

# Chapter 10

## The Characteristic Matrix of Nonuniqueness for First-Kind Equations

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### 10.1 Introduction

Let  $S$  be a finite domain in  $\mathbb{R}^2$ , bounded by a simple, closed,  $C^2$  curve  $\partial S$ . We denote by  $x$  and  $y$  generic points in  $S \cup \partial S$  and by  $|x - y|$  the distance between  $x$  and  $y$  in the Cartesian metric. Also, let  $C^{0,\alpha}(\partial S)$  and  $C^{1,\alpha}(\partial S)$ ,  $\alpha \in (0, 1)$ , be, respectively, the spaces of Hölder continuous and Hölder continuously differentiable functions on  $\partial S$ . In what follows, Greek and Latin indices take the values 1, 2 and 1, 2, 3, respectively, and a superscript T denotes matrix transposition.

For any function  $f$  continuous on  $\partial S$ , we define the ‘calibration’ functional  $p$  by

$$pf = \int_{\partial S} f ds.$$

Using the fundamental solution for the two-dimensional Laplacian

$$g(x, y) = -\frac{1}{2\pi} \ln|x - y|,$$

we define the single-layer harmonic potential of density  $\varphi$  by

$$(V\varphi)(x) = \int_{\partial S} g(x, y)\varphi(y) ds(y).$$

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The proof of the following assertion can be found, for example, in [Co94] or [Co00].

**Theorem 1.** *For any  $\alpha \in (0, 1)$ , there are a unique nonzero function  $\Phi \in C^{0,\alpha}(\partial S)$  and a unique number  $\omega$  such that*

$$V\Phi = \omega \quad \text{on } \partial S, \quad p\Phi = 1.$$

It is easy to see that  $\Phi$  and  $\omega$  depend on  $g$  and  $\partial S$ .

The numbers  $2\pi\omega$  and  $e^{-2\pi\omega}$  are called *Robin's constant* and the *logarithmic capacity* of  $\partial S$ .

For a circle with the center at the origin and radius  $R$ , both  $\Phi$  and  $\omega$  can be determined explicitly:

$$\Phi = \frac{1}{2\pi R}, \quad \omega = -\frac{1}{2\pi} \ln R.$$

For other boundary curves,  $\Phi$  and  $\omega$  are practically impossible to determine analytically and must be computed by numerical methods.

If the solution of the Dirichlet problem in  $S$  with data function  $f$  on  $\partial S$  is sought as  $u = V\varphi$ , then  $\varphi$  is a solution of the (weakly singular) first-kind boundary integral equation

$$V\varphi = f \quad \text{on } \partial S.$$

This is a well-posed problem if and only if  $\omega \neq 0$ . If  $\omega = 0$ , the above equation has infinitely many solutions, which are expressed in terms of  $\Phi$ .

## 10.2 Plane Elastic Strain

Consider a plate made of a homogeneous and isotropic material with Lamé constants  $\lambda$  and  $\mu$ , which undergoes deformations in the  $(x_1, x_2)$ -plane. If the body forces are negligible, then its (static) displacement vector  $u = (u_1, u_2)^T$  satisfies the equilibrium system of equations [Co00]

$$Au = 0 \quad \text{in } S,$$

where

$$A(\partial_1, \partial_2) = \begin{pmatrix} \mu\Delta + (\lambda + \mu)\partial_1^2 & (\lambda + \mu)\partial_1\partial_2 \\ (\lambda + \mu)\partial_1\partial_2 & \mu\Delta + (\lambda + \mu)\partial_2^2 \end{pmatrix}.$$

It is not difficult to show [Co00] that the columns  $F^{(i)}$  of the matrix

$$F = \begin{pmatrix} 1 & 0 & x_2 \\ 0 & 1 & -x_1 \end{pmatrix}$$

form a basis for the space of rigid displacements.

The ‘calibrating’ vector-valued functional  $p$  is defined for continuous  $2 \times 1$  vector functions  $f$  by

$$pf = \int_{\partial S} F^T f ds.$$

A matrix of fundamental solutions for  $A$  is [Co00]

$$D(x, y) = -\frac{1}{4\pi\mu(\gamma+1)} \\ \times \begin{pmatrix} 2\gamma \ln|x-y| + 2\gamma + 1 - \frac{2(x_1-y_1)^2}{|x-y|^2} & -\frac{2(x_1-y_1)(x_2-y_2)}{|x-y|^2} \\ -\frac{2(x_1-y_1)(x_2-y_2)}{|x-y|^2} & 2\gamma \ln|x-y| + 2\gamma + 1 - \frac{2(x_2-y_2)^2}{|x-y|^2} \end{pmatrix}, \\ \gamma = \frac{\lambda + 3\mu}{\lambda + \mu}.$$

The single-layer potential of density  $\varphi$  is defined by

$$(V\varphi)(x) = \int_{\partial S} D(x, y)\varphi(y) ds(y).$$

The proof of the following assertion can be found in [Co00].

**Theorem 2.** *There is a unique  $2 \times 3$  matrix function  $\Phi \in C^{0,\alpha}(\partial S)$  and a unique  $3 \times 3$  constant symmetric matrix  $\mathcal{C}$  such that the columns  $\Phi^{(i)}$  of  $\Phi$  are linearly independent and*

$$V\Phi = F\mathcal{C} \quad \text{on } \partial S, \quad p\Phi = I,$$

where  $I$  is the identity matrix.

Clearly,  $\Phi$  and  $\mathcal{C}$  depend on  $A$ ,  $D$ , and  $\partial S$ .

In the so-called alternative indirect method [Co00], the solution of the Dirichlet problem in  $S$  with data function  $f$  on  $\partial S$  is sought in the form

$$u = V\varphi. \tag{10.1}$$

Then the problem reduces to the (weakly singular) boundary integral equation

$$V\varphi = f \quad \text{on } \partial S. \quad (10.2)$$

**Theorem 3.** Equation (10.2) has a unique solution  $\varphi \in C^{0,\alpha}(\partial S)$ ,  $\alpha \in (0, 1)$ , for any  $f \in C^{1,\alpha}(\partial S)$  if and only if  $\det \mathcal{C} \neq 0$ . In this case, (10.1) is the unique solution of the Dirichlet problem.

If  $\det \mathcal{C} = 0$ , then the unique solution of the Dirichlet problem is obtained by solving an ill-posed modified boundary integral equation whose infinitely many solutions are constructed with  $\Phi$  and  $\mathcal{C}$ .

In the so-called refined indirect method [Co00], the solution of the Dirichlet problem is sought as a pair  $\{\varphi, c\}$  such that

$$u = V\varphi - Fc \quad \text{in } S, \quad p\varphi = s,$$

where  $s$  a constant  $3 \times 1$  vector chosen (arbitrarily) a priori and  $c$  is a constant  $3 \times 1$  vector. This leads to the system of boundary integral equations

$$V\varphi - Fc = f \quad \text{on } \partial S, \quad p\varphi = s. \quad (10.3)$$

**Theorem 4.** System (10.3) has a unique solution  $\{\varphi, c\}$  with  $\varphi \in C^{0,\alpha}(\partial S)$  for any  $f \in C^{1,\alpha}(\partial S)$ ,  $\alpha \in (0, 1)$ , and any  $s$ .

It is important to evaluate the arbitrariness in the representation of the solution with respect to the prescribed ‘calibration’  $s$ .

**Theorem 5.** If  $\{\varphi^{(1)}, c^{(1)}\}$ ,  $\{\varphi^{(2)}, c^{(2)}\}$  are two solutions of (10.3) constructed with  $s^{(1)}$  and  $s^{(2)}$ , respectively, then

$$\begin{aligned} \varphi^{(2)} &= \varphi^{(1)} + \Phi(s^{(2)} - s^{(1)}), \\ c^{(2)} &= c^{(1)} + \mathcal{C}(s^{(2)} - s^{(1)}). \end{aligned}$$

It is not easy to compute  $\Phi$  and  $\mathcal{C}$  analytically, or even numerically, in arbitrary domains  $S$ , but this can be accomplished if  $S$  is a circular disk. Let  $\partial S$  be the circle with center at the origin and radius  $R$ . In this case,  $\Phi$  and  $\mathcal{C}$  can be determined analytically as

$$\Phi = \frac{1}{2\pi R} \begin{pmatrix} 1 & 0 & R^{-2}x_2 \\ 0 & 1 & -R^{-2}x_1 \end{pmatrix},$$

$$\mathcal{C} = -\frac{1}{4\pi\mu(\lambda + 2\mu)R^2} \times \begin{pmatrix} (\lambda + 3\mu)R^2(\ln R + 1) & 0 & 0 \\ 0 & (\lambda + 3\mu)R^2(\ln R + 1) & 0 \\ 0 & 0 & -(\lambda + \mu) \end{pmatrix}.$$

Clearly,  $\det \mathcal{C} = 0$  if and only if  $R = e^{-1}$ .

Analytic computation of  $\Phi$  and  $\mathcal{C}$  is practically impossible for non-circular domains, and must be performed numerically.

We choose four  $3 \times 1$  constant vectors  $s^{(0)}, s^{(i)}$  such that the set  $\{s^{(i)} - s^{(0)}\}$  is linearly independent, and form the  $3 \times 3$  matrix  $\Sigma$  with columns  $s^{(i)} - s^{(0)}$ . Also, we choose an arbitrary function  $f$ . Next, we compute the solutions  $\{\varphi^{(0)}, c^{(0)}\}, \{\varphi^{(i)}, c^{(i)}\}$  of (10.3) corresponding to  $s^{(0)}, s^{(i)}$ , respectively, and  $f$ , by the refined indirect method, then form the  $2 \times 3$  matrix function  $\Psi$  with columns  $\varphi^{(i)} - \varphi^{(0)}$  and the constant  $3 \times 3$  matrix  $\Gamma$  with columns  $c^{(i)} - c^{(0)}$ .

From Theorem 4 it follows that

$$\begin{aligned} \varphi^{(i)} - \varphi^{(0)} &= \Phi(s^{(i)} - s^{(0)}), \\ c^{(i)} - c^{(0)} &= \mathcal{C}(s^{(i)} - s^{(0)}), \end{aligned}$$

or, what is the same,

$$\Phi \Sigma = \Psi, \quad \mathcal{C} \Sigma = \Gamma;$$

hence,

$$\Phi = \Psi \Sigma^{-1}, \quad \mathcal{C} = \Gamma \Sigma^{-1}.$$

A similar analysis can be performed for other two-dimensional linear elliptic systems with constant coefficients—for example, the system modeling bending of elastic plates with transverse shear deformation [Co00]. No apparent connection exists between the matrix  $\mathcal{C}$  and the characteristic constant  $\omega$  of  $\partial S$ .

### 10.3 Numerical Examples

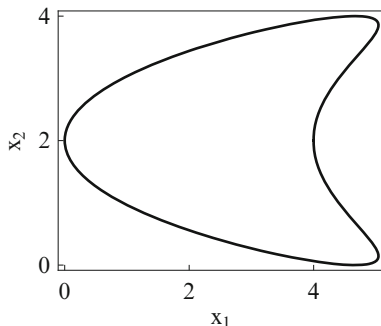
Consider a steel plate with scaled Lamé coefficients

$$\lambda = 11.5, \quad \mu = 7.69,$$

and let  $\partial S$  (see Figure 10.1) be the curve of parametric equations

$$\begin{aligned} x_1(t) &= 2 \cos(\pi t) - \frac{4}{3} \cos(2\pi t) + \frac{10}{3}, \\ x_2(t) &= 2 \sin(\pi t) + 2, \quad 0 \leq t \leq 2. \end{aligned}$$

**Fig. 10.1** The boundary curve  $\partial S$ .



We choose the vectors

$$s^{(0)} = \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix}, \quad s^{(1)} = \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}, \quad s^{(2)} = \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix}, \quad s^{(3)} = \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix},$$

$$f(x) = \begin{pmatrix} 1 \\ 0 \end{pmatrix}.$$

The approximating functions for computing  $\varphi^{(0)}(t)$  and  $\varphi^{(i)}(t)$  are piecewise cubic Hermite splines on 12 knots; that is, the interval  $0 \leq t \leq 2$  is divided into 12 equal subintervals. Then the characteristic matrix (with entries rounded off to 5 decimal places) is

$$\mathcal{C} = \begin{pmatrix} -0.01627 & -0.01083 & -0.00370 \\ -0.01083 & -0.00892 & 0.00542 \\ -0.00370 & 0.00542 & 0.00185 \end{pmatrix}.$$

Here,

$$\det \mathcal{C} = 1.08273 \times 10^{-6}.$$

The graphs of the components  $\Phi_{\alpha i}$  of  $\Phi$  are shown in Figure 10.2.

As a second example, consider the ‘expanding’ ellipse  $\partial S$  of parametric equations

$$x_1(t) = 2k \cos(\pi t),$$

$$x_2(t) = k \sin(\pi t), \quad 0 \leq t \leq 2.$$

The graph of  $\det \mathcal{C}$  as a function of  $k$  is shown in Figure 10.3.

Here,  $\det \mathcal{C} = 0$  for  $k = 0.22546$  and  $k = 0.26934$ .

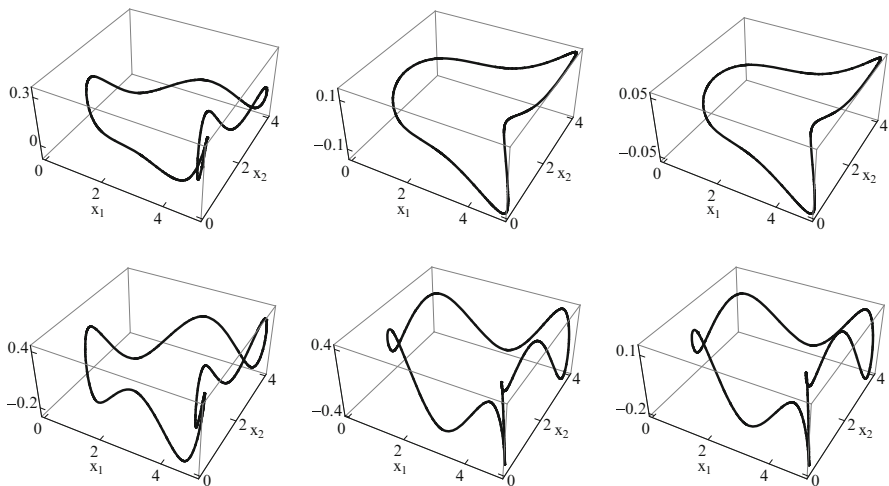


Fig. 10.2 Graphs of the  $\Phi_{\alpha_i}$ .

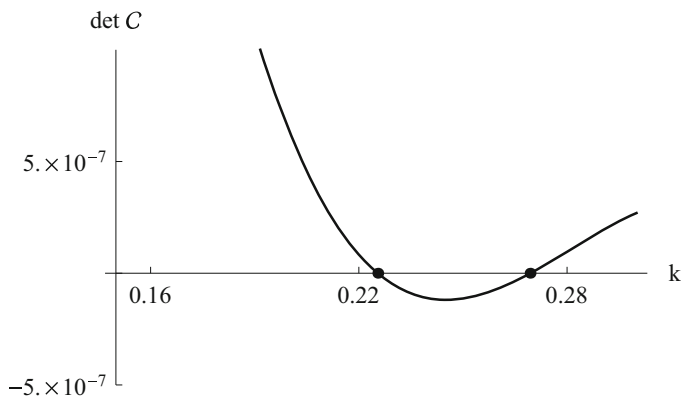


Fig. 10.3 Graph of  $\det \mathcal{C}$ .

### References

[Co94] Constanda, C.: On integral solutions of the equations of thin plates. *Proc. Roy. Soc. London Ser. A*, **444**, 261–268 (1994).  
 [Co00] Constanda, C.: *Direct and Indirect Boundary Integral Equation Methods*, Chapman & Hall/CRC, Boca Raton, FL (2000).