

Analytical Calculation of Beam Supports Stiffness to Provide the First Eigenfrequency and Critical Load

I. V. Kudryavtsev $^{1(\boxtimes)}$, M. V. Brungardt¹, M. P. Khomutov², O. I. Rabetskaya³, and A. E. Mityaev $¹$ </sup>

¹ Siberian Federal University, 79, Svobodny pr., Krasnoyarsk 660041, Russia ikudryavcev@sfu-kras.ru

² Kuban State Technological University, 2, Moskovskaya, Krasnodar 350072, Russia ³ Reshetnev Siberian State University of Science and Technology, 31, Krasnoyarskij rabochij, Krasnoyarsk 660037, Russia

Abstract. This paper is concerned with the problem of ensuring both the first eigenfrequency of the a two-support beam and its first critical force by determining the required stiffness of supports. This is a quite complex mathematical problem and well-known scientific literature usually offers a solution only in graphic or tabular form. One of the main problems is the highly nonlinear dependence of results on supports stiffness at free vibrations and stability loss. In the paper, these difficulties are overcome by approximation of nonlinear dependencies using the least-squares method and getting analytical quadratic approximating functions instead. As a result, we get a closed-form solution for the problem in the form of a fourth-degree resolving equation that has an analytical solution. This solution allows determining support stiffness which provides the first eigenfrequency of a two-support beam and its first critical force. Replacing the strongly nonlinear dependencies with simpler quadratic functions, however, adversely affected the calculation error which can reach 10%. To reduce this error, it is recommended that the stiffness of both supports be equal or of the same order.

Keywords: Euler-Bernoulli beam · Elastic supports · Free vibrations · Buckling · First eigenfrequency · Axial force · Analytical solution

1 Introduction

There are many beam structures subject to forced vibration along with compressive axial force. In this case, these structures have to provide the values of the first eigenfrequency and the first critical force. These conditions can be written as:

$$
f_1 \ge [f], \ P_{cr1} \ge [P] \tag{1}
$$

where $[f]$ and $[P]$ are the allowable frequency and force values, respectively.

The dynamic behavior of the Euler-Bernoulli beam is governed by two 4th-order differential equations with specified boundary conditions, which reflect beam support

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features. Simple support conditions, like a hinge, fixed-ended, and their combinations, lead to trivial solutions of the governing equations that can be found in the numerous scientific literature and dynamics guides $[1-4]$ $[1-4]$. In fact, the real supports have some finite stiffness, which affects both the value of the first vibration frequency and the first critical force.

Taking into account the finite stiffness of the supports complicates the solution since it has a form of the transcendent equation without an explicit analytical solution. In dynamic, different types of support are taken into account in the form of so-called support coefficients, which are obtained from the numerical solution of governing equations. In simple cases of fixation, when the support stiffness is zero or infinity, the values of these coefficients are given in the scientific literature. Finite support stiffness with corresponding values of support coefficients usually are presented in the literature in the form of graphs or tables on the base of their separate values. This method of representation is less convenient in calculations than analytical expressions, since the required stiffness value usually does not coincide with the presented data.

This paper proposes a method of solution for the governing equations through approximation of these nonlinear dependencies using the least-squares method and getting analytical quadratic functions. As a result of which we can solve reverse problems for the determination of the required supports stiffness, which provides both values of the first eigenfrequency and critical force for a two-support beam.

2 Methods

Let's consider the mathematical description of the free vibration and stability problem for an Euler-Bernoulli beam on elastic supports and get the fundamental equations.

2.1 Governing Equations

Let us consider a straight beam fixed in hinged supports with some stiffness k_1 , k_2 and exposed to external action in the form of axial force *P* and uniform temperature Δ*T* (Fig. [1\)](#page-1-0). The dynamic behavior of the beam is evaluated by the first eigenfrequency of bending vibrations and the first critical force.

Fig. 1. Beam with elastic supports.

The equation of free vibration of the beam taking into account the action of the compressive axial force *P* has the form $[1-4]$ $[1-4]$:

$$
EJ_{\min}\frac{\partial^4 y}{\partial x^4} + P\frac{\partial^2 y}{\partial x^2} + m\frac{\partial^2 y}{\partial t^2} = 0,
$$
 (2)

where $y = y(x, t)$ is the deflection function; *E* is the elastic modulus; *J_{min}* is the minimum moment of inertia of cross section; *m* is the unit weight, $m = \rho S$, ρ is the density; *S* is the cross-sectional area; *P* is the axial force.

If axial force *P* reaches a certain critical value, it can lead to a beam buckling. The stability equation has the form $[5-16]$ $[5-16]$:

$$
EJ_{\min}\frac{\partial^4 y}{\partial x^4} + P\frac{\partial^2 y}{\partial x^2} = 0,\tag{3}
$$

The solution of the Eq. [\(2\)](#page-2-0) for the first eigenfrequency at $P = 0$ has the form:

$$
f_1 = \frac{\alpha^2}{2\pi l^2} \cdot \sqrt{\frac{E J_{\text{min}}}{m}},\tag{4}
$$

where α is the support coefficient at vibrations, which reflects support stiffness.

The first critical force can be found from Eq. [\(3\)](#page-2-1) in the form:

$$
P_{cr1} = \frac{\pi^2 E J_{\text{min}}}{\mu^2 \cdot l^2},\tag{5}
$$

where μ is the support coefficient at beam buckling or the so-called column effective length factor, which also determined by the beam supports.

With the simultaneous action of axial force and free vibration, the obtained solutions [\(4,](#page-2-2) [5\)](#page-2-3) are connected to each other by Galef's formula [\[16\]](#page-8-3):

$$
f_{1(P>0)} = f_{1(P=0)} \cdot \sqrt{1 - \frac{P}{P_{cr1}}},
$$
\n(6)

where $f_i(P \le \theta)$ is the first eigenfrequency under action of axial force $P; f_1(P = 0)$ is the first eigenfrequency in absence of axial force P ; P_{cr1} is the first critical force.

Let's decompose the axial force P into two components according to the principle of superposition:

$$
P = P_{\Delta T} + P_F = \alpha_t \cdot \Delta T \cdot ES + P_F,\tag{7}
$$

where $P_{\Delta T}$ is the axial force from temperature change; P_F is the external force action; α_t is the coefficient of thermal expansion.

Substituting [\(4,](#page-2-2) [5,](#page-2-3) [7\)](#page-2-4) in Eq. [\(6\)](#page-2-5) gives the equation for the first eigenfrequency, taking into account temperature and force effects, as well as the method of fixation in the form of:

$$
f_{\parallel (P \sim 0)} = \left(\frac{\alpha}{\pi l}\right)^2 \cdot \sqrt{\frac{1}{4m} \left[\pi^2 E J_{\min} - \mu^2 \cdot l^2 \cdot \left(\alpha_i \cdot \Delta T \cdot ES + P_F\right)\right]} \,,\tag{8}
$$

Values of support coefficients α and μ for common cases of supports can be found in dynamics guides $[17–19]$ $[17–19]$ or calculated from governing equations solution $(2, 3)$ $(2, 3)$ $(2, 3)$. However, existing reference books offer a small number of these coefficients for simple types of beam supports, which have zero or infinity stiffness.

2.2 Functions of Support Coefficients

Consider the case of the stiffness of the supports k_1, k_2 between zero and infinity, then we get that the coefficients of the supports will be some functions of their stiffness (Fig. [1\)](#page-1-0):

$$
\alpha = \alpha(k_1, k_2), \mu = \mu(k_1, k_2). \tag{9}
$$

Express the support coefficients through the relative stiffness *Ci*:

$$
C_1 = k_1 \frac{l}{E J_{\min}}, C_2 = k_2 \frac{l}{E J_{\min}}.
$$
 (10)

Then we get new functions:

$$
\alpha = \alpha(C_1, C_2), \mu = \mu(C_1, C_2). \tag{11}
$$

Functions [\(11\)](#page-3-0) are strongly non-linear and cannot be expressed by simple analytical dependences; this is a reason that in the reference literature they are presented in tabular or graphical form only. In our previous papers [\[20,](#page-8-6) [21\]](#page-9-0) analytical functions [\(11\)](#page-3-0) were obtained for three ranges of stiffness *C1*, *C2*.

3 Determination of Required Support Stiffness

3.1 Governing Equation

Let's express Eq. (9) in a condition form:

$$
f_1 = \left(\frac{\alpha(C_1, C_2)}{\pi l}\right)^2 \cdot \sqrt{\frac{1}{4m} \left[\pi^2 E J_{\min} - \mu^2(C_1, C_2) \cdot l^2 \cdot (\alpha_l \cdot \Delta T \cdot ES + P_F)\right]} \ge [f_1].
$$
\n(12)

Next, we rewrite an Eq. [\(12\)](#page-2-0) in abbreviated form:

$$
A - \mu^2(C_1, C_2) \cdot B = \frac{C}{\alpha^4(C_1, C_2)},
$$
\n(13)

where

$$
A = \pi^2 E J_{\min}; \quad B = l^2 \cdot (\alpha_l \cdot \Delta T \cdot ES + P_F); \quad C = 4m[f_1]^2 \pi^4 l^4. \tag{14}
$$

The solution of Eq. (13) regarding the stiffnesses of C_i requires not the functions (11) themselves, but their 2nd and 4th degrees, which we denote as μ_2 and α_4 , respectively. We get these new functions in a few steps. Firstly, we calculated two sets of initial data from the obtained functions [\(11\)](#page-3-0) for the stiffness range $C = 0...1000$. Then we increase the values in the data sets till the required power. Finally, we approximate the new dataset using the least squares method and check the quadratic functions as

$$
\mu_2(C_1, C_2) = a_1 \cdot \left(C_1^2 + C_2^2 \right) + a_2 \cdot (C_1 + C_2) + a_3,\tag{15}
$$

$$
\alpha_4(C_1, C_2) = b_1 \cdot \left(C_1^2 + C_2^2 \right) + b_2 \cdot (C_1 + C_2) + b_3. \tag{16}
$$

The calculated coefficient values in the functions $(15, 16)$ $(15, 16)$ $(15, 16)$ are given in Table [1.](#page-4-1) Equation [\(13\)](#page-2-1) takes the form:

$$
A - \mu_2(C_1, C_2) \cdot B = \frac{C}{\alpha_4(C_1, C_2)}.
$$
 (17)

Table 1. Coefficient values.

Parameter	μ_2			α_4		
	a ₁	a ₂	a_3	b_1	b ₂	b_3
Zone I: $C = 010$	$1.9E - 04$	$1.5E - 0.5$	$2.5E - 07$	-1.4	-0.023	$-4E-0.5$
Zone II: $C = 10100$	-0.01381	-0.00592	-0.000515	34	4.36	0.08
Zone III: $C = 1001000$		0.938	0.545	97.5	257	456,5

The solution of Eq. (17) has the uncertainty in the form of an infinite number of C_1 and C_2 combinations, which satisfy a problem. To eliminate this uncertainty, we take $C_2 = n \cdot C_1$ then instead of [\(15,](#page-3-2) [16\)](#page-4-0) we get:

$$
\mu_2(C_1) = a_1 \cdot \left(C_1^2 + n^2 \cdot C_1^2 \right) + a_2 \cdot (C_1 + n \cdot C_1) + a_3,\tag{18}
$$

$$
\alpha_4(C_1) = b_1 \cdot \left(C_1^2 + n^2 \cdot C_1^2 \right) + b_2 \cdot (C_1 + n \cdot C_1) + b_3. \tag{19}
$$

For the short, we enter new coefficients:

$$
a_{1n} = a_1 \cdot \left(1 + n^2\right); \quad a_{2n} = a_2 \cdot (1 + n), \tag{20}
$$

$$
b_{1n} = b_1 \cdot \left(1 + n^2\right); \quad b_{2n} = b_2 \cdot (1 + n). \tag{21}
$$

Then Eq. [\(17\)](#page-2-4) takes the form:

$$
A - \left(a_{1n}x^2 + a_{2n}x + a_3\right) \cdot B = \frac{C}{b_{1n}x^2 + b_{2n}x + b_3}.\tag{22}
$$

Let's rewrite:

$$
c_1 = -Ba_1 \cdot \left(1 + n^2\right); c_2 = -Ba_2 \cdot \left(1 + n^2\right); c_3 = -Ba_3 + A;
$$

\n
$$
d_1 = \frac{b_1 \cdot \left(1 + n^2\right)}{C}; d_2 = \frac{b_2 \cdot \left(1 + n^2\right)}{C}; d_3 = \frac{b_3}{C}.
$$
\n(23)

Finally, we get the 4th degree governing equation in the form:

$$
c_1x^2 + c_2x + c_3 = \frac{1}{d_1x^2 + d_2x + d_3}.\tag{24}
$$

or

$$
c_1d_1x^4 + (c_1d_2 + c_2d_1)x^3 + (c_1d_3 + c_2d_2 + c_3d_1)x^2
$$

+
$$
(c_2d_3 + c_3d_2)x + (c_3d_3 - 1) = 0
$$
 (25)

This algebraic equation of the 4th degree [\(31\)](#page-5-0) determines the required value of the support stiffness, which is one of its 4 roots. Graphically, this solution can be shown as one of the intersection points of two 4th degree curves.

3.2 Analytical Solution of the Governing Equation

The 4th degree is the highest order of algebraic equation for which analytical solution exists [\[22\]](#page-9-1). The 4th degree equation has several analytical methods of solution. In this paper, we use method developed by Nesmeev that presented in his papers [\[22,](#page-9-1) [23\]](#page-9-2). To do so, let's transform Eq. [\(31\)](#page-4-2) so that the coefficient at the highest degree becomes equal to one:

$$
x4 + a3x3 + a2x2 + a1x + a0 = 0.
$$
 (26)

where

$$
a_3 = \frac{c_1 d_2 + c_2 d_1}{c_1 d_1}; \quad a_2 = \frac{c_1 d_3 + c_2 d_2 + c_3 d_1}{c_1 d_1}; \quad a_1 = \frac{c_2 d_3 + c_3 d_2}{c_1 d_1}; \quad a_0 = c_3 d_3 - 1.
$$
\n
$$
(27)
$$

Next, we find the coefficients of the auxiliary cubic equation:

$$
a = 1;
$$
 $b = -a_2;$ $c = a_1a_3 - 4a_0;$ $d = -a_1^2 - a_0a_3^2 + 4a_0a_2.$ (28)

Solution of the auxiliary cubic equation in canonical form requires finding its coefficients:

$$
p = \frac{3ac - b^2}{9a^2}; \quad q = \frac{(ab)^3}{27} - \frac{bc}{6a^2} + \frac{d}{2a}; \quad k = p^3 + q^2. \tag{29}
$$

Check of the coefficient number signs (29) shows that for the considered range of support stiffness we get $p < 0$, $q > 0$, $k < 0$. In this case, characteristic angle φ can be defined as:

$$
\varphi = \arctg\left(\sqrt{\frac{|p|^3}{q^2} - 1}\right). \tag{30}
$$

From the three roots of the auxiliary cubic equation, the maximum root is correct only. This root can be found from the equation:

$$
u_1 = 2r \cdot \cos\left(\frac{\pi - \varphi}{3}\right) - \frac{b}{3}.\tag{31}
$$

Next solution involves two square equations. Let's find the auxiliary coefficients for the square equations:

$$
d_1 = \frac{a_3^2}{4} + u_1 - a_2; \quad d_2 = \frac{u_1^2}{4} - a_0; \quad d_3 = a_3 u_1 - 2a_1. \tag{32}
$$

$$
s_{re} = \frac{0, 5a_3 + \sqrt{d_1}}{4} - \frac{u_1}{2} + \sqrt{d_2}.
$$
 (33)

The first square equation has a form:

$$
k_1 x^2 + k_2 x + k_3 = 0. \tag{34}
$$

where

$$
k_1 = 1;
$$
 $k_2 = \frac{a_3}{2} + \sqrt{d_1};$ $k_3 = \frac{u_1}{2} - \sqrt{d_2};$ (35)

Its real root is:

$$
x_1 = \frac{1}{2k_1} \left(-k_2 + \sqrt{k_2^2 - 4k_1k_3} \right). \tag{36}
$$

Similarly, we solve the second square equation:

$$
l_1x^2 + l_2x + l_3 = 0.\t\t(37)
$$

where

$$
l_1 = 1;
$$
 $l_2 = \frac{a_3}{2} - \sqrt{d_1};$ $l_3 = \frac{u_1}{2} + \sqrt{d_2};.$ (38)

Its real root is:

$$
x_4 = \frac{1}{2l_1} \left(-l_2 - \sqrt{l_2^2 - 4l_1l_3} \right). \tag{39}
$$

The required support stiffness is the smallest of the previously found real roots:

$$
C_1 = x = \min(x_1, x_4); \quad C_2 = C_1 \cdot n. \tag{40}
$$

Finally, the beam support stiffness can be defined as:

$$
k_1 = C_1 \cdot \frac{E J_{\min}}{l}; \quad k_2 = C_2 \cdot \frac{E J_{\min}}{l}.
$$
 (41)

Calculated support stiffness [\(41\)](#page-6-0) provides the required value of the beam first eigenfrequency at the specified temperature or axial force.

4 Discussion

Here, the solution of this problem is reduced to an algebraic equation of the fourth-order [\(24\)](#page-4-3) or [\(25\)](#page-5-1). The solution of this governing equation can be graphically represented as the intersection of two approximating quadratic functions (Fig. [2\)](#page-7-0). Intersection point determines required stiffness of supports C_1 [\(40\)](#page-6-1).

The error of the stiffness calculation depends mainly on the accuracy of the approximate functions $(18, 19)$ $(18, 19)$ $(18, 19)$. It is quite difficult to provide high accuracy by using quadratic functions $(18, 19)$ $(18, 19)$ $(18, 19)$, since the original curves (11) are highly non-linear. In this paper, the reasonable accuracy is achieved by dividing the full stiffness range into sub-bands, in the same way as it was done in [\[20,](#page-8-6) [21\]](#page-9-0). Comparative calculations showed that acceptable accuracy of calculation, up to 10%, is achieved with the support stiffness of one order; otherwise, error can greatly increase.

Fig. 2. Graphical interpretation of the solution.

5 Conclusion

In this paper, we proposed the solution for the problem of determining the required stiffness of the beam supports, which provide specified values of the first eigenfrequency and first critical force. The least-squares method allowed to approximate the non-linear dependencies by quadratic functions and made it possible to deduce the problem to the fourth-order algebraic equation. We used analytical method to solve the fourth-order equation and got calculation error less 10%, which is quite acceptable for engineering calculations as the first approximation.

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450 I. V. Kudryavtsev et al.

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