Local Anomaly Cancellation and Equivariant Cohomology of Jet Bundles

Roberto Ferreiro Pérez

Dedicated to Jaime Muñoz Masqué on the occasion of his 65th birthday, with my best wishes.

Abstract We study the problem, suggested by Singer in [\[17\]](#page-5-0), and consisting in determining a notion of "local cohomology" adequate to deal with the problem of locality in those approaches to local anomalies based on the Atiyah–Singer index theorem.

Keywords Local cohomology · Equivariant cohomology · Jet bundle · Anomaly cancellation

1 Introduction

An anomaly appears in a theory when a classical symmetry is broken at the quantum level. As we consider only local anomalies, we can assume that the group *G* is connected. Let $\mathcal{L}(\psi, s)$ be a *G*-invariant Lagrangian density depending on bosonic fields $s \in \Gamma(E)$ and fermionic fields ψ . At the quantum level, the corresponding effective action $W(s)$, defined in terms of the fermionic path integral by $\exp(-W(s)) = \int \mathcal{D}\psi \mathcal{D}\psi \exp(-\int_M \mathcal{L}(\psi, s))$ could fail to be \mathcal{G} -invariant. We define a form $A \in \Omega^1(\text{Lie } \mathcal{G}, \Omega^0(\Gamma(E)))$ by $A = \delta W$, i.e. $A(X)(s) = L_X W(s)$ for $X \in \text{Lie } \mathcal{G}, s \in \Gamma(E)$. Although *W* is clearly a non-local functional, *A* is local in *X* and *s*, i.e. we have $A \in \Omega^1_{loc}(\text{Lie } \mathcal{G}, \Omega^0_{loc}(\Gamma(E)))$. It is clear that *A* satisfies the

R. Ferreiro Pérez (B)

Facultad de Ciencias Económicas y Empresariales, Departamento de Economa

Financiera y Contabilidad I, Universidad Complutense de Madrid, Campus de Somosaguas,

e-mail: roferreiro@ccee.ucm.es

²⁸²²³ Pozuelo de Alarcón, Spain

[©] Springer International Publishing Switzerland 2016

M. Castrillón López et al. (eds.), *Geometry, Algebra and Applications: From Mechanics to Cryptography*, Springer Proceedings

in Mathematics & Statistics 161, DOI 10.1007/978-3-319-32085-4_9

condition $\delta A = 0$ (the Wess–Zumino consistency condition). Moreover, if $A = \delta \Lambda$ for a *local* functional $\Lambda = \int_M \lambda \in \Omega^0_{loc}(\Gamma(E))$ then we can define a new lagrangian density $\hat{\mathcal{L}} = \mathcal{L} + \lambda$, such that the new effective action \hat{W} is *G*-invariant, and in that case the anomaly cancels. If $A \neq \delta \Lambda$ for every $\Lambda \in \Omega_{\text{loc}}^0(\Gamma(E))$ then we say that there exists an anomaly in the theory. Hence the anomaly is measured by the cohomology class of *A* in the BRST cohomology $H_{loc}^1(\text{Lie } \mathcal{G}, \Omega_{loc}^0(\Gamma(E)))$ (e.g. see [\[4,](#page-4-0) [6](#page-5-1)[–8,](#page-5-2) [16](#page-5-3)]).

Local anomalies also admit a nice geometrical interpretation in terms of the Atiyah–Singer index theorem for families of elliptic operators (see [\[1,](#page-4-1) [2](#page-4-2), [4,](#page-4-0) [17](#page-5-0)]). The first Chern class c_1 (det Ind*D*) $\in H^2(\Gamma(E)/\mathcal{G})$ of the determinant line bundle det Ind $D \rightarrow \Gamma(E)/\mathcal{G}$ represents an obstruction for anomaly cancellation. However, the condition c_1 (det Ind*D*) = 0 is a necessary but not a sufficient condition for local anomaly cancellation due to the problem of locality. In [\[17\]](#page-5-0) (see also [\[1\]](#page-4-1)) Singer proposes the problem of defining a notion of "local cohomology of Γ (*E*)/*G*", $H^2_{\text{loc}}(\Gamma(E)/\mathcal{G})$, adequate to study local anomaly cancellation. The principal difficulty is the fact that the expression of the curvature of det Ind*D* itself contains non-local terms (Green operators).

Moreover, we recall (see [\[2](#page-4-2), [5,](#page-4-3) [15\]](#page-5-4)) that the BRST and index theory approaches are related by means of the transgression map t (see Sect. [2\)](#page-1-0), i.e., we have $[A]$ = $t(c_1$ (det Ind*D*)). As *t* is injective, the condition c_1 (det Ind*D*) = 0 on $H^2(\Gamma(E)/\mathcal{G})$ is equivalent to $[A] = 0$ on $H^1(\text{Lie }\mathcal{G}, \Omega^0(\Gamma(E)))$. However, the condition for local anomaly cancellation is $[A] = 0$ on the BRST cohomology $H_{loc}^1(\text{Lie } \mathcal{G}, \Omega_{loc}^0(\Gamma(E)))$. We define $H_{\text{loc}}^{\bullet}(\Gamma(E)/\mathcal{G})$ in such a way that the preceding condition is equivalent to the vanishing of the class of *c*₁ (det Ind *D*) on $H_{\text{loc}}^2(\Gamma(E)/\mathcal{G})$, hence solving Singer's problem.

2 The Transgression Maps

First we recall the definition of equivariant cohomology in the Cartan model (*e.g.* see [\[3](#page-4-4)]). We consider a left action of a connected Lie group G on a manifold *N*. We have an induced Lie algebra homomorphism Lie $G \rightarrow \mathfrak{X}(\mathcal{N})$, $X \mapsto$ $X_N = \frac{d}{dt}\Big|_{t=0}$ $\rho(\exp(-tX))$. We denote by $\mathcal{P}^k(\text{Lie } \mathcal{G}, \Omega^r(\mathcal{N}))^{\mathcal{G}}$ the space of degree *k G*-invariant polynomials on Lie *G* with values in $\Omega^r(\mathcal{N})$. We recall that $\alpha \in \mathbb{R}$ \mathcal{P}^k (Lie $\mathcal{G}, \Omega^r(\mathcal{N})$)^{\mathcal{G}} if and only if $\alpha(\text{Ad}_{g}X) = \rho(g^{-1})^*(\alpha(X))$ $\forall X \in \text{Lie } \mathcal{G}, \forall g \in \mathcal{G}$.

The space of *G*-equivariant differential *q*-forms is defined by

$$
\Omega_{\mathcal{G}}^q(\mathcal{N}) = \bigoplus_{2k+r=q} (\mathcal{P}^k(\text{Lie }\mathcal{G}, \Omega^r(\mathcal{N})))^{\mathcal{G}}.
$$
 (1)

The Cartan differential d_c : $\Omega_g^q(\mathcal{N}) \to \Omega_g^{q+1}(\mathcal{N}), (d_c \alpha)(X) = d(\alpha(X)) - \iota_{X_N} \alpha(X)$ for $X \in \text{Lie } \mathcal{G}$, satisfies $(d_c)^2 = 0$, and the \mathcal{G} -equivariant cohomology of \mathcal{N} , $H^q_{\mathcal{G}}(\mathcal{N})$, is the cohomology of this complex.

Local Anomaly Cancellation … 103

We recall (e.g. see [\[3\]](#page-4-4)) that if $\mathcal{N} \to \mathcal{N}/\mathcal{G}$ is a principal \mathcal{G} -bundle we have the (generalized) Chern–Weil homomorphism ChW: $H^{\bullet}_{\mathcal{G}}(\mathcal{N}) \to H^{\bullet}(\mathcal{N}/\mathcal{G})$. If *A* is an arbitrary connection on $\mathcal{N} \to \mathcal{N}/\mathcal{G}$ with curvature F_A , and $\alpha \in \Omega_g^q(\mathcal{N})$, then we have $ChW([\alpha]) = [hor_A(\alpha(F_A))]$, where hor_A is the horizontalization with respect to the connection *A*. We also use the notation $\alpha = \text{ChW}(\alpha)$.

If $\omega \in \Omega^2_{\mathcal{G}}(\mathcal{N})$ is a closed *G*-equivariant 2-form, then we have $\omega = \omega_0 + \mu$ where $\omega_0 \in \Omega^2(\mathcal{N})$ is closed, and $\mu:$ Lie $\mathcal{G} \to C^\infty(\mathcal{N})$, is a \mathcal{G} -equivariant moment map for ω_0 , i.e., $\iota_{X_N} \omega_0 = d(\mu(X))$ for $X \in \text{Lie } \mathcal{G}$. A direct computation shows that we have the following

Proposition 1 *Assume that* $\mathcal{N} \rightarrow \mathcal{N}/\mathcal{G}$ *is a principal* \mathcal{G} *-bundle, and let* $A \in$ $\Omega^1(\mathcal{N}, \text{Lie }\mathcal{G})$ *be a connection form. If* $\omega = \omega_0 + \mu \in \Omega^2_{\mathcal{G}}(\mathcal{N})$ *is a closed* \mathcal{G} *equivariant* 2*-form and we define* $\alpha \in \Omega^1(\mathcal{N})^{\mathcal{G}}$ *by* $\alpha(X) = \mu(A(X))(x)$ *for* $X \in T_x \mathcal{N}$, then we have $ChW_A(\omega) = \omega + d_c \alpha$.

Corollary 1 *The map* ChW: $H^2_{\mathcal{G}}(\mathcal{N}) \to H^2(\mathcal{N}/\mathcal{G})$ *is an isomorphism.*

Let us assume now that $H^1(\mathcal{N}) = H^2(\mathcal{N}) = 0$. The cohomology of the Lie algebra Lie *G* with values in $\Omega^0(\mathcal{N})$ is denoted by $H^{\bullet}(\text{Lie } \mathcal{G}, \Omega^0(\mathcal{N}))$.

Proposition 2 *Let* $\omega = \omega_0 + \mu \in \Omega^2_{\mathcal{G}}(\mathcal{N})$ *be a closed G-equivariant form. If* $\rho \in \Omega$ $\Omega^1(\mathcal{N})$ *satisfies* $\omega_0 = d\rho$ *, then the map* $\tau_\rho \in \Omega^1(\text{Lie } \mathcal{G}, \Omega^0(\mathcal{N}))$ *given by* $\tau_\rho(X) =$ $\rho(X_N) + \mu(X)$ determines a linear map $\tau : H^2_G(\mathcal{N}) \to H^1(\text{Lie } \mathcal{G}, \Omega^0(\mathcal{N}))$ which is *independent of the form* ρ *chosen, and that we call the transgression map* τ *. If* $\mathcal G$ *is connected, then* τ *is injective.*

Proof The first part of the Proposition easily follows using that $L_{Y\wedge K}(X)$ = $\mu([Y, X])$ by the invariance of μ . We restrict ourselves to prove that τ is injective. By definition $[\tau_{\rho}] = 0$ on $H^1(\text{Lie } \mathcal{G}, \Omega^0(\mathcal{N}))$ if and only if there exists $\beta \in \Omega^0(\mathcal{N})$ such that for every $X \in \text{Lie } \mathcal{G}$ we have $\tau_{\rho}(X) = L_{X \wedge \rho} \beta = \iota_{X \wedge \rho} d\beta$. If we set $\rho' = \rho - d\beta$ then for every $X \in \text{Lie } \mathcal{G}$ we have $d\rho' = \omega_0$, $\iota_{X \wedge \rho} \rho' = -\mu(X)$, $L_{X \wedge \rho'} \rho' = 0$, *i.e.*, $\rho' \in \Omega^1(\mathcal{N})^\mathcal{G}$ and $d_c \rho' = \omega$.

Now we assume that $\pi: \mathcal{N} \to \mathcal{N}/\mathcal{G}$ is a principal \mathcal{G} -bundle. Then we can consider the more familiar transgression map defined as follows

Proposition 3 *Let* $\underline{\omega} \in \Omega^2(\mathcal{N}/\mathcal{G})$ *be a closed* 2*-form. If* $\eta \in \Omega^1(\mathcal{N})$ *is a form such that* $\pi^* \omega = d\eta$ *, then the map* t_n : Lie $\mathcal{G} \to \Omega^0(\mathcal{N})$ *,* $t_n(X) = \eta(X_N)$ determines a *linear map t* : $H^2(\mathcal{N}/\mathcal{G}) \to H^1(\text{Lie }\mathcal{G}, \Omega^0(\mathcal{N}))$, which is independent of the form η *chosen, and that we call the transgression map t. If G is connected, then t is injective.*

Proof Again we only prove that *t* is injective. If $t_n = \delta v$ for certain $v \in \Omega^0(\mathcal{N})$, then $\eta(X_N) = L_{X_N} v$. We define $\eta' = \eta - dv$, and we have $d\eta' = \pi^* \omega$, $\iota_{X_N} \eta' = 0$, *L* $\chi_N \eta' = 0$. Hence η' is projectable onto a form $\eta' \in \Omega^1(\mathcal{N}/\mathcal{G})$ and $d\eta' = \omega$.

Proposition 4 *If* $\omega \in H^2_G(\mathcal{N})$ *and* $\underline{\omega} = \text{ChW}(\omega)$ *then we have* $\tau(\omega) = t(\underline{\omega})$ *.*

Proof If $\omega = \omega_0 + \mu$, by Proposition [1](#page-2-0) we have $\omega = \pi^* \omega + d_c \alpha$ for some $\alpha \in$ $\Omega_{\mathcal{G}}^1(\mathcal{N}) = \Omega^1(\mathcal{N})^{\mathcal{G}}$, i.e. $\omega_0 = \pi^* \underline{\omega} + d\alpha$ and $\mu(X) = -\alpha(X_N)$.

Let $\eta \in \Omega^1(\mathcal{N})$ be a form such that $\pi^* \underline{\omega} = d\eta$. If we set $\rho = \eta + \alpha$ then $\omega_0 = d\rho$ and for every $X \in \text{Lie } \mathcal{G}$ we have $\tau_o(X) = t_n(X)$.

3 Local Equivariant Cohomology

Let $p: E \to M$ be a bundle over a compact, oriented *n*-manifold M without boundary. We denote by *J^r E* its *r*-jet bundle, by $J^{\infty}E$ the infinite jet bundle and by $\Gamma(E)$ be the manifold of global sections of *E* (assumed to be not empty).

We denote by Proj E the space of projectable diffeomorphism of E , and by Proj⁺ E the subgroup of elements preserving the orientation of *M*. The space of projectable vector fields on *E* is denoted by proj*E*. We consider the natural actions of Proj*E* on $J^{\infty}E$ and $\Gamma(E)$.

Let j^{∞} : $M \times \Gamma(E) \to J^{\infty}E$, $j^{\infty}(x, s) = j^{\infty}_x s$ be the evaluation map. We define a $\text{map } \mathfrak{S} \colon \Omega^{n+k}(J^{\infty}E) \to \Omega^{k}(\Gamma(E)), \text{by } \mathfrak{I}[\alpha] = \int_{M} (\mathfrak{j}^{\infty})^{*} \alpha, \text{ for } \alpha \in \Omega^{n+k}(J^{\infty}E).$ The map \Im commutes with the exterior differential and is Proj⁺ *E*-equivariant (see [\[9\]](#page-5-5)). We define the space of local *k*-forms on $\Gamma(E)$, as the image of the map \Im , i.e. $\Omega_{\text{loc}}^k(\Gamma(E)) = \Im(\Omega^{n+k}(J^{\infty}E)) \subset \Omega^k(\Gamma(E))$. The cohomology $H_{\text{loc}}^{\bullet}(\Gamma(E))$ of the complex $(\Omega^{\bullet}_{loc}(\Gamma(E)), d)$ is called the local cohomology of $\Gamma(E)$. We have $H_{\text{loc}}^k(\Gamma(E)) \cong H^{n+k}(E)$ for $k > 0$ (see [\[10](#page-5-6)]).

Let G be a Lie group acting on E by elements $Proj^+E$. In order to define an adequate notion of local equivariant cohomology we made the following

Assumption 1 We assume that Lie G is isomorphic to the space of sections of a Lie algebroid $V \to M$, i.e. Lie $\mathcal{G} \cong \Gamma(V)$. We also assume that the map Lie $\mathcal{G} \cong \Gamma(V)$ → projE, $X \mapsto X_E$ is a differential operator. Finally, in the definition of *G*-equivariant cohomology $H_{\mathcal{G}}^{n+k}(J^{\infty}E)$, we assume that the polynomial maps α : Lie $\mathcal{G} \to \Omega^{\bullet}(J^{\infty}E)$ are differential operators.

We extend the integration operator to a map $\mathfrak{S}: \Omega_{\mathcal{G}}^{n+k}(J^{\infty}E) \to \Omega_{\mathcal{G}}^{k}(\Gamma(E))$, by setting $(\Im[\alpha])(X) = \Im[\alpha(X)]$ for every $\alpha \in \Omega_{\mathcal{G}}^{n+k}(J^{\infty}E)$, $X \in \text{Lie }\mathcal{G}$. The map \Im commutes with the Cartan differential and induces a homomorphism in equivariant cohomology $\Im: H_{\mathcal{G}}^{n+k}(J^{\infty}E) \to H_{\mathcal{G}}^{k}(T(E))$ (see [\[9\]](#page-5-5)).

We define the space of local *G*-equivariant *k*-forms by

$$
\Omega_{\mathcal{G},\text{loc}}^k(\Gamma(E)) = \mathfrak{F}(\Omega_{\mathcal{G}}^{n+k}(J^{\infty}E)) \subset \Omega_{\mathcal{G}}^k(\Gamma(E)).
$$
\n(2)

The local *G*-equivariant cohomology of $\Gamma(E)$, $H_{\mathcal{G},loc}^{\bullet}(\Gamma(E))$, is defined as the cohomology of the complex $(\Omega_{\mathcal{G},\text{loc}}^{\bullet}(\Gamma(E)), d_c)$.

4 Application to Local Anomaly Cancellation

Let us define the BRST cohomology (see [\[6](#page-5-1), [16](#page-5-3)]). Recall (see Assumption [1\)](#page-3-0) that we assume Lie $G \cong \Gamma(V)$ for some vector bundle $V \to M$. A map $\alpha: \bigwedge_{k=1}^{k} \text{Lie } \mathcal{G} \to$ $\Omega_{\text{loc}}^0(\Gamma(E))$ is said to be local if there exists a differential operator *A*: \bigwedge^k Lie $\mathcal{G} \to$ $\Omega^n(J^\infty E)$ such that $\alpha(X_1,\ldots,X_k) = \Im[A(X_1,\ldots,X_k)]$ for every $X_1,\ldots,X_k \in$ Lie *G*. We denote by Ω^k_{loc} (Lie *G*, Ω^0_{loc} ($\Gamma(E)$) the space of local *k*-forms on Lie *G* with

values on $\Omega^0_{\text{loc}}(\Gamma(E))$. The differential δ on the complex $\Omega^{\bullet}(\text{Lie }\mathcal{G}, \Omega^0_{\text{loc}}(\Gamma(E)))$ induces a differential on $\Omega^{\bullet}_{loc}(\text{Lie } \mathcal{G}, \Omega^0_{loc}(\Gamma(E)))$. The corresponding cohomology $H_{loc}^{\bullet}(\text{Lie } \mathcal{G}, \Omega_{loc}^{0}(T(E)))$ is called the BRST cohomology. We assume that $H^2(\Gamma(E)) = H^1(\Gamma(E)) = 0$ and also that $H^2_{loc}(\Gamma(E)) = H^1_{loc}(\Gamma(E)) = 0$.

Proposition 5 *The restriction of the transgression map* τ *to* $H^2_{G,\text{loc}}(\Gamma(E))$ *takes values* on the BRST cohomology $I_{\text{loc}}^1(\text{Lie } \mathcal{G}, \Omega_{\text{loc}}^0(\Gamma(E)))$ *. The map* $\tau: H^2_{\mathcal{G},loc}(\Gamma(E)) \to H^1_{loc}(\text{Lie } \mathcal{G}, \Omega^0_{loc}(\Gamma(E)))$ is injective for $\mathcal G$ connected.

Proof Let $\omega = \omega_0 + \mu \in \Omega^2_{\mathcal{G},loc}(\Gamma(E))$ be a closed local *G*-equivariant 2-form. As $H^2_{loc}(\Gamma(E)) = 0$, we have $\omega_0 = d\rho$, for certain $\rho \in \Omega^1_{loc}(\Gamma(E))$. By our assumption in the definition of local equivariant cohomology (Assumption [1\)](#page-3-0), the map τ_{ρ} : Lie $\mathcal{G} \to \Omega_{\text{loc}}^0(\Gamma(E)), \tau_{\rho}(X) = \rho(X_{\Gamma(E)}) + \mu(X) \in \Omega_{\text{loc}}^0(\Gamma(E))$ is a local map. The injectiveness of τ follows from Proposition [2.](#page-2-1)

Assume that $\Gamma(E) \to \Gamma(E)/\mathcal{G}$ is a principal \mathcal{G} -bundle. Then we define the local cohomology by $H^k_{loc}(\Gamma(E)/\mathcal{G}) = \text{ChW}(H^k_{\mathcal{G},loc}(\Gamma(E)))$. By Proposition [4](#page-2-2) we have the following

Proposition 6 *Let* $\omega \in \Omega^2_{\mathcal{G},loc}(\Gamma(E))$ *be a closed local G-equivariant* 2*-form and let* ω = ChW(ω). Then we have $\tau(\omega) = t(\omega)$, and in particular we conclude that $t(\underline{\omega}) \in H^1_{loc}(\text{Lie } \mathcal{G}, \Omega^0_{loc}(\Gamma(E)))$ *. Moreover, the following conditions are equivalent*

(a) $[\omega] = 0$ *on H*²_{*G*,loc}($\Gamma(E)$). (b) $[\underline{\omega}] = 0$ *on* $H_{\text{loc}}^2(\Gamma(E)/\mathcal{G})$. (c) $[\tau(\omega)] = [t(\omega)] = 0$ *on* $H_{\text{loc}}^1(\text{Lie } \mathcal{G}, \Omega_{\text{loc}}^0(\Gamma(E))).$

Proposition [6](#page-4-5) applied to $\omega = c_1(\det \mathrm{Ind} D)$ shows that our definition of $H^2_{\text{loc}}(\Gamma(E)/\mathcal{G})$ solves Singer's problem. We also note that if $\omega \in \Omega^2_{\mathcal{G},\text{loc}}(\Gamma(E))$ is closed, the form $\omega \in \Omega^2(\Gamma(E)/\mathcal{G})$ determining the class ChW([ω]) could contain non-local terms, as ω depends on the curvature of a connection Θ on the principal *G*-bundle $\Gamma(E) \to \Gamma(E)/\mathcal{G}$, and Θ usually contains non-local terms. This fact explains the appearance of non-local terms on the expression of the curvature of det Ind*D* commented on the Introduction.

References

- 1. Álvarez, O., Singer, I., Zumino, B.: Gravitational anomalies and the family's index theorem. Commun. Math. Phys. **96**, 409–417 (1984)
- 2. Atiyah, M.F., Singer, I.: Dirac operators coupled to vector potentials. Proc. Natl. Acad. Sci. USA **81**, 2597–2600 (1984)
- 3. Berline, N., Getzler, E., Vergne, M.: Heat Kernels and Dirac Operators, Springer, Berlin (1992)
- 4. Bertlmann, R.A.: Anomalies in Quantum Field Theory. Oxford University Press, Oxford (2000)
- 5. Blau, M.: Wess-Zumino terms and the geometry of the determinant line bundle. Phys. Lett. **209** B, 503–506 (1988)
- 6. Bonora, L., Cotta-Ramusino, P.: Some remarks on BRS transformations, anomalies and the cohomology of the Lie algebra of the group of gauge transformations. Commun. Math. Phys. **87**, 589–603 (1983)
- 7. Bonora, L., Cotta-Ramusino, P.: Consistent and covariant anomalies and local cohomology. Phys. Rev. D **33**(3), 3055–3059 (1986)
- 8. Dubois-Violette, M., Henneaux, M., Talon, M., Viallet, C.: General solution of the consistency equation. Phys. Lett. B **289**, 361–367 (1992)
- 9. Pérez, R.F.: Equivariant characteristic forms in the bundle of connections. J. Geom. Phys. **54**, 197–212 (2005)
- 10. Pérez, R.F.: Local cohomology and the variational bicomplex. Int. J. Geom. Methods Mod. Phys. **5** (2008), 587–604
- 11. Pérez, R.F.: Local anomalies and local equivariant cohomology. Comm. Math. Phys. **286**, 445–458 (2009)
- 12. Pérez, R.F., Masqué, J,M.: Natural connections on the bundle of Riemannian metrics. Monatsh. Math. **155**, 67–78 (2008)
- 13. Guillemin, V., Sternberg, S.: Supersymmetry and Equivariant de Rham Theory. Springer, Berlin (1999)
- 14. Mañes, J., Stora, R., Zumino, B.: Algebraic study of chiral anomalies. Comm. Math. Phys. **102**, 157–174 (1985)
- 15. Martellini, M., Reina, C.: Some remarks on the index theorem approach to anomalies. Ann. Inst. H. Poincarè **113**, 443–458 (1985)
- 16. Schmid, R.: Local cohomology in gauge theories, BRST transformations and anomalies. Differential Geom. Appl. **4**(2), 107–116 (1994)
- 17. Singer, I.M.: Families of Dirac operators with applications to physics, The mathematical heritage of Élie Cartan (Lyon, 1984). Astérisque (1985), Numero Hors Serie, 323–340