Coordinate and Time Systems

GPS satellites are orbiting around the Earth with time. GPS surveys are made mostly on the Earth. To describe the GPS observation (distance) as a function of the GPS orbit (satellite position) and the measuring position (station location), suitable coordinate and time systems have to be defined.

2.1 Geocentric Earth-Fixed Coordinate Systems

It is convenient to use the Earth-Centred Earth-Fixed (ECEF) coordinate system to describe the location of a station on the Earth's surface. The ECEF coordinate system is a right-handed Cartesian system (x, y, z). Its origin and the Earth's centre of mass coincide, while its *z*-axis and the mean rotational axis of the Earth coincide; the *x*-axis is pointing to the mean Greenwich meridian, while the *y*-axis is directed to complete a right-handed system (cf., Fig. 2.1). In other words, the *z*-axis is pointing to a mean pole of the Earth's rotation. Such a mean pole, defined by international convention, is called the Conventional International Origin (CIO). Then the *xy*-plane is called mean equatorial plane, and the *xz*-plane is called mean zero-meridian.





The ECEF coordinate system is also known as the Conventional Terrestrial System (CTS). The mean rotational axis and mean zero-meridian used here are necessary. The true rotational axis of the Earth changes its direction with respect to the Earth's body all the time. If such a pole would be used to define a coordinate system, then the coordinates of the station would also change all the time. Because the surveying is made in our true world, so it is obvious that the polar motion has to be taken into account and will be discussed later.

The ECEF coordinate system can, of course, be represented by a spherical coordinate system (r, ϕ, λ) , where *r* is the radius of the point (x, y, z), ϕ and λ are the geocentric latitude and longitude, respectively (cf., Fig. 2.2). λ is counted eastward from the zero-meridian. The relationship between (x, y, z) and (r, ϕ, λ) is obvious:

$$\begin{pmatrix} x \\ y \\ z \end{pmatrix} = \begin{pmatrix} r \cos \phi \cos \lambda \\ r \cos \phi \sin \lambda \\ r \sin \phi \end{pmatrix}, \text{ or } \begin{cases} r = \sqrt{x^2 + y^2 + z^2} \\ \tan \lambda = y / x \\ \tan \phi = z / \sqrt{x^2 + y^2} \end{cases}.$$
 (2.1)

An ellipsoidal coordinate system (φ , λ , h) may be also defined based on the ECEF coordinates; however, geometrically, two additional parameters are needed to define the shape of the ellipsoid (cf., Fig. 2.3). φ , λ and h are geodetic latitude, longitude and height, respectively. The ellipsoidal surface is a rotational ellipse. The ellipsoidal system is also called the geodetic coordinate system. Geocentric longitude and geodetic longitude are identical. The two geometric parameters could be the semi-major radius (denote by a) and the semi-minor radius (denote by b) of the rotating ellipse, or the semi-major radius and the flattening (denote by f) of the ellipsoid. They are equivalent sets of parameters. The relationship between (x, y, z) and (φ , λ , h) is (cf., e.g., Torge 1991):

$$\begin{pmatrix} x \\ y \\ z \end{pmatrix} = \begin{pmatrix} (N+h)\cos\varphi\cos\lambda \\ (N+h)\cos\varphi\sin\lambda \\ (N(1-e^2)+h)\sin\varphi \end{pmatrix},$$
(2.2)

or

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where

$$N = \frac{a}{\sqrt{1 - e^2 \sin^2 \varphi}} \quad . \tag{2.4}$$

N is the radius of curvature in the prime vertical, and *e* is the first eccentricities. The geometric meaning of *N* is shown in Fig. 2.4. In Eq. 2.3, the φ and *h* have to be solved by iteration; however, the iteration process converges quickly, since $h \ll N$. The flattening and the first eccentricities are defined as:

$$f = \frac{a-b}{a}$$
, and $e = \frac{\sqrt{a^2 - b^2}}{a}$. (2.5)

In cases where $\varphi = \pm 90^{\circ}$ or *h* is very large, the iteration formulas of Eq. 2.3 could be instable. Alternatively, using (cf., Lelgemann 2002)

$$\begin{aligned} & \operatorname{ctan} \varphi = \frac{\sqrt{x^2 + y^2}}{z + \Delta z} , \\ & \Delta z = e^2 N \sin \varphi = \frac{a e^2 \sin \varphi}{\sqrt{1 - e^2 \sin^2 \varphi}} , \end{aligned}$$





may lead to a stably iterated result of φ . Δz and e^2N are the lengths of \overline{OB} and \overline{AB} (cf., Fig. 2.4) respectively. *h* can be obtained by using Δz , i.e.,

$$h = \sqrt{x^2 + y^2 + (z + \Delta z)^2} - N$$

The two geometric parameters used in the World Geodetic System 1984 (WGS-84) are (a = 6378137 m, f = 1/298.2572236). In International Terrestrial Reference Frame 1996 (ITRF-96), the two parameters are (a = 6378136.49 m, f = 1/298.25645). ITRF uses the International Earth Rotation Service (IERS) Conventions (cf., McCarthy 1996). In PZ-90 (Parameters of the Earth Year 1990) coordinate system of GLONASS, the two parameters are (a = 6378136 m, f = 1/298.2578393).

The relation between the geocentric and geodetic latitude ϕ and ϕ may be given by (cf., Eqs. 2.1 and 2.3):

$$\tan\phi = \left(1 - e^2 \frac{N}{N+h}\right) \tan\varphi \quad . \tag{2.6}$$

2.2 Coordinate System Transformations

Any Cartesian coordinate system can be transformed to another Cartesian coordinate system through three succeeded rotations if their origins are the same and if they are both right-handed or left-handed coordinate systems. These three rotational matrices are:

$$R_{1}(\alpha) = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos \alpha & \sin \alpha \\ 0 & -\sin \alpha & \cos \alpha \end{pmatrix},$$

$$R_{2}(\alpha) = \begin{pmatrix} \cos \alpha & 0 & -\sin \alpha \\ 0 & 1 & 0 \\ \sin \alpha & 0 & \cos \alpha \end{pmatrix},$$

$$R_{3}(\alpha) = \begin{pmatrix} \cos \alpha & \sin \alpha & 0 \\ -\sin \alpha & \cos \alpha & 0 \\ 0 & 0 & 1 \end{pmatrix},$$
(2.7)

where α is the rotating angle, which has a positive sign for a counter-clockwise rotation as viewed from the positive axis to the origin. R_1 , R_2 , and R_3 are called the rotating matrix around the *x*, *y*, and *z*-axis, respectively. For any rotational matrix *R*, there are $R^{-1}(\alpha) = R^T(\alpha)$ and $R^{-1}(\alpha) = R(-\alpha)$; that is, the rotational matrix is an orthogonal one, where R^{-1} and R^T are the inverse and transpose of the matrix *R*.

For two Cartesian coordinate systems with different origins and different length units, the general transformation can be given in vector (matrix) form as

$$X_{n} = X_{0} + \mu R X_{old} , \text{ or}$$

$$\begin{pmatrix} x_{n} \\ y_{n} \\ z_{n} \end{pmatrix} = \begin{pmatrix} x_{0} \\ y_{0} \\ z_{0} \end{pmatrix} + \mu R \begin{pmatrix} x_{old} \\ y_{old} \\ z_{old} \end{pmatrix} ,$$
(2.8)

where μ is the scale factor (or the ratio of the two length units), and *R* is a transformation matrix that can be formed by three suitably succeeded rotations. x_n and x_{old} denote the new and old coordinates, respectively; x_0 denotes the translation vector and is the coordinate vector of the origin of the old coordinate system in the new one.

If rotational angle α is very small, then one has $\sin \alpha \approx \alpha$ and $\cos \alpha \approx 0$. In such a case, the rotational matrix can be simplified. If the three rotational angles $\alpha_1, \alpha_2, \alpha_3$ in *R* of Eq. 2.8 are very small, then *R* can be written as (cf., e.g., Lelgemann and Xu 1991):

$$R = \begin{pmatrix} 1 & \alpha_3 & -\alpha_2 \\ -\alpha_3 & 1 & \alpha_1 \\ \alpha_2 & -\alpha_1 & 1 \end{pmatrix},$$
 (2.9)

where α_1 , α_2 , α_3 are small rotating angles around the *x*, *y* and *z*-axis, respectively. Using the simplified *R*, the transformation 2.8 is called the Helmert transformation.

As an example, the transformation from WGS-84 to ITRF-90 is given by (McCarthy 1996):

x _{ITRF-90}		0.060		1	-0.0070''	-0.0003"	$\left(x_{WGS-84}\right)$	
$y_{\rm ITRF-90}$	=	-0.517	$+\mu$	0.0070"	1	-0.0183"	$y_{\rm WGS-84}$	
z _{ITRF-90}		-0.223)	0.0003″	0.0183″	1)	$z_{\rm WGS-84}$)

where $\mu = 0.999999989$, the translation vector has the unit of meter.

The transformations between the coordinate systems of GPS, GLONASS and Galileo can be generally represented by Eq. 2.8 with the scale factor $\mu = 1$ (i.e., the length units used in the three systems are the same). A formula of velocity transformations between different coordinate systems can be obtained by differentiating the Eq. 2.8 with respect to the time.

2.3 Local Coordinate System

The local left-handed Cartesian coordinate system (x', y', z') can be defined by placing the origin to the local point $P_1(x_1, y_1, z_1)$, whose z'-axis is pointed to the vertical, x'-axis is directed to the north, and y' is pointed to the east (cf., Fig. 2.5). The x'y'-plane is called the horizontal plane; the vertical is defined perpendicular to the ellipsoid.



Such a coordinate system is also called a local horizontal coordinate system. For any point P_2 , whose coordinates in the global and local coordinate system are (x_2, y_2, z_2) and (x', y', z'), respectively, one has relations of

$$\begin{pmatrix} x' \\ y' \\ z' \end{pmatrix} = d \begin{pmatrix} \cos A \sin Z \\ \sin A \sin Z \\ \cos Z \end{pmatrix}, \text{ and } \begin{pmatrix} d = \sqrt{x'^2 + {y'}^2 + {z'}^2} \\ \tan A = {y'}/{x'} \\ \cos Z = {z'}/{d} \end{pmatrix},$$
(2.10)

where A is the azimuth, Z is the zenith distance and d is the radius of the P_2 in the local system. A is measured from the north clockwise; Z is the angle between the vertical and the radius d.

The local coordinate system (x', y', z') can indeed be obtained by two succeeded rotations of the global coordinate system (x, y, z) by $R_2(90^\circ - \phi)R_3(\lambda)$ and then by changing the *x*-axis to a right-handed system. In other words, the global system has to be rotated around the *z*-axis with angle λ , then around the *y*-axis with angle $90^\circ - \phi$, and then change the sign of the *x*-axis. The total transformation matrix *R* is then

$$R = \begin{pmatrix} -\sin\varphi\cos\lambda & -\sin\varphi\sin\lambda & \cos\varphi \\ -\sin\lambda & \cos\lambda & 0 \\ \cos\varphi\cos\lambda & \cos\varphi\sin\lambda & \sin\varphi \end{pmatrix},$$
(2.11)

and there are:

$$X_{\text{local}} = RX_{\text{global}} \text{ and } X_{\text{global}} = R^T X_{\text{local}}$$
, (2.12)

where X_{local} and X_{global} are the same vector represented in local and global coordinate systems. (φ , λ) are the geodetic latitude and longitude of the local point.

If the vertical direction is defined as the plump line of the gravitational field at the local point, then such a local coordinate system is called an astronomic horizontal system (its *x*'-axis is pointed to the north, left-handed system). The plump line of gravity *g*

and the vertical line of the ellipsoid at the point *p* are generally not coinciding with each other; however, the difference is very small. The difference is omitted in GPS practice.

Combining Eqs. 2.10 and 2.12, the zenith angle and azimuth of a point P_2 (satellite) related to the station P_1 can be directly computed by using the global coordinates of the two points by

$$\cos Z = \frac{z'}{d} \quad \text{and} \quad \tan A = \frac{y'}{x'}, \tag{2.13}$$

where

$$\begin{split} d &= \sqrt{(x_2 - x_1)^2 + (y_2 - y_1)^2 + (z_2 - z_1)^2} ,\\ x' &= -(x_2 - x_1)\sin\varphi\cos\lambda - (y_2 - y_1)\sin\varphi\sin\lambda + (z_2 - z_1)\cos\varphi ,\\ y' &= -(x_2 - x_1)\sin\lambda + (y_2 - y_1)\cos\lambda \text{ and}\\ z' &= (x_2 - x_1)\cos\varphi\cos\lambda + (y_2 - y_1)\cos\varphi\sin\lambda + (z_2 - z_1)\sin\varphi . \end{split}$$

2.4 Earth-Centred Inertial Coordinate System

To describe the motion of the GPS satellites, an inertial coordinate system has to be defined. The motion of the satellites follows the Newtonian mechanics, and the Newtonian mechanics is valid and expressed in an inertial coordinate system. For reasons, the Conventional Celestial Reference Frame (CRF) is suitable for our purpose. The *xy*-plane of the CRF is the plane of the Earth's equator; the coordinates are celestial longitude, measured eastward along the equator from the vernal equinox, and celestial latitude. The vernal equinox is a crossover point of the ecliptic and the equator. So the right-handed Earth-centred inertial (ECI) system uses the Earth centre as the origin, CIO (Conventional International Origin) as the *z*-axis, and its *x*-axis is directed to the equinox of J2000.0 (Julian Date of 12^{h} 1^{st} January 2000). Such a coordinate system is also called equatorial coordinates of date. Because of the motion (acceleration) of the Earth's centre, ECI is indeed a quasi-inertial system, and the general relativistic effects have to be taken into account in this system. The system moves around the Sun, however, without rotating with respect to the CIO. This system is also called the Earth-centred space-fixed (ECSF) coordinate system.

An excellent figure has been given by Torge (1991) to illustrate the motion of the Earth's pole with respect to the ecliptic pole (cf., Fig. 2.6). The Earth's flattening, combined with the obliquity of the ecliptic, results in a slow turning of the equator on the ecliptic due to the differential gravitational effect of the Moon and the Sun. The slow circular motion with a period of about 26 000 years is called precession, and the other quicker motion with periods from 14 days to 18.6 years is called nutation. Taking the precession and nutation into account, the Earth's mean pole (related to the mean equator) is transformed to the Earth's true pole (related to the true equator). The *x*-axis of the ECI is pointed to the vernal equinox of date.

The angle of the Earth's rotation from the equinox of date to the Greenwich meridian is called Greenwich Apparent Sidereal Time (GAST). Taking GAST into account (called the Earth's rotation), the ECI of date is transformed to the true equatorial co-



ordinate system. The difference between the true equatorial system and the ECEF system is the polar motion. So we have transformed the ECI system with a geometric way to the ECEF system. Such a transformation process can be written as

$$X_{\rm ECEF} = R_{\rm M} R_{\rm S} R_{\rm N} R_{\rm P} X_{\rm ECI} \quad , \tag{2.14}$$

where $R_{\rm p}$ is the precession matrix, $R_{\rm N}$ is the nutation matrix, $R_{\rm S}$ is the Earth rotation matrix, $R_{\rm M}$ is the polar motion matrix, X is the coordinate vector, and indices ECEF and ECI denote the related coordinate systems.

Precession

The precession matrix consists of three succeeded rotational matrices, i.e., (cf., e.g., Hofman-Wellenhof et al. 1997; Leick 1995; McCarthy 1996)

$$R_{\rm P} = R_3(-z)R_2(\theta)R_3(-\zeta)$$

$$= \begin{pmatrix} \cos z \cos \theta \cos \zeta - \sin z \sin \zeta & -\cos z \cos \theta \sin \zeta - \sin z \cos \zeta & -\cos z \sin \theta \\ \sin z \cos \theta \cos \zeta + \cos z \sin \zeta & -\sin z \cos \theta \sin \zeta + \cos z \cos \zeta & -\sin z \sin \theta \\ \sin \theta \cos \zeta & -\sin \theta \sin \zeta & \cos \theta \end{pmatrix}, (2.15)$$

where z, θ , ζ are precession parameters and

$$z = 2306."2181T + 1."09468T^{2} + 0."018203T^{3},$$

$$\theta = 2004."3109T - 0."42665T^{2} - 0."041833T^{3} \text{ and}$$

$$\zeta = 2306."2181T + 0."30188T^{2} + 0."017998T^{3},$$

(2.16)

where *T* is the measuring time in Julian centuries (36 525 days) counted from J2000.0 (cf., Sect. 2.6 time systems).

Nutation

The nutation matrix consists of three succeeded rotational matrices, i.e., (cf., e.g., Hoffman-Wellenhof et al. 1997; Leick 1995; McCarthy 1996)

$$R_{N} = R_{1}(-\varepsilon - \Delta\varepsilon)R_{3}(-\Delta\psi)R_{1}(\varepsilon)$$

$$= \begin{pmatrix} \cos\Delta\psi & -\sin\Delta\psi\cos\varepsilon & -\sin\Delta\psi\sin\varepsilon \\ \sin\Delta\psi\cos\varepsilon_{t} & \cos\Delta\psi\cos\varepsilon_{t}\cos\varepsilon + \sin\varepsilon_{t}\sin\varepsilon & \cos\Delta\psi\sin\varepsilon_{t}\sin\varepsilon \\ \sin\Delta\psi\sin\varepsilon_{t} & \cos\Delta\psi\sin\varepsilon_{t}\cos\varepsilon - \cos\varepsilon_{t}\sin\varepsilon & \cos\Delta\psi\sin\varepsilon_{t}\sin\varepsilon + \cos\varepsilon_{t}\cos\varepsilon \end{pmatrix}$$

$$\approx \begin{pmatrix} 1 & -\Delta\psi\cos\varepsilon & -\Delta\psi\sin\varepsilon \\ \Delta\psi\cos\varepsilon_{t} & 1 & -\Delta\varepsilon \\ \Delta\psi\sin\varepsilon_{t} & \Delta\varepsilon & 1 \end{pmatrix}, \qquad (2.17)$$

where ε is the mean obliquity of the ecliptic angle of date, $\Delta \psi$ and $\Delta \varepsilon$ are nutation angles in longitude and obliquity, $\varepsilon_t = \varepsilon + \Delta \varepsilon$, and

$$\varepsilon = 84381.''448 - 46.''8150T - 0.''00059T^2 + 0.''001813T^3 .$$
(2.18)

The approximation is made by letting $\cos \Delta \psi = 1$ and $\sin \Delta \psi = \Delta \psi$ for very small $\Delta \psi$. For precise purposes, the exact rotation matrix shall be used. The nutation parameters $\Delta \psi$ and $\Delta \varepsilon$ can be computed by using the International Astronomical Union (IAU) theory or IERS theory:

$$\Delta \Psi = \sum_{i=1}^{106} (A_i + A_i'T) \sineta$$
, $\Delta arepsilon = \sum_{i=1}^{106} (B_i + B_i'T) \coseta$

or

$$\Delta \Psi = \sum_{i=1}^{263} (A_i + A_i'T) \sin \beta + A_i' \cos \beta ,$$

$$\Delta \varepsilon = \sum_{i=1}^{263} (B_i + B_i'T) \cos \beta + B_i' \cos \beta ,$$

where argument

$$\beta = N_{1i}l + N_{2i}l' + N_{3i}F + N_{4i}D + N_{5i}\Omega$$
 ,

where *l* is the mean anomaly of the Moon, *l'* is the mean anomaly of the Sun, $F = L - \Omega$, *D* is the mean elongation of the Moon from the Sun, Ω is the mean longitude of the ascending node of the Moon, and *L* is the mean longitude of the Moon. The formulas of *l*, *l'*, *F*, *D*, and Ω , are given in Sect. 11.2.8. The coefficient values of N_{1i} , N_{2i} , N_{3i} , N_{4i} ,

 N_{5i} , A_i , B_i , A'_i , B'_i , A''_i , and B''_i can be found in, e.g., McCarthy (1996). The updated formulas and tables can be found in updated IERS conventions. For convenience, the coefficients of the IAU 1980 nutation model are given in Appendix 1.

Earth Rotation

The Earth rotation matrix can be represented as

$$R_{\rm S} = R_3({\rm GAST}) , \qquad (2.19)$$

where GAST is Greenwich Apparent Sidereal Time and

$$GAST = GMST + \Delta \Psi \cos \varepsilon + 0.''00264 \sin \Omega + 0.''000063 \sin 2\Omega , \qquad (2.20)$$

where GMST is Greenwich Mean Sidereal Time. Ω is the mean longitude of the ascending node of the Moon; the second term on the right-hand side is the nutation of the equinox. Furthermore,

GMST = GMST₀ +
$$\alpha$$
UT1 , (2.21)
GMST₀ = 6 × 3 600."0 + 41 × 60."0 + 50."54841
+ 8 640 184 ."812866 T_0 + 0."093104 T_0^2 - 6."2 × 10⁻⁶ T_0^3 ,
 α = 1.0027379093 50795 + 5.9006 × 10⁻¹¹ T_0 - 5.9 × 10⁻¹⁵ T_0^2 ,

where GMST_0 is Greenwich Mean Sidereal Time at midnight on the day of interest. α is the rate of change. UT1 is the polar motion corrected Universal Time (cf., Sect. 2.6). T_0 is the measuring time in Julian centuries (36525 days) counted from J2000.0 to 0^hUT1 of the measuring day. By computing GMST, UT1 is used (cf., Sect. 2.6).



Polar Motion

As shown in Fig. 2.7, the polar motion is defined as the angles between the pole of date and the CIO pole. The polar motion coordinate system is defined by *xy*-plane coordinates, whose *x*-axis is pointed to the south and is coincided to the mean Greenwich meridian, and whose *y*-axis is pointed to the west. x_p and y_p are the angles of the pole of date, so the rotation matrix of polar motion can be represented as

$$R_{\rm M} = R_2(-x_{\rm p})R_{\rm I}(-y_{\rm p}) = \begin{pmatrix} \cos x_p & \sin x_p \sin y_p & \sin x_p \cos y_p \\ 0 & \cos y_p & -\sin y_p \\ -\sin x_p & \cos x_p \sin y_p & \cos x_p \cos y_p \end{pmatrix}$$

$$\approx \begin{pmatrix} 1 & 0 & x_p \\ 0 & 1 & -y_p \\ -x_p & y_p & 1 \end{pmatrix}$$
(2.22)

The IERS determined x_p and y_p can be obtained from the home pages of IERS.

2.5 Geocentric Ecliptic Inertial Coordinate System

As discussed above, ECI used the CIO pole in the space as the *z*-axis (through consideration of the polar motion, nutation and precession). If the ecliptic pole is used as the *z*-axis, then an ecliptic coordinate system is defined, and it may be called the Earth Centred Ecliptic Inertial (ECEI) coordinate system. ECEI places the origin at the mass centre of the Earth, its *z*-axis is directed to ecliptic pole (or, *xy*-plane is the mean ecliptic), and its *x*-axis is pointed to the vernal equinox of date. The coordinate transformation between the ECI and ECEI systems can be represented as

 $X_{\rm ECEI} = R_1(-\varepsilon)X_{\rm ECI}$,

where ε is the ecliptic angle (mean obliquity) of the ecliptic plane related to the equatorial plane. The formula for ε is given in Sect. 2.4. Usually, coordinates of the Sun and the Moon as well as planets are given in the ECEI system.

2.6 Time Systems

Three time systems are used in satellite surveying. They are sidereal time, dynamic time and atomic time (cf., e.g., Hofman-Wellenhof et al. 1997; Leick 1995; McCarthy 1996; King et al. 1987).

Sidereal time is a measure of the Earth's rotation and is defined as the hour angle of the vernal equinox. If the measure is counted from the Greenwich meridian, the sidereal time is called Greenwich Sidereal Time. Universal Time (UT) is the Greenwich hour angle of the apparent Sun, which is orbiting uniformly in the equatorial plane. Because the angular velocity of the Earth's rotation is not a constant, sidereal time is not a uniformly-scaled time. The oscillation of UT is also partly caused by the polar motion of the Earth. The universal time corrected for the polar motion is denoted by UT1.

Dynamical time is a uniformly-scaled time used to describe the motion of bodies in a gravitational field. Barycentric Dynamic Time (TDB) is applied in an inertial coordinate system (its origin is located at the centre-of-mass (Barycentre)). Terrestrial Dynamic Time (TDT) is used in a quasi-inertial coordinate system (such as ECI). Because of the motion of the Earth around the Sun (or say, in the Sun's gravitational field), TDT will have a variation with respect to TDB. However, both the satellite and the Earth are subject to almost the same gravitational perturbations. TDT may be used for describing the satellite motion without taking into account the influence of the gravitational field of the Sun. TDT is also called Terrestrial Time (TT).

Atomic Time is a time system kept by atomic clocks such as International Atomic Time (TAI). It is a uniformly-scaled time used in the ECEF coordinate system. TDT is realised by TAI in practice with a constant offset (32.184 sec). Because of the slowing down of the Earth's rotation with respect to the Sun, Coordinated Universal Time (UTC) is introduced to keep the synchronisation of TAI to the solar day (by inserting the leap seconds). GPS Time (GPST) is also an atomic time.

The relationships between different time systems are given as follows:

 $TAI = GPST + 19.0 \sec$ $TAI = TDT - 32.184 \sec$ $TAI = UTC + n \sec$ UT1 = UTC + dUT1(2.23)

where dUT1 can be obtained by IERS, (dUT1 < 0.7 sec, cf., Zhu et al. 1996), (dUT1 is also broadcasted with the navigation data), n is the number of leap seconds of date and is inserted into UTC on the 1st of January and 1st of July of the years. The actual n can be found in the IERS report.

Time argument T (Julian centuries) is used in the formulas given in Sect. 2.4. For convenience, T is denoted by TJD, and TJD can be computed from the civil date (Year, Month, Day, and Hour) as follows:

$$JD = INT(365.25Y) + INT(30.6001(M + 1)) + Day + Hour / 24 + 1720981.5 and$$

(2.24)

TJD = JD / 36525,

where

Y =Year - 1, M =Month + 12, if Month ≤ 2 , Y =Year, M =Month, if Month > 2,

where JD is the Julian Date, Hour is the time of UT and INT denotes the integer part of a real number. The Julian Date counted from JD2000.0 is then JD2000 = JD – JD2000.0,

where JD2000.0 is the Julian Date of 2000 January 1st 12^h and has the value of 2451 545.0 days. One Julian century is 36 525 days.

Inversely, the civil date (Year, Month, Day and Hour) can be computed from the Julian Date (JD) as follows:

b = INT(JD + 0.5) + 1537,

c = INT((b - 122.1) / 365.25),

d = INT(365.25c),

e = INT((b - d) / 30.6001),

Hour = JD + 0.5 - INT(JD + 0.5),

Day = b - d - INT(30.6001e),

Month = e - 1 - 12INT(e/14) and

 $Year = c - 4715 - INT((7 + Month) / 10) , \qquad (2.25)$

where *b*, *c*, *d*, and *e* are auxiliary numbers.

Because the GPS standard epoch is defined as $JD = 2\,444\,244.5$ (1980 January 6, 0^h), GPS week and the day of week (denoted by Week and *N*) can be computed by

N = modulo(INT(JD + 1.5), 7) and Week = INT((JD - 2444244.5)/7), (2.26)

where *N* is the day of week (N = 0 for Monday, N = 1 for Tuesday, and so on).

For saving digits and counting the date from midnight instead of noon, the Modified Julian Date (MJD) is defined as

 $MJD = (JD - 2\,400\,000.5)\,. \tag{2.27}$

GLONASS time (GLOT) is defined by Moscow time UTC_{SU} , which equals UTC plus three hours (corresponding to the offset of Moscow time to Greenwich time), theoretically. GLOT is permanently monitored and adjusted by the GLONASS Central Synchroniser (cf. Roßbach 2000). UTC and GLOT then has a simple relation

UTC=GLOT+ τ_c -3h,

where τ_c is the system time correction with respect to UTC_{SU}, which is broadcasted by the GLONASS ephemerides and is less than one microsecond. Therefore there is approximately

GPST = GLOT + m - 3h,

where *m* is called a number of "leap seconds" between GPS and GLONASS (UTC) time and is given in the GLONASS ephemerides. *m* is indeed the leap seconds since GPS standard epoch (1980 January 6, 0^{h}).

Galileo system time (GST) will be maintained by a number of UTC laboratory clocks. GST and GPST are time systems of various UTC laboratories. After the offset of GST and GPST is made available to the user, the interoperability will be ensured.