Multifractal Analysis of Bioenergy Transport in a Protein Nanomotor

Narmin Sefidkar, Samira Fathizadeh, and Fatemeh Nemati

Abstract The potential application of biological molecules as functional devices has been heralded as the dawn of a new field in biotechnology and medicine. In this regard, molecular motors have attracted the most attention for decades. In the current study, we have studied the bioenergy transfer in a protein chain as a self-powered nanomotor. The effect of different factors on the energy transfer in protein is studied to obtain the best functional condition for the protein machine. The temperature plays the critical role in the control of energy transfer in the system. On the other hand, the external mechanical tension as a vibrator can increase the energy flowing in our system. The chaos theory tools can verify and estimate the results. Generally, one can engineer a self-powered nanomotor based on protein chains and control bioenergy transfer.

Keywords Nanomotor · Protein · Biological energy · Energy transfer · Multifractal analysis

1 Introduction

Biological molecules, an exciting field with continuous and robust growth for about half a century of existence, encompass the medical, biological, chemical, and material sciences [\[1\]](#page-7-0). In the current study, we have studied a protein chain as a biological molecule and its energy transfer. Protein, as the largest macromolecule in the body, is one of the most important components of the life systems that plays vital role in the body. Proteins are composed of amino acid chain sequences that the number of the amino acids can reach several thousand depending on the type of protein [\[2\]](#page-7-1).

In the current study, we investigate the bioenergy transfer in protein molecules as an automated nanomotor. Automated nanomotor uses the bioenergy in the protein as self-burning energy and converts it into autonomous motion [\[3\]](#page-7-2). These nanomotors, which derive their kinetic energy from the biological energy contained in proteins,

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are called protein nanomotors. Protein motors have potential medical applications and can collect, transport, and release drug carriers of various sizes. Internal order in eukaryotic cells is created by protein motors that transport organs and molecules along the cytoskeletal pathways of self-assembled proteins such as tubulin and shuttle actin. Three known families of cytoskeletal protein motors include kinesin, dynein, and myosin [\[4\]](#page-7-3). Kinesins and dyneins, which are the first and second types of protein motors, respectively, move in tubules or microtubules [\[5\]](#page-7-4). But Myosins, the third family of protein motors, move on actin filaments and are responsible for muscle contraction in the heart and skeletal muscles [\[6\]](#page-7-5).

2 Model and Methods

In the current work, the Hamiltonian model used to study the energy transfer in protein molecules is based on the Peng model as follows:

$$
H = H_1 + H_2 + H_3 \tag{1}
$$

where H_1 indicates that a boson-type Frankit exciton is excited in protein molecules using the energy released in ATP hydrolysis written as follows:

$$
H_1 = \sum_{n} \left[\varepsilon_0 a_n^{\dagger} a_n - J \left(a_n^{\dagger} a_{n+1} + a_n a_{n+1}^{\dagger} \right) \right]
$$
 (2)

 H_2 defines the harmonic properties of the remaining amino acids:

$$
H_2 = \sum_{n} \left[\frac{P_n^2}{2m} + \frac{1}{2} W (u_n - u_{n-1})^2 \right]
$$
 (3)

 H_3 introduces the interaction between the two modes of motion [\[7\]](#page-7-6):

$$
H_3 = \sum_{n} \left[\chi_1(u_{n+1} - u_{n-1}) a_n^{\dagger} a_n + \chi_2(u_{n+1} - u_n) \left(a_{n+1}^{\dagger} a_n + a_n^{\dagger} a_{n+1} \right) \right]
$$
 (4)

The parameters used in the calculations are shown in Table [1](#page-2-0) [\[8\]](#page-7-7):

Here, a_n and a_n^{\dagger} are the creation and annihilation operators for exciton. u_n and p_n are the displacement and momentum operators for the amino acid residue at site n, respectively. ε_0 is the energy of the exciton. χ_1 and χ_2 are nonlinear coupling constants. *m* is the mass of amino acid residue. *W* is the elasticity constant of the proteins, and *J* is the dipole–dipole interaction energy between neighboring amino acids [\[9\]](#page-8-0).

In this work, we have used the classical chaos theory. In this case, the evolution equations for the classical part are derived using the Hamiltonian equation **Table 1** Parameters used

 $(p_n = -\frac{\partial H}{\partial q_n})$. Also, the evolution equations for the quantum part are analyzed using the Heisenberg equation $(\dot{a}_n = -\frac{i}{\hbar}[a_n, H])$ as follows:

$$
\ddot{u}_n = \frac{1}{m} \sum_n \left[W(u_{n+1} - 2u_n + u_{n-1}) + \chi_1 \left(a_{n+1}^\dagger a_{n+1} - a_{n-1}^\dagger a_{n-1} \right) + \chi_2 \left(a_{n+1}^\dagger a_n + a_n^\dagger a_{n+1} - a_n^\dagger a_{n-1} - a_{n-1}^\dagger a_n \right) \right]
$$
(5)

$$
\dot{a}_n = -\frac{i}{\hbar} \sum_n \left\{ -J(a_{n-1} + a_{n+1}) + \varepsilon_0 a_n + \chi_1 (u_{n+1} - u_{n-1}) a_n \right. \\ \left. + \chi_2 \left[(u_n - u_{n-1}) a_{n-1} + (u_{n+1} - u_n) a_{n+1} \right] \right\} \tag{6}
$$

To study the energy transfer in protein molecules, we have extracted the energy flux. The energy flux obtains using the continuity equation $(I = -\frac{i}{\hbar} [a_n^{\dagger} a_n, H]$ as follows:

$$
I = -\frac{i}{\hbar} \sum_{n} \left[-J \left(a_{n}^{\dagger} - a_{n-1} - a_{n-1}^{\dagger} a_{n} + a_{n}^{\dagger} a_{n+1} - a_{n+1}^{\dagger} a_{n} \right) - \chi_{2} (u_{n} - u_{n-1}) a_{n-1}^{\dagger} a_{n} + \chi_{2} (u_{n+1} - u_{n}) \left(a_{n}^{\dagger} a_{n+1} - a_{n+1}^{\dagger} a_{n} \right) \right]
$$
\n(7)

3 Results and Discussion

Different factors effect on energy transfer in biological systems. One of the vital parameters on energy transfer in the protein system is the effect of temperature. We have used the Nosé Hoover thermostat to apply the temperature to the system. The evolution equation of thermostat is written as follows [\[10\]](#page-8-1):

$$
\dot{\xi} = \frac{1}{M} \left(\sum_{n} m \dot{u}_n^2 - N K_B T_0 \right) \tag{8}
$$

where ξ describes a thermodynamic coefficient of friction. T_0 is the temperature of the system. K_B is the constant of Boltzmann, and M is the thermostat constant.

A biological system can be affected by mechanical shocks. Therefore, we have investigated the effect of mechanical stress on energy transfer in protein molecules. The Hamiltonian of mechanical stress is written as follows [\[11\]](#page-8-2):

$$
H_{str} = -\delta_{1,n}(u_{n+1} - u_n)F_0 \sin \omega t \tag{9}
$$

where F_0 and ω are the amplitude and frequency of the external mechanical stress.

We can also examine the best range of parameters affecting energy transfer in protein molecules using multifractal system analysis. Multifractal analysis of the system can confirm the obtained results. In addition, it can classify system parameters to determine the desired results.

In this regard, we use the R*e*nyi dimension spectrum, analogous of the free energy and analogous of the thermodynamic specific heat.

3.1 Simultaneous Effect of Ambient Temperature and the Range of Force Exerting Mechanical Stress

Temperature and external stress as the very influential factors in the performance of biological systems can be considered as the control parameters in the energy transfer in the protein system. We have considered the simultaneous effect of ambient temperature and the range of force exerting mechanical stress on the energy flux in the protein system (Fig. [1\)](#page-3-0). The variation of the ambient temperature from 300 to

Fig. 1 The energy flux through the system concerning the simultaneous variation of the amplitude of applied stress and the ambient temperature

330 K and at the same time changing the amplitude of the force from 0 to 2 pN has negligible effect on the amount of energy flux transmitted in the system. In this region, the energy flux fluctuates at a minimum of about 10 to 30 μ eV. But, when the temperature rises above 330 K, in the force range of about 1.2 pN, we can see a considerable peak in the energy flux, so that the energy flux reaches about $100 \mu eV$. Similar to this phenomenon is also observed at a temperature of about 345 K and a force range of about 0.8 pN. Therefore, it can be said that temperatures above 330 and the applied stress are two critical parameters in the transfer of bioenergy in our protein system.

3.2 Simultaneous Effect of Temperature and Applied Stress Frequency

The frequency of applied stress simultaneously with the ambient temperature can be another compelling factor in regulating the energy flux through the system. As shown in Fig. [2,](#page-4-0) at 300–350 K, the energy flux is more than 50 μ V. But there are characteristic frequencies that can significantly increase the energy flux through the system. A frequency of about 0.09 THz at 310 K causes the energy flux of the system to reach about over 200 μeV, and this shows that energy transport in the protein system is highly sensitive to several frequencies that can be used to optimize the performance of the biological system.

Fig. 2 Energy flux relative to a simultaneous variation of applied stress frequency and temperature $(N = 150, F₀ = 3 pN)$

3.3 Multifractal Analysis

We have studied the multifractal spectrum of the energy flux through the protein system by using R*e*nyi dimension. In this regard, we consider the d dimensional phase space of system and divide it into cubes of size land calculate the R*e*nyi dimension (Dq) as follows [\[12\]](#page-8-3):

$$
D_q = \frac{1}{q - 1} \lim_{l \to 0} \frac{\sum_{i}^{M} P_i^q}{\ln l}
$$
 (10)

Figure [3](#page-5-0) shows the Renyi dimension spectrum for four different frequencies of applied stress. D_q shows the descending behavior by increasing q which indicates that the system is multifractal. If we go back to the previous results, we observe that the temperature of 300 K for a force amplitude of 3 pN, the frequencies $\omega = 0.01$ THz, $\omega = 0.03$ THz, $\omega = 0.05$ THz are in the blue region where the energy flux crossing the system is minimal (Fig. [2\)](#page-4-0). Also, the frequency $\omega = 0.09$ THz is in the red region where the energy flux passing through the system is maximum. As shown in Fig. [3,](#page-5-0) for negative q, the R*e*nyi dimensions at 300 K and the force amplitude of 3 pN corresponds to the frequencies $\omega = 0.01$ THz, $\omega = 0.03$ THz, $\omega = 0.05$ THz are coincident with almost the same amount. But the R*e*nyi dimension of the frequency ω = 0.09THz shows a greater value for negative values of q. The R*e*nyi dimension distinguishes the regions with the highest energy flux from the regions with the lowest energy flux.

To have a similarity between a multifractal system and a thermodynamic state, relations are usually equated with their thermodynamic equivalents. We consider $\tau(q)$ as analogous the thermodynamic free energy of the system, and check the

Fig. 3 Renyi dimension spectrum in different values of the stress frequency (T = 310 K, F_0 = 3 pN, $N = 150$

multifractality of the system. The analogous of the free energy is written as follows [\[13\]](#page-8-4):

$$
\tau(q) \equiv (q-1)D_q \tag{11}
$$

The system is a homogenous fractal when $\tau(q)$ have a linear dependence to q. On the other hand, $\tau(q)$ shows a deviation from the linear state when the system is multifractal and the higher deviation from the linear state shows the more multifrac-tality of the system. According the Fig. [4,](#page-6-0) $\tau(q)$ is deviated from the linear state in terms of *q* for all frequencies at the point $q = 0$, and this confirms the multifractality of the system. On the other hand, for negative q , $\tau(q)$ is the same for the frequencies in the blue region. But the frequency $\omega = 0.09$ THz, which is in the red area, has the highest deviation from linear behavior. Thus, the analogous of the thermodynamic free energy shows the distinct regions for negative q at 300 K and for the force range of 3 pN.

On the other hand, the analogous of the specific heat *C(q)* which is obtained from a second-order derivative of the analogous of free energy concerning *q* can be analyze the multifractal behavior of system through the following equation:

$$
\frac{\partial^2 \tau}{\partial q^2} \approx \tau(q+1) - 2\tau(q) + \tau(q-1)
$$
 (12)

Analogous of specific heat is one of the measures to check the multifractality of a system. As shown in Fig. [5,](#page-7-8) the frequencies of the blue region overlap in the peak area. The frequency $\omega = 0.09$ THz, which is in the red region, shows a higher peak. Therefore, a analogous of thermodynamic specific heat also shows the distinct regions and confirms the previous results.

Fig. 4 Analogous of the thermodynamic free energy in multifractal systems for different frequencies of the external stress ($N = 150$, $F_0 = 3$ pN, $T = 310$ K)

Fig. 5 Analogous of thermodynamic specific heat in multifractal systems for different frequencies $(N = 150, F₀ = 3 pN, T = 310 K)$

4 Conclusions

We have studied the bioenergy transfer in a protein system to design a nanomotor. In this regard, we have investigated the effect of various parameters such as temperature and mechanical stress on the energy transfer in protein nanomotors.We have analyzed the multifractal nature of the system using the multifractal analysis methods. Using the multifractal analysis, we are able to confirm the previous results and show the distinct areas.

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