

# Non-equilibrium spin-spin interactions in strongly correlated systems

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**Abstract** We develop a theory for magnetism of strongly correlated systems driven out of equilibrium by an external time-dependent electric field. We provide expressions for computing the effective interaction parameters between electronic spins, including a new interaction that we name *twist exchange*. Our theory is suitable for laser-induced ultrafast magnetization dynamics.

## Introduction and Method

The theoretical study of magnetization in realistic systems is a challenging problem. In equilibrium, magnetic interactions in magnetic metals and semiconductors are non-Heisenberg [1-4], i.e., the lengths of magnetic moments and values of exchange parameters depend on the magnetic configuration for which they are computed, which calls for an accurate *ab initio* formulation. The expressions for the exchange-interaction parameters were given years ago [5, 6], and they can be used within a classical Heisenberg model to simulate spin dynamics. However, this treatment is not expected to be appropriate for ultrafast dynamics [7], since the time scale of the excitations is faster than the typical scale of exchange interactions (10÷100 fs). Exchange parameters are thus time-dependent. Our purpose is to derive their expressions in a general non-equilibrium framework, including external time-dependent fields [8]. The resulting formulas are valid for any time scale, thus they are applicable to ultrafast spin dynamics.

We consider a system described by the time-dependent multi-band Hubbard Hamiltonian,

$$\hat{H}(t) \equiv \sum_a \sum_b \sum_{\sigma=\uparrow,\downarrow} \hat{\phi}_{a\sigma}^+ T_{ab}(t) \hat{\phi}_{b\sigma} + \hat{H}_V, \quad 1$$

where italic letters  $a, b$  denote sets of site and orbital indexes;  $\hat{H}_V$  accounts for on-site interactions between the electrons. The matrix element  $T_{ab}(t)$  accounts for in-

ter-site hopping and orbital hybridization due to both the equilibrium structure and an external time-dependent electric field, switched on at time  $t_0$ . Here we do not consider relativistic effects nor external magnetic fields, so  $T_{ab}(t)$  is spin-independent.

We write the partition function for the electronic system in the Kadanoff-Baym formalism and we derive an effective action that describes the system in equilibrium for  $t \leq t_0$ , as well as out of equilibrium for  $t > t_0$ . We assume spontaneous symmetry breaking (SSB) along the direction  $z$ , specified by the unit vector  $\vec{u}_z$ , and we study the spin excitations on the top of the collinear equilibrium ground state. Thus, we apply time- and site- dependent Holstein-Primakoff spin rotations to map the old fermionic fields  $\phi_{a\sigma}$  into new fields  $\psi_{a\sigma}$  having their spins aligned with  $t$ -dependent vectors  $\vec{e}_a(t)$ , which are the unit vectors of the (classical) magnetic moments. The deviations of the  $\vec{e}_a(t)$ 's from  $\vec{u}_z$  are described by auxiliary bosonic fields; in the low-energy sector we assume such deviations to be small, obtaining an action quadratic in the bosons. Integrating out the fermionic fields  $\psi_{a\sigma}$  gives an effective bosonic action, and we finally map the bosons to the  $\vec{e}_a(t)$ 's to obtain the spin-spin interactions.

## Results

Our results are presented in Ref.[8]. Here we summarize some of the main points.

We find that there are two forms of quadratic interactions between magnetic moments. The first form,  $\propto \vec{e}_a \cdot \vec{e}_b$ , is exchange. The second form,  $\propto (\vec{e}_a \times \vec{e}_b) \cdot \vec{u}_z$ , looks similar to an effective Dzyaloshinskii-Moriya interaction (DMI), but this interpretation is not correct since relativistic effects are absent in our model. We note, instead, that an *effective three-body* interaction of the form  $(\vec{e}_a \times \vec{e}_b) \cdot \vec{e}_c$  generates precisely this term when up to second-order deviations of the magnetic moments from  $\vec{u}_z$  are considered. We have proposed the name *twist exchange* for this new interaction [9]. Our method provides both the equilibrium and the non-equilibrium formulas for the coefficients describing exchange and twist-exchange, which are obtained after the fermionic field integration as convolutions of electronic Green functions and self-energies. These can be computed by the means of non-equilibrium dynamical mean-field theory. The formulas for exchange parameters reproduce, in the particular case of the single-band Hubbard model in equilibrium, the results of Ref.[6].

We find that the non-equilibrium spin-spin interactions are not, in general, local in time, i.e., the system has *memory*: the parameters depend on two times. We here present the formula for non-equilibrium exchange parameters in the Hartree-Fock

approximation (Ref.[8] contains the more general formulas, also for twist exchange):

$$\begin{aligned}
J_{ab}(t, t') = & \frac{1}{4} \delta(t - t') \operatorname{Re} \left\{ \int_{-\infty}^{+\infty} \frac{d\omega}{2\pi} f(\omega, t) \left[ \bar{\Sigma}_{ab}^{\downarrow}(t) A_{ba}^{\uparrow}(\omega, t) + A_{ab}^{\downarrow}(\omega, t) \bar{\Sigma}_{ba}^{\uparrow}(t) \right] \right\} \\
& + \frac{\operatorname{sign}(t_r)}{8} \operatorname{Im} \left\{ \int_{-\infty}^{+\infty} \frac{d\omega}{2\pi} \int_{-\infty}^{+\infty} \frac{d\omega'}{2\pi} e^{-i(\omega - \omega')t_r} [f(\omega, t_c) - f(\omega', t_c)] \right. \\
& \cdot \left\{ \left[ \bar{\Sigma}^{\downarrow}(t) \cdot A^{\downarrow}(\omega, t_c) \right]_{ab} \left[ \bar{\Sigma}^{\uparrow}(t') \cdot A^{\uparrow}(\omega', t_c) \right]_{ba} - A_{ab}^{\downarrow}(\omega, t_c) \left[ \bar{\Sigma}^{\uparrow}(t') \cdot A^{\uparrow}(\omega', t_c) \cdot \bar{\Sigma}^{\uparrow}(t) \right]_{ba} \right. \\
& \left. \left. + \left[ A^{\downarrow}(\omega, t_c) \cdot \bar{\Sigma}^{\downarrow}(t') \right]_{ab} \left[ A^{\uparrow}(\omega', t_c) \cdot \bar{\Sigma}^{\uparrow}(t) \right]_{ba} - \left[ \bar{\Sigma}^{\downarrow}(t) \cdot A^{\downarrow}(\omega, t_c) \cdot \bar{\Sigma}^{\downarrow}(t') \right]_{ab} A_{ba}^{\uparrow}(\omega', t_c) \right\} \right\}.
\end{aligned} \tag{2}$$

Here  $t_r = t - t'$ ,  $t_c = (t + t')/2$ , and  $\bar{\Sigma}^{\sigma}(t)$  is the Hartree-Fock self-energy (because of SSB,  $\bar{\Sigma}^{\downarrow} \neq \bar{\Sigma}^{\uparrow}$ ), while  $f(\omega, t)$  and  $A^{\sigma}(\omega, t)$  are, respectively, the time-dependent occupation number and spectral function, related to the non-equilibrium Green functions via the relation  $G_{ab}^{\lessdot}(\omega, t) \equiv i f(\omega, t) A_{ab}(\omega, t)$ . As an application, in Ref.[8] we also derive the time-dependent spin stiffness tensor for a ferromagnet, which may be relevant for the interpretation of experiments where spin dynamics is induced by a laser beam of macroscopic size [7].

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