

Summation Formulas of Euler–Maclaurin and Abel–Plana: Old and New Results and Applications

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Abstract Summation formulas of the Euler–Maclaurin and Abel–Plana and their connections with several kinds of quadrature rules are studied. Besides the history of these formulas, several of their modifications and generalizations are considered. Connections between the Euler–Maclaurin formula and basic quadrature rules of Newton–Cotes type, as well as the composite Gauss–Legendre rule and its Lobatto modification are presented. Besides the basic Plana summation formula a few integral modifications (the midpoint summation formula, the Binet formula, Lindelöf formula) are introduced and analysed. Starting from the moments of their weight functions and applying the recent MATHEMATICA package `OrthogonalPolynomials`, recursive coefficients in the three-term recurrence relation for the corresponding orthogonal polynomials are constructed, as well as the parameters (nodes and Christoffel numbers) of the corresponding Gaussian quadrature formula.

Keywords Summation • Euler–Maclaurin formula • Abel–Plana formula • Gaussian quadrature formula • Orthogonal polynomial • Three-term recurrence relation

Mathematics Subject Classification (2010): 33C45, 33C47, 41A55, 65B15, 65D30, 65D32

1 Introduction and Preliminaries

A summation formula was discovered independently by Leonhard Euler [18, 19] and Colin Maclaurin [35] plays an important role in the broad area of numerical analysis, analytic number theory, approximation theory, as well as in many applications in

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N.K. Govil et al. (eds.), *Progress in Approximation Theory and Applicable Complex Analysis*, Springer Optimization and Its Applications 117, DOI 10.1007/978-3-319-49242-1_20

429

other fields. This formula, today known as the Euler–Maclaurin summation formula,

$$\sum_{k=0}^n f(k) = \int_0^n f(x) dx + \frac{1}{2}(f(0) + f(n)) + \sum_{\nu=1}^r \frac{B_{2\nu}}{(2\nu)!} [f^{(2\nu-1)}(n) - f^{(2\nu-1)}(0)] + E_r(f), \quad (1)$$

was published first time by Euler in 1732 (without proof) in connection with the problem of determining the sum of the reciprocal squares,

$$1 + \frac{1}{2^2} + \frac{1}{3^2} + \dots, \quad (2)$$

which is known as the *Basel problem*. The brothers Johann and Jakob Bernoulli, Leibnitz, Stirling, etc. also dealt intensively by such a kind of problems. In modern terminology, the sum (2) represents the zeta function of 2, where more generally

$$\zeta(s) = 1 + \frac{1}{2^s} + \frac{1}{3^s} + \dots \quad (s > 1).$$

Although at that time the theory of infinite series was not exactly based, it was observed a very slow convergence of this series, e.g. in order to compute directly the sum with an accuracy of six decimal places it requires taking into account at least a million first terms, because

$$\frac{1}{n+1} < \sum_{k=n+1}^{+\infty} \frac{1}{k^2} < \frac{1}{n}.$$

Euler discovered the remarkable formula with much faster convergence

$$\zeta(2) = \log^2 2 + \sum_{k=1}^{+\infty} \frac{1}{2^{k-1}k^2},$$

and obtained the value $\zeta(2) = 1.644944\dots$ (with seven decimal digits). But the discovery of a general summation procedure (1) enabled Euler to calculate $\zeta(2)$ to 20 decimal places. For details see Gautschi [25, 26] and Varadarajan [61].

Using a generalized Newton identity for polynomials (when their degree tends to infinity), Euler [19] proved the exact result $\zeta(2) = \pi^2/6$. Using the same method he determined $\zeta(s)$ for even $s = 2m$ up to 12,

$$\zeta(4) = \frac{\pi^4}{90}, \quad \zeta(6) = \frac{\pi^6}{945}, \quad \zeta(8) = \frac{\pi^8}{9450}, \quad \zeta(10) = \frac{\pi^{10}}{93555}, \quad \zeta(12) = \frac{691\pi^{12}}{638512875}.$$

Sometime later, using his own partial fraction expansion of the cotangent function, Euler obtained the general formula

$$\zeta(2\nu) = (-1)^{\nu-1} \frac{2^{2\nu-1} B_{2\nu}}{(2\nu)!} \pi^{2\nu},$$

where $B_{2\nu}$ are the Bernoulli numbers, which appear in the general Euler–Maclaurin summation formula (1). Detailed information about Euler’s complete works can be found in *The Euler Archive* (<http://eulerarchive.maa.org>).

We return now to the general Euler–Maclaurin summation formula (1) which holds for any $n, r \in \mathbb{N}$ and $f \in C^{2r}[0, n]$. As we mentioned before this formula was found independently by Maclaurin. While in Euler’s case the formula (1) was applied for computing slowly converging infinite series, in the second one Maclaurin used it to calculate integrals. A history of this formula was given by Barnes [5], and some details can be found in [3, 8, 25, 26, 38, 61].

Bernoulli numbers B_k ($B_0 = 1, B_1 = -1/2, B_2 = 1/6, B_3 = 0, B_4 = -1/30, \dots$) can be expressed as values at zero of the corresponding Bernoulli polynomials, which are defined by the generating function

$$\frac{te^{xt}}{e^t - 1} = \sum_{k=0}^{+\infty} B_k(x) \frac{t^k}{k!}.$$

Similarly, Euler polynomials can be introduced by

$$\frac{2e^{xt}}{e^t + 1} = \sum_{k=0}^{+\infty} E_k(x) \frac{t^k}{k!}.$$

Bernoulli and Euler polynomials play a similar role in numerical analysis and approximation theory like orthogonal polynomials. First few Bernoulli polynomials are

$$B_0(x) = 1, \quad B_1(x) = x - \frac{1}{2}, \quad B_2(x) = x^2 - x + \frac{1}{6}, \quad B_3(x) = x^3 - \frac{3x^2}{2} + \frac{x}{2},$$

$$B_4(x) = x^4 - 2x^3 + x^2 - \frac{1}{30}, \quad B_5(x) = x^5 - \frac{5x^4}{2} + \frac{5x^3}{3} - \frac{x}{6}, \quad \text{etc.}$$

Some interesting properties of these polynomials are

$$B'_n(x) = nB_{n-1}(x), \quad B_n(1-x) = (-1)^n B_n(x), \quad \int_0^1 B_n(x) \, dx = 0 \quad (n \in \mathbb{N}).$$

The error term $E_r(f)$ in (1) can be expressed in the form (cf. [8])

$$E_r(f) = (-1)^r \sum_{k=1}^{+\infty} \int_0^n \frac{e^{i2\pi kt} + e^{-i2\pi kt}}{(2\pi k)^{2r}} f^{(2r)}(x) \, dx,$$

or in the form

$$E_r(f) = - \int_0^n \frac{B_{2r}(x - [x])}{(2r)!} f^{(2r)}(x) dx, \tag{3}$$

where $[x]$ denotes the largest integer that is not greater than x . Supposing $f \in C^{2r+1}[0, n]$, after an integration by parts in (3) and recalling that the odd Bernoulli numbers are zero, we get (cf. [28, p. 455])

$$E_r(f) = \int_0^n \frac{B_{2r+1}(x - [x])}{(2r + 1)!} f^{(2r+1)}(x) dx. \tag{4}$$

If $f \in C^{2r+2}[0, n]$, using Darboux's formula one can obtain (1), with

$$E_r(f) = \frac{1}{(2r + 2)!} \int_0^1 [B_{2r+2} - B_{2r+2}(x)] \left(\sum_{k=0}^{n-1} f^{(2r+2)}(k + x) \right) dx \tag{5}$$

(cf. Whittaker and Watson [65, p. 128]). This expression for $E_r(f)$ can be also derived from (4), writing it in the form

$$\begin{aligned} E_r(f) &= \int_0^1 \frac{B_{2r+1}(x)}{(2r + 1)!} \left(\sum_{k=0}^{n-1} f^{(2r+1)}(k + x) \right) dx \\ &= \int_0^1 \frac{B'_{2r+2}(x)}{(2r + 2)!} \left(\sum_{k=0}^{n-1} f^{(2r+1)}(k + x) \right) dx, \end{aligned}$$

and then by an integration by parts,

$$\begin{aligned} E_r(f) &= \left[\frac{B_{2r+2}(x)}{(2r + 2)!} \left(\sum_{k=0}^{n-1} f^{(2r+1)}(k + x) \right) \right]_0^1 \\ &\quad - \int_0^1 \frac{B_{2r+2}(x)}{(2r + 2)!} \left(\sum_{k=0}^{n-1} f^{(2r+2)}(k + x) \right) dx. \end{aligned}$$

Because of $B_{2r+2}(1) = B_{2r+2}(0) = B_{2r+2}$, the last expression can be represented in the form (5).

Since

$$(-1)^r [B_{2r+2} - B_{2r+2}(x)] \geq 0, \quad x \in [0, 1],$$

and

$$\int_0^1 [B_{2r+2} - B_{2r+2}(x)] dt = B_{2r+2},$$

according to the *Second Mean Value Theorem for Integrals*, there exists $\eta \in (0, 1)$ such that

$$E_r(f) = \frac{B_{2r+2}}{(2r+2)!} \left(\sum_{k=0}^{n-1} f^{(2r+2)}(k+\eta) \right) = \frac{nB_{2r+2}}{(2r+2)!} f^{(2r+2)}(\xi), \quad 0 < \xi < n. \tag{6}$$

The Euler–Maclaurin summation formula can be considered on an arbitrary interval (a, b) instead of $(0, n)$. Namely, taking $h = (b - a)/n$, $t = a + xh$, and $f(x) = f((t - a)/h) = \varphi(t)$, formula (1) reduces to

$$h \sum_{k=0}^n \varphi(a + kh) = \int_a^b \varphi(t) dt + \frac{h}{2} [\varphi(a) + \varphi(b)] + \sum_{v=1}^r \frac{B_{2v} h^{2v}}{(2v)!} [\varphi^{(2v-1)}(b) - \varphi^{(2v-1)}(a)] + E_r(\varphi), \tag{7}$$

where, according to (6),

$$E_r(\varphi) = (b - a) \frac{B_{2r+2} h^{2r+2}}{(2r+2)!} \varphi^{(2r+2)}(\xi), \quad a < \xi < b. \tag{8}$$

Remark 1. An approach in the estimate of the remainder term of the Euler–Maclaurin formula was given by Ostrowski [47].

Remark 2. The Euler–Maclaurin summation formula is implemented in MATHEMATICA as the function NSum with option Method -> Integrate.

2 Connections Between Euler–Maclaurin Summation Formula and Some Basic Quadrature Rules of Newton–Cotes Type

In this section we first show a direct connection between the Euler–Maclaurin summation formula (1) and the well-known composite trapezoidal rule,

$$T_n f := \sum_{k=0}^n {}'f(k) = \frac{1}{2} f(0) + \sum_{k=1}^{n-1} f(k) + \frac{1}{2} f(n), \tag{9}$$

for calculating the integral

$$I_n f := \int_0^n f(x) dx. \tag{10}$$

This rule for integrals over an arbitrary interval $[a, b]$ can be presented in the form

$$h \sum_{k=0}^n \prime \varphi(a + kh) = \int_a^b \varphi(t) dt + E^T(\varphi), \tag{11}$$

where, as before, the sign $\sum \prime$ denotes summation with the first and last terms halved, $h = (b - a)/n$, and $E^T(\varphi)$ is the remainder term.

Remark 3. In general, the sequence of the composite trapezoidal sums converges very slowly with respect to step refinement, because of $|E^T(\varphi)| = O(h^2)$. However, the trapezoidal rule is very attractive in numerical integration of analytic and periodic functions, for which $\varphi(t + b - a) = \varphi(t)$. In that case, the sequence of trapezoidal sums

$$T_n(\varphi; h) := h \sum_{k=0}^n \prime \varphi(a + kh) = h \sum_{k=1}^n \varphi(a + kh) \tag{12}$$

converges geometrically when applied to analytic functions on periodic intervals or the real line. A nice survey on this subject, including history of this phenomenon, has been recently given by Trefethen and Weideman [59] (see also [64]). For example, when φ is a $(b - a)$ -periodic and analytic function, such that $|\varphi(z)| \leq M$ in the half-plane $\text{Im } z > -c$ for some $c > 0$, then for each $n \geq 1$, the following estimate

$$|E^T(\varphi)| = \left| T_n(\varphi; h) - \int_a^b \varphi(t) dt \right| \leq \frac{(b - a)M}{e^{2\pi cn/(b-a)} - 1}$$

holds. A similar result holds for integrals over \mathbb{R} .

It is well known that there are certain types of integrals which can be transformed (by changing the variable of integration) to a form suitable for the trapezoidal rule. Such transformations are known as *Exponential* and *Double Exponential Quadrature Rules* (cf. [44–46, 57, 58]). However, the use of these transformations could introduce new singularities in the integrand and the analyticity strip may be lost. A nice discussion concerning the error theory of the trapezoidal rule, including several examples, has been recently given by Waldvogel [63].

Remark 4. In 1990 Rahman and Schmeisser [51] gave a specification of spaces of functions for which the trapezoidal rule converges at a prescribed rate as $n \rightarrow +\infty$, where a correspondence is established between the speed of convergence and regularity properties of integrands. Some examples for these spaces were provided in [64].

In a general case, according to (1), it is clear that

$$T_n f - I_n f = \sum_{v=1}^r \frac{B_{2v}}{(2v)!} [f^{(2v-1)}(n) - f^{(2v-1)}(0)] + E_r^T(f), \tag{13}$$

where $T_n f$ and $I_n f$ are given by (9) and (10), respectively, and the remainder term $E_r^T(f)$ is given by (6) for functions $f \in C^{2r+2}[0, n]$.

Similarly, because of (7), the corresponding formula on the interval $[a, b]$ is

$$h \sum_{k=0}^n \varphi(a + kh) - \int_a^b \varphi(t) dt = \sum_{\nu=1}^r \frac{B_{2\nu} h^{2\nu}}{(2\nu)!} [\varphi^{(2\nu-1)}(b) - \varphi^{(2\nu-1)}(a)] + E_r^T(\varphi),$$

where $E_r^T(\varphi)$ is the corresponding remainder given by (8). Comparing this with (11) we see that $E^T(\varphi) = E_0^T(\varphi)$.

Notice that if $\varphi^{(2r+2)}(x)$ does not change its sign on (a, b) , then $E_r^T(\varphi)$ has the same sign as the first neglected term. Otherwise, when $\varphi^{(2r+2)}(x)$ is not of constant sign on (a, b) , then it can be proved that (cf. [14, p. 299])

$$|E_r^T(\varphi)| \leq h^{2r+2} \frac{|2B_{2r+2}|}{(2r+2)!} \int_a^b |\varphi^{(2r+2)}(t)| dt \approx 2 \left(\frac{h}{2\pi}\right)^{2r+2} \int_a^b |\varphi^{(2r+2)}(t)| dt,$$

i.e., $|E_r^T(\varphi)| = O(h^{2r+2})$. Supposing that $\int_a^{+\infty} |\varphi^{(2r+2)}(x)| dx < +\infty$, this holds also in the limit case as $b \rightarrow +\infty$. This limit case enables applications of the Euler–Maclaurin formula in summation of infinite series, as well as for obtaining asymptotic formulas for a large b .

A standard application of the Euler–Maclaurin formula is in numerical integration. Namely, for a small constant h , the trapezoidal sum can be dramatically improved by subtracting appropriate terms with the values of derivatives at the endpoints a and b . In such a way, the corresponding approximations of the integral can be improved to $O(h^4)$, $O(h^6)$, etc.

Remark 5. Rahman and Schmeisser [52] obtained generalizations of the trapezoidal rule and the Euler–Maclaurin formula and used them for constructing quadrature formulas for functions of exponential type over infinite intervals using holomorphic functions of exponential type in the right half-plane, or in a vertical strip, or in the whole plane. They also determined conditions which equate the existence of the improper integral to the convergence of its approximating series.

Remark 6. In this connection an interesting question can be asked. Namely, what happens if the function $\varphi \in C^\infty(\mathbb{R})$ and its derivatives are $(b - a)$ -periodic, i.e., $\varphi^{(2\nu-1)}(a) = \varphi^{(2\nu-1)}(b)$, $\nu = 1, 2, \dots$? The conclusion that $T_n(\varphi; h)$, in this case, must be exactly equal to $\int_a^b \varphi(t) dt$ is wrong, but the correct conclusion is that $E^T(\varphi)$ decreases faster than any finite power of h as n tends to infinity.

Remark 7. Also, the Euler–Maclaurin formula was used in getting an extrapolating method well-known as Romberg’s integration (cf. [14, pp. 302–308 and 546–523] and [39, pp. 158–164]).

In the sequel, we consider a quadrature sum with values of the function f at the points $x = k + \frac{1}{2}$, $k = 0, 1, \dots, n - 1$, i.e., the so-called midpoint rule

$$M_n f := \sum_{k=0}^{n-1} f\left(k + \frac{1}{2}\right).$$

Also, for this rule there exists the so-called second Euler–Maclaurin summation formula

$$M_n f - I_n f = \sum_{\nu=1}^r \frac{(2^{1-2\nu} - 1)B_{2\nu}}{(2\nu)!} [f^{(2\nu-1)}(n) - f^{(2\nu-1)}(0)] + E_r^M(f), \tag{14}$$

for which

$$E_r^M(f) = n \frac{(2^{-1-2r} - 1)B_{2r+2}}{(2r + 2)!} f^{(2r+2)}(\xi), \quad 0 < \xi < n,$$

when $f \in C^{2r+2}[0, n]$ (cf. [39, p. 157]). As before, $I_n f$ is given by (10).

The both formulas, (13) and (14), can be unified as

$$Q_n f - I_n f = \sum_{\nu=1}^r \frac{B_{2\nu}(\tau)}{(2\nu)!} [f^{(2\nu-1)}(n) - f^{(2\nu-1)}(0)] + E_r^Q(f),$$

where $\tau = 0$ for $Q_n \equiv T_n$ and $\tau = 1/2$ for $Q_n \equiv M_n$. It is true, because of the fact that [50, p. 765] (see also [10])

$$B_\nu(0) = B_\nu \quad \text{and} \quad B_\nu\left(\frac{1}{2}\right) = (2^{1-\nu} - 1)B_\nu.$$

If we take a combination of $T_n f$ and $M_n f$ as

$$Q_n f = S_n f = \frac{1}{3}(T_n f + 2M_n f),$$

which is, in fact, the well-known classical composite Simpson rule,

$$S_n f := \frac{1}{3} \left[\frac{1}{2} f(0) + \sum_{k=1}^{n-1} f(k) + 2 \sum_{k=0}^{n-1} f\left(k + \frac{1}{2}\right) + \frac{1}{2} f(n) \right],$$

we obtain

$$S_n f - I_n f = \sum_{\nu=2}^r \frac{(4^{1-\nu} - 1)B_{2\nu}}{3(2\nu)!} [f^{(2\nu-1)}(n) - f^{(2\nu-1)}(0)] + E_r^S(f). \tag{15}$$

Notice that the summation on the right-hand side in the previous equality starts with $\nu = 2$, because the term for $\nu = 1$ vanishes. For $f \in C^{2r+2}[0, n]$ it can be proved that there exists $\xi \in (0, n)$, such that

$$E_r^S(f) = n \frac{(4^{-r} - 1)B_{2r+2}}{3(2r + 2)!} f^{(2r+2)}(\xi).$$

For some modification and generalizations of the Euler–Maclaurin formula, see [2, 7, 20–22, 37, 55, 60]. In 1965 Kalinin [29] derived an analogue of the Euler–Maclaurin formula for C^∞ functions, for which there is Taylor series at each positive integer $x = \nu$,

$$\int_a^b f(x) \, dx = \sum_{k=0}^{+\infty} \frac{\theta^{k+1} - (\theta - 1)^{k+1}}{(k + 1)!} h^{k+1} \sum_{\nu=1}^n f^{(k)}(a + (\nu - \theta)h),$$

where $h = (b - a)/n$, and used it to find some new expansions for the gamma function, the ψ function, as well as the Riemann zeta function.

Using Bernoulli and Euler polynomials, $B_n(x)$ and $E_n(x)$, in 1960 Keda [30] established a quadrature formula similar to the Euler–Maclaurin,

$$\int_0^1 f(x) \, dx = T_n + \sum_{k=0}^{n-1} A_k [f^{(2k+2)}(0) + f^{(2k+2)}(1)] + R_n,$$

where

$$T_n = \frac{1}{n} \sum_{k=0}^n f\left(\frac{k}{n}\right), \quad A_k = \sum_{\nu=1}^{2k+2} \frac{B_\nu E_{2k+3-\nu}}{\nu!(2k + 3 - \nu)!n^\nu} \quad (k = 0, 1, \dots, n - 1),$$

and

$$R_n = f^{(2n+2)}(\xi) \sum_{m=1}^{n+1} \frac{2B_{2m}E_{2n-2m+3}}{(2m)!(2n - 2m + 3)!n^{2m}} \quad (0 \leq \xi \leq 1)$$

for $f \in C^{2n+2}[0, 1]$, where $B_n = B_n(0)$ and $E_n = E_n(0)$. The convergence of Euler–Maclaurin quadrature formulas on a class of smooth functions was considered by Vaskevič [62].

Some periodic analogues of the Euler–Maclaurin formula with applications to number theory have been developed by Berndt and Schoenfeld [6]. In the last section of [6], they showed how the composite Newton–Cotes quadrature formulas (Simpson’s parabolic and Simpson’s three-eighths rules), as well as various other quadratures (e.g., Weddle’s composite rule), can be derived from special cases of their periodic Euler–Maclaurin formula, including explicit formulas for the remainder term.

3 Euler–Maclaurin Formula Based on the Composite Gauss–Legendre Rule and Its Lobatto Modification

In the papers [15, 48, 56], the authors considered generalizations of the Euler–Maclaurin formula for some particular Newton–Cotes rules, as well as for 2- and 3-point Gauss–Legendre and Lobatto formulas (see also [4, 17, 33, 34]).

Recently, we have done [40] the extensions of Euler–Maclaurin formulas by replacing the quadrature sum Q_n by the composite Gauss–Legendre shifted formula, as well as by its Lobatto modification. In these cases, several special rules have been obtained by using the MATHEMATICA package `OrthogonalPolynomials` (cf. [9, 43]). Some details on construction of orthogonal polynomials and quadratures of Gaussian type will be given in Sect. 5.

We denote the space of all algebraic polynomials defined on \mathbb{R} (or some its subset) by \mathcal{P} , and by $\mathcal{P}_m \subset \mathcal{P}$ the space of polynomials of degree at most m ($m \in \mathbb{N}$).

Let $w_\nu = w_\nu^G$ and $\tau_\nu = \tau_\nu^G$, $\nu = 1, \dots, m$, be weights (Christoffel numbers) and nodes of the Gauss–Legendre quadrature formula on $[0, 1]$,

$$\int_0^1 f(x) dx = \sum_{\nu=1}^m w_\nu^G f(\tau_\nu^G) + R_m^G(f). \tag{16}$$

Note that the nodes τ_ν are zeros of the shifted (monic) Legendre polynomial

$$\pi_m(x) = \binom{2m}{m}^{-1} P_m(2x - 1).$$

Degree of its algebraic precision is $d = 2m - 1$, i.e., $R_m^G(f) = 0$ for each $f \in \mathcal{P}_{2m-1}$. The quadrature sum in (16) we denote by $Q_m^G f$, i.e.,

$$Q_m^G f = \sum_{\nu=1}^m w_\nu^G f(\tau_\nu^G).$$

The corresponding composite Gauss–Legendre sum for approximating the integral $I_n f$, given by (10), can be expressed in the form

$$G_m^{(n)} f = \sum_{k=0}^{n-1} Q_m^G f(k + \cdot) = \sum_{\nu=1}^m w_\nu^G \sum_{k=0}^{n-1} f(k + \tau_\nu^G). \tag{17}$$

In the sequel we use the following expansion of a function $f \in C^s[0, 1]$ in Bernoulli polynomials for any $x \in [0, 1]$ (see Krylov [31, p. 15])

$$f(x) = \int_0^1 f(t) dt + \sum_{j=1}^{s-1} \frac{B_j(x)}{j!} [f^{(j-1)}(1) - f^{(j-1)}(0)] - \frac{1}{s!} \int_0^1 f^{(s)}(t) L_s(x, t) dt, \tag{18}$$

where $L_s(x, t) = B_s^*(x-t) - B_s^*(x)$ and $B_s^*(x)$ is a function of period one, defined by

$$B_s^*(x) = B_s(x), \quad 0 \leq x < 1, \quad B_s^*(x+1) = B_s^*(x). \tag{19}$$

Notice that $B_0^*(x) = 1$, $B_1^*(x)$ is a discontinuous function with a jump of -1 at each integer, and $B_s^*(x)$, $s > 1$, is a continuous function.

Suppose now that $f \in C^{2r}[0, n]$, where $r \geq m$. Since the all nodes $\tau_\nu = \tau_\nu^G$, $\nu = 1, \dots, m$, of the Gaussian rule (16) belong to $(0, 1)$, using the expansion (18), with $x = \tau_\nu$ and $s = 2r + 1$, we have

$$\begin{aligned} f(\tau_\nu) = I_1 f + \sum_{j=1}^{2r} \frac{B_j(\tau_\nu)}{j!} [f^{(j-1)}(1) - f^{(j-1)}(0)] \\ - \frac{1}{(2r+1)!} \int_0^1 f^{(2r+1)}(t) L_{2r+1}(\tau_\nu, t) dt, \end{aligned}$$

where $I_1 f = \int_0^1 f(t) dt$.

Now, if we multiply it by $w_\nu = w_\nu^G$ and then sum in ν from 1 to m , we obtain

$$\begin{aligned} \sum_{\nu=1}^m w_\nu f(\tau_\nu) = \left(\sum_{\nu=1}^m w_\nu \right) I_1 f + \sum_{j=1}^{2r} \frac{1}{j!} \left(\sum_{\nu=1}^m w_\nu B_j(\tau_\nu) \right) [f^{(j-1)}(1) - f^{(j-1)}(0)] \\ - \frac{1}{(2r+1)!} \int_0^1 f^{(2r+1)}(t) \left(\sum_{\nu=1}^m w_\nu L_{2r+1}(\tau_\nu, t) \right) dt, \end{aligned}$$

i.e.,

$$Q_m^G f = Q_m^G(1) \int_0^1 f(t) dt + \sum_{j=1}^{2r} \frac{Q_m^G(B_j)}{j!} [f^{(j-1)}(1) - f^{(j-1)}(0)] + E_{m,r}^G(f),$$

where

$$E_{m,r}^G(f) = -\frac{1}{(2r+1)!} \int_0^1 f^{(2r+1)}(t) Q_m^G(L_{2r+1}(\cdot, t)) dt.$$

Since

$$\int_0^1 B_j(x) dx = \begin{cases} 1, & j = 0, \\ 0, & j \geq 1, \end{cases}$$

and

$$Q_m^G(B_j) = \sum_{\nu=1}^m w_\nu B_j(\tau_\nu) = \begin{cases} 1, & j = 0, \\ 0, & 1 \leq j \leq 2m - 1, \end{cases}$$

because the Gauss–Legendre formula is exact for all algebraic polynomials of degree at most $2m - 1$, the previous formula becomes

$$Q_m^G f - \int_0^1 f(t) dt = \sum_{j=2m}^{2r} \frac{Q_m^G(B_j)}{j!} [f^{(j-1)}(1) - f^{(j-1)}(0)] + E_{m,r}^G(f). \quad (20)$$

Notice that for Gauss–Legendre nodes and the corresponding weights the following equalities

$$\tau_\nu + \tau_{m-\nu+1} = 1, \quad w_\nu = w_{m-\nu+1} > 0, \quad \nu = 1, \dots, m,$$

hold, as well as

$$w_\nu B_j(\tau_\nu) + w_{m-\nu+1} B_j(\tau_{m-\nu+1}) = w_\nu B_j(\tau_\nu)(1 + (-1)^j),$$

which is equal to zero for odd j . Also, if m is odd, then $\tau_{(m+1)/2} = 1/2$ and $B_j(1/2) = 0$ for each odd j . Thus, the quadrature sum

$$Q_m^G(B_j) = \sum_{\nu=1}^m w_\nu B_j(\tau_\nu) = 0$$

for odd j , so that (20) becomes

$$Q_m^G f - \int_0^1 f(t) dt = \sum_{j=m}^r \frac{Q_m^G(B_{2j})}{(2j)!} [f^{(2j-1)}(1) - f^{(2j-1)}(0)] + E_{m,r}^G(f). \quad (21)$$

Consider now the error of the (shifted) composite Gauss–Legendre formula (17). It is easy to see that

$$\begin{aligned} G_m^{(n)} f - I_n f &= \sum_{k=0}^{n-1} \left[Q_m^G f(k + \cdot) - \int_k^{k+1} f(t) dt \right] \\ &= \sum_{k=0}^{n-1} \left[Q_m^G f(k + \cdot) - \int_0^1 f(k + x) dx \right]. \end{aligned}$$

Then, using (21) we obtain

$$\begin{aligned} G_m^{(n)} f - I_n f &= \sum_{k=0}^{n-1} \left\{ \sum_{j=m}^r \frac{Q_m^G(B_{2j})}{(2j)!} [f^{(2j-1)}(k+1) - f^{(2j-1)}(k)] + E_{m,r}^G(f(k + \cdot)) \right\} \\ &= \sum_{j=m}^r \frac{Q_m^G(B_{2j})}{(2j)!} [f^{(2j-1)}(n) - f^{(2j-1)}(0)] + E_{n,m,r}^G(f), \end{aligned}$$

where $E_{n,m,r}^G(f)$ is given by

$$E_{n,m,r}^G(f) = -\frac{1}{(2r+1)!} \int_0^1 \left(\sum_{k=0}^{n-1} f^{(2r+1)}(k+t) \right) Q_m^G(L_{2r+1}(\cdot, t)) dt. \tag{22}$$

Since $L_{2r+1}(x, t) = B_{2r+1}^*(x-t) - B_{2r+1}^*(x)$ and

$$B_{2r+1}^*(\tau_v) = B_{2r+1}(\tau_v), \quad B_{2r+1}^*(\tau_v - t) = -\frac{1}{2r+2} \frac{d}{dt} B_{2r+2}^*(\tau_v - t),$$

we have

$$\begin{aligned} Q_m^G(L_{2r+1}(\cdot, t)) &= Q_m^G(B_{2r+1}^*(\cdot - t)) - Q_m^G(B_{2r+1}^*(\cdot)) \\ &= -\frac{1}{2r+2} Q_m^G\left(\frac{d}{dt} B_{2r+2}^*(\cdot - t)\right), \end{aligned}$$

because $Q_m^G(B_{2r+1}(\cdot)) = 0$. Then for (22) we get

$$(2r+2)!E_{n,m,r}^G(f) = \int_0^1 \left(\sum_{k=0}^{n-1} f^{(2r+1)}(k+t) \right) Q_m^G\left(\frac{d}{dt} B_{2r+2}^*(\cdot - t)\right) dt.$$

By using an integration by parts, it reduces to

$$(2r+2)!E_{n,m,r}^G(f) = F(t)Q_m^G(B_{2r+2}^*(\cdot - t)) \Big|_0^1 - \int_0^1 Q_m^G(B_{2r+2}^*(\cdot - t)) F'(t) dt,$$

where $F(t)$ is introduced in the following way

$$F(t) = \sum_{k=0}^{n-1} f^{(2r+1)}(k+t).$$

Since $B_{2r+2}^*(\tau_v - 1) = B_{2r+2}^*(\tau_v) = B_{2r+2}(\tau_v)$, we have

$$\begin{aligned} F(t)Q_m^G(B_{2r+2}^*(\cdot - t)) \Big|_0^1 &= (F(1) - F(0))Q_m^G(B_{2r+2}^*(\cdot)) \\ &= Q_m^G(B_{2r+2}(\cdot)) \int_0^1 F'(t) dt, \end{aligned}$$

so that

$$(2r+2)!E_{n,m,r}^G(f) = \int_0^1 [Q_m^G(B_{2r+2}(\cdot)) - Q_m^G(B_{2r+2}^*(\cdot - t))] F'(t) dt.$$

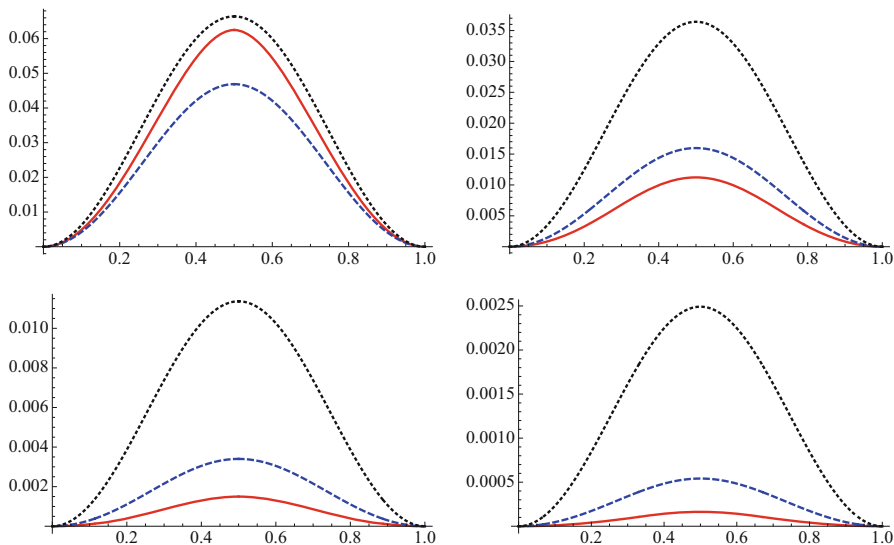


Fig. 1 Graphs of $t \mapsto g_{m,r}^G(t)$, $r = m$ (solid line), $r = m + 1$ (dashed line), and $r = m + 2$ (dotted line), when $m = 1, m = 2$ (top), and $m = 3, m = 4$ (bottom)

Since

$$g_{m,r}^G(t) := (-1)^{r-m} Q_m^G [B_{2r+2}(\cdot) - B_{2r+2}^*(\cdot - t)] > 0, \quad 0 < t < 1, \quad (23)$$

there exists an $\eta \in (0, 1)$ such that

$$(2r + 2)! E_{n,m,r}^G(f) = F'(\eta) \int_0^1 Q_m^G [B_{2r+2}(\cdot) - B_{2r+2}^*(\cdot - t)] dt.$$

Typical graphs of functions $t \mapsto g_{m,r}^G(t)$ for some selected values of $r \geq m \geq 1$ are presented in Fig. 1.

Because of continuity of $f^{(2r+2)}$ on $[0, n]$ we conclude that there exists also $\xi \in (0, n)$ such that $F'(\eta) = n f^{(2r+2)}(\xi)$.

Finally, because of $\int_0^1 Q_m^G [B_{2r+2}^*(\cdot - t)] dt = 0$, we obtain that

$$(2r + 2)! E_{n,m,r}^G(f) = n f^{(2r+2)}(\xi) \int_0^1 Q_m^G [B_{2r+2}(\cdot)] dt.$$

In this way, we have just proved the Euler–Maclaurin formula for the composite Gauss–Legendre rule (17) for approximating the integral $I_n f$, given by (10) (see [40]):

Theorem 1. For $n, m, r \in \mathbb{N}$ ($m \leq r$) and $f \in C^{2r}[0, n]$ we have

$$G_m^{(n)}f - I_n f = \sum_{j=m}^r \frac{Q_m^G(B_{2j})}{(2j)!} [f^{(2j-1)}(n) - f^{(2j-1)}(0)] + E_{n,m,r}^G(f), \tag{24}$$

where $G_m^{(n)}f$ is given by (17), and $Q_m^G B_{2j}$ denotes the basic Gauss–Legendre quadrature sum applied to the Bernoulli polynomial $x \mapsto B_{2j}(x)$, i.e.,

$$Q_m^G(B_{2j}) = \sum_{v=1}^m w_v^G B_{2j}(\tau_v^G) = -R_m^G(B_{2j}), \tag{25}$$

where $R_m^G(f)$ is the remainder term in (16).

If $f \in C^{2r+2}[0, n]$, then there exists $\xi \in (0, n)$, such that the error term in (24) can be expressed in the form

$$E_{n,m,r}^G(f) = n \frac{Q_m^G(B_{2r+2})}{(2r+2)!} f^{(2r+2)}(\xi). \tag{26}$$

We consider now special cases of the formula (24) for some typical values of m . For a given m , by $G^{(m)}$ we denote the sequence of coefficients which appear in the sum on the right-hand side in (24), i.e.,

$$G^{(m)} = \{Q_m^G(B_{2j})\}_{j=m}^\infty = \{Q_m^G(B_{2m}), Q_m^G(B_{2m+2}), Q_m^G(B_{2m+4}), \dots\}.$$

These Gaussian sums we can calculate very easily by using MATHEMATICA Package `OrthogonalPolynomials` (cf. [9, 43]). In the sequel we mention cases when $1 \leq m \leq 6$.

Case $m = 1$. Here $\tau_1^G = 1/2$ and $w_1^G = 1$, so that, according to (25),

$$Q_1^G(B_{2j}) = B_{2j}(1/2) = (2^{1-2j} - 1)B_{2j},$$

and (24) reduces to (14). Thus,

$$G^{(1)} = \left\{ -\frac{1}{12}, \frac{7}{240}, -\frac{31}{1344}, \frac{127}{3840}, -\frac{2555}{33792}, \frac{1414477}{5591040}, -\frac{57337}{49152}, \frac{118518239}{16711680}, \dots \right\}.$$

Case $m = 2$. Here we have

$$\tau_1^G = \frac{1}{2} \left(1 - \frac{1}{\sqrt{3}} \right), \quad \tau_2^G = \frac{1}{2} \left(1 + \frac{1}{\sqrt{3}} \right) \quad \text{and} \quad w_1^G = w_2^G = \frac{1}{2},$$

so that $Q_2^G(B_{2j}) = \frac{1}{2} (B_{2j}(\tau_1^G) + B_{2j}(\tau_2^G)) = B_{2j}(\tau_1^G)$. In this case, the sequence of coefficients is

$$G^{(2)} = \left\{ -\frac{1}{180}, \frac{1}{189}, -\frac{17}{2160}, \frac{97}{5346}, -\frac{1291411}{21228480}, \frac{16367}{58320}, -\frac{243615707}{142767360}, \dots \right\}.$$

Case $m = 3$. In this case

$$\tau_1^G = \frac{1}{10} (5 - \sqrt{15}), \quad \tau_2^G = \frac{1}{2}, \quad \tau_3^G = \frac{1}{10} (5 + \sqrt{15})$$

and

$$w_1^G = \frac{5}{18}, \quad w_2^G = \frac{4}{9}, \quad w_3^G = \frac{5}{18},$$

so that

$$Q_3^G(B_{2j}) = \frac{5}{9} B_{2j}(\tau_1^G) + \frac{4}{9} B_{2j}(\tau_2^G)$$

and

$$G^{(3)} = \left\{ -\frac{1}{2800}, \frac{49}{72000}, -\frac{8771}{5280000}, \frac{4935557}{873600000}, -\frac{15066667}{576000000}, \frac{3463953717}{21760000000}, \dots \right\}.$$

Cases $m = 4, 5, 6$. The corresponding sequences of coefficients are

$$G^{(4)} = \left\{ -\frac{1}{44100}, \frac{41}{565950}, -\frac{3076}{11704875}, \frac{93553}{75631500}, -\frac{453586781}{60000990000}, \frac{6885642443}{117354877500}, \dots \right\},$$

$$G^{(5)} = \left\{ -\frac{1}{698544}, \frac{205}{29719872}, -\frac{100297}{2880541440}, \frac{76404959}{352578272256}, -\frac{839025422533}{496513166929920}, \dots \right\},$$

$$G^{(6)} = \left\{ -\frac{1}{11099088}, \frac{43}{70436520}, -\frac{86221}{21074606784}, \frac{147502043}{4534139665440}, -\frac{1323863797}{4200045163776}, \dots \right\}.$$

The Euler–Maclaurin formula based on the composite Lobatto formula can be considered in a similar way. The corresponding Gauss-Lobatto quadrature formula

$$\int_0^1 f(x) dx = \sum_{\nu=0}^{m+1} w_\nu^L f(\tau_\nu^L) + R_m^L(f), \tag{27}$$

with the endnodes $\tau_0 = \tau_0^L = 0, \tau_{m+1} = \tau_{m+1}^L = 1$, has internal nodes $\tau_\nu = \tau_\nu^L, \nu = 1, \dots, m$, which are zeros of the shifted (monic) Jacobi polynomial,

$$\pi_m(x) = \binom{2m+2}{m}^{-1} P_m^{(1,1)}(2x-1),$$

orthogonal on the interval $(0, 1)$ with respect to the weight function $x \mapsto x(1-x)$. The algebraic degree of precision of this formula is $d = 2m + 1$, i.e., $R_m^L(f) = 0$ for each $f \in \mathcal{P}_{2m+1}$.

For constructing the Gauss-Lobatto formula

$$Q_m^L(f) = \sum_{\nu=0}^{m+1} w_\nu^L f(\tau_\nu^L), \tag{28}$$

we use parameters of the corresponding Gaussian formula with respect to the weight function $x \mapsto x(1 - x)$, i.e.,

$$\int_0^1 g(x)x(1 - x) dx = \sum_{\nu=1}^m \widehat{w}_\nu^G g(\widehat{\tau}_\nu^G) + \widehat{R}_m^G(g).$$

The nodes and weights of the Gauss-Lobatto quadrature formula (27) are (cf. [36, pp. 330–331])

$$\tau_0^L = 0, \quad \tau_\nu^L = \widehat{\tau}_\nu^G \quad (\nu = 1, \dots, m), \quad \tau_{m+1}^L = 1,$$

and

$$w_0^L = \frac{1}{2} - \sum_{\nu=1}^m \frac{\widehat{w}_\nu^G}{\widehat{\tau}_\nu^G}, \quad w_\nu^L = \frac{\widehat{w}_\nu^G}{\widehat{\tau}_\nu^G(1 - \widehat{\tau}_\nu^G)} \quad (\nu = 1, \dots, m), \quad w_{m+1}^L = \frac{1}{2} - \sum_{\nu=1}^m \frac{\widehat{w}_\nu^G}{1 - \widehat{\tau}_\nu^G},$$

respectively. The corresponding composite rule is

$$\begin{aligned} L_m^{(n)}f &= \sum_{k=0}^{n-1} Q_m^L f(k + \cdot) = \sum_{\nu=0}^{m+1} w_\nu^L \sum_{k=0}^{n-1} f(k + \tau_\nu^L), \\ &= (w_0^L + w_{m+1}^L) \sum_{k=0}^{n//} f(k) + \sum_{\nu=1}^m w_\nu^L \sum_{k=0}^{n-1} f(k + \tau_\nu^L). \end{aligned} \tag{29}$$

As in the Gauss–Legendre case, there exists a symmetry of nodes and weights, i.e.,

$$\tau_\nu^L + \tau_{m+1-\nu}^L = 1, \quad w_\nu^L = w_{m+1-\nu}^L > 0 \quad \nu = 0, 1, \dots, m + 1,$$

so that the Gauss-Lobatto quadrature sum

$$Q_m^L(B_j) = \sum_{\nu=0}^{m+1} w_\nu^L B_j(\tau_\nu^L) = 0$$

for each odd j .

By the similar arguments as before, we can state and prove the following result.

Theorem 2. For $n, m, r \in \mathbb{N}$ ($m \leq r$) and $f \in C^{2r}[0, n]$ we have

$$L_m^{(n)}f - I_n f = \sum_{j=m+1}^r \frac{Q_m^L(B_{2j})}{(2j)!} [f^{(2j-1)}(n) - f^{(2j-1)}(0)] + E_{n,m,r}^L(f), \tag{30}$$

where $L_m^{(n)}f$ is given by (29), and $Q_m^L B_{2j}$ denotes the basic Gauss-Lobatto quadrature sum (28) applied to the Bernoulli polynomial $x \mapsto B_{2j}(x)$, i.e.,

$$Q_m^L(B_{2j}) = \sum_{v=0}^{m+1} w_v^L B_{2j}(\tau_v^L) = -R_m^L(B_{2j}),$$

where $R_m^L(f)$ is the remainder term in (27).

If $f \in C^{2r+2}[0, n]$, then there exists $\xi \in (0, n)$, such that the error term in (30) can be expressed in the form

$$E_{n,m,r}^L(f) = n \frac{Q_m^L(B_{2r+2})}{(2r+2)!} f^{(2r+2)}(\xi).$$

In the sequel we give the sequence of coefficients $L^{(m)}$ which appear in the sum on the right-hand side in (30), i.e.,

$$L^{(m)} = \{Q_m^L(B_{2j})\}_{j=m+1}^\infty = \{Q_m^L(B_{2m+2}), Q_m^L(B_{2m+4}), Q_m^L(B_{2m+6}), \dots\},$$

obtained by the Package `OrthogonalPolynomials`, for some values of m .

Case $m = 0$. This is a case of the standard Euler–Maclaurin formula (1), for which $\tau_0^L = 0$ and $\tau_1^L = 1$, with $w_0^L = w_1^L = 1/2$. The sequence of coefficients is

$$L^{(0)} = \left\{ \frac{1}{6}, -\frac{1}{30}, \frac{1}{42}, -\frac{1}{30}, \frac{5}{66}, -\frac{691}{2730}, \frac{7}{6}, -\frac{3617}{510}, \frac{43867}{798}, -\frac{174611}{330}, \frac{854513}{138}, \dots \right\},$$

which is, in fact, the sequence of Bernoulli numbers $\{B_{2j}\}_{j=1}^\infty$.

Case $m = 1$. In this case $\tau_0^L = 0$, $\tau_1^L = 1/2$, and $\tau_2 = 1$, with the corresponding weights $w_0^L = 1/6$, $w_1^L = 2/3$, and $w_2^L = 1/6$, which is, in fact, the Simpson formula (15). The sequence of coefficients is

$$L^{(1)} = \left\{ \frac{1}{120}, -\frac{5}{672}, \frac{7}{640}, -\frac{425}{16896}, \frac{235631}{2795520}, -\frac{3185}{8192}, \frac{19752437}{8355840}, -\frac{958274615}{52297728}, \dots \right\}.$$

Case $m = 2$. Here we have

$$\tau_0^L = 0, \quad \tau_1^L = \frac{1}{10}(5 - \sqrt{5}), \quad \tau_2^L = \frac{1}{10}(5 + \sqrt{5}), \quad \tau_3^L = 1$$

and $w_0^L = w_3^L = 1/12$, $w_1^L = w_2^L = 5/12$, and the sequence of coefficients is

$$L^{(2)} = \left\{ \frac{1}{2100}, -\frac{1}{1125}, \frac{89}{41250}, -\frac{25003}{3412500}, \frac{3179}{93750}, -\frac{2466467}{11953125}, \frac{997365619}{623437500}, \dots \right\}.$$

Case $m = 3$. Here the nodes and the weight coefficients are

$$\tau_0^L = 0, \quad \tau_1^L = \frac{1}{14}(7 - \sqrt{31}), \quad \tau_2^L = \frac{1}{2}, \quad \tau_3^L = \frac{1}{14}(7 + \sqrt{31}), \quad \tau_4^L = 1$$

and

$$w_0^L = \frac{1}{20}, \quad w_1^L = \frac{49}{180}, \quad w_2^L = \frac{16}{45}, \quad w_3^L = \frac{49}{180}, \quad w_4^L = \frac{1}{20},$$

respectively, and the sequence of coefficients is

$$L^{(3)} = \left\{ \frac{1}{35280}, -\frac{65}{724416}, \frac{38903}{119857920}, -\frac{236449}{154893312}, \frac{1146165227}{122882027520}, \dots \right\}.$$

Cases $m = 4, 5$. The corresponding sequences of coefficients are

$$L^{(4)} = \left\{ \frac{1}{582120}, -\frac{17}{2063880}, \frac{173}{4167450}, -\frac{43909}{170031960}, \frac{160705183}{79815002400}, -\frac{76876739}{3960744480}, \dots \right\},$$

$$L^{(5)} = \left\{ \frac{1}{9513504}, -\frac{49}{68999040}, \frac{5453}{1146917376}, -\frac{671463061}{17766424811520}, \frac{1291291631}{3526568534016}, \dots \right\}.$$

Remark 8. Recently Dubeau [16] has shown that an Euler–Maclaurin like formula can be associated with any interpolatory quadrature rule.

4 Abel–Plana Summation Formula and Some Modifications

Another important summation formula is the so-called Abel–Plana formula, but it is not so well known like the Euler–Maclaurin formula. In 1820 Giovanni (Antonio Amedea) Plana [49] obtained the summation formula

$$\sum_{k=0}^{+\infty} f(k) - \int_0^{+\infty} f(x) dx = \frac{1}{2}f(0) + i \int_0^{+\infty} \frac{f(iy) - f(-iy)}{e^{2\pi y} - 1} dy, \tag{31}$$

which holds for analytic functions f in $\Omega = \{z \in \mathbb{C} : \operatorname{Re} z \geq 0\}$ which satisfy the conditions:

- 1° $\lim_{|y| \rightarrow +\infty} e^{-|2\pi y|} |f(x \pm iy)| = 0$ uniformly in x on every finite interval,
- 2° $\int_0^{+\infty} |f(x + iy) - f(x - iy)| e^{-|2\pi y|} dy$ exists for every $x \geq 0$ and tends to zero when $x \rightarrow +\infty$.

This formula was also proved in 1823 by Niels Henrik Abel [1]. In addition, Abel also proved an interesting “*alternating series version*”, under the same conditions,

$$\sum_{k=0}^{+\infty} (-1)^k f(k) = \frac{1}{2} f(0) + i \int_0^{+\infty} \frac{f(iy) - f(-iy)}{2 \sinh \pi y} dy. \tag{32}$$

Otherwise, this formula can be obtained only from (31). Note that, by subtracting (31) from the same formula written for the function $z \mapsto 2f(2z)$, we get (32).

For the finite sum $S_{n,m}f = \sum_{k=m}^n (-1)^k f(k)$, (32) the *Abel summation formula* becomes

$$S_{n,m}f = \frac{1}{2} [(-1)^m f(m) + (-1)^n f(n + 1)] - \int_{-\infty}^{+\infty} [(-1)^m \psi_m(y) + (-1)^n \psi_{n+1}(y)] w^A(y) dy, \tag{33}$$

where the *Abel weight* on \mathbb{R} and the function $\phi_m(y)$ are given by

$$w^A(x) = \frac{x}{2 \sinh \pi x} \quad \text{and} \quad \phi_m(y) = \frac{f(m + iy) - f(m - iy)}{2iy}. \tag{34}$$

The moments for the Abel weight can be expressed in terms of Bernoulli numbers as

$$\mu_k = \begin{cases} 0, & k \text{ odd,} \\ (2^{k+2} - 1) \frac{(-1)^{k/2} B_{k+2}}{k + 2}, & k \text{ even.} \end{cases} \tag{35}$$

A general Abel–Plana formula can be obtained by a contour integration in the complex plane. Let $m, n \in \mathbb{N}$, $m < n$, and $C(\varepsilon)$ be a closed rectangular contour with vertices at $m \pm ib$, $n \pm ib$, $b > 0$ (see Fig. 2), and with semicircular indentations of radius ε round m and n . Let f be an analytic function in the strip $\Omega_{m,n} = \{z \in \mathbb{C} : m \leq \text{Re } z \leq n\}$ and suppose that for every $m \leq x \leq n$,

$$\lim_{|y| \rightarrow +\infty} e^{-|2\pi y|} |f(x \pm iy)| = 0 \quad \text{uniformly in } x,$$

and that

$$\int_0^{+\infty} |f(x + iy) - f(x - iy)| e^{-|2\pi y|} dy$$

exists.

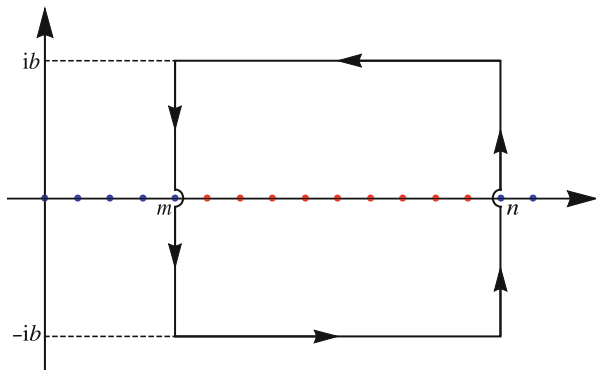


Fig. 2 Rectangular contour $C(\varepsilon)$

The integration

$$\int_{C(\varepsilon)} \frac{f(z)}{e^{-i2\pi z} - 1} dz,$$

with $\varepsilon \rightarrow 0$ and $b \rightarrow +\infty$, leads to the *Plana formula* in the following form (cf. [42])

$$T_{m,n}f - \int_m^n f(x) dx = \int_{-\infty}^{+\infty} (\phi_n(y) - \phi_m(y))w^P(y) dy, \tag{36}$$

where

$$\phi_m(y) = \frac{f(m + iy) - f(m - iy)}{2iy} \quad \text{and} \quad w^P(y) = \frac{|y|}{e^{2\pi|y|} - 1}. \tag{37}$$

Practically, the Plana formula (36) gives the error of the composite trapezoidal formula (like the Euler–Maclaurin formula). As we can see the formula (36) is similar to the Euler–Maclaurin formula, with the difference that the sum of terms

$$\frac{B_{2j}}{(2j)!} (f^{(2j-1)}(n) - f^{(2j-1)}(m))$$

replaced by an integral. Therefore, in applications this integral must be calculated by some quadrature rule. It is natural to construct the Gaussian formula with respect to the *Plana* weight function $x \mapsto w^P(x)$ on \mathbb{R} (see the next section for such a construction).

In order to find the moments of this weight function, we note first that if k is odd, the moments are zero, i.e.,

$$\mu_k(w^P) = \int_{\mathbb{R}} x^k w^P(x) dx = \int_{\mathbb{R}} x^k \frac{|x|}{e^{2\pi|x|} - 1} dx = 0.$$

For even k , we have

$$\mu_k(w^P) = 2 \int_0^{+\infty} \frac{x^{k+1}}{e^{2\pi x} - 1} dx = \frac{2}{(2\pi)^{k+2}} \int_0^{+\infty} \frac{t^{k+1}}{e^t - 1} dt,$$

which can be exactly expressed in terms of the Riemann zeta function $\zeta(s)$,

$$\mu_k(w^P) = \frac{2(k+1)!\zeta(k+2)}{(2\pi)^{k+2}} = (-1)^{k/2} \frac{B_{k+2}}{k+2},$$

because the number $k+2$ is even. Thus, in terms of Bernoulli numbers, the moments are

$$\mu_k(w^P) = \begin{cases} 0, & k \text{ is odd,} \\ (-1)^{k/2} \frac{B_{k+2}}{k+2}, & k \text{ is even.} \end{cases} \tag{38}$$

Remark 9. By the Taylor expansion for $\phi_m(y)$ (and $\phi_n(y)$) on the right-hand side in (36),

$$\phi_m(y) = \frac{f(m+iy) - f(m-iy)}{2iy} = \sum_{j=1}^{+\infty} \frac{(-1)^{j-1} y^{2j-2}}{(2j-1)!} f^{(2j-1)}(m),$$

and using the moments (38), the Plana formula (36) reduces to the Euler–Maclaurin formula,

$$\begin{aligned} T_{m,n}f - \int_m^n f(x) dx &= \sum_{j=1}^{+\infty} \frac{(-1)^{j-1}}{(2j-1)!} \mu_{2j-2}(w^P) (f^{(2j-1)}(n) - f^{(2j-1)}(m)) \\ &= \sum_{j=1}^{+\infty} \frac{B_{2j}}{(2j)!} (f^{(2j-1)}(n) - f^{(2j-1)}(m)), \end{aligned}$$

because of $\mu_{2j-2}(w^P) = (-1)^{j-1} B_{2j}/(2j)$. Note that $T_{m,n}f$ is the notation for the composite trapezoidal sum

$$T_{m,n}f := \sum_{k=m}^n f(k) = \frac{1}{2}f(m) + \sum_{k=m+1}^{n-1} f(k) + \frac{1}{2}f(n). \tag{39}$$

For more details see Rahman and Schmeisser [53, 54], Dahlquist [11–13], as well as a recent paper by Butzer, Ferreira, Schmeisser, and Stens [8].

A similar summation formula is the so-called *midpoint summation formula*. It can be obtained by combining two Plana formulas for the functions $z \mapsto f(z-1/2)$

and $z \mapsto f((z + m - 1)/2)$. Namely,

$$T_{m,2n-m+2}f\left(\frac{z + m - 1}{2}\right) - T_{m,n+1}f\left(z - \frac{1}{2}\right) = \sum_{k=m}^n f(k),$$

i.e.,

$$\sum_{k=m}^n f(k) - \int_{m-1/2}^{n+1/2} f(x) \, dx = \int_{-\infty}^{+\infty} [\phi_{m-1/2}(y) - \phi_{n+1/2}(y)]w^M(y) \, dy, \tag{40}$$

where the *midpoint weight function* is given by

$$w^M(x) = w^P(x) - w^P(2x) = \frac{|x|}{e^{2\pi|x|} + 1}, \tag{41}$$

and $\phi_{m-1/2}$ and $\phi_{n+1/2}$ are defined in (37), taking $m := m - 1/2$ and $m := n + 1/2$, respectively. The moments for the midpoint weight function can be expressed also in terms of Bernoulli numbers as

$$\mu_k(w^M) = \int_{\mathbb{R}} x^k \frac{|x|}{e^{2\pi|x|} + 1} \, dx = \begin{cases} 0, & k \text{ is odd,} \\ (-1)^{k/2}(1 - 2^{-(k+1)})\frac{B_{k+2}}{k+2}, & k \text{ is even.} \end{cases} \tag{42}$$

An interesting weight function and the corresponding summation formula can be obtained from the Plana formula, if we integrate by parts the right side in (36) (cf. [13]). Introducing the so-called *Binet* weight function $y \mapsto w^B(y)$ and the function $y \mapsto \psi_m(y)$ by

$$w^B(y) = -\frac{1}{2\pi} \log(1 - e^{-2\pi|y|}) \quad \text{and} \quad \psi_m(y) = \frac{f'(m + iy) + f'(m - iy)}{2}, \tag{43}$$

respectively, we see that $dw^B(y)/dy = -w^P(y)/y$ and

$$\begin{aligned} \frac{d}{dy}\{[\phi_n(y) - \phi_m(y)]y\} &= \frac{1}{2i} \frac{d}{dy} \{[f(n + iy) - f(n - iy)] - [f(m + iy) - f(m - iy)]\} \\ &= \psi_n(y) - \psi_m(y), \end{aligned}$$

so that

$$\begin{aligned} \int_{-\infty}^{+\infty} [\phi_n(y) - \phi_m(y)]w^P(y) \, dy &= \int_{-\infty}^{+\infty} [\phi_n(y) - \phi_m(y)](-y) \, dw^B(y) \\ &= \int_{-\infty}^{+\infty} [\psi_n(y) - \psi_m(y)]w^B(y) \, dy, \end{aligned}$$

because $w^B(y) = O(e^{-2\pi|y|})$ as $|y| \rightarrow +\infty$. Thus, the *Binet summation formula* becomes

$$T_{m,n}f - \int_m^n f(x) dx = \int_{-\infty}^{+\infty} [\psi_n(y) - \psi_m(y)]w^B(y) dy. \tag{44}$$

Such a formula can be useful when $f'(z)$ is easier to compute than $f(z)$.

The moments for the Binet weight can be obtained from ones for w^P . Since

$$\mu_k(w^P) = \int_{\mathbb{R}} y^k w^P(y) dy = \int_{\mathbb{R}} y^k (-y) dw^B(y) = (k + 1)\mu_k(w^B),$$

according to (38),

$$\mu_k(w^B) = \begin{cases} 0, & k \text{ is odd,} \\ (-1)^{k/2} \frac{B_{k+2}}{(k + 1)(k + 2)}, & k \text{ is even.} \end{cases} \tag{45}$$

There are also several other summation formulas. For example, the *Lindelöf formula* [32] for alternating series is

$$\sum_{k=m}^{+\infty} (-1)^k f(k) = (-1)^m \int_{-\infty}^{+\infty} f(m - 1/2 + iy) \frac{dy}{2 \cosh \pi y}, \tag{46}$$

where the *Lindelöf* weight function is given by

$$w^L(x) = \frac{1}{2 \cosh \pi y} = \frac{1}{e^{\pi x} + e^{-\pi x}}. \tag{47}$$

Here, the moments

$$\mu_k(w^L) = \int_{\mathbb{R}} \frac{x^k}{e^{\pi x} + e^{-\pi x}} dx$$

can be expressed in terms of the generalized Riemann zeta function $z \mapsto \zeta(z, a)$, defined by

$$\zeta(z, a) = \sum_{v=0}^{+\infty} (v + a)^{-z}.$$

Namely,

$$\mu_k(w^L) = \begin{cases} 0, & k \text{ odd,} \\ 2(4\pi)^{-k-1} k! [\zeta(k + 1, \frac{1}{4}) - \zeta(k + 1, \frac{3}{4})], & k \text{ even.} \end{cases} \tag{48}$$

5 Construction of Orthogonal Polynomials and Gaussian Quadratures for Weights of Abel–Plana Type

The weight functions $w \in \{w^P, w^M, w^B, w^A, w^L\}$ which appear in the summation formulas considered in the previous section are even functions on \mathbb{R} . In this section we consider the construction of (monic) orthogonal polynomials $\pi_k \equiv \pi_k(w; \cdot)$ and corresponding Gaussian formulas

$$\int_{\mathbb{R}} f(x)w(x) dx = \sum_{\nu=1}^n A_{\nu}f(x_{\nu}) + R_n(w;f), \tag{49}$$

with respect to the inner product $(p, q) = \int_{\mathbb{R}} p(x)q(x)w(x) dx$ ($p, q \in \mathcal{P}$). We note that $R_n(w;f) \equiv 0$ for each $f \in \mathcal{P}_{2n-1}$.

Such orthogonal polynomials $\{\pi_k\}_{k \in \mathbb{N}_0}$ and Gaussian quadratures (49) exist uniquely, because all the moments for these weights $\mu_k \equiv \mu_k(w)$, $k = 0, 1, \dots$, exist, are finite, and $\mu_0 > 0$.

Because of the property $(xp, q) = (p, xq)$, these (monic) orthogonal polynomials π_k satisfy the fundamental *three-term recurrence relation*

$$\pi_{k+1}(x) = x\pi_k(x) - \beta_k\pi_{k-1}(x), \quad k = 0, 1, \dots, \tag{50}$$

with $\pi_0(x) = 1$ and $\pi_{-1}(x) = 0$, where $\{\beta_k\}_{k \in \mathbb{N}_0} (= \{\beta_k(w)\}_{k \in \mathbb{N}_0})$ is a sequence of recursion coefficients which depend on the weight w . The coefficient β_0 may be arbitrary, but it is conveniently defined by $\beta_0 = \mu_0 = \int_{\mathbb{R}} w(x) dx$. Note that the coefficients α_k in (50) are equal to zero, because the weight function w is an even function! Therefore, the nodes in (49) are symmetrically distributed with respect to the origin, and the weights for symmetrical nodes are equal. For odd n one node is at zero.

A characterization of the Gaussian quadrature (49) can be done via an eigenvalue problem for the symmetric tridiagonal Jacobi matrix (cf. [36, p. 326]),

$$J_n = J_n(w) = \begin{bmatrix} \alpha_0 & \sqrt{\beta_1} & & & \mathbf{0} \\ \sqrt{\beta_1} & \alpha_1 & \sqrt{\beta_2} & & \\ & \sqrt{\beta_2} & \alpha_2 & \ddots & \\ & & \ddots & \ddots & \sqrt{\beta_{n-1}} \\ \mathbf{0} & & & \sqrt{\beta_{n-1}} & \alpha_{n-1} \end{bmatrix},$$

constructed with the coefficients from the three-term recurrence relation (50) (in our case, $\alpha_k = 0, k = 0, 1, \dots, n - 1$).

The nodes x_{ν} are the eigenvalues of J_n and the weights A_{ν} are given by $A_{\nu} = \beta_0 v_{\nu,1}^2, \nu = 1, \dots, n$, where β_0 is the moment $\mu_0 = \int_{\mathbb{R}} w(x) dx$, and $v_{\nu,1}$ is the first component of the normalized eigenvector $\mathbf{v} = [v_{\nu,1} \ \dots \ v_{\nu,n}]^T$ (with $\mathbf{v}_{\nu}^T \mathbf{v}_{\nu} = 1$)

corresponding to the eigenvalue x_ν ,

$$J_n \mathbf{v}_\nu = x_\nu \mathbf{v}_\nu, \quad \nu = 1, \dots, n.$$

An efficient procedure for constructing the Gaussian quadrature rules was given by Golub and Welsch [27], by simplifying the well-known QR algorithm, so that only the first components of the eigenvectors are computed.

The problems are very sensitive with respect to small perturbations in the data.

Unfortunately, the recursion coefficients are known explicitly only for some narrow classes of orthogonal polynomials, as e.g. for the classical orthogonal polynomials (Jacobi, the generalized Laguerre, and Hermite polynomials). However, for a large class of the so-called *strongly non-classical polynomials* these coefficients can be constructed numerically, but procedures are very sensitive with respect to small perturbations in the data. Basic procedures for generating these coefficients were developed by Walter Gautschi in the eighties of the last century (cf. [23, 24, 36, 41]).

However, because of progress in symbolic computations and variable-precision arithmetic, recursion coefficients can be today directly generated by using the original Chebyshev method of moments (cf. [36, pp. 159–166]) in symbolic form or numerically in sufficiently high precision. In this way, instability problems can be eliminated. Respectively symbolic/variable-precision software for orthogonal polynomials and Gaussian and similar type quadratures is available. In this regard, the MATHEMATICA package `OrthogonalPolynomials` (see [9] and [43]) is downloadable from the web site <http://www.mi.sanu.ac.rs/~gvm/>. Also, there is Gautschi's software in MATLAB (packages `OPQ` and `SOPQ`). Thus, all that is required is a procedure for the symbolic calculation of moments or their calculation in variable-precision arithmetic.

In our case we calculate the first $2N$ moments in a symbolic form (list `mom`), using corresponding formulas (for example, (38) in the case of the Plana weight w^P), so that we can construct the Gaussian formula (49) for each $n \leq N$. Now, in order to get the first N recurrence coefficients $\{a_l, b_e\}$ in a symbolic form, we apply the implemented function `aChebyshevAlgorithm` from the Package `OrthogonalPolynomials`, which performs construction of these coefficients using Chebyshev algorithm, with the option `Algorithm->Symbolic`. Thus, it can be implemented in the MATHEMATICA package `OrthogonalPolynomials` in a very simple way as

```
<<orthogonalPolynomials`
mom=Table[<expression for moments>,{k,0,199}];
{a1,b1}=aChebyshevAlgorithm[mom,Algorithm->Symbolic]
pq[n_]:=aGaussianNodesWeights[n,a1,b1,
WorkingPrecision->65,Precision->60]
xA = Table[pq[n],{n,5,40,5}];
```

where we put $N = 100$ and the `WorkingPrecision->65` in order to obtain very precisely quadrature parameters (nodes and weights) with `Precision->60`. These parameters are calculated for $n = 5(5)40$, so that `xA[[k]][[1]]` and `xA[[k]][[2]]` give lists of nodes and weights for five-point formula when $k=1$, for ten-point formula when $k=2$, etc. Otherwise, here we can calculate the n -point Gaussian quadrature formula for each $n \leq N = 100$.

All computations were performed in `MATHEMATICA`, Ver. 10.3.0, on MacBook Pro (Retina, Mid 2012) OS X 10.11.2. The calculations are very fast. The running time is evaluated by the function `Timing` in `MATHEMATICA` and it includes only CPU time spent in the `MATHEMATICA` kernel. Such a way may give different results on different occasions within a session, because of the use of internal system caches. In order to generate worst-case timing results independent of previous computations, we used also the command `ClearSystemCache[]`, and in that case the running time for the Plana weight function w^P has been 4.2 ms (calculation of moments), 0.75 s (calculation of recursive coefficients), and 8 s (calculation quadrature parameters for $n = 5(5)40$).

In the sequel we mention results for different weight functions, whose graphs are presented in Fig. 3.

1. Abel and Lindelöf Weight Functions w^A and w^L These weight functions are given by (34) and (47), and their moments by (35) and (48), respectively. It is interesting that their corresponding coefficients in the three-term recurrence relation (50) are known explicitly (see [36, p. 159])

$$\beta_0^A = \mu_0^A = \frac{1}{4}, \quad \beta_k^A = \frac{k(k+1)}{4}, \quad k = 1, 2, \dots,$$

and

$$\beta_0^L = \mu_0^L = \frac{1}{2}, \quad \beta_k^L = \frac{k^2}{4}, \quad k = 1, 2, \dots$$

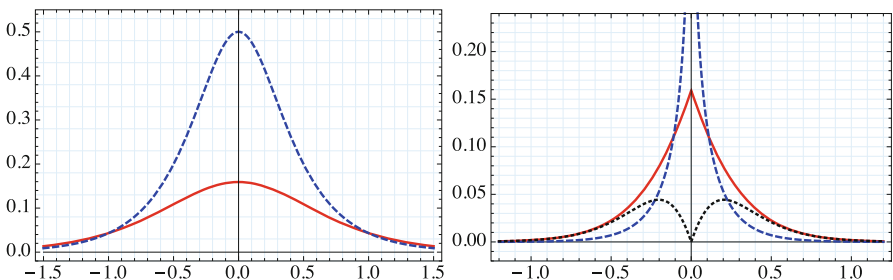


Fig. 3 Graphs of the weight functions: (left) w^A (solid line) and w^L (dashed line); (right) w^P (solid line), w^B (dashed line) and w^M (dotted line)

Thus, for these two weight functions we have recursive coefficients in the explicit form, so that we go directly to construction quadrature parameters.

2. Plana Weight Function w^P This weight function is given by (37), and the corresponding moments by (38). Using the Package `OrthogonalPolynomials` we obtain the sequence of recurrence coefficients $\{\beta_k^P\}_{k \geq 0}$ in the rational form:

$$\begin{aligned} \beta_0^P &= \frac{1}{12}, \quad \beta_1^P = \frac{1}{10}, \quad \beta_2^P = \frac{79}{210}, \quad \beta_3^P = \frac{1205}{1659}, \quad \beta_4^P = \frac{262445}{209429}, \quad \beta_5^P = \frac{33461119209}{18089284070}, \\ \beta_6^P &= \frac{361969913862291}{137627660760070}, \quad \beta_7^P = \frac{85170013927511392430}{24523312685049374477}, \\ \beta_8^P &= \frac{1064327215185988443814288995130}{236155262756390921151239121153}, \\ \beta_9^P &= \frac{286789982254764757195675003870137955697117}{51246435664921031688705695412342990647850}, \\ \beta_{10}^P &= \frac{15227625889136643989610717434803027240375634452808081047}{2212147521291103911193549528920437912200375980011300650}, \\ \beta_{11}^P &= \frac{587943441754746283972138649821948554273878447469233852697401814148410885}{71529318090286333175985287358122471724664434392542372273400541405857921}, \\ &\text{etc.} \end{aligned}$$

As we can see, the fractions are becoming more complicated, so that already β_{11}^P has the “form of complexity” $\{72/71\}$, i.e., it has 72 decimal digits in the numerator and 71 digits in the denominator. Further terms of this sequence have the “form of complexity” $\{88/87\}$, $\{106/05\}$, $\{129/128\}$, $\{152/151\}$, ..., $\{13451/13448\}$.

Thus, the last term β_{99}^P has more than 13 thousand digits in its numerator and denominator. Otherwise, its value, e.g. rounded to 60 decimal digits, is

$$\beta_{99}^P = 618.668116294139071216871819412846078447729830182674784697227.$$

3. Midpoint Weight Function w^M This weight function is given by (41), and the corresponding moments by (42). As in the previous case, we obtain the sequence of recurrence coefficients $\{\beta_k^M\}_{k \geq 0}$ in the rational form:

$$\begin{aligned} \beta_0^M &= \frac{1}{24}, \quad \beta_1^M = \frac{7}{40}, \quad \beta_2^M = \frac{2071}{5880}, \quad \beta_3^M = \frac{999245}{1217748}, \quad \beta_4^M = \frac{21959166635}{18211040276}, \\ \beta_5^M &= \frac{108481778600414331}{55169934195679160}, \quad \beta_6^M = \frac{2083852396915648173441543}{813782894744588335008520}, \\ \beta_7^M &= \frac{25698543837390957571411809266308135}{7116536885169433586426285918882662}, \\ \beta_8^M &= \frac{202221739836050724659312728605015618097349555485}{45788344599633183797631374444694817538967629598}, \\ \beta_9^M &= \frac{14077564493254853375144075652878384268409784777236869234539068357}{2446087170499983327141705915330961521888001335934900402777402200}, \end{aligned}$$

etc. In this case, the last term β_{99}^M has slightly complicated the “form of complexity” {16401/16398} than one in the previous case, precisely. Otherwise, its value (rounded to 60 decimal digits) is

$$\beta_{99}^M = 619.562819405146668677971154899553589896235540274133472854031.$$

4. Binet Weight Function w^B The moments for this weight function are given in (38), and our Package `OrthogonalPolynomials` gives the sequence of recurrence coefficients $\{\beta_k^B\}_{k \geq 0}$ in the rational form:

$$\beta_0^B = \frac{1}{12}, \quad \beta_1^B = \frac{1}{30}, \quad \beta_2^B = \frac{53}{210}, \quad \beta_3^B = \frac{195}{371}, \quad \beta_4^B = \frac{22999}{22737}, \quad \beta_5^B = \frac{29944523}{19733142},$$

$$\beta_6^B = \frac{109535241009}{48264275462}, \quad \beta_7^B = \frac{29404527905795295658}{9769214287853155785},$$

$$\beta_8^B = \frac{455377030420113432210116914702}{113084128923675014537885725485},$$

$$\beta_9^B = \frac{26370812569397719001931992945645578779849}{5271244267917980801966553649147604697542},$$

$$\beta_{10}^B = \frac{152537496709054809881638897472985990866753853122697839}{24274291553105128438297398108902195365373879212227726},$$

$$\beta_{11}^B = \frac{100043420063777451042472529806266909090824649341814868347109676190691}{13346384670164266280033479022693768890138348905413621178450736182873},$$

etc. Numerical values of coefficients β_k^B for $k = 12, \dots, 39$, rounded to 60 decimal digits, are presented in Table 1.

For this case we give also quadrature parameters x_v^B and A_v^B , $v = 1, \dots, n$, for $n = 10$ (rounded to 30 digits in order to save space). Numbers in parenthesis indicate the decimal exponents (Table 2).

Table 1 Numerical values of the coefficients $\beta_k^B, k = 12, \dots, 39$

k	β_k^B
12	9.04066023436772669953113936026048174933621963537072222675357
13	10.4893036545094822771883713045926295220972379893834049993209
14	12.2971936103862058639894371400919176597365509004516453610177
15	13.9828769539924301882597606512787300859080333154700506431789
16	16.0535514167049354697156163650062601783515764970917711361702
17	17.9766073998702775925694723076715543993147838556500117187847
18	20.3097620274416537438054147204948968937016485345196881526453
19	22.4704716399331324955179415715079221089953862901823520893038
20	25.0658465489459720291634003225063053682385176354570207084270
21	27.4644518250291336091755589826462226732286473857913864921713
22	30.3218212316730471268825993064057869944873787313809977426698
23	32.9585339299729872199940664514120882069601000999724796349878
24	36.0776989312992426451533209008554523367760033115543468301504
25	38.9527066823115557345443904104810462991593233805616588397077
26	42.3334900435769572113818539488560973399147861411953446717663
27	45.4469608500616210144241757375414510828484368311407665782656
28	49.0892031290125977081648833502750872924491998898068036677541
29	52.4412887514153373125698560469961084271478607455930155529787
30	56.3448453453418435384203659474761135421333046623523607025848
31	59.9356839071658582078525834927521121101345464090376940621335
32	64.1004227559203545279066118922379177529092202107679570943670
33	67.9301407880182211863677027451985358165225510069351193013587
34	72.3559405552117019696800529632362179107517585345562462880100
35	76.4246546268296897525850904222875264035700459112308348153069
36	81.1114032372479654848142309856834609745026942246296395824649
37	85.4192212764109726145856387173486827269888223681684704599999
38	90.3668147238641085955135745816833777807870911939721581625005
39	94.9138371000098879530762312919869274587678241868936940165561

Table 2 Gaussian quadrature parameters x_ν^B and $A_\nu^B, \nu = 1, \dots, n$, for ten-point rule

ν	$x_{\nu+5}^B (= -x_{6-\nu}^B)$	$A_{\nu+5}^B (= A_{6-\nu}^B)$
1	1.19026134410869931041299717296(-1)	3.95107541334705577733788440045(-2)
2	5.98589257742219693357956162107(-1)	2.10956883221363967243739596594(-3)
3	1.25058028819024934653033542222	4.60799503427397559669146065886(-5)
4	2.12020925569172605355904853247	2.63574272352001106479781030329(-7)
5	3.34927819645835833349223106504	1.76367377463777032308587486531(-10)

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