Chapter 5 Mafic Lower-Mantle Mineral Association

Abstract Mafic mineral association in the lower mantle is subordinate to the ultramafic one. It includes bridgmanite, CaSi-perovskite, $SiO₂$ and anhydrous aluminous phases. The former three are the same as that observed in the ultramafic association; but their chemical compositions differ from those in the ultramafic association, mainly in the significant enrichment of Al. Among aluminous phases the NAL phase occurs at low-pressure conditions and is replaced by a CF phase at a depth of 800–1200 km depth. NAL phase is also concentrated in Na and K, while CF phase does not contain K. The partition coefficient of aluminium between bridgmanite and the NAL phase vary from 0.10 to 0.26, demonstrating that the Al enrichment in bridgmanite occurs at the expense of the Al decrease in the NAL phases. The Al concentration in the CF phase remains constant and the Al concentration in bridgmanite, after reaching maximal concentrations (24–25 wt%) with disintegration of the NAL phase, remains constant as well. In addition to the major minerals, phase Egg, δ -AlOOH, and a series of dense hydrous magnesium silicates (DHMS) are expected to be present in the mafic association. Among these DHMS, Phase D and Phase H are most likely to occur in the subducting slabs within the lower mantle. Some of these minerals (phase Egg and δ -AlOOH) are observed in natural geological materials; the others have only been synthesized in laboratory experiments.

5.1 General

Mafic association in the lower mantle appears as a consequence of recrystallization from subducting slab materials, predominantly of the middle-ocean ridge basalt (MORB) composition during the course of the subduction of lithospheric plates through the 660 km discontinuity. MORB contains more silica (48–51 wt% $SiO₂$) and aluminium (15–16 wt% Al_2O_3) than the bulk mantle, and therefore is composed of a different mineral association than the juvenile ultramafic composition. Four major mineral phases complete the mafic mineral association in the lower mantle: bridgmanite, CaSi-perovskite, $SiO₂$, and anhydrous aluminous phase (NAL, CF and/or CAS) (Funamori et al. [2000;](#page-36-0) Hirose et al. [2005](#page-36-0); Akaogi [2007;](#page-35-0) Ricolleau et al.

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[2010\)](#page-40-0). The first three minerals above are the same as those observed in the ultramafic association; they experienced similar phase transformations with pressure. However, the mafic association is not as commonly observed as the ultramafic counterparts; only a few samples have been identified as inclusions in lower-mantle diamond to date (Wirth et al. [2007](#page-41-0); Walter et al. [2011](#page-41-0); Thomson et al. [2014\)](#page-41-0).

The chemical compositions of the mafic association, established from high pressure–temperature $(P-T)$ experiments on MORB materials, differ from those in the ultramafic association, mainly in the significant enrichment in Al (Kesson et al. [1994;](#page-37-0) Hirose et al. [1999,](#page-36-0) [2005;](#page-36-0) Ono et al. [2001;](#page-39-0) Hirose and Fei [2002](#page-36-0); Litasov et al. [2004;](#page-38-0) Litasov and Ohtani [2005](#page-38-0); Ricolleau et al. [2010\)](#page-40-0). In addition to the major mineral phases mentioned above, high P–T experimental results have showed the occurrence of various minor hydrous minerals at relevant $P-T$ conditions of the subducting slabs, such as hydrous aluminium phases (Phase Egg and δ -AlOOH) and a series of DHMS. Phase Egg and δ -AlOOH have been observed in geological environments, but other hydrous minerals are known only from experiments and theoretical calculations.

5.2 Major Minerals

5.2.1 Bridgmanite

According to high-pressure experiments with dry MORB compositions, the occurrence of bridgmanite is associated with CaSi-perovskite, stishovite and Al-rich phase(s) at 27–80 GPa. It has a significantly lower magnesium index $mg = 0.567$ – 0.771 (average 0.658 \pm 0.124), higher concentration of Al (Al₂O₃ = 11.1–24.0 wt %; average 16.23 \pm 4.15 wt%), as well as higher concentrations of Ti and Ca than low-Al bridgmanite from the ultramafic association with an average $mg = 0.916$ and $Al_2O_3 = 1.55$ wt% (Kesson et al. [1994](#page-37-0); Hirose et al. [1999,](#page-36-0) [2005;](#page-36-0) Ono et al. [2001;](#page-39-0) Hirose and Fei [2002;](#page-36-0) Litasov et al. [2004;](#page-38-0) Litasov and Ohtani [2005](#page-38-0); Ricolleau et al. [2010](#page-40-0)) (Table [5.1;](#page-3-0) Fig. [5.1](#page-2-0)). Experimental synthesis of bridgmanite using a hydrated mafic composition at 25–28 GPa and \sim 1500 °C shows that bridgmanite has a similar composition with $mg = 0.630$ and may be associated with Al-rich Phase D (Litasov and Ohtani [2005;](#page-38-0) Pamato et al. [2015](#page-40-0)).

Bridgmanite from the natural mafic association identified in Juina-5 kimberlite pipe, Brazil, has similar compositions with $fe = 0.346 - 0.401$ and $Al_2O_3 = 16.26 -$ 18.37 wt% (Table [5.1\)](#page-3-0) (Thomson et al. [2014](#page-41-0)).

The concentration of Al in bridgmanite of the mafic association depends on the pressure at which it forms. The Al concentration is doubled from 11 to 13 wt% Al_2O_3 at 26 GPa to 24 wt% Al_2O_3 at 100 GPa, while the mg value remains almost constant (Fig. [5.2\)](#page-4-0). Considering this correlation as a possible geobarometer, natural 'mafic' bridgmanite from Juina-5 pipe with $16.26-18.37$ wt% Al_2O_3 have been formed at \sim 65–75 GPa. Such a drastic increase of wt% Al₂O₃ with pressure is most likely correlated to the redistribution of Al in bridgmanite from an Al-rich phase with increasing pressure. On the other hand, while Ca and Na concentrations in bridgmanite show a positive correlation with Al concentration (Fig. $5.3a,b$ $5.3a,b$) that is similar to that in

Fig. 5.1 Distribution of mg index (a) and Al_2O_3 (b) in bridgmanite from the mafic association. Data from Kesson et al. ([1994\)](#page-37-0), Hirose et al. [\(1999](#page-36-0), [2005\)](#page-36-0), Ono et al. ([2001\)](#page-39-0), Hirose and Fei ([2002\)](#page-36-0), Litasov et al. [\(2004](#page-38-0)), Litasov and Ohtani [\(2005](#page-38-0)), Ricolleau et al. [\(2010](#page-40-0)), Thomson et al. ([2014\)](#page-41-0)

bridgmanite of the ultramafic association, the correlation of Al with iron index fe is negative (Fig. [5.3](#page-5-0)c). That is, the most Al-rich bridgmanite are more Mg-rich.

Among bridgmanite varieties of the ultramafic association, there is also a group with a high Al concentration (see Sect. 4.2.2). However, those high-Al 'ultramafic' bridgmanite samples, in comparison with bridgmanite from the mafic association, are significantly richer in Mg, Ca, Na and Cr $(0.64-5.32 \text{ wt\%}$ against $0.02-1.19 \text{ wt}$ % CaO; 0.82–6.21 wt% against 0–0.11 wt% Na2O; 1.19–3.14 wt% against 0.04– 0.45 wt% Cr_2O_3 respectively) (Fig. [5.3\)](#page-5-0).

Note n.a. not analysed Note n.a. not analysed

Fig. 5.2 Al_2O_3 concentration as a function of pressure in bridgmanite synthesised using MORB compositions at high pressures and temperatures. Data from Kesson et al. ([1994](#page-37-0)), Hirose et al. ([1999,](#page-36-0) [2005](#page-36-0)), Ono et al. ([2001\)](#page-39-0), Hirose and Fei ([2002\)](#page-36-0), Litasov et al. ([2004\)](#page-38-0), Litasov and Ohtani ([2005\)](#page-38-0), Ricolleau et al. ([2010\)](#page-40-0)

5.2.2 $CaSiO₃$

Representative chemical compositions of CaSi-perovskite from experimental studies on mafic associations are shown in Table [5.2](#page-6-0). Similar to the compositions of bridgmanite discussed above, it is enriched in Al: $0.89-4.49$ wt% Al₂O₃ (average: 2.49 wt%) as compared with the Al concentration of 0.1–1.52 wt% (average: 0.29 wt%) in CaSi-perovskite from the ultramafic association (Fig. [5.4a](#page-7-0)). The Al concentration in CaSi-perovskite synthesized from MORB materials remains almost constant within the entire pressure range of 25–113 GPa where the mineral was synthesized. CaSi-perovskite from natural association identified in Juina-5 pipe (Brazil), by contrast, has very low concentrations of Al_2O_3 (0.14–0.20 wt%; Thomson et al. [2014\)](#page-41-0), similar to concentrations of $A₁O₃$ in CaSi-perovskite from the ultramafic association (see Figs. 4.21 and [5.4b](#page-7-0)).

Although Fe and Mg concentrations in CaSi-perovskite from the mafic association are significantly higher than in the ultramafic association, the magnesium indices in CaSi-perovskites from the both associations are similar (averages: 0.519 and 0.521, respectively) (Fig. [5.4b](#page-7-0)).

There are, among natural CaSi-perovskite from the mafic association, high-Ti varieties with $TiO_2 = 6.64-23.09$ wt% TiO_2 (Thomson et al. [2014\)](#page-41-0). Similar to the ultramafic association, $Ca(Si, Ti)O₃$ grains have dissolution structures (Fig. 4e in Thomson et al. [2014](#page-41-0)).

Concentrations of Fe, Mg and minor elements, such as Na, are significantly greater in CaSi-perovskite from the mafic association in comparison to this mineral from the ultramafic association (Fig. 5.5). Fe and Mg show positive correlation with Al.

The Al-concentration in $CaSiO₃$ is found to influence the transition boundary of the tetragonal to cubic structure transition, increasing it from 550 to 1900 K at

Fig. 5.3 Chemical characteristics of bridgmanite from the mafic association $(n = 20)$ in comparison to bridgmanite from the ultramafic association $(n = 34)$. **a** CaO versus Al₂O₃. **b** Na₂O versus Al₂O₃. c Iron index fe versus Al_2O_3 . Data from Kesson et al. ([1994\)](#page-37-0), Hutchison [\(1997](#page-36-0)), Hirose et al. [\(1999](#page-36-0), [2005](#page-36-0)), Stachel et al. ([2000\)](#page-41-0), Kaminsky et al. ([2001\)](#page-37-0), Ono et al. [\(2001](#page-39-0)), Hirose and Fei ([2002\)](#page-36-0), Davies et al. ([2004\)](#page-35-0), Litasov et al. [\(2004](#page-38-0)), Hayman et al. [\(2005](#page-36-0)), Litasov and Ohtani [\(2005](#page-38-0)), Tappert et al. ([2009\)](#page-41-0), Ricolleau et al. ([2010\)](#page-40-0), Zedgenizov et al. ([2014,](#page-42-0) [2015](#page-42-0)). Star shows the composition of bridgmanite from the Tenham meteorite after Tschauner et al. [\(2014](#page-41-0))

60 GPa with increasing Al_2O_3 concentration from 0 to 5.9 wt% (Kurashina et al. [2004\)](#page-38-0). The tetragonal \rightarrow cubic transition may thus occur in relatively cold subducted lithospheric slabs where $CaSiO₃$ perovskite can contain a higher proportion of Al_2O_3 (Hirose et al. [2005](#page-36-0)).

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Note n.a. not analysed

Fig. 5.4 Distribution of Al₂O₃ (a) and mg index (b) in CaSi-perovskite from the mafic association. Data from Kesson et al. [\(1994](#page-37-0)), Hirose et al. ([1999,](#page-36-0) [2005](#page-36-0)), Ono et al. ([2001\)](#page-39-0), Hirose and Fei [\(2002](#page-36-0)), Litasov et al. [\(2004](#page-38-0)), Litasov and Ohtani [\(2005](#page-38-0)), Ricolleau et al. [\(2010](#page-40-0))

5.2.3 $SiO₂$

 $SiO₂$ from the mafic association is synthesized within a range of 25–100 GPa and, like perovskitic minerals, also has an elevated Al concentration: $0.35-4.2$ wt% Al_2O_3 with an average of 2.59 wt% Al_2O_3 ; 10.3 wt% Al_2O_3 as was reported for one case (Table 4.3). It increases up to 12.6 wt% at 113 GPa in seifertite (Hirose et al. [2005\)](#page-36-0). The only known natural sample of $SiO₂$ from the mafic association in Juina-5 pipe, Brazil, does not have high $A₁O₃$ (0.09 wt%; Thomson et al. [2014](#page-41-0)). This may be caused by its association with kyanite, which is probably formed as a result of dissociation and collecting the entire amount of Al from the initial high-pressure composition.

Stishovite reported in experimental works under pressure values of 80–100 GPa (Kesson et al. [1994\)](#page-37-0) is doubtful because at such $P-T$ conditions CaCl₂-structured SiO₂ is stable. Moreover, experiments by Hirose et al. [\(2005](#page-36-0)) with MORB compositions demonstrated that the phase transition boundary between Al-bearing stishovite and CaCl₂-structured SiO₂ is lower than that of pure SiO₂ by \sim 25 GPa: it is located at \sim 62 GPa under 2000 K, which corresponds to a depth of \sim 1500 km (Table [5.3\)](#page-9-0).

Fig. 5.5 Chemical characteristics of CaSi-perovskite from the mafic association. a fe versus Al₂O₃; **b** mg versus Al₂O₃; c Na versus $Al₂O₃$

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ا ś $\overline{}$ \overrightarrow{a} Table 5.2

Fig. 5.6 Phase diagram for $SiO₂$ in subducted slabs. Temperature profiles in slabs after Kirby et al. ([1996\)](#page-37-0). Data from Hirose et al. ([2005](#page-36-0)), Lakshtanov et al. [\(2007](#page-38-0)), Bolfan-Casanova et al. ([2009\)](#page-35-0), Andrault et al. [\(2014](#page-35-0)), and Umemoto et al. ([2016\)](#page-41-0). Mantle adiabate after Katsura et al. [\(2010](#page-37-0))

Concentration of H_2O in stishovite has been shown to be related to the Al admixture in silica. While Al-free stishovite contains only $16-30$ wt. ppm H_2O , synthetic stishovite with 4.4 wt% Al_2O_3 contains \sim 3010 wt. ppm H₂O (Litasov et al. $2007a$). Most hydrogen in stishovite is associated with Al^{3+} substitutional defects on the octahedral (Si^{4+}) site. The hydrogen can occupy 40% of vacancies created by the incorporation of Al^{3+} at 20 GPa (Litasov et al. [2007a](#page-38-0)). It is possible that Al actually forms an AlOOH phase with H_2O , which is highly soluble in stishovite (e.g., Chung and Kagi [2002;](#page-35-0) Panero and Stixrude [2004](#page-40-0)). The incorporation of Al and H_2O into SiO_2 significantly reduces the transition pressure for the stishovite to post-stishovite transition boundary (Fig. 5.6). Specifically, hydrous Al-bearing stishovite transforms to post-stishovite at 24 GPa and 298 K, which is far lower than the transition pressure of 50–60 GPa for pure silica (Lakshtanov et al. [2007](#page-38-0)). First-principles static calculations and molecular dynamics demonstrated that hydrogen bonds and hydrogen mobility play a crucial role in lowering the post-stishovite transition pressure (Umemoto et al. [2016](#page-41-0)). A co-operative redistribution of hydrogen atoms is the main mechanism responsible for the reduction of the transition pressure in hydrous Al-bearing stishovite. Stishovite, containing 6.25 mol% AlOOH, transforms into $CaCl₂$ -structured phase at pressures from 30 to 32 GPa (1000 K) to 50 GPa (1300 K), i.e., within the range of 800– 1200 km depth. Since $SiO₂$ is a candidate composition in the subducted slabs, it can undergo a series of the aforementioned phase transitions that are influenced by their Al concentrations, causing stratification in the subducting slabs in the lower mantle: at 1200–1300 km depth with stishovite \rightarrow CaCl₂-structure transition and at 2200– 2300 km with CaCl₂-structure \rightarrow seifertite transition (Fig. 5.6). This is similar to the stratification in juvenile ultramafic mantle, except for the presence of the aluminous mineral phase. In addition, phase boundaries in subducted slabs are located higher compared with the juvenile ultramafic host.

5.3 Anhydrous Aluminous Phases

5.3.1 General

High $P-T$ experiments on anhydrous oceanic basalts have demonstrated that an aluminous phase with the calcium–ferrite (CF) structure (space group *Pbnm*, $Z = 4$) appears at pressures above 25 GPa and with temperatures above $1200 \degree C$ (Irifune and Ringwood [1993;](#page-37-0) Kesson et al. [1994](#page-37-0); Ono 2001). The CF phase occurs in association with bridgmanite, CaSi-perovskite and $SiO₂$ at high P–T. Subsequently, another aluminous phase, named a 'new hexagonal aluminous phase' (NAL), which has a similar chemical composition to the CF phase, was observed in the garnet-to-perovskite transformation in the $MgAl₂O₄$ -CaAl₂O₄ system at relevant conditions of the lower mantle (Akaogi et al. [1999;](#page-35-0) Miyajima et al. [1999,](#page-39-0) [2001\)](#page-39-0). The NAL and CF phases can comprise 10–30% of the subducted MORB materials in the lower mantle (Ricolleau et al. [2010;](#page-40-0) Mookherjee et al. [2012](#page-39-0)).

Studies of diamonds from the Juina-5 kimberlite in Brazil indicated that their mineral inclusions have compositions consistent with the phase assemblage expected to crystallize from basaltic materials at lower-mantle conditions: bridgmanite, CaSi-perovskite, $SiO₂$, and NAL and CF phases (Walter et al. [2011\)](#page-41-0). These results provide geological evidence for the possible existence of the NAL and CF phases in the lower mantle.

5.3.2 New Hexagonal Aluminous Phase (NAL)

The NAL phase has a general formula of ${}^{IX}A^{VIII}B_2^{VI}C_6O_{12}$, where ${}^{IX}A$ sites are comprised by large monovalent or divalent cations (Na⁺, K⁺ and/or Ca^{2+}), ^{VIII}B sites are represented by smaller cations (Mg^{2+}) and Fe^{2+}), and ^{VI}C octahedral sites are occupied by Al^{3+} (mostly), Si^{4+} and Fe^{3} . The full chemical formula of NAL is (Na, K, Ca)(Mg, Fe²⁺)₂(Al, Fe³⁺, Si)₆O₁₂. Both Na- and K-rich NAL phases were synthesized at pressures of 22–25 GPa and a temperature of 1500 \degree C (Kojitani et al. [2011\)](#page-37-0). Hence, the NAL phase concentrates not only Al but also alkalis, while the CF phase does not accommodate potassium (Miyajima et al. [2001](#page-39-0); Guignot and Andrault [2004](#page-36-0)).

Representative compositions of NAL are shown in Table [5.4](#page-12-0). Concentrations of both Na (4.93–5.9 wt% Na₂O; average 5.46 wt%) and K (0.54–1.94 wt% K₂O; average 1.00 wt%) are quite constant, and are independent of other components and pressure conditions. Al_2O_3 has a strong linear negative correlation with SiO_2 , demonstrating their substitution in the C site (Fig. $5.7a$). The Al/Si_{at} strongly decreases from 4.7 at 25 GPa to 1.73 at 33 GPa (Fig. [5.7](#page-13-0)b), which coincides with the increase of the Al concentration with increasing pressure in the associated bridgmanite.

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> Note n.a. not analysed Note n.a. not analysed

Fig. 5.7 Chemical characteristics of the NAL phase. a Al_2O_3 versus SiO_2 . b Al/Si versus pressure. Data from Hirose and Fei ([2002\)](#page-36-0), Litasov et al. [\(2004](#page-38-0)), Litasov and Ohtani ([2005\)](#page-38-0), Ricolleau et al. [\(2010](#page-40-0))

The NAL phase has a hexagonal crystal structure with the space group of $P6₃/m$ (Gasparik et al. [2000](#page-39-0); Miura et al. 2000; Miyajima et al. [2001](#page-39-0); Kojitani et al. [2011;](#page-37-0) Mookherjee et al. [2012;](#page-39-0) Wicks and Duffy [2016](#page-41-0); Wu et al. [2016a\)](#page-41-0) (Fig. [5.8\)](#page-14-0). The ^{IX}A site is a nine-fold coordinated tunnel with a hexagonal cross-section and is occupied by a large monovalent (Na^+, K^+) ^{IX}A site are likely to be disordered owing to its Wyckoff symmetry being 2*a* and the fact that only one atom occupies either of the two equivalent sites, i.e., there is half occupancy (Mookherjee et al. [2012\)](#page-39-0). The B site is six-fold co-ordinated but trigonal prismatic rather than octahedral and is typically occupied by Mg^{2+} (or replacing with Fe²⁺ in a Fe-rich variety). Each Mg^{2+} ion is surrounded by three double chains of AIO_6 (SiO₆) octahedra connected by corner-sharing oxygen ions. They form the C site. Large channels are formed by these octahedra in three double chains and are occupied typically by Na^+ or K^+ ions at half occupancy. The distorted C octahedra form edge-sharing double chains extending along the c axis that are corner linked to form tunnels. Within the tunnels lie the larger A sites while the smaller B sites are

Fig. 5.8 Hexagonal crystal structure of Fe-free NAL (space group $P6_3/m$) in different axis orientations. Lattice parameters from Pamato et al. [\(2014](#page-39-0))

surrounded by three double chains. With multiple sites that can each accept more than one cation, the NAL structure can accommodate a wide range of compositional variation. The large cation sites in the NAL structure allow them to be potential hosts for the alkali elements (Wicks and Duffy [2016\)](#page-41-0).

The hexagonal NAL structure is stable under pressure conditions of up to 40– 50 GPa, after which it transforms to the orthorhombic CF phase (Ono et al. [2009;](#page-39-0) Ricolleau et al. [2010;](#page-40-0) Imada et al. [2011\)](#page-36-0) in accordance with the later theoretical results obtained using the first-principles computational method (Kawai and Tsuchiya [2012](#page-37-0)). The enrichment of the NAL phase in potassium significantly increases its stability up to 141 GPa (Guignot and Andrault [2004](#page-36-0); Kato et al. [2013\)](#page-37-0).

Density functional theory calculations showed that the NAL phase has lower seismic velocities compared with the CF phase, and NAL could contribute to low‐ velocity provinces in the deep mantle (Mookherjee et al. [2012\)](#page-39-0). However, calculations performed by Kawai and Tsuchiya ([2012\)](#page-37-0) concluded that NAL and CF phases would have comparable shear wave velocities, producing less than 1% velocity difference across the phase transition; this conclusion is in agreement with experiments carried out by Pamato et al. ([2014\)](#page-39-0).

Since the NAL phase probably contains up to 12 mol% of iron (Litasov and Ohtani [2005;](#page-38-0) Ricolleau et al. [2010;](#page-40-0) see Table [5.4](#page-12-0)), the valence and spin states of iron in the NAL phase can potentially affect its elasticity, which in turn can affect our understanding of the mineral physics of the MORB assemblage in the lower mantle (Lin et al. [2013;](#page-38-0) Wu et al. [2016b\)](#page-42-0). Recently, two single-crystal samples (iron-free $Na_{1.14}Mg_{1.83}Al_{4.74}Si_{1.23}O_{12}$ and iron-bearing $Na_{0.71}Mg_{2.05}Al_{4.62}Si_{1.16}Fe_{0.09}^{2+}$ $\text{Fe}_{0.17}^{3+} \text{O}_{12}$) were studied by Wu et al. [\(2016a\)](#page-41-0) using synchrotron nuclear forward scattering (NFS) and X-ray diffraction (XRD) combined with diamond-anvil cells (DACs) up to 86 GPa at room temperature to investigate the potential effects of the spin transition on its equation of state at high pressures. Based on the evaluation of the hyperfine quadrupole splitting from NFS results, a pressure-induced high-spin (HS) to low-spin (LS) transition of the octahedral Fe^{3+} in the iron-bearing NAL was observed at approximately 30 GPa. Compared with the equation of state for the iron-free NAL phase using XRD data, the iron-bearing NAL undergoes a volume reduction of 1.0% $({\sim}1.2 \text{ Å}^3)$ at 33–47 GPa, which is associated with the spin transition of the octahedral Fe³⁺. The bulk sound velocities (V_{Φ}) of the iron-free and iron-bearing NAL phase are approximately 6% larger than those of Al, Fe-bearing bridgmanite and CaSi-perovskite in the lower mantle, except for the spin transition region where a notable softening of V_{Φ} with a maximum reduction of 9.4% occurs in the iron-bearing NAL phase at 41 GPa. Considering the high volume proportion of the NAL phase in subducted MORB, the distinct elastic properties of the iron-bearing NAL phase across the spin transition may provide an explanation for the observed seismic heterogeneities of subducted slabs in the lower mantle at depths below 1200 km.

5.3.3 Aluminous Phase with the Calcium–Ferrite Structure (CF Phase)

The CF phase, which is considered as a high-pressure modification of NAL, has a general formula on a 12 oxygen basis to make it better comparable with the NAL phase, of $\mathrm{VIII} \text{B}_{3}^{\text{VI}} \text{C}_{6}\text{O}_{12}$, where $\mathrm{VIII} \text{B}$ sites are occupied by Na^+ , Ca^{2+} and Mg^{2+} cations and the ^{VI}C sites are framework-forming cations Al^{3+} (mostly) and Si^{4+} , similar to that in the NAL phase. However, unlike the NAL phase, it is significantly enriched (almost doubled) in sodium (average $10.55 \text{ wt\% Na}_2\text{O}$) and does not contain potassium. Its full chemical formula is $(Na, Ca, Mg)_{3}(Al, Si)_{6}O_{12}$ (Irifune and Ringwood [1993;](#page-37-0) Kesson et al. [1994](#page-37-0); Ono et al. [2001;](#page-39-0) Hirose and Fei [2002\)](#page-36-0). Al concentration in the CF phase varies from 32.8 to 45.1 wt% Al_2O_3 (av.: 39.4 ± 3.6 wt% Al₂O₃), independently of Si and pressure values. Representative compositions of the CF phase from experimental data are presented in Table [5.5.](#page-16-0)

The CF phase has an orthorhombic space group *Pbnm* and four formula units $(Z = 4$ on a four oxygen primitive basis) in the unit cell (Decker and Kasper [1957\)](#page-35-0). Such CF structure is a common high‐pressure structure type (Wicks and Duffy [2016\)](#page-41-0) experimentally adopted by several compositions including $MgAl₂O₄$ (Irifune et al. [1991\)](#page-37-0), CaAl₂O₄ (Reid and Ringwood [1969](#page-40-0)), and NaAlSiO₄ (Liu [1977](#page-38-0)). It is one of a series of similar structures ($CaFe₂O₄$, $CaTi₂O₄$, $CaMn₂O₄$) that are com-mon high-pressure polymorphs of spinel-structured phases (Yamanaka et al. [2008\)](#page-42-0). Single-crystal diffraction experiments have shown that forsterite, $Mg_2SiO₄$, metastably adopts a related structure upon room temperature compression above 58 GPa (Finkelstein et al. [2014\)](#page-36-0). CF structure was established for $FeCr₂O₄$ composition in meteorites (Chen et al. 2003) and suggested a new CaCr₂O₄ oxide found in ultramafic lower-mantle microxenolite (Kaminsky et al. [2015;](#page-37-0) see Sect. 4.7.3).

The CF phase has a chain framework structure formed by edge-sharing $AIO₆$ octahedra along the c-axis direction similar to the NAL phase, but it only forms one kind of 'tunnel' formed by the linked double chains being in eight-fold coordination, where sodium atoms are located (Yamada et al. [1983](#page-42-0)) (Fig. [5.9\)](#page-17-0). This model is

Table 5.5 Representative compositions of the CE phase from experimental matic associations ($\frac{arg}{n}$) Table 5.5 Representative compositions of the CF phase from experimental mafic associations (wt%)

Note n.a. not analysed Note n.a. not analysed

based on structural parameters for the MgAl₂O₄ end-member (Kojitani et al. [2007\)](#page-37-0), whereas the CF-type phase in basaltic compositions of the lower mantle can exist within the NaAlSiO₄–MgAl₂O₄ system.

Experimental studies on the NaAlSiO₄–MgAl₂O₄ system in a laser-heated DAC coupled with in situ synchrotron XRD measurements demonstrated that the CF phase is the high pressure form of the NAL phase, with a transition pressure between 34 and 43 GPa (Ono et al. [2009;](#page-39-0) Imada et al. [2011](#page-36-0)). On the other hand, the CF phase exists in association with the NAL phase at pressures of up to 50 GPa (Ono et al. [2009;](#page-39-0) Ricolleau et al. [2010;](#page-40-0) Imada et al. [2011](#page-36-0)), and at much higher pressures of up to 120 GPa and 2300 K, in the potassium-rich system. The calcium-bearing CF phase is also found to coexist with the NAL phase (Kato et al. [2013\)](#page-37-0). Thus far, these studies have shown that the relative stability of NAL and CF phases is a complex, sensitive function of $P-T$ and composition. While pressure tends to favour stability of the CF phase, enrichment in Mg or alkalis tends to favour the NAL phase (Mookherjee et al. [2012\)](#page-39-0).

The NAL to CF phase transition causes an increase in density and bulk modulus (Guignot and Andrault [2004](#page-36-0)) and a change in compressional wave velocity by −0.2% and shear wave velocity by +0.9% (Kawai and Tsuchiya [2012\)](#page-37-0); the shear wave velocity increases by 2.5% across phase transition at \sim 40 GPa (Dai et al. [2013\)](#page-35-0). The change in the elastic properties and shear wave anisotropy across the NAL to CF phase transition is proposed to be significant enough to be seismically detectable (Kawai and Tsuchiya [2012;](#page-37-0) Dai et al. [2013](#page-35-0)). A study on the physical properties of the aluminous phases performed by Mookherjee et al. ([2012\)](#page-39-0), demonstrated that the sound wave velocities of the NAL phase is significantly lower than the CF phase and other major lower mantle phases. The existence of such a lower-velocity NAL phase in the deeply subducted MORB materials can compensate for the faster sound velocity of silica $(SiO₂)$ and its high-pressure polymorphic phase(s) in the lower mantle.

5.3.4 Calcium Aluminium Silicate with Ba-Ferrite Structure (CAS Phase)

A Ca–Al silicate (CAS) with an end-member chemical formula of $CaAl₄Si₂O₁₁$ containing some amount of water was first reported by Irifune et al. [\(1994](#page-37-0)) in experiments performed at 17.5–24 GPa and 1500–1900 °C, and later by Gautron et al. ([1996\)](#page-36-0) at 14 GPa and 1400–1500 °C. The CAS phase has a hexagonal unit cell with lattice parameters $a = 5.4 \text{ Å}$ and $c = 12.7 \text{ Å}$ and a space group, most likely, $P6_3/mmc$ with $Z = 2$, analogous to hexagonal Ba-ferrite BaTi₂Fe₄O₁₁. It has a six-layer, close-packed structure, reasonable for a phase stable under the conditions of the transition zone (Gautron et al. [1997](#page-36-0)).

Subsequently, the CAS phase was synthesized as an Al-rich phase using anhydrous MORB materials in multi-anvil apparatus at 25.5 GPa (Hirose and Fei [2002;](#page-36-0) Litasov and Ohtani [2005](#page-38-0)). In these experiments, the NAL and CF phases appeared at higher pressures, above 27–27.5 GPa, coexisting with bridgmanite, stishovite and CaSi-perovskite. The NAL phase synthesized at 26 GPa was stable below 2030 °C but transformed to the CAS phase at higher temperatures. The CAS phase does not appear below the solidus temperature in the presence of bridgmanite. It remains unclear whether the CAS phase can present in the lower mantle.

In the natural environment, a single sample of CAS in association with kyanite was reported from an inclusion in Juina-5 kimberlite pipe, Brazil (Thomson et al. [2014\)](#page-41-0). The natural CAS is poorer in Mg and Al (with $Al/Si = 1.027$) than when obtained from experiments (Table [5.6\)](#page-19-0).

5.4 Partitioning of Aluminium Between Bridgmanite and Aluminous Phases

Aluminium in the lower-mantle mafic association has been observed to be concentrated in three major phases: bridgmanite, NAL phas, and CF phase. The concentration of Al in bridgmanite of the mafic association varies from 11 to 13 wt% A_2O_3 at 26 GPa to 24 wt% A_2O_3 at 100 GPa (see Fig. [5.1\)](#page-2-0), independently of the magnesium index value (mg), which is rather constant. Al-rich NAL phase with 38.9–53.8 wt% Al₂O₃ (average: 46.2 ± 4.9 wt% Al₂O₃) probably incorporated its aluminium from majorite in the lower mantle, and shows a strong correlation of its $A₂O₃$ concentration with SiO₂ (Fig. [5.7a](#page-13-0)). The ratio Al/Si_{at} in the NAL phase strongly decreases from 4.7 at 25 GPa to 1.73 at 33 GPa (Fig. [5.7](#page-13-0)b). Al concentration in the CF phase varies from 32.8 to 45.1 wt% Al₂O₃ (average 39.4 \pm 3.6 wt $\%$ Al₂O₃), independently of Si and pressure values. This implies that with redistribution of Al between bridgmanite and NAL occurs with increasing pressure, as subducting slabs penetrate deeper into the lower mantle.

Table [5.7](#page-20-0) shows Si/Al ratio values in coexisting bridgmanite in the NAL phase from the experimental solidus associations, as well as Al/Si ratios and aluminium partitioning coefficient values. The partition coefficient of aluminium between bridgmanite and the

Table 5.6 Compositions of synthetic and natural CAS phase $(wt\%)$ Table 5.6 Compositions of synthetic and natural CAS phase (wt%)

Note n.a. not analysed ž.

Table 5.7 SI/AI ratio values (atomic) and partition coefficient of AI in coexisting bridgenamite and NAI. Table 5.7 Si/Al ratio values (atomic) and partition coefficient of Al in coexisting bridgmanite and NAL

Note *The associations with liquid are not included Note *The associations with liquid are not included

Fig. 5.10 Dependence of aluminium partition coefficient $K_{D}^{\text{Brd-NAL}}$ of pressure

NAL phase defined as $K_{\text{D}}^{\text{Brd-NAL}} = (Al/Si)^{\text{Brd}}/(Al/Si)^{\text{NAL}}_{at}$ varies from 0.10 to 0.26, asymptotically increasing within a pressure range of 25–44 GPa (Fig. 5.10). This confirms that the Al enrichment in bridgmanite with increasing pressure is a result of the dissolution of the NAL phase, while the Al concentration in the CF phase, which replaces the NAL phases at pressures of 40–50 GPa, remains constant. The Al concentration in bridgmanite, after reaching a maximal concentration of \sim 24–25 wt%, also remains constant.

5.5 Hydrous Aluminous Phases

5.5.1 General

Subducting slabs can contain significant amounts of water in hydrous minerals in hydrothermally altered basalts (e.g., serpentine) and sediments (e.g., clay minerals). High-pressure investigations of the $Al_2O_3-SiO_2-H_2O$ system relevant to subducting slab materials have found that three hydrous mineral phases, phase Egg, δ -AlOOH and topaz-OH can possibly become stable at $P-T$ conditions corresponding to the lower mantle. Initially, however, these three phases were reported to be stable at relatively lower *P–T* conditions at 5.5–21 GPa and temperatures up to 1500 °C (Wunder et al. [1993](#page-42-0); Schmidt [1995](#page-40-0); Ono [1999](#page-39-0); Suzuki et al. [2000](#page-41-0)). Subsequently, they were experimentally proven to be stable at greater depths relevant to the lower-mantle conditions. Phase Egg and δ -AlOOH are found as inclusions in natural diamonds.

5.5.2 Phase Egg

A hydrous aluminous silicate $AISiO₃OH$ was first observed in high-pressure experiments on hydrous basaltic systems at pressures exceeding 10 GPa. The first water-containing high-pressure mineral, initially suggested as $Al_5Si_5O_{17}OH$ (Egglton et al. [1978](#page-36-0)):

$$
AIOOH + SiO2 \rightarrow AISiO3OH.
$$

diaspore + stishovite \rightarrow phase Egg

This phase was subsequently synthesized from topaz-OH + stishovite mixture at 12 GPa, and has been later termed as phase Egg to give credit to Professor R.A. Egglton who first discovered the phase (Egglton et al. [1978;](#page-36-0) Schmidt [1995\)](#page-40-0). High P –T stability of the phase was experimentally extended to 20 GPa and 1700 $^{\circ}$ C (Ono [1999\)](#page-39-0), then to 23 GPa and 1200 °C (Sano et al. [2004](#page-40-0)), then to 25 GPa and \sim 1500 °C (Boffa Ballaran et al. [2010\)](#page-35-0), and later to 26 GPa and 1460 °C (Pamato et al. [2015](#page-40-0)). Most recently Fukuyama et al. ([2017\)](#page-36-0) experimentally established the stability limit of phase Egg at 23.5 GPa, suggesting that differences in previous results were caused by the use of different pressure scales. They concluded that phase Egg is stable only within the transition zone (Fukuyama et al. [2017](#page-36-0)). However, other experiments and equation of state study of this phase demonstrate that, despite the presence of hydrogen in its composition, phase Egg is stable up to 40 GPa (Vanpeteghem et al. [2003](#page-41-0)). For this reason, the presence of phase Egg and stishovite in natural samples may be suggestive not only of the origin in the

transition zone but in the lower mantle as well, and phase Egg may be considered as a major carrier of water in the deep Earth. A natural association of stishovite and phase Egg was observed in a diamond from the Rio Soriso, Juina area, Brazil (Wirth et al. [2007](#page-41-0)) (Fig. [5.11\)](#page-23-0). Both minerals were

identified with the observation of in situ Raman spectroscopy and electron diffractometry, such that original mineral structures were not altered during the examination. Phase Egg forms an aggregate of euhedral, cubic and cubo-octahedral grains that

are 20–30 to 200–300 nm in size (Fig. [5.11](#page-23-0)b). The bulk Si/Al ratios of the two analyzed natural grains are approximately 1:1 (46.74–53.84 at.% and 45.63– 52.60 at.%). A minor admixture of Fe $(0.29-1.52 \text{ at.} %)$ is present in its composition. The presence of OH in phase Egg was identified using the pre-peak at approximately 528 eV in the electron energy-loss spectrum and the OH-stretching band between 1800 and 3000 cm−¹ in the Raman spectra. Numerous irregularly shaped pores are present in the stishovite matrix (Fig. [5.11](#page-23-0)a). They contained a fluid, which was released to the vacuum during the focused ion beam milling of the specimen for the transmission electron microscopy analysis. Quench products of the fluid contain minor amounts of K, Ca, Ba, P, S, Cl and F that were also detected along the walls of the pores. This porosity possibly reflects the initial stage of degassing of phase Egg under conditions close to its pressure-dependent disintegration. In some of the Raman spectra taken from phase Egg, characteristic spectral features of δ -AlOOH

Fig. 5.11 Electron-microscope images of phase Egg. a Porous aggregate of stishovite and phase Egg. b Single cubic crystal of phase Egg. From Wirth et al. [2007.](#page-41-0) © Earth and Planetary Science Letters (Elsevier)

phase located near 2150 and 2500 cm−¹ were detected, analogous to the spectral characteristics of experimentally produced δ -AlOOH phase (Ohtani et al. [2001b\)](#page-39-0). The coexistence of δ -AlOOH phase with phase Egg indicates the most likely pressure interval for their joint formation to be 23–27 GPa, i.e., in the uppermost regions of the lower mantle (Sano et al. [2004](#page-40-0)).

The crystal structure of phase Egg, obtained in experimental studies, was deter-mined, using powder XRD, to be monoclinic (Fig. [5.12\)](#page-24-0) with a space group $P2_1/n$, a unit cell volume $V_0 = 212.99(1)$ Å³, and zero-pressure density of 3.74 g/cm³ (Schmidt et al. [1998\)](#page-40-0). Most grains of natural phase Egg are monoclinic as well, but

some analyzed natural grains are tetragonal ($a = 0.7435$ nm and $c = 0.706$ nm), suggesting that the tetragonal polymorph may be a relic of the original phase Egg.

The monoclinic structure of phase Egg resembles that of stishovite and δ -AlOOH, containing an ordered arrangement of SiO₆ and AlO₆ octahedra linked by both edge- and corner-sharing (Fig. 5.12). Similar to δ -AlOOH (see Sect. 5.5.3), the crystal structure of phase Egg is characterized by strong hydrogen bonding, which can have great effects on many of the properties in the mineral including elasticity, rheology and electrical conductivity. Hydrogen bonding, which is the electrostatic attraction between two polar groups (in this case hydrogen and oxygen atoms), also plays an important role in stabilizing hydrous and nominally anhydrous phases by lowering the enthalpy, or increasing the entropy, or both (e.g., Prewitt and Parise [2000\)](#page-40-0). The strong hydrogen bonding in phase Egg and δ -AlOOH (as well as in DHMSs) may be an important factor in stabilizing these structures at high pressures and temperatures.

5.5.3 Hydrous d-AlOOH and e-FeOOH Phases

Aluminium oxyhydroxide, the hydrous δ -AlOOH phase, which is a high-pressure polymorph of diaspore (α -AlOOH) and boehmite (γ -AlOOH), was initially synthesized in association with majorite and phase Egg at \sim 21 GPa and 1000 °C (Suzuki et al. [2000\)](#page-41-0). Subsequently, it has been found to be stable at pressures above 30 GPa and temperatures of 1000–1200 °C, and thus it may serve (like phase Egg) as a water carrier in the cold subducted slabs, transporting water into the lower mantle (Ohtani et al. [2001a](#page-39-0), [b](#page-39-0); Sano et al. [2004](#page-40-0)). Later, in situ XRD experiments in the diamond anvil cell were conducted to determine the stability field of δ -AlOOH over the entire P–T conditions of the lower mantle. δ -AlOOH was observed to exist at pressures from 33 to 134 GPa at 1350–2300 K, indicating that this hydrous mineral in the sediment layer of a subducting slab could transport water into the

core-mantle boundary (Sano et al. [2008](#page-40-0); Mashino et al. [2016\)](#page-38-0). The experiments by Ohira et al. [\(2014](#page-39-0)) also confirmed that a hydrous δ -AlOOH phase could coexist with Al-depleted bridgmanite or post-perovskite under the P–T conditions relevant to the cold slab geotherm. δ -AlOOH was identified in geological samples to exist in association with phase Egg (Wirth et al. [2007\)](#page-41-0).

Theoretical calculations further predicted the pressure-induced phase transition of δ -AlOOH into a new cubic phase at approximately 170 GPa (Tsuchiya and Tsuchiya [2011](#page-41-0)). Therefore, δ -AlOOH phase is likely to be the highest pressure phase among all hydrous mineral compounds, since these hydrous materials tend to dissociate and release H_2O at pressures of around a few tens of GPa (Ohtani et al. [2001b\)](#page-39-0).

Pure δ -AlOOH has an orthorhombic symmetry with a non-centrosymmetric space group $P2_1$ nm that is very close to that of $SiO₂$ in the CaCl₂-type (Suzuki et al. [2000;](#page-41-0) Komatsu et al. [2006](#page-37-0); Vanpeteghem et al. [2007\)](#page-41-0) (Fig. 5.13a). The structure contains both $AIO₆$ and $HO₆$ octahedra, and the oxygen anions form a slightly distorted hexagonal close-packed arrangement perpendicular to the b direction

Fig. 5.13 Orthorhombic crystal structure models of δ -AlOOH based on $P2_1$ _{nm} and Pnnm space groups. Dotted lines are partially ordered H bonds. Unit cell parameters from Komatsu et al. ([2011\)](#page-37-0)

(Wicks and Duffy [2016](#page-41-0)). A single-crystal XRD study indicates that δ -AlOOH transforms to the CaCl₂-type structure (space group *Pnnm*) at \sim 8 GPa (Kuribayashi et al. [2013](#page-38-0)) (Fig. [5.13](#page-25-0)). This structure consists of corner-sharing chains of edge-sharing octahedra $(AIO₆,$ in this case). There is no discontinuity in unit cell volume across this displaceable transition. The main difference between the two structures lies in the disordering of hydrogen atoms in the high-pressure structure. The substitution of Mg^{2+} and Si^{4+} into δ -AlOOH also promotes the transformation of $P2_1$ nm structure (Suzuki et al. [2000](#page-41-0); Kudoh et al. [2004](#page-37-0); Komatsu et al. [2011\)](#page-37-0). The disordered cations, Mg^{2+} and Si^{4+} , with different valences than Al^{3+} may fluctuate hydrogen positions, causing the disordered hydrogen to change its symmetry. As a result, Mg–Si-containing δ -AlOOH phase exhibits the *Pnnm* structure (Komatsu et al. [2011](#page-37-0)) or Pnn2 structure (Kudoh et al. [2004](#page-37-0)).

Pressure-induced hydrogen bond symmetrization starts to occur in δ -AlOOH at 28 GPa (Tsuchiya et al. [2002\)](#page-41-0), similar to that reported for phase Egg. Since the crystal structure with the symmetric hydrogen bonding has a smaller volume than the asymmetric hydrogen bond counterpart, the occurrence of the hydrogen bond symmetrisation can further stabilize hydrous materials to higher-pressure conditions. It is conceivable that higher-pressure hydrous phases would also be expected to have symmetric hydrogen bonds. The details of the hydrogen bonding in δ -AlOOH phases has been extensively discussed, especially with regards to the formation of symmetric hydrogen bonds, but there is currently no consensus on this issue (Tsuchiya et al. [2002](#page-41-0); Panero and Stixrude [2004](#page-40-0); Li et al. [2006;](#page-38-0) Xue et al. [2006;](#page-42-0) Vanpeteghem et al. [2007](#page-41-0); Sano‐Furukawa et al. [2008\)](#page-40-0).

The predicted ultrahigh-pressure cubic phase has a cubic pyrite-type $AIO₂$ framework with interstitial H atoms forming symmetric hydrogen bonds, whose symmetry is assigned to the space group $Pa\overline{3}$.

Other cation substitutions, such as Fe and Ti in the oxyhydroxide, have also been suggested based on high $P-T$ experiments corresponding to the transition zone conditions: at pressures of 16 GPa and temperatures of 1000–1600 °C (Nishihara and Matsukage [2016](#page-39-0)). In the studied conditions, two stable iron-titanium oxyhydroxide phases were identified whose estimated composition is $(FeH)_{1-x}Ti_xO_2$. One is the Fe-rich solid solution $(x < 0.23)$ with ε -FeOOH-type crystal structure (orthorhombic, P21 nm); and the other is the more Ti-rich solid solution ($x > 0.35$) with α -PbO₂-type structure (orthorhombic, *Pbcn*). The ϵ -FeOOH-type phase is stable up to \sim 1100 °C, whereas the Ti-rich α -PbO₂-type phase is stable up to 1500 °C for a composition of $x = 0.5$ and at least to 1600 °C for $x = 0.75$. These phases may be stable within subducting basaltic slabs into the transition zone, but their stability in the lower mantle remains uncertain and requires further clarifications in future studies.

In iron-rich compositions, a complete Fe–Al substitution may occur. e-FeOOH, the high-pressure polymorph of goetite (α -FeOOH), is isostructural with δ -AlOOH, and has a slightly distorted rutile structure with corner-linked single bands of edge-shared octahedra parallel to the c-axis (Otte et al. [2009\)](#page-39-0). e-FeOOH undergoes a spin transition beginning at 49 GPa and completing at \sim 65 GPa, which is

assisted by symmetrization of hydrogen bonds during the transition from $P2_1$ _{nm} to Pnnm. (Gleason et al. [2008](#page-36-0), [2013;](#page-36-0) Reagan et al. [2016](#page-40-0)).

5.5.4 Topaz-OH

Topaz-OH $Al_2SiO_4(OH)_2$, the OH end-member of the F-OH topaz, has been synthesized at a pressure of 6–10 GPa from diaspore $(\alpha$ -AlOOH) and phase Π $A₁₃S₁₂O₇(OH)₃$ starting materials as a result of their reaction with SiO₂. This phase is shown to be stable up to 13 GPa and 1500 °C (Wunder et al. [1993](#page-42-0); Xue et al. 2006), after which it reacts with SiO₂ to form phase Egg AlSiO₃OH (Egglton et al. [1978;](#page-36-0) Schmidt [1995,](#page-40-0) 1998; Ono [1999](#page-39-0)). At high temperatures, it undergoes a dehydration reaction to form kyanite (Wunder et al. [1993](#page-42-0)). At pressures above 14 GPa, a high-pressure topaz-OH II variety, with significant cation disorder occurs (Kanzaki [2010](#page-37-0)).

The structure of topaz-OH shows close agreement with that of F-rich topaz, containing tetrahedral $SiO₄$ and linked $AlO₄(OH)₂$ octahedra. It was first refined with the space group *Pbnm* from single-crystal XRD by Wunder et al. ([1993\)](#page-42-0), although the hydrogen positions could not be located. A subsequent single-crystal XRD study (Northrup et al. [1994](#page-39-0)) reported that the single H position found in F-rich topaz is split into two nonequivalent, half-occupied H positions in topaz-OH. They also suggested that the structure is non-centrosymmetric and might have a long-range ordered H distribution with a reduction of symmetry to $Pbn2_1$. Each H is associated with three O atoms in an irregular trifurcated H-bond arrangement. A lattice energy calculation and molecular dynamics simulation study (Churakov and Wunder [2004\)](#page-35-0) predicted four nonequivalent crystallographic H positions in topaz-OH that are in dynamic exchange at ambient conditions and suggested instantaneous configuration of H violates in all symmetry elements, whereas the space- and time-averaged structure possesses *Pbnm* symmetry. Monoclinic structure $P2₁/c$ is also suggested for topaz-OH (Mookherjee et al. [2016\)](#page-39-0).

The stability of topaz-OH under high P–T conditions has not been studied.

5.6 Dense Hydrous Magnesium Silicates (DHMSs)

5.6.1 General

Numerous experimental high-pressure works were carried out on DHMSs of relevant mantle compositions, showing that they can be stable with increasing pressure at characteristic slab temperatures. Ringwood and Major ([1967\)](#page-40-0) recognized the first of these DHMSs in experiments performed in the $MgO-SiO-H₂O$ system under pressures of up to 18 GPa, and denoted these phases as A $[Mg_7Si_2O_8(OH)_6]$, B $[Mg_{12}Si_4O_{19}(OH)_6]$, and C, known also as superhydrous Phase B

 $[Mg_{10}Si_3O_{14}(OH)_4]$. The naming system for the DHMS phases proposed by them was later adopted for naming a whole sequence of 'alphabetic' DHMS phases. Phases termed D (Yamamoto and Akimoto [1977](#page-42-0); Liu [1986](#page-38-0), [1987\)](#page-38-0), E, F (Kanzaki [1989,](#page-37-0) [1991;](#page-37-0) Kudoh et al. [1993](#page-37-0), [1995\)](#page-38-0), and G (Ohtani et al. [1997\)](#page-39-0) were subsequently reported; however the latter two, F and G, appear to be similar if not identical to Phase D (Yang et al. [1997;](#page-42-0) Frost and Fei [1998](#page-36-0), [1999\)](#page-36-0). These phases contain 2.14– 18 wt% H_2O and may act as water carriers in subducting hydrated plates to >35 GPa (Yang et al. [1997;](#page-42-0) Ohtani et al. [2004\)](#page-39-0), in addition to hydrous aluminium phases. Most of these phases appeared to be unstable under lower-mantle pressure conditions (e.g., Kanzaki [1991\)](#page-37-0), with the exception of the following three: superhydrous Phase B, Phase D, and the recently discovered Phase H.

5.6.2 Superhydrous Phase B

'Superhydrous' Phase B $[Mg_{12}Si_4O_{19}(OH)_6]$, containing as much as 5.8 wt% H₂O, is stable at $P-T$ conditions relevant to the top of the lower mantle and decomposes into periclase, bridgmanite and Phase D at approximately 30 GPa and high temperatures (Shieh et al. [1998](#page-40-0)). On the other hand, the following reaction sequences were observed with increasing temperature from 1200 to 1600 °C in the pressure range from 18–20 GPa: superhydrous Phase $B \rightarrow$ Phase B + liquid \rightarrow wadsleyite + periclase + liquid \rightarrow anhydrous Phase B + periclase + liquid (Ohtani et al. 2001a).

5.6.3 Phase D

Initially, the term Phase D was used by Yamamoto and Akimoto [\(1977](#page-42-0)) to describe the hydroxyl chondrodite, which was synthesized at pressures between 3 and 12 GPa. However, Liu ([1987\)](#page-38-0) reapplied the name Phase D to describe a phase discovered as a dissociation product of $Mg_3Si_2O_5(OH)_4$ serpentine at pressures above \sim 20 GPa (an equivalent to phases F and G). It has the ideal formula $MgSi₂O₆H₂ (Mg-Phase D)$, but is typically non-stoichiometric. Phase D contains nominally 10.1 wt% H_2O (varying from 10 to 18 wt%) and varies significantly in stoichiometry depending on synthesis conditions (Frost and Fei [1998](#page-36-0)) with Mg/Fe ratio from 0.56 to 0.71 within the range of 16–27 GPa and 1073–1573 K; the Mg/Si ratio and H₂O concentration decrease with increasing temperature and increase with increasing pressure (Shinmei et al. [2008\)](#page-41-0). Along with most other DHMS phases, Phase D synthesized from the starting samples containing complex peridotitic bulk compositions have been found to contain up to 7 wt% $\mathrm{Al}_2\mathrm{O}_3$ and 5 wt% FeO (Frost [1999;](#page-36-0) Litasov et al. [2007b](#page-38-0), [2008;](#page-38-0) Rosa et al. [2012;](#page-40-0) Ganskow and Langenhorst [2014\)](#page-36-0), suggesting that natural Phase D samples contain not only water but also some Al and Fe. Al^{3+} substitutes for Si^{4+} , Fe^{2+} substitutes for Mg, and Fe^{3+} may occupy Si and Mg sites (Ganskow and Langenhorst

Fig. 5.14 Crystal structure of Al-free trigonal phase D (space group $P-31m$) in different axis orientations. Lattice parameters from Yang et al. [\(1997](#page-42-0))

[2014\)](#page-36-0). Available data demonstrate that the $Fe^{3+}/\Sigma Fe$ ratio of Fe-bearing phase D determined by conventional Mössbauer spectroscopy and EEL spectroscopy is 0.60–0.94, and iron is not involved in water-substitution mechanisms (Chang et al. [2013](#page-35-0); Ganskow and Langenhorst [2014](#page-36-0)).

Phase D has a trigonal structure with the space group $P-31$ m, based on a hexagonal close-packed array of oxygen atoms (Kudoh et al. [1997](#page-38-0); Yang et al. [1997;](#page-42-0) Xue et al. [2008\)](#page-42-0), similar to stishovite (Fig. 5.14). All cations except H are in the octahedrally coordinated sites. The MgO_6 and SiO_6 octahedra occur in separated, alternating layers stacked along the c direction (Fig. 5.14c), with Mg and Si at the M1 and M2 sites, respectively, and with the octahedral sites corresponding to the $2c$ (M3) and 1b Wyckoff positions being vacant (Yang et al. [1997\)](#page-42-0). The MgO₆ and $SiO₆$ octahedra share edges with neighbours to form rings in layers perpendicular to the c axis (Fig. $5.14b$). The SiO₆ octahedra form gibbsite-like layers, in which each octahedron shares three edges, leaving one third of the octahedral sites vacant. The MgO₆ octahedra lie above and below vacant sites in the $SiO₆$ layer and are corner-linked to the Si octahedra, and so two thirds of the Mg octahedral sites are empty (Wicks and Duffy [2016\)](#page-41-0). Charge balance is achieved by hydrogen anions. The hydrogen positions are located in the $MgO₆$ layers, with O–H bonds facing away from $SiO₆$ octahedra.

Phase D is the only DHMS structure reported to date that contains Si exclusively in octahedral co-ordination. Owing to this, the calculated density of Phase D is 3.50 $g/cm³$, which is substantially greater than any other DHMS found thus far, including Phase A (2.96 g/cm³), Phase B (3.37 g/cm³), superhydrous Phase B (3.21 g/cm³), Phase E (2.82 g/cm³), and Phase F (Yang et al. [1997](#page-42-0) and references therein). As a result of the layered nature of the structure, with strong $SiO₆$ layers alternating with weaker $MgO₆$ octahedra resulting in enhanced compressibility along the c axis, Phase D exhibits anisotropic compression (Frost and Fei [1999\)](#page-36-0). Theoretical calculations of the elastic constants indicate that the anisotropy decreases with depth but significant anisotropy in seismic wave velocities is retained to the highest pressures (Mainprice et al. [2007\)](#page-38-0).

At pressures exceeding 35–40 GPa Phase D, like phase Egg and δ -AlOOH (see above), experiences hydrogen bond symmetrization (Tsuchiya et al. [2005](#page-41-0); Shinmei et al. [2008](#page-41-0)). A symmetric hydrogen bond is one in which the hydrogen atom is located at the midpoint between the two neighbouring oxygen atoms, rather than the asymmetric O–H distances that characterize a conventional hydrogen bond. This serves stabilizing hydrous Phase D (as well as phase Egg and δ -AlOOH) under ultrahigh-pressure conditions. This pressure-induced hydrogen bond symmetrization has a significant effect on the compression behaviour of Phase D. The bulk modulus increases by about 20% with this structural change (Tsuchiya et al. [2005](#page-41-0)). A powder XRD study showed evidence of such an anomaly (Hushur et al. [2011\)](#page-36-0), but a more recent high-resolution single-crystal XRD study to 65 GPa found no evidence for a bulk modulus anomaly, at least to this pressure (Rosa et al. [2013a,](#page-40-0) [b\)](#page-40-0). No evidence for hydrogen bond symmetrization was found in an infrared spectroscopic study (Shieh et al. [2009](#page-40-0)) or in other theoretical calculations (Mainprice et al. [2007](#page-38-0)).

Fe-bearing Phase D was reported to undergo the electron spin transition of $Fe³⁺$ at high pressure, which is shown to cause a softening in the bulk modulus and bulk sound velocity of the crystal (Chang et al. [2013](#page-35-0)). Recently, a two-stage HS-LS spin transition involving both Fe^{2+} and Fe^{3+} was reported to occur in Phase D at high pressures (Wu et al. [2016b](#page-42-0)). This transition of iron occurs for Fe^{2+} at 37–41 GPa and for $Fe³⁺$ at 64–68 GPa. These transitions are accompanied by an increase in density and a significant softening in the bulk modulus and bulk velocity at their respective pressure range. The occurrence of the dense low-spin FeAl-bearing phase D with relatively high velocity anisotropies in deep-subducted slabs can potentially contribute to small-scale seismic heterogeneities. However, hydrogen symmetrization and spin transitions in Phase D may overlap with each other, and their interplay at high pressures is not yet fully understood (Wu et al. [2016b\)](#page-42-0).

The stability field of Phase D extends from 20 to 50 GPa, independent of temperature from \sim 1300–1700 K (Shieh et al. [1998,](#page-40-0) [2000;](#page-41-0) Frost and Fei [1998,](#page-36-0) [1999;](#page-36-0) Tsuchiya et al. [2005;](#page-41-0) Litasov et al. [2008](#page-38-0); Shinmei et al. [2008;](#page-41-0) Hushur et al. [2011;](#page-36-0) Ghosh and Schmidt [2014;](#page-36-0) Nishi et al. [2014](#page-39-0); Walter et al. [2015](#page-41-0)). Pure Phase D is stable only at relatively lower temperatures and can break down at temperatures lower than those of the typical mantle geotherm (e.g., Nishi et al. [2014\)](#page-39-0). In experiments, Phase D decomposes at 44–50 GPa, corresponding to a depth of 1250–1400 km, into an assemblage of nominally anhydrous phases consisting of low-Al bridgmanite and H_2O (Shieh et al. [1998](#page-40-0); Walter et al. [2015\)](#page-41-0).

The more stable, Al-rich variety of Phase D with up to 50 wt% $A1_2O_3$ was synthesized in association with phase Egg in a simplified basaltic bulk composition at lower mantle conditions at 25 GPa and \sim 1500 °C, which is significantly higher than temperature stability for pure Phase D, allowing the Al-rich Phase D to remain stable and host H_2O in the subducting oceanic crust in the Earth's lower mantle (Boffa Ballaran et al. [2010\)](#page-35-0). This idea was confirmed by experiments that demonstrated that Al-rich Phase D $A_2SiO_6H_2$ in association with phase Egg and stishovite, has very high thermal stability extending to over 2000 \degree C at 26 GPa, i.e., higher than the expected geotherm of the normal lower mantle (Pamato et al. [2015\)](#page-40-0). In the subsolidus, the Al-rich Phase D is stable to \sim 55 GPa (Walter et al. [2015\)](#page-41-0). The Al-rich Phase D may have an admixture of Fe, which actually decreases its stability (Ghosh and Schmidt [2014](#page-36-0)). According to experimental data, Al/Fe ratios in Phase D and in coexisting bridgmanite are roughly equal (Pamato et al. [2015](#page-40-0)).

The Al-rich Phase D was suggested initially to have the same structure as Mg-Phase D (space group $P-31m$), but with a more disordered cation distribution (Boffa Ballaran et al. [2010\)](#page-35-0). A more recent study by Pamato et al. ([2015\)](#page-40-0) attributes its structure to hexagonal with the space group $P6\sqrt{mcm}$, where Al replaces Mg and all six octahedral sites become partially occupied by a random and disordered distribution of Si and Al, such that M1 and M4 become equivalent (as do M2 and M3), resulting in an increase in symmetry. One of the main differences between the structures is that the Al–Si disorder in Al-rich Phase D results in essentially an undistorted octahedra of similar size, whereas in Mg-Phase D, large MgO distances cause the octahedra to be strongly distorted. The lesser extent of the octahedral distortion is likely to stabilize the Al-rich Phase D relative to its Mg-bearing counterpart. On the other hand, Bindi et al. ([2015\)](#page-35-0) reported no change in the unit-cell volume in Al-rich Phase D with respect to pure $MgSi₂O₆H₂$. Earlier, Boffa Ballaran et al. [\(2010](#page-35-0)) and Komatsu et al. [\(2011](#page-37-0)) observed that the Tscher-mak's $Si^{4+} + Mg^{2+} \rightarrow 2Al^{3+}$ substitution, in hydrous silicates, does not induce either an enlargement of the unit cell or a lengthening of the octahedral bond distances.

An important factor in the charge distribution in the Al-rich Phase D is the large degree of Al–Si disorder. Although the contribution to the configurational entropy due to Al–Si disorder has a stabilizing influence on low-pressure minerals, such as feldspar, the preference of Al for the octahedral site over the tetrahedral coordination at pressures of a few gigapascals means that this effect plays only a minor role in much of the upper mantle and transition zone, such disorder seems to become important again in the lower mantle, where Si and Al both exist exclusively in the octahedral coordination. The Al-rich Phase D may be the first member of a new class of completely disordered hydrous aluminosilicates that consequently have high thermal stabilities (Pamato et al. [2015\)](#page-40-0) (Fig. [5.15\)](#page-32-0).

Fig. 5.15 Crystal structure of the Al-rich hexagonal phase D (space group $P6_3/mcm$) in different axis orientations. Lattice parameters from Pamato et al. ([2015](#page-40-0))

5.6.4 Phase H

Using theoretical calculations, Tsuchiya ([2013\)](#page-41-0) predicted that Phase D should undergo a phase transition to a new hydrous (with \sim 15% of water) Phase H with an ideal formula of MgSiO₄H₂, in association with SiO₂ (MgSi₂O₆H₂ \rightarrow $MgSiO_4H_2 + SiO_2$) at pressures exceeding \sim 40 GPa. Subsequently, multi-anvil apparatus experiments using sintered diamond anvils and in situ X-ray measurements showed that Phase D transforms to an assemblage containing Phase Hat pressures above 44–50 GPa (Nishi et al. [2014;](#page-39-0) Ohtani et al. [2014;](#page-39-0) Walter et al. [2015\)](#page-41-0).

The crystal structure of Phase H was initially suggested as non-centrosymmetric orthorhombic with a space group $P2_1$ _{nm} similar to the one of the high-pressure hydrous phase δ -AlOOH (Suzuki et al. [2000](#page-41-0)). Later, Tsuchiya [\(2013](#page-41-0)) suggested that it should have a monoclinic symmetry with a space group $P2/m$. However, single-crystal XRD study on Phase D synthesized at 45 GPa unequivocally showed the orthorhombic symmetry with the CaCl₂-structure (space group $Pnnm$) (Bindi

Fig. 5.16 Orthorhombic crystal structure model of phase H (space group *Pnnm*) in different axis orientations. Lattice parameters from Bindi et al. ([2014\)](#page-35-0)

et al. [2014](#page-35-0)). In this structure, Mg^{2+} and Si^{4+} are disordered over the octahedral sites. Distorted Mg and Si octahedra are charge compensated for by H^+ ions, which may occupy one of two positions depicted between the octahedra (Fig. [5.16\)](#page-32-0). When cations with different valence states (i.e., Mg^{2+} , Al^{3+} , and Si^{4+}) become disordered at the octahedral sites, a fluctuation of the hydrogen position occurs to maintain the charge neutrality. Similar to δ -AlOOH, phase Egg and Phase D, the hydrogen bond symmetrization also occurs in Phase H at high pressure, and the hydrogen bond of this phase is already symmetric at the transition from Phase D at 40 GPa (Tsuchiya [2013\)](#page-41-0). Like in Phase D, the hydrogen bond symmetrization probaby helps to stabilize hydrous materials under high-pressure conditions.

Tsuchiya [\(2013](#page-41-0)) predicted that pure Phase H had a narrower stability field near 52 GPa at static 0 K, above which it would dissociate into bridgmanite and $H₂O$ (ice VIII). This coincides with experiments performed by Ohtani et al. [\(2014](#page-39-0)), in which the stability field of pure Phase H was limited at $P-T$ conditions up to 60 GPa and \sim 1600 K, above which Phase H decomposes into bridgmanite and H₂O.

Considering that Mg- and Si-containing δ -AlOOH has the same as Phase H *Pnnm* structure (Komatsu et al. [2011](#page-37-0)), one may suggest that Phase H (MgSiO₄H₂) and δ -AlOOH ($AIAIO₄H₂$) could form a complete solid solution at high pressures (Bindi et al. [2014](#page-35-0)). Experiments have shown that intermediate Al-rich compositions can be synthesized together with perovskite and post-perovskite under slab geotherm conditions (Ohira et al. 2014). In association with Phase D, Phase H contains little Al, which strongly partitions into Phase D. At pressure above \sim 50 GPa, where Phase D is no longer stable, Phase H has the ability to incorporate a significant amount of Al, present in hydrous compositions in the deep lower mantle (Bindi et al. [2015\)](#page-35-0). In this case, the effect of the Al component in Phase H is expected to extend its stability field, which is similar to that in Phase D. Phase H may be considered as an end-member of hydrous-phase solid solution $AIOOH-MgSiO₄H₂$. Experiments by Ohira et al. [\(2014](#page-39-0)) in a high-pressure DAC demonstrated that at 68 GPa and 2010 K, hydrous δ -phase contains at least 44 mol% of MgSiO₄H₂ component; at 128 GPa and 2190 K, it contains 23 mol% of Phase H component. In other series of the DAC experiments in system $MgO-AI_2O_3-SiO_2-H_2O$, Phase H was stable between \sim 45 and 115 GPa in all bulk compositions studied (Walter et al. 2015). Most recent experiments in the MgSiO₄H₂-AlOOH–SiO₂ system within the 20-60 GPa pressure range demonstrated that phase H and δ -AlOOH form an ideal mixture and that hydrogen in phase H is bonded in a single well at all pressures, and only the pure phase results in pressure-induced symmetrization, while the solid solution maintains asymmetric bonding to high pressure (Panero and Caracas [2017\)](#page-40-0). Thus, Phase H and δ -AlOOH are likely to be stable hydrous phases in the slabs and transport significant amounts of water into the deep lower mantle.

The presence of hydrogen in the lower mantle may influence other phase transitions. It sharpens the width of the bridgmanite to post-perovskite transition. The unit cell volume of bridgmanite at ambient conditions increases systematically with increasing synthesis pressure above \sim 55 GPa, possibly representing an increase in Al concentration and potentially hydrogen concentration with depth (Inoue et al. [2010\)](#page-37-0).

Bridgmanite in equilibrium with phases D and H has a relatively low Al concentration, and Al partitions preferentially into the hydrous phases (Walter et al. [2015](#page-41-0)).

The stability of all DHMS phases that are expected to be major inventories of water in subducting slabs are summarized in Fig. 5.17. Phase H in this composition is formed by the reaction of Phase D and periclase, which is assumed to occur at the same pressure as the reaction of Phase D and brucite, although this boundary could be shifted towards higher pressures, defined by the reaction from Phase D to Phase H plus stishovite (Nishi et al. [2014\)](#page-39-0).

High-pressure aluminium-rich phases D and H are more likely to exist in the mafic association because the ultramafic association does not contain enough Al to make these phases stable in the lower-mantle conditions. Phases D and H can survive only in cold subducting slabs because, along an adiabatic geotherm, solid hydrous silicates are not stable. However, to date, none of the DHMS phases were identified in the natural environment.

5.7 Conclusions

The mafic association plays a less important role in the lower mantle compared with the ultramafic one. Firstly, it is much rarer because of the relatively small volume of subducted lithospheric slabs within the lower mantle. Secondly, the structural changes are not so noticeable compared with that in the ultramafic association. For example, a possible NAL to CF phase transition, accompanied by an increase in density, bulk modulus and share wave anisotropy, could be seismically detectable.

Depth, km

However, it is compensated for, in the bulk slab composition, by the faster sound velocity of $SiO₂$ and its high-pressure polymorphic phase(s). The transitions may be observed only locally within subducting slabs. This means that the mafic association does not play a major role in the general structure of the lower mantle. On the other hand, subducting slabs may cause local seismic anomalies and deliver elements such as Si, Al and water to the deep mantle, forming local water reservoirs within the lower mantle. This may result in local seismic and chemical heterogeneities that can influence deep-mantle geodynamic processes.

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