# <span id="page-0-0"></span>An Undetermined Coefficients Method for a Class of Ordinary Differential Equation with Initial Values

Li Gao $^{(\boxtimes)}$ 

Liberal Arts Experimental Teaching Center, Neijiang Normal University, Neijiang, Sichuan 641112, People's Republic of China 249435911@qq.com

Abstract. In this paper, based on the Taylor series, we present a method of undetermined coefficients to solve a class of ordinary differential equation with initial values. Theoretical analysis and examples show this method can achieve accuracy  $O(h^{m+1})$  where m is the order of Taylor series of the right function  $f(x)$ . Furthermore, compared with traditional methods such as the finite difference method and the finite element method, this method can avoid solving complicated and large linear systems.

Keywords: Ordinary differential equation  $\cdot$  Euler type  $\cdot$  Undetermined coefficient coefficient

## 1 Introduction

For numerically solving ordinary differential equations (ODEs) and partial differential equations (PDEs) in various engineering problems, some typical and traditional methods have been presented so far, including the finite difference method (FDM), the boundary element method (BEM) and so on, too many results about their developments, accuracy, convergence, stability have arisen since they were presented. These can be partly seen in  $[1-11]$  $[1-11]$  $[1-11]$  $[1-11]$  and references therein.

In this paper, based on the Taylor series, we present a new numerical method for the ordinary differential equation

$$
\sum_{k=0}^{n} a_{n-k+1} x^{k} y^{(k)} = x^{n} f(x), x \in (0, 1),
$$
  

$$
y^{(k)}(0) = 0, k = 0, 1, \dots, n-1,
$$
 (1)

where  $a_k$ ,  $k = 1, 2, \dots, n + 1$  are all constants,  $y(x)$  is the function to be determined. This method is quite different from such the existed methods as the FDM. FFM and RFM method is quite different from such the existed methods as the FDM, FEM and BEM.

In the next sections, by some theoretical analyses and experimental demonstrations, we find this method has some evident characteristics:

- It can provide the unique numerical solution of  $(1)$ ;
- It can achieve accuracy  $O(h^{m+1})$  when  $f(x) \in C^{m+1}[0,1]$ ;

<sup>©</sup> Springer International Publishing AG 2018

F. Qiao et al. (eds.), Recent Developments in Mechatronics and Intelligent Robotics, Advances in Intelligent Systems and Computing 690, DOI 10.1007/978-3-319-65978-7\_20

<span id="page-1-0"></span>• It does not lead to complicated and large linear systems.

The remainder of the paper is organized as follows. In Sect. 2, we describe our method in details, including the construction of this method, the existence, uniqueness, the accuracy of the numerical solution and the computational complexity of this method. In Sect. [3](#page-4-0), some examples are given to check this presented method in Sect. 2. Finally, we draw some conclusions about our method, point out some shortages which are expected to be overcome.

## 2 The Method of Undetermined Coefficients

Supposing the equation

$$
y^{(2)}(x) + y'(x) + y(x) = d(x), x = [0, 1]
$$
 (2)

has a unique solution  $y(x) \in C^2[0, 1]$ .

**Case 1:** If the function  $d(x)$  is a given 2 order polynomial, namely,

$$
d(x) = d_0 + d_1 x + d_2 x^2,
$$
\t(3)

we can guess that  $y(x)$  is also a 2 order polynomial of the form

$$
y(x) = t_0 + t_1 x + t_2 x^2.
$$
 (4)

Substituting  $(4)$ ,  $(3)$  into  $(2)$ , we get the equality

$$
(t_0 + t_1 + 2t_2) + (t_1 + 2t_2)x + t_2x^2 = d_0 + d_1x + d_2x^2
$$

for all  $x = [0, 1]$ . Consequently, by comparing the coefficients, one immediately gets

$$
t_2 = d_2, t_1 = d_1 - 2t_2, t_0 = d_0 - t_1 - 2t_2,
$$
\n<sup>(5)</sup>

and thus the solution  $y(x)$  is exactly found.

**Case 2:** When the  $d(x)$  is not a polynomial, we can subdivide the domain  $(0, 1)$  into  $\bigcup_{i=1}^{N} \Omega_i$ . On each sub-domain  $\Omega_i$ ,  $(i = 1, 2, \dots, N)$ , we replace  $d(x)$  with its 2 order<br>Faylor series and the remaining work is repeating the above course from (3) to (5) for Taylor series, and the remaining work is repeating the above course from (3) to (5) for seeking an approximate solution of  $(2)$  in  $\Omega_i$ .

For convenience, we assume that  $y(x) \in C^{m+1}(\overline{\Omega}), f(x) \in C^{m+1}(\overline{\Omega}), m \ge 1$ . For  $\overline{\Omega} = [0, 1]$ , let  $\Delta_x = \{0 = x_0 < x_1 < \cdots < x_M = 1\}$  be uniform partitions of  $\Omega$  with means sizes  $h - h$ . Throughout this paper, we denote by  $\tau, i = 1, 2, \ldots, M$  the midmesh sizes  $h_x = h$ . Throughout this paper, we denote by  $\tau_i$ ,  $i = 1, 2, ..., M$  the midpoints of  $\Delta_x$ , and by  $\Omega_i = (x_{i-1}, x_i)$ .

According to the conditions  $y^{(k)}(0) = 0, k = 0, 1, \ldots, n - 1$ , the solution  $y(x)$  can be expressed as  $y = x^n u(x)$  with  $u(x)$  the new unknown function. Substituting this into [\(1](#page-0-0)), by some computations, we obtain the equivalent system

<span id="page-2-0"></span>130 L. Gao

$$
\sum_{k=0}^{n} b_{k+1} x^{k} u^{(k)} = f(x), x \in (0, 1),
$$
 (6)

where

$$
b_{k+1} = \sum_{t=1}^{n-k+1} a_t C_{n-t+1}^k n! / (t+k-1)!, k = 0, 1, ..., n.
$$

In every sub-interval  $\Omega_i$ , let

$$
\hat{f}(x) = \sum_{k=0}^{m} f^{(k)}(\tau_i) (x - \tau_i)^k / k! \tag{7}
$$

after expanding all the polynomials  $(x - \tau_i)^k$ ,  $k = 0, 1, ..., m$  and combining like terms (7) reads terms, (7) reads

$$
\hat{f}(x) = \sum_{k=0}^{m} g_{k+1} x^{k},
$$
\n(8)

where

$$
g_{k+1} = \sum_{t=k}^{m} f^{(t)}(\tau_i) C_t^{t-k} (-\tau_i)^{t-k} / t!.
$$

Then, for the right function  $f(x)$  in (6), we can easily write its *m* order Taylor series on the point  $\tau_i$  in sub-domain  $\Omega_i$  (i = 1,2,..., M) as

$$
f(x) = \hat{f}(x) + (x - \tau_i)^{m+1} f^{(m+1)}(\theta_i) / (m+1)!
$$
 (9)

With  $\theta_i \in \overline{\Omega}_i$ .

By using the idea similar to ([4\)](#page-1-0), taking approximate solution  $u_h$  as

$$
u_h = \sum_{t=0}^{m} c_t x^t, \qquad (10)
$$

with  $c_t$ ,  $t = 0, 1, \ldots, m$  undetermined coefficients, we get

$$
\sum_{k=0}^{n} b_{k+1} x^{k} u_{h}^{(k)} = \sum_{k=0}^{n} s_{k} x^{k},
$$
\n(11)

<span id="page-3-0"></span>where

$$
s_k = \begin{cases} \sum_{t=0}^k k! b_{t+1} c_k / (k-t)!, & \text{if } k = 0, 1, ..., n, \\ \sum_{t=0}^n k! b_{t+1} c_k / (k-t)!, & \text{if } k = n+1, ..., m. \end{cases}
$$
(12)

Substituting  $(8)$  $(8)$ – $(12)$  into  $(6)$  $(6)$ , and dropping the remainder term

$$
(x - \tau_i)^{m+1} f(\theta_i)^{(m+1)}/(m+1)!
$$

in sub-domain  $\Omega_i$ , we obtain the approximate system

$$
\sum_{k=0}^{n} b_{k+1} x^{k} u_{h}^{(k)} = \hat{f}(x), x \in (0, 1),
$$
\n(13)

namely

$$
\sum_{k=0}^{m} s_k x^k = \sum_{k=0}^{m} g_{k+1} x^k
$$

in  $\Omega_i$ . Let the corresponding coefficients be equal to each other in this equality, then the coefficients  $c_k$ ,  $k = 0,..., m$  can be expressed as:

$$
c_k = \begin{cases} g_{k+1} / \sum_{t=0}^k k! b_{t+1} / (k-t)!, & \text{if } k = 0, 1, ..., n, \\ g_{k+1} / \sum_{t=0}^n k! b_{t+1} / (k-t)!, & \text{if } k = n+1, ..., m. \end{cases}
$$
(14)

From (14), we immediately have the following result about existence and uniqueness for approximate system (13).

**Theorem 1.** Supposing  $\sum_{t=0}^{k} k!b_{t+1}/(k-t)! \neq 0$  when  $k \leq n$ ,  $\sum_{t=0}^{n} k!b_{t+1}/(k-t)!$ <br> $\neq 0$  when  $n+1 \leq k \leq m$  and  $f(x) \in C^{m+1}(\overline{0})$  then the approximate solution  $y$ .  $\neq 0$  when  $n+1 \leq k \leq m$ , and  $f(x) \in C^{m+1}(\overline{\Omega})$ , then the approximate solution  $u_h$ <br>defined by (10) can be uniquely solved by the Eq. (13) defined by  $(10)$  $(10)$  can be uniquely solved by the Eq.  $(13)$ .

Now, we analyze the convergence of this method, we denote by

$$
Lu = \sum_{k=0}^{n} b_{k+1} x^{k} u^{(k)}, u \in \Omega,
$$
\n(15)

then the following theorem is true.

**Theorem 2.** Assuming that Eq. ([1\)](#page-0-0) has unique solution  $y(x)$ , and  $y(x)$ ,  $f(x) \in$  $C^{m+1}(\bar{\Omega})$ . Let  $e \equiv y - y_h$  with  $y_h = x^n u_h$ . Then

$$
|e| \leq C_1 C_2 (h/2)^{m+1}, \tag{16}
$$

<span id="page-4-0"></span>where h is the step length,  $C_1$  is a constant corresponding to the operator  $L^{-1}$ , and

$$
C_2 = \max_{i=1}^{M} f(\theta_i)^{(m+1)} / (m+1)!, \theta_i \in \bar{\Omega}_i, i = 1, 2, ..., M.
$$

Furthermore, this error satisfies

$$
\lim_{|a_n+1|\to+\infty}|e|=0.\tag{17}
$$

**Proof.** In fact, by  $(6)$  $(6)$ ,  $(13)$  $(13)$ , we can get

$$
L(u - u_h) = f - \hat{f},
$$

on each sub-domain  $\Omega_i$ ,  $i = 1, 2, \ldots, M$ .

Because Eq. ([1\)](#page-0-0) has unique solution  $y(x) \in C^{m+1}(\overline{\Omega})$ , by the relation  $y \to x^n u(x)$  we know (1) is equivalent to (6) and the operator *I* is an invertible  $y(x) = x<sup>n</sup>u(x)$ , we know [\(1](#page-0-0)) is equivalent to [\(6\)](#page-2-0), and the operator L is an invertible bounded linear operator on  $C^{m+1}(\overline{\Omega}_i)$ ,  $i = 1, 2,..., M$ , which shows that ([16\)](#page-3-0) is true.

Furthermore, by  $b_{k+1} = \sum_{t=1}^{n-k+1} a_t C_{n-t+1}^k n! / (t+k-1)!$ ,  $k = 0, 1, ..., n$ , we know  $\lim_{a_n+1\to+\infty} b_1 = \infty$ . Combining with ([13\)](#page-3-0), ([6\)](#page-2-0), we have  $u - u_h =$  $\left(\frac{f}{f} - \hat{f}\right)/b_1 - \sum_{k=1}^n b_{k+1} x^k u^{(k)}/b_1$  in  $\overline{\Omega}_i$ ,  $i = 1, 2,..., M$ , which leads to (17), and the proof of Theorem [2](#page-3-0) is completed.

### 3 Numerical Examples

In this section, we give some numerical examples to show the performance of our method. In these examples, we mainly check the result  $(16)$  $(16)$  and  $(17)$  in Theorem [2](#page-3-0): the relation of accuracy with the order of Taylor series of the function  $f(x)$  and  $a_{n+1}$ . We always take the step length  $h = 0.1$ , and in each sub-domain  $\Omega_i$ ,  $i = 1, 2, ..., N$ , we compute Taylor series of  $f(x)$  in the center point  $\tau_i$  of this sub-domain.

In the following tables, for convenience, the notation  $x.y_1y_2 - p$  means  $x.y_1y_2 \times 10^{-p}$ . We test errors in the center and all endpoints of all sub-domains:

- $E_n$  the maximum absolute errors at the centers  $\{\tau_i\}_{i=1}^N$ ;
- $E_v$  the maximum absolute errors at the endpoints  $\{x_i\}_{i=0}^N$ .

The tested equations have respectively the following information:

#### Example 1.

$$
a_1 = a_2 = 1, n = 2, y = x^2 e^x;
$$

a <sub>3</sub>	$m = 2$		$m = 3$		$m = 4$		$m = 5$	
	$E_n$	$E_{\nu}$	$E_n$	$E_v$	$E_n$	$E_{v}$	$E_n$	$E_v$
50					$5.41-3$   $5.69-3$   $5.69-4$   $4.06-4$   $1.56-5$   $1.09-5$   $3.74e-6$   $5.82-6$			
100			$2.00-3$   $1.58-3$   $6.52-5$   $2.58-5$   $8.66-6$				$1.38 - 5 \mid 1.25 - 6 \mid 1.14 - 6$	
500					$1.10 - 4 \times 07 - 5 \times 14.56 - 6 \times 7.38 - 6 \times 2.95 - 7$		$ 2.37-8 6.40-11 1.18-8$	
					$1000$   2.87-5   5.14-5   1.50-6   1.13-6   4.31-8		$ 3.87 - 8 7.23 - 10 1.45 - 9$	
					$5000$   1.18-6   2.68-5   7.26-8   1.66-7   3.91-10   1.61-9   1.19-11   2.47-11			

Table 1. Results of Example 1





#### Example 2.

$$
a_1 = a_2 = a_3 = a_4 = a_5 = 1, n = 5, y = x^5 \sin(0.2x + 1).
$$

From Tables 1 and 2, we can clearly see that the results are in accordance with the theoretical analysis in Sect. [2:](#page-1-0) the method basically achieves accuracy of  $O(h^{m+1})$ , in the same time, just as we expected, the errors of  $u$  and  $u_h$  is inversely proportional to  $|a_{n+1}|$ .

# 4 Conclusions

In this paper, we introduced a new numerical method for solving a class of ordinary differential equation. By giving direct formulas of the undetermined coefficients, we showed this method can avoid solving complicated and large linear systems. Theoretically analysis and numerical experiments demonstrated this method can achieve accuracy  $O(h^{m+1})$  when  $f(x) \in C^{m+1}(\overline{\Omega})$ .

# <span id="page-6-0"></span>References

- 1. Morton, K.W., Mayers, D.F.: Numerical Solution of Partial Differential Equations. Cambridge University Press, Cambridge (2005)
- 2. LeVeque, J.R.: Finite Difference Methods for Ordinary and Partial Differential Equations. Society for Industrial and Applied Mathematics, Philadelphia (2007)
- 3. Fairweathe, G., Karageorghis, A., Maack, J.: Compact optimal quadratic spline collocation methods for the Helmholtz equation. J. Comput. Phys. 230, 2880–2895 (2011)
- 4. Christara, C.C.: Quadratic spline collocation methods for elliptic partial differential equations. BIT 34, 33–61 (1994)
- 5. Abushama, A.A., Bialecki, B.: Modified nodal cubic spline collocation for Poisson's equation. SIAM J. Numer. Anal. 46, 397–418 (2008)
- 6. Knabner, P., Angermann, L.: Numerical Methods for Elliptic and Parabolic Partial Differential Equations. Springer, New York (2003)
- 7. Cecka, C., Darve, E.: Fourier-based fast multipole method for the Helmholtz equation. SIAM J. Sci. Comput. 35, A79–A103 (2013)
- 8. Hewett, D.P., Langdon, S., Melenk, J.M.: A high frequency hp boundary element method for scattering by convex polygons. SIAM J. Numer. Anal. 51, 629–653 (2013)
- 9. Barnett, A.H., Betcke, T.: An exponentially convergent nonpolynomial finite element method for time-harmonic scattering from polygons. SIAM J. Sci. Comput. 32, 1417–1441 (2010)
- 10. Chen, K.: Matrix Preconditioning Techniques and Applications. Cambridge University Press, Cambridge (2005)
- 11. Saad, Y.: Iterative Methods for Sparse Linear Systems. Society for Industrial and Applied Mathematics, Philadelphia (2003)