# **Entropic Isoperimetric Inequalities**



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#### 2020 Mathematics Subject Classification: 60E15; 94A17

# 1 Introduction

The entropic isoperimetric inequality asserts that

$$N(X) I(X) \ge 2\pi e n \tag{1.1}$$

for any random vector X in  $\mathbb{R}^n$  with a smooth density. Here

$$N(X) = \exp\left\{-\frac{2}{n}\int p(x)\log p(x)\,dx\right\} \quad \text{and} \quad I(X) = \int \frac{|\nabla p(x)|^2}{p(x)}\,dx$$

denote the Shannon entropy power and the Fisher information of *X* with density *p*, respectively (with integration with respect to Lebesgue measure dx on  $\mathbb{R}^n$  which may be restricted to the supporting set supp $(p) = \{x : p(x) > 0\}$ ).

This inequality was discovered by Stam [15] where it was treated in dimension one. It is known to hold in any dimension, and the standard normal distribution on  $\mathbb{R}^n$  plays an extremal role in it. Later on, Costa and Cover [6] pointed out a remarkable analogy between (1.1) and the classical isoperimetric inequality relating the surface of an arbitrary body A in  $\mathbb{R}^n$  to its volume  $vol_n(A)$ . The terminology "isoperimetric inequality for entropies" goes back to Dembo, Costa, and Thomas [8].

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<sup>©</sup> The Author(s), under exclusive license to Springer Nature Switzerland AG 2023 R. Adamczak et al. (eds.), *High Dimensional Probability IX*, Progress in Probability 80, https://doi.org/10.1007/978-3-031-26979-0\_3

As Rényi entropies have become a focus of numerous investigations in the recent time, it is natural to explore more general relations of the form

$$N_{\alpha}(X) I(X) \ge c_{\alpha,n} \tag{1.2}$$

for the functional

$$N_{\alpha}(X) = \left(\int p(x)^{\alpha} dx\right)^{-\frac{2}{n(\alpha-1)}}.$$
(1.3)

It is desirable to derive (1.2) with optimal constants  $c_{\alpha,n}$  independent of the density p, where  $\alpha \in [0, \infty]$  is a parameter called the order of the Rényi entropy power  $N_{\alpha}(X)$ . Another representation

$$N_{\alpha}(X)^{-\frac{n}{2}} = \|p\|_{L^{\alpha-1}(p(x)\,dx)}$$

shows that  $N_{\alpha}$  is non-increasing in  $\alpha$ . This allows one to define the Rényi entropy power for the two extreme values by the monotonicity to be

$$N_{\infty}(X) = \lim_{\alpha \to \infty} N_{\alpha}(X) = \|p\|_{\infty}^{-\frac{2}{n}},$$

$$N_{0}(X) = \lim_{\alpha \to 0} N_{\alpha}(X) = \operatorname{vol}_{n}(\operatorname{supp}(p))^{\frac{2}{n}},$$
(1.4)

where  $||p||_{\infty} = \operatorname{ess sup} p(x)$ . As a standard approach, one may also put  $N_1(X) = \lim_{\alpha \downarrow 1} N_{\alpha}(X)$  which returns us to the usual definition of the Shannon entropy power  $N_1(X) = N(X)$  under mild moment assumptions (such as  $N_{\alpha}(X) > 0$  for some  $\alpha > 1$ ).

Returning to (1.1)–(1.2), the following two natural questions arise.

**Question 1** Given *n*, for which range  $\mathfrak{A}_n$  of the values of  $\alpha$  does (1.2) hold with some positive constant?

**Question 2** What is the value of the optimal constant  $c_{\alpha,n}$  and can the extremizers in (1.2) be described?

The entropic isoperimetric inequality (1.1) answers both questions for the order  $\alpha = 1$  with an optimal constant  $c_{1,n} = 2\pi e n$ . As for the general order, let us first stress that, by the monotonicity of  $N_{\alpha}$  with respect to  $\alpha$ , the function  $\alpha \mapsto c_{\alpha,n}$  is also non-increasing. Hence, the range in Question 1 takes necessarily the form  $\mathfrak{A}_n = [0, \alpha_n)$  or  $\mathfrak{A}_n = [0, \alpha_n]$  for some critical value  $\alpha_n \in [0, \infty]$ . The next assertion specifies these values.

**Theorem 1.1** We have

$$\mathfrak{A}_n = \begin{cases} [0,\infty] & \text{for } n = 1, \\ [0,\infty) & \text{for } n = 2, \\ [0,\frac{n}{n-2}] & \text{for } n \ge 3. \end{cases}$$

.

Entropic Isoperimetric Inequalities

Thus, in the one dimensional case there is no restriction on  $\alpha$  (the range is full). In fact, this already follows from the elementary sub-optimal inequality

$$N_{\infty}(X)I(X) \ge 1,\tag{1.5}$$

implying that  $c_{\alpha,1} \ge 1$  for all  $\alpha$ . To see this, assume that I(X) is finite, so that X has a (locally) absolutely continuous density p, thus differentiable almost everywhere. Since p is non-negative, any point  $y \in \mathbb{R}$  such that p(y) = 0 is a local minimum, and necessarily p'(y) = 0 (as long as p is differentiable at y). Hence, applying the Cauchy inequality, we have

$$\begin{split} \int_{-\infty}^{\infty} |p'(y)| \, dy &= \int_{p(y)>0} \frac{|p'(y)|}{\sqrt{p(y)}} \sqrt{p(y)} \, dy \\ &\leq \Big( \int_{p(y)>0} \frac{p'(y)^2}{p(y)} \, dy \Big)^{1/2} \left( \int_{p(y)>0} p(y) \, dy \right)^{1/2} = \sqrt{I(X)}. \end{split}$$

It follows that p has a bounded total variation not exceeding  $\sqrt{I(X)}$ , so  $p(x) \le \sqrt{I(X)}$  for every  $x \in \mathbb{R}$ . This amounts to (1.5) according to (1.4) for n = 1.

Turning to Question 2, we will see that the optimal constants  $c_{\alpha,1}$  together with the extremizers in (1.2) may be explicitly described in the one dimensional case for every  $\alpha$  using the results due to Nagy [13]. Since the transformation of these results in the information-theoretic language is somewhat technical, we discuss this case in detail in the next three sections (Sects. 2, 3, and 4). Let us only mention here that

$$4 \le c_{\alpha,1} \le 4\pi^2,$$

where the inequalities are sharp for  $\alpha = \infty$  and  $\alpha = 0$ , respectively, with extremizers

$$p(x) = \frac{1}{2} e^{-|x|}$$
 and  $p(x) = \frac{2}{\pi} \cos^2(x) \mathbf{1}_{\{|x| \le \frac{\pi}{2}\}}.$ 

The situation in higher dimensions is more complicated, and only partial answers to Question 2 will be given here. Anyway, in order to explore the behavior of the constants  $c_{\alpha,n}$ , one should distinguish between the dimensions n = 2 and  $n \ge 3$  (which is also suggested by Theorem 1.1). In the latter case, these constants can be shown to satisfy

$$4\pi n(n-2)\left(\frac{\Gamma(\frac{n}{2})}{\Gamma(n)}\right)^{\frac{2}{n}} \le c_{\alpha,n} \le 4\pi^2 n, \qquad 0 \le \alpha \le \frac{n}{n-2},$$

where the left inequality is sharp and corresponds to the critical order  $\alpha = \frac{n}{n-2}$ . With respect to the growing dimension, these constants are asymptotically  $2\pi en + O(1)$ , which exhibits nearly the same behavior as for the order  $\alpha = 1$ . However (which is rather surprising), the extremizers for the critical order exist for  $n \ge 5$  only and are described as densities of the (generalized) Cauchy distributions on  $\mathbb{R}^n$ . We discuss these issues in Sect. 7, while Sect. 6 deals with dimension n = 2, where some description of the constants  $c_{\alpha,2}$  will be given for the range  $\alpha \in [\frac{1}{2}, \infty)$ .

We end this introduction by giving an equivalent formulation of the isoperimetric inequalities (1.2) in terms of functional inequalities of Sobolev type. As was noticed by Carlen [5], in the classical case  $\alpha = 1$ , (1.1) is equivalent to the logarithmic Sobolev inequality of Gross [9], cf. also [4]. However, when  $\alpha \neq 1$ , a different class of inequalities should be involved. Namely, using the substitution  $p = f^2 / \int f^2$  (here and in the sequel integrals are understood with respect to the Lebesgue measure on  $\mathbb{R}^n$ ), we have

$$N_{\alpha}(X) = \left(\int f^{2\alpha}\right)^{-\frac{2}{n(\alpha-1)}} \left(\int f^2\right)^{\frac{2\alpha}{n(\alpha-1)}}$$

and

$$I(X) = 4 \int |\nabla f|^2 / \int f^2$$

Therefore (provided that f is square integrable), (1.2) can be equivalently reformulated as a homogeneous analytic inequality

$$\left(\int |f|^{2\alpha}\right)^{\frac{2}{n(\alpha-1)}} \le \frac{4}{c_{\alpha,n}} \int |\nabla f|^2 \left(\int f^2\right)^{\frac{\alpha(2-n)+n}{n(\alpha-1)}},\tag{1.6}$$

where we can assume that f is smooth and has gradient  $\nabla f$  (however, when speaking about extremizers, the function f should be allowed to belong to the Sobolev class  $W_1^2(\mathbb{R}^n)$ ). Such inequalities were introduced by Moser [11, 12] in the following form

$$\left(\int |f|^{2+\frac{4}{n}}\right) \le B_n \int |\nabla f|^2 \left(\int f^2\right)^{\frac{2}{n}}.$$
(1.7)

More precisely, (1.7) corresponds to (1.6) for the specific choice  $\alpha = 1 + \frac{2}{n}$ . Here, the one dimensional case is covered by Nagy's paper with the optimal factor  $B_1 = \frac{4}{\pi^2}$ . This corresponds to  $\alpha = 3$  and n = 1, and therefore  $c_{3,1} = \pi^2$  which complements the picture depicted above. To the best of our knowledge, the best constants  $B_n$  for  $n \ge 2$  are not known. However, using the Euclidean log-Sobolev inequality and the optimal Sobolev inequality, Beckner [2] proved that asymptotically  $B_n \sim \frac{2}{\pi e_n}$ .

Both Moser's inequality (1.7) and (1.6) with a certain range of  $\alpha$  enter the general framework of Gagliardo–Nirenberg's inequalities

$$\left(\int |f|^r\right)^{\frac{1}{r}} \leq \kappa_n(q, r, s) \left(\int |\nabla f|^q\right)^{\frac{\theta}{q}} \left(\int |f|^s\right)^{\frac{1-\theta}{s}}$$
(1.8)

with  $1 \le q, r, s \le \infty$ ,  $0 \le \theta \le 1$ , and  $\frac{1}{r} = \theta \left(\frac{1}{q} - \frac{1}{n}\right) + (1 - \theta) \frac{1}{s}$ . We will make use of the knowledge on Gagliardo–Nirenberg's inequalities to derive information on (1.2).

In the sequel, we denote by  $||f||_r = (\int |f|^r)^{\frac{1}{r}}$  the  $L^r$ -norm of f with respect to the Lebesgue measure on  $\mathbb{R}^n$  (and use this functional also in the case 0 < r < 1).

# 2 Nagy's Theorem

In the next three sections we focus on dimension n = 1, in which case the entropic isoperimetric inequality (1.2) takes the form

$$N_{\alpha}(X) I(X) \ge c_{\alpha,1} \tag{2.1}$$

for the Rényi entropy

$$N_{\alpha}(X) = \left(\int p(x)^{\alpha} dx\right)^{-\frac{2}{\alpha-1}}$$

and the Fisher information

$$I(X) = \int \frac{p'(x)^2}{p(x)} dx = 4 \int \left(\frac{d}{dx}\sqrt{p(x)}\right)^2 dx.$$

In dimension one, our basic functional space is the collection of all (locally) absolutely continuous functions on the real line whose derivatives are understood in the Radon–Nikodym sense. We already know that (2.1) holds for all  $\alpha \in [0, \infty]$ .

According to (1.6), the family (2.1) takes now the form

$$\int |f|^{2\alpha} \le \left(\frac{4}{c_{\alpha,1}}\right)^{\frac{\alpha-1}{2}} \left(\int f'^2\right)^{\frac{\alpha-1}{2}} \left(\int f^2\right)^{\frac{\alpha+1}{2}}$$
(2.2)

when  $\alpha > 1$ , and

$$\int f^2 \leq \left(\frac{4}{c_{\alpha,1}}\right)^{\frac{1-\alpha}{1+\alpha}} \left(\int f^{\prime 2}\right)^{\frac{1-\alpha}{1+\alpha}} \left(\int |f|^{2\alpha}\right)^{\frac{2}{1+\alpha}}$$
(2.3)

when  $\alpha \in (0, 1)$ .

In fact, these two families of inequalities can be seen as sub-families of the following one, studied by Nagy [13],

$$\int |f|^{\gamma+\beta} \le D\Big(\int |f'|^p\Big)^{\frac{\beta}{pq}} \Big(\int |f|^{\gamma}\Big)^{1+\frac{\beta(p-1)}{pq}}$$

with

$$p > 1, \quad \beta, \gamma > 0, \quad q = 1 + \frac{\gamma(p-1)}{p},$$
 (2.4)

and some constants  $D = D_{\gamma,\beta,p}$  depending on  $\gamma, \beta$ , and p, only. For such parameters, introduce the functions  $y_{p,\gamma} = y_{p,\gamma}(t)$  defined for  $t \ge 0$  by

$$y_{p,\gamma}(t) = \begin{cases} (1+t)^{\frac{p}{p-\gamma}} & \text{if } p < \gamma, \\ e^{-t} & \text{if } p = \gamma, \\ (1-t)^{\frac{p}{p-\gamma}} \mathbf{1}_{[0,1]}(t) & \text{if } p > \gamma. \end{cases}$$

To involve the parameter  $\beta$ , define additionally  $y_{p,\gamma,\beta}$  implicitly as follows. Put  $y_{p,\gamma,\beta}(t) = u, 0 \le u \le 1$ , with

$$t = \int_{u}^{1} \left( s^{\gamma} (1 - s^{\beta}) \right)^{-\frac{1}{p}} ds$$

if  $p \le \gamma$ . If  $p > \gamma$ , then  $y_{p,\gamma,\beta}(t) = u, 0 \le u \le 1$ , is the solution of the above equation for

$$t \le t_0 = \int_0^1 \left( s^{\gamma} (1 - s^{\beta}) \right)^{-\frac{1}{p}} ds$$

and  $y_{p,\gamma,\beta}(t) = 0$  for all  $t > t_0$ . With these notations, Nagy established the following result.

**Theorem 2.1 ([13])** Under the constraint (2.4), for any (locally) absolutely continuous function  $f : \mathbb{R} \to \mathbb{R}$ ,

(i)

$$\|f\|_{\infty} \le \left(\frac{q}{2}\right)^{\frac{1}{q}} \left(\int |f'|^{p}\right)^{\frac{1}{pq}} \left(\int |f|^{\gamma}\right)^{\frac{p-1}{pq}}.$$
(2.5)

Moreover, the extremizers take the form  $f(x) = ay_{p,\gamma}(|bx + c|)$  with a, b, c constants  $(b \neq 0)$ .

(ii)

$$\int |f|^{\beta+\gamma} \leq \left(\frac{q}{2} H\left(\frac{q}{\beta}, \frac{p-1}{p}\right)\right)^{\frac{\beta}{q}} \left(\int |f'|^p\right)^{\frac{\beta}{pq}} \left(\int |f|^\gamma\right)^{1+\frac{\beta(p-1)}{pq}},$$
(2.6)

where

$$H(u,v) = \frac{\Gamma(1+u+v)}{\Gamma(1+u)\Gamma(1+v)} \left(\frac{u}{u+v}\right)^u \left(\frac{v}{u+v}\right)^v, \quad u,v \ge 0.$$

*Moreover, the extremizers take the form*  $f(x) = ay_{p,\gamma,\beta}(|bx + c|)$  *with* a, b, c *constants*  $(b \neq 0)$ .

Here,  $\Gamma$  denotes the classical Gamma function, and we use the convention that H(u, 0) = H(0, v) = 1 for  $u, v \ge 0$ . It was mentioned by Nagy that H is monotone in each variable. Moreover, since  $H(u, 1) = (1 + \frac{1}{u})^{-u}$  is between 1 and  $\frac{1}{e}$ , one has  $1 > H(u, v) > (1 + \frac{1}{u})^{-u} > \frac{1}{e}$  for all 0 < v < 1. This gives a two-sided bound

$$1 \ge H\left(\frac{q}{\beta}, \frac{p-1}{p}\right) > \left(1 + \frac{\beta}{q}\right)^{-\frac{q}{\beta}} > \frac{1}{e}.$$

#### **3** One Dimensional Isoperimetric Inequalities for Entropies

The inequalities (2.2) and (2.3) correspond to (2.6) with parameters

$$p = \gamma = q = 2, \ \beta = 2(\alpha - 1)$$
 in the case  $\alpha > 1$ 

and

$$p = 2, \ \beta = 2(1 - \alpha), \ \gamma = 2\alpha, \ q = 1 + \alpha$$
 in the case  $\alpha \in (0, 1),$ 

respectively. Hence, as a corollary from Theorem 2.1, we get the following statement which solves Question 2 when n = 1. Note that, by Theorem 2.1, the extremal distributions (their densities p) in (2.1) are determined in a unique way up to nondegenerate affine transformations of the real line. So, it is sufficient to indicate just one specific extremizer for each admissible collection of the parameters. Recall the definition of the optimal constants  $c_{\alpha,1}$  from (2.1).

#### Theorem 3.1

(i) In the case  $\alpha = \infty$ , we have

$$c_{\infty,1} = 4.$$

Moreover, the density  $p(x) = \frac{1}{2} e^{-|x|}$  ( $x \in \mathbb{R}$ ) of the two-sided exponential distribution represents an extremizer in (2.1).

(ii) In the case  $1 < \alpha < \infty$ , we have

$$c_{\alpha,1} = \frac{2\pi}{\alpha - 1} \left(\frac{2}{\alpha + 1}\right)^{\frac{\alpha - 3}{\alpha - 1}} \left(\frac{\Gamma(\frac{1}{\alpha - 1})}{\Gamma(\frac{\alpha + 1}{2(\alpha - 1)})}\right)^2.$$

Moreover, the density  $p(x) = a \cosh(x)^{-\frac{2}{\alpha-1}}$  with a normalization constant  $a = \frac{1}{\sqrt{\pi}} \frac{\Gamma(\frac{\alpha+1}{2(\alpha-1)})}{\Gamma(\frac{1}{\alpha-1})}$  represents an extremizer in (2.1). (iii) In the case  $0 < \alpha < 1$ ,

$$c_{\alpha,1} = \frac{2\pi}{1-\alpha} \left(\frac{2}{1+\alpha}\right)^{\frac{1+\alpha}{1-\alpha}} \left(\frac{\Gamma(\frac{1+\alpha}{2(1-\alpha)})}{\Gamma(\frac{1}{1-\alpha})}\right)^2.$$

Moreover, the density  $p(x) = a \cos(x)^{\frac{2}{1-\alpha}} \mathbf{1}_{\left[-\frac{\pi}{2}, \frac{\pi}{2}\right]}(x)$  with constant  $a = \frac{1}{\sqrt{\pi}} \frac{\Gamma(\frac{3-\alpha}{1-\alpha})}{\Gamma(\frac{3-\alpha}{2(3-\alpha)})}$  represents an extremizer in (2.1).

To prove the theorem, we need a simple technical lemma.

#### Lemma 3.2

(i) Given a > 0 and  $t \ge 0$ , the (unique) solution  $y \in (0, 1]$  to the equation  $\int_{y}^{1} \frac{ds}{s\sqrt{1-s^{a}}} = t$  is given by

$$y = \left[\cosh\left(\frac{at}{2}\right)\right]^{-\frac{2}{a}}.$$

(ii) Given a, b > 0 and  $c \in \mathbb{R}$ , we have

$$\int_{-\infty}^{\infty} \cosh(|bx+c|)^{-a} dx = \frac{\sqrt{\pi}}{b} \frac{\Gamma(\frac{a}{2})}{\Gamma(\frac{a+1}{2})}.$$

(iii) Given  $a \in (0, 1)$  and  $u \in [0, 1]$ , we have

$$\int_{u}^{1} \frac{ds}{s^{a}\sqrt{1-s^{2(1-a)}}} = \frac{1}{1-a} \arccos(u^{1-a}).$$

*Remark 3.3* Since  $\Gamma(\frac{a+1}{2}) = \Gamma(m + \frac{1}{2}) = \frac{(2m)!}{4^m m!} \sqrt{\pi}$  for a = 2m with an integer  $m \ge 1$ , for such particular values of a, we have

$$\int_{-\infty}^{\infty} \cosh(|bx+c|)^{-a} \, dx = \frac{1}{b} \cdot \frac{4^m m! \, (m-1)!}{(2m)!}.$$

**Proof of Lemma 3.2** Changing the variable  $u = \sqrt{1 - s^a}$ , we have

$$\int_{y}^{1} \frac{ds}{s\sqrt{1-s^{a}}} = \frac{2}{a} \int_{0}^{\sqrt{1-y^{a}}} \frac{du}{1-u^{2}} = \frac{1}{a} \log\left(\frac{1+\sqrt{1-y^{a}}}{1-\sqrt{1-y^{a}}}\right).$$

Inverting this equality leads to the desired result of item (i).

For item (ii) we use the symmetry of the cosh-function together with the change of variables u = bx + c and then  $t = \sinh(u)^2$  to get

$$\int_{-\infty}^{\infty} \cosh(|bx+c|)^{-a} dx = \frac{1}{b} \int_{-\infty}^{\infty} \cosh(|u|)^{-a} du$$
$$= \frac{2}{b} \int_{0}^{\infty} \cosh(u)^{-a} du = \frac{1}{b} \int_{0}^{\infty} t^{-\frac{1}{2}} (1+t)^{-\frac{a+1}{2}} dt.$$

To obtain the result, we need to perform a final change of variables  $v = \frac{1}{1+t}$ . This turns the last integral into

$$\int_0^1 (1-v)^{-\frac{1}{2}} v^{\frac{a}{2}-1} dv = B\left(\frac{1}{2}, \frac{a}{2}\right) = \sqrt{\pi} \frac{\Gamma(\frac{a}{2})}{\Gamma(\frac{a+1}{2})},$$

where we used the beta function  $B(x, y) = \int_0^1 (1 - v)^{x-1} v^{y-1} dv = \frac{\Gamma(x)\Gamma(y)}{\Gamma(x+y)}, x, y > 0.$ 

Finally, in item (iii), a change of variables leads to

$$\int_{u}^{1} \frac{ds}{s^{a}\sqrt{1-s^{2(1-a)}}} = \frac{1}{1-a} \int_{u}^{1} \frac{ds^{1-a}}{\sqrt{1-s^{2(1-a)}}}$$
$$= \frac{1}{1-a} \int_{u^{1-a}}^{1} \frac{dv}{\sqrt{1-v^{2}}} = \frac{1}{1-a} \arccos(u^{1-a}).$$

**Proof of Theorem 3.1** When  $\alpha = \infty$  as in the case (i), (2.2) with  $\int f^2 = 1$  becomes

$$||f||_{\infty} \le \left(\frac{4}{c_{\infty,1}}\int f'^2\right)^{\frac{1}{4}}.$$

This corresponds to (2.5) with parameters  $p = q = \gamma = 2$ . Therefore, item (i) of Theorem 2.1 applies and leads to

$$\|f\|_{\infty} \le \left(\int f'^2\right)^{\frac{1}{4}},$$

that is,  $c_{\infty,1} = 4$ . Moreover, the extremizers in (2.5) are given by

$$f(x) = ay_{2,2}(|bx + c|) = a e^{-|bx+c|}, \quad b \neq 0, \ a, c \in \mathbb{R}.$$

But, the extremizers in (2.1) are of the form  $p = f^2 / \int f^2$  with f an extremizer in (2.5). The desired result then follows after a change of variables.

Next, let us turn to the case (ii), where  $1 < \alpha < \infty$ . Here (2.1) is equivalent to (2.2) and corresponds to (2.6) with  $p = \gamma = q = 2$  and  $\beta = 2(\alpha - 1)$ . Therefore, by Theorem 2.1,  $(\frac{4}{c_{\alpha,1}})^{\frac{\alpha-1}{2}} = H(\frac{1}{\alpha-1}, \frac{1}{2})^{\alpha-1}$ , so that

$$c_{\alpha,1} = \frac{4}{H(\frac{1}{\alpha-1}, \frac{1}{2})^2} = 4 \frac{\Gamma(1 + \frac{1}{\alpha-1})^2 \Gamma(\frac{3}{2})^2}{\Gamma(\frac{3}{2} + \frac{1}{\alpha-1})^2} \left(\frac{\frac{1}{\alpha-1} + \frac{1}{2}}{\frac{1}{\alpha-1}}\right)^{\frac{2}{\alpha-1}} \left(\frac{\frac{1}{\alpha-1} + \frac{1}{2}}{\frac{1}{2}}\right)$$
$$= \pi \left(\frac{\frac{1}{\alpha-1}}{\frac{\alpha+1}{2(\alpha-1)}}\right)^2 \frac{\Gamma(\frac{1}{\alpha-1})^2}{\Gamma(\frac{\alpha+1}{2(\alpha-1)})^2} \left(\frac{\alpha+1}{2}\right)^{\frac{2}{\alpha-1}} \left(\frac{\alpha+1}{\alpha-1}\right),$$

where we used the identities  $\Gamma(3/2) = \sqrt{\pi}/2$  and  $\Gamma(1+z) = z\Gamma(z)$ . This leads to the desired expression for  $c_{\alpha,1}$ .

As for extremizers, item (ii) of Theorem 2.1 applies and asserts that the equality cases in (2.2) are reached, up to numerical factors, for functions f(x) = y(|bx+c|), with  $b \neq 0, c \in \mathbb{R}$ , and  $y: [0, \infty) \rightarrow \mathbb{R}$  defined implicitly for  $t \in [0, \infty)$  by  $y(t) = u, 0 \le u \le 1$ , with

$$t = \int_{u}^{1} \left( s^{2} (1 - s^{2(\alpha - 1)}) \right)^{-\frac{1}{2}} ds = \int_{u}^{1} \frac{1}{s\sqrt{1 - s^{2(\alpha - 1)}}} ds.$$

Now, Lemma 3.2 provides the solution  $y(t) = (\cosh((\alpha - 1)t))^{-\frac{1}{\alpha-1}}$ . Therefore, the extremizers in (2.2) are reached, up to numerical factors, for functions of the form

$$f(x) = (\cosh(|bx + c|))^{-\frac{1}{\alpha - 1}}, \quad b \neq 0, \ c \in \mathbb{R}.$$

Similarly to the case (i), the extremizers in (2.1) are of the form  $p = f^2 / \int f^2$  with f an extremizer in (2.2). Therefore, by Lemma 3.2, with some b > 0 and  $c \in \mathbb{R}$ ,

$$p(x) = \frac{\cosh(|bx+c|)^{-\frac{2}{\alpha-1}}}{\int \cosh(|bx+c|)^{-\frac{2}{\alpha-1}} dx} = \frac{b}{\sqrt{\pi}} \frac{\Gamma(\frac{\alpha+1}{2(\alpha-1)})}{\Gamma(\frac{1}{\alpha-1})} \cosh(bx+c)^{-\frac{2}{\alpha-1}}$$

as announced.

Finally, let us turn to item (iii), when  $\alpha \in (0, 1)$ . As already mentioned, (2.1) is equivalent to (2.3) and therefore corresponds to (2.6) with p = 2,  $\beta = 2(1 - \alpha)$ ,  $\gamma = 2\alpha$ , and  $q = 1 + \alpha$ . An application of Theorem (2.1) leads to the desired conclusion after some algebra (which we leave to the reader) concerning the explicit value of  $c_{\alpha,1}$ . In addition, the extremizers are of the form  $p(x) = ay^2(|bx+c|)$ , with a a normalization constant,  $b \neq 0$ , and  $c \in \mathbb{R}$ . Here y = y(t) is defined implicitly by the equation

$$t = \int_{y}^{1} \frac{1}{s^{\alpha}\sqrt{1 - s^{2(1-\alpha)}}} \, ds$$

for  $t \le t_0 = \int_0^1 \frac{1}{s^\alpha \sqrt{1 - s^{2(1-\alpha)}}} ds$  and y(t) = 0 for  $t > t_0$ . Item (iii) of Lemma 3.2 asserts that

$$t_0 = \frac{\pi}{2(1-\alpha)}$$
 and  $y(t) = \left(\cos((1-\alpha)t)\right)^{\frac{1}{1-\alpha}} \mathbf{1}_{[0,\frac{\pi}{2(1-\alpha)}]}(t).$ 

This leads to the desired conclusion.

### 4 Special Orders

As an illustration, here we briefly mention some explicit values of  $c_{\alpha,1}$  and extremizers for specific values of the parameter  $\alpha$  in the one dimensional entropic isoperimetric inequality

$$N_{\alpha}(X) I(X) \ge c_{\alpha,1}. \tag{4.1}$$

The order  $\alpha = 0$  The limit in item (iii) of Theorem 3.1 leads to the optimal constant

$$c_{0,1} = \lim_{\alpha \to 0} c_{\alpha,1} = 4\pi^2.$$

Since all explicit expressions are continuous with respect to  $\alpha$ , the limits of the extremizers in (2.1) for  $\alpha \to 0$  represent extremizers in (2.1) for  $\alpha = 0$ . Therefore, the densities

$$p(x) = \frac{2b}{\pi} \cos^2(bx+c) \, \mathbb{1}_{\left[-\frac{\pi}{2}; \frac{\pi}{2}\right]}(bx+c), \quad b > 0, \ c \in \mathbb{R},$$

are extremizers in (2.1) with  $\alpha = 0$ .

The order  $\alpha = \frac{1}{2}$  Direct computation leads to  $c_{\frac{1}{2},1} = (4/3)^3 \pi^2$ . Moreover, the extremizers in (2.1) are of the form

$$p(x) = \frac{8b}{3\pi} \cos^4(bx+c) \,\mathbf{1}_{\left[-\frac{\pi}{2};\frac{\pi}{1}\right]}(bx+c), \quad b > 0, \ c \in \mathbb{R}.$$

The order  $\alpha = 1$  This case corresponds to Stam's isoperimetric inequality for entropies. Here  $c_{1,1} = 2\pi e$ , and, using the Stirling formula, one may notice that indeed

$$c_{1,1} = \lim_{\alpha \to 1} c_{\alpha,1} = 2\pi e.$$

Moreover, Gaussian densities can be obtained from the extremizers  $p(x) = \cosh(bx + c)^{-\frac{2}{\alpha-1}}$  with  $b = b'\sqrt{\alpha-1}$ ,  $c = c'\sqrt{\alpha-1}$  in the limit as  $\alpha \downarrow 1$ . (Note that the limit  $\alpha \uparrow 1$  would lead to the same conclusion.)

The order  $\alpha = 2$  A direct computation leads to  $c_{2,1} = 12$  with extremizers of the form

$$p(x) = \frac{b}{2\cosh^2(bx+c)}, \quad b > 0, \ c \in \mathbb{R}.$$

In this case, the entropic isoperimetric inequality may equivalently be stated in terms of the Fourier transform  $\hat{p}(t) = \int e^{itx} p(x)$ ,  $t \in \mathbb{R}$ , of the density *p*. Indeed, thanks to Plancherel's identity, we have

$$N_2(X)^{-1/2} = \int p^2 = \frac{1}{2\pi} \int |\hat{p}|^2.$$

Therefore, the (optimal) isoperimetric inequality for entropies yields the relation

$$\int |\hat{p}|^2 \le \pi \sqrt{\frac{I(X)}{3}}$$

which is a global estimate on the  $L^2$ -norm of  $\hat{p}$ . In [18], Zhang derived the following pointwise estimate: If the random variable X with density p has finite Fisher information I(X), then (see also [3] for an alternative proof)

$$|\hat{p}(t)| \le \frac{I(X)}{I(X) + t^2}, \quad t \in \mathbb{R}.$$

The latter leads to some bounds on  $c_{2,1}$ , namely

$$N_2(X)^{-1/2} = \frac{1}{2\pi} \int |\hat{p}|^2 \le \frac{1}{2\pi} \int \frac{I(X)^2}{(I(X) + t^2)^2} dt = \frac{1}{2} \sqrt{I(X)}.$$

Hence  $N_2(X)I(X) \ge 4$  that should be compared to  $N_2(X)I(X) \ge 12$ .

The order  $\alpha = 3$  Then  $c_{3,1} = \pi^2$ , and the extremizers are of the form

$$p(x) = \frac{b}{\pi \cosh(bx+c)}, \quad b > 0, \ c \in \mathbb{R}.$$

The order  $\alpha = \infty$  From Theorem 3.1,  $c_{\infty,1} = 4$ , and the extremizers are of the form

$$p(x) = b e^{-|bx+c|}, \quad b > 0, \ c \in \mathbb{R}.$$

### 5 Fisher Information in Higher Dimensions

In order to perform the transition from the entropic isoperimetric inequality (1.2) to the form of the Gagliardo–Nirenberg inequality such as (1.8) via the change of functions  $p = f^2 / \int f^2$  and back, and to justify the correspondence of the constants in the two types of inequalities, let us briefly fix some definitions and recall some approximation properties of the Fisher information. This is dictated by the observation that in general f in (1.8) does not need to be square integrable, and then p will not be defined as a probability density.

The Fisher information of a random vector X in  $\mathbb{R}^n$  with density p may be defined by means of the formula

$$I(X) = I(p) = 4 \int |\nabla \sqrt{p}|^2.$$
 (5.1)

This functional is well-defined and finite if and only if  $f = \sqrt{p}$  belongs to the Sobolev space  $W_1^2(\mathbb{R}^n)$ . There is the following characterization: A function f belongs to  $W_1^2(\mathbb{R}^n)$ , if and only if it belongs to  $L^2(\mathbb{R}^n)$  and

$$\sup_{h \neq 0} \left[ \frac{1}{|h|} \| f(x+h) - f(x) \|_2 \right] < \infty.$$

In this case, there is a unique vector-function  $g = (g_1, \ldots, g_n)$  on  $\mathbb{R}^n$  with components in  $L^2(\mathbb{R}^n)$ , called a weak gradient of f and denoted  $g = \nabla f$ , with the property that

$$\int gv = -\int f\nabla v \quad \text{for all } v \in C_0^\infty(\mathbb{R}^n).$$
(5.2)

As usual,  $C_0^{\infty}(\mathbb{R}^n)$  denotes the class of all  $C^{\infty}$ -smooth, compactly supported functions on  $\mathbb{R}^n$ . Still equivalently, there is a representative  $\bar{f}$  of f which is absolutely continuous on almost all lines parallel to the coordinate axes and whose partial derivatives  $\partial_{x_k} \bar{f}$  belong to  $L^2(\mathbb{R}^n)$ . In particular,  $g_k(x) = \partial_{x_k} \bar{f}(x)$  for almost all  $x \in \mathbb{R}^n$  (cf. [19], Theorems 2.1.6 and 2.1.4).

Applied to  $f = \sqrt{p}$  with a probability density p on  $\mathbb{R}^n$ , the property that  $f \in W_1^2(\mathbb{R}^n)$  ensures that p has a representative  $\bar{p}$  which is absolutely continuous on almost all lines parallel to the coordinate axes and such that the functions  $\partial_{x_k} \bar{p} / \sqrt{p}$  belong to  $L^2(\mathbb{R}^n)$ . Moreover,

$$I(p) = \sum_{k=1}^{n} \left\| \frac{\partial_{x_k} \bar{p}}{\sqrt{p}} \right\|_2^2.$$

Note that  $W_1^2(\mathbb{R}^n)$  is a Banach space for the norm defined by

$$\|f\|_{W_1^2}^2 = \|f\|_2^2 + \|\nabla f\|_2^2$$
  
=  $\|f\|_2^2 + \|g_1\|_2^2 + \dots + \|g_n\|_2^2$   $(g = \nabla f).$ 

We use the notation  $N_{\alpha}(X) = N_{\alpha}(p)$  when a random vector X has density p.

**Proposition 5.1** Given a (probability) density p on  $\mathbb{R}^n$  such that I(p) is finite, there exists a sequence of densities  $p_k \in C_0^{\infty}(\mathbb{R}^n)$  satisfying as  $k \to \infty$ 

(a)  $I(p_k) \to I(p)$ , and (b)  $N_{\alpha}(p_k) \to N_{\alpha}(p)$  for any  $\alpha \in (0, \infty)$ ,  $\alpha \neq 1$ .

**Proof** Let us recall two standard approximation arguments. Fix a non-negative function  $\omega \in C_0^{\infty}(\mathbb{R}^n)$  supported in the closed unit ball  $\overline{B}_n(0, 1) = \{x \in \mathbb{R}^n : |x| \le 1\}$  and such that  $\int \omega = 1$ , and put  $\omega_{\varepsilon}(x) = \varepsilon^{-n} \omega(x/\varepsilon)$  for  $\varepsilon > 0$ . Given a locally integrable function f on  $\mathbb{R}^n$ , one defines its regularization (mollification) as the convolution

$$f_{\varepsilon}(x) = (f * \omega_{\varepsilon})(x) = \int \omega_{\varepsilon}(x - y) f(y) dy$$
$$= \int f(x - \varepsilon y) \omega(y) dy, \quad x \in \mathbb{R}^{n}.$$
(5.3)

It belongs to  $C^{\infty}(\mathbb{R}^n)$ , has gradient  $\nabla f_{\varepsilon} = f * \nabla \omega_{\varepsilon}$ , and is non-negative, when f is non-negative. From the definition it follows that, if  $f \in L^2(\mathbb{R}^n)$ , then

$$||f_{\varepsilon}||_{2} \le ||f||_{2}, \qquad \lim_{\varepsilon \to 0} ||f_{\varepsilon} - f||_{2} = 0.$$

Moreover, if  $f \in W_1^2(\mathbb{R}^n)$ , then, by (5.2)–(5.3), we have  $\nabla f_{\varepsilon} = \nabla f * \omega_{\varepsilon}$ . Hence

$$\|\nabla f_{\varepsilon}\|_{2} \leq \|\nabla f\|_{2}, \qquad \lim_{\varepsilon \to 0} \|\nabla f_{\varepsilon} - \nabla f\|_{2} = 0,$$

so that

$$\|f_{\varepsilon}\|_{W_{1}^{2}} \leq \|f\|_{W_{1}^{2}}, \quad \lim_{\varepsilon \to 0} \|f_{\varepsilon} - f\|_{W_{1}^{2}} = 0.$$
(5.4)

Thus,  $C^{\infty}(\mathbb{R}^n) \cap W_1^2(\mathbb{R}^n)$  is dense in  $W_1^2(\mathbb{R}^n)$ .

To obtain (a), define  $f = \sqrt{p}$ . Given  $\delta \in (0, \frac{1}{2})$ , choose  $\varepsilon > 0$  such that  $||f_{\varepsilon} - f||_{W_1^2} < \delta$ . Let us take a non-negative function  $w \in C_0^{\infty}(\mathbb{R}^n)$  with w(0) = 1 and consider a sequence

$$u_l(x) = f_{\varepsilon}(x)w(x/l).$$

These functions belong to  $C_0^{\infty}(\mathbb{R}^n)$ , and by the Lebesgue dominated convergence theorem,  $u_l \to f_{\varepsilon}$  in  $W_1^2(\mathbb{R}^n)$  as  $l \to \infty$ . Hence

$$||u - f||_{W_1^2} < \delta$$

for some  $u = u_l$ , which implies

$$| ||u||_2 - 1| = | ||u||_2 - ||f||_2| \le ||u - f||_2 < \delta$$

and thus  $||u||_2 > \frac{1}{2}$ . As a result, the normalized function  $\tilde{f} = u/||u||_2$  satisfies

$$\|\tilde{f} - f\|_{W_1^2} = \frac{\|u - \|u\|_2 f\|_{W_1^2}}{\|u\|_2} \le \frac{\delta + \delta \|f\|_{W_1^2}}{\|u\|_2} < 4\delta \|f\|_{W_1^2},$$

where we used  $||f||_{W_1^2} \ge ||f||_2 = 1$ . This gives

$$\|\nabla \tilde{f}\|_{2} - \|\nabla f\|_{2}\| < 4\delta \|f\|_{W_{1}^{2}} \le 2 \|f\|_{W_{1}^{2}}$$

and hence

$$\begin{split} \left| \|\nabla \tilde{f}\|_{2}^{2} - \|\nabla f\|_{2}^{2} \right| &\leq 4\delta \|f\|_{W_{1}^{2}} \left( \|\nabla \tilde{f}\|_{2} + \|\nabla f\|_{2} \right) \\ &\leq 4\delta \|f\|_{W_{1}^{2}} \left( 2\|f\|_{W_{1}^{2}} + 2\|\nabla f\|_{2} \right) \\ &= 8\delta \left( \|f\|_{W_{1}^{2}}^{2} + \|f\|_{W_{1}^{2}} \|\nabla f\|_{2} \right). \end{split}$$

Here  $||f||_{W_1^2}^2 = 1 + I(p)$  and

$$\|f\|_{W_1^2} \|\nabla f\|_2 \le \frac{1}{2} \|f\|_{W_1^2}^2 + \frac{1}{2} \|\nabla f\|_2^2 \le \frac{1}{2} + I(p).$$

Eventually, the probability density  $\tilde{p} = \tilde{f}^2$  satisfies

$$|I(\tilde{p}) - I(p)| \le 4\delta \,(3 + 4I(p)). \tag{5.5}$$

With  $\delta = \delta_k \to 0$ , we therefore obtain a sequence  $p_k = \tilde{p}$  such that  $I(p_k) \to I(p)$  as  $k \to \infty$ , thus proving (*a*).

Let us see that similar functions  $p_k$  may be used in (b) when

$$\int p(x)^{\alpha} dx = \int f(x)^{2\alpha} dx = \infty$$

which corresponds to the case where  $N_{\alpha}(p) = 0$  for  $\alpha > 1$  and  $N_{\alpha}(p) = \infty$  for  $0 < \alpha < 1$ . Returning to the previously defined functions  $u_l$ , we observe that  $||u_l||_{2\alpha} \rightarrow 0$ 

 $||f_{\varepsilon}||_{2\alpha}$  as  $l \to \infty$ . Hence, it is sufficient to check that  $||f_{\varepsilon}||_{2\alpha} \to ||f||_{2\alpha} = \infty$  for some sequence  $\varepsilon = \varepsilon_k \to 0$ . Indeed, since  $||f||_2 = 1$ , the function f is locally integrable, implying that  $f_{\varepsilon}(x) \to f(x)$  as  $\varepsilon \to 0$  for almost all points  $x \in \mathbb{R}$ . This follows from (5.2) and the Lebesgue differentiation theorem which yields

$$\begin{split} |f_{\varepsilon}(x) - f(x)| &\leq \int \omega_{\varepsilon}(x - y) |f(y) - f(x)| \, dy \\ &\leq \|\omega\|_{\infty} \, \varepsilon^{-n} \int_{|y - x| < \varepsilon} |f(y) - f(x)| \, dy \to 0 \quad \text{a.e.} \end{split}$$

Hence, by Fatou's lemma,  $||f||_{2\alpha} \leq \liminf_{\varepsilon \to 0} ||f_{\varepsilon}||_{2\alpha}$ , and we are done.

Now, let us turn to the basic case where  $\int p(x)^{\alpha} dx < \infty$ ,  $\alpha \in (0, \infty)$ . To prove (b), we borrow arguments from the proof of Theorem 2.3.2 in [19]. Consider a partition  $\{w_i\}_{i=0}^{\infty}$  of unity of  $\mathbb{R}^n$  subordinate to the covering  $G_i = B_n(0, i + 1) \setminus \overline{B_n(0, i - 1)}$ , in which  $B_n(0, -1) = B_n(0, 0) = \emptyset$ . Every function  $w_i$  is supposed to be in  $C_0^{\infty}(\mathbb{R}^n)$  with a support lying in  $G_i$ , to be non-negative, and all of them satisfy

$$\sum_{i=0}^{\infty} w_i(x) = 1, \quad x \in \mathbb{R}^n.$$
(5.6)

As before, let  $f = \sqrt{p}$ . Given  $0 < \delta < \frac{1}{2}$ , for each  $i \ge 0$  choose  $\varepsilon_i > 0$  small enough such that  $(w_i f)_{\varepsilon_i}$  is still supported in  $G_i$  and

$$\|(w_i f)_{\varepsilon_i} - w_i f\|_{W_1^2} < 2^{-i-1}\delta.$$
(5.7)

The latter is possible due to the property (5.3) applied to  $w_i f$ .

By the integrability assumption on p, we have  $||w_i f||_{2\alpha} < \infty$ , implying

$$\|(w_i f)_{\varepsilon} - w_i f\|_{2\alpha} \to 0 \quad \text{as } \varepsilon \to 0 \tag{5.8}$$

as long as  $2\alpha \ge 1$ . Since  $f \in L^2(\mathbb{R}^n)$ , we similarly have  $||(w_i f)_{\varepsilon} - w_i f||_2 \to 0$ . The latter implies that (5.8) holds in the case  $2\alpha < 1$  as well, since  $w_i f$  is supported on a bounded set. Therefore, in addition to (5.7), we may require that

$$\int |(w_i f)_{\varepsilon_i} - w_i f|^{2\alpha} \, dx < (2^{-i-1}\delta)^{\max(2\alpha,1)}.$$
(5.9)

Now, by (5.6),  $f(x) = \sum_{i=0}^{\infty} w_i(x) f(x)$ , where the series contains only finitely many non-zero terms. More precisely,

$$f(x) = \sum_{i=0}^{m} w_i(x) f(x), \quad |x| < m+1.$$

Similarly, for the function  $u(x) = \sum_{i=0}^{\infty} (w_i(x) f(x))_{\varepsilon_i}$ , we have

$$u(x) = \sum_{i=0}^{m} (w_i(x) f(x))_{\varepsilon_i}, \quad |x| < m+1.$$

This equality shows that u is non-negative and belongs to the class  $C_0^{\infty}(\mathbb{R}^n)$ . In addition, by (5.7),

$$\|u - f\|_{W_1^2} \le \sum_{i=0}^{\infty} \|(w_i f)_{\varepsilon_i} - w_i f\|_{W_1^2} < \delta.$$

Hence

$$\|u - f\|_2 < \delta, \tag{5.10}$$

and repeating the arguments from the previous step, we arrive at the bound (5.5) for the density  $\tilde{p} = \tilde{f}^2$  with  $\tilde{f} = u/||u||_2$ .

Next, if  $\alpha \ge \frac{1}{2}$ , by the triangle inequality in  $L^{2\alpha}$ , from (5.9) we also get  $||u - f||_{2\alpha} < \delta$ , so

$$| \| u \|_{2\alpha} - \| f \|_{2\alpha} | < \delta.$$
(5.11)

If  $\alpha < \frac{1}{2}$ , then, applying the inequality  $(a_1 + \cdots + a_N)^{2\alpha} \le a_1^{2\alpha} + \cdots + a_N^{2\alpha}$   $(a_k \ge 0)$ , from (5.9) we deduce that

$$\int |u-f|^{2\alpha} dx \leq \sum_{i=1}^{\infty} \int |u-w_i f|^{2\alpha} dx < \delta.$$

This yields

$$\left|\int u^{2\alpha}\,dx - \int f^{2\alpha}\,dx\right| < \delta$$

and therefore, by Jensen's inequality,

$$| \| u \|_{2\alpha} - \| f \|_{2\alpha} | < (2\delta)^{1/(2\alpha)}.$$
(5.12)

In view of (5.10), inequalities similar to (5.11)–(5.12) hold also true for the function  $\tilde{f} = u/||u||_2$  in place of u. Applying this with  $\delta = \delta_k \to 0$ , we obtain a sequence  $\tilde{f}_k$  such that the probability densities  $\tilde{p} = \tilde{f}^2$  satisfy (a) - (b) for any  $\alpha \neq 1$ .

**Corollary 5.2** For any  $\alpha > 0$ ,  $\alpha \neq 1$ , the infimum

$$\inf_{I(p)<\infty} \left[ N_{\alpha}(p)I(p) \right]$$

may be restricted to the class of compactly supported,  $C^{\infty}$ -smooth densities p on  $\mathbb{R}^n$  with finite Fisher information.

# 6 Two Dimensional Isoperimetric Inequalities for Entropies

In this section we deal with dimension n = 2. As will be clarified, the entropic isoperimetric inequality

$$N_{\alpha}(X)I(X) \ge c_{\alpha,2} \tag{6.1}$$

holds true for any  $\alpha \in [0, \infty)$  with a positive constant  $c_{\alpha,2}$  and does not hold for  $\alpha = \infty$  which answers Question 1 in the introduction. In addition, we will give a certain description of the optimal constants  $c_{\alpha,2}$  in (6.1) for the range  $\alpha \in [\frac{1}{2}, \infty)$ , thus answering partially Question 2.

When n = 2, the family of inequalities (1.6) takes now the form

$$\left(\int |f|^{2\alpha}\right)^{\frac{1}{2\alpha}} \le \left(\frac{4}{c_{\alpha,2}}\right)^{\frac{\alpha-1}{2\alpha}} \left(\int |\nabla f|^2\right)^{\frac{\theta}{2}} \left(\int f^2\right)^{\frac{1-\theta}{2}} \tag{6.2}$$

with  $\theta = \frac{\alpha - 1}{\alpha}$  when  $\alpha > 1$ , and

$$\left(\int f^2\right)^{\frac{1}{2}} \le \left(\frac{4}{c_{\alpha,2}}\right)^{\frac{1-\alpha}{2}} \left(\int |\nabla f|^2\right)^{\frac{\theta}{2}} \left(\int |f|^{2\alpha}\right)^{\frac{1-\theta}{2\alpha}} \tag{6.3}$$

with  $\theta = 1 - \alpha$  when  $\alpha \in (0, 1)$ .

Both inequalities enter the framework of Gagliardo–Nirenberg's inequality (1.8). The best constants and extremizers in (1.8) are not known for all admissible parameters. The most recent paper on this topic is due to Liu and Wang [10] (see references therein and historical comments). The case q = s = 2 in (1.8) that corresponds to (6.2) with  $r = 2\alpha$  goes back to Weinstein [17] who related the best constants to the solutions of non-linear Schrödinger equations.

We present now part of the results of [10] that are useful for us. Since we will use them for any dimension  $n \ge 2$ , the next statement does not deal only with the case n = 2. Also, since all the inequalities of interest for us deal with the  $L^2$ -norm of the gradient only, we may restrict ourselves to q = 2 for simplicity, when (1.8) becomes

$$\left(\int |f|^r\right)^{\frac{1}{r}} \le \kappa_n(2, r, s) \left(\int |\nabla f|^2\right)^{\frac{\theta}{2}} \left(\int |f|^s\right)^{\frac{1-\theta}{s}}$$
(6.4)

with parameters satisfying  $1 \le r, s \le \infty, 0 \le \theta \le 1$ , and  $\frac{1}{r} = \theta(\frac{1}{2} - \frac{1}{n}) + (1 - \theta)\frac{1}{s}$ . This inequality may be restricted to the class of all smooth, compactly supported functions  $f \ge 0$  on  $\mathbb{R}^n$ . Once (6.4) holds in  $C_0^{\infty}(\mathbb{R}^n)$ , this inequality is extended by a regularization and density arguments to the Sobolev space of functions  $f \in L^s(\mathbb{R}^n)$  such that  $|\nabla f| \in L^2(\mathbb{R}^n)$  (the gradients in this space are understood in a weak sense).

The next statement relates the optimal constant in (6.4) to the solutions of the ordinary non-linear equation

$$u''(t) + \frac{n-1}{t}u'(t) + u(t)^{r-1} = u(t)^{s-1}$$
(6.5)

on the positive half-axis. Put

$$\sigma = \begin{cases} \frac{n+2}{n-2} & \text{if } n \ge 3, \\ \infty & \text{if } n = 2. \end{cases}$$

We denote by |x| the Euclidean norm of a vector  $x \in \mathbb{R}^n$ .

**Theorem 6.1 ([10])** *In the range*  $1 \le s < \sigma$ ,  $s < r < \sigma + 1$ ,

$$\kappa_n(2, r, s) = \theta^{-\frac{\theta}{2}} (1 - \theta)^{\frac{\theta}{2} - \frac{1}{r}} M_s^{-\frac{\theta}{n}}, \quad M_s = \int_{\mathbb{R}^n} u_{r,s}^s(|x|) \, dx,$$

where the functions  $u_{r,s} = u_{r,s}(t)$  are defined for  $t \ge 0$  as follows.

- (i) If s < 2, then  $u_{r,s}$  is the unique positive decreasing solution to the equation (6.5) in 0 < t < T (for some T), satisfying u'(0) = 0, u(T) = u'(T) = 0, and u(t) = 0 for all  $t \ge T$ .
- (ii) If  $s \ge 2$ , then  $u_{r,s}$  is the unique positive decreasing solution to (6.5) in t > 0, satisfying u'(0) = 0 and  $\lim_{t\to\infty} u(t) = 0$ .

*Moreover, the extremizers in* (6.4) *exist and have the form*  $f(x) = au_{r,s}(|bx+c|)$  *with*  $a \in \mathbb{R}$ *,*  $b \neq 0$ *,*  $c \in \mathbb{R}^n$ .

Note that (6.2) corresponds to Gagliardo–Nirenberg's inequality (6.4) with s = 2,  $r = 2\alpha$ , and  $\theta = \frac{\alpha - 1}{\alpha}$  for  $\alpha > 1$ , while (6.3) with  $\alpha \in [\frac{1}{2}, 1)$  corresponds to (6.4) with r = 2,  $s = 2\alpha$ , and  $\theta = 1 - \alpha$ . Applying Corollary 5.2, we therefore conclude that

$$\kappa_2(2, r, s) = (4/c_{\alpha,2})^{\frac{\alpha-1}{2\alpha}} \quad \text{when } \alpha > 1,$$
  
 $\kappa_2(2, r, s) = (4/c_{\alpha,2})^{\frac{1-\alpha}{2}} \quad \text{when } \alpha \in [1/2, 1).$ 

Together with Liu–Wang's theorem, we immediately get the following corollary, where we put as before

$$M_s = \int_{\mathbb{R}^2} u^s(|x|) \, dx = 2\pi \int_0^\infty u^s(t) \, t \, dt.$$

#### **Corollary 6.2**

(i) For any  $\alpha > 1$ , we have

$$c_{\alpha,2} = 4(\alpha - 1) \,\alpha^{-\frac{1}{\alpha - 1}} M_2,$$

where  $M_2$  is defined for the unique positive decreasing solution u(t) on  $(0, \infty)$  to the equation  $u''(t) + \frac{u'(t)}{t} + u(t)^{2\alpha-1} = u(t)$  with u'(0) = 0 and  $\lim_{t\to\infty} u(t) = 0$ .

(ii) For any  $\alpha \in [\frac{1}{2}, 1)$ , we have

$$c_{\alpha,2} = 4(1-\alpha) \,\alpha^{\frac{\alpha}{1-\alpha}} M_{2\alpha}$$

where  $M_{2\alpha}$  is defined for the unique positive decreasing solution u(t) to  $u''(t) + \frac{1}{t}u'(t) + u(t) = u(t)^{2\alpha-1}$  in 0 < t < T with u'(0) = 0, u(T) = u'(T) = 0, and u(t) = 0 for all  $t \ge T$ .

In both cases the extremizers in (6.1) represent densities of the form  $p(x) = \frac{b}{M}u^2(|bx + c|), x \in \mathbb{R}^2$ , with b > 0 and  $c \in \mathbb{R}^2$ .

So far, we have seen that (6.1) holds for any  $\alpha \in [1/2, \infty)$ . Since, as observed in the introduction,  $\alpha \mapsto c_{\alpha,n}$  is non-increasing, (6.1) holds also for  $\alpha < 1/2$  and therefore for any  $\alpha \in [0, \infty)$ . Note that the case  $\alpha = 1$ , which is formally not contained in the results above, is the classical isoperimetry inequality for entropies (1.1). Let us now explain why (6.1) cannot hold for  $\alpha = \infty$ . The functional form for (6.1) should be the limit case of (6.2) as  $\alpha \to \infty$ , when it becomes

$$\|f\|_{\infty}^{2} \le D \int |\nabla f|^{2} dx \tag{6.6}$$

with  $D = 4/c_{\infty,2}$ . To see that (6.6) may not hold with any constant D, we reproduce Example 1.1.1 in [14]. Let, for  $x \in \mathbb{R}^2$ ,

$$f(x) = \begin{cases} \log |\log |x|| & \text{if } |x| \le 1/e, \\ 0 & \text{otherwise.} \end{cases}$$

Then, passing to radial coordinates, we have

$$\int |\nabla f|^2 = 2\pi \int_0^{1/e} \frac{dr}{r |\log r|^2} = 2\pi,$$

while f is not bounded. In fact, (6.6) is also violated for a sequence of smooth bounded approximations of f.

#### 7 **Isoperimetric Inequalities for Entropies in Dimension** n = 3 and Higher

One may exhibit two different behaviors between  $n = 3, 4, \text{ and } n \ge 5$  in the entropic isoperimetric inequality

$$N_{\alpha}(X)I(X) \ge c_{\alpha,n}.$$
(7.1)

Let us rewrite the inequality (1.6) separately for the three natural regions, namely as

$$\left(\int |f|^{2\alpha}\right)^{\frac{1}{2\alpha}} \le \left(\frac{4}{c_{\alpha,n}}\right)^{\frac{n(\alpha-1)}{4\alpha}} \left(\int |\nabla f|^2\right)^{\frac{\theta}{2}} \left(\int f^2\right)^{\frac{1-\theta}{2}}$$
(7.2)

with  $\theta = \frac{n(\alpha-1)}{2\alpha}$  when  $1 < \alpha \leq \frac{n}{n-2}$ ,

$$\left(\int |f|^{2\alpha}\right)^{\frac{\theta}{2\alpha}} \left(\int f^2\right)^{\frac{1-\theta}{2}} \le \frac{2}{\sqrt{c_{\alpha,n}}} \left(\int |\nabla f|^2\right)^{\frac{1}{2}}$$
(7.3)

with  $\theta = \frac{2\alpha}{n(\alpha-1)}$  when  $\alpha > \frac{n}{n-2}$  (observe that  $\theta \in (0, 1)$  in this case), and finally

$$\left(\int f^2\right)^{\frac{1}{2}} \le \left(\frac{4}{c_{\alpha,n}}\right)^{\frac{n(1-\alpha)}{2[\alpha(2-n)+n]}} \left(\int |\nabla f|^2\right)^{\frac{\theta}{2}} \left(\int |f|^{2\alpha}\right)^{\frac{1-\theta}{2\alpha}}$$
(7.4)

with  $\theta = \frac{n(1-\alpha)}{\alpha(2-n)+n}$  when  $\alpha \in (0, 1)$ . Both (7.2) and (7.4) enter the framework of Gagliardo–Nirenberg's inequality (1.8). As for (7.3), we will show that such an inequality cannot hold. To that aim, we need to introduce the limiting case  $\theta = 1$  in (7.2), which corresponds to  $\alpha = \frac{n}{n-2}$ . It amounts to the classical Sobolev inequality

$$\left(\int |f|^{\frac{2n}{n-2}}\right)^{\frac{n-2}{2n}} \le S_n \left(\int |\nabla f|^2\right)^{\frac{1}{2}}$$
(7.5)

which is known to hold true with best constant

$$S_n = \frac{1}{\sqrt{\pi n(n-2)}} \left(\frac{\Gamma(n)}{\Gamma(\frac{n}{2})}\right)^{\frac{1}{n}}.$$

Moreover, the only extremizers in (7.5) have the form

$$f(x) = \frac{a}{(1+b|x-x_0|^2)^{\frac{n-2}{2}}}, \quad a \in \mathbb{R}, \ b > 0, \ x_0 \in \mathbb{R}^n$$
(7.6)

(sometimes called the Barenblatt profile), see [1, 7, 16]. If  $f \in L^2(\mathbb{R}^n)$  and  $|\nabla f| \in L^2(\mathbb{R}^n)$ , then, by (7.3), we would have that  $f \in L^p(\mathbb{R}^n)$  with  $p = 2\alpha > \frac{2n}{n-2}$  which contradicts the Sobolev embeddings. Therefore (7.3) cannot be true, so that (7.1) holds only for  $\alpha \in [0, \frac{n}{n-2}]$ .

As for the value of the best constant  $c_{\alpha,n}$  in (7.1) and the form of the extremizers, we need to use again Theorem 6.1 which can, however, be applied only for  $n \le 5$ . As in Corollary 6.2, we adopt the notation

$$M_s = \int_{\mathbb{R}^n} u^s(|x|) \, dx$$

for a function *u* satisfying the non-linear ordinary differential equation

$$u''(t) + \frac{n-1}{t}u'(t) + u(t)^{2\alpha - 1} = u(t), \quad 0 < t < \infty,$$
(7.7)

or (in a different scenario)

$$u''(t) + \frac{n-1}{t}u'(t) + u(t) = u(t)^{2\alpha - 1}, \quad 0 < t < T.$$
(7.8)

**Corollary 7.2** *Let*  $3 \le n \le 5$ *.* 

(i) For any  $1 < \alpha < \frac{n}{n-2}$ , we have

$$c_{\alpha,n} = \frac{2n(\alpha-1)}{\alpha} \left(\frac{2\alpha}{\alpha(2-n)+n}\right)^{\frac{n(\alpha-1)-2}{n(\alpha-1)}} M_2^{\frac{2}{n}},$$

where  $M_2$  is defined for the unique positive decreasing solution u(t) to (7.7) on  $(0, \infty)$  with u'(0) = 0 and  $\lim_{t\to\infty} u(t) = 0$ .

(ii) For any  $\alpha \in [\frac{1}{2}, 1)$ ,

$$c_{\alpha,n} = 4 \frac{n(1-\alpha)}{\alpha(2-n)+n} \left(\frac{2\alpha}{\alpha(2-n)+n}\right)^{\frac{2\alpha}{n(1-\alpha)}} M_{2\alpha}^{\frac{2}{n}}$$

where  $M_{2\alpha}$  is defined for the unique positive decreasing solution u(t) to (7.8) with u'(0) = 0, u(T) = u'(T) = 0, and u(t) = 0 for all  $t \ge T$ .

In both cases, the extremizers in (7.1) are densities of the form  $p(x) = \frac{b}{M}u^2(|bx+c|), x \in \mathbb{R}^n$ , with b > 0 and  $c \in \mathbb{R}^n$ .

For the critical value of  $\alpha$ , the picture is more complete but is different.

**Corollary 7.3** Let  $n \ge 3$  and  $\alpha = \frac{n}{n-2}$ . Then

$$c_{\alpha,n} = 4\pi n(n-2) \left(\frac{\Gamma(\frac{n}{2})}{\Gamma(n)}\right)^{\frac{2}{n}}.$$

- (i) For n = 3 and n = 4, (7.1) has no extremizers, i.e., there does not exist any density p for which equality holds in (7.1) with the optimal constant.
- (ii) For  $n \ge 5$ , the extremizers in (7.1) exist and have the form

$$p(x) = \frac{a}{(1+b|x-x_0|^2)^{n-2}}, \quad a, b > 0, \ x_0 \in \mathbb{R}^n.$$
(7.9)

*Remark 7.4* Recall that  $c_{1,n} = 2\pi en$ . Using the Stirling formula, it is easy to see that, for  $\alpha = \frac{n}{n-2}$ ,

$$c_{\alpha,n} \sim 2\pi en - 2\pi e \left(2 + \log 2\right) + O\left(\frac{1}{n}\right) \text{ as } n \to \infty$$

In particular,  $c_{\alpha,n} \ge 2\pi en - c_0$  for all  $0 \le \alpha \le \frac{n}{n-2}$  with some absolute constant  $c_0 > 0$ . To get a similar upper bound, it is sufficient to test (7.1) with  $\alpha = 0$  on some specific probability distributions. In this case, this inequality becomes

$$\operatorname{vol}_{n}(\operatorname{supp}(p))^{\frac{2}{n}}I(X) \ge c_{0,n}.$$
 (7.10)

Suppose that the random vector  $X = (X_1, ..., X_n)$  in  $\mathbb{R}^n$  has independent components such that every  $X_k$  has a common density  $w(s) = \frac{2}{\pi} \cos^2(s)$ ,  $|s| \le \frac{\pi}{2}$ . As we already mentioned in Sect. 4, this one dimensional probability distribution appears as an extremal one in the entropic isoperimetric inequality (1.2) for the parameter  $\alpha = 0$ . The random vector X has density

$$p(x) = w(x_1) \dots w(x_n), \quad x = (x_1, \dots, x_n) \in \mathbb{R}^n,$$

so that

$$N_0(X) = N_0(X_1) = \pi^2$$
,  $I(X) = nI(X_1) = 4n$ .

Therefore, from (7.10) we may conclude that  $c_{0,n} \leq 4\pi^2 n$ .

*Proof of Corollaries* **7.2–7.3** The first corollary is obtained by a straight forward application of Theorem 6.1 with

$$s = 2, r = 2\alpha, \ \theta = \frac{n(\alpha - 1)}{2\alpha}, \quad \kappa_n(2, r, s) = \left(\frac{4}{c_{\alpha, n}}\right)^{\frac{n(\alpha - 1)}{4\alpha}}$$

when  $1 < \alpha < \frac{n}{n-2}$ , and with

$$q = r = 2, \ s = 2\alpha, \ \theta = \frac{n(1-\alpha)}{\alpha(2-n)+n}, \quad \kappa_n(2,r,s) = \left(\frac{4}{c_{\alpha,n}}\right)^{\frac{n(1-\alpha)}{2[\alpha(2-n)+n]}}$$

when  $\alpha \in (0, 1)$ . Details are left to the reader.

For the second corollary, we first observe that (7.2) can be recast for  $n \ge 3$  and  $\alpha = \frac{n}{n-2}$  as

$$\left(\int |f|^{\frac{2n}{n-2}}\right)^{\frac{n-2}{2n}} \le \left(\frac{4}{c_{\alpha,n}}\right)^{1/2} \left(\int |\nabla f|^2\right)^{\frac{1}{2}}.$$
(7.11)

Therefore  $\frac{4}{c_{\alpha,n}} = S_n^2$  from which the explicit value of  $c_{\alpha,n}$  follows (recalling Corollary 5.2).

Now, in order to analyze the question about the extremizers in (7.1), suppose that we have an equality in it for a fixed (probability) density p on  $\mathbb{R}^n$ . In particular, we should assume that the function  $f = \sqrt{p}$  belongs to  $W_1^2(\mathbb{R}^n)$ . Rewriting (7.1) in terms of f, we then obtain an equality in (7.11), which is the same as (7.5). As mentioned earlier, this implies that f must be of the form (7.6), thus leading to (7.9). However, whether or not this function p is integrable depends on the dimension. Using polar coordinates, one immediately realizes that

$$\int \frac{dx}{(1+b|x-x_0|^2)^{n-2}}$$

has the same behavior as  $\int_1^\infty \frac{1}{r^{n-3}} dr$ . But, the latter integral converges only if  $n \ge 5$ .

**Acknowledgments** Research of the first author was partially supported by the NSF grant DMS-2154001. Research of the second author was partially supported by the grants ANR-15-CE40-0020-03 – LSD – Large Stochastic Dynamics, ANR 11-LBX-0023-01 – Labex MME-DII and Fondation Simone et Cino del Luca in France. This research has been conducted within the FP2M federation (CNRS FR 2036).

#### References

- 1. T. Aubin, Problèmes isopérimétriques et espaces de Sobolev. J. Differ. Geom. 11, 573–598 (1976)
- W. Beckner, Asymptotic estimates for Gagliardo-Nirenberg embedding constants. Potential Anal. 17, 253–266 (2002)
- S.G. Bobkov, G.P. Chistyakov, F. Götze, Fisher information and the central limit theorem. Probab. Theory Relat. Fields 159(1–2), 1–59 (2014)
- S.G. Bobkov, N. Gozlan, C. Roberto, P.-M. Samson, Bounds on the deficit in the logarithmic Sobolev inequality. J. Funct. Anal. 267(11), 4110–4138 (2014)

- 5. E.A. Carlen, Superadditivity of Fisher's information and logarithmic Sobolev inequalities. J. Funct. Anal. **101**(1), 194–211 (1991)
- M.H.M. Costa, T.M. Cover, On the similarity of the entropy power inequality and the Brunn-Minkowski inequality. IEEE Trans. Inform. Theory 30, 837–839 (1984)
- M. Del Pino, J. Dolbeault, Best constants for Gagliardo-Nirenberg inequalities and applications to nonlinear diffusions. J. Math. Pures Appl. (9) 81(9), 847–875 (2002)
- A. Dembo, T.M. Cover, J. Thomas, Information theoretic inequalities. IEEE Trans. Inform. Theory 37(6), 1501–1518 (1991)
- 9. L. Gross, Logarithmic Sobolev inequalities. Am. J. Math. 97, 1061-1083 (1975)
- 10. J.-G. Liu, J. Wang, On the best constant for Gagliardo-Nirenberg interpolation inequalities. Preprint (2017). Available at http://arxiv.org/abs/1712.10208v1
- J. Moser, On Harnack's theorem for elliptic differential equations. Commun. Pure Appl. Math. 14, 577–591 (1961)
- J. Moser, A Harnack inequality for parabolic differential equations. Commun. Pure Appl. Math. 17, 101–134 (1964); correction in 20, 231–236 (1967)
- B. Nagy, Über integralungleichungen zwischen einer Funktion und ihrer Ableitung. Acta Univ. Szeged. Sect. Sci. Math. 10, 64–74 (1941)
- L. Saloff-Coste, Aspects of Sobolev-Type Inequalities, vol. 289 (Cambridge University Press, Cambridge/New York, 2002)
- A.J. Stam, Some inequalities satisfied by the quantities of information of Fisher and Shannon. Inform. Control 2, 101–112 (1959)
- 16. G. Talenti, Best constant in Sobolev inequality. Ann. Mat. Pura Appl. 110, 353-372 (1976)
- M.I. Weinstein, Nonlinear Schrödinger equations and sharp interpolation estimates. Commun. Math. Phys. 87, 567–576 (1983)
- Z. Zhang, Inequalities for characteristic functions involving Fisher information. C. R. Math. Acad. Sci. Paris 344(5), 327–330 (2007)
- W.P. Ziemer, Weakly Differentiable Functions. Sobolev Spaces and Functions of Bounded Variation. Graduate Texts in Mathematics, vol. 120 (Springer, New York, 1989), xvi+308pp