



A General Method of Finding New Symplectic Schemes for Hamiltonian Mechanics

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Abstract. The explicit symplectic difference schemes are considered for the numerical solution of molecular dynamics problems described by systems with separable Hamiltonians. A general method for finding symplectic schemes of high order of accuracy using parametric Gröbner bases, resultants, and permutations of variables is proposed. The implementation of the method is described by the example of four-stage partitioned Runge–Kutta (PRK) schemes of the Forest–Ruth family. All required symbolic calculations are performed using the computer algebra system *Mathematica*. 96 new PRK schemes of Forest–Ruth family have been obtained.

Keywords: Molecular dynamics · Symplectic difference schemes · Partitioned schemes · Gröbner bases · Resultants · Permutations

1 Introduction

The equations of molecular dynamics (MD) are ordinary Hamilton differential equations describing the interaction of material particles. MD equations have an exact analytical solution in a very limited number of cases [8]. Therefore, in the general case, these equations are solved numerically using difference schemes in which the differential operator is replaced by a difference operator.

When solving the Hamilton equations, it is natural to use difference schemes that preserve the symplectic properties of these equations. Violation of this condition leads to non-conservation of Poincaré invariants and the appearance of non-physical instability in numerical calculations. It follows that the difference operator of the numerical scheme must have the properties of a canonical transformation. The corresponding difference schemes are called symplectic.

This work is a continuation of the works of the authors [15, 16], which provide a detailed overview of the methods for constructing symplectic schemes. In [15, 16], symplectic schemes were considered in relation to the MD method, in which the Hamiltonian of interacting particles splits into the sum of kinetic

and potential energy. It was shown that the problem of constructing symplectic schemes of a given order of accuracy with a fixed number of stages is reduced to the problem of finding the roots of a polynomial system of equations that arise when the coefficients in the expression for the error of the scheme are turned to zero. It turned out that the use of the Gröbner basis technique [1] implemented in the *Mathematica* computer algebra system makes it possible to find the roots of a system of polynomial equations in the case when its Vandermonde determinant is zero. With the help of this technique, 21 new four-stage schemes of the 4th order of accuracy were obtained [16].

This paper presents a generalization of the method proposed in [15, 16] to the more general case when the Vandermonde determinant of a system of polynomial equations is nonzero. The effectiveness of the proposed method is demonstrated by the example of a four-stage scheme of the 4th order of accuracy. 96 new symplectic four-stage schemes of the 4th order of accuracy of the Forest–Ruth type were obtained. The error functional for several of the best new schemes is two orders of magnitude smaller than for the Forest–Ruth scheme [6].

Verification of schemes is carried out by comparing the numerical solutions with the exact solution of the test problem. It is shown that the new Forest–Ruth four-stage schemes with the smallest norm of the leading error term provide a more accurate preservation of a balance in total energy of the particle system than the schemes of the same formal approximation order with a larger approximation error.

2 Governing Equations

For simplicity, consider the Hamilton's equations for a single particle in the one-dimensional case:

$$dx/dt = p/m, \quad dp/dt = f(x), \quad f(x) = -\partial H(x)/\partial x, \quad H = \frac{p^2}{2m} + V(x), \quad (1)$$

where t is the time, x is the coordinate, p is the momentum, m is the particle mass, f is the force acting on the particle, $V(x)$ is the potential energy, $p^2/(2m)$ is the kinetic energy, H is the Hamiltonian.

Hamilton's equations (1) generate a one-parameter group of diffeomorphisms $x^0 \rightarrow x(t, x^0)$ [5]. The smooth functions $z(x, p)$ form the algebra of the Lie group with respect to the Poisson bracket [5]

$$\frac{dz}{dt} = \{z, H\}, \quad \{z, H\} = \left(\frac{\partial H}{\partial p} \frac{\partial z}{\partial x} - \frac{\partial H}{\partial x} \frac{\partial z}{\partial p} \right). \quad (2)$$

Let us rewrite these equations in the operator form

$$\frac{dz}{dt} = \hat{L}z, \quad \hat{L} = \left(\frac{\partial H}{\partial p} \frac{\partial}{\partial x} - \frac{\partial H}{\partial x} \frac{\partial}{\partial p} \right), \quad (3)$$

where \hat{L} is the Liouville operator. Write the solution of Eq. (3) in the form

$$z(t) = \hat{U}z(0), \quad \hat{U} = \exp(\hat{L}t), \quad \hat{L} = \hat{L}_1 + \hat{L}_2, \quad \hat{L}_1 = \frac{\partial H}{\partial p} \frac{\partial}{\partial x}, \quad \hat{L}_2 = -\frac{\partial H}{\partial x} \frac{\partial}{\partial p}. \quad (4)$$

Propagators \hat{U} form a one-parameter Lie group and satisfy the group relations

$$\hat{U}(t_1)\hat{U}(t_2) = \hat{U}(t_1 + t_2), \quad \hat{U}^{-1} = \exp(-\hat{L}t), \quad \hat{U}(0) = I, \quad \hat{U}\hat{U}^{-1} = I. \quad (5)$$

Equations (4) underlie the propagator method of constructing symplectic schemes [3, 7, 9, 18].

3 Symplectic Partitioned Runge–Kutta Schemes

PRK schemes are based on the propagator method, which is determined by relations (4). In the case of a separable Hamiltonian, they have the following form:

$$\begin{aligned} z(t) &= \hat{U}z(0), \quad \hat{U} = \exp[(\hat{L}_1 + \hat{L}_2)t], \\ \hat{L}_1 &= \frac{p}{m} \frac{\partial}{\partial x}, \quad \hat{L}_2 = f(x) \frac{\partial}{\partial p}, \quad f(x) = -\frac{\partial V}{\partial x}. \end{aligned} \quad (6)$$

If the values $z^n = (x^n, p^n)$ are known at the n th time layer t^n , then we obtain at the $(n + 1)$ th time layer $t^{n+1} = t^n + h$

$$z^{n+1} = \exp((\hat{L}_1 + \hat{L}_2)h)z^n. \quad (7)$$

Let us subdivide the passage from t^n to t^{n+1} into K stages. Let us approximate the propagator (7) by a product of propagators with the accuracy $O(h^{\lambda+1})$:

$$\exp[(\hat{L}_1 + \hat{L}_2)h] = \prod_{s=1}^K \exp(d_s h \hat{L}_1) \exp(c_s h \hat{L}_2) + O(h^{\lambda+1}), \quad \sum_{s=1}^K c_s = 1, \quad \sum_{s=1}^K d_s = 1. \quad (8)$$

Thus, the solution at the $(n + 1)$ th time layer is given with the accuracy of $O(h^\lambda)$ by the equation

$$\begin{aligned} z^{n+1} &= \prod_{s=1}^K \exp(d_s h \hat{L}_1) \exp(c_s h \hat{L}_2) z^n, \\ z^K &= z^{n+1} = (x^{n+1}, p^{n+1}), \quad z^{(0)} = z^n = (x^n, p^n). \end{aligned} \quad (9)$$

We obtain from Eqs. (7) and (9) that at the passage from the $(i - 1)$ th stage to the i th stage, the following relations are valid [6]:

$$p^{(i)} = p^{(i-1)} + c_i h f(x^{(i-1)}), \quad x^{(i)} = x^{(i-1)} + d_i \frac{h}{m} p^{(i)}, \quad i = 1, \dots, K, \quad (10)$$

which satisfy the condition $J = \frac{\partial(p^i, x^i)}{\partial(p^{i-1}, x^{i-1})} = 1$. Note that one can use instead of (9) the expansion (in the notation of the work [18]) of the form

$$z^{n+1} = \prod_{s=1}^K \exp(d_s h \hat{L}_2) \exp(c_s h \hat{L}_1) z^n, \quad (11)$$

which leads to the relations [18]

$$x^{(i)} = x^{(i-1)} + c_i \frac{h}{m} p^{(i-1)}, \quad p^{(i)} = p^{(i-1)} + d_i h f(x^{(i)}), \quad i = 1, \dots, K, \quad (12)$$

which also satisfy the condition $J = 1$. Equations (10) and (12) define the class of explicit symplectic schemes the accuracy of which is determined by the number of stages K and coefficients c_i, d_i .

Below we show how new sets of parameters of a four-stage multiparameter family of Forest–Ruth symplectic schemes (10) can be obtained analytically using Gröbner bases. Some of these schemes have a much smaller norm of the leading error term than the two four-stage schemes obtained in [6].

4 Forest–Ruth Scheme

Consider scheme (10) with $K = 4$. Using symbolic calculations, it is not difficult to obtain an expression for the error δp_n in the form [16]

$$\begin{aligned} \delta p_n = & hP_1 f(x) + (h^2/2)P_2 u(t)f'(x) + h^3[P_{31}f(x)f'(x)/m + P_{32}u^2 f''(x)]/6 \\ & + (h^4 u)/(24m)\{P_{41} \cdot [f'(x)]^2 + 3P_{42}f(x)f''(x) + P_{43}mu^2 f^{(3)}(x)\} \\ & - [h^5/(120m^2)](3P_{51}f^2(x)f''(x) + f(x)(P_{52} \cdot [f'(x)]^2 - 6P_{53}mu^2 f^{(3)}(x) \\ & - mu^2(5P_{54}f'(x)f''(x) + P_{55}mu^2 f^{(4)}(x))), \end{aligned} \quad (13)$$

where $u(t)$ is the particle velocity. The expressions for $P_1, P_2, P_{31}, P_{32}, P_{41}, P_{42}, P_{43}, P_{51}, P_{52}, P_{53}, P_{54}$, and P_{55} are available in [16], and we give here only the expression for P_1 for the sake of brevity: $P_1 = 1 - c_1 - c_2 - c_3 - c_4$.

Let us calculate the weighted root-mean-square value of five polynomials P_{5j} , $j = 1, \dots, 5$:

$$P_{5,rms}^{(l)} = [(1/5) \sum_{j=1}^5 (\sigma_j P_{5j})^2]^{1/2}. \quad (14)$$

Here $\sigma_1, \dots, \sigma_5$ are problem-independent factors affecting the polynomials P_{5j} in (13), $\sigma_1 = -3, \sigma_2 = -1, \sigma_3 = 6, \sigma_4 = 5, \sigma_5 = 1$. The error δx_n obtained when calculating the coordinate x^{n+1} according to the scheme under consideration is given by the formula

$$\begin{aligned} \delta x_n = & \frac{hp}{m} R_1 + \frac{h^2}{2m} R_2 f(x) + \frac{h^3}{6m} R_3 u(t)f'(x) + \frac{h^4}{24m^2} [R_{41}f(x)f'(x) \\ & + R_{42}mu^2 f''(x)] \frac{h^5 u(t)}{120m^2} \{R_{51}[f'(x)]^2 + 3R_{52}f(x)f''(x) + R_{53}mu^2 f^{(3)}(x)\}. \end{aligned} \quad (15)$$

The expressions for $R_1, R_2, R_3, R_{41}, R_{42}, R_{51}, R_{52}$, and R_{53} are available in [16], and we give here only the expression for R_1 for the sake of brevity: $R_1 = 1 - d_1 - d_2 - d_3 - d_4$.

The system of equations $P_1 = 0, P_2 = 0, P_{32} = 0, P_{43} = 0$ is linear in c_i , $i = 1, \dots, 4$. It is easy to rewrite it in the form $\mathcal{P}\mathbf{c} = \mathbf{b}$, where $\mathbf{c} = (c_1, c_2, c_3, c_4)^T$,

Table 1. The values of the parameters of the Forest–Ruth schemes FRl ($l = 1, \dots, 6$) obtained in the case of the zero determinant (17)

FRl		c_i	d_i
FR1	$d_2 = -d_1$	1.681562217889354644	1.689032314592042095
		0.014940193405374902	-1.689032314592042095
		-1.471849481432093948	0.612326464931317799
FR2	$d_2 = -d_1$	0.462663328009943900	0.754092747422918106
		0.582858838825948413	-0.754092747422918106
		-1.983103761290466971	0.031209202772712671
FR3	$d_3 = -d_2$	0	φ_1^+
		φ_2^-	$-z/3$
		1/2	$z/3$
FR4	$d_3 = -d_2$	0	φ_1^-
		φ_2^+	$z/3$
		1/2	$-z/3$
FR5	$d_4 = 1$	-0.298786423182561709	0.092128124809906823
		0.781829095984937063	0.671234501323664155
		0.560639906662391246	-0.763362626133570977
FR6	$d_4 = 1$	0.447302683845444312	0.747066325751166375
		0.599527283811444124	-0.771392086019630084
		-2.142311455850704292	0.024325760268463710

Table 2. The values of the parameters of the Forest–Ruth schemes FRl, $l = 7, \dots, 10$ in the particular cases of $c_1 = 0$ and $d_4 = 0$

FRl	c_1	c_2	c_3	c_4	d_1	d_2	d_3	d_4	
FR7	$c_1 = 0$	0	φ_2^+	$\frac{1}{2}$	φ_2^-	φ_1^-	$\frac{z}{3}$	$-\frac{z}{3}$	φ_1^+
FR8	$c_1 = 0$	0	φ_2^-	$\frac{1}{2}$	φ_2^+	φ_1^+	$-\frac{z}{3}$	$\frac{z}{3}$	φ_1^-
FR9*	$c_1 = 0$	0	$2\varphi_3$	φ_4	$2\varphi_3$	φ_3	$\frac{1}{2} - \varphi_3$	$\frac{1}{2} - \varphi_3$	φ_3
FR10*	$d_4 = 0$	φ_3	$\frac{1}{2} - \varphi_3$	$\frac{1}{2} - \varphi_3$	φ_3	$2\varphi_3$	φ_4	$2\varphi_3$	0

$\mathbf{b} = (1, \frac{1}{2}, \frac{1}{3}, \frac{1}{4})^T$, and

$$\mathcal{P} = \begin{pmatrix} 1 & 1 & 1 & 1 \\ 0 & d_1 & d_1 + d_2 & d_1 + d_2 + d_3 \\ 0 & d_1^2 & (d_1 + d_2)^2 & (d_1 + d_2 + d_3)^2 \\ 0 & d_1^3 & (d_1 + d_2)^3 & (d_1 + d_2 + d_3)^3 \end{pmatrix}. \tag{16}$$

The determinant of this matrix is as follows:

$$\det \mathcal{P} = d_1 d_2 d_3 (d_1 + d_2) (d_2 + d_3) (1 - d_4). \tag{17}$$

The values c_i obtained from the system $P_1 = 0, P_2 = 0, P_{32} = 0, P_{43} = 0$ have the form of fractional rational functions whose numerators are polynomials in d_1, \dots, d_4 , and denominator is the determinant (17). If $\det \mathcal{P} = 0$, then it will be problematic to find expressions for c_i . In this connection, all six particular cases of the vanishing determinant (17) were considered in detail in [16]. The following polynomial system was handled therein:

$$R_1 = 0, P_1 = 0, P_2 = 0, P_{31} = 0, P_{32} = 0, P_{41} = 0, P_{42} = 0, P_{43} = 0. \quad (18)$$

The number of equations in (18) is equal to the number of unknowns $c_1, c_2, c_3, c_4, d_1, d_2, d_3, d_4$. We found with the aid of Gröbner bases six new analytical solutions of the above system, which are summarized in Table 1. The values $z, \zeta, \varphi_1, \varphi_2^\pm, \varphi_3$ were calculated as follows:

$$\begin{aligned} z &= \sqrt{3}, \quad \zeta = 2^{1/3}, \quad \varphi_1^\pm = \frac{1}{6}(3 \pm z), \quad \varphi_2^\pm = \frac{1}{12}(3 \pm 2z), \\ \varphi_3 &= \frac{1}{12}(4 + 2\zeta + \zeta^2), \quad \varphi_4 = -\frac{1}{3}(1 + \zeta)^2. \end{aligned} \quad (19)$$

In Table 1 and in further similar tables, we give for brevity only the first three values of the parameters $c_i, d_i, i = 1, 2, 3$. The numerical values of the parameters c_4 and d_4 can be calculated using the formulas: $c_4 = 1 - c_1 - c_2 - c_3, d_4 = 1 - d_1 - d_2 - d_3$.

Earlier in the work [6], cases were considered when either $c_1 = 0$ or $d_4 = 0$; at the same time, additional symmetry conditions of the scheme were imposed: $d_2 = d_3, d_1 = d_4$ (at $c_1 = 0$) and $c_2 = c_3, c_1 = c_4$ (at $d_4 = 0$). We also considered in [16] special cases when either $c_1 = 0$ or $d_4 = 0$, but did not impose symmetry conditions. At $c_1 = 0$, two solutions coincide with those obtained for $d_3 = -d_2$, and the third solution coincides with the one given in [6]. For $d_4 = 0$, a unique real solution was obtained that coincides with the solution (4.9) from [6]. A more detailed description of these solutions is given in [16], see also Table 2, where symmetric schemes are marked with an asterisk.

Along with (14), we also introduce the functional $X_{5,rms}^{(l)}$, which represents the weighted root-mean square value of three polynomials $R_{5j}, j = 1, 2, 3$:

$$X_{5,rms}^{(l)} = \{ [R_{51}^2 + (3R_{52})^2 + R_{53}^2] / 3 \}^{1/2}. \quad (20)$$

This formula accounts for multiplier “3” before R_{52} in (15). In the following, we will not give the magnitudes of the functionals $P_{5,rms}^{(l)}$ and $X_{5,rms}^{(l)}$ separately for each scheme, but we will give them in a summary Table 19.

One can see from Table 19 that in the cases of symmetric schemes FR9 and FR10, the value $P_{5,rms}^{(l)}$ obtained from the leading term of the scheme error is greater than in the cases of nonsymmetric schemes. Thus, although symmetric schemes allow calculations in the direction of decreasing time ($h < 0$), they are less accurate than nonsymmetric schemes (according to the latter schemes, problems can be solved only at $h > 0$, i.e., moving in the direction of increasing time).

4.1 The General Case

In general, when the determinant (17) is nonzero, expressions for polynomials (18) become more complicated. The *Mathematica* command

$$\text{GroebnerBasis}[\{\text{P1, R1, P2, P31, P32, P41, P42, P43}\}, \{\text{c1, c2, c3, c4, d1, d2, d3, d4}\}] \tag{21}$$

has enabled the obtaining of the Gröbner basis $\{G_1, \dots, G_{19}\}$, which consists of 19 polynomials. The polynomial system is termed zero-dimensional if it has a finite number of complex solutions. A zero-dimensional system with as many equations as variables is sometimes said to be well-behaved [10]. The Bézout theorem asserts that a well-behaved system of n equations, which have degrees d_1, \dots, d_n , has at most $d_1 \cdots d_n$ solutions. In our case, the polynomials $P_1, R_1, P_2, P_{31}, P_{32}, P_{41}, P_{42}$, and P_{43} have according to [16] the degrees 1, 1, 2, 3, 3, 4, 4, 4, respectively. Thus, system (18) has at most 1152 complex and real roots counted with their multiplicities.

Unfortunately, in the Gröbner basis obtained from (21), there is no univariate polynomial. The shortest polynomials in the resulting basis are the polynomials G_1 and G_2 , moreover, $G_1 = \mathcal{G}_1(d_3, d_4)$, $G_2 = \mathcal{G}_2(d_2, d_3, d_4)$. In order for the system $G_1 = 0$, $G_2 = 0$ to have a common root, it is necessary that its resultant, based on the Sylvester matrix, be zero [2, 4]. The polynomials G_1 and G_2 depend collectively on three variables d_2, d_3 , and d_4 . It follows that one needs to consider three options for excluding variables: excluding the variable d_2 , excluding the variable d_3 , and excluding the variable d_4 . Therefore, we have considered three resultants: $\text{Res}(G_1, G_2, d_4)$, $\text{Res}(G_1, G_2, d_3)$, $\text{Res}(G_1, G_2, d_2)$.

In the CAS *Mathematica*, the symbolic computation of the resultant is implemented in the function `Resultant[...]`. The command `Resultant[G1, G2, d4]` eliminates the variable d_4 . As a result, it has been found that

$$\text{Res}(G_1, G_2, d_4) = A_{54}d_3\mathcal{F}_4(d_3)\mathcal{F}_8(d_3)\mathcal{F}_{13}(d_2, d_3)[\mathcal{F}_{14}(d_3)]^2\mathcal{F}_{56}(d_3). \tag{22}$$

Here A_{54} is a constant consisting of 54 decimal digits, $\mathcal{F}_4(d_3) = -1 - 16d_3 - 96d_3^2 - 36d_3^3 + 324d_3^4$. The value $d_3 = 0$ is one of the solutions of the Eq. (22). However, it was mentioned above that with this value of the parameter d_3 , the system (18) is incompatible. The solution of the equation $\mathcal{F}_4(d_3) = 0$ is in radicals, but its substitution into system (18) leads to giant symbolic expressions.

On the other hand, the roots of this equation can be found numerically with very high accuracy. But all the same, the real roots found are approximate. Substituting the root d_3 , given as a floating-point machine number, into the original polynomial system (18) turns this system into a polynomial system with approximate coefficients. In the review paper [11], problems related to obtaining approximate Gröbner bases from polynomial systems with approximate coefficients were discussed. Here we note the following two problems.

- In principle, approximate solutions of the original polynomial system can be obtained from the approximate Gröbner basis. But substituting the found

solution into the original system leads to the residuals that are several decimal orders of magnitude greater than machine rounding errors; this was shown in a number of examples in [11].

- There exists the problem of ill-conditioned polynomials [12, 17]. The presence of errors in the coefficients of a poorly conditioned polynomial can even lead to a change in the number of real roots of the polynomial equation [17].

In cases where it is not possible to find solutions to a polynomial system in an analytical form, it is more reliable to search for these solutions by numerically solving an optimization problem formulated directly for the polynomials included in the original system. This methodology was described in [15, 16].

Table 3. The values of the parameters of the Forest–Ruth schemes FR l ($l = 11, \dots, 14$) obtained in the case of the zero resultant (23)

FR l	d_4	c_i	d_i
FR11	$d_4 = 1/3$	-3.793578045247588829	-2.015338965874309019
		-0.002556494312384837	2.010225977249539346
		4.023008448811463529	0.671779655291436340
FR12	$d_4 = 1/3$	0.215291546204744457	1.145093955234638990
		0.524182325872439832	-0.096729303489759326
		-0.717640932851958484	-0.381697985078212997
FR13	$d_4 = 1/4$	-0.456387075637887056	0.559356790721160223
		-0.065655505000061227	-0.478888139152354456
		0.903328999195627837	0.669531348431194233
FR14	$d_4 = 1/4$	0.313736873232085596	0.817539342386075480
		-0.440874365834237127	-1.966375102129151814
		-0.007941192169999430	1.898835759743076334

For the resultant $\text{Res}(G_1, G_2, d_3)$, the following expression is obtained:

$$\text{Res}(G_1, G_2, d_3) = 103129560704(3d_4 - 1)^9(4d_4 - 1)^5(1 - 9d_4 + 24d_4^2)^9 \times \mathcal{F}_{12}(d_2, d_4)[\mathcal{F}_{14}(d_4)]^6. \tag{23}$$

Expressions for polynomials $\mathcal{F}_{12}(d_2, d_4)$ and $\mathcal{F}_{14}(d_4)$ are not given here because of their bulky appearance.

It follows from (23) that the value $d_4 = 1/3$ is one of the roots of the equation $\text{Res}(G_1, G_2, d_3) = 0$. Substituting this value into the polynomial system (18) significantly simplifies the appearance of the polynomials included in the system. This makes it easy to find the corresponding Gröbner basis consisting of seven polynomials G_1, \dots, G_7 . The Gröbner polynomial G_1 yields the following quartic equation for c_1 : $G(1) = -25 + 340c_1 - 1920c_1^2 + 3840c_1^3 + 1152c_1^4 = 0$. Its roots are expressed in radicals, we omit them for brevity. Two of these roots are real. Since the Gröbner basis is triangular, the parameters $d_3, d_2, d_1, c_4, c_3, c_2$ are easily found. The corresponding schemes in Table 3 are the schemes FR11 and

FR12. The case of $d_4 = 1/4$ was handled similarly to the foregoing case. The corresponding schemes in Table 3 are the schemes FR13 and FR14. The roots of the equation $1 - 9d_4 + 24d_4^2 = 0$ are complex.

And finally, $\text{Res}(G_1, G_2, d_2) = 1$. It is a polynomial of degree zero, that is, it does not contain any of the variables d_2, d_3, d_4 . Therefore, this resultant is useless when searching for solutions for these three variables.

Below we will show that with the help of parametric Gröbner bases, it is possible to obtain several dozen of further exact solutions of polynomial systems arising from the study of PRK schemes. At the same time, we can declare any of the eight variables $c_j, d_j, j = 1, \dots, 4$ as a parameter. Following [10], we will call the quantities that are not included in the set of parameters, the indeterminates.

Our goal is to search for schemes in the Forest and Ruth family for which the functionals $P_{5,\text{rms}}$ and $X_{5,\text{rms}}$ are smaller than in the cases of symplectic schemes described in [15, 16]. Taking into account the previous considerations for working with parametric Gröbner bases, we have developed and implemented a procedure consisting of seven steps in the CAS *Mathematica* language.

Step 1. For certainty, consider the case when the variable c_4 is declared as a parameter. Then the Gröbner basis of system (18) is computed using the call

$$\text{GroebnerBasis}[\{\text{P1, R1, P2, P31, P32, P41, P42, P43}\}, \quad (24)$$

$$\{\text{c1, c2, c3, d1, d2, d3, d4}\}]$$

By default, in all subsequent calls to the function `GroebnerBasis[...]`, lexicographic ordering of monomials is used. We will show in the following that the Gröbner basis of the same polynomial system also depends on the order of the indeterminates in the call to the *Mathematica* function `GroebnerBasis[...]`.

Step 2. The Gröbner basis $\{G_1, G_2, \dots, G_N\}$ ($N > 2$) obtained with the aid of the call (24) has no univariate polynomial. The first two polynomials G_1 and G_2 have the smallest length. We introduce for further convenience the notation: $\alpha_j = c_j, \alpha_{j+4} = d_j, j = 1, \dots, 4$. As a rule, $G_1 = \mathcal{F}_1(\alpha_{i_1}, \alpha_{i_2}), G_2 = \mathcal{F}_2(\alpha_{i_1}, \alpha_{i_2}, \alpha_{i_3})$, where $1 \leq i_1, i_2, i_3 \leq 8, i_1 \neq i_2, i_2 \neq i_3, i_1 \neq i_3$. Then we calculate the following three resultants:

$$\text{Res}(G_1, G_2, \alpha_{i_k}), \quad k = 1, 2, 3. \quad (25)$$

The elimination of the variable α_{i_k} from the basis polynomials G_1 and G_2 is carried out with the aid of the call `Res[G1, G2, α_{i_k}]`.

Step 3. All three resultants (25) are considered. First, using the *Mathematica* function `Factor[...]` we find out whether a particular resultant from the set (25) is factorizable. Some of the resultants (25) are irreducible over the field \mathbb{Q} of rational numbers. Below we will briefly write “irreducible polynomials”.

Step 4. All integer and rational roots of the equations $\text{Res}(G_1, G_2, \alpha_{i_k}) = 0, k = 1, 2, 3$ are found.

Step 5. Each of the found roots $\alpha_j \in \mathbb{Q}$ is substituted into system (18) and then the corresponding Gröbner basis is found. As a rule, this basis turns out to be triangular. In this case, the polynomial G_1 is univariate. If its degree is 2,

3 or 4, then the roots of the equation $G_1 = 0$ are in radicals. In order to avoid the appearance of giant symbolic expressions when substituting these closed-form roots into the remaining basic polynomials, the roots are translated into machine numbers with a mantissa length of 50 decimal digits $0, 1, \dots, 9$. To do this, we use the *Mathematica* function `SetAccuracy[root, 50]`.

If the degree of G_1 is higher than 4, then the solution of the equation $G_1 = 0$ is found with high accuracy using the command `NSolve[G[[1]]==0, c4, WorkingPrecision-> 50]`. Due to the fact that the Gröbner basis is triangular, all the required unknowns $\alpha_j, j = 1, \dots, 8$ are then easily found.

Step 6. Among the roots of each of the three resultants (25), there may be extraneous roots [13]. Therefore, verification of all found roots of system (18) is carried out by substituting them into this system.

Step 7. The functionals $P_{5,rms}^{(l)}$ and $X_{5,rms}^{(l)}$ are calculated according to (14) and (20) as criteria for evaluating the accuracy of the obtained real symplectic schemes.

The c_1 Parameter. In this case, (21) is replaced with the call

$$\text{GroebnerBasis}[\{P1, R1, P2, P31, P32, P41, P42, P43\}, \{c2, c3, c4, d1, d2, d3, d4\}] \tag{26}$$

The Gröbner basis $\{G_1, G_2, \dots\}$ is obtained as 18 polynomials, with $G_1 = \mathcal{F}_1(c_1, d_4)$ and $G_2 = \mathcal{F}_2(c_1, d_3, d_4)$. Therefore, we have considered three resultants: $\text{Res}(G_1, G_2, d_4)$, $\text{Res}(G_1, G_2, c_1)$, and $\text{Res}(G_1, G_2, d_3)$. The first of them is

$$\text{Res}(G_1, G_2, d_4) = A(c_1 - 1)^{10}(3c_1 - 1)^{17}(4c_1 - 1)^6(24c_1^2 - 9c_1 + 1)^{18}[\mathcal{F}_{12}(c_1)]^7 \mathcal{F}_{16}(c_1, d_3). \tag{27}$$

Here $A = 8077716527296737705984$. The system of equations $G_1 = 0, G_2 = 0$ has a common root if and only if the resultant $\text{Res}(G_1, G_2, d_4) = 0$. It follows from (27) that this equality holds, in particular, in the following cases: $c_1 = 1, c_1 = 1/3$, and $c_1 = 1/4$. Below we will look at all these cases.

Substituting the value $c_1 = 1$ into the polynomial system (18) makes it easy to find the corresponding Gröbner basis $\{G_1, \dots, G_7\}$, which proves to be triangular. It is then easy to find numerical values of all parameters from it. The equation $G_1 = 0$ has three real roots and two complex conjugate roots. The schemes FR15, FR16, and FR17 in Table 4 correspond to three real roots.

In the case when $c_1 = 1/3$, a triangular Gröbner basis is also obtained. The Gröbner polynomial G_1 is univariate and it yields the following quartic equation for d_4 : $960d_4^4 - 3552d_4^3 + 1778d_4^2 - 340d_4 + 25 = 0$. Its roots are expressed in radicals, we omit them for brevity. Two of these roots are real. The corresponding schemes in Table 4 are the schemes FR18 and FR19 (see also [16]). The case of $c_1 = 1/4$ is handled similarly to the case of $c_1 = 1/3$, the corresponding schemes in Table 4 are the schemes FR20 and FR21 (see also [16]).

The equation $24c_1^2 - 9c_1 + 1 = 0$ has no real roots. The equation $\mathcal{F}_{12}(c_1) = 0$ has two real roots and five pairs of complex conjugate roots. At the same time, there

Table 4. The values of the parameters of the Forest–Ruth schemes FR*l* ($l = 15, \dots, 23$) obtained in the case of the zero resultants (27) and (30)

FR <i>l</i>	c_1	c_i	d_i
FR15	$c_1 = 1$	1.000000000000000000	-0.289496692562112655
		0.480967226556355638	0.095557008516098051
		-1.051024192239418045	0.957728476679205942
FR16	$c_1 = 1$	1.000000000000000000	1.166062538245970983
		0.094779042188058442	-1.130360771387392979
		-0.969092132852183332	0.449341881021738872
FR17	$c_1 = 1$	1.000000000000000000	-0.418994738647895922
		-0.517065975450449051	1.310130928129584835
		0.421590609795226645	-1.858323611577820880
FR18	$c_1 = 1/3$	1/3	0.854618769602128809
		-0.340323030871210685	-1.755412024452683413
		-0.013979395075754704	1.663063053955831823
FR19	$c_1 = 1/3$	1/3	0.717790337479662637
		0.676663198257488871	-2.867617920531312327
		2.019993063181644409	0.003543429387338222
FR20	$c_1 = 1/4$	1/4	1.053722261159673065
		-0.120984256091361380	-0.341956531133635430
		0.867350020601588225	2.076656572336811881
FR21	$c_1 = 1/4$	1/4	1.627745859492881698
		0.381426792344745990	-0.082614832312547523
		-0.458605339881521336	-0.834595847450495149
FR22	$c_1 = 1/6$	1/6	0.842603502826414075
		0.947753813816637058	-0.103841775357201984
		-1.284460514173413455	-0.182943493689021194
FR23	$c_1 = 1/6$	1/6	0.460642622151786919
		0.660127726674692834	0.491578202509501365
		0.215313914255060158	-0.735889167477648790

is no rational root among the real roots. Obtaining an approximate Gröbner basis leads to the difficulties discussed above at the beginning of this subsection.

In an effort to get several further roots of the equation $\text{Res}(G_1, G_2, d_4) = 0$, we substitute the value of $c_1 = 1$ into the polynomial $\mathcal{F}_{16}(c_1, d_3)$. As a result, we obtain the equation $\mathcal{F}_{16}(1, d_3) = -36(4418 - 34584d_3 + 93477d_3^2 - 87708d_3^3 - 11520d_3^4 + 36864d_3^5) = 0$. This equation has three real roots and two complex conjugate roots. The real roots are as follows: -1.858323611577820880, 0.4493418810217388717, and 0.9577284766792059415. It is easy to see that these values of the parameter d_3 coincide with those given in Table 4 for schemes FR17, FR16, and FR15. The substitution of the value $c_1 = 1/3$ in $\mathcal{F}_{16}(c_1, d_3)$ has enabled us to obtain an equation $\mathcal{F}_{16}(1/3, d_3) = 0$ whose roots give the schemes FR18 and FR19. Finally, the equation $\mathcal{F}_{16}(1/4, d_3) = 0$ gives as the output the familiar schemes FR20 and FR21. Thus, the equation in question did not add new roots.

The resultant $\text{Res}(G_1, G_2, c_1) = 0$ is as follows:

$$\begin{aligned} \text{Res}(G_1, G_2, c_1) &= A_{50}(d_4 - 1)^2(3d_4 - 1)(4d_4 - 1)(24d_4^2 - 9d_4 + 1)^3 \\ &\times \mathcal{F}_3(d_4)\mathcal{F}_4(d_4)\mathcal{F}_6(d_4)[\mathcal{F}_8(d_4)]^2\mathcal{F}_{10}(d_3, d_4)[\mathcal{F}_{12}(d_4)]^2\mathcal{F}_{60}(d_4), \end{aligned} \tag{28}$$

where A_{50} is a large integer number containing 50 decimal digits. The polynomials $\mathcal{F}_3(d_4)$, $\mathcal{F}_4(d_4)$, $\mathcal{F}_6(d_4)$, $\mathcal{F}_8(d_4)$, $\mathcal{F}_{10}(d_3, d_4)$, $\mathcal{F}_{12}(d_4)$, and $\mathcal{F}_{60}(d_4)$ are not presented here for the sake of brevity. The case when $d_4 = 1$ was already considered above, see the schemes FR5 and FR6 in Table 1. The cases when $d_4 = 1/3$ and $d_4 = 1/4$ were already considered above, see Table 3.

The third resultant $\text{Res}(G_1, G_2, d_3)$ is the 13th degree irreducible polynomial in two variables c_1 and d_4 . It adds no new solutions to system (18).

We have used above in the call (26) to the *Mathematica* function `GroebnerBasis[...]` the sequence of the indeterminates in the form `c2, c3, c4, d1, d2, d3, d4`. Below we show that permutation of elements in this sequence can lead to new solutions of the original polynomial system. To do this, consider instead of (26) the following call to the function `GroebnerBasis[...]`:

$$\begin{aligned} \text{GroebnerBasis}[\{P1, R1, P2, P31, P32, P41, P42, P43\}, \\ \{d1, d2, d3, d4, c2, c3, c4\}] \end{aligned} \tag{29}$$

In this case, a Gröbner basis is obtained in which $G_1 = \mathcal{G}_1(c_1, c_4)$, $G_2 = \mathcal{G}_2(c_1, c_3, c_4)$. Therefore, we have considered three resultants: $\text{Res}(G_1, G_2, c_4)$, $\text{Res}(G_1, G_2, c_3)$, and $\text{Res}(G_1, G_2, c_1)$. The first of them is as follows:

$$\begin{aligned} \text{Res}(G_1, G_2, c_4) &= A_{20}(4c_1 - 1)^6(6c_1 - 1)^{12}(24c_1^2 - 9c_1 + 1)^7 \\ &\times (27c_1^2 - 10c_1 + 1)^5(96c_1^3 - 60c_1^2 + 12c_1 - 1)^5[\mathcal{F}_{25}(c_1)]^7. \end{aligned} \tag{30}$$

Table 5. Successful permutations

Permutation No.	Permutation	New solutions
8	$\{d_1, d_2, d_3, c_2, d_4, c_4, c_3\}$	$c_3 = 0, c_3 = 1/2$
32	$\{d_1, d_2, d_4, c_2, d_3, c_4, c_3\}$	$c_3 = 0, c_3 = 1/2$
40	$\{d_1, d_2, d_4, c_3, c_2, c_4, d_3\}$	$d_3 = 0, d_3 = 3/4$
48	$\{d_1, d_2, d_4, c_4, c_3, c_2, d_3\}$	$d_3 = 0$
56	$\{d_1, d_2, c_2, d_4, d_3, c_4, c_3\}$	$c_3 = 1/2, c_1 = 1/6, c_3 = 0$
64	$\{d_1, d_2, c_2, c_3, d_4, c_4, d_3\}$	$d_3 = 0, d_3 = 3/4$
80	$\{d_1, d_2, c_3, d_4, d_3, c_4, c_2\}$	$c_2 = 0$
88	$\{d_1, d_2, c_3, c_2, d_4, c_4, d_3\}$	$d_3 = 0, d_3 = 3/4$
104	$\{d_1, d_2, c_4, d_4, d_3, c_3, c_2\}$	$c_2 = 0$
1000	$\{d_2, d_4, d_3, c_3, c_2, c_4, d_1\}$	$d_1 = 0$
3000	$\{c_2, d_1, c_4, c_3, d_4, d_3, d_2\}$	$d_2 = 0$
4000	$\{c_3, d_4, d_2, c_2, d_3, c_4, d_1\}$	$d_1 = 0$
5000	$\{c_4, c_3, d_4, d_2, d_1, c_2, d_3\}$	$c_1 = 1/6$

Here A_{20} is an integer consisting of 20 decimal digits. The equation $\text{Res}(G_1, G_2, c_4) = 0$ has two rational roots $c_1 = 1/4$ and $c_1 = 1/6$. Root $c_1 = 1/6$ is a new root that is missing in (27) and (28). Its consideration leads to two new schemes FR22 and FR23, see Table 4.

The resultant $\text{Res}(G_1, G_2, c_3)$ is an irreducible polynomial of the 13th degree in two variables c_1 and c_4 . The third resultant has the following form: $\text{Res}(G_1, G_2, c_1) = A_{72}c_4^2(8c_4^2 - 7c_4 + 2)^2(64c_4^3 - 40c_4^2 + 8c_4 + 1)[\mathcal{F}_6(c_4)]^2\mathcal{F}_8(c_4) \times \mathcal{F}_{12}(c_3, c_4)\mathcal{F}_{125}(c_4)$. The root $c_4 = 0$ of the equation $\text{Res}(G_1, G_2, c_1) = 0$ has proved to be an extraneous root of the polynomial system (18).

The above two examples of using different sequences of seven indeterminates point to the following: if we want to find as many solutions of the original polynomial system as possible, then we must consider all permutations of the indeterminates. In the case of seven indeterminates, the number of permutations is equal to $7! = 5040$. For each particular permutation, one needs to get three resultants based on the first two polynomials included in the Gröbner basis.

To do this, one first needs to get expressions for G_1 and G_2 and see which specific indeterminates and parameters they depend on. As a rule, G_1 and G_2 depend on two indeterminates and one variable selected as a parameter. Thus, we need to consider $5040 \cdot 3 = 15120$ resultants. The whole procedure could be implemented in the language of the *Mathematica* system. However, this requires considerable programming effort. Instead of such a complete analysis of all 5040 permutations, we considered the permutations P_j with the following numbers j : $j = 8, 16, 24, 32, 40, 48, 56, 64, 72, 80, 88, 96, 104, 1000, 2000, 3000, 4000, 5000$. In Table 5, not all of these permutations are given, but only those that led to new solutions for the eight unknowns $c_j, d_j, j = 1, \dots, 4$. To obtain all permutations of seven elements, the *Mathematica* function `Permutations[perm0]` was used. Here `perm0` is the initial permutation, we took `perm0 = {d1, d2, d3, d4, c2, c3, c4}`.

Table 5 contains a number of new solutions for the unknowns c_j and d_j . We have at first considered six cases when $c_j = 0, j = 2, 3, 4$ or $d_j = 0, j = 1, 2, 3$. In all these cases, the polynomial system (18) proved to be incompatible.

At $c_3 = 1/2$, a triangular Gröbner basis $\{G_1, \dots, G_7\}$ is obtained, where $G_1 = (1 - 3c_4 + 6c_4^2)(-1 - 24c_4 + 48c_4^2)(1 - 32c_4 + 364c_4^2 - 4000c_4^3 + 10288c_4^4 - 8064c_4^5 - 29696c_4^6 + 73728c_4^7)$. The roots of the equation $1 - 3c_4 + 6c_4^2 = 0$ are complex. The roots of the equation $48c_4^2 - 24c_4 - 1 = 0$ are real: $(c_4)_{1,2} = \frac{1}{12}(3 \pm 2\sqrt{3})$. This case was handled above, see the schemes FR3 and FR4. The equation $1 - 32c_4 + 364c_4^2 - 4000c_4^3 + 10288c_4^4 - 8064c_4^5 - 29696c_4^6 + 73728c_4^7 = 0$ has three real roots and two pairs of complex conjugate roots. Real roots lead to three new schemes FR24, FR25, and FR26.

At $d_3 = 3/4$, a triangular Gröbner basis $\{G_1, \dots, G_7\}$ is also obtained. The equation $G_1 = 0$ has only two real roots. The corresponding new schemes are presented in Table 6 as the schemes FR27 and FR28.

So, above we considered only 18 permutations, or 3.57 % of 5040 permutations. Therefore, it is natural that considering all 5040 permutations will allow one to obtain several further new solutions to the system (18).

The Remaining Variables as Parameters. The cases when one of seven variables $c_2, c_3, c_4, d_1, d_2, d_3, d_4$ is declared as a parameter were handled similarly to the above-considered case of the variable c_1 as a parameter, see Table 7, 8, 10, 11, 12, 13, 14, 15, and 16. Table 17 shows the number of new schemes obtained at the consideration of each of the variables $c_i, d_i, i = 1, \dots, 4$ as a parameter.

Table 18 presents the orders of smallness of the residuals r_{\min} and r_{\max} for several new symplectic schemes of the Forest–Ruth family, which were obtained above. Naturally, in cases where the analytical expressions for c_i and $d_i, i = 1, \dots, 4$ are relatively simple, it is possible to check their correctness by substituting them into the original polynomial system without switching to machine floating-point numbers. If the solution found in the analytical form is correct, then all the equations of the polynomial system are exactly satisfied. Therefore, in these cases, the equalities $r_{\min} = 0$ and $r_{\max} = 0$ take place. Formally, the schemes FR3, FR4, FR7, FR8, FR9, and FR10 are the best in terms of the smallness of the residuals r_{\min} and r_{\max} . However, they are far from optimal in terms of the smallness of the functionals $P_{5,rms}^{(l)}$ and $X_{5,rms}^{(l)}$, see Table 19.

In other cases, numerical values of c_i and $d_i, i = 1, \dots, 4$ were substituted into the initial system in the form of machine floating-point numbers and with a mantissa length of 18 decimal digits. As can be seen from Table 18, schemes FR27, FR50, FR51, FR52, and FR73 have the smallest residuals $r_{\min} = O(10^{-50})$ and $r_{\max} = O(10^{-49})$. The schemes FR42, FR47, FR50, FR51, FR52, FR55, and FR59 have the least approximation errors, these are the best found schemes (see Table 19). These conclusions partially correlate with the conclusions from Table 18 concerning the schemes with the least residuals r_{\min} and r_{\max} . The magnitude of the error functional $P_{5,rms}^{(l)}$ of scheme FR10 obtained previously in [6] is 121 times larger than in the cases of new schemes FR47, FR50, FR51, and FR52.

5 Kepler’s Problem

The problem of the motion of a system consisting of two interacting particles (the problem of two bodies, the Kepler’s problem) admits a complete analytical solution in the general form [8]. This solution was presented in [14] for the case of equal masses of the both bodies, and in [15], it was given for the bodies with different masses.

This problem is described by a system of eight ordinary differential equations (ODEs), in which four equations describe the time evolution of the coordinates (x_k, y_k) of each of the two bodies, and the quantities $(p_{kx}, p_{ky}), k = 1, 2$ describe the evolution of the impulses of the bodies. The mentioned system of ODEs is solved at the given initial positions of the bodies and their impulses. With equal masses of bodies $m_1 = m_2 = 1$, the total initial energy E_0 of both bodies $|E_0| = |v_0^2 - 1/(2a_0)|$, where $2a_0$ is the given initial distance between the bodies ($a_0 > 0$), v_0 is the given absolute value of the initial velocity of each body.

Table 6. The values of the parameters of the Forest–Ruth schemes FRl ($l = 24, \dots, 28$) obtained in the cases when $c_3 = 1/2$ or $d_3 = 3/4$

FRl		c_i	d_i
FR24	$c_3 = 1/2$	0.391025572210943674	0.967398350433141497
		0.678124844465627657	2.202427111618248827
		1/2	-0.111003864212194508
FR25	$c_3 = 1/2$	0.835311238373158736	-0.350920339289419778
		-0.378958977995565850	1.171732061213387849
		1/2	-1.814913008091942417
FR26	$c_3 = 1/2$	0.094132364293527189	0.723711816433472820
		-0.046380082078234786	-0.399429463535403217
		1/2	0.497003465452825098
FR27	$d_3 = 3/4$	0.200343862556644811	0.645182676558732474
		-1.010914108245682741	-0.024497356766309670
		1.772663714786654245	3/4
FR28	$d_3 = 3/4$	-1.742065593683020234	-0.631721835302486836
		-0.296628561098109581	0.410158533684560698
		1.724192327420640529	3/4

Table 7. The values of the parameters of the Forest–Ruth schemes FRl ($l = 29, \dots, 36$) obtained in the case when c_2 is the parameter

FRl		c_i	d_i
FR29	$c_2 = 2/3$	0.364254344284583814	0.869516211247767583
		2/3	2.792690380329816512
		0.587980436123375884	0.054227787182224768
FR30	$c_2 = 2/3$	1.878392245341283130	-0.148939912684807651
		2/3	0.072050469147999406
		-2.119991445360319150	0.835740764826119230
FR31	$c_2 = 2/3$	0.197431266262198278	0.997479482407853902
		2/3	-0.097322556987957231
		-0.933855594163194543	-0.268583735242126327
FR32	$c_2 = 2/3$	0.168148869366041101	0.464264589751261379
		2/3	0.493986641393048766
		0.2088349346809210962	-0.737725321402639611
FR33	$c_2 = 2/3$	-0.170487789371228258	0.124713923957057081
		2/3	0.644865752568700276
		0.5457035136145351390	-0.695444714544033987
FR34	$c_2 = 2/3$	0.374170872627683649	0.717310607339155385
		2/3	-1.033722587512287717
		-6.140624051165247365	0.001454385616574405
FR35	$c_2 = 1/3$	0.305530612248020483	4.937738540440601523
		1/3	-0.027239637450445254
		-0.339218512202854752	-4.168275877672941372
FR36	$c_2 = 1/3$	0.680077031402771998	0.881570343081302874
		1/3	-0.790588680725968301
		-1.186453739372546212	0.176773784101247734

Table 8. The values of the parameters of the Forest–Ruth schemes FR*l* ($l = 37, \dots, 46$) obtained in the case when c_2 is the parameter

FR <i>l</i>	c_i	d_i
FR37	$c_2 = (1 + \sqrt{5})/3$	3.782386995165049639
		-0.076051961619127312
		0.043960674924039808
FR38	$c_2 = (1 + \sqrt{5})/3$	0.144891076593792282
		0.784982662463530111
		-0.115777773963705551
FR39	$c_2 = (1 + \sqrt{5})/3$	0.213103271189846452
		0.573068437077953834
		1.002219472013933223
FR40	$c_2 = (1 + \sqrt{5})/3$	0.021085719326262061
		-1.091461204437575932
		1.078689325833263232
FR41	$c_2 = (1 - \sqrt{5})/3$	0.219311236012542665
		0.760452067359764565
		-0.091641127721282675
FR42	$c_2 = (1 - \sqrt{5})/3$	0.1394284522075851837
		0.486323892019975997
		-0.042011883595088195
FR43	$c_2 = (1 - \sqrt{5})/3$	0.967231066343295769
		0.441924899417734389
		0.318257154554292683
FR44	$c_2 = (1 - \sqrt{5})/3$	0.0318257154554292683
		0.826446621342308992
		-1.905814461246662400
FR45	$c_2 = (1 - \sqrt{5})/3$	-0.009200018152718083
		1.832414525155164045
		-2.671203027064244664
FR46	$c_2 = (1 - \sqrt{5})/3$	0.127955739058414071
		-0.099395301108287731
		3.473342068227190026
FR47	$c_2 = (1 - \sqrt{5})/3$	0.715057655495029390
		-1.873117657996741733
		-0.649164960959432341
FR48	$c_2 = (1 - \sqrt{5})/3$	-0.412022659166596565
		0.364748420484928815
		1.780573723084865883
FR49	$c_2 = (1 - \sqrt{5})/3$	0.775557276840459310
		-1.873117657996741733
		-0.649164960959432341
FR50	$c_2 = (1 - \sqrt{5})/3$	-0.412022659166596565
		0.364748420484928815
		1.780573723084865883
FR51	$c_2 = (1 - \sqrt{5})/3$	0.775557276840459310

When numerically solving the Kepler problem using symplectic difference schemes, the energy e^n at $t = t_n > 0$ is calculated with an error. Let $\delta E^n = (E^n - E_0)/E_0$, where $E^n = (1/2)[(p_{1x}^n)^2 + (p_{1y}^n)^2 + (p_{2x}^n)^2 + (p_{2y}^n)^2] - 1/r^n$, $r^n = [(x_1^n - x_2^n)^2 + (y_1^n - y_2^n)^2]^{1/2}$. It was shown in [16] that for ensuring the zero eccentricity e of the particles orbits it is sufficient to set $v_0 = 0.5/\sqrt{a_0}$. Along

Table 9. The values of the parameters of the Forest–Ruth schemes FR*l* ($l = 47, \dots, 53$) obtained in the cases when $c_4 = 3/4$ or $d_1 = 1/2$

FR <i>l</i>		c_i	d_i
FR47	$c_4 = 3/4$	4.290878395464025169	3.434400270740872001
		0.000963763146154897	-3.435707719128205871
		-4.041842158610180066	0.656514850491044033
FR48	$c_4 = 3/4$	0.277948154623496169	2.459832693693553053
		0.343888647781502973	-0.056699486586081620
		-0.371836802404999142	-1.672913435532665059
FR49	$c_4 = 3/4$	0.520981254725582490	1.113980205862570658
		-0.161927364012237208	-1.472633501812066773
		-0.109053890713345282	1.213681750558400100
FR50	$d_1 = 1/2$	0.136825942475053071	1/2
		-0.295364245574992759	-0.062976137694193308
		0.837926908632179149	0.440948891327570560
FR51	$d_1 = 1/2$	-0.586438190415940380	1/2
		-0.075359198960140901	-0.430061540610022957
		1.038753638680670670	0.676446958625960263
FR52	$d_1 = 1/2$	0.182572224329682333	1/2
		0.744884110251947176	0.538588670363744296
		0.141395996554118555	-0.758360900822698969
FR53	$d_1 = 1/2$	-3.530966905753254684	1/2
		6.547681802370795675	-0.068515429495817908
		-7.087388112997812705	-0.375425403877225227

with the error in the energy δE^n , the values of $\delta r_{1,\max}$ and $\delta y_{1,\text{mean}}$ were also calculated, where $\delta r_{1,\max} = \max_j((x_{1,j}^2 + y_{1,j}^2)^{1/2} - a)/a$ for the case of the zero eccentricity. The quantity $\delta y_{1,\text{mean}}$ was computed as the arithmetic mean of the quantities $\delta y_{1j} = y_{1j} - y_{1,\text{ex}}$. Here $y_{1,\text{ex}}$ is the exact value of the y coordinate at the point of the intersection of line $x = x_{1j}$ with the ellipse of the first particle.

As one can see in Table 20, the symmetric Forest–Ruth scheme FR9 produces much larger errors $|\delta E|_{\text{mean}}$ than the schemes FR3 and FR50. This is consistent with the behavior of the norm of the leading error term of these schemes.

Table 21 shows the results obtained at the nonzero eccentricity by the time $t = 500$ (10^5 time steps). As in the case of the zero eccentricity, the schemes FR3 and FR50 demonstrate a higher accuracy than the scheme FR9. The error in energy δE_{mean} amounts in the case of the scheme FR50 to only about 1/8 of the error δE_{mean} obtained in the case of the scheme FR3. This means that the scheme FR50 produces a much smaller error in energy than the scheme FR3.

Table 10. The values of the parameters of the Forest–Ruth schemes FR*l* (*l* = 54, ..., 63) obtained in the case when c_3 is the parameter

FR <i>l</i>	c_i	d_i	
FR54	$c_3 = 1$	0.232552203620762051	0.864933853224711811
		-0.240171870993110794	-0.173521835714879970
		1	1.450514067565065171
FR55	$c_3 = 1$	0.139963046708770169	0.484042373117847543
		-0.442222496742616455	-0.038391737134385132
		1	0.442340759342235601
FR56	$c_3 = 1$	-0.549628408322132739	0.516200660144340638
		-0.072468674149454992	-0.443663538175265205
		1	0.674727683742161504
FR57	$c_3 = 1$	2.302601066465332119	0.112742296458862205
		-1.983691962779933225	0.511421634072704643
		1	-0.936108794035271761
FR58	$c_3 = 2/3$	0.351122270415848032	0.822180431133326332
		0.663430432875993088	3.675251576576031878
		2/3	-0.029344841230348257
FR59	$c_3 = 2/3$	0.130496314992084781	0.545153172452857850
		-0.152112316952245843	-0.129275520560538766
		2/3	0.445295236274126733
FR60	$c_3 = 2/3$	-0.214714867139629529	0.684249379826096300
		-0.050505925257735998	-0.566790506272220428
		2/3	0.644799229296194684
FR61	$c_3 = 2/3$	12.985132914037646926	0.000735492767526182
		-12.559557711764700905	0.717113943685514226
		2/3	-1.050357583178864997
FR62	$c_3 = -1/3$	0.323096342154017994	13.50376205182945044
		0.332614327860069729	-0.009534841164435389
		-1/3	-12.74671531967306675
FR63	$c_3 = -1/3$	1.549722230386746638	-3.514818811551034895
		1.429077640949028648	-1.457098486909627231
		-1/3	2.622648492232692156

The absolute value of the error $\delta y_{1,mean}$ amounts in the case of the scheme FR50 to only about 1/5 of the error $\delta y_{1,mean}$ obtained in the case of the scheme FR3. Similar errors $\delta y_{2,mean}$ obtained for the second particle have the same absolute values as in the case of the first particle, but their signs are opposite to the signs of $\delta y_{1,mean}$.

Table 11. The values of the parameters of the Forest–Ruth schemes FR l ($l = 64, \dots, 71$) obtained in the case when c_4 is the parameter

FR l	c_4	c_i	d_i
FR64	$c_4 = -1$	0.923752664541586022	16.023236171683106527
		1.028858031528064696	10.667066118735512987
		0.047389303930349282	-9.439832216877909698
FR65	$c_4 = -1$	0.230401316001649608	0.614993313027275731
		1.767802001353397688	2.349801988037760294
		0.001796682644952704	-2.372282095213200527
FR66	$c_4 = -1$	2.693266252232378788	0.272765858595981272
		-2.687864699530211844	0.258660862508406001
		1.994598447297833056	-0.704601530894185303
FR67	$c_4 = 2/3$	0.693911807455204925	1.389601984542431907
		-0.180750033711638009	-1.751837222444102726
		-0.179828440410233583	1.391280849389145194
FR68	$c_4 = (-1 + \sqrt{5})/3$	0.114099408934045299	0.649850775109876324
		-0.066472584291344747	-0.282718024274440564
		0.540350516190702883	0.469755418612118481
FR69	$c_4 = (-1 - \sqrt{5})/3$	0.231348624383101124	0.617413961551446993
		1.845849356904791962	2.505944087496861072
		0.001491344545370146	-2.526048748369218988
FR70	$c_4 = (-1 - \sqrt{5})/3$	2.763558393992365262	0.282057553866780170
		-2.795071176619652551	0.244409846615326875
		2.110202108460550521	-0.690943228859052633
FR71	$c_4 = (-1 - \sqrt{5})/3$	0.332267944750488842	0.743399694345549828
		0.667975398001819257	-21.981819840398356626
		1.078445983080955133	0.001613923746367732

Table 12. The values of the parameters of the Forest–Ruth schemes FR l ($l = 72, \dots, 75$) obtained in the case when $c_4 = 1/6$ or $c_4 = 16/27$

FR l	c_4	c_i	d_i
FR72	$c_4 = 1/6$	1.075394753302521569	-0.638191122739963413
		-0.562921915373837090	1.672633703407438461
		0.320860495404648855	-2.181423749195410509
FR73	$c_4 = 16/27$	0.818672678175588272	1.833608220246540205
		-0.239249580730862964	-2.402016364561991597
		-0.172015690037317901	1.987452002208430171
FR74	$c_4 = 16/27$	-0.176216464958336337	0.705396326757519899
		-0.048272736660086400	-0.577957534527034398
		0.631896609025830144	0.637881726660246176
FR75	$c_4 = 16/27$	0.494971671137937628	-0.948190157578799342
		0.369511495225417267	0.079013950429035923
		-0.457075758955947487	1.633761849965980780

Table 13. The values of the parameters of the Forest–Ruth schemes FR*l* (*l* = 76, . . . , 85) obtained in the case when *d*₁ is the parameter

FR <i>l</i>	<i>d</i> ₁	<i>c</i> _{<i>i</i>}	<i>d</i> _{<i>i</i>}
FR76	<i>d</i> ₁ = 2/3	0.204275083815869887	2/3
		−0.821240613630061829	−0.034251150245736271
		1.586167530241972297	0.808482764907095971
FR77	<i>d</i> ₁ = 2/3	0.110529614169578658	2/3
		−0.060536883821488756	−0.308218556075319167
		0.529316569535254555	0.475003634825955049
FR78	<i>d</i> ₁ = 2/3	−0.246981683746760770	2/3
		−0.052395107294725361	−0.556352049684461481
		0.696641970456244612	0.649688395520743066
FR79	<i>d</i> ₁ = 2/3	−0.029918619502392475	2/3
		1.391696101106459054	−0.179948437412172961
		−1.753316304811621968	−0.180888740802033290
FR80	<i>d</i> ₁ = 3/4	0.217797377255593747	3/4
		−0.439474733154072297	−0.084286494754561427
		1.206367201598754224	1.065889648917041605
FR81	<i>d</i> ₁ = 3/4	0.082220658475439553	3/4
		−0.042043274315474989	−0.446288824714853250
		0.490110712558433011	0.511001742424794030
FR82	<i>d</i> ₁ = 3/4	−0.093660943038827783	3/4
		−0.043524035940067826	−0.593590536527777054
		0.562491408233619701	0.617331657398234449
FR83	<i>d</i> ₁ = 3/4	0.282751320628652862	3/4
		−0.875414602261540111	−2.866262183797889275
		−0.001749298330001449	2.839324076485054340
FR84	<i>d</i> ₁ = 3/4	0.122962367199072599	3/4
		1.154710480553913358	−2.866262183797889275
		−1.483263553884396613	2.839324076485054340
FR85	<i>d</i> ₁ = 3/4	0.453772840879245976	3/4
		0.592568295889989923	−0.763378994607420991
		−2.069852675023231005	0.027157631288323441

Table 14. The values of the parameters of the Forest–Ruth schemes FR*l* (*l* = 86, . . . , 90) obtained in the cases when $d_1 = 1/3$ or $d_2 = 1/6$

FR <i>l</i>		c_i	d_i
FR86	$d_1 = 1/3$	-1.061210383861491179	1/3
		-0.124764893770052270	-0.283051205857396058
		1.561210383861491179	0.690936406092486640
FR87	$d_1 = 1/3$	0.105100678126442684	1/3
		0.529988715863071753	0.502653125702352770
		0.394899321873557316	-0.609454883379762706
FR88	$d_1 = 1/3$	3.383237864740253990	1/3
		-3.727448818616042444	1/6
		3.082935363183809675	-0.615610764345786215
FR89	$d_2 = 1/6$	-4.276220596064195618	-0.859926774390249674
		-1.952856967612454919	1/6
		3.168223353300535282	0.943721911449059608
FR90	$d_2 = 1/6$	3.383237864740253990	1/3
		-3.727448818616042444	1/6
		3.082935363183809675	-0.615610764345786215

Table 15. The values of the parameters of the Forest–Ruth schemes FR*l* (*l* = 91, . . . , 94) obtained in the case when d_2 is the parameter

FR <i>l</i>		c_i	d_i
FR91	$d_2 = 1/2$	-1.656461349913587036	-0.636061115269224549
		-0.161698671975507331	1/2
		1.728130415131394716	0.716119776115948032
FR92	$d_2 = 1/2$	-0.927682804781144334	0.821495974995660855
		0.567079873156935371	1/2
		0.999351869998952014	-0.741437314148937372
FR93	$d_2 = 1/2$	0.108591275285244788	0.339219840428187569
		0.532722086799126630	1/2
		0.388349357158891234	-0.616031036746662437
FR94	$d_2 = 1/2$	2.293935572964878270	0.119753909827438190
		-1.982466778384829536	1/2
		1.023930342194913976	-0.926323914261688727

Table 16. The values of the parameters of the Forest–Ruth schemes FR*l* (*l* = 95, . . . , 98) obtained in the case when d_3 is the parameter

FR <i>l</i>		c_i	d_i
FR95	$d_3 = 2/3$	0.193635546310194233	0.610720278963038252
		-1.514131240893582026	-0.012201087465723859
		2.266721524338190638	2/3
FR96	$d_3 = 2/3$	-0.414351004195240352	0.579788510045530680
		-0.062798023340654284	-0.494940129503867069
		0.860535223163841919	2/3
FR97	$d_3 = 1/2$	0.091689601485633373	0.729954875905782486
		-0.045246631531655207	-0.410133860075899430
		0.497436187438423795	1/2
FR98	$d_3 = 1/2$	1.095187697120513955	1.250031937811706533
		0.066378351994972941	-1.227061470621614945
		-1.003038424429919768	1/2

Table 17. The number of new schemes found in the Forest–Ruth family at the consideration of each of the variables $c_i, d_i, i = 1, \dots, 4$ as a parameter

Parameter	c_1	c_2	c_3	c_4	d_1	d_2	d_3	d_4
Number of new schemes	28	25	10	12	15	4	4	0

Table 18. The residuals r_{\min} and r_{\max} for several schemes from the Forest–Ruth family

l	r_{\min}	r_{\max}	l	r_{\min}	r_{\max}	l	r_{\min}	r_{\max}
9	0	0	50	$O(10^{-50})$	$O(10^{-49})$	73	$O(10^{-50})$	$O(10^{-49})$
10	0	0	51	$O(10^{-50})$	$O(10^{-49})$			
27	$O(10^{-50})$	$O(10^{-49})$	52	$O(10^{-50})$	$O(10^{-49})$			

Table 19. The values of the error functionals $P_{5,rms}^{(l)}$ and $X_{5,rms}^{(l)}$ for several schemes from the Forest–Ruth family

l	$P_{5,rms}^{(l)}$	$X_{5,rms}^{(l)}$	l	$P_{5,rms}^{(l)}$	$X_{5,rms}^{(l)}$	l	$P_{5,rms}^{(l)}$	$X_{5,rms}^{(l)}$
9	6.3431	2.5624	47	0.0386	0.0708	52	0.0386	0.0708
10	4.6743	8.3036	50	0.0386	0.0708	55	0.0498	0.0915
42	0.0471	0.0876	51	0.0386	0.0708	59	0.0559	0.0550

Table 20. Errors $\delta E_{mean}, |\delta E|_{mean},$ and $\delta r_{1,max}$ at $e = 0$ and $t = 7140h, h = 0.005$ for the fourth-order Forest–Ruth methods from Tables 1, 2, and 9

Forest–Ruth scheme	δE_{mean}	$ \delta E _{mean}$	$\delta r_{1,max}$
FR9	$-1.878e - 14$	$1.878e - 14$	$4.636e - 13$
FR3	$-5.369e - 15$	$6.693e - 15$	$5.462e - 14$
FR50	$3.533e - 15$	$7.654e - 15$	$1.954e - 14$

Table 21. Errors $\delta E_{mean}, |\delta E|_{mean},$ and $\delta y_{1,mean}$ at $v_0 = 0.15$ for the fourth-order Forest–Ruth methods from Tables 1, 2, and 9

Forest–Ruth scheme	δE_{mean}	$ \delta E _{mean}$	$\delta y_{1,mean}$
FR9	$3.292e - 9$	$3.292e - 9$	$-3.542e - 6$
FR3	$9.684e - 10$	$9.684e - 10$	$-8.455e - 7$
FR50	$1.226e - 10$	$1.284e - 10$	$1.714e - 7$

6 Conclusions

The problem of constructing higher-order symplectic integration techniques for molecular dynamics problems with separable Hamiltonians is considered. A general method for finding symplectic schemes of high order of accuracy using parametric Gröbner bases, resultants, and permutations of variables is proposed. The implementation of the method is described by the example of four-stage partitioned Runge–Kutta (PRK) schemes of the Forest–Ruth family. This method has enabled us to find 96 new symplectic four-stage schemes in the Forest–Ruth family. Among these schemes, several schemes have been found that are the best in terms of the smallness of the leading term of the approximation error. It turned out that the value of the error functional of the best of the two schemes obtained earlier in [6] is 121 times greater than in the case of the new best PRK schemes found in this paper.

All required symbolic calculations are performed using the computer algebra system *Mathematica*. When searching for new schemes, it turned out to be effective to combine the technique of Gröbner bases with Sylvester’s resultants and with permutations in the order of variables in the call of *Mathematica* function that calculates the Gröbner basis.

We emphasize that the real solutions of polynomial systems found using the technique described above do not exhaust the entire variety of solutions of these systems, since in this paper we focused mainly on finding solutions to the original polynomial system obtained using integer and rational roots of resultants, as well as roots of quadratic multipliers of resultants. Further solutions can be found with sufficient accuracy for applications by numerically solving the problems of minimizing the functional, which takes into account the value of one of the desired parameters, found using the resultants. This work is quite feasible and can lead to even more accurate PRK schemes compared to those schemes that were discovered in the framework of the study described above.

The presented study shows that before increasing the number of stages, it is advisable to conduct a detailed search for optimal parameters of the scheme at a fixed number of stages.

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