{-17}$  N m<sup>-1</sup>; friction of graphene shows a trend that it monotonically increases as the number of layers decreases for suspended graphene and epitaxial graphene on S<sub>i</sub>C; fracture behavior of graphene depends on its nature of crystalline. Simulation results show that polycrystalline graphene exhibits 'flaw tolerance'. However, debates over properties such as wear and shear modulus remain unsolved. In addition to the techniques this chapter described, MEMS devices which have micrometersized features can also be applied to study single-layer and multilayer graphene as they can bridge investigation of materials at the macro- and nanoscales. Although a number of MEMS devices have been developed for mechanical characterizing individual 1D nanomaterials, such as carbon nanotubes [\[52](#page-14-0)] and silicon nanowires [\[53\]](#page-14-0), few devices exist for characterizing 2D nanomaterials such as graphene. In future work, the development of advanced experimental techniques and methodologies for interpreting data will produce more thorough mechanical characterization results and enable additional in-depth understanding of this important nanomaterial.

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