# **Decomposition and Limit Theorems for a Class of Self-Similar Gaussian Processes**

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**Abstract** We introduce a new class of self-similar Gaussian stochastic processes, where the covariance is defined in terms of a fractional Brownian motion and another Gaussian process. A special case is the solution in time to the fractional-colored stochastic heat equation described in Tudor (Analysis of variations for self-similar processes: a stochastic calculus approach. Springer, Berlin, 2013). We prove that the process can be decomposed into a fractional Brownian motion (with a different parameter than the one that defines the covariance), and a Gaussian process first described in Lei and Nualart (Stat Probab Lett 79:619–624, 2009). The component processes can be expressed as stochastic integrals with respect to the Brownian sheet. We then prove a central limit theorem about the Hermite variations of the process.

**Keywords** Fractional Brownian motion • Hermite variations • Self-similar processes • Stochastic heat equation

AMS 2010 Classification 60F05, 60G18, 60H07

## 1 Introduction

The purpose of this paper is to introduce a new class of Gaussian self-similar stochastic processes related to stochastic partial differential equations, and to establish a decomposition in law and a central limit theorem for the Hermite variations of the increments of such processes.

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F. Baudoin, J. Peterson (eds.), *Stochastic Analysis and Related Topics*, Progress in Probability 72, DOI 10.1007/978-3-319-59671-6\_5

Consider the *d*-dimensional stochastic heat equation

$$\frac{\partial u}{\partial t} = \frac{1}{2}\Delta u + \dot{W}, \ t \ge 0, \ x \in \mathbb{R}^d,$$
(1.1)

with zero initial condition, where  $\dot{W}$  is a zero mean Gaussian field with a covariance of the form

$$\mathbb{E}\left[\dot{W}^{H}(t,x)\dot{W}^{H}(s,y)\right] = \gamma_{0}(t-s)\Lambda(x-y), \quad s,t \ge 0, \ x,y \in \mathbb{R}^{d}.$$

We are interested in the process  $U = \{U_t, t \ge 0\}$ , where  $U_t = u(t, 0)$ .

Suppose that  $\dot{W}$  is white in time, that is,  $\gamma_0 = \delta_0$  and the spatial covariance is the Riesz kernel, that is,  $\Lambda(x) = c_{d,\beta}|x|^{-\beta}$ , with  $\beta < \min(d, 2)$  and  $c_{d,\beta} = \pi^{-d/2} 2^{\beta-d} \Gamma(\beta/2) / \Gamma((d-\beta)/2)$ . Then U has the covariance (see [14])

$$\mathbb{E}[U_t U_s] = D\left((t+s)^{1-\frac{\beta}{2}} - |t-s|^{1-\frac{\beta}{2}}\right), \quad s, t \ge 0,$$
(1.2)

for some constant

$$D = (2\pi)^{-d} (1 - \beta/2)^{-1} \int_{\mathbb{R}^d} e^{-\frac{|\xi|^2}{2}} \frac{d\xi}{|\xi|^{d-\beta}}.$$
 (1.3)

Up to a constant, the covariance (1.2) is the covariance of the *bifractional Brownian* motion with parameters  $H = \frac{1}{2}$  and  $K = 1 - \frac{\beta}{2}$ . We recall that, given constants  $H \in (0, 1)$  and  $K \in (0, 1)$ , the bifractional Brownian motion  $B^{H,K} = \{B_t^{H,K}, t \ge 0\}$ , introduced in [4], is a centered Gaussian process with covariance

$$R_{H,K}(s,t) = \frac{1}{2^K} \left( (t^{2H} + s^{2H})^K - |t-s|^{2HK} \right), \quad s,t \ge 0.$$

When K = 1, the process  $B^H = B^{H,1}$  is simply the fractional Brownian motion (fBm) with Hurst parameter  $H \in (0, 1)$ , with covariance  $R_H(s, t) = R_{H,1}(s, t)$ . In [5], Lei and Nualart obtained the following decomposition in law for the bifractional Brownian motion

$$B^{H,K} = C_1 B^{HK} + C_2 Y^K_{,2H}$$

where  $B^{HK}$  is a fBm with Hurst parameter HK, the process  $Y^{K}$  is given by

$$Y_t^K = \int_0^\infty y^{-\frac{1+K}{2}} (1 - e^{-yt}) dW_y, \qquad (1.4)$$

with  $W = \{W_y, y \ge 0\}$  a standard Brownian motion independent of  $B^{H,K}$ , and  $C_1, C_2$  are constants given by  $C_1 = 2^{\frac{1-K}{2}}$  and  $C_2 = \sqrt{\frac{2^{-K}}{\Gamma(1-K)}}$ . The process  $Y^K$ 

has trajectories which are infinitely differentiable on  $(0, \infty)$  and Hölder continuous of order  $HK - \epsilon$  in any interval [0, T] for any  $\epsilon > 0$ . In particular, this leads to a decomposition in law of the process U with covariance (1.2) as the sum of a fractional Brownian motion with Hurst parameter  $\frac{1}{2} - \frac{\beta}{4}$  plus a regular process.

The classical one-dimensional space-time white noise can also be considered as an extension of the covariance (1.2) if we take  $\beta = 1$ . In this case the covariance corresponds, up to a constant, to that of a bifractional Brownian motion with parameters  $H = K = \frac{1}{2}$ .

The case where the noise term  $\dot{W}$  is a fractional Brownian motion with Hurst parameter  $H \in (\frac{1}{2}, 1)$  in time and a spatial covariance given by the Riesz kernel, that is,

$$\mathbb{E}\left[\dot{W}^{H}(t,x)\dot{W}^{H}(s,y)\right] = \alpha_{H}c_{d,\beta}|s-t|^{2H-2}|x-y|^{-\beta},$$

where  $0 < \beta < \min(d, 2)$  and  $\alpha_H = H(2H - 1)$ , has been considered by Tudor and Xiao in [14]. In this case the corresponding process *U* has the covariance

$$\mathbb{E}[U_t U_s] = D\alpha_H \int_0^t \int_0^s |u - v|^{2H-2} (t + s - u - v)^{-\gamma} du dv,$$
(1.5)

where *D* is given in (1.3) and  $\gamma = \frac{d-\beta}{2}$ . This process is self-similar with parameter  $H - \frac{\gamma}{2}$  and it has been studied in a series of papers [1, 8, 12–14]. In particular, in [14] it is proved that the process *U* can be decomposed into the sum of a scaled fBm with parameter  $H - \frac{\gamma}{2}$ , and a Gaussian process *V* with continuously differentiable trajectories. This decomposition is based on the stochastic heat equation. As a consequence, one can derive the exact uniform and local moduli of continuity and Chung-type laws of the iterated logarithm for this process. In [12], assuming that d = 1, 2 or 3, a central limit theorem is obtained for the renormalized quadratic variation

$$V_n = n^{2H-\gamma-\frac{1}{2}} \sum_{j=0}^{n-1} \left\{ (U_{(j+1)T/n} - U_{jT/n})^2 - \mathbb{E} \left[ (U_{(j+1)t/n} - U_{jT/n})^2 \right] \right\},$$

assuming  $\frac{1}{2} < H < \frac{3}{4}$ , extending well-known results for fBm (see for example [6, Theorem 7.4.1]).

The purpose of this paper is to establish a decomposition in law, similar to that obtained by Lei and Nualart in [5] for the bifractional Brownian motion, and a central limit theorem for the Hermite variations of the increments, for a class of self-similar processes that includes the covariance (1.5). Consider a centered Gaussian process { $X_t$ ,  $t \ge 0$ } with covariance

$$R(s,t) = \mathbb{E}[X_s X_t] = \mathbb{E}\left[\left(\int_0^t Z_{t-r} dB_r^H\right)\left(\int_0^s Z_{s-r} dB_r^H\right)\right],\tag{1.6}$$

where

- (i)  $B^H = \{B_t^H, t \ge 0\}$  is a fBm with Hurst parameter  $H \in (0, 1)$ .
- (ii)  $Z = \{Z_t, t > 0\}$  is a zero-mean Gaussian process, independent of  $B^H$ , with covariance

$$\mathbb{E}[Z_s Z_t] = (s+t)^{-\gamma}, \qquad (1.7)$$

where  $0 < \gamma < 2H$ .

In other words, X is a Gaussian process with the same covariance as the process  $\{\int_0^t Z_{t-r} dB_r^H, t \ge 0\}$ , which is not Gaussian.

When  $H \in (\frac{1}{2}, 1)$ , the covariance (1.6) coincides with (1.5) with D = 1. However, we allow the range of parameters 0 < H < 1 and  $0 < \gamma < 2H$ . In other words, up to a constant, *X* has the law of the solution in time of the stochastic heat equation (1.1), when  $H \in (0, 1)$ ,  $d \ge 1$  and  $\beta = d - 2\gamma$ . Also of interest is that *X* can be constructed as a sum of stochastic integrals with respect to the Brownian sheet (see the proof of Theorem 1).

## 1.1 Decomposition of the Process X

Our first result is the following decomposition in law of the process X as the sum of a fractional Brownian motion with Hurst parameter  $\frac{\alpha}{2} = H - \frac{\gamma}{2}$  plus a process with regular trajectories.

**Theorem 1** The process X has the same law as  $\{\sqrt{\kappa}B_t^{\frac{\alpha}{2}} + Y_t, t \ge 0\}$ , where here and in what follows,  $\alpha = 2H - \gamma$ ,

$$\kappa = \frac{1}{\Gamma(\gamma)} \int_0^\infty \frac{z^{\gamma-1}}{1+z^2} dz, \qquad (1.8)$$

 $B^{\frac{\alpha}{2}}$  is a fBm with Hurst parameter  $\frac{\alpha}{2}$ , and Y (up to a constant) has the same law as the process  $Y^{K}$  defined in (1.4), with  $K = 2\alpha + 1$ , that is, Y is a centered Gaussian process with covariance given by

$$\mathbb{E}\left[Y_t Y_s\right] = \lambda_1 \int_0^\infty y^{-\alpha - 1} (1 - e^{-yt}) (1 - e^{-ys}) \, dy,$$

where

$$\lambda_1 = \frac{4\pi}{\Gamma(\gamma)\Gamma(2H+1)\sin(\pi H)} \int_0^\infty \frac{\eta^{1-2H}}{1+\eta^2} d\eta.$$

The proof of this theorem is given in Sect. 3.

#### **1.2 Hermite Variations of the Process**

For each integer  $q \ge 0$ , the *q*th Hermite polynomial is given by

$$H_q(x) = (-1)^q e^{\frac{x^2}{2}} \frac{d^q}{dx^q} e^{-\frac{x^2}{2}}.$$

See [6, Sect. 1.4] for a discussion of properties of these polynomials. In particular, it is well known that the family  $\{\frac{1}{\sqrt{q!}}H_q, q \ge 0\}$  constitutes an orthonormal basis of the space  $L^2(\mathbb{R}, \gamma)$ , where  $\gamma$  is the N(0, 1) measure.

Suppose  $\{Z_n, n \ge 1\}$  is a stationary, Gaussian sequence, where each  $Z_n$  follows the N(0, 1) distribution with covariance function  $\rho(k) = \mathbb{E}[Z_n Z_{n+k}]$ . If  $\sum_{k=1}^{\infty} |\rho(k)|^q < \infty$ , it is well known that as *n* tends to infinity, the Hermite variation

$$V_n = \frac{1}{\sqrt{n}} \sum_{j=1}^n H_q(Z_j)$$
(1.9)

converges in distribution to a Gaussian random variable with mean zero and variance given by  $\sigma^2 = q! \sum_{k=1}^{\infty} \rho(k)^q$ . This result was proved by Breuer and Major in [3]. In particular, if  $B^H$  is a fBm, then the sequence  $\{Z_{j,n}, 0 \le j \le n-1\}$  defined by

$$Z_{j,n} = n^H \left( B_{\frac{j+1}{n}}^H - B_{\frac{j}{n}}^H \right)$$

is a stationary sequence with unit variance. As a consequence, if  $H < 1 - \frac{1}{q}$ , we have that

$$\frac{1}{\sqrt{n}}\sum_{j=0}^{n-1}H_q\left(n^H\left(B_{\frac{j+1}{n}}^H-B_{\frac{j}{n}}^H\right)\right)$$

converges to a normal law with variance given by

$$\sigma_q^2 = \frac{q!}{2^q} \sum_{m \in \mathbb{Z}} \left( |m+1|^{2H} - 2|m|^{2H} + |m-1|^{2H} \right)^q.$$
(1.10)

See [3] and Theorem 7.4.1 of [6].

The above Breuer-Major theorem can not be applied to our process because *X* is not necessarily stationary. However, we have a comparable result.

**Theorem 2** Let  $q \ge 2$  be an integer and fix a real T > 0. Suppose that  $\alpha < 2 - \frac{1}{q}$ , where  $\alpha$  is defined in Theorem 1. For  $t \in [0, T]$ , define,

$$F_n(t) = n^{-\frac{1}{2}} \sum_{j=0}^{\lfloor nt \rfloor - 1} H_q \left( \frac{\Delta X_{\frac{j}{n}}}{\left\| \Delta X_{\frac{j}{n}} \right\|_{L^2(\Omega)}} \right),$$

where  $H_q(x)$  denotes the qth Hermite polynomial. Then as  $n \to \infty$ , the stochastic process  $\{F_n(t), t \in [0, T]\}$  converges in law in the Skorohod space D([0, T]), to a scaled Brownian motion  $\{\sigma B_t, t \in [0, T]\}$ , where  $\{B_t, t \in [0, T]\}$  is a standard Brownian motion and  $\sigma = \sqrt{\sigma^2}$  is given by

$$\sigma^{2} = \frac{q!}{2^{q}} \sum_{m \in \mathbb{Z}} \left( |m+1|^{\alpha} - 2|m|^{\alpha} + |m-1|^{\alpha} \right)^{q}.$$
(1.11)

The proof of this theorem is given in Sect. 4.

#### 2 Preliminaries

#### 2.1 Analysis on the Wiener Space

The reader may refer to [6, 7] for a detailed coverage of this topic. Let  $Z = \{Z(h), h \in \mathcal{H}\}$  be an *isonormal Gaussian process* on a probability space  $(\Omega, \mathcal{F}, P)$ , indexed by a real separable Hilbert space  $\mathcal{H}$ . This means that Z is a family of Gaussian random variables such that  $\mathbb{E}[Z(h)] = 0$  and  $\mathbb{E}[Z(h)Z(g)] = \langle h, g \rangle_{\mathcal{H}}$  for all  $h, g \in \mathcal{H}$ .

For integers  $q \ge 1$ , let  $\mathcal{H}^{\otimes q}$  denote the *q*th tensor product of  $\mathcal{H}$ , and  $\mathcal{H}^{\odot q}$  denote the subspace of symmetric elements of  $\mathcal{H}^{\otimes q}$ .

Let  $\{e_n, n \ge 1\}$  be a complete orthonormal system in  $\mathcal{H}$ . For elements  $f, g \in \mathcal{H}^{\odot q}$ and  $p \in \{0, \ldots, q\}$ , we define the *p*th-order contraction of f and g as that element of  $\mathcal{H}^{\otimes 2(q-p)}$  given by

$$f \otimes_p g = \sum_{i_1,\dots,i_p=1}^{\infty} \langle f, e_{i_1} \otimes \dots \otimes e_{i_p} \rangle_{\mathcal{H}^{\otimes p}} \otimes \langle g, e_{i_1} \otimes \dots \otimes e_{i_p} \rangle_{\mathcal{H}^{\otimes p}}, \qquad (2.1)$$

where  $f \otimes_0 g = f \otimes g$ . Note that, if  $f, g \in \mathcal{H}^{\odot q}$ , then  $f \otimes_q g = \langle f, g \rangle_{\mathcal{H}^{\odot q}}$ . In particular, if f, g are real-valued functions in  $\mathcal{H}^{\otimes 2} = L^2(\mathbb{R}^2, \mathcal{B}^2, \mu^2)$  for a non-atomic measure  $\mu$ , then we have

$$f \otimes_1 g = \int_{\mathbb{R}} f(s, t_1) g(s, t_2) \ \mu(ds).$$

$$(2.2)$$

Let  $\mathcal{H}_q$  be the *q*th Wiener chaos of *Z*, that is, the closed linear subspace of  $L^2(\Omega)$  generated by the random variables  $\{H_q(Z(h)), h \in \mathcal{H}, \|h\|_{\mathcal{H}} = 1\}$ , where  $H_q(x)$  is the *q*th Hermite polynomial. It can be shown (see [6, Proposition 2.2.1]) that if  $Z, Y \sim N(0, 1)$  are jointly Gaussian, then

$$\mathbb{E}\left[H_p(Z)H_q(Y)\right] = \begin{cases} p! \left(\mathbb{E}\left[ZY\right]\right)^p & \text{if } p = q\\ 0 & \text{otherwise} \end{cases}.$$
 (2.3)

For  $q \ge 1$ , it is known that the map

$$I_q(h^{\otimes q}) = H_q(Z(h)) \tag{2.4}$$

provides a linear isometry between  $\mathcal{H}^{\odot q}$  (equipped with the modified norm  $\sqrt{q!} \| \cdot \|_{\mathcal{H}^{\otimes q}}$ ) and  $\mathcal{H}_q$ , where  $I_q(\cdot)$  is the generalized Wiener-Itô stochastic integral (see [6, Theorem 2.7.7]). By convention,  $\mathcal{H}_0 = \mathbb{R}$  and  $I_0(x) = x$ .

We use the following integral multiplication theorem from [7, Proposition 1.1.3]. Suppose  $f \in \mathcal{H}^{\odot p}$  and  $g \in \mathcal{H}^{\odot q}$ . Then

$$I_p(f)I_q(g) = \sum_{r=0}^{p \wedge q} r! \binom{p}{r} \binom{q}{r} I_{p+q-2r}(f \widetilde{\otimes}_r g),$$
(2.5)

where  $f \otimes_r g$  denotes the symmetrization of  $f \otimes_r g$ . For a product of more than two integrals, see Peccati and Taqqu [9].

#### 2.2 Stochastic Integration and fBm

We refer to the 'time domain' and 'spectral domain' representations of fBm. The reader may refer to [10, 11] for details. Let  $\mathcal{E}$  denote the set of real-valued step functions on  $\mathbb{R}$ . Let  $B^H$  denote fBm with Hurst parameter H. For this case, we view  $B^H$  as an isonormal Gaussian process on the Hilbert space  $\mathfrak{H}$ , which is the closure of  $\mathcal{E}$  with respect to the inner product  $\langle f, g \rangle_{\mathfrak{H}} = \mathbb{E}[I(f)I(g)]$ . Consider also the inner product space

$$\tilde{\Lambda}_{H} = \left\{ f : f \in L^{2}(\mathbb{R}), \int_{\mathbb{R}} |\mathcal{F}f(\xi)|^{2} |\xi|^{1-2H} d\xi < \infty \right\},\$$

where  $\mathcal{F}f = \int_{\mathbb{R}} f(x)e^{i\xi x} dx$  is the Fourier transform, and the inner product of  $\tilde{\Lambda}_H$  is given by

$$\langle f, g \rangle_{\tilde{\Lambda}_H} = \frac{1}{C_H^2} \int_{\mathbb{R}} \mathcal{F}f(\xi) \overline{\mathcal{F}g(\xi)} |\xi|^{1-2H} d\xi,$$
 (2.6)

where  $C_H = \left(\frac{2\pi}{\Gamma(2H+1)\sin(\pi H)}\right)^{\frac{1}{2}}$ . It is known (see [10, Theorem 3.1]) that the space  $\tilde{\Lambda}_H$  is isometric to a subspace of  $\mathfrak{H}$ , and  $\tilde{\Lambda}_H$  contains  $\mathcal{E}$  as a dense subset. This inner product (2.6) is known as the 'spectral measure' of fBm. In the case  $H \in (\frac{1}{2}, 1)$ , there is another isometry from the space

$$|\Lambda_H| = \left\{ f: \int_0^\infty \int_0^\infty |f(u)| |f(v)| |u-v|^{2H-2} du \, dv < \infty \right\}$$

to a subspace of  $\mathfrak{H}$ , where the inner product is defined as

$$\langle f,g \rangle_{|\Lambda_H|} = H(2H-1) \int_0^\infty \int_0^\infty f(u)g(v)|u-v|^{2H-2} du \, dv,$$

see [10] or [7, Sect. 5.1].

## **3 Proof of Theorem 1**

For any  $\gamma > 0$  and  $\lambda > 0$ , we can write

$$\lambda^{-\gamma} = \frac{1}{\Gamma(\gamma)} \int_0^\infty y^{\gamma-1} e^{-\lambda y} dy,$$

where  $\Gamma$  is the Gamma function defined by  $\Gamma(\gamma) = \int_0^\infty y^{\gamma-1} e^{-y} dy$ . As a consequence, the covariance (1.7) can be written as

$$\mathbb{E}[Z_s Z_t] = \frac{1}{\Gamma(\gamma)} \int_0^\infty y^{\gamma - 1} e^{-(t+s)y} dy.$$
(3.1)

Notice that this representation implies the covariance (1.7) is positive definite. Taking first the expectation with respect to the process Z, and using formula (3.1), we obtain

$$R(s,t) = \frac{1}{\Gamma(\gamma)} \int_0^\infty \mathbb{E}\left[\left(\int_0^t e^{yu} dB_u^H\right) \left(\int_0^s e^{yu} dB_u^H\right)\right] y^{\gamma-1} e^{-(t+s)y} dy$$
$$= \frac{1}{\Gamma(\gamma)} \int_0^\infty \left\langle e^{yu} \mathbf{1}_{[0,t]}(u), e^{yv} \mathbf{1}_{[0,s]}(v) \right\rangle_{\mathfrak{H}} y^{\gamma-1} e^{-(t+s)y} dy.$$

Using the isometry between  $\tilde{\Lambda}_H$  and a subspace of  $\mathfrak{H}$  (see Sect. 2.2), we can write

$$\begin{split} \left\langle e^{yu} \mathbf{1}_{[0,t]}(u), e^{yv} \mathbf{1}_{[0,s]}(v) \right\rangle_{s_{5}} &= C_{H}^{-2} \int_{\mathbb{R}} |\xi|^{1-2H} (\mathcal{F} \mathbf{1}_{[0,t]} e^{y\cdot}) (\overline{\mathcal{F} \mathbf{1}_{[0,s]} e^{y\cdot}}) \, d\xi \\ &= C_{H}^{-2} \int_{\mathbb{R}} \frac{|\xi|^{1-2H}}{y^{2} + \xi^{2}} \left( e^{yt + i\xi t} - 1 \right) \left( e^{ys - i\xi s} - 1 \right) \, d\xi, \end{split}$$

where  $(\mathcal{F}\mathbf{1}_{[0,t]}e^{x})$  denotes the Fourier transform and  $C_H = \left(\frac{2\pi}{\Gamma(2H+1)\sin(\pi H)}\right)^{\frac{1}{2}}$ . This allows us to write, making the change of variable  $\xi = \eta y$ ,

$$R(s,t) = \frac{1}{\Gamma(\gamma)C_{H}^{2}} \int_{0}^{\infty} \int_{\mathbb{R}} y^{\gamma-1} \frac{|\xi|^{1-2H}}{y^{2}+\xi^{2}} \left(e^{i\xi t} - e^{-yt}\right) \left(e^{-i\xi s} - e^{-ys}\right) d\xi dy$$

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$$= \frac{1}{\Gamma(\gamma)C_{H}^{2}} \int_{0}^{\infty} \int_{\mathbb{R}} y^{-\alpha-1} \frac{|\eta|^{1-2H}}{1+\eta^{2}} \left( e^{i\eta yt} - e^{-yt} \right) \left( e^{-i\eta ys} - e^{-ys} \right) \, d\eta \, dy,$$
(3.2)

where  $\alpha = 2H - \gamma$ . By Euler's identity, adding and subtracting 1 to compensate the singularity of  $y^{-\alpha-1}$  at the origin, we can write

$$e^{i\eta yt} - e^{-yt} = (\cos(\eta yt) - 1 + i\sin(\eta yt)) + (1 - e^{-yt}).$$
(3.3)

Substituting (3.3) into (3.2) and taking into account that the integral of the imaginary part vanishes because it is an odd function, we obtain

$$\begin{aligned} R(s,t) &= \frac{2}{\Gamma(\gamma)C_H^2} \int_0^\infty \int_0^\infty y^{-\alpha-1} \frac{\eta^{1-2H}}{1+\eta^2} \Big( (1-\cos(\eta yt))(1-\cos(\eta ys)) \\ &+ \sin(\eta yt) \sin(\eta ys) + (\cos(\eta ys) - 1)(1-e^{-yt}) + (\cos(\eta yt) - 1)(1-e^{-ys}) \\ &+ (1-e^{-yt})(1-e^{-ys}) \Big) \, d\eta \, dy. \end{aligned}$$

Let  $B^{(j)} = \{B^{(j)}(\eta, t), \eta \ge 0, t \ge 0\}, j = 1, 2$  denote two independent Brownian sheets. That is, for  $j = 1, 2, B^{(j)}$  is a continuous Gaussian field with mean zero and covariance given by

$$\mathbb{E}\left[B^{(j)}(\eta,t)B^{(j)}(\xi,s)\right] = \min(\eta,\xi) \times \min(t,s).$$

We define the following stochastic processes:

$$U_t = \frac{\sqrt{2}}{\sqrt{\Gamma(\gamma)}C_H} \int_0^\infty \int_0^\infty y^{-\frac{\alpha}{2} - \frac{1}{2}} \sqrt{\frac{\eta^{1-2H}}{1 + \eta^2}} \left(\cos(\eta y t) - 1\right) B^{(1)}(d\eta, dy), \quad (3.4)$$

$$V_t = \frac{\sqrt{2}}{\sqrt{\Gamma(\gamma)}C_H} \int_0^\infty \int_0^\infty y^{-\frac{\alpha}{2} - \frac{1}{2}} \sqrt{\frac{\eta^{1-2H}}{1+\eta^2}} \left(\sin(\eta yt)\right) B^{(2)}(d\eta, dy), \tag{3.5}$$

$$Y_t = \frac{\sqrt{2}}{\sqrt{\Gamma(\gamma)}C_H} \int_0^\infty \int_0^\infty y^{-\frac{\alpha}{2} - \frac{1}{2}} \sqrt{\frac{\eta^{1-2H}}{1+\eta^2}} \left(1 - e^{-yt}\right) B^{(1)}(d\eta, dy), \tag{3.6}$$

where the integrals are Wiener-Itô integrals with respect to the Brownian sheet. We then define the stochastic process  $X = \{X_t, t \ge 0\}$  by  $X_t = U_t + V_t + Y_t$ , and we have  $\mathbb{E}[X_sX_t] = R(s, t)$  as given in (3.2). These processes have the following properties:

(I) The process  $W_t = U_t + V_t$  is a fractional Brownian motion with Hurst parameter  $\frac{\alpha}{2}$  scaled with the constant  $\sqrt{\kappa}$ . In fact, the covariance of this process is

$$\mathbb{E}[W_t W_s] = \frac{2}{\Gamma(\gamma) C_H^2} \int_0^\infty \int_0^\infty y^{-\alpha - 1} \frac{\eta^{1 - 2H}}{1 + \eta^2} \Big( (\cos(\eta yt) - 1) (\cos(\eta ys) - 1) \\ + \sin(\eta yt) \sin(\eta ys) \Big) d\eta dy \\ = \frac{1}{\Gamma(\gamma) C_H^2} \int_0^\infty \int_\mathbb{R} y^{\gamma - 1} \frac{|\xi|^{1 - 2H}}{y^2 + \xi^2} (e^{i\xi t} - 1) (e^{-i\xi s} - 1) d\xi dy.$$

Integrating in the variable y we finally obtain

$$\mathbb{E}[W_t W_s] = \frac{c_1}{\Gamma(\gamma) C_H^2} \int_{\mathbb{R}} \frac{(e^{i\xi t} - 1)(e^{-i\xi s} - 1)}{|\xi|^{\alpha+1}} d\xi,$$

where  $c_1 = \int_0^\infty \frac{z^{\gamma-1}}{1+z^2} dz = \kappa \Gamma(\gamma)$ . Taking into account the Fourier transform representation of fBm (see [11, p. 328]), this implies  $\kappa^{-\frac{1}{2}}W$  is a fractional Brownian motion with Hurst parameter  $\frac{\alpha}{2}$ .

(II) The process Y coincides, up to a constant, with the process  $Y^K$  introduced in (1.4) with  $K = 2\alpha + 1$ . In fact, the covariance of this process is given by

$$\mathbb{E}[Y_t Y_s] = \frac{2c_2}{\Gamma(\gamma)C_H^2} \int_0^\infty y^{-\alpha - 1} (1 - e^{-yt}) (1 - e^{-ys}) dy, \qquad (3.7)$$

where

$$c_2 = \int_0^\infty \frac{\eta^{1-2H}}{1+\eta^2} d\eta.$$

Notice that the process X is self-similar with exponent  $\frac{\alpha}{2}$ . This concludes the proof of Theorem 1.

## 4 Proof of Theorem 2

Along the proof, the symbol *C* denotes a generic, positive constant, which may change from line to line. The value of *C* will depend on parameters of the process and on *T*, but not on the increment width  $n^{-1}$ .

For integers  $n \ge 1$ , define a partition of  $[0, \infty)$  composed of the intervals  $\{[\frac{j}{n}, \frac{j+1}{n}), j \ge 0\}$ . For the process X and related processes U, V, W, Y defined in Sect. 3, we introduce the notation

$$\Delta X_{\underline{j}} = X_{\underline{j+1}} - X_{\underline{j}} \text{ and } \Delta X_0 = X_{\underline{1}},$$

with corresponding notation for U, V, W, Y. We start the proof of Theorem 2 with two technical results about the components of the increments.

## 4.1 Preliminary Lemmas

**Lemma 3** Using above notation with integers  $n \ge 2$  and  $j, k \ge 0$ , we have

- (a)  $\mathbb{E}\left[\Delta W_{\frac{j}{n}}\Delta W_{\frac{k}{n}}\right] = \frac{\kappa}{2}n^{-\alpha} \left(|j-k-1|^{\alpha}-2|j-k|^{\alpha}+|j-k-1|^{\alpha}\right), \text{ where } \kappa \text{ is } defined in (1.8).$
- (b) For  $j + k \ge 1$ ,

$$\left|\mathbb{E}\left[\Delta Y_{\underline{i}} \Delta Y_{\underline{k}} \right]\right| \le Cn^{-\alpha} (j+k)^{\alpha-2}$$

for a constant C > 0 that is independent of j, k and n.

*Proof* Property (a) is well-known for fractional Brownian motion. For (b), we have from (3.7):

$$\mathbb{E}\left[\Delta Y_{\frac{j}{n}}\Delta Y_{\frac{k}{n}}\right] = \frac{2c_2}{\Gamma(\gamma)C_H^2 n^{\alpha}} \int_0^{\infty} y^{-\alpha-1} \left(e^{-yj} - e^{-y(j+1)}\right) \left(e^{-yk} - e^{-y(k+1)}\right) dy$$
$$= \frac{2c_2}{\Gamma(\gamma)C_H^2 n^{\alpha}} \int_0^{\infty} y^{-\alpha+1} \int_0^1 \int_0^1 e^{-y(j+k+u+v)} du \, dv \, dy.$$

Note that the above integral is nonnegative, and we can bound this with

$$\left| \mathbb{E} \left[ \Delta Y_{\frac{j}{n}} \Delta Y_{\frac{k}{n}} \right] \right| \le C n^{-\alpha} \int_0^\infty y^{-\alpha+1} e^{-y(j+k)} \, dy$$
$$= C n^{-\alpha} (j+k)^{\alpha-2} \int_0^\infty u^{-\alpha+1} e^{-u} du$$
$$\le C n^{-\alpha} (j+k)^{\alpha-2}.$$

**Lemma 4** For  $n \ge 2$  fixed and integers  $j, k \ge 1$ ,

$$\left|\mathbb{E}\left[\Delta W_{\frac{j}{n}}\Delta Y_{\frac{k}{n}}\right]\right| \leq Cn^{-\alpha}j^{2H-2}k^{-\gamma}$$

for a constant C > 0 that is independent of j, k and n.

*Proof* From (3.4)–(3.6) in the proof of Theorem 1, observe that

$$\mathbb{E}\left[\Delta W_{\frac{j}{n}}\Delta Y_{\frac{k}{n}}\right] = \mathbb{E}\left[\left(\Delta U_{\frac{j}{n}} + \Delta V_{\frac{j}{n}}\right)\Delta Y_{\frac{k}{n}}\right] = \mathbb{E}\left[\Delta U_{\frac{j}{n}}\Delta Y_{\frac{k}{n}}\right].$$

Assume s, t > 0. By self-similarity we can define the covariance function  $\psi$  by  $\mathbb{E}[U_t Y_s] = s^{\alpha} \mathbb{E}[U_{t/s} Y_1] = s^{\alpha} \psi(t/s)$ , where, using the change-of-variable  $\theta = \eta x$ ,

$$\psi(x) = \int_0^\infty \int_0^\infty y^{-\alpha - 1} \frac{\eta^{1 - 2H}}{1 + \eta^2} \left( \cos(y\eta x) - 1 \right) \left( 1 - e^{-y} \right) d\eta \, dy$$
$$= \int_0^\infty y^{-\alpha - 1} (1 - e^{-y}) \int_0^\infty \frac{\theta^{1 - 2H} x^{2H}}{x^2 + \theta^2} \left( \cos(y\theta) - 1 \right) \, d\theta \, dy.$$

Then using the fact that

$$\left|\frac{\theta^{1-2H}x^{2H}}{x^2+\theta^2}\right| \le |\theta^{-2H}| \, |x|^{2H-1},\tag{4.1}$$

we see that  $|\psi(x)| \leq Cx^{2H-1}$ , and

$$\psi'(x) = 2H \int_0^\infty y^{-\alpha - 1} (1 - e^{-y}) \int_0^\infty \frac{\theta^{1 - 2H} x^{2H - 1}}{x^2 + \theta^2} (\cos(y\theta) - 1) \ d\theta \ dy$$
$$-2 \int_0^\infty y^{-\alpha - 1} (1 - e^{-y}) \int_0^\infty \frac{\theta^{1 - 2H} x^{2H + 1}}{(x^2 + \theta^2)^2} (\cos(y\theta) - 1) \ d\theta \ dy.$$

Using (4.1) and similarly

$$\left|\frac{\theta^{1-2H}x^{2H+1}}{(x^2+\theta^2)^2}\right| \le |\theta^{-2H}| \ |x|^{2H-2},\tag{4.2}$$

we can write

$$\left|\psi'(x)\right| \le x^{2H-2} |2H-2| \int_0^\infty y^{-\alpha-1} (1-e^{-y}) \int_0^\infty \theta^{-2H} \left(\cos(y\theta) - 1\right) \ d\theta \ dy \le C x^{2H-2}.$$

By continuing the computation, we can find that  $|\psi''(x)| \leq Cx^{2H-3}$ . We have for  $j, k \geq 1$ ,

$$\mathbb{E}\left[\Delta U_{\frac{j}{n}}\Delta Y_{\frac{k}{n}}\right] = n^{-\alpha}(k+1)^{\alpha}\left(\psi\left(\frac{j+1}{k+1}\right) - \psi\left(\frac{j}{k+1}\right)\right)$$
$$-n^{-\alpha}k^{\alpha}\left(\psi\left(\frac{j+1}{k}\right) - \psi\left(\frac{j}{k}\right)\right)$$
$$= n^{-\alpha}\left((k+1)^{\alpha} - k^{\alpha}\right)\left(\psi\left(\frac{j+1}{k+1}\right) - \psi\left(\frac{j}{k+1}\right)\right)$$
$$+n^{-\alpha}k^{\alpha}\left(\psi\left(\frac{j+1}{k+1}\right) - \psi\left(\frac{j}{k+1}\right) - \psi\left(\frac{j+1}{k}\right) + \psi\left(\frac{j}{k}\right)\right).$$

With the above bounds on  $\psi$  and its derivatives, the first term is bounded by

$$\begin{split} n^{-\alpha} \left| (k+1)^{\alpha} - k^{\alpha} \right| \left| \psi \left( \frac{j+1}{k+1} \right) - \psi \left( \frac{j}{k+1} \right) \right| \\ & \leq \alpha n^{-\alpha} \int_0^1 (k+u)^{\alpha-1} du \int_0^{\frac{1}{k+1}} \left| \psi' \left( \frac{j}{k+1} + v \right) \right| \, dv \\ & \leq C n^{-\alpha} k^{\alpha-2} \left( \frac{j}{k} \right)^{2H-2} \leq C n^{-\alpha} k^{-\gamma} j^{2H-2}, \end{split}$$

and

$$\begin{split} n^{-\alpha}k^{\alpha} \left| \psi\left(\frac{j+1}{k+1}\right) - \psi\left(\frac{j}{k+1}\right) - \psi\left(\frac{j+1}{k}\right) + \psi\left(\frac{j}{k}\right) \right| \\ &= n^{-\alpha}k^{\alpha} \left| \int_{0}^{\frac{1}{k+1}} \psi'\left(\frac{j}{k+1} + u\right) \, du - \int_{0}^{\frac{1}{k}} \psi'\left(\frac{j}{k} + u\right) \, du \right| \\ &\leq n^{-\alpha}k^{\alpha} \int_{\frac{1}{k+1}}^{\frac{1}{k}} \left| \psi'\left(\frac{j}{k} + u\right) \right| \, du + \int_{0}^{\frac{1}{k+1}} \int_{\frac{j}{k+1}}^{\frac{j}{k}} \left| \psi''(u+v) \right| \, dv \, du \\ &\leq Cn^{-\alpha}k^{\alpha-2} \left(\frac{j}{k}\right)^{2H-2} + Cn^{-\alpha}k^{\alpha-3}j\left(\frac{j}{k}\right)^{2H-3} \leq Cn^{-\alpha}k^{-\gamma}j^{2H-2}. \end{split}$$

This concludes the proof of the lemma.

## 4.2 Proof of Theorem 2

We will make use of the notation  $\beta_{j,n} = \left\| \Delta X_{\frac{j}{n}} \right\|_{L^2(\Omega)}$ . We have for integer  $j \ge 1$ ,

$$\beta_{j,n}^{2} = \mathbb{E}\left[\Delta W_{\frac{j}{n}}^{2}\right] + \mathbb{E}\left[\Delta Y_{\frac{j}{n}}^{2}\right] + 2\mathbb{E}\left[\Delta W_{\frac{j}{n}}\Delta Y_{\frac{j}{n}}\right] = \kappa n^{-\alpha}(1+\theta_{j,n}),$$

where

$$\kappa n^{-\alpha} \theta_{j,n} = \mathbb{E}\left[\Delta Y_{\underline{i}}^{2}\right] + 2\mathbb{E}\left[\Delta W_{\underline{i}}\Delta Y_{\underline{i}}\right].$$

It follows from Lemmas 3 and 4 that  $|\theta_{j,n}| \le Cj^{\alpha-2}$  for some constant C > 0. Notice that, in the definition of  $F_n(t)$ , it suffices to consider the sum for  $j \ge n_0$  for a fixed  $n_0$ . Then, we can choose  $n_0$  in such a way that  $Cn_0^{\alpha-2} \le \frac{1}{2}$ , which implies

$$\beta_{j,n}^2 \ge \kappa n^{-\alpha} (1 - C j^{\alpha - 2}) \tag{4.3}$$

for any  $j \ge n_0$ . By (2.4),

$$\beta_{j,n}^{q}H_{q}\left(\beta_{j,n}^{-1}\Delta X_{\frac{j}{n}}\right) = I_{q}^{X}\left(\left(\mathbf{1}_{\left[\frac{j}{n},\frac{j+1}{n}\right]}\right)^{\otimes q}\right),$$

where  $I_q^X$  denotes the multiple stochastic integral of order q with respect to the process X. Thus, we can write

$$F_n(t) = n^{-\frac{1}{2}} \sum_{j=n_0}^{\lfloor nt \rfloor - 1} \beta_{j,n}^{-q} I_q^X \left( \mathbf{1}_{\lfloor \frac{i}{n}, \frac{i+1}{n}}^{\otimes q} \right) \right).$$

The decomposition X = W + Y leads to

$$I_q^X\left(\mathbf{1}_{\left[\frac{j}{n},\frac{j+1}{n}\right]}^{\otimes q}\right) = \sum_{r=0}^q \binom{q}{r} I_r^W\left(\mathbf{1}_{\left[\frac{j}{n},\frac{j+1}{n}\right]}^{\otimes r}\right) I_{q-r}^Y\left(\mathbf{1}_{\left[\frac{j}{n},\frac{j+1}{n}\right]}^{\otimes q-r}\right).$$

We are going to show that the terms with r = 0, ..., q - 1 do not contribute to the limit. Define

$$G_n(t) = n^{-\frac{1}{2}} \sum_{j=n_0}^{\lfloor nt \rfloor - 1} \beta_{j,n}^{-q} I_q^W \left( \mathbf{1}_{\left[\frac{j}{n}, \frac{j+1}{n}\right]}^{\otimes q} \right)$$

and

$$\widetilde{G}_n(t) = n^{-\frac{1}{2}} \sum_{j=n_0}^{\lfloor nt \rfloor - 1} \left\| \Delta W_{j/n} \right\|_{L^2(\Omega)}^{-q} I_q^W \left( \mathbf{1}_{\lfloor \frac{j}{n}, \frac{j+1}{n} \rfloor}^{\otimes q} \right) \right\}.$$

Consider the decomposition

$$F_n(t) = (F_n(t) - G_n(t)) + (G_n(t) - \widetilde{G}_n(t)) + \widetilde{G}_n(t)$$

Notice that all these processes vanish at t = 0. We claim that for any  $0 \le s < t \le T$ , we have

$$\mathbb{E}[|F_n(t) - G_n(t) - (F_n(s) - G_n(s))|^2] \le \frac{(\lfloor nt \rfloor - \lfloor ns \rfloor)^{\delta}}{n}$$
(4.4)

and

$$\mathbb{E}[|G_n(t) - \widetilde{G}_n(t) - (G_n(s) - \widetilde{G}_n(s))|^2] \le \frac{(\lfloor nt \rfloor - \lfloor ns \rfloor)^{\delta}}{n}, \tag{4.5}$$

where  $0 \leq \delta < 1$ . By Lemma 3,  $\|\Delta W_{j/n}\|_{L^2(\Omega)}^2 = \kappa n^{-\alpha}$  for every *j*. As a consequence, using (2.4) we can also write

$$\widetilde{G}_n(t) = n^{-\frac{1}{2}} \sum_{j=n_0}^{\lfloor nt \rfloor - 1} H_q\left(\kappa^{-\frac{1}{2}} n^{\frac{\alpha}{2}} \Delta W_{\frac{j}{n}}\right).$$

Since  $\kappa^{-\frac{1}{2}}W$  is a fractional Brownian motion, the Breuer-Major theorem implies that the process  $\widetilde{G}$  converges in D([0, T]) to a scaled Brownian motion  $\{\sigma B_t, t \in [0, T]\}$ , where  $\sigma^2$  is given in (1.11). By the fact that all the *p*-norms are equivalent on a fixed Wiener chaos, the estimates (4.4) and (4.5) lead to

$$\mathbb{E}[|F_n(t) - G_n(t) - (F_n(s) - G_n(s))|^{2p}] \le \frac{(\lfloor nt \rfloor - \lfloor ns \rfloor)^{\delta p}}{n^p}$$
(4.6)

and

$$\mathbb{E}[|G_n(t) - \widetilde{G}_n(t) - (G_n(s) - \widetilde{G}_n(s))|^{2p}] \le \frac{(\lfloor nt \rfloor - \lfloor ns \rfloor)^{\delta p}}{n^p}, \tag{4.7}$$

for all  $p \ge 1$ . Letting *n* tend to infinity, we deduce from (4.6) and (4.7) that for any  $t \in [0, T]$  the sequences  $F_n(t) - G_n(t)$  and  $G_n(t) - \widetilde{G}_n(t)$  converge to zero in  $L^{2p}(\Omega)$  for any  $p \ge 1$ . This implies that the finite dimensional distributions of the processes  $F_n - G_n$  and  $G_n - \widetilde{G}_n$  converge to zero in law. Moreover, by Billingsley [2, Theorem 13.5], (4.6) and (4.7) also imply that the sequences  $F_n - G_n$  and  $G_n - \widetilde{G}_n$  are tight in D([0, T]). Therefore, these sequences converge to zero in the topology of D([0, T]).

*Proof of* (4.4) We can write

$$\mathbb{E}\left[|F_n(t) - G_n(t) - (F_n(s) - G_n(s))|^2\right] \le C \sum_{r=0}^{q-1} \mathbb{E}[\Phi_{r,n}^2],$$

where

$$\Phi_{r,n} = n^{-\frac{1}{2}} \sum_{j=\lfloor ns \rfloor \lor n_0}^{\lfloor nt \rfloor - 1} \beta_{j,n}^{-q} I_r^W \left( \mathbf{1}_{\left[\frac{j}{n}, \frac{j+1}{n}\right]}^{\otimes r} \right) I_{q-r}^Y \left( \mathbf{1}_{\left[\frac{j}{n}, \frac{j+1}{n}\right]}^{\otimes q-r} \right).$$

We have, using (4.3),

$$\mathbb{E}[\Phi_{r,n}^{2}] \leq n^{-1+q\alpha} \\ \times \sum_{j,k=\lfloor ns \rfloor \lor n_{0}}^{\lfloor nt \rfloor -1} \left| \mathbb{E}\left[ I_{r}^{W}\left(\mathbf{1}_{\left[\frac{i}{n},\frac{i+1}{n}\right]}^{\otimes r}\right) I_{q-r}^{Y}\left(\mathbf{1}_{\left[\frac{i}{n},\frac{i+1}{n}\right]}^{\otimes q-r}\right) I_{r}^{W}\left(\mathbf{1}_{\left[\frac{k}{n},\frac{k+1}{n}\right]}^{\otimes r}\right) I_{q-r}^{Y}\left(\mathbf{1}_{\left[\frac{k}{n},\frac{k+1}{n}\right]}^{\otimes q-r}\right) \right] \right|.$$

Using a diagram method for the expectation of four stochastic integrals (see [9]), we find that, for any j, k, the above expectation consists of a sum of terms of the form

$$\left(\mathbb{E}\left[\Delta W_{\frac{j}{n}}\Delta W_{\frac{k}{n}}\right]\right)^{a_1}\left(\mathbb{E}\left[\Delta Y_{\frac{j}{n}}\Delta Y_{\frac{k}{n}}\right]\right)^{a_2}\left(\mathbb{E}\left[\Delta W_{\frac{j}{n}}\Delta Y_{\frac{k}{n}}\right]\right)^{a_3}\left(\mathbb{E}\left[\Delta Y_{\frac{j}{n}}\Delta W_{\frac{k}{n}}\right]\right)^{a_4},$$

where the  $a_i$  are nonnegative integers such that  $a_1+a_2+a_3+a_4 = q$ ,  $a_1 \le r \le q-1$ , and  $a_2 \le q-r$ . First, consider the case with  $a_3 = a_4 = 0$ , so that we have the sum

$$n^{-1+q\alpha} \sum_{j,k=\lfloor ns \rfloor \lor n_0}^{\lfloor nt \rfloor -1} \left( \mathbb{E} \left[ \Delta W_{\frac{j}{n}} \Delta W_{\frac{k}{n}} \right] \right)^{a_1} \left( \mathbb{E} \left[ \Delta Y_{\frac{j}{n}} \Delta Y_{\frac{k}{n}} \right] \right)^{q-a_1},$$

where  $0 \le a_1 \le q - 1$ . Applying Lemma 3, we can control each of the terms in the above sum by

$$n^{-q\alpha}(|j-k+1|^{\alpha}-2|j-k|^{\alpha}+|j-k-1|^{\alpha})^{a_1}(j+k)^{(q-a_1)(\alpha-2)},$$

which gives

$$n^{-1+q\alpha} \sum_{j,k=\lfloor ns \rfloor \lor n_0}^{\lfloor nt \rfloor -1} \left| \mathbb{E} \left[ \Delta W_{\frac{j}{n}} \Delta W_{\frac{k}{n}} \right] \right|^{a_1} \left| \mathbb{E} \left[ \Delta Y_{\frac{j}{n}} \Delta Y_{\frac{k}{n}} \right] \right|^{q-a_1}$$

$$\leq Cn^{-1} \left( \sum_{j=\lfloor ns \rfloor \lor n_0}^{\lfloor nt \rfloor -1} j^{\alpha-2} + \sum_{j,k=\lfloor ns \rfloor \lor n_0, j \neq k}^{\lfloor nt \rfloor -1} |j-k|^{(q-1)(\alpha-2)} (j+k)^{\alpha-2} \right)$$

$$\leq Cn^{-1} \sum_{j=\lfloor ns \rfloor \lor n_0}^{\lfloor nt \rfloor -1} (j^{\alpha-2} + j^{q(\alpha-2)+1})$$

$$\leq Cn^{-1} \left( \lfloor nt \rfloor - \lfloor ns \rfloor \right)^{(\alpha-1)\lor 0} + (\lfloor nt \rfloor - \lfloor ns \rfloor)^{[q(\alpha-2)+2]\lor 0} \right).$$
(4.8)

Next, we consider the case where  $a_3 + a_4 \ge 1$ . By Lemma 3, we have that, up to a constant *C*,

$$\left|\mathbb{E}\left[\Delta Y_{\underline{i}} \Delta Y_{\underline{k}}\right]\right| \leq C \left|\mathbb{E}\left[\Delta W_{\underline{i}} \Delta W_{\underline{k}}\right]\right|,$$

so we may assume  $a_2 = 0$ , and have to handle the term

$$n^{-1+q\alpha} \sum_{j,k=\lfloor ns \rfloor \lor n_0}^{\lfloor nt \rfloor -1} \left| \mathbb{E} \left[ \Delta W_{\frac{j}{n}} \Delta W_{\frac{k}{n}} \right] \right|^{q-a_3-a_4} \left| \mathbb{E} \left[ \Delta W_{\frac{j}{n}} \Delta Y_{\frac{k}{n}} \right] \right|^{a_3} \left| \mathbb{E} \left[ \Delta Y_{\frac{j}{n}} \Delta W_{\frac{k}{n}} \right] \right|^{a_4}$$

$$(4.9)$$

for all allowable values of  $a_3$ ,  $a_4$  with  $a_3 + a_4 \ge 1$ . Consider the decomposition

$$n^{-1+q\alpha} \sum_{j,k=\lfloor ns \rfloor \lor n_0}^{\lfloor nt \rfloor -1} \left| \mathbb{E} \left[ \Delta W_{\frac{j}{n}} \Delta W_{\frac{k}{n}} \right] \right|^{q-a_3-a_4} \left| \mathbb{E} \left[ \Delta W_{\frac{j}{n}} \Delta Y_{\frac{k}{n}} \right] \right|^{a_3} \left| \mathbb{E} \left[ \Delta Y_{\frac{j}{n}} \Delta W_{\frac{k}{n}} \right] \right|^{a_4}$$

$$= n^{q\alpha-1} \sum_{j=\lfloor ns \rfloor \lor n_0}^{\lfloor nt \rfloor -1} \left| \mathbb{E} \left[ \Delta W_{\frac{j}{n}}^2 \right] \right|^{q-a_3-a_4} \left| \mathbb{E} \left[ \Delta W_{\frac{j}{n}} \Delta Y_{\frac{j}{n}} \right] \right|^{a_3+a_4}$$

$$+ n^{q\alpha-1} \sum_{j=\lfloor ns \rfloor \lor n_0}^{\lfloor nt \rfloor -1} \sum_{k=\lfloor ns \rfloor \lor n_0}^{j-1} \left| \mathbb{E} \left[ \Delta W_{\frac{j}{n}} \Delta W_{\frac{k}{n}} \right] \right|^{q-a_3-a_4} \left| \mathbb{E} \left[ \Delta W_{\frac{j}{n}} \Delta Y_{\frac{k}{n}} \right] \right|^{a_3} \left| \mathbb{E} \left[ \Delta Y_{\frac{j}{n}} \Delta W_{\frac{k}{n}} \right] \right|^{a_4}$$

$$+ n^{q\alpha-1} \sum_{k=\lfloor ns \rfloor \lor n_0}^{\lfloor nt \rfloor -1} \sum_{k=\lfloor ns \rfloor \lor n_0}^{k-1} \left| \mathbb{E} \left[ \Delta W_{\frac{j}{n}} \Delta W_{\frac{k}{n}} \right] \right|^{q-a_3-a_4} \left| \mathbb{E} \left[ \Delta W_{\frac{j}{n}} \Delta Y_{\frac{k}{n}} \right] \right|^{a_3} \left| \mathbb{E} \left[ \Delta Y_{\frac{j}{n}} \Delta W_{\frac{k}{n}} \right] \right|^{a_4}.$$

We have, by Lemmas 3 and 4,

$$n^{-1+q\alpha} \sum_{j,k=\lfloor ns \rfloor \lor n_0}^{\lfloor nt \rfloor -1} \left| \mathbb{E} \left[ \Delta W_{\frac{j}{n}} \Delta W_{\frac{k}{n}} \right] \right|^{q-a_3-a_4} \left| \mathbb{E} \left[ \Delta W_{\frac{j}{n}} \Delta Y_{\frac{k}{n}} \right] \right|^{a_3} \left| \mathbb{E} \left[ \Delta Y_{\frac{j}{n}} \Delta W_{\frac{k}{n}} \right] \right|^{a_4} \\ \leq Cn^{-1} \sum_{j=\lfloor ns \rfloor \lor n_0}^{\lfloor nt \rfloor -1} j^{(a_3+a_4)(\alpha-2)} \\ + Cn^{-1} \sum_{j=\lfloor ns \rfloor \lor n_0}^{\lfloor nt \rfloor -1} j^{a_3(2H-2)-a_4\gamma} \sum_{k=\lfloor ns \rfloor \lor n_0}^{j-1} k^{-a_3\gamma+a_4(2H-2)} \left| j-k \right|^{(q-a_3-a_4)(\alpha-2)} \\ + Cn^{-1} \sum_{k=\lfloor ns \rfloor \lor n_0}^{\lfloor nt \rfloor -1} k^{-a_3\gamma+a_4(2H-2)} \sum_{j=\lfloor ns \rfloor \lor n_0}^{k-1} j^{a_3(2H-2)-a_4\gamma} \left| k-j \right|^{(q-a_3-a_4)(\alpha-2)} \\ \leq Cn^{-1} \left( (\lfloor nt \rfloor - \lfloor ns \rfloor)^{[(a_3+a_4)(\alpha-2)+1]\lor 0} + (\lfloor nt \rfloor - \lfloor ns \rfloor)^{[q(\alpha-2)+2]\lor 0} \\ + (\lfloor nt \rfloor - \lfloor ns \rfloor)^{[a_3(2H-2)-a_4\gamma+1]\lor 0} + (\lfloor nt \rfloor - \lfloor ns \rfloor)^{[a_4(2H-2)-a_3\gamma+1]\lor 0} \right).$$

$$(4.10)$$

Then (4.8) and (4.10) imply (4.4) because  $\alpha < 2 - \frac{1}{q}$ . *Proof of (4.5)* We have

$$G_n(t) - \widetilde{G}_n(t) = n^{-\frac{1}{2}} \sum_{j=n_0}^{\lfloor nt \rfloor - 1} \left( \beta_{j,n}^{-q} - \left\| \Delta W_{\frac{j}{n}} \right\|_{L^2(\Omega)}^{-q} \right) I_q^W \left( \mathbf{1}_{\lfloor \frac{j}{n}, \frac{j+1}{n}}^{\otimes q} \right)$$

and we can write, using (4.3) for any  $j \ge n_0$ ,

$$\left|\beta_{j,n}^{-q} - \left\|\Delta W_{\frac{j}{n}}\right\|_{L^{2}(\Omega)}^{-q}\right| = (\kappa^{-1}n^{\alpha})^{\frac{q}{2}} \left|(1+\theta_{j,n})^{-\frac{q}{2}} - 1\right| \le C \left(\kappa^{-1}n^{\alpha}j^{\alpha-2}\right)^{\frac{q}{2}}.$$

This leads to the estimate

$$\mathbb{E}\left[\left|G_{n}(t)-\widetilde{G}_{n}(t)-(G_{n}(s)-\widetilde{G}_{n}(s))\right|^{2}\right] \leq Cn^{-1}$$

$$\times \left(\sum_{j=\lfloor ns \rfloor \lor n_{0}}^{\lfloor nt \rfloor-1} j^{\alpha-2} + \sum_{j,k=\lfloor ns \rfloor \lor n_{0}, j \neq k}^{\lfloor nt \rfloor-1} |j-k|^{q(\alpha-2)}\right)$$

$$\leq Cn^{-1} \left(\lfloor nt \rfloor - \lfloor ns \rfloor\right)^{(\alpha-1)\lor 0} + \left(\lfloor nt \rfloor - \lfloor ns \rfloor\right)^{[q(\alpha-2)+2]\lor 0}\right),$$

which implies (4.5).

This concludes the proof of Theorem 2.

Acknowledgements D. Nualart is supported by NSF grant DMS1512891 and the ARO grant FED0070445

## References

- 1. R.M. Balan, C.A. Tudor, The stochastic heat equation with fractional-colored noise: existence of the solution. Latin Am. J. Probab. Math. Stat. 4, 57–87 (2008)
- 2. P. Billingsley, Convergence of Probability Measures, 2nd edn. (Wiley, New York, 1999)
- P. Breuer, P. Major, Central limit theorems for nonlinear functionals of Gaussian fields. J. Multivar. Anal. 13(3), 425–441 (1983)
- C. Houdré, J. Villa, An example of infinite dimensional quasi-helix. Stoch. Models Contemp. Math. 366, 195–201 (2003)
- P. Lei, D. Nualart, A decomposition of the bifractional Brownian motion and some applications. Stat. Probab. Lett. 79, 619–624 (2009)
- 6. I. Nourdin, G. Peccati, Normal Approximations with Malliavin Calculus: From Stein's Method to Universality (Cambridge University Press, Cambridge, 2012)
- 7. D. Nualart, The Malliavin Calculus and Related Topics, 2nd edn. (Springer, Berlin, 2006)
- 8. H. Ouahhabi, C.A. Tudor, Additive functionals of the solution to the fractional stochastic heat equation. J. Fourier Anal. Appl. **19**(4), 777–791 (2012)
- 9. G. Peccati, M.S. Taqqu, Wiener Chaos: Moments, Cumulants and Diagrams (Springer, Berlin, 2010)
- V. Pipiras, M.S. Taqqu, Integration questions related to fractional Brownian motion. Probab. Theory Relat. Fields 118 251–291 (2000)
- 11. G. Samorodnitsky, M.S. Taqqu, Stable Non-Gaussian Random Processes: Stochastic Models with Infinite Variance (Chapman & Hall, New York, 1994)
- S. Torres, C.A. Tudor, F. Viens, Quadratic variations for the fractional-colored stochastic heat equation. Electron. J. Probab. 19(76), 1–51 (2014)
- 13. C.A. Tudor, Analysis of Variations for Self-Similar Processes: A Stochastic Calculus Approach (Springer, Berlin, 2013)
- 14. C.A. Tudor, Y. Xiao, Sample paths of the solution to the fractional-colored stochastic heat equation. Stochastics Dyn. **17**(1) (2017)