

Chapter 13

Topological Properties of Strong Solutions for the 3D Navier-Stokes Equations

Pavlo O. Kasyanov, Luisa Toscano and Nina V. Zadoianchuk

Abstract In this chapter we give a criterion for the existence of global strong solutions for the 3D Navier-Stokes system for any regular initial data.

13.1 Introduction

Let $\Omega \subseteq \mathbb{R}^3$ be a bounded open set with sufficiently smooth boundary $\partial\Omega$ and $0 < T < +\infty$. We consider the incompressible Navier-Stokes equations

$$\begin{cases} y_t + (y \cdot \nabla)y = \nu \Delta y - \nabla p + f \text{ in } Q = \Omega \times (0, T), \\ \operatorname{div} y = 0 \text{ in } Q, \\ y = 0 \text{ on } \partial\Omega \times (0, T), \quad y(x, 0) = y_0(x) \text{ in } \Omega, \end{cases} \quad (13.1)$$

where $\nu > 0$ is a constant. We define the usual function spaces

$$\mathcal{V} = \{u \in (C_0^\infty(\Omega))^3 : \operatorname{div} u = 0\},$$

$$H = \text{closure of } \mathcal{V} \text{ in } (L^2(\Omega))^3, \quad V = \{u \in (H_0^1(\Omega))^3 : \operatorname{div} u = 0\}.$$

P. O. Kasyanov (✉)

Institute for Applied System Analysis, National Technical University of Ukraine “Kyiv Polytechnic Institute”, Peremogy ave., 37, build, 35, Kyiv 03056, Ukraine
e-mail: kasyanov@i.ua

L. Toscano

Department of Mathematics and Applications R. Caccioppoli, University of Naples “Federico II”, via Claudio 21, 80125 Naples, Italy
e-mail: luisatoscano@libero.it

N. V. Zadoianchuk

Department of Computational Mathematics, Taras Shevchenko National University of Kyiv, Volodimirska Street 64, Kyiv 03601, Ukraine
e-mail: ninelllll@i.ua

We denote by V^* the dual space of V . The spaces H and V are separable Hilbert spaces and $V \subset H \subset V^*$ with dense and compact embedding when H is identified with its dual H^* . Let (\cdot, \cdot) , $\|\cdot\|_H$ and $((\cdot, \cdot))$, $\|\cdot\|_V$ be the inner product and the norm in H and V , respectively, and let $\langle \cdot, \cdot \rangle$ be the pairing between V and V^* . For $u, v, w \in V$, the equality

$$b(u, v, w) = \int_{\Omega} \sum_{i,j=1}^3 u_i \frac{\partial v_j}{\partial x_i} w_j dx$$

defines a trilinear continuous form on V with $b(u, v, v) = 0$ when $u \in V$ and $v \in (H_0^1(\Omega))^3$. For $u, v \in V$, let $B(u, v)$ be the element of V^* defined by $\langle B(u, v), w \rangle = b(u, v, w)$ for all $w \in V$.

We say that the function y is a *weak solution* of Problem (13.1) on $[0, T]$, if $y \in L^\infty(0, T; H) \cap L^2(0, T; V)$, $\frac{dy}{dt} \in L^1(0, T; V^*)$, if

$$\frac{d}{dt}(y, v) + \nu((y, v)) + b(y, y, v) = \langle f, v \rangle \quad \text{for all } v \in V, \tag{13.2}$$

in the sense of distributions on $(0, T)$, and if y satisfies the energy inequality

$$V(y)(t) \leq V(y)(s) \quad \text{for all } t \in [s, T], \tag{13.3}$$

for a.e. $s \in (0, T)$ and for $s = 0$, where

$$V(y)(t) := \frac{1}{2} \|y(t)\|_H^2 + \nu \int_0^t \|y(\tau)\|_V^2 d\tau - \int_0^t \langle f(\tau), y(\tau) \rangle d\tau. \tag{13.4}$$

This class of solutions is called Leray–Hopf or physical one. If $f \in L^2(0, T; V^*)$, and if y satisfies (13.2), then $y \in C([0, T]; H_w)$, $\frac{dy}{dt} \in L^{\frac{4}{3}}(0, T; V^*)$, where H_w denotes the space H endowed with the weak topology. In particular, the initial condition $y(0) = y_0$ makes sense for any $y_0 \in H$.

Let $A : V \rightarrow V^*$ be the linear operator associated to the bilinear form $((u, v)) = \langle Au, v \rangle$. Then A is an isomorphism from $D(A)$ onto H with $D(A) = (H^2(\Omega))^3 \cap V$. We recall that the embedding $D(A) \subset V$ is dense and continuous. Moreover, we assume $\|Au\|_H$ as the norm on $D(A)$, which is equivalent to the one induced by $(H^2(\Omega))^3$. The Problem (13.1) can be rewritten as

$$\begin{cases} \frac{dy}{dt} + \nu Ay + B(y, y) = f \text{ in } V^*, \\ y(0) = y_0, \end{cases} \tag{13.5}$$

where the first equation we understand in the sense of distributions on $(0, T)$. Now we write

$\mathcal{D}(y_0, f) = \{ y : y \text{ is a weak solution of Problem (13.1) on } [0, T] \}$.

It is well known (cf. [1]) that if $f \in L^2(0, T; V^*)$, and if $y_0 \in H$, then $\mathcal{D}(y_0, f)$ is not empty.

A weak solution y of Problem (13.1) on $[0, T]$ is called a *strong* one, if it additionally belongs to Serrin’s class $L^8(0, T; (L^4(\Omega))^3)$. We note that any strong solution y of Problem (13.1) on $[0, T]$ belongs to $C([0, T]; V) \cap L^2(0, T; D(A))$ and $\frac{dy}{dt} \in L^2(0, T; H)$ (cf. [2, Theorem 1.8.1, p. 296] and references therein).

For any $f \in L^\infty(0, T; H)$ and $y_0 \in V$ it is well known the only local existence of strong solutions for the 3D Navier-Stokes equations (cf. [1–4] and references therein). Here we provide a criterion for existence of strong solutions for Problem (13.1) on $[0, T]$ for any initial data $y_0 \in V$ and $0 < T < +\infty$. Presented results were announced in [5].

13.2 Topological Properties of Strong Solutions

The main result of this note has the following form.

Theorem 13.1 *Let $f \in L^2(0, T; H)$ and $y_0 \in V$. Then either for any $\lambda \in [0, 1]$ there is an $y_\lambda \in C([0, T]; V) \cap L^2(0, T; D(A))$ such that $y_\lambda \in \mathcal{D}(\lambda y_0, \lambda f)$, or the set*

$$\{ y \in C([0, T]; V) \cap L^2(0, T; D(A)) : y \in \mathcal{D}(\lambda y_0, \lambda f), \lambda \in (0, 1) \} \tag{13.6}$$

is unbounded in $L^8(0, T; (L^4(\Omega))^3)$.

In the proof of Theorem 13.1 we use an auxiliary statement connected with continuity property of strong solutions on parameters of Problem (13.1) in Serrin’s class $L^8(0, T; (L^4(\Omega))^3)$.

Theorem 13.2 *Let $f \in L^2(0, T; H)$ and $y_0 \in V$. If y is a strong solution for Problem (13.1) on $[0, T]$, then there exist $L, \delta > 0$ such that for any $z_0 \in V$ and $g \in L^2(0, T; H)$, satisfying the inequality*

$$\|z_0 - y_0\|_V^2 + \|g - f\|_{L^2(0, T; H)}^2 < \delta, \tag{13.7}$$

the set $\mathcal{D}(z_0, g)$ is one-point set $\{z\}$ which belongs to $C([0, T]; V) \cap L^2(0, T; D(A))$, and

$$\|z - y\|_{C([0, T]; V)}^2 + \frac{\nu}{4} \|z - y\|_{D(A)}^2 \leq L \left(\|z_0 - y_0\|_V^2 + \|g - f\|_{L^2(0, T; H)}^2 \right). \tag{13.8}$$

Remark 13.1 We note that from Theorem 13.2 with $z_0 \in V$ and $g \in L^2(0, T; H)$ with $\|z_0\|_V^2 + \|g\|_{L^2(0, T; H)}^2$ sufficiently small, Problem (13.1) has only one global strong solution.

Remark 13.2 Theorem 13.2 provides that, if for any $\lambda \in [0, 1]$ there is an $y_\lambda \in L^8(0, T; (L^4(\Omega))^3)$ such that $y_\lambda \in \mathcal{D}(\lambda y_0, \lambda f)$, then the set

$$\{y \in C([0, T]; V) \cap L^2(0, T; D(A)) : y \in \mathcal{D}(\lambda y_0, \lambda f), \lambda \in (0, 1)\}$$

is bounded in $L^8(0, T; (L^4(\Omega))^3)$.

If Ω is a C^∞ -domain and if $f \in C_0^\infty(\overline{(0, T) \times \Omega})^3$, then any strong solution y of Problem (13.1) on $[0, T]$ belongs to $C^\infty((0, T) \times \Omega)^3$ and $p \in C^\infty((0, T) \times \Omega)$ (cf. [2, Theorem 1.8.2, p. 300] and references therein). This fact directly provides the next corollary of Theorems 13.1 and 13.2.

Corollary 13.1 *Let Ω be a C^∞ -domain, $f \in C_0^\infty(\overline{(0, T) \times \Omega})^3$. Then either for any $y_0 \in V$ there is a strong solution of Problem (13.1) on $[0, T]$, or the set*

$$\{y \in C^\infty((0, T) \times \Omega)^3 : y \in \mathcal{D}(\lambda y_0, \lambda f), \lambda \in (0, 1)\}$$

is unbounded in $L^8(0, T; (L^4(\Omega))^3)$ for some $y_0 \in C_0^\infty(\Omega)^3$.

13.3 Proof of Theorem 13.2

Let $f \in L^2(0, T; H)$, $y_0 \in V$, and $y \in C([0, T]; V) \cap L^2(0, T; D(A))$ be a strong solution of Problem (13.1) on $[0, T]$. Due to [6], [1, Chap. 3] the set $\mathcal{D}(y_0, f) = \{y\}$. Let us now fix $z_0 \in V$ and $g \in L^2(0, T; H)$ satisfying (13.7) with

$$\delta = \min \left\{ 1; \frac{\nu}{4} \right\} e^{-2TC}, \quad C = \max \left\{ \frac{27c^4}{2\nu^3}; \frac{7^7 c^8}{2^9 \nu^7} \right\} \left(\|y\|_{C([0, T]; V)}^4 + 1 \right)^2, \quad (13.9)$$

$c > 0$ is a constant from the inequalities (cf. [2, 1])

$$|b(u, v, w)| \leq c \|u\|_V \|v\|_V^{\frac{1}{2}} \|v\|_{D(A)}^{\frac{1}{2}} \|w\|_H \quad \forall u \in V, v \in D(A), w \in H; \quad (13.10)$$

$$|b(u, v, w)| \leq c \|u\|_{D(A)}^{\frac{3}{4}} \|u\|_V^{\frac{1}{4}} \|v\|_V \|w\|_H \quad \forall u \in D(A), v \in V, w \in H. \quad (13.11)$$

The auxiliary Problem

$$\begin{cases} \frac{d\eta}{dt} + \nu A\eta + B(\eta, \eta) + B(y, \eta) + B(\eta, y) = g - f \text{ in } V^*, \\ \eta(0) = z_0 - y_0, \end{cases} \quad (13.12)$$

has a strong solution $\eta \in C([0, T]; V) \cap L^2(0, T; D(A))$ with $\frac{d\eta}{dt} \in L^2(0, T; H)$, i.e.

$$\frac{d}{dt}(\eta, v) + \nu((\eta, v)) + b(\eta, \eta, v) + b(y, \eta, v) + b(\eta, y, v) = \langle g - f, v \rangle \quad \text{for all } v \in V,$$

in the sense of distributions on $(0, T)$. In fact, let $\{w_j\}_{j \geq 1} \subset D(A)$ be a special basis (cf. [7, p. 56]), i.e. $Aw_j = \lambda_j w_j$, $j = 1, 2, \dots$, $0 < \lambda_1 \leq \lambda_2 \leq \dots$, $\lambda_j \rightarrow +\infty$, $j \rightarrow +\infty$. We consider Galerkin approximations $\eta_m : [0, T] \rightarrow \text{span}\{w_j\}_{j=1}^m$ for solutions of Problem (13.12) satisfying

$$\frac{d}{dt}(\eta_m, w_j) + \nu((\eta_m, w_j)) + b(\eta_m, \eta_m, w_j) + b(y, \eta_m, w_j) + b(\eta_m, y, w_j) = \langle g - f, w_j \rangle,$$

with $(\eta_m(0), w_j) = (z_0 - y_0, w_j)$, $j = \overline{1, m}$. Due to (13.10), (13.11) and Young's inequality we get

$$\begin{aligned} 2\langle g - f, A\eta_m \rangle &\leq 2\|g - f\|_H \|\eta_m\|_{D(A)} \leq \frac{\nu}{4} \|\eta_m\|_{D(A)}^2 + \frac{4}{\nu} \|f - g\|_H^2; \\ -2b(\eta_m, \eta_m, A\eta_m) &\leq 2c \|\eta_m\|_V^{\frac{3}{2}} \|\eta_m\|_{D(A)}^{\frac{3}{2}} \leq \frac{\nu}{2} \|\eta_m\|_{D(A)}^2 + \frac{27c^4}{2\nu^3} \|\eta_m\|_V^6; \\ -2b(y, \eta_m, A\eta_m) &\leq 2c \|y\|_V \|\eta_m\|_V^{\frac{1}{2}} \|\eta_m\|_{D(A)}^{\frac{3}{2}} \leq \frac{\nu}{2} \|\eta_m\|_{D(A)}^2 + \frac{27c^4}{2\nu^3} \|y\|_{C([0, T]; V)}^4 \|\eta_m\|_V^2; \\ -2b(\eta_m, y, A\eta_m) &\leq 2c \|\eta_m\|_{D(A)}^{\frac{7}{4}} \|\eta_m\|_V^{\frac{1}{4}} \|y\|_V \leq \frac{\nu}{2} \|\eta_m\|_{D(A)}^2 + \frac{7^7 c^8}{2^9 \nu^7} \|y\|_{C([0, T]; V)}^8 \|\eta_m\|_V^2. \end{aligned}$$

Thus,

$$\frac{d}{dt} \|\eta_m\|_V^2 + \frac{\nu}{4} \|\eta_m\|_{D(A)}^2 \leq C(\|\eta_m\|_V^2 + \|\eta_m\|_V^6) + \frac{4}{\nu} \|g - f\|_H^2,$$

where $C > 0$ is a constant from (13.9). Hence, the absolutely continuous function $\varphi = \min\{\|\eta_m\|_V^2, 1\}$ satisfies the inequality $\frac{d}{dt}\varphi \leq 2C\varphi + \frac{4}{\nu}\|g - f\|_H^2$, and therefore $\varphi \leq L(\|z_0 - y_0\|_V^2 + \|g - f\|_{L^2(0, T; H)}^2) < 1$ on $[0, T]$, where $L = \delta^{-1}$. Thus, $\{\eta_n\}_{n \geq 1}$ is bounded in $L^\infty(0, T; V) \cap L^2(0, T; D(A))$ and $\{\frac{d}{dt}\eta_n\}_{n \geq 1}$ is bounded in $L^2(0, T; H)$. In a standard way we get that the limit function η of η_n , $n \rightarrow +\infty$, is a strong solution of Problem (13.12) on $[0, T]$. Due to [6], [1, Chap. 3] the set $\mathcal{D}(z_0, g)$ is one-point $z = y + \eta \in L^8(0, T; (L^4(\Omega))^3)$. So, z is strong solution of Problem (13.1) on $[0, T]$ satisfying (13.8).

The theorem is proved.

13.4 Proof of Theorem 13.1

We provide the proof of Theorem 13.1. Let $f \in L^2(0, T; H)$ and $y_0 \in V$. We consider the 3D controlled Navier-Stokes system (cf. [8, 9])

$$\begin{cases} \frac{dy}{dt} + \nu Ay + B(z, y) = f, \\ y(0) = y_0, \end{cases} \tag{13.13}$$

where $z \in L^8(0, T; (L^4(\Omega))^3)$.

By using standard Galerkin approximations (see [1]) it is easy to show that for any $z \in L^8(0, T; (L^4(\Omega))^3)$ there exists a unique weak solution $y \in L^\infty(0, T; H) \cap L^2(0, T; V)$ of Problem (13.13) on $[0, T]$, that is,

$$\frac{d}{dt} (y, v) + \nu((y, v)) + b(z, y, v) = \langle f, v \rangle, \text{ for all } v \in V, \tag{13.14}$$

in the sense of distributions on $(0, T)$. Moreover, by the inequality

$$|b(u, v, Av)| \leq c_1 \|u\|_{(L^4(\Omega))^3} \|v\|_V^{\frac{1}{4}} \|v\|_{D(A)}^{\frac{7}{4}} \leq \frac{\nu}{2} \|v\|_{D(A)}^2 + c_2 \|u\|_{(L^4(\Omega))^3}^8 \|v\|_V^2, \tag{13.15}$$

for all $u \in (L^4(\Omega))^3$ and $v \in D(A)$, where $c_1, c_2 > 0$ are some constants that do not depend on u, v (cf. [1]), we find that $y \in C([0, T]; V) \cap L^2(0, T; D(A))$ and $B(z, y) \in L^2(0, T; H)$, so $\frac{dy}{dt} \in L^2(0, T; H)$ as well. We add that, for any $z \in L^8(0, T; (L^4(\Omega))^3)$ and corresponding weak solution $y \in C([0, T]; V) \cap L^2(0, T; D(A))$ of (13.13) on $[0, T]$, by using Gronwall inequality, we obtain

$$\begin{aligned} \|y(t)\|_V^2 &\leq \|y_0\|_V^2 e^{2c_2 \int_0^t \|z(t)\|_{(L^4(\Omega))^3}^8 dt}, \quad \forall t \in [0, T]; \\ \int_0^T \|y(t)\|_{D(A)}^2 dt &\leq \|y_0\|_V^2 \left[1 + 2c_2 e^{2c_2 \int_0^T \|z(t)\|_{(L^4(\Omega))^3}^8 dt} \|z\|_{L^8(0,T;(L^4(\Omega))^3)}^8 \right]. \end{aligned} \tag{13.16}$$

Let us consider the operator $F : L^8(0, T; (L^4(\Omega))^3) \rightarrow L^8(0, T; (L^4(\Omega))^3)$, where $F(z) \in C([0, T]; V) \cap L^2(0, T; D(A))$ is the unique weak solution of (13.13) on $[0, T]$ corresponded to $z \in L^8(0, T; (L^4(\Omega))^3)$.

Let us check that F is a compact transformation of Banach space $L^8(0, T; (L^4(\Omega))^3)$ into itself (cf. [10]). In fact, if $\{z_n\}_{n \geq 1}$ is a bounded sequence in $L^8(0, T; (L^4(\Omega))^3)$, then, due to (13.15) and (13.16), the respective weak solutions $y_n, n = 1, 2, \dots$, of Problem (13.13) on $[0, T]$ are uniformly bounded in $C([0, T]; V) \cap L^2(0, T; D(A))$ and their time derivatives $\frac{dy_n}{dt}, n = 1, 2, \dots$, are uniformly bounded in $L^2(0, T; H)$. So, $\{F(z_n)\}_{n \geq 1}$ is a precompact set in $L^8(0, T; (L^4(\Omega))^3)$. In a standard way we deduce that $F : L^8(0, T; (L^4(\Omega))^3) \rightarrow L^8(0, T; (L^4(\Omega))^3)$ is continuous mapping.

Since F is a compact transformation of $L^8(0, T; (L^4(\Omega))^3)$ into itself, Schaefer's Theorem (cf. [10, p. 133] and references therein) and Theorem 13.2 provide the statement of Theorem 13.1. We note that Theorem 13.2 implies that the set

$\{z \in L^8(0, T; (L^4(\Omega))^3) : z = \lambda F(z), \lambda \in (0, 1)\}$ is bounded in $L^8(0, T; (L^4(\Omega))^3)$ iff the set defined in (13.6) is bounded in $L^8(0, T; (L^4(\Omega))^3)$.

The theorem is proved.

Acknowledgments The authors thank Professors J.M. Ball, V.V. Chepyzhov, and M.Z. Zgurovsky for useful suggestions during the preparation of this manuscript. The first author was partially supported by the Ukrainian State Fund for Fundamental Researches under grants GP/F44/076, GP/F49/070, and by the NAS of Ukraine under grant 2273/13.

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