

# Benchmark 3D: The Cell-Centered Finite Volume Method Using Least Squares Vertex Reconstruction (“Diamond Scheme”)

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## 1 Presentation of the scheme

We consider, for this contribution, the cell-centered finite volume method based on least squares vertex reconstruction. This method, which is also popularly known as “*the diamond scheme*”, was originally presented for the advection-diffusion equation in two-dimensions and then extended in 3-D. The discretization of the diffusive term in 2-D and 3-D is found in in [1, 3–5]. The scalar solution of the diffusion problem  $u$  is numerically approximated by a piecewise constant function  $u_T$  on the cells  $K$  of mesh  $T$ . The numerical approximation  $u_T$  is defined as  $u_T(x) = \sum_{K \in T} u_K \chi_K(x)$  ( $\chi_K(x)$  being the characteristic function of cell  $K$ ) through the values  $(u_K)_{K \in T}$ . To define the numerical diffusive flux through the interface  $f$  of the mesh, a polyhedral cell is built around this interface. This polyhedral cell, which has a quadrilateral shape in two dimensions, is named after its shape as “*diamond cell*”, which also motivates the name of the method. Specifically, let  $x_K \in K$  be the center of gravity of the cell  $K$  of mesh  $T$ . The diamond cell  $D$  associated to the interface  $f$  between two cells  $K$  and  $L$  in  $T$  is the convex hull  $D = \text{hull}(f, x_K, x_L)$ . If  $f$  is a boundary face, thus defined by  $f = \partial K \cap \partial \Omega$  where  $\partial \Omega$  is the boundary of the computational domain  $\Omega$ , then the diamond cell associated to  $f$  is the convex hull  $D = \text{hull}(f, x_K)$ .

The numerical diffusive flux is built by using a constant approximation of the solution gradient on each diamond cell. Let  $(x_1, x_2, \dots, x_m)$  denote the vertices of face  $f$ ,  $x_K$  and  $x_L$  the centers of gravity of the cells  $K$  and  $L$  that share this face, and  $D$  the convex hull of these points. For any function  $u \in H^1(D)$ , the Green-Gauss formula yields the relation  $\int_D \nabla u(x) dx = \int_{\partial D} u(x) n(x) d\sigma(x)$ , where  $n$  is

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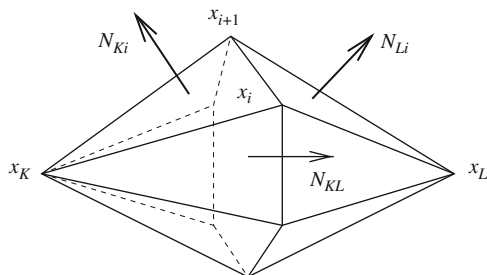


Fig. 1 Geometry of the diamond cell

the unit vector orthogonal to  $\partial D$  and pointing out of  $D$ . If the restriction of  $u$  to the face  $f$  of  $\partial D$  is an affine function, the boundary integral only depends on the values of  $u$  at the vertices of  $D$ . In this case, the Green-Gauss divergence theorem yields

$$\begin{aligned} \frac{1}{|D|} \int_D \nabla u(x) dx &= \frac{1}{|D|} \int_{\partial D} u(x)n(x) dx = \frac{1}{|D|} \sum_{f \in \partial D} \int_f u(x)n(x) dx \\ &= \frac{1}{3|D|} \sum_{i=1}^m (N_{Ki}(u(x_i) + u(x_{i+1}) + u(x_K)) + N_{Li}(u(x_i) + u(x_{i+1}) + u(x_L))) \\ &= \frac{1}{3|D|} \left( (u(x_L) - u(x_K)) N_{KL} + \sum_{i=1}^m u(x_i) N_i \right), \end{aligned}$$

where  $N_{Ki}$  and  $N_{Li}$  are the vectors orthogonal to the triangular facets hull( $x_K, x_i, x_{i+1}$ ) and hull( $x_L, x_i, x_{i+1}$ ), respectively, and having lengths equal to the measure of the facets; specifically,  $N_{Ki} = \frac{1}{2}(x_i - x_K) \times (x_{i+1} - x_K)$  and  $N_{Li} = -\frac{1}{2}(x_i - x_L) \times (x_{i+1} - x_L)$ , see Fig. 1. The vectors  $N_{KL}$  and  $N_i$  actually used during the computations are

$$N_{KL} = \sum_{i=1}^m N_{Li} = -\sum_{i=1}^m N_{Ki} \quad \text{and} \quad N_i = N_{Ki} + N_{Li} + N_{K_{i-1}} + N_{L_{i-1}}.$$

To derive the gradient formula, we consider the right-hand side of the above equality with the unknown  $u_K$  and  $u_L$  replacing  $u(x_K)$  and  $u(x_L)$  and some values  $u_i$  replacing the values  $u(x_i)$ . These values,  $u_1, u_2, \dots, u_m$ , are linearly interpolated from the values  $(u_K)_{K \in T}$  as follows. For any vertex  $x_i$  of the mesh  $T$ , we consider  $u_i = \sum_{K \in T_i} \omega_{iK} u_K$  where  $T_i = \{K : x_i \in K\}$  denotes the subset of the mesh cells which share the vertex  $x_i$ . The interpolation weights  $\omega_{iK}$  are assumed to verify the consistency relations [4]:

$$\sum_{K \in T_i} \omega_{iK} = 1 \quad \text{and} \quad \sum_{K \in T_i} \omega_{iK} (x_i - x_K) = 0.$$

The interpolation weights  $\omega_{iK}$  are obtained by solving the reconstruction problem that approximates the cell-averaged data set  $\{(x_K, u_K)$  for  $K \in T_i\}$  by the affine function

$$\tilde{u}_i(x) = \alpha + \beta \cdot (x - x_i) \text{ for } x \in \mathcal{V}_i$$

on the co-volume  $\mathcal{V}_i = \bigcup_{K \in T_i} K$  and in a least square sense, cf. [2, 4]. The reconstructed value at vertex  $x_i$  is now given by taking  $u_i = \tilde{u}_i(x_i) = \alpha$ . The coefficients  $(\alpha, \beta)^T$  are the minimizers of the least squares functional

$$\mathcal{J}(\alpha, \beta^T) = \sum_{K: x_i \in K} (\alpha + \beta \cdot (x - x_i) - u_K)^2.$$

Imposing the zero gradient condition, i.e.,  $\nabla_{\alpha, \beta} \mathcal{J}(\alpha, \beta^T) = 0$ , yields a linear system for the coefficients  $(\alpha, \beta)$ , whose solution returns the interpolation weights. The values  $u_i$  at the vertices  $x_i \in \partial\Omega$  on the Dirichlet boundary are constrained to the boundary data, for instance  $u_i = 0$  for a homogeneous condition. Other kinds of boundary conditions, e.g., Neumann or Robin, can be taken into account by extending to the 3-D case the technique investigated in [2]. Finally, the scheme reads as

$$\forall K \in T, \quad - \sum_{f \subset \partial K} \Lambda_f \nabla_{DU_T} \cdot N_{KL} = f_K |K| := \int_K f(x) dx,$$

where  $\Lambda_f$  is an arithmetic average of the diffusion tensor  $\Lambda$  over the diamond cell located around face  $f$  and  $N_{KL}$  is exactly the normal from above.

## 2 Numerical results

- **Test 1 Mild anisotropy**,  $u(x, y, z) = 1 + \sin(\pi x) \sin(\pi(y + \frac{1}{2})) \sin(\pi(z + \frac{1}{3}))$   
 min = 0, max = 2, **Tetrahedral meshes**

i	nu	nmat	umin	uemin	umax	uemax	normg
0	215	6985	3.02E-02	3.15E-02	1.949	1.948	1.627
1	2003	107331	2.03E-02	1.13E-02	1.989	1.995	1.730
2	3898	227618	6.84E-03	4.21E-03	1.989	1.990	1.750
3	7711	476645	9.13E-03	8.18E-03	1.994	1.995	1.767
4	15266	994892	5.52E-03	4.10E-03	1.997	1.997	1.776
5	30480	2072944	1.49E-03	2.57E-04	1.997	1.999	1.784
6	61052	4292073	1.83E-03	1.20E-03	1.997	1.998	1.789

i	nu	erl2	ratio12	ergrad	ratiograd	ener	ratioener
0	215	3.750E-02	-	3.503E-01	-	2.636E-01	-
1	2003	9.173E-03	1.892	1.568E-01	1.081	1.071E-01	1.210
2	3898	5.897E-03	1.991	1.215E-01	1.149	8.159E-02	1.225
3	7711	3.551E-03	2.230	9.410E-02	1.122	6.016E-02	1.339
4	15266	2.255E-03	1.994	7.387E-02	1.063	4.648E-02	1.132
5	30480	1.412E-03	2.032	5.768E-02	1.073	3.565E-02	1.152
6	61052	8.882E-04	2.001	4.502E-02	1.070	2.733E-02	1.147

Name of the solver: BiCG-stab with Jacobi preconditioner (in-house implementation).

- **Test 1 Mild anisotropy**,  $u(x, y, z) = 1 + \sin(\pi x) \sin(\pi(y + \frac{1}{2})) \sin(\pi(z + \frac{1}{3}))$   
min = 0, max = 2, **Voronoi meshes**

i	nu	nmat	umin	uemin	umax	uemax	normg
1	29	257	8.51E-02	-6.18E+00	1.870	7.968	20.435
2	66	660	1.43E-01	2.06E-01	1.854	1.846	1.887
3	130	1410	3.85E-02	7.00E-03	1.925	1.941	1.855
4	228	2620	1.74E-02	2.37E-02	1.914	1.920	2.067
5	356	4424	2.84E-03	-2.18E+00	1.979	3.546	3.274

i	nu	erl2	ratio12	ergrad	ratiograd	ener	ratioener
1	29	2.639E+00	-	3.631E+01	-	9.728E+00	-
2	66	9.077E-02	12.292	9.457E-01	13.307	3.751E-01	11.876
3	130	5.508E-02	2.210	7.434E-01	1.065	3.055E-01	0.907
4	228	6.650E-02	-1.006	1.163E+00	-2.391	3.224E-01	-0.287
5	356	3.674E-01	-11.507	5.071E+00	-9.912	1.287E+00	-9.321

Name of the solver: BiCG-stab with Jacobi preconditioner (in-house implementation).

- **Test 1 Mild anisotropy**,  $u(x, y, z) = 1 + \sin(\pi x) \sin(\pi(y + \frac{1}{2})) \sin(\pi(z + \frac{1}{3}))$   
min = 0, max = 2, **Kershaw meshes**

i	nu	nmat	umin	uemin	umax	uemax	normg
1	512	10648	3.03E-02	8.74E-02	1.958	1.916	1.768
2	4096	97336	1.06E-02	3.00E-02	1.993	1.973	1.700
3	32768	830584	1.75E-03	5.87E-03	1.997	1.991	1.726
4	262144	6859000	7.14E-04	9.88E-04	1.999	1.998	1.765

i	nu	erl2	ratiol2	ergrad	ratiograd	ener	ratioener
1	512	6.846E-02	-	6.798E-01	-	4.901E-01	-
2	4096	4.715E-02	0.537	3.403E-01	0.998	2.715E-01	0.851
3	32768	2.866E-02	0.718	1.831E-01	0.894	1.532E-01	0.825
4	262144	1.315E-02	1.123	8.289E-02	1.143	6.942E-02	1.142

Name of the solver: BiCG-stab with Jacobi preconditioner (in-house implementation).

- **Test 1 Mild anisotropy**,  $u(x, y, z) = 1 + \sin(\pi x) \sin(\pi(y + \frac{1}{2})) \sin(\pi(z + \frac{1}{3}))$   
min = 0, max = 2, **Checkerboard meshes**

i	nu	nmat	umin	uemin	umax	uemax	normg
1	36	424	1.54E-01	1.33E-01	1.846	1.833	1.588
2	288	4528	4.01E-02	3.47E-02	1.960	1.958	1.721
3	2304	41896	1.01E-02	8.74E-03	1.990	1.990	1.773
4	18432	360280	2.54E-03	1.98E-03	1.997	1.998	1.791
5	147456	2987704	6.36E-04	5.26E-04	1.999	1.999	1.796

i	nu	erl2	ratiol2	ergrad	ratiograd	ener	ratioener
1	36	1.356E-01	-	2.488E-01	-	3.406E-01	-
2	288	4.427E-02	1.615	1.471E-01	0.758	1.346E-01	1.339
3	2304	1.191E-02	1.894	7.031E-02	1.065	4.678E-02	1.524
4	18432	3.112E-03	1.936	3.410E-02	1.043	1.687E-02	1.471
5	147456	7.976E-04	1.964	1.695E-02	1.008	7.003E-03	1.268

Name of the solver: BiCG-stab with Jacobi preconditioner (in-house implementation).

- **Test 2 Heterogeneous anisotropy**,  $u(x, y, z) = x^3 y^2 z + x \sin(2\pi xz) \sin(2\pi xy) \sin(2\pi z)$ , min = -0.862, max = 1.0487, **Prism meshes**

i	nu	nmat	umin	uemin	umax	uemax	normg
1	1210	21308	-8.42E-01	-8.57E-01	0.978	0.977	1.481
2	8820	169418	-8.38E-01	-8.41E-01	1.010	1.011	1.638
3	28830	570328	-8.58E-01	-8.60E-01	1.032	1.033	1.676
4	67240	1350038	-8.57E-01	-8.58E-01	1.033	1.034	1.690

i	nu	erl2	ratio12	ergrad	ratiograd	ener	ratioener
1	1210	9.551E-02	-	2.356E-01	-	2.404E-01	-
2	8820	2.403E-02	2.084	8.174E-02	1.598	7.974E-02	1.666
3	28830	1.067E-02	2.057	4.167E-02	1.706	3.947E-02	1.781
4	67240	6.013E-03	2.030	2.562E-02	1.722	2.371E-02	1.805

Name of the solver: BiCG-stab with Jacobi preconditioner (in-house implementation).

• **Test 3 Flow on random meshes,  $u(x, y, z) = \sin(2\pi x) \sin(2\pi y) \sin(2\pi z)$ , min = 0, max = 1, Random meshes**

i	nu	nmat	umin	uemin	umax	uemax	normg
1	64	1000	-7.56E-01	-7.11E-01	0.711	0.525	1.650
2	512	10648	-9.39E-01	-8.32E-01	0.926	0.933	2.674
3	4096	97336	-9.86E-01	-9.77E-01	0.982	0.978	3.330
4	32768	830584	-9.96E-01	-9.92E-01	0.996	0.990	3.527

i	nu	erl2	ratio12	ergrad	ratiograd	ener	ratioener
1	64	5.548E-01	-	7.060E-01	-	7.651E-01	-
2	512	1.427E-01	1.958	3.067E-01	1.202	3.264E-01	1.229
3	4096	2.967E-02	2.266	9.532E-02	1.685	9.569E-02	1.770
4	32768	7.166E-03	2.049	3.253E-02	1.551	2.529E-02	1.919

Name of the solver: BiCG-stab with Jacobi preconditioner (in-house implementation).

• **Test 4 Flow around a well, Well meshes, min = 0, max = 5.415**

i	nu	nmat	umin	uemin	umax	uemax	normg
1	890	18876	4.57E-01	5.26E-01	5.317	5.318	1573.020
2	2232	51800	2.61E-01	2.89E-01	5.329	5.329	1600.780
3	5016	121584	1.62E-01	1.73E-01	5.329	5.329	1613.840
4	11220	280868	1.23E-01	1.29E-01	5.330	5.330	1619.520
5	23210	592448	9.28E-02	9.66E-02	5.339	5.339	1620.960
6	42633	1100865	7.42E-02	7.67E-02	5.345	5.345	1621.200
7	74679	1942619	5.75E-02	5.91E-02	5.361	5.361	1621.930

i	nu	erl2	ratio12	ergrad	ratiograd	ener	ratioener
1	890	9.562E-03	-	8.767E-02	-	5.372E-02	-
2	2232	3.699E-03	3.098	3.903E-02	2.640	2.305E-02	2.761
3	5016	1.676E-03	2.932	1.916E-02	2.636	1.104E-02	2.727
4	11220	1.190E-03	1.275	1.270E-02	1.531	7.205E-03	1.589
5	23210	7.545E-04	1.882	8.919E-03	1.458	5.053E-03	1.463
6	42633	4.601E-04	2.439	6.148E-03	1.835	3.522E-03	1.781
7	74679	3.402E-04	1.616	5.400E-03	0.693	3.174E-03	0.556

Name of the solver: BiCG-stab with Jacobi preconditioner (in-house implementation).

- **Test 5 Discontinuous permeability**,  $u(x, y, z) = \sin(\pi x) \sin(\pi y) \sin(\pi z)$ ,  $\min = 0$ ,  $\max = 1$ , **Locally refined meshes**

i	nu	nmat	umin	uemin	umax	uemax	normg
1	22	252	-1.00E+02	-5.24E+01	100.000	52.359	58.097
2	176	3220	-3.54E+01	-2.62E+01	35.355	26.180	43.055
3	1408	31524	-7.89E+01	-7.30E+01	78.858	73.021	76.757
4	11264	277396	-9.43E+01	-9.25E+01	94.346	92.545	92.422
5	90112	2324532	-9.86E+01	-9.81E+01	98.562	98.089	97.247

i	nu	erl2	ratio12	ergrad	ratiograd	ener	ratioener
1	22	9.831E-01	-	7.196E-01	-	3.176E+02	-
2	176	5.072E-01	0.954	7.376E-01	-0.035	8.184E-01	8.600
3	1408	1.376E-01	1.882	3.770E-01	0.968	6.058E-01	0.433
4	11264	3.347E-02	2.039	1.874E-01	1.008	4.685E-01	0.370
5	90112	9.731E-03	1.782	1.159E-01	0.693	3.448E-01	0.442

Name of the solver: BiCG-stab with Jacobi preconditioner (in-house implementation).

### 3 Comments

This finite volume method is truly cell-centered and, for this reason, it has a relatively small number of degrees of freedom with respect to other finite volume discretizations which introduce face unknowns to approximate the scalar variable. The coercivity was proved only for simple cases (see [3] for details), so very few can be said from a theoretical standpoint about the convergence properties of this scheme and the literature misses a general convergence analysis. Despite this fact, the resulting finite volume method generally show second order of accuracy in all numerical experiments where the exact solution is enough regular and on

**Table 1** LS-FVM method, test 1 using Kershaw mesh with grid resolution  $32 \times 32 \times 32$ ; CPU times are measured in seconds.

solver	precond	CPU time	# iters	Rel. resid.
UMFPACK	none	28.712	0	7.287e-15
ISTL-BiCGstab	Jacobi	36.048	2990	1.126e-10
ISTL-GMRES	Jacobi	64.775	9186	3.931e-10
ISTL-BiCGstab	none	78.880	11417	3.431e-10
ISTL-BiCGstab	ILU(4)	1888.48	2	1.191e-12
ISTL-GMRES	ILU(4)	1692.49	4	9.103e-10

“reasonable” meshes. Moreover, it can be easily applied to complex, distorted meshes and anisotropic permeabilities for which it provides a reliable numerical approximation. It is also generally robust even if a locking phenomenon for the convergence has been reported in the literature [6].

The linear system for the cell-centered unknowns that is originated by this scheme on a general polyhedral mesh leads to an *asymmetric sparse matrix*. Therefore, this system can be solved efficiently by standard preconditioned Krylov methods (BiCG-stab or GMRES) or by direct solvers for general asymmetric systems (UMFPACK). An example of a typical behavior is reported in Table 1 for a subset of the combinations solvers and preconditioners available on the benchmark site. The comparison among these results reveals that the BiCG-stab solver using a diagonal Jacobi preconditioner seems to be the more efficient choice in most of the cases. The performance is usually comparable with that offered by the direct solver (UMFPACK), but the memory storage required by this latter may be from 2 to 60 times greater.

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The paper is in final form and no similar paper has been or is being submitted elsewhere.