Lecture Notes in Networks and Systems 168

Zakia Hammouch Hemen Dutta Said Melliani Michael Ruzhansky *Editors*

Nonlinear Analysis: Problems, Applications and Computational Methods



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Nonlinear Analysis: Problems, Applications and Computational Methods



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Existence Results for Impulsive Partial Functional Fractional Differential Equation with State Dependent Delay

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Abstract. In this paper, we study the existence of mild solutions of impulsive fractional semilinear differential equation with state dependent delay of order $0 < \alpha < 1$. We shall rely on fixed point theorem for the sum of completely continuous and contraction operators due to Burton and Kirk. An example is given to illustrate the theory.

1 Introduction

Fractional calculus is a generalization of classical differentiation and integration to an arbitrary real order. Fractional calculus is the most well known and valuable branch of mathematics which gives a good framework for biological and physical phenomena, mathematical modeling of engineering, etc. Numerous writings have showed that fractional-order differential equation could provide more methods to deal with complex problem in statistical physics and environmental issues; see the monographs of Abbas et al. [ABN12, ABN15], A. Kilbas et al. [KST06], Podlubny [P93] and Zhou [Z14] and the references therein. On the other hand, the theory of impulsive differential equations has undergone rapid development over the years and played a very important role in modern applied mathematical models of real processes rising in phenomena studied in physics, chemistry, engineering, etc.

Recently, the study of fractional differential equations with impulses has been studied by many authors (see [BHN06, HAM10, LCX12, WFZ11]).

Modivated by work [HGBA13], in this paper, we study the existence of mild solutions for fractional semilinear differential equation of the equation of the form

$${}^{c}D_{t_{k}}^{\alpha}y(t) - Ay(t) = f(t, y_{\rho(t, y_{t})}), t \in J_{k} := (t_{k}, t_{k+1}], k = 0, 1, ..m,$$
(1)

$$\Delta y |_{y=y_k} = I_k(y_{t_k}) \quad k = 1, ..., m,$$
(2)

$$y(t) = \phi(t), \ t \in (-\infty, 0].$$
 (3)

where ${}^{c}D_{t_{k}}^{\alpha}$ is caputo fractional derivative of order $0 < \alpha < 1, A : D(A) \subset E \to E$ is the bounded linear operator of an α - resolvent family $S_{\alpha}(t) : t \geq 0$ defined on a Banach

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space $E, f: J \times \mathcal{D} \to E$ is a given function, $\mathcal{D} = \{ \psi : (-\infty, 0] \to E, \psi \text{ is continuous} \text{ every where except for a finite number of points$ *s* $at which <math>\psi(s^-), \psi(s^+) \text{ exist and } \psi(s^-) = \psi(s) \}$, $\phi \in D, (0 < r < \infty)$, $I_k : E \to E$, $(k = 0, 1, ..., m + 1), 0 = t_0 < t_1 < ..., < t_m < t_{m+1} = b$, $\Delta y|_{y=y_k} = y(t_k^+) - y(t_k^-)$, where $y(t_k^+) = \lim_{h\to 0^+} y(t_k + h)$ and $y(t_k^-) = \lim_{h\to 0^+} y(t_k - h)$ represent the right and left limits of y(t) at $t = t_k$, respectively. We denote by y_t the element of \mathcal{D} defined by $y_t(\theta) = y(t + \theta), \theta \in (-\infty, 0]$. Here y_t represents the history up to the present time *t* of the state y(.). We assume that the histories y_t belongs to some abstract phase \mathcal{D} , to specified later, and $\phi \in \mathcal{D}$. This paper is organized as follow, in Sect. 2 we introduce some preliminaries that will be used in the sequel, in Sect. 3 we give definition to the mild solution of problem 1–3 result inspired by works [HGBA13, HL20], also the proof of our main results is given. Finally, an example is included in Sect. 4.

2 Preliminaries

In this Section, we state some notations, definitions and properties which be used throughout this paper.

Let *E* be a Banach space endowed with the norm $\|.\|$, and L(E) represents the Banach space of all bounded linear operators from *E* into *E* and the corresponding norm $\|.\|_{L(E)}$.

C(J,E) is the Banach space of all continuous functions from J to E with the norm

$$||u||_{C(J,E)} = \sup\{|u(t)| : t \in J\},\$$

 $L^1[J, E]$ is the Banach space of measurable functions $u: J \longrightarrow E$ which are Bochner integrable normed by

$$||u||_{L^1} = \int_0^b |u(t)| dt.$$

Definition 1. A family $(S_{\alpha}(t))_{t>0} \subset l(E)$ of bounded linear operators in *E* is called an α - resolvent operator function generating by *A* if the following conditions hold:

a) $(S_{\alpha}(t))_{t>0}$ is strong continuous on R_+ and $S_{\alpha}(0) = I$;

b) $S_{\alpha}(t)D(A) \subset D(A)$ and $AS_{\alpha}(t)x = S_{\alpha}(t)Ax$ for all $x \in D(A)$ and t > 0;

c) For all $x \in E$, $I_t^{\alpha} S_{\alpha}(t) x \in D(A)$ and

$$S_{\alpha}(t)x = x + AI_t^{\alpha}S_{\alpha}(t)x, \ t > 0;$$

d) $x \in D(A)$ and Ax = y if and only if

$$S_{\alpha}(t)x = x + AI_t^{\alpha}S_{\alpha}(t)x, \ t > 0;$$

e) A is closed and densely defined

The generator A of $(S_{\alpha}(t))_{t>0}$ is defined by:

$$D(A) := \{ x \in E : \lim_{t \to 0^+} \frac{S_{\alpha}(t)x - x}{\psi_{\alpha+1}(t)} \ exists \},\$$

and

$$Ax = \lim_{t \to 0^+} \frac{S_{\alpha}(t)x - x}{\psi_{\alpha+1}(t)}, \ x \in D(A),$$

where $\psi_{\alpha}(t) = \frac{t^{\alpha-1}}{\Gamma(\alpha)}$ for t > 0 and $\psi_{\alpha}(t) = 0$ for $t \le 0$ and $\psi_{\alpha}(t) \longrightarrow \delta(t)$ as $\alpha \longrightarrow 0$, where the function delta is defined by:

$$\delta_a: D(\Omega) \longrightarrow R; \phi \to \phi(a),$$

and

$$D(\Omega) = \{ \phi \in C^{\infty}(\Omega) : supp \phi \subset \Omega \text{ is compact} \}$$

Definition 2. An α -ROF $(S_{\alpha}(t))_{t>0}$ is called analytic, if the function $S_{\alpha}(): \mathbb{R} \to \mathbb{R}$ l(X) admits analytic extension to a sector $\Sigma(0, \theta_0)$ for some $0 < \theta_0 \le \frac{\pi}{2}$. An analytic α -ROF (S_{α}) is said to be of analyticity type (ω_0, θ_0) if for each $\theta < \overline{\theta_0}$ and $\omega > \omega_0$ there exists $M_1 = M_1(\omega, \theta)$ such that $||S_{\alpha}(z)|| \leq M_1 e^{\omega Rez}$ for $z \in \Sigma(0, \theta)$ where Rezdenotes the real part of z and $\Sigma(\omega, \theta) := \{\lambda \in C : |arg(\lambda - \omega)| < \theta, \quad \omega, \theta \in R\}$

Definition 3. An α -ROF $(S_{\alpha}(t))_{t\geq 0}$ is called compact for t > 0 if for every t > 0, $S_{\alpha}(t)$ is a compact operator.

Theorem 1. Let A generate a compact analytic semigroup $T(t)_{t\geq 0}$ then for any α it also generates a compact analytic resolvent family $(S_{\alpha}(t))_{t\geq 0}$.

Lemma 1. Assume that α -ROF $(S_{\alpha}(t))_{t\geq 0}$ is compact for t > 0 and analytic of type (ω_0, θ_0) . Then the following assertions hold:

- 1. $\lim_{h \to 0} ||S_{\alpha}(t+h) S_{\alpha}(t)|| = 0$, for t > 0. 2. $\lim_{h \to 0^+} ||S_{\alpha}(t+h) S_{\alpha}(h)S_{\alpha}(t)|| = 0$, for t > 0.

Definition 4. An α -ROF $(S_{\alpha}(t))_{t>0}$ is said to be exponentially bounded if there exist constants $M \ge 1$, $\omega \ge 0$ such that

$$||S_{\alpha}(t)|| \leq M e^{\omega t} \quad for \ t \geq 0.$$

in this case we write $A \in C_{\alpha}(M, \omega)$.

Definition 5. The fractional integral operator I^{α} of order $\alpha > 0$ of a continuous function f(t) is defined by

$$I_t^{\alpha} f(t) := \frac{1}{\Gamma(\alpha)} \int_0^t (t-s)^{\alpha-1} f(s) ds$$

Observe that $I_t^{\alpha} f(t) = f(t) * \psi_{\alpha}(t)$, where $\psi_{\alpha}(t) = \frac{t^{\alpha-1}}{\Gamma(\alpha)}$ for t > 0 and $\psi_{\alpha}(t) = 0$ for $t \leq 0$ and $\psi_{\alpha}(t) \longrightarrow \delta(t)$ as $\alpha \longrightarrow 0$.

Definition 6. The α -Riemann-Liouville fractional-order derivative of the function f, is defined by

$$D_a^{\alpha}f(t) = \frac{1}{\Gamma(n-\alpha)} \frac{d^n}{dt^n} \int_a^t (t-s)^{n-\alpha-1} f(s) ds.$$

where $n = [\alpha] + 1$ and $[\alpha]$ denotes the integer part of α .

Definition 7 [P93]. For a function f defined on the interval [a,b], the Caputo fractional order derivative of order α of f, is defined by

$$\binom{c}{a+}D_t^{\alpha}f)(t) = \frac{1}{\Gamma(n-\alpha)}\int_0^t (t-s)^{n-\alpha-1}f^{(n)}(s)ds,$$

Where $n = [\alpha] + 1$.

Therefore, for $0 < \alpha < 1$, $n = [\alpha] + 1 = 1$ and for a = 0, the Caputo's fractional derivative for $t \in [0, b]$ is given by

$$\binom{c}{_{0}}D_{t}^{\alpha}f(t) = \frac{1}{\Gamma(1-\alpha)}\int_{0}^{t}(t-s)^{-\alpha}f'(s)ds.$$

In this paper, we will employ an axiomatic definition for the phase space D which is similar to those introduced by Hale and Kato [HK78]. Specifically, D will be a linear space of functions mapping $] - \infty, b]$ into E endowed with a semi-norm $\|.\|_D$, and satisfies the following axioms:

- (A1) There exist a positive constant *H* and functions $K(\cdot)$, $M(\cdot) : R^+ \to R^+$ with *K* continuous and *M* locally bounded, such that for any b > 0, if $x : (-\infty, b] \to E$, $x \in D$, and $x(\cdot)$ is continuous on [0,b], then for every $t \in [0,b]$ the following conditions hold:
 - (i) x_t is in D;
 - (ii) $|x(t)| \le H ||x_t||_D$;
 - (iii) $||x_t||_D \le K(t) \sup\{|x(s)|: 0 \le s \le t\} + M(t) ||x_0||_D$, and H, K and M are independent of $x(\cdot)$.

Denote

$$K_b = \sup\{K(t): t \in J\}$$
 and $M_b = \sup\{M(t): t \in J\}.$

 (A_2) The space *D* is complete.

Example 1. Let $h(.): (-\infty, -r] \to R$ be a positive Lebesgue integrable function and $D := PC_r \times L^2(h; E), r \ge 0$, be the space formed of all classes of functions $\varphi: (-\infty, 0] \to E$ such that $\varphi|_{[-r,0]} \in PC([-r,0], E), \varphi(.)$ is Lebesgue-measurable on $(-\infty, -r]$ and $h|\varphi|^p$ is Lebesgue integrable on $(-\infty, -r]$. the semi-norm in $\|.\|_D$ is defined by

$$\|\varphi\|_{D} = \sup_{\theta \in [-r,0]} \|\varphi(\theta)\| + \left(\int_{-\infty}^{-r} h(\theta) \|\varphi(\theta)\|^{p} d\theta\right)^{1/p},$$
(4)

Assume that h(.) satisfies conditions (g–6) and (g–7) in the terminology of [HMN91]. proceeding as in the proof of [[HMN91]. Theorem 1.3.8] it follows that \mathscr{D} is a phase space which verifies the axioms (A1)–(A2) and (A3). Moreover, when r = 0 this space coincides with $C^0 \times L^2(h, E)$ and the parameters $H = 1; M(t) = \gamma(-t)^{1/2}$ and $K(t) = 1 + \left(\int_{-r}^0 h(\xi) d\xi\right)^{1/2}$, for $t \ge 0$ (see [HMN91]).

Definition 8. A map $f: [0,b] \times \mathscr{D} \to E$ is said to be carathéodory if

- 1. the function $t \mapsto f(t, y)$ is measurable for each $y \in \mathcal{D}$;
- 2. the function $t \mapsto f(t, y)$ is continuous for almost all $t \in J_k := (t_k, t_{k+1}], k = 0, 1, ...m$.

In order to define the mild solution 1-3, we consider the following space

$$PC(J,E) = \{y : [0,b] \longrightarrow E : y \text{ is continous at } t \neq t_k, y(t_k^-) = y(t_k),$$

and $y(t_k^+)$ exists, for all k = 1, ..., m

which is a Banach space with the norm

$$||y|| = max\{||y_k||_{\infty}; k = 1, 2, .., m\},\$$

and

$$D_b = \{y:] -\infty, b] \longrightarrow E: y|_{]-\infty,0]} \in D \text{ and } y|_J \in PC(J, E)\}.$$

Let $\|.\|_b$ be the semi norm in D_b defined by

$$||y||_b = ||y_0||_{D_b} + \sup\{|y(s)|: 0 \le s \le b\}, y \in D_b.$$

Let us introduce the definition of Caputo's derivative in each interval $(t_k, t_{k+1}], k = 0, ..., m$,

$$(^{c}D_{t_{k}}^{\alpha}f)(t)=\frac{1}{\Gamma(1-\alpha)}\int_{t_{k}}^{t}(t-s)^{-\alpha}f'(s)ds.$$

3 Main Result

Before starting and proving our main result, we give the meaning of mild solution of our problem 1-3.

Definition 9. A function $y \in PC((-\infty, b], E)$ is said to be mild solution of our problem if $y(t) = \phi(t)$, for all $t \in (-\infty, 0], \Delta y|_{y=y_k} = I_k(y_{t_i}), k = 1, 2, ..., m$ and such that *y* satisfies the following integral equation:

$$y(t) = \begin{cases} S_{\alpha}(t)\phi(0) + \int_{0}^{t} S_{\alpha}(t-s)f(s,y_{\rho(s,y_{s})})ds; & \text{if } t \in [0,t_{1}], \\ S_{\alpha}(t-t_{k})\prod_{i=1}^{k} S_{\alpha}(t_{i}-t_{i-1})\phi(0) \\ + \sum_{i=1}^{k} \int_{t_{i-1}}^{t_{i}} S_{\alpha}(t-t_{k})\prod_{j=i}^{k-1} S_{\alpha}(t_{j+1}-t_{j})S_{\alpha}(t_{i}-s)f(s,y_{\rho(s,y_{s})}) \\ + \int_{t_{k}}^{t} S_{\alpha}(t-s)f(s,y_{\rho(s,y_{s})})ds \\ + \sum_{i=1}^{k} S_{\alpha}(t-t_{k})\prod_{j=i}^{k-1} S_{\alpha}(t_{j+1}-t_{j})I_{i}(y_{t_{i}}); & \text{if } t \in (t_{k},t_{k+1}]. \end{cases}$$

Set

$$\mathscr{R}(\rho^{-}) = \{\rho(s,\varphi) : (s,\varphi) \in J \times \mathscr{D}, \rho(s,\varphi) \le 0\}$$

We always assume that $\rho: J \times \mathscr{D} \to (-\infty, b]$ is continuous. Additionally, we introduce the following hypothesis:

• (H_{φ}) The function $t \to \varphi_t$ is continuous from $\mathscr{R}(\rho^-)$ into \mathscr{D} and there exists a continuous and bounded function $L^{\phi} : \mathscr{R}(\rho^-) \to (0, \infty)$ such that

 $\|\phi_t\|_{\mathscr{D}} \leq L^{\phi}(t)\|\phi\|_{\mathscr{D}}, \qquad ext{for every} \quad t\in \mathscr{R}(
ho^-).$

Lemma 2 ([*HPL06*]). If $y:] -\infty, b] \longrightarrow E$ is a function such that $y_0 = \phi$, then

$$\|y_t\|_D \le (M_b + L^{\phi}) \|\phi\|_{\mathscr{D}} + K_b \sup\{|y(s)|; s \in [0, \max\{0, t\}]\},\$$

where $M_b = \sup_{t \in J} M(t)$, $K_b = \sup_{t \in J} K(t)$ and $L^{\phi} = \sup_{t \in \mathscr{R}(\rho^-)} L^{\phi}(t)$.

Our main result in this section is based upon the following fixed point theorem due to Burton and Kirk [BK98].

Theorem 2 ([*BK98*]). Let X be a Banach space and $A, B : X \longrightarrow X$ be two operators satisfying:

- 1. A is a contraction,
- 2. B is completely continuous,

Then, either;

- 1. the operator equation y = Ay + By has a solution, or
- 2. the set $\Upsilon = \{u \in X : \lambda A(\frac{u}{\lambda}) + \lambda B(u) = u, \lambda \in (0, 1)\}$ is unbounded.

We introduce the following hypotheses:

(*H*₁) A generate a compact and analytic α -ROF $(S_{\alpha}(t))_{t\geq 0}$ which is exponentially bounded i.e there exist constants $M \geq 1, \omega \geq 0$ such that

$$\|S_{\alpha}(t)\| \leq Me^{\omega t}; \quad t \geq 0.$$

 (H_2) The functions $I_k: E \longrightarrow E$ are Lipschitz. Let M_k , for k = 1, 2, 3, ..., m, be such that

$$||I_k(y) - I_k(x)|| \le M_k ||y - x||$$
; for each $y, x \in E$.

- (H_3) the function $f: J \times D \longrightarrow E$ is Caratheodory.
- (*H*₄) There exists a function $p \in L^1(J, \mathbb{R}_+)$ and a continuous nondeceasing function $\psi : [0, +\infty[\longrightarrow [0, +\infty[$ such that

$$|f(t,y)| \le p(t)\psi(||y||)_D,$$

a.e, $t \in J$, for all $y \in D$, with

$$\int_{C_0}^{\infty} \frac{du}{\psi(u)} = \infty$$

and

$$\int_{C_3}^{\infty} \frac{du}{\psi(u)} = \infty,$$

•

where

$$C_0 = C$$
, $C_3 = min(C_1, C_2)$,
 $C = (M_b + L^{\phi} + K_b M e^{\omega b}) \|\phi\|_{\mathcal{D}_b^0}$,

$$\begin{split} C_1 &= \frac{K_b \Big(M^{k+1} e^{\omega b} |\phi(0)| + \sum_{i=1}^k M^{k-i+1} e^{\omega(b-t_i)} (|I_i(0)| + C) \Big)}{1 - K_b \sum_{i=1}^k M^{k-i+1} e^{\omega(b-t_i)} M_i} \\ &+ \frac{K_b \sum_{i=1}^k M^{k-i+2} e^{\omega(b-t_{k-1})} \int_{t_{i-1}}^{t_i} e^{-\omega s} p(s) \psi(\mu(s)) ds}{1 - K_b \sum_{i=1}^k M^{k-i+1} e^{\omega(b-t_i)} M_i} + C \quad, \end{split}$$

$$C_2 = \frac{K_b M e^{\omega b}}{(1 - K_b \sum_{i=1}^k M^{k-i+1} e^{\omega(b-t_i)} M_i)}$$

Theorem 3. Assume that Hypotheses $(H_{\varphi}), (A_1), (A_2), (H_1), (H_4)$ are satisfied with

$$K_b \sum_{i=1}^k M^{k-i+1} e^{\omega(b-t_i)} M_i < 1,$$

then the problem (1.1)–(1.3) has at least one mild solution on $] - \infty, b]$.

Proof. Transform the problem (1.1)–(1.3) into a fixed point problem. Consider the operator $N : \mathscr{D}_b \to \mathscr{D}_b$ defined by

$$N(y)(t) = \begin{cases} \phi(t); & t \in (-\infty, 0], \\ S_{\alpha}(t)\phi(0) + \int_{0}^{t} S_{\alpha}(t-s)f(s, y_{\rho(s, y_{s})})ds; & t \in [0, t_{1}], \\ S_{\alpha}(t-t_{k})\prod_{i=1}^{k} S_{\alpha}(t_{i}-t_{i-1})\phi(0) \\ + \sum_{i=1}^{k} \int_{t_{i-1}}^{t_{i}} S_{\alpha}(t-t_{k})\prod_{j=i}^{k-1} S_{\alpha}(t_{j+1}-t_{j})S_{\alpha}(t_{i}-s)f(s, y_{\rho(s, y_{s})}) \\ + \int_{t_{k}}^{t} S_{\alpha}(t-s)f(s, y_{\rho(s, y_{s})})ds \\ + \sum_{i=1}^{k} S_{\alpha}(t-t_{k})\prod_{j=i}^{k-1} S_{\alpha}(t_{j+1}-t_{j})I_{i}(y_{t_{i}}); & t \in (t_{k}, t_{k+1}]. \end{cases}$$

Let $x(.):] - \infty, b] \longrightarrow E$, be the function defined by

$$x(t) = \begin{cases} \phi(t), & \text{if } t \in] -\infty, 0], \\ S_{\alpha}(t)\phi(0), & \text{if } t \in [0, t_1], \\ 0, & \text{if } t \in (t_k, t_{k+1}]. \end{cases}$$

Then $x_0 = \phi$. For each $z \in \mathscr{D}_b$ with z(0) = 0, we denote by \overline{z} the function defined by

$$\bar{z}(t) = \begin{cases} 0, & if \ t \in] -\infty, 0], \\ \int_0^t S_\alpha(t-s) f(s, x_{\rho(s,x_s+\bar{z}_s)} + \bar{z}_{\rho(s,x_s+\bar{z}_s)}) + \bar{z}_{\rho(s,x_s+\bar{z}_s)}) ds, & if \ t \in [0,t_1], \\ S_\alpha(t-t_k) \prod_{i=1}^k S_\alpha(t_i-t_{i-1}) \phi(0) \\ + \sum_{i=1}^k \int_{t_{i-1}}^{t_i} S_\alpha(t-t_k) \prod_{j=i}^{k-1} S_\alpha(t_{j+1}-t_j) S_\alpha(t_i-s) f(s, x_{\rho(s,x_s+\bar{z}_s)} + \bar{z}_{\rho(s,x_s+\bar{z}_s)}) ds \\ + \int_{t_k}^t S_\alpha(t-s) f(s, y_{\rho(s,y_s)}) ds \\ + \sum_{i=1}^k S_\alpha(t-t_k) \prod_{j=i}^{k-1} S_\alpha(t_{j+1}-t_j) I_i(x_{t_i} + \bar{z}_{t_i}), & if \ t \in (t_k, t_{k+1}]. \end{cases}$$
(5)

If y(.) satisfies (3), we can decompose it as y(t) = x(t) + z(t), $0 \le t \le b$, which implies $y_t = z_t + x_t$ for every $0 \le t \le b$ and the function z(.) satisfies

$$z^{*}(t) = \begin{cases} 0, & \text{if } t \in] -\infty, 0], \\ z(t), & \text{if } t \in [0, b]. \end{cases}$$

where

$$z(t) = \begin{cases} \int_0^t S_\alpha(t-s) f(s, x_{\rho(s, x_s + \overline{z}_s)} + \overline{z}_{\rho(s, x_s + \overline{z}_s)}) ds, & \text{if } t \in [0, t_1], \\ S_\alpha(t-t_k) \prod_{i=1}^k S_\alpha(t_i - t_{i-1}) \phi(0) \\ + \sum_{i=1}^k \int_{t_{i-1}}^{t_i} S_\alpha(t-t_k) \prod_{j=i}^{k-1} S_\alpha(t_{j+1} - t_j) S_\alpha(t_i - s) f(s, x_{\rho(s, x_s + \overline{z}_s)} + \overline{z}_{\rho(s, x_s + \overline{z}_s)}) \\ + \int_{t_k}^t S_\alpha(t-s) f(s, x_{\rho(s, x_s + \overline{z}_s)} + \overline{z}_{\rho(s, x_s + \overline{z}_s)}) ds \\ + \sum_{i=1}^k S_\alpha(t-t_k) \prod_{j=i}^{k-1} S_\alpha(t_{j+1} - t_j) I_i(x_{t_i} + \overline{z}_{t_i}), & \text{if } t \in (t_k, t_{k+1}]. \end{cases}$$

Set

$$\mathscr{D}_b^0 := \{ z \in \mathscr{D}_b : z_0 = 0 \}.$$

and let $\|.\|_b$ be the seminorm in \mathscr{D}_b^0 defined by

$$||z||_b = ||z_0||_+ \sup\{|z(t)| : 0 \le t \le b\}$$

= sup{|z(t)| : 0 \le t \le b}.

$$\mathscr{D}_{b}^{0}$$
 is Banach space with the norm $\|.\|_{b}$.

Transform the problem 1-3 into a fixed point problem. Consider the two operators

$$\mathscr{A}, \mathscr{B}: \mathscr{D}_b^0 \longrightarrow \mathscr{D}_b^0$$

defined by

$$\mathscr{A}z(t) = \begin{cases} 0, & \text{if } t \in [0, t_1], \\ S_{\alpha}(t - t_k) \prod_{i=1}^k S_{\alpha}(t_i - t_{i-1})\phi(0) \\ + \sum_{i=1}^k S_{\alpha}(t - t_k) \prod_{j=i}^{k-1} S_{\alpha}(t_{j+1} - t_j) I_i(x_{t_i} + \overline{z}_{t_i}), & \text{if } t \in (t_k, t_{k+1}] \end{cases}$$

and

$$\mathscr{B}z(t) = \begin{cases} \int_{0}^{t} S_{\alpha}(t-s) f(s, x_{\rho(s, x_{s} + \overline{z}_{s})} + \overline{z}_{\rho(s, x_{s} + \overline{z}_{s})}) ds, & \text{if } t \in [0, t_{1}], \\ \sum_{i=1}^{k} \int_{t_{i-1}}^{t_{i}} S_{\alpha}(t-t_{k}) \prod_{j=i}^{k-1} S_{\alpha}(t_{j+1} - t_{j}) S_{\alpha}(t_{i} - s) f(s, x_{\rho(s, x_{s} + \overline{z}_{s})} + \overline{z}_{\rho(s, x_{s} + \overline{z}_{s})}) \\ + \int_{t_{k}}^{t} S_{\alpha}(t-s) f(s, x_{\rho(s, x_{s} + \overline{z}_{s})} + \overline{z}_{\rho(s, x_{s} + \overline{z}_{s})}) ds, & \text{if } t \in (t_{k}, t_{k+1}]. \end{cases}$$

Then the problem of finding the solution of the problem 1–3 is reduced to finding the solution of operator equation $\mathscr{A}z(t) + \mathscr{B}z(t) = z(t), t \in (-\infty, b]$, we shall that the operators \mathscr{A} and \mathscr{B} satisfy all the conditions of theorem 3.

We give the proof into a sequence of steps.

Step 1: *B* is continuous.

Let $(z^n)_{n\geq 0}$ be a sequence such that $z^n \longrightarrow z$ in \mathscr{D}_b^0 . Since f satisfies (H3), we get

$$f(s, x_s + \overline{z}_s^n) \to f(s, x_s + \overline{z}_s)$$
 as $n \to \infty$.

Then

1. For $t \in [0, t_1]$, we have

$$\begin{split} |\mathscr{B}(z^{n})(t) - \mathscr{B}(z)(t)| \\ &= |\int_{0}^{t} S_{\alpha}(t-s) [f(s,x_{\rho(s,x_{s}+\overline{z}_{s}^{n})} + \overline{z}_{\rho(s,x_{s}+\overline{z}_{s}^{n})}^{n}) - f(s,x_{\rho(s,x_{s}+\overline{z}_{s})} + \overline{z}_{\rho(s,x_{s}+\overline{z}_{s})}^{n})] ds \\ &\leq \int_{0}^{t} ||S_{\alpha}(t-s)|| |f(s,x_{\rho(s,x_{s}+\overline{z}_{s}^{n})} + \overline{z}_{\rho(s,x_{s}+\overline{z}_{s}^{n})}^{n}) - f(s,x_{\rho(s,x_{s}+\overline{z}_{s})} + \overline{z}_{\rho(s,x_{s}+\overline{z}_{s})})] ds \\ &\leq Me^{\omega t} \int_{0}^{t} e^{-\omega s} |f(s,x_{\rho(s,x_{s}+\overline{z}_{s}^{n})} + \overline{z}_{\rho(s,x_{s}+\overline{z}_{s}^{n})}^{n}) - f(s,x_{\rho(s,x_{s}+\overline{z}_{s})} + \overline{z}_{\rho(s,x_{s}+\overline{z}_{s})})| ds \longrightarrow 0. \end{split}$$

2. For
$$t \in (t_k, t_{k+1}]$$
,

$$|\mathscr{B}(z^n)(t) - \mathscr{B}(z)(t)|$$

$$\begin{split} &= |\sum_{i=1}^{k} \int_{l_{i-1}}^{l_{i}} S_{\alpha}(t-t_{k}) \prod_{j=i}^{k-1} S_{\alpha}(t_{j+1}-t_{j}) S_{\alpha}(t_{i}-s) [f(s,x_{\rho(s,x_{s}+\overline{z}_{s}^{n})} + \overline{z}_{\rho(s,x_{s}+\overline{z}_{s}^{n})}^{n}) - f(s,x_{\rho(s,x_{s}+\overline{z}_{s})} + \overline{z}_{\rho(s,x_{s}+\overline{z}_{s})}^{n})] ds + \int_{l_{k}}^{l} S_{\alpha}(t-s) [f(s,x_{\rho(s,x_{s}+\overline{z}_{s}^{n})} + \overline{z}_{\rho(s,x_{s}+\overline{z}_{s}^{n})}^{n}) - f(s,x_{\rho(s,x_{s}+\overline{z}_{s})} + \overline{z}_{\rho(s,x_{s}+\overline{z}_{s})}^{n})] ds \\ &\leq \sum_{i=1}^{k} \int_{l_{i-1}}^{l_{i}} ||S_{\alpha}(t-t_{k})|| \prod_{j=i}^{k-1} ||S_{\alpha}(t_{j+1}-t_{j})|| ||S_{\alpha}(t_{i}-s)|| \times ||f(s,x_{\rho(s,x_{s}+\overline{z}_{s}^{n})} + \overline{z}_{\rho(s,x_{s}+\overline{z}_{s}^{n})}^{n}) \\ &- f(s,x_{\rho(s,x_{s}+\overline{z}_{s})} + \overline{z}_{\rho(s,x_{s}+\overline{z}_{s})})] ds + \int_{l_{k}}^{l} ||S_{\alpha}(t-s)|| |[f(s,x_{\rho(s,x_{s}+\overline{z}_{s}^{n})} + \overline{z}_{\rho(s,x_{s}+\overline{z}_{s}^{n})}^{n}) \\ &- f(s,x_{\rho(s,x_{s}+\overline{z}_{s})} + \overline{z}_{\rho(s,x_{s}+\overline{z}_{s})})] ds \\ &\leq \sum_{i=1}^{k} \int_{l_{i-1}}^{l_{i}} Me^{\omega(t-l_{k})} \prod_{j=i-1}^{k-1} Me^{\omega(l_{j+1}-l_{j})} Me^{\omega(t-s)} |[f(s,x_{\rho(s,x_{s}+\overline{z}_{s}^{n})} + \overline{z}_{\rho(s,x_{s}+\overline{z}_{s}^{n})}^{n}) \\ &- f(s,x_{\rho(s,x_{s}+\overline{z}_{s})} + \overline{z}_{\rho(s,x_{s}+\overline{z}_{s})})] ds + \int_{l_{k}}^{l} Me^{\omega(t-s)} |[f(s,x_{\rho(s,x_{s}+\overline{z}_{s}^{n}) + \overline{z}_{\rho(s,x_{s}+\overline{z}_{s}^{n})}^{n}) \\ &- f(s,x_{\rho(s,x_{s}+\overline{z}_{s})} + \overline{z}_{\rho(s,x_{s}+\overline{z}_{s})})] ds \\ &\leq \sum_{i=1}^{k} \int_{l_{i-1}}^{l_{i-1}} Me^{\omega(t-l_{k})} [Me^{\omega(t_{i-1}-l_{i})} \times Me^{\omega(t_{i-2}-l_{i+1})} \times Me^{\omega(t_{i+3}-t_{i+2})} \\ &\times \dots \times Me^{\omega(t-k_{k})} [Me^{\omega(t_{i-1}s)} |[f(s,x_{\rho(s,x_{s}+\overline{z}_{s})} + \overline{z}_{\rho(s,x_{s}+\overline{z}_{s}^{n})}] ds \\ &+ \int_{l_{k}}^{l} Me^{\omega(t-s)} |[f(s,x_{\rho(s,x_{s}+\overline{z}_{s})} + \overline{z}_{\rho(s,x_{s}+\overline{z}_{s}^{n})}) \\ &- f(s,x_{\rho(s,x_{s}+\overline{z}_{s})} + \overline{z}_{\rho(s,x_{s}+\overline{z}_{s}^{n})})]ds \\ &+ \int_{l_{k}}^{l} Me^{\omega(t-s)} |[f(s,x_{\rho(s,x_{s}+\overline{z}_{s})} + \overline{z}_{\rho(s,x_{s}+\overline{z}_{s}^{n})})] ds \\ &+ Me^{\omega(t-s)} [f(s,x_{\rho(s,x_{s}+\overline{z}_{s})} + \overline{z}_{\rho(s,x_{s}+\overline{z}_{s}^{n})}) \\ &- f(s,x_{\rho(s,x_{s}+\overline{z}_{s})} + \overline{z}_{\rho(s,x_{s}+\overline{z}_{s}^{n})}) - f(s,x_{\rho(s,x_{s}+\overline{z}_{s}^{n})} + \overline{z}_{\rho(s,x_{s}+\overline{z}_{s}^{n})})] ds \\ &+ \int_{l_{k}}^{l} Me^{\omega(t-s)} [H^{s}(s,x_{\rho(s,x_{s}+\overline{z}_{s})} + \overline{z}_{\rho(s,x_{$$

We get

$$\|\mathscr{B}(z^n)(t) - \mathscr{B}(z)(t)\|_{\mathscr{D}^0_b} \longrightarrow 0.$$

as $n \longrightarrow +\infty$.

This means that \mathcal{B} is continuous.

Step 2: \mathscr{B} maps bounded sets into bounded sets in \mathscr{D}_b^0 .

A linear operator $\mathscr{B}: \mathscr{D}_b^0 \longrightarrow \mathscr{D}_b^0$ is bounded if only it maps bounded sets into bounded sets; i.e it is enough to show that for any q > 0, there exists a positive constant $l_k; k = 1, 2, ..., m$ such that for each $z \in B_q = \{z \in \mathscr{D}_b^0 : ||z|| \le q\}$, we have $||\mathscr{B}(z)|| \le l_k$.

Let $z \in B_q$. Then,

$$\begin{split} |\mathscr{B}z(t)| &\leq \begin{cases} \int_{0}^{t} \|S_{\alpha}(t-s)\| \|f(s,x_{\rho(s,x_{s}+\overline{z}_{s})}+\overline{z}_{\rho(s,x_{s}+\overline{z}_{s})})\|ds, & \text{if } t \in [0,t_{1}], \\ \sum_{i=1}^{k} \int_{t_{i-1}}^{t_{i}} \|S_{\alpha}(t-t_{k})\| \\ &\times \prod_{j=i}^{k-1} \|S_{\alpha}(t_{j+1}-t_{j})\| \|S_{\alpha}(t_{i}-s)\| \|f(s,x_{\rho(s,x_{s}+\overline{z}_{s})}+\overline{z}_{\rho(s,x_{s}+\overline{z}_{s})})\|ds \\ &+ \int_{t_{k}}^{t} \|S_{\alpha}(t-s)\| \|f(s,x_{\rho(s,x_{s}+\overline{z}_{s})}+\overline{z}_{\rho(s,x_{s}+\overline{z}_{s})})\|ds, & \text{if } t \in (t_{k},t_{k+1}]. \end{cases} \\ |\mathscr{B}z(t)| &\leq \begin{cases} \int_{0}^{t} \|S_{\alpha}(t-s)\| \|p(s)\psi(\|x_{\rho(s,x_{s}+\overline{z}_{s})}+\overline{z}_{\rho(s,x_{s}+\overline{z}_{s})}\|)ds, & \text{if } t \in [0,t_{1}], \\ \sum_{i=1}^{k} \int_{t_{i-1}}^{t_{i}} \|S_{\alpha}(t-t_{k})\| \\ &\times \prod_{j=i}^{k-1} \|S_{\alpha}(t_{j+1}-t_{j})\| \|S_{\alpha}(t_{i}-s)\| \|p(s)\psi(\|x_{\rho(s,x_{s}+\overline{z}_{s})}+\overline{z}_{\rho(s,x_{s}+\overline{z}_{s})}\|)ds, & \text{if } t \in (t_{k},t_{k+1}]. \end{cases} \end{cases}$$

Using Lemma 3.1, we get

$$\|x_{\rho(s,x_s+\bar{z}_s)}+\bar{z}_{\rho(s,x_s+\bar{z}_s)}\|_{\mathscr{D}_b^0} \leq K_b M e^{\omega t_1} |\phi(0)| + (M_b + L^{\phi}) \|\phi\|_{\mathscr{D}_b^0} + K_b |z(s)|.$$

Then

$$\|x_{\rho(s,x_{s}+\bar{z}_{s})}+\bar{z}_{\rho(s,x_{s}+\bar{z}_{s})}\|_{\mathscr{D}_{b}^{0}} \leq (M_{b}+L^{\phi}+K_{b}Me^{\omega b})\|\phi\|_{\mathscr{D}_{b}^{0}}+K_{b}q=q^{*}$$

Set $C = (M_b + L^{\phi} + K_b M e^{\omega b}) \|\phi\|_{\mathscr{D}_b^0}$. Then we obtain

$$\|x_{\rho(s,x_s+\overline{z}_s)}+\overline{z}_{\rho(s,x_s+\overline{z}_s)}\|_{\mathscr{D}_b^0}\leq K_b|z(s)|+C.$$

$$|\mathscr{B}z(t)| \leq \begin{cases} Me^{\omega t_1} \psi(q_1^*) \int_0^t e^{-\omega s} p(s) ds, & \text{if } t \in [0, t_1], \\ \sum_{i=1}^k Me^{\omega(t-t_k)} [Me^{\omega(t_{i-1}-t_i)} Me^{\omega(t_{i+2}-t_{i+1})} \dots Me^{\omega(t_{k-1}-t_{k-2})} \\ \times Me^{\omega(t_k-t_{k-1})}] Me^{\omega t_i} \times \psi(q^*) \int_{t_{i-1}}^{t_i} p(s) e^{-\omega s} ds \\ + Me^{\omega t} \psi(q^*) \int_{t_k}^t p(s) e^{-\omega s} ds, & \text{if } t \in (t_k, t_{k+1}]. \end{cases}$$

Then

$$|\mathscr{B}_{z}(t)| \leq \begin{cases} Me^{\omega t_{1}}\psi(q^{*}))\int_{0}^{t}e^{-\omega s}p(s)ds, & \text{if } t \in [0,t_{1}], \\ \sum_{i=1}^{k}M^{k-i+2}e^{\omega(t-t_{k})+t_{i-1}-t_{i}+t_{i+2}-t_{i+1}\dots+t_{k-1}-t_{k-2}+t_{k}-t_{k-1}+t_{i}}) \\ \times\psi(q^{*})\int_{t_{i-1}}^{t}p(s)e^{-\omega s}ds + Me^{\omega t}\psi(q^{*})\int_{t_{k}}^{t}p(s)e^{-\omega s}ds, & \text{if } t \in (t_{k},t_{k+1}]. \end{cases}$$

Using characteristic of the exponential function, we get

$$|\mathscr{B}z(t)| \leq \begin{cases} Me^{\omega t_1} \psi(q_1^*) \int_0^t e^{-\omega s} p(s) ds, & \text{if } t \in [0, t_1], \\ \sum_{i=1}^k M^{k-i+2} e^{\omega(t-t_{k-1})} \times \psi(q_2^*) \int_{t_{i-1}}^{t_i} p(s) e^{-\omega s} ds \\ + Me^{\omega t} \psi(q_2^*) \int_{t_k}^t p(s) e^{-\omega s} ds, & \text{if } t \in (t_k, t_{k+1}]. \end{cases}$$

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Finally, we obtain

$$|\mathscr{B}z(t)| \leq \begin{cases} Me^{\omega t_1} \psi(q_1^*) \int_0^t e^{-\omega s} p(s) ds = l_1, & \text{if } t \in [0, t_1], \\ \sum_{i=1}^k M^{k-i+2} e^{\omega (t_{k+1}-t_{k-1})} \times \psi(q_2^*) \int_{t_{i-1}}^{t_i} p(s) e^{-\omega s} ds \\ + Me^{\omega t_{k+1}} \psi(q_2^*) \int_{t_k}^t p(s) e^{-\omega s} ds = l_k, \quad k = 2, 3, ..., m, \quad \text{if } t \in (t_k, t_{k+1}]. \end{cases}$$

$$|\mathscr{B}z(t)| \leq \begin{cases} Me^{\omega t_1} \psi(q^*) \int_0^t e^{-\omega s} p(s) ds = l_1, & \text{if } t \in [0, t_1], \\ \sum_{i=1}^k M^{k-i+2} e^{\omega (t_{k+1}-t_{k-1})} \times \psi(q^*) \int_{t_{i-1}}^{t_i} p(s) e^{-\omega s} ds \\ + Me^{\omega t_{k+1}} \psi(q^*) \int_{t_k}^t p(s) e^{-\omega s} ds = l_k, & k = 2, 3, ..., m, & \text{if } t \in (t_k, t_{k+1}]. \end{cases}$$

Step 3: \mathscr{B} maps bounded sets into equicontinuous sets of \mathscr{D}_h^0 .

Let $\tau_1, \tau_2 \in J \setminus \{t_1, t_2, ..., t_m\}$ with $\tau_1 < \tau_2$, let B_q be a bounded set in \mathscr{D}_b^0 , and let $z \in B_q$. • If $\tau_1, \tau_2 \in [0, t_1]$, we have

$$|\mathscr{B}z(\tau_2) - \mathscr{B}z(\tau_1)|$$

$$= |\int_0^{\tau_2} S_{\alpha}(\tau_2 - s) f(s, x_{\rho(s, x_s + \overline{z}_s)} + \overline{z}_{\rho(s, x_s + \overline{z}_s)}) ds - \int_0^{\tau_1} S_{\alpha}(\tau_1 - s) f(s, x_{\rho(s, x_s + \overline{z}_s)} + \overline{z}_{\rho(s, x_s + \overline{z}_s)}) ds|.$$

Using the linearity of integral operator and hypotheses H_4 , we get

$$\begin{split} |\mathscr{B}z(\tau_{2}) - \mathscr{B}z(\tau_{1})| \\ = |\int_{0}^{\tau_{1}} S_{\alpha}(\tau_{2} - s)f(s, x_{\rho(s, x_{s} + \bar{z}_{s})} + \bar{z}_{\rho(s, x_{s} + \bar{z}_{s})})ds + \int_{\tau_{1}}^{\tau_{2}} S_{\alpha}(\tau_{2} - s)f(s, x_{\rho(s, x_{s} + \bar{z}_{s})} + \bar{z}_{\rho(s, x_{s} + \bar{z}_{s})})ds \\ &- \int_{0}^{\tau_{1}} S_{\alpha}(\tau_{1} - s)f(s, x_{\rho(s, x_{s} + \bar{z}_{s})} + \bar{z}_{\rho(s, x_{s} + \bar{z}_{s})})ds| \\ = |\int_{0}^{\tau_{1}} (S_{\alpha}(\tau_{2} - s) - S_{\alpha}(\tau_{1} - s))f(s, \bar{z}_{t}(s) + x_{t}(s))ds + \int_{\tau_{1}}^{\tau_{2}} S_{\alpha}(\tau_{2} - s)f(s, x_{\rho(s, x_{s} + \bar{z}_{s})} + \bar{z}_{\rho(s, x_{s} + \bar{z}_{s})})ds| \\ \leq \int_{0}^{\tau_{1}} ||S_{\alpha}(\tau_{2} - s) - S_{\alpha}(\tau_{1} - s)|||f(s, x_{\rho(s, x_{s} + \bar{z}_{s})} + \bar{z}_{\rho(s, x_{s} + \bar{z}_{s})})|ds \\ &+ \int_{\tau_{1}}^{\tau_{2}} ||S_{\alpha}(\tau_{2} - s) - S_{\alpha}(\tau_{1} - s)|||f(s, x_{\rho(s, x_{s} + \bar{z}_{s})} + \bar{z}_{\rho(s, x_{s} + \bar{z}_{s})})|ds \\ \leq \Psi(q_{1}^{*}) \int_{0}^{\tau_{1}} ||S_{\alpha}(\tau_{2} - s) - S_{\alpha}(\tau_{1} - s)||p(s)ds + Me^{\omega\tau_{2}}\Psi(q^{*}) \int_{\tau_{1}}^{\tau_{2}} e^{-\omega s}p(s)ds. \end{split}$$

If $\tau_1 = 0$, the right-hand side of previous inequality tends to zero as $\tau_2 \longrightarrow 0$ uniformly for $z \in \mathscr{D}_b^0$. If $0 < \tau_1 < \tau_2$, for $\varepsilon > 0$ whit $\varepsilon < \tau_1 < \tau_2$, we have

$$|\mathscr{B}(z(\tau_2)) - \mathscr{B}(z(\tau_1))| \leq \int_0^{\tau_1 - \varepsilon} \|S_{\alpha}(\tau_2 - s) - S_{\alpha}(\tau_1 - s)\| |f(s, x_{\rho(s, x_s + \overline{z}_s)} + \overline{z}_{\rho(s, x_s + \overline{z}_s)})| ds$$

$$\begin{split} + \int_{\tau_{1}-\varepsilon}^{\tau_{1}} \|S_{\alpha}(\tau_{2}-s) - S_{\alpha}(\tau_{1}-s)\| \|f(s,x_{\rho(s,x_{s}+\bar{z}_{s})} + \bar{z}_{\rho(s,x_{s}+\bar{z}_{s})})\|ds \\ + \int_{\tau_{1}}^{\tau_{2}} \|S_{\alpha}(\tau_{2}-s)\| \|f(s,x_{\rho(s,x_{s}+\bar{z}_{s})} + \bar{z}_{\rho(s,x_{s}+\bar{z}_{s})})\|ds \\ &\leq \psi(q_{1}^{*}) \int_{0}^{\tau_{1}-\varepsilon} \|S_{\alpha}(\tau_{2}-s) - S_{\alpha}(\tau_{1}-s)\| p(s)ds \\ &+ \psi(q_{1}^{*}) \int_{\tau_{1}-\varepsilon}^{\tau_{1}} \|S_{\alpha}(\tau_{2}-s) - S_{\alpha}(\tau_{1}-s)\| p(s)ds \\ &+ Me^{\omega\tau_{2}} \psi(q^{*}) \int_{\tau_{1}}^{\tau_{2}} e^{-\omega s} p(s)ds. \end{split}$$

From lemma 1, the operator $S_{\alpha}(t)$ is a uniformly continuous operator for $t \in [\varepsilon, t_1]$. Combining this and the arbitrariness of ε with the above estimation on $|\mathscr{B}(z(\tau_2)) - \mathscr{B}(z(\tau_1))|$, we can conclude that

$$\lim_{[\tau_1,\tau_2]\longrightarrow 0} |\mathscr{B}(z(\tau_2)) - B(z(\tau_1))| = 0.$$

Thus the operator \mathscr{B} is equicontinous on $[0, t_1]$.

• If $\tau_1, \tau_2 \in (t_k, t_{k+1}]$,

Then

$$\begin{aligned} |\mathscr{B}z(\tau_{2}) - \mathscr{B}z(\tau_{1})| \\ \leq \sum_{i=1}^{k} \int_{t_{i-1}}^{t_{i}} \|S_{\alpha}(\tau_{2} - t_{k}) - S_{\alpha}(\tau_{1} - t_{k})\| \prod_{j=i}^{k-1} \|S_{\alpha}(t_{j+1} - t_{j})\| \\ \times \|S_{\alpha}(t_{i} - s)\| \|f(s, x_{\rho(s, x_{s} + \bar{z}_{s})} + \bar{z}_{\rho(s, x_{s} + \bar{z}_{s})})| ds + \|\int_{t_{k}}^{\tau_{1}} S_{\alpha}(\tau_{2} - s)f(s, x_{\rho(s, x_{s} + \bar{z}_{s})} + \bar{z}_{\rho(s, x_{s} + \bar{z}_{s})}) ds \end{aligned}$$

$$+\int_{\tau_1}^{\tau_2} S_{\alpha}(\tau_2-s)f(s,x_{\rho(s,x_s+\overline{z}_s)}+\overline{z}_{\rho(s,x_s+\overline{z}_s)})ds - \int_{t_k}^{\tau_1} S_{\alpha}(\tau_1-s)f(s,x_{\rho(s,x_s+\overline{z}_s)}+\overline{z}_{\rho(s,x_s+\overline{z}_s)})ds \|.$$

Which gives

$$|\mathscr{B}z(\tau_2) - \mathscr{B}z(\tau_1)| \le \sum_{i=1}^k \int_{t_{i-1}}^{t_i} \|S_{\alpha}(\tau_2 - t_k) - S_{\alpha}(\tau_1 - t_k)\| \prod_{j=i}^{k-1} \|S_{\alpha}(t_{j+1} - t_j)\|$$

$$\times \|S_{\alpha}(t_{i}-s)\| \|f(s,x_{\rho(s,x_{s}+\overline{z}_{s})}+\overline{z}_{\rho(s,x_{s}+\overline{z}_{s})})\|ds$$

$$+ \int_{t_{k}}^{\tau_{1}} \|S_{\alpha}(\tau_{2}-s)-S_{\alpha}(\tau_{1}-s)\| \|f(s,x_{\rho(s,x_{s}+\overline{z}_{s})}+\overline{z}_{\rho(s,x_{s}+\overline{z}_{s})})\|ds$$

$$+ \int_{\tau_{1}}^{\tau_{2}} \|S_{\alpha}(\tau_{2}-s)\| \|f(s,x_{\rho(s,x_{s}+\overline{z}_{s})}+\overline{z}_{\rho(s,x_{s}+\overline{z}_{s})})\|ds.$$

Under the hypothesis H_4 , and lemma, we obtain

$$\begin{aligned} |\mathscr{B}z(\tau_{2}) - \mathscr{B}z(\tau_{1})| &\leq \sum_{i=1}^{k} \psi(q_{2}^{*}) \int_{t_{i-1}}^{t_{i}} \|S_{\alpha}(\tau_{2} - t_{k}) - S_{\alpha}(\tau_{1} - t_{k})\| \\ \prod_{j=i}^{k-1} \|S_{\alpha}(t_{j+1} - t_{j})\| \times \|S_{\alpha}(t_{i} - s)\| \|p(s)ds \\ &+ \psi(q^{*}) \int_{t_{k}}^{\tau_{1} - \varepsilon} \|S_{\alpha}(\tau_{2} - s) - S_{\alpha}(\tau_{1} - s)\| p(s)ds \\ &+ \psi(q^{*}) \int_{\tau_{1} - \varepsilon}^{\tau_{1}} \|S_{\alpha}(\tau_{2} - s) - S_{\alpha}(\tau_{1} - s)\| p(s)ds \\ &+ M\psi(q^{*}) e^{\omega\tau_{2}} \int_{\tau_{1}}^{\tau_{2}} e^{-\omega s} p(s)ds. \end{aligned}$$

As $\tau_1 \longrightarrow \tau_2$ and ε becomes sufficiently small, the right-hand side of the above inequality tends to zero, since S_{α} is analytic operator and the compactness of $S_{\alpha}(t)$ for t > 0 implies the continuity in the uniform operator topology. This proves the equicontinuity for the case where $t \neq t_i, i = 1, ..., m + 1$.

Now, it remains to examine equicontinuity at $t = t_l$. We have for $z \in B_q$, for each $t \in J$.

First, we prove equicontinuity at $t = t_l^-$. Fix $\delta_1 > 0$ such that $\{t_k, k \neq l\} \cap [t_l - \delta_1, t_l - \delta_1] = \emptyset$. For $0 < h < \delta_1$, we have • if l = 1 i.e $t_1 - h, t_1 \in [0, t_1]$,

$$|\mathscr{B}(z(t_1-h)) - \mathscr{B}(z(t_1))| \le \psi(q_1^*) \int_0^{t_1-h} ||S_{\alpha}(t_1-s) - S_{\alpha}(t_1-h-s)|| p(s) ds + Me^{\omega t_1} \psi(q^*) \int_{t_1-h}^{t_1} e^{-\omega s} p(s) ds$$

Which tends to zero as $h \longrightarrow 0$ since $S_{\alpha}(t)$ is uniformly continuous operator for $t \in [0, t_1]$ thus the operator *B* is equicontinuous at $t = t_1^-$.

• if
$$t_l - h, t_l \in [t_k, t_{k+1}]$$
.
en:

$$\begin{aligned} |\mathscr{B}(z)(t_{l}-h) - \mathscr{B}(z)(t_{l})| &\leq \sum_{i=1}^{k} \psi(q^{*}) \int_{t_{i-1}}^{t_{i}} \|S_{\alpha}(t_{l}-t_{k}) - S_{\alpha}(t_{l}-h-t_{k})\| \\ &\times \prod_{j=i}^{k-1} \|S_{\alpha}(t_{j+1}-t_{j})\| \|S_{\alpha}(t_{i}-s)\| p(s) ds \\ &+ \psi(q^{*}) \int_{t_{k}}^{t_{l}-h} \|S_{\alpha}(t_{l}-s) - S_{\alpha}(t_{l}-h-s)\| p(s) ds \\ &+ M \psi(q^{*}) e^{\omega t_{l}} \int_{t_{l}-h}^{t_{l}} e^{-\omega s} p(s) ds. \end{aligned}$$

The right-hand side of the previous inequality tends to zero as $h \longrightarrow 0$. So the operator \mathscr{B} is equicontinuous at t_l^- . Now, define

$$\mathscr{B}_0(z)(t) = \mathscr{B}(z)(t), if \ t \in [0, t_1],$$

and

$$\widehat{\mathscr{B}}_i(z)(t) = \begin{cases} \mathscr{B}(z)(t), & \text{if } t \in (t_i, t_{i+1}), \\ \mathscr{B}(z)(t_i^+), & \text{if } t = t_i. \end{cases}$$

Next, we prove equicontinuity at $t = t_i^+$. Fix $\delta_2 > 0$ such that $\{t_k, k \neq i\} \cap [t_i - \delta_2, t_i + \delta_2] = \emptyset$. First, we study the equicontinuity at $t = 0^+$. If $t \in [0, t_1]$, we have

$$\widehat{\mathscr{B}}_0(z)(t) = \begin{cases} \mathscr{B}z(t), & \text{if } t \in [0, t_1], \\ 0, & \text{if } t = 0. \end{cases}$$

For $0 < h < \delta_2$, we have

$$\begin{split} |\widehat{\mathscr{B}}_{0}(z)(h) - \widehat{\mathscr{B}}_{0}(z)(0)| &= |\mathscr{B}(z)(h)| \\ &= \|\int_{0}^{h} S_{\alpha}(h-s)f(s, x_{\rho(s,x_{s}+\bar{z}_{s})} + \bar{z}_{\rho(s,x_{s}+\bar{z}_{s})})ds\| \\ &\leq \int_{0}^{h} \|S_{\alpha}(h-s)\||f(s, x_{\rho(s,x_{s}+\bar{z}_{s})} + \bar{z}_{\rho(s,x_{s}+\bar{z}_{s})})|ds| \\ &\leq \psi(q^{*})e^{\omega h}\int_{0}^{h} e^{-\omega s}p(s)ds. \end{split}$$

The right-hand side tends to zero as $h \longrightarrow 0$.

Now, we study the equicontinuity at $t_1^+, t_2^+, ..., t_m^+(t_l^+, 1 \le l \le m)$. For $0 < h < \delta_2$, we have

$$\begin{aligned} |\mathscr{B}(z)(t_l+h) - \mathscr{B}(z)(t_l)| &\leq \sum_{i=1}^k \psi(q^*) \int_{t_i-1}^{t_i} \|S_{\alpha}(h) - S_{\alpha}(0)\| \\ &\times \prod_{j=i}^{k-1} \|S_{\alpha}(t_{j+1}-t_j)\| \|S_{\alpha}(t_i-s)\| p(s) ds \\ &+ M \psi(q^*) e^{\omega(t_l+h)} \int_{t_i}^{t_l+h} e^{-\omega s} p(s) ds. \end{aligned}$$

It is clear that the right-hand side tends to zero as $h \longrightarrow 0$.

Then \mathscr{B} is equicontinuous at t_l^+ , $(1 \le l \le m)$. The equicontinuity for the cases $\tau_1 < \tau_2 \le 0$ and $\tau_1 \le 0 \le \tau_2$ follows from the uniform continuity of ϕ on the interval $] -\infty, 0]$. As a consequence of steps 1 and 3 toghether with Arzel-Ascoli Theorem it suffices to show that \mathscr{B}_z maps B_q into a precompact set in *E* i.e.: we show that the set $\{\mathscr{B}_z(t), z \in B_q\}$ is precompact in *E* for every $t \in [0, b]$.

Now, let $x \in B_q$ and let ε be a positive real number satisfying $0 < \varepsilon < t \le b$. For $z \in B_q$ and $t \in [0, t_1]$. we have if t = 0 the set $\{\mathscr{B}z(0); z \in B_q\} = \{0\}$ which is precompact as a finite set.

For $0 < \varepsilon < t \le t_1$, we have

$$\begin{aligned} \mathscr{B}(z)(t) &= \int_0^t S_\alpha(t-s) f(s, x_{\rho(s, x_s + \overline{z}_s)} + \overline{z}_{\rho(s, x_s + \overline{z}_s)}) ds \\ &= \int_0^{t-\varepsilon} S_\alpha(t-s) f(s, x_{\rho(s, x_s + \overline{z}_s)} + \overline{z}_{\rho(s, x_s + \overline{z}_s)}) ds + \int_{t-\varepsilon}^t S_\alpha(t-s) f(s, x_{\rho(s, x_s + \overline{z}_s)} + \overline{z}_{\rho(s, x_s + \overline{z}_s)}) ds. \end{aligned}$$

Set $F_0 := \{S_{\alpha}(t-\theta)f(\theta, x_{\rho(\theta, x_{\theta}+\overline{z}_{\theta})} + \overline{z}_{\rho(\theta, x_{\theta}+\overline{z}_{\theta})}); \theta \in [0, t-\varepsilon], z \in B_q\}$, from the mean value Theorem for the Bochner integral, we have

$$\int_{0}^{t-\varepsilon} S_{\alpha}(t-s) f(s, \overline{z}_{t}(s) + x_{t}(s)) ds \in (t-\varepsilon) Conv(\overline{F_{0}}).$$
(6)

On the other hand, using hypotheses (H_1) and (H_4) , we obtain

$$\int_{t-\varepsilon}^{t} |S_{\alpha}(t-s)f(s,\overline{z}_{t}(s)+x_{t}(s))| ds \leq Me^{\omega t} \psi(q^{*}) \int_{t-\varepsilon}^{t} e^{-\omega s} p(s) ds.$$

Let C^0_{ε} the circle who's diameter d^0_{ε} is such that

$$d_{\varepsilon}^{0} \leq M e^{\omega t} \psi(q_{1}^{*}) \int_{t-\varepsilon}^{t} e^{-\omega s} p(s) ds.$$
⁽⁷⁾

As a consequence of (6) and (7), we conclude that

$$Bz(t) \in (t - \varepsilon)Conv(\overline{F_0}) + C_{\varepsilon}^0, \quad \forall 0 < \varepsilon < t \le t_1.$$
(8)

For $t_k < \varepsilon < t < t_{k+1}$ and $z \in B$, we have

$$Bz(t) = \sum_{i=1}^{k} \int_{t_{i-1}}^{t_i} S_{\alpha}(t-t_k) \prod_{j=i}^{k-1} S_{\alpha}(t_{j+1}-t_j) S_{\alpha}(t_i-s) f(s,\overline{z}_t(s)+x_t(s)) ds.$$
(9)

$$+\int_{t_k}^{t-\varepsilon}S_{\alpha}(t-s)f(s,\overline{z}_t(s)+x_t(s))ds+\int_{t-\varepsilon}^{t}S_{\alpha}(t-s)f(s,\overline{z}_t(s)+x_t(s))ds$$

Set $F_k := \{S_{\alpha}(t-\theta)f(\theta, \overline{z}_t(\theta) + x_t(\theta)); \theta \in (t_k, t_{k+1}), z \in B_q\}$, from the mean value theorem for the Bochner integral, we have

$$\int_{t_k}^{t-\varepsilon} S_{\alpha}(t-s) f(s, x_{\rho(s, x_s+\overline{z}_s)} + \overline{z}_{\rho(s, x_s+\overline{z}_s)}) ds \in (t-t_k-\varepsilon) Conv(\overline{F_k}).$$
(10)

From $(H_1), (H_4)$ we obtain

$$\begin{split} \sum_{i=1}^{k} \int_{t_{i-1}}^{t_i} S_{\alpha}(t-t_k) \prod_{j=i}^{k-1} S_{\alpha}(t_{j+1}-t_j) S_{\alpha}(t_i-s) f(s, x_{\rho(s,x_s+\overline{z}_s)} + \overline{z}_{\rho(s,x_s+\overline{z}_s)}) ds \\ &+ \int_{t_k}^{t-\varepsilon} S_{\alpha}(t-s) f(s, \overline{z}_t(s) + x_t(s)) ds \\ \leq \Psi(q^*) \sum_{i=1}^{k} M^{k-i+2} e^{\omega(t-t_{k-1}} \int_{t_{i-1}}^{t_i} e^{-\omega s} p(s) ds + M \Psi(q^*) e^{\omega t} \int_{t-\varepsilon}^{t} e^{-\omega s} p(s) ds. \end{split}$$

Let C_{ε}^{k} the circle who's diameter d_{ε}^{k} is such that

$$d_{\varepsilon}^{k} \leq \psi(q^{*}) \sum_{i=1}^{k} M^{k-i+2} e^{\omega(t-t_{k-1})} \int_{t_{i-1}}^{t_{i}} e^{-\omega s} p(s) ds + M \psi(q^{*}) e^{\omega t} \int_{t-\varepsilon}^{t} e^{-\omega s} p(s) ds.$$
(11)

From (9) and (11), it follows that

$$\mathscr{B}_{z}(t) \in (t - t_{k} - \varepsilon)Conv(\overline{F_{k}}) + C_{\varepsilon}^{k}, \ \forall t_{k} < \varepsilon < t < t_{k+1}.$$

$$(12)$$

From 8 and 12, we conclude that $\Re z(t)$ is precompact in *E*. From Step1– Step 3, we deduce that \Re is completely continuous.

Step 4: *A* is a contraction.

For $t \in]-\infty, t_1]$, we have

$$|\mathscr{A}z_1(t) - \mathscr{A}z_2(t)| = 0.$$

which implies that \mathscr{A} is contraction for all $t \in]-\infty, t_1]$. It remains to prove that A is a contraction operator for $t \in [t_k, t_{k+1}], k \ge 1$

$$\begin{split} |\mathscr{A}z_{1}(t) - \mathscr{A}z_{2}(t)| &= |\sum_{i=1}^{k} S_{\alpha}(t-t_{k}) \prod_{j=i}^{k-1} S_{\alpha}(t_{j+1}-t_{j}) \\ \times [(I_{i}(x_{t_{i}} + \overline{z}_{t_{i}}^{1}) - I_{i}(x_{t_{i}} + \overline{z}_{t_{i}}^{2})] \\ &\leq \sum_{i=1}^{k} \|S_{\alpha}(t-t_{k})\| \prod_{j=i}^{k-1} \|S_{\alpha}(t_{j+1}-t_{j})\| |[I_{i}(\overline{z}_{t_{i}}^{1}) - I_{i}(\overline{z}_{t_{i}}^{2})| \\ &\leq \sum_{i=1}^{k} M e^{\omega(t-t_{k})} \prod_{j=i}^{k-1} M e^{\omega(t_{j+1}-t_{j})}\| |[I_{i}(\overline{z}_{t_{i}}^{1}) - I_{i}(\overline{z}_{t_{i}}^{2})| \\ &\leq \sum_{i=1}^{k} M e^{\omega(t-t_{k})} [M e^{\omega(t_{i+1}-t_{i})} M e^{\omega(t_{i+2}-t_{i+1})} ... M e^{\omega(t_{k-1}-t_{k-2})} \\ &\times M e^{\omega(t_{k}-t_{k-1})}] \times \| |[I_{i}(\overline{z}_{t_{i}}^{1}) - I_{i}(\overline{z}_{t_{i}}^{2})| \\ &\leq \sum_{i=1}^{k} M^{k-i+1} e^{\omega(t-t_{i})} \| |[I_{i}(\overline{z}_{t_{i}}^{1}) - I_{i}(\overline{z}_{t_{i}}^{2})|. \end{split}$$

Since $t \in J := [0, b]$ and the functions I_k ; k = 1, 2, ..., m. Lipschitz; Then

$$|\mathscr{A}z_{1}(t) - \mathscr{A}z_{2}(t)| \leq \sum_{i=1}^{k} M^{k-i+1} e^{\omega(b-t_{i})} M_{i} ||z_{t_{i}}^{1} - z_{t_{i}}^{2} ||_{\mathscr{D}}$$

It fallows that

$$\|\mathscr{A} z_1 - \mathscr{A} z_2\| \leq K_b \sum_{i=1}^k M^{k-i+1} e^{\omega(b-t_i)} M_i \| z^1 - z^2 \|_{\mathscr{D}}.$$

Thus the operator \mathscr{A} is a contraction, since

$$K_b \sum_{i=1}^k M^{k-i+1} e^{\omega(b-t_i)} M_i < 1.$$

Step 5: A priori bounds.

Now it remains to show that the set

$$\Upsilon = \{z \in PC(] - \infty, b], E) : z = \lambda \mathscr{B}(z) + \lambda \mathscr{A}(\frac{z}{\lambda}), \quad \text{for some} \quad 0 < \lambda < 1\}.$$

is bounded.

Let $z \in \Upsilon$ be any element, then $z = \lambda \mathscr{B}(z) + \lambda \mathscr{A}(\frac{z}{\lambda})$, for some $0 < \lambda < 1$. First, for each $t \in [0, t_1]$,

$$\begin{aligned} |z(t)| &= |\lambda \int_0^t S_\alpha(t-s) f(s, x_{\rho(s,x_s+z_s)} + \bar{z}_{\rho(s,x_s+z_s)}) ds| \\ &\leq M e^{\omega t} \int_0^t e^{-\omega s} |f(s, x_{\rho(s,x_s+z_s)} + \bar{z}_{\rho(s,x_s+z_s)})| ds \\ &\leq M e^{\omega t_1} \int_0^{t_1} e^{-\omega s} p(s) \psi(||x_{\rho(s,x_s+z_s)} + \bar{z}_{\rho(s,x_s+z_s)}||) ds \\ &\leq M e^{\omega t_1} \int_0^{t_1} e^{-\omega s} p(s) \psi(K_b|z(s)| + (M_b + MK_b e^{\omega t_1} + L^{\phi}) ||\phi||) ds. \end{aligned}$$

On the other hand, for each $t \in (t_k, t_{k+1}]$, we have

$$\begin{aligned} |z(t)| &= \|\lambda \left(\sum_{i=1}^{k} \int_{t_{i-1}}^{t_i} S_\alpha(t-t_k) \prod_{j=i}^{k-1} S_\alpha(t_{j+1}-t_j) S_\alpha(t_i-s) f(s, x_{\rho(s, x_s+z_s)} + \bar{z}_{\rho(s, x_s+z_s)}) ds \\ &+ \int_{t_k}^{t} S_\alpha(t-s) f(s, x_{\rho(s, x_s+z_s)} + \bar{z}_{\rho(s, x_s+z_s)}) + \lambda \left(S_\alpha(t-t_k) \prod_{i=1}^{k} S_\alpha(t_i-t_{i-1}) \phi(0) \right) \\ &+ \sum_{i=1}^{k} S_\alpha(t-t_k) \prod_{j=i}^{k-1} S_\alpha(t_{j+1}-t_j) I_i(\frac{z_{i_j}}{\lambda} + x_{t_i}) \end{aligned}$$

From (H_1) , (H_2) and since $\lambda < 1$, we obtain

$$\begin{aligned} |z(t)| &\leq M^{k+1} e^{\omega t} |\phi(0)| + \sum_{i=1}^{k} M^{k-i+1} e^{\omega(t-t_i)} |I_i(0)| \\ &+ \sum_{i=1}^{k} M^{k-i+2} e^{\omega(t-t_{k-1})} \int_{t_{i-1}}^{t_i} e^{-\omega s} p(s) \psi(K_b | z(s)| + C) ds \\ &+ M e^{\omega t} \int_{t_k}^{t} e^{-\omega s} p(s) \psi(K_b | z(s)| + C) ds \\ &+ \sum_{i=1}^{k} M^{k-i+1} e^{\omega(t-t_i)} |I_i(\frac{z_{t_i}}{\lambda} + x_{t_i}) - I_i(0)|. \end{aligned}$$

Since I_i are Lipschitz, then

$$\begin{split} |z(t)| &\leq M^{k+1}e^{\omega t}|\phi(0)| + \sum_{i=1}^{k} M^{k-i+1}e^{\omega (t-t_i)}|I_i(0)| \\ &+ \sum_{i=1}^{k} M^{k-i+2}e^{\omega (t-t_{k-1})} \int_{t_{i-1}}^{t_i} e^{-\omega s} p(s) \psi(K_b|z(s)|+C) ds \\ &+ Me^{\omega t} \int_{t_k}^{t} e^{-\omega s} p(s) \psi(K_b|z(s)|+C) ds \\ &+ \sum_{i=1}^{k} M^{k-i+1}e^{\omega (t-t_i)} M_i ||z(t_i^-)|| \\ &\leq M^{k+1}e^{\omega t} |\phi(0)| + \sum_{i=1}^{k} M^{k-i+1}e^{\omega (t-t_i)}|I_i(0)| \\ &+ \sum_{i=1}^{k} M^{k-i+2}e^{\omega (t-t_{k-1})} \int_{t_{i-1}}^{t_i} e^{-\omega s} p(s) \psi(K_b|z(s)|+C) ds \\ &+ Me^{\omega t} \int_{t_k}^{t} e^{-\omega s} p(s) \psi(K_b|z(s)|+C) ds \\ &+ \sum_{i=1}^{k} M^{k-i+1}e^{\omega (t-t_i)} M_i(K_b|z(t)|+C). \end{split}$$

Therefore

$$\begin{split} [1 - K_b \sum_{i=1}^k M^{k-i+1} e^{\omega(b-t_i)} M_i] |z(t)| &\leq M^{k+1} e^{\omega t} \|\phi(0)\| + \sum_{i=1}^k M^{k-i+1} e^{\omega(b-t_i)} |I_i(0)| \\ &+ \sum_{i=1}^k M^{k-i+2} e^{\omega(t-t_{k-1})} \int_{t_{i-1}}^{t_i} e^{-\omega s} p(s) \psi(K_b |z(s)| + C) ds \\ &+ M e^{\omega b} \int_{t_k}^t e^{-\omega s} p(s) \psi(K_b |z(s)| + C) ds + C \sum_{i=1}^k M^{k-i+1} e^{\omega(b-t_i)} M_i. \end{split}$$

Thus we have

$$\begin{aligned} |z(t)| &\leq \left(M^{k+1} e^{\omega b} |\phi(0)| + \sum_{i=1}^{k} M^{k-i+1} e^{\omega(b-t_i)} (|I_i(0)| + C) \right. \\ &+ \sum_{i=1}^{k} M^{k-i+2} e^{\omega(b-t_{k-1})} \int_{t_{i-1}}^{t_i} e^{-\omega s} p(s) \psi(K_b | z(s)| + C) ds \right) \Big/ \left(1 - K_b \sum_{i=1}^{k} M^{k-i+1} e^{\omega(b-t_i)} M_i \right) \\ &+ \frac{M e^{\omega b}}{(1 - K_b \sum_{i=1}^{k} M^{k-i+1} e^{\omega(b-t_i)} M_i)} \int_{t_k}^{t} e^{-\omega s} p(s) \psi(K_b | z(s)| + C) ds \right). \end{aligned}$$

We consider the function $\mu(t)$ defined by

$$\mu(t) = \sup\{K_b|z(s)| + C; 0 \le s \le t\}, \quad 0 \le t \le b.$$

Let $t^* \in [0, t]$ be such that $\mu(t) = K_b |z(t^*)| + C$. If $t^* \in J$, by the previous inequality, we have for $t \in J$.

• if $t \in [0, t_1]$,

$$\mu(t) \leq K_b M e^{\omega b} \int_0^t e^{-\omega s} p(s) \psi(\mu(s)) ds + C.$$

• if
$$t \in (t_k, t_{k+1}]$$
,
 $|\mu(t)| \leq K_b \Big(M^{k+1} e^{\omega b} |\phi(0)| + \sum_{i=1}^k M^{k-i+1} e^{\omega(b-t_i)} (|I_i(0)| + C) + \sum_{i=1}^k M^{k-i+2} e^{\omega(b-t_{k-1})} \int_{t_{i-1}}^{t_i} e^{-\omega s} p(s) \psi(|\mu(s)|) ds \Big) \Big/ \Big(1 - K_b \sum_{i=1}^k M^{k-i+1} e^{\omega(b-t_i)} M_i \Big) + \frac{K_b M e^{\omega b}}{(1 - K_b \sum_{i=1}^k M^{k-i+1} e^{\omega(b-t_i)} M_i)} \int_{t_k}^t e^{-\omega s} p(s) \psi(|\mu(s)|) ds \Big) + C.$

Then

$$\mu(t) \leq C_1 + C_2 \int_{t_k}^t e^{-\omega s} p(s) \psi(\mu(s)) ds.$$

Where

$$C_{1} = \frac{K_{b} \Big(M^{k+1} e^{\omega b} |\phi(0)| + \sum_{i=1}^{k} M^{k-i+1} e^{\omega(b-t_{i})} (|I_{i}(0)| + C) \Big)}{1 - K_{b} \sum_{i=1}^{k} M^{k-i+1} e^{\omega(b-t_{i})} M_{i}} \\ + \frac{K_{b} \sum_{i=1}^{k} M^{k-i+2} e^{\omega(b-t_{k-1})} \int_{t_{i-1}}^{t_{i}} e^{-\omega s} p(s) \psi(\mu(s)) ds}{1 - K_{b} \sum_{i=1}^{k} M^{k-i+1} e^{\omega(b-t_{i})} M_{i}} + C. \\ C_{2} = \frac{K_{b} M e^{\omega b}}{(1 - K_{b} \sum_{i=1}^{k} M^{k-i+1} e^{\omega(b-t_{i})} M_{i})}.$$

It follows that

$$\mu(t) \leq \begin{cases} C + K_b M e^{\omega b} \int_0^t e^{-\omega s} p(s) \psi(\mu(s)) ds, & \text{if } t \in [0, t_1] \\ C_1 + C_2 \int_{t_k}^t e^{-\omega s} p(s) \psi(\mu(s)) ds, & \text{if } t \in (t_k, t_{k+1}). \end{cases}$$

Let us take the right-hand side of the above inequality as $\vartheta(t)$,

$$\mu(t) \leq \vartheta(t).$$

and

$$\vartheta(t) := \begin{cases} C + K_b M e^{\omega b} \int_0^t e^{-\omega s} p(s) \psi(\mu(s)) ds, & \text{if } t \in [0, t_1] \\ C_1 + C_2 \int_{t_k}^t e^{-\omega s} p(s) \psi(\mu(s)) ds, & \text{if } t \in (t_k, t_{k+1}). \end{cases}$$

$$\begin{cases} \vartheta(0) = C, \\ \vartheta(t_k) = C_1, \quad k = 1, 2, ..., m. \end{cases}$$

And differentiating both sides of the above equality, we obtain

$$\vartheta'(t) := \begin{cases} K_b M e^{\omega(b-t)} p(t) \psi(\mu(t)), & \text{if } t \in [0,t_1], \\ C_2 e^{-\omega t} p(t) \psi(\mu(t)), & \text{if } t \in (t_k, t_{k+1}]. \end{cases}$$

Using the non decreasing character of function ψ , i.e

$$\boldsymbol{\mu}(t) \leq \boldsymbol{\vartheta}(t) \Rightarrow \boldsymbol{\psi}(\boldsymbol{\mu}(t)) \leq \boldsymbol{\psi}(\boldsymbol{\vartheta}(t))$$

We have

$$\vartheta'(t) \leq \begin{cases} K_b M e^{\omega(b-t)} p(t) \psi(\vartheta(t)), & \text{if } t \in [0,t_1], \\ C_2 e^{-\omega t} p(t) \psi(\vartheta(t)), & \text{if } t \in (t_k,t_{k+1}]. \end{cases}$$

It gives

$$\frac{\vartheta'(t)}{\psi(\vartheta(t))} \leq \begin{cases} K_b M e^{\omega(b-t)} p(t), & \text{if } t \in [0,t_1], \\ C_2 e^{-\omega t} p(t), & \text{if } t \in (t_k, t_{k+1}). \end{cases}$$

• Integrating from 0 to t, if $t \in [0, t_1]$, we get

$$\int_0^t \frac{\vartheta'(s)}{\psi(\vartheta(s))} ds \leq K_b M e^{\omega b} \int_0^t e^{-\omega s} p(s) ds.$$

By change of variable $(\vartheta(s) = u)(s : 0 \longrightarrow t, u : C \longrightarrow \vartheta(t))$:

$$\int_C^{\nu(t)} \frac{du}{\psi(u)} \le M e^{\omega b} \int_C^t e^{-\omega s} p(s) ds \le \int_0^\infty \frac{du}{\psi(u)}$$

Hence, there exists a constant η_1 such that

$$\mu(t) \leq \vartheta(t) \leq \eta_1.$$

• Now, integrating from t_k to t if $t \in (t_k, t_{k+1}]$, we get

$$\int_{t_k}^t \frac{\vartheta'(s)}{\psi(\vartheta(s))} ds \leq C_2 \int_{t_k}^t e^{-\omega s} p(s) ds.$$

By change of variable $(\vartheta(s) = u)(s : t_k \longrightarrow t, u : C_1 \longrightarrow \vartheta(t)):$

$$\int_0^{v(t)} \frac{du}{\psi(u)} \leq C_2 \int_{t_k}^t e^{-\omega s} p(s) ds \leq \int_{C_3}^{v(t)} \frac{du}{\psi(u)}.$$

Where $C_3 = min(C, C_2)$. Henc, there exists a constant η_2 such that

$$\mu(t) \leq \vartheta(t) \leq \eta_2, \quad t \in (t_k, t_{k+1})$$

In conclusion, there exists $\eta = \min(\eta_1, \eta_2)$ such that

$$\mu(t) \leq \vartheta(t) \leq \eta$$
, for all $t \in J$.

Now from the definition of μ it follows that, there exist $\eta^* > 0$ such that

$$||z||_{\mathscr{D}^0_h} \le \eta^*, \qquad \forall z \in \Upsilon.$$

This shows that the set Υ is bounded.

As a consequence of theorem, we deduce that $\mathscr{A} + \mathscr{B}$ has a fixed point z^* . Then $y^*(t) = z^*(t) + x(t), t \in (-\infty, b]$ is a fixed point of the operator N and hence the problem have a mild solution on interval $(-\infty, b]$. This completes the proof.

4 Application

We consider the following impulsive fractional differential equation of the form:

$$\frac{\partial_t^q}{\partial t^q} v(t,x) = \frac{\partial^2}{\partial x^2} v(t,x) + \int_{-\infty}^t a_1(s-t)v(s-\rho_1(t)\rho_2(|v(t)|,\xi)ds, \ x \in [0,\pi], \ t \in [0,b] \setminus \{t_1,...,t_m\},$$
(13)

$$\Delta v(t_i)(x) = \int_{-\infty}^t d_i(t_i - s)v(s, x)ds, \ x \in [0, \pi], i = 1, ..., m,$$
(14)

$$v(t,0) = v(t,\pi) = 0, t \in [0,b],$$
(15)

$$v(t,x) = v_0(\theta, x), \ \theta \in] -\infty, 0], x \in [0,\pi].$$
(16)

where 0 < q < 1, $d_i : R \to R, i = 1, 2, ..., m$, and $a_1 : (-\infty, 0] \to R$, $\rho_i : [0, +\infty) \to [0, +\infty), i = 1, 2$ are continuous functions.

Set $E = L^2([0,\pi])$ and let $D(A) \subset E \to E$ be the operator Au = u'' with the domaine

$$D(A) = \{ u \in H^1_0(0,\pi) \cap H^2(0,\pi) \}.$$

The operator A is the infinitesimal generator of analytic semi-group S(t).

Set $\gamma > 0$. For the phase space, we choose \mathcal{D} to defined by:

$$\mathscr{D} = PC^{\gamma} = \{ \Phi \in PC((-\infty, 0], E) : \lim_{\theta \in (-\infty, 0]} \exp(\gamma^{\theta}) \Phi(\theta) \quad \text{ exists in } E \}.$$

with norm

$$\|\phi\|_{\gamma} = \sup_{\theta \in (-\infty,0]} \exp(\gamma^{\theta}) |\phi(\theta)|, \qquad \phi \in PC^{\gamma}.$$

For this space, axioms (A1), (A2) are satisfied. the problem (4.1)–(4.4) takes the abstract form (1.1)–(1.3) by making the following change of variables.

$$y(t)(x) = v(t,x), \quad x \in [0,\pi], t \in J = [0,1].$$

$$\phi(\theta)(x) = v_0(\theta,x), \quad x \in [0,\pi], \theta \le 0.$$

$$f(t,\varphi)(x) = \int_{-\infty}^t a_1\varphi(s,x)ds.$$

$$\rho(t,\varphi) = s - \rho_1(t)\rho_2(|\varphi(0)|).$$

$$I_i(\varphi)(x) = \int_{-\infty}^0 d_i(-\theta)\varphi(\theta)(x)d\theta$$

Theorem 4. Let $\varphi \in \mathcal{B}$ such that H_{φ} holds, the problem (4.1)–(4.4) has at least one mild solution.

5 Conclusion

In this work, we provided the existence of mild solutions and with sufficient conditions for some differential fractional equations. The main tool of this paper is the fixed point theory combined with resolvent famillies. To our knowlege, there are few works using this technique. The obtained results have a contribution to the related literature and extend the results in [HL20, HGBA13].

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A Novel Method for Solving Nonlinear Jerk Equations

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Abstract. In this article, reproducing kernel method for solving Jerk equations is given. Convergence of the solution is shown. This method is applied to the equation for chosen values of the parameters that seem in the model and some numerical experiments prove that the reproducing kernel method is very effective method.

Keywords: Jerk equations \cdot Reproducing kernel functions \cdot Bounded linear operator

2000 Mathematics Subject Classification: 47B32 · 46E22 · 30E25

1 Introduction

Although most of the efforts on dynamical systems are related to second-order differential equations, the behavior of some dynamical systems is governed by nonlinear jerk (third-order) differential equations. Jerk is the rate of acceleration change in physics; that is, the time derivative of acceleration, and as such the second velocity derivative, or the third time position derivative. The jerk is significant in some mechanics and acoustics implementations. Many geometric features of the Jerk vector are founded for plane motion utilizing the aberrancy features of curves [1]. Nonlinear third-order differential equations, known as nonlinear Jerk equations, including the third temporal displacement derivative, are of great interest in investigating the structures which exhibit rotating and translating movements, such as robots or machine tools, where excessive Jerk leads to accelerated wear of transmissions and bearing elements, noisy operations and large contouring errors in discontinuities (such as corners) in the machining path [2]. The jerk equations are the minimal setting for solutions showing chaotic behaviour. The numerical solutions of the Jerk equation have been worked by many investigators [3].

Hu et al. [4] have investigated the iteration calculations of periodic solutions to nonlinear Jerk equations. Liu et al. [5] have obtained the periods and periodic solutions of nonlinear Jerk equations by an iterative algorithm based on a shape function method. Rahman et al. [6] have worked on modified harmonic balance method for the solution of nonlinear Jerk equations.

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We investigate the Jerk equation by reproducing kernel method in this paper. Reproducing kernel method (RKM) is very accurate and reliable method. There are many papers related to the reproducing kernel method in the literature. We apply this method to a new problem in this work. Akgül [7-9] has worked on reproducing kernel Hilbert space method based on reproducing kernel functions for investigating boundary layer flow of a Powell-Eyring non-Newtonian fluid, new reproducing kernel functions and the solutions of variable-order fractional differential equations by reproducing kernel method. Akgül et al. [10] have investigated the numerical solutions of fractional differential equations of Lane-Emden type by an accurate technique. Aronszajn [11] has studied the theory of reproducing kernels. Arqub [12-14] has investigated the approximate solutions of DASs with nonclassical boundary conditions using novel reproducing kernel algorithm, the reproducing kernel algorithm for handling differential algebraic systems of ordinary differential equations and the fitted reproducing kernel Hilbert space method for the solutions of some certain classes of time-fractional partial differential equations subject to initial and Neumann boundary conditions. Azarnavid et al. [15] have worked on an iterative reproducing kernel method in Hilbert space for the multi-point boundary value problems.

In this work, we studied a Jerk equation as:

(1.1)
$$y''' = J(y, y', y''),$$

with initial conditions

(1.2)
$$y(0) = 0, y(0) = B, y(0) = 0.$$

We define the most general function of Jerk as;

(1.3)
$$y^{\prime\prime\prime} + \alpha y y^{\prime} y^{\prime\prime} + \beta y^{\prime} y^{\prime\prime2} + \delta y^2 y^{\prime} + \varepsilon y^3 + \gamma y^{\prime} = 0,$$

where the parameters α , β , δ , ε and γ are constants.

This paper is organized as follows. In Sect. 2, the RKM is discussed with some preliminary concept and definition. The solution procedure and approximate solutions of Eqs. (1.2)-(1.3) are presented in this section. Numerical experiments are demonstrated in Sect. 3. Conclusion is given in the last section.

2 Reproducing Kernel Method

First of all we will construct the reproducing kernel Hilbert spaces that we need to solve our problem.

Definition 2.1. The first reproducing kernel Hilbert space that we will use is $M_2^3[0,1]$

$$M_2^3[0,1] = \{a \in AC[0,1]: a', a'' \in AC[0,1], a^{(3)} \in L^2[0,1]\}.$$

We have the inner product and the norm for this space as:

$$\langle a,b\rangle_{M_2^3} = a(0)b(0) + a'(0)b'(0) + a''(0)b''(0) + \int_0^1 a'''(x)b'''(x)dx, \quad a,b \in M_2^3[0,1]$$

and

$$\|a\|_{M_2^3} = \sqrt{\langle a, a \rangle_{M_2^3}}, \quad a \in M_2^3[0, 1].$$

Lemma 2.2. $M_2^3[0,1]$ is a reproducing kernel Hilbert space. We get the reproducing kernel function K_y by [16]:

$$K_{y}(x) = \begin{cases} 1 + xy + \frac{x^{2}y^{2}}{4} + \frac{x^{3}y^{2}}{12} - \frac{x^{4}y}{24} + \frac{x^{5}}{120}, & x \leq y, \\ \\ 1 + yx + \frac{y^{2}x^{2}}{4} + \frac{y^{3}x^{2}}{12} - \frac{y^{4}x}{24} + \frac{y^{5}}{120}, & x > y. \end{cases}$$

Definition 2.3. We construct the reproducing kernel Hilbert space $M_2^4[0,1]$ as:

$$M_2^4[0,1] = \{a \in AC[0,1]: \ a',a'',a''' \in AC[0,1], \ a^{(4)} \in L^2[0,1], \ a(0) = a'(0) = a''(0) = 0\}.$$

We have the inner product and the norm for this special Hilbert space by:

$$\langle a,b\rangle_{M_2^4} = \sum_{i=0}^3 a^{(i)}(0)b^{(i)}(0) + \int_0^1 a^{(4)}(x)b^{(4)}(x)\mathrm{d}x, \quad a,b \in M_2^4[0,1]$$

and

$$\|a\|_{M_2^4} = \sqrt{\langle a, a \rangle_{M_2^4}}, \quad a \in M_2^4[0, 1].$$

Theorem 2.4. We find the reproducing kernel function for the reproducing kernel Hilbert space $M_2^4[0,1]$ as:

$$L_y(x) = \begin{cases} h_y(x), & x \le y, \\ \\ g_y(x), & x > y. \end{cases}$$

where,

$$h_y(x) = -\frac{x^7}{5040} + \frac{x^6y}{720} - \frac{x^5y^2}{240} + \frac{x^4y^3}{144} + \frac{x^3y^3}{36},$$
$$g_y(x) = -\frac{y^7}{5040} + \frac{y^6x}{720} - \frac{y^5x^2}{240} + \frac{y^4x^3}{144} + \frac{y^3x^3}{36}.$$

Proof. We have

$$\langle b, L_y \rangle_{M_2^4} = \sum_{i=0}^3 L_y^{(i)}(0) b^{(i)}(0) + \int_0^1 L_y^{(4)}(x) b^{(4)}(x) \mathrm{d}x,$$

by Definition 2.3. We obtain

$$\begin{split} \langle b, L_y \rangle_{M_2^4} &= L_y(0)b(0) + L_y'(0)b'(0) + L_y''(0)b''(0) + L_y'''(0)b'''(0) \\ &+ L_y^{(4)}(1)b'''(1) - L_y^{(4)}(0)b'''(0) - L_y^{(5)}(1)b''(1) \\ &+ L_y^{(5)}(0)b''(0) + L_y^{(6)}(1)b'(1) - L_y^{(6)}(0)b'(0) \\ &- \int_0^1 L_y^{(7)}(x)b'(x)\mathrm{d}x, \end{split}$$

by integration by parts. Since $b(0)=b^{\prime}(0)=b^{\prime\prime}(0)=0$ we get

$$\langle b, L_y \rangle_{M_2^4} = L_y'''(0)b'''(0) + L_y^{(4)}(1)b'''(1) - L_y^{(4)}(0)b'''(0) - L_y^{(5)}(1)b''(1)$$
$$L_y^{(6)}(1)b'(1) - \int_0^1 L_y^{(7)}(x)b'(x)dx.$$

We have

$$\begin{split} L_y'''(0) &= \frac{y^3}{6}, \\ L_y^{(4)}(0) &= \frac{y^3}{6}, \\ L_y^{(4)}(1) &= 0, \\ L_y^{(5)}(1) &= 0, \end{split}$$

$$L_y^{(6)}(1) = 0,$$

Therefore, we obtain

$$\langle b, L_y \rangle_{M_2^4} = -\int_0^y L_y^{(7)}(x) b'(x) \mathrm{d}x - \int_y^1 L_y^{(7)}(x) b'(x) \mathrm{d}x.$$

We know

$$L_y^{(7)}(x) = \begin{cases} -1, & x < y, \\ \\ 0, & x > y. \end{cases}$$

Then, we reach

$$\langle b, L_y \rangle_{M_2^4} = \int_0^y b'(x) \mathrm{d}x.$$

Thus, we get

$$\langle b, L_y \rangle_{M_2^4} = b(y).$$

This completes the proof.

We consider the solutions of the problem (1.3) in the reproducing kernel Hilbert space $M_2^4[0,1]$. We denote the bounded linear operator $T: M_2^4[0,1] \rightarrow M_2^3[0,1]$ as:

$$(2.1) \ Tu = u'''(x) + \alpha x B^2 u''(x) + 2\delta x B^2 u(x) + \delta x^2 B^2 u'(x) + 3\varepsilon x^2 B^2 u(x) + T\gamma u'(x),$$

we have the following problem.

$$(2.2) Tu = D(r, u),$$

with the initial conditions

(2.3)
$$u(0) = u'(0) = u''(0) = 0$$

where

$$\begin{split} D(r,u) &= -\alpha u(x) u'(x) u''(x) - \alpha B u(x) u''(x) - \alpha B x u'(x) u''(x) - \alpha B^2 x u''(x) \\ &- \beta u'(x) u''(x)^2 - \beta B u''(x)^2 - \delta u'(x) u(x)^2 - B \delta u(x)^2 \\ &- 2\delta B x u'(x) u(x) - \delta B^3 x^2 - \varepsilon u(x)^3 - 3\varepsilon B x u(x)^2 - \varepsilon B^3 x^3 - \gamma B. \end{split}$$

Lemma 2.5. T is a bounded linear operator.

Proof. We need to prove

$$||Tu||^2_{M^3_2[0,1]} \le A ||u||^2_{M^4_2[0,1]},$$

where A is a positive constant. We have (2.4)

$$\|Tu\|_{M_{2}^{3}[0,1]}^{2} = \langle Tu, Tu \rangle_{M_{2}^{3}[0,1]} = [Tu(0)]^{2} + [Tu(0)]^{2} + [Tu(0)]^{2} + \int_{0}^{1} [Tu(y)]^{2} dy.$$

By reproducing property, we have

$$u(y) = \langle u(\cdot), L_y(\cdot) \rangle_{M_2^4[0,1]}$$

and

$$Tu(y) = \langle u(\cdot), TL_y(\cdot) \rangle_{M_2^4[0,1]},$$

$$(Tu(y))' = \langle u(\cdot), (TL_y(\cdot))' \rangle_{M_2^4[0,1]},$$

$$(Tu(y))'' = \langle u(\cdot), (TL_y(\cdot))'' \rangle_{M_2^4[0,1]},$$

so

$$|Tu| \le ||u||_{M_2^4[0,1]} ||TL_y||_{M_2^4[0,1]} = A_1 ||u||_{M_2^4[0,1]}$$

thus

 $[Tu]^{2} \leq A_{1}^{2} \|u\|_{M_{2}^{4}[0,1]}^{2}.$

Since

$$(Tu)'(y) = \langle u(\cdot), (TL_y)'(\cdot) \rangle_{M_2^4[0,1]},$$

we get

$$|(Tu)'| \le ||u||_{M_2^4[0,1]} ||(TF_y)'||_{M_2^4[0,1]} = A_2 ||u||_{M_2^4[0,1]}$$

so, we have

$$\begin{split} \left[Tu\right]'^2 &\leq A_2^2 \left\|u\right\|_{M_2^4[0,1]}^2, \\ \left[Tu\right]''^2 &\leq A_3^2 \left\|u\right\|_{M_2^4[0,1]}^2, \\ \left[Tu\right]'''^2 &\leq A_4^2 \left\|u\right\|_{M_2^4[0,1]}^2, \end{split}$$

that is,

$$\begin{aligned} \|Tu\|_{M_{2}^{3}[0,1]}^{2} &= [Tu(0)]^{2} + [Tu(0)]^{\prime 2} + [Tu(0)]^{\prime \prime 2} + \int_{0}^{1} \left[(Tu)^{\prime \prime \prime}(y) \right]^{2} dy \\ &\leq \left(A_{1}^{2} + A_{2}^{2} + A_{3}^{2} + A_{4}^{2}\right) \left\| u \right\|_{M_{2}^{4}[0,1]}^{2}, \end{aligned}$$

where $A = A_1^2 + A_2^2 + A_3^2 + A_4^2$ is a positive constant.

We construct $\varsigma_i(x) = K_{x_i}(x)$ and $\psi_i(x) = T^* \varsigma_i(x)$, where T^* is conjugate operator of T. The orthonormal system $\left\{\widehat{\psi}_i(x)\right\}_{i=1}^{\infty}$ of $M_2^4[0,1]$ can be acquired by Gram-Schmidt orthogonalization operation of $\{\psi_i(x)\}_{i=1}^{\infty}$,

(2.5)
$$\widehat{\psi}_i(x) = \sum_{k=1}^i \beta_{ik} \psi_k(x), \quad (\beta_{ii} > 0, \quad i = 1, 2, \ldots).$$

Theorem 2.6. Let $\{x_i\}_{i=1}^{\infty}$ be dense in [0,1] and $\psi_i(x) = T_y L_x(y)|_{y=u_i}$. Then the sequence $\{\psi_i(x)\}_{i=1}^{\infty}$ is a complete system in $M_2^4[0,1]$.

Proof. We get

$$\psi_i(x) = (T^*\varsigma_i)(x) = \langle (T^*\varsigma_i)(y), L_x(y) \rangle = \langle (\varsigma_i)(y), T_y L_x(y) \rangle = T_y L_x(y)|_{y=x_i}.$$

Let $\langle u(x), \psi_i(x) \rangle = 0$, (i = 1, 2, ...), which means that,

$$\langle u(x), (T^*\varsigma_i)(x) \rangle = \langle Tu(\cdot), \varsigma_i(\cdot) \rangle = (Tu)(x_i) = 0.$$

 $\{x_i\}_{i=1}^{\infty}$ is dense in [0, 1]. Therefore, (Tu)(x) = 0. $u \equiv 0$ by T^{-1} .

Theorem 2.7. If u(x) is the exact solution of (2.2), then we acquire

(2.6)
$$u(x) = \sum_{i=1}^{\infty} \sum_{k=1}^{i} \beta_{ik} D(x_k, u_k) \widehat{\psi}_i(x).$$

where $\{x_i\}_{i=1}^{\infty}$ is dense in [0, 1].

Proof. We get

$$\begin{split} u(x) &= \sum_{i=1}^{\infty} \left\langle u(x), \widehat{\psi}_{i}(x) \right\rangle_{M_{2}^{4}[0,1]} \widehat{\psi}_{i}(x) \\ &= \sum_{i=1}^{\infty} \sum_{k=1}^{i} \beta_{ik} \left\langle u(x), \psi_{k}(x) \right\rangle_{M_{2}^{4}[0,1]} \widehat{\psi}_{i}(x) \\ &= \sum_{i=1}^{\infty} \sum_{k=1}^{i} \beta_{ik} \left\langle u(x), T^{*}\varsigma_{k}(x) \right\rangle_{M_{2}^{4}[0,1]} \widehat{\psi}_{i}(x) \\ &= \sum_{i=1}^{\infty} \sum_{k=1}^{i} \beta_{ik} \left\langle Tu(x), \varsigma_{k}(x) \right\rangle_{M_{2}^{3}[0,1]} \widehat{\psi}_{i}(x) \\ &= \sum_{i=1}^{\infty} \sum_{k=1}^{i} \beta_{ik} \left\langle D(r,u) K_{x_{k}} \right\rangle_{M_{2}^{3}[0,1]} \widehat{\psi}_{i}(x) \\ &= \sum_{i=1}^{\infty} \sum_{k=1}^{i} \beta_{ik} D(x_{k}, u_{k}) \widehat{\psi}_{i}(x). \end{split}$$

by uniqueness of solution of (2.2). This completes the proof.

The approximate solution $u_n(x)$ can be obtained as:

(2.7)
$$u_n(x) = \sum_{i=1}^n \sum_{k=1}^i \beta_{ik} D(x_k, u_k) \widehat{\psi}_i(x).$$

3 Numerical Experiments

We apply the RKM for the approximate analytical solution of the Jerk Equation (1.1)-(1.2).
Example 3.1. We take into consideration

(3.1)
$$y''' = -y' + yy'y'',$$

with initial conditions y(0) = 0, y'(0) = B, y''(0) = 0. We have the exact solution as:

(3.2)
$$y(x) = \frac{B}{\Omega}\sin(\Omega x) + \frac{B}{96\Omega^3}((-9B^2 - 48 + 48\Omega^2)\sin(\Omega x) - B^2\sin(3\Omega x))$$

where $\Omega = \frac{1}{2}\sqrt{B^2 + 4}$.

Table 1 shows the absolute errors for B = 0.2, 0.3, 0.4 respectively. Table 2 shows the relative errors for B = 0.2, 0.3, 0.4 respectively.

Example 3.2. We investigate

(3.3)
$$y''' = -y' - y'(y'')^2,$$

with initial conditions y(0) = 0, y'(0) = B, y''(0) = 0. We get exact solution as:

$$y(x) = \frac{B}{\Omega}\sin(\Omega x) + \frac{B}{96\Omega^3}((-9B^2\Omega^2 - 48 + 48\Omega^2)\sin(\Omega x) + ((12B^2(\Omega)^3 - 48\Omega + 48(\Omega)^3)x\cos(\Omega x) - B^2\Omega^2\sin(3\Omega x))$$

where $\Omega = 2\sqrt{\frac{1}{4-B^2}}$.

Table 3 shows the absolute errors for B = 0.2, 0.3, 0.4 respectively. Table 4 shows the relative errors for B = 0.2, 0.3, 0.4 respectively.

 Table 1. Absolute Errors for the first example.

x	B = 0.2	<i>B</i> = 0.3	<i>B</i> = 0.4
0.125	2.8088×10^{-7}	4.2077×10^{-7}	5.6041×10^{-7}
0.250	0.0000011935	0.0000017943	0.0000024143
0.375	0.0000027482	0.0000041708	0.0000057546
0.500	0.0000049504	0.0000076345	0.0000109688
0.625	0.0000078006	0.0000122943	0.0000185799
0.750	0.0000112879	0.0000182374	0.0000290702
0.875	0.0000153830	0.0000254769	0.0000426444
1.00	0.0000200317	0.0000339035	0.0000590332
1.125	0.0000233859	0.0000405546	0.0000736859
1.250	0.0001115761	0.0000089211	0.0000352990
1.375	0.0001115761	0.0001651283	0.0002024070

x	B = 0.2	B = 0.3	B = 0.4
0.125	0.00001126451361	0.00001124981105	0.00001123744895
0.250	0.00002412134243	0.00002417518187	0.00002439700359
0.375	0.00003751627443	0.00003795830294	0.00003928026318
0.500	0.00005163110050	0.00005308586566	0.00005720732970
0.625	0.00006666731728	0.00007005716478	0.00007942052378
0.750	0.00008281647176	0.00008922519004	0.00010670764110
0.875	0.00010024618200	0.00011073531080	0.00013910935150
1.00	0.00011910109620	0.00013449058700	0.00017583060030
1.125	0.00012972006220	0.00015015336190	0.00020497726770
1.250	0.00000985305713	0.00003143790375	0.00009353950721
1.375	0.00056991392360	0.00056376716640	0.00052021705710

 Table 2. Relative Errors for the first example.

 Table 3. Absolute Errors for the second example.

x	<i>B</i> = 0.2	<i>B</i> = 0.3	<i>B</i> = 0.4
0.125	2.8088×10^{-7}	4.2077×10^{-7}	5.6041×10^{-7}
0.250	0.0000011935	0.0000017943	0.0000024143
0.375	0.0000027482	0.0000041708	0.0000057546
0.500	0.0000049504	0.0000076345	0.0000109688
0.625	0.0000078006	0.0000122943	0.0000185799
0.750	0.0000112879	0.0000182374	0.0000290702
0.875	0.0000153830	0.0000254769	0.0000426444
1.00	0.0000200317	0.0000339035	0.0000590332
1.125	0.0000233859	0.0000405546	0.0000736859
1.250	0.0001115761	0.0000089211	0.0000352990
1.375	0.0001115761	0.0001651283	0.0002024070

x	B = 0.2	B = 0.3	B = 0.4
0.125	0.00001126451361	0.00001124981105	0.00001123744895
0.250	0.00002412134243	0.00002417518187	0.00002439700359
0.375	0.00003751627443	0.00003795830294	0.00003928026318
0.500	0.00005163110050	0.00005308586566	0.00005720732970
0.625	0.00006666731728	0.00007005716478	0.00007942052378
0.750	0.00008281647176	0.00008922519004	0.00010670764110
0.875	0.00010024618200	0.00011073531080	0.00013910935150
1.00	0.00011910109620	0.00013449058700	0.00017583060030
1.125	0.00012972006220	0.00015015336190	0.00020497726770
1.250	0.00000985305713	0.00003143790375	0.00009353950721
1.375	0.00056991392360	0.00056376716640	0.00052021705710

 Table 4. Relative Errors for the second example.

4 Conclusions

In this paper, we investigated the nonlinear Jerk equations by the reproducing kernel method. We constructed very useful reproducing kernel Hilbert spaces and we found some important reproducing kernel functions in these spaces. We found a bounded linear operator to get the results for the problems. We obtained absolute errors and relative errors for some numerical experiments. We proved the efficiency of the proposed method in the work.

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Solving a New Type of Fractional Differential Equation by Reproducing Kernel Method

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Abstract. The aim of this work is to get the solutions of the fractional counterpart of a boundary value problem by implementing the reproducing kernel Hilbert space method. Convergence of the solution problem discussed has been shown. The efficiency of the proposed technique is demonstrated by some tables.

1 Introduction

We consider the following problem [2]:

(1.1)
$$\frac{d^2\omega}{dr^2} + \frac{1}{r}\frac{d\omega}{dr} + Ha^2\left(1 - \frac{\omega}{1 - \alpha\omega}\right) = 0, \quad 0 \ \left(\ r \left(1, \right) \right)$$

where $\omega(r)$ is the velocity of the fluid, r is the radial distance from the cylindrical conduit centre, Ha is the Hartmann electric number and α is the magnitude of the power of non-linearity. We have the boundary conditions as:

(1.2)
$$\omega'(0) = 0, \quad \omega(1) = 0.$$

The existence and uniqueness of a solution to the problem have been investigated in [3]. Mastroberardino [4] has investigated the problem by the homotopy analysis method. Moghtadaei et al. [5] have applied a spectral method to investigate the problem. Chebyshev spectral collocation method has been used to solve the problem in [6]. Alomari et al. [7] have investigated fractional version of a singular boundary value problem.

In this paper, we consider the following problem.

(1.3)
$$\frac{d^{\gamma}\omega}{dr^{\gamma}} + \frac{1}{r}\frac{d^{\beta}\omega}{dr^{\beta}} + Ha^{2}\left(1 - \frac{\omega}{1 - \alpha\omega}\right) = 0,$$

where d^{γ}/dr^{γ} and d^{β}/dr^{β} are the fractional derivative operators in the Caputo sense. $\gamma \in (1, 2]$ and $\beta \in (0, 1]$ are parameters defining the order of the fractional © Springer Nature Switzerland AG 2021

Z. Hammouch et al. (Eds.): SM2A 2019, LNNS 168, pp. 34–43, 2021. https://doi.org/10.1007/978-3-030-62299-2_3 derivative with the property $\gamma - \beta \ge 1$ and subject to the boundary conditions (1.2). Abbas et al. [8] have presented some cases that show the fractional models present better approximate results. Iyiolaa et al. [9] have worked the cancer tumor model of fractional order which demonstrates better approximate results.

We apply reproducing kernel method (RKM) to get the approximate solutions of Eq. (1.3). Reproducing kernel space is a special Hilbert space. Many investigators have applied the RKM to many problems [10]. Arqub et al. [11,12] have investigated some interesting problems by RKM.

This paper is organized as follows. In Sect. 2, RKM is discussed with some preliminary concepts and definitions. The solution procedure and approximate solutions of Eqs. (1.2)-(1.3) are presented in this section. Numerical experiments are demonstrated in Sect. 3. Conclusion is given in the last section.

2 Reproducing Kernel Method

First of all we will construct the reproducing kernel Hilbert spaces that we need to solve our problem.

Definition 2.1. The first reproducing kernel Hilbert space that we will use is $E_2^1 [0,1]$

$$E_2^1 [0,1] = \{ s \in AC [0,1] : s' \in L^2 [0,1] \}.$$

We have the inner product and the norm for this space as:

$$\langle s, p \rangle_{E_2^1} = s(0)p(0) + \int_0^1 s'(\tau)p'(\tau)d\tau, \quad s, p \in E_2^1[0,1]$$

and

$$\|s\|_{E_2^1} = \sqrt{\langle s,s\rangle_{E_2^1}}, \quad s \in E_2^1 \left[\begin{array}{c} 0,1 \end{array} \right].$$

Lemma 2.2. $E_2^1[0,1]$ is a reproducing kernel Hilbert space. We get the reproducing kernel function G_z by [10]:

$$G_z(\tau) = \begin{cases} 1+\tau, & \tau \leq z, \\ 1+z, & \tau \rangle z. \end{cases}$$

Definition 2.3. We construct the reproducing kernel Hilbert space $E_2^3 [0,1]$ as:

$$E_2^3 [0,1] = \{ s \in AC [0,1] : s', s'' \in AC [0,1], s^{(3)} \in L^2 [0,1], s'(0) = 0 = s(1) \}.$$

We have the inner product and the norm for this special Hilbert space by:

$$\langle s, p \rangle_{E_2^3} = \sum_{i=0}^2 s^{(i)}(0) p^{(i)}(0) + \int_0^1 s^{(3)}(\tau) p^{(3)}(\tau) d\tau, \quad s, p \in E_2^3 [0, 1]$$

and

$$\|s\|_{E_2^3} = \sqrt{\langle s,s\rangle_{E_2^3}}, \quad s \in E_2^3 \left[\ 0,1 \ \right].$$

Theorem 2.4. We find the reproducing kernel function for the reproducing kernel Hilbert space $E_2^3 [0,1]$ as:

$$F_z(\tau) = \begin{cases} h_z(\tau), & \tau \le z, \\ \\ g_z(\tau), & \tau \rangle z. \end{cases}$$

Where,

$$\begin{split} h_z(\tau) &= -\frac{\tau^5 z^2}{624} + \frac{\tau^5 z^4}{3744} - \frac{\tau^5 z^3}{1872} + \frac{5\tau^4 z^3}{1872} + \frac{5\tau^4 z^2}{624} - \frac{\tau^5 z^5}{18720} + \frac{\tau^4 z^5}{3744} \\ &- \frac{5\tau^4 z^4}{3744} + \frac{5\tau^3 z^4}{1872} - \frac{5\tau^3 z^3}{936} - \frac{\tau^4 z}{24} - \frac{5\tau^2 z^3}{312} + \frac{7\tau^3 z^2}{104} - \frac{\tau^3 z^5}{1872} \\ &- \frac{\tau^2 z^5}{624} + \frac{5\tau^2 z^4}{624} + \frac{21\tau^2 z^2}{104} - \frac{5z^2}{26} + \frac{\tau^5}{520} + \frac{5z^5}{156} - \frac{5z^3}{78} + \frac{5\tau^4}{156} \\ &- \frac{z^5}{156} - \frac{5\tau^3}{78} - \frac{5\tau^2}{26} + \frac{3}{13}, \end{split}$$

$$g_{z}(\tau) = -\frac{z^{5}\tau^{2}}{624} + \frac{z^{5}\tau^{4}}{3744} - \frac{z^{5}\tau^{3}}{1872} + \frac{5z^{4}\tau^{3}}{1872} + \frac{5z^{4}\tau^{2}}{624} - \frac{z^{5}\tau^{5}}{18720} + \frac{z^{4}\tau^{5}}{3744}$$
$$-\frac{5z^{4}\tau^{4}}{3744} + \frac{5z^{3}\tau^{4}}{1872} - \frac{5z^{3}\tau^{3}}{936} - \frac{z^{4}\tau}{24} - \frac{5z^{2}\tau^{3}}{312} + \frac{7z^{3}\tau^{2}}{104} - \frac{z^{3}\tau^{5}}{1872}$$
$$-\frac{z^{2}\tau^{5}}{624} + \frac{5z^{2}\tau^{4}}{624} + \frac{21z^{2}\tau^{2}}{104} - \frac{5\tau^{2}}{26} + \frac{z^{5}}{520} + \frac{5\tau^{5}}{156} - \frac{5\tau^{3}}{78} + \frac{5z^{4}}{156}$$
$$-\frac{\tau^{5}}{156} - \frac{5z^{3}}{78} - \frac{5z^{2}}{26} + \frac{3}{13}.$$

Proof. We have

$$\langle p, F_z \rangle_{E_2^3} = \sum_{i=0}^2 F_z^{(i)}(0) p^{(i)}(0) + \int_0^1 F_z^{(3)}(\tau) p^{(3)}(\tau) d\tau,$$

by Definition 2.3. We obtain

$$\begin{split} \langle p, F_z \rangle_{E_2^3} &= F_z(0) p(0) + F_z'(0) p'(0) + F_z''(0) p''(0) \\ &+ F_z'''(1) p''(1) - F_z'''(0) p''(0) - F_z^{(4)}(1) p'(1) \\ &+ F_z^{(4)}(0) p'(0) + \int_0^1 F_z^{(5)}(\tau) p'(\tau) \mathrm{d}\tau, \end{split}$$

by integration by parts. Since p'(0) = 0 = p(1), we get

$$\langle p, F_z \rangle_{E_2^3} = F_z(0)p(0) + F_z''(0)p''(0) + F_z'''(1)p''(1) - F_z'''(0)p''(0) - F_z^{(4)}(1)p'(1) + \int_0^1 F_z^{(5)}(\tau)p'(\tau)d\tau.$$

We have

$$F_z(0) = -\frac{z^5}{156} + \frac{5z^4}{156} - \frac{5z^3}{178} - \frac{5z^2}{26} + \frac{3}{13},$$

$$F_z''(0) = \frac{21z^2}{52} - \frac{z^5}{312} + \frac{5z^4}{312} - \frac{5z^3}{156} - \frac{5}{13},$$

$$F_{z}^{\prime\prime\prime}(0) = \frac{21z^{2}}{52} - \frac{z^{5}}{312} + \frac{5z^{4}}{312} - \frac{5z^{3}}{156} - \frac{5}{13},$$

$$F_z^{\prime\prime\prime}(1) = F_z^{(4)}(1) = 0,$$

Therefore, we obtain

$$\langle p, F_z \rangle_{E_2^3} = \left(-\frac{z^5}{156} + \frac{5z^4}{156} - \frac{5z^3}{178} - \frac{5z^2}{26} + \frac{3}{13} \right) p(0) + \int_0^z F_z^{(5)}(\tau) p'(\tau) d\tau + \int_z^1 F_z^{(5)}(\tau) p'(\tau) d\tau.$$

We know

$$F_{z}^{(5)}(\tau) = \begin{cases} -\frac{z^{5}}{156} + \frac{5z^{4}}{156} - \frac{5z^{3}}{178} - \frac{5z^{2}}{26} + \frac{3}{13}, & \tau \langle z, \\ -\frac{z^{5}}{156} + \frac{5z^{4}}{156} - \frac{5z^{3}}{178} - \frac{5z^{2}}{26} - \frac{10}{13}, & \tau \rangle z. \end{cases}$$

Then, we reach

$$\langle p, F_z \rangle_{E_2^3} = \left(-\frac{z^5}{156} + \frac{5z^4}{156} - \frac{5z^3}{178} - \frac{5z^2}{26} + \frac{3}{13} \right) p(0) + \int_0^z \left(-\frac{z^5}{156} + \frac{5z^4}{156} - \frac{5z^3}{178} - \frac{5z^2}{26} + \frac{3}{13} \right) p'(\tau) \mathrm{d}\tau + \int_z^1 \left(-\frac{z^5}{156} + \frac{5z^4}{156} - \frac{5z^3}{178} - \frac{5z^2}{26} - \frac{10}{13} \right) p'(\tau) \mathrm{d}\tau.$$

Thus, we obtain

$$\begin{split} \langle p,F_z\rangle_{E_2^3} &= \big(-\frac{z^5}{156} + \frac{5z^4}{156} - \frac{5z^3}{178} - \frac{5z^2}{26} + \frac{3}{13}\big)p(0) \\ &\quad + \big(-\frac{z^5}{156} + \frac{5z^4}{156} - \frac{5z^3}{178} - \frac{5z^2}{26} + \frac{3}{13}\big)p(z) \\ &\quad - \big(-\frac{z^5}{156} + \frac{5z^4}{156} - \frac{5z^3}{178} - \frac{5z^2}{26} + \frac{3}{13}\big)p(0) \\ &\quad + \big(-\frac{z^5}{156} + \frac{5z^4}{156} - \frac{5z^3}{178} - \frac{5z^2}{26} - \frac{10}{13}\big)p(1) \\ &\quad - \big(-\frac{z^5}{156} + \frac{5z^4}{156} - \frac{5z^3}{178} - \frac{5z^2}{26} - \frac{10}{13}\big)p(z). \end{split}$$

Therefore, we obtain

$$\langle p, F_z \rangle_{E_2^3} = p(z).$$

This completes the proof.

We consider the solutions of the problem (1.3) in the reproducing kernel Hilbert space $E_2^3 [0,1]$. We denote the bounded linear operator $X : E_2^3 [0,1] \rightarrow E_2^1 [0,1]$ as:

(2.1)
$$X\omega = \frac{d^{\gamma}\omega}{dr^{\gamma}} + \frac{1}{r}\frac{d^{\beta}\omega}{dr^{\beta}}.$$

Then, we have the following problem.

(2.2)
$$X\omega = D(r,\omega),$$

with the boundary conditions

$$(2.3) \qquad \qquad \omega'(0) = 0 = \omega(1),$$

where

(2.4)
$$D(r,\omega) = -Ha^2 \left(1 - \frac{\omega}{1 - \alpha\omega}\right)$$

Lemma 2.5. X is a bounded linear operator.

Proof. We need to prove

$$\|X\omega\|_{E_2^1[0,1]}^2 \le K \|\omega\|_{E_2^3[0,1]}^2,$$

where K is a positive constant. We have

(2.5)
$$\|X\omega\|_{E_2^1[0,1]}^2 = \langle X\omega, X\omega \rangle_{E_2^1[0,1]} = [X\omega(0)]^2 + \int_0^1 [X\omega'(z)]^2 dz.$$

By reproducing property, we have

$$\omega(z) = \langle \omega(\cdot), F_z(\cdot) \rangle_{E_2^3[0,1]}$$

and

$$X\omega(z) = \langle \omega(\cdot), XF_z(\cdot) \rangle_{E_2^3[0,1]},$$

 \mathbf{SO}

$$|X\omega| \le \|\omega\|_{E_2^3[0,1]} \|XF_z\|_{E_2^3[0,1]} = K_1 \|\omega\|_{E_2^3[0,1]},$$

thus

$$[X\omega(0)]^2 \le K_1^2 \|\omega\|_{E_2^3[0,1]}^2$$

Since

$$(X\omega)'(z) = \langle \omega(\cdot), (XF_z)'(\cdot) \rangle_{E_2^3[0,1]},$$

we get

$$|(X\omega)'| \le ||\omega||_{E_2^3[0,1]} ||(XF_z)'||_{E_2^3[0,1]} = K_2 ||\omega||_{E_2^3[0,1]},$$

so, we have

$$[X\omega]^2 \le K_2^2 \|\omega\|_{E_2^3[0,1]}^2,$$

that is,

$$\|X\omega\|_{E_{2}^{1}[0,1]}^{2} = [X\omega(0)]^{2} + \int_{0}^{1} [(X\omega)'(z)]^{2} dz \le (K_{1}^{2} + K_{2}^{2}) \|\omega\|_{E_{2}^{3}[0,1]}^{2},$$

where $K = K_1^2 + K_2^2$ is a positive constant.

We construct $\zeta_i(\tau) = G_{\tau_i}(\tau)$ and $\psi_i(\tau) = X^* \zeta_i(\tau)$, where X^* is conjugate operator of X. The orthonormal system $\{\widehat{\psi}_i(\tau)\}_{i=1}^{\infty}$ of $E_2^3[0,1]$ can be acquired by Gram-Schmidt orthogonalization operation of $\{\psi_i(\tau)\}_{i=1}^{\infty}$,

(2.6)
$$\widehat{\psi}_i(\tau) = \sum_{k=1}^i \beta_{ik} \psi_k(\tau), \quad (\beta_{ii}) 0, \quad i = 1, 2, \ldots).$$

Theorem 2.6. Let $\{\tau_i\}_{i=1}^{\infty}$ be dense in [0,1] and $\psi_i(\tau) = X_z F_\tau(z)|_{z=\tau_i}$. Then the sequence $\{\psi_i(\tau)\}_{i=1}^{\infty}$ is a complete system in $E_2^3 [0,1]$.

Proof. We get

$$\psi_i(\tau) = (X^*\varsigma_i)(\tau) = \langle (X^*\varsigma_i)(z), F_\tau(z) \rangle = \langle (\varsigma_i)(z), X_z F_\tau(z) \rangle = X_z F_\tau(z)|_{z=\tau_i}.$$

Let $\langle \omega(\tau), \psi_i(\tau) \rangle = 0$, (i = 1, 2, ...), which means that,

$$\langle \omega(\tau), (X^*\varsigma_i)(\tau) \rangle = \langle X\omega(\cdot), \varsigma_i(\cdot) \rangle = (X\omega)(\tau_i) = 0.$$

 $\{\tau_i\}_{i=1}^{\infty}$ is dense in [0,1]. Therefore, $(X\omega)(\tau) = 0$. $\omega \equiv 0$ by X^{-1} .

Theorem 2.7. If $\omega(r)$ is the exact solution of (2.2), then we acquire

(2.7)
$$\omega(r) = \sum_{i=1}^{\infty} \sum_{k=1}^{i} \beta_{ik} D(r_k, \omega_k) \widehat{\psi}_i(r).$$

where $\{r_i\}_{i=1}^{\infty}$ is dense in [0,1].

Proof. We get

$$\omega(r) = \sum_{i=1}^{\infty} \left\langle \omega(r), \widehat{\psi}_i(r) \right\rangle_{E_2^3[0,1]} \widehat{\psi}_i(r)$$
$$= \sum_{i=1}^{\infty} \sum_{k=1}^{i} \beta_{ik} \left\langle \omega(r), \psi_k(r) \right\rangle_{E_2^3[0,1]} \widehat{\psi}_i(r)$$
$$= \sum_{i=1}^{\infty} \sum_{k=1}^{i} \beta_{ik} \left\langle \omega(r), X^* \varsigma_k(r) \right\rangle_{E_2^3[0,1]} \widehat{\psi}_i(r)$$
$$= \sum_{i=1}^{\infty} \sum_{k=1}^{i} \beta_{ik} \left\langle X \omega(r), \varsigma_k(r) \right\rangle_{E_2^1[0,1]} \widehat{\psi}_i(r)$$

$$= \sum_{i=1}^{\infty} \sum_{k=1}^{i} \beta_{ik} \langle D(r,\omega), G_{r_k} \rangle_{E_2^1[0,1]} \widehat{\psi}_i(r)$$
$$= \sum_{i=1}^{\infty} \sum_{k=1}^{i} \beta_{ik} D(r_k, \omega_k) \widehat{\psi}_i(r).$$

by uniqueness of solution of (2.2). This completes the proof.

The approximate solution $\omega_n(r)$ can be obtained as:

(2.8)
$$\omega_n(r) = \sum_{i=1}^n \sum_{k=1}^i \beta_{ik} D(r_k, \omega_k) \widehat{\psi}_i(r).$$

3 Numerical Experiments

In Table 1, we show the solution when $\gamma = 1.9$, $\beta = 0.9$, $\alpha = 0.5$ and vary the Hartmann electric number. In Table 2, we fixed the fractional derivatives as $\gamma = 1.9$, $\beta = 0.9$, $Ha^2 = 1.0$ and vary the α . In Table 3, we give the solution with the fractional derivatives $\gamma = 1.3$, $\beta = 0.3$, $Ha^2 = 1.0$ and vary the α .

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r	$Ha^2 = 1$	$Ha^2 = 2$
0.0	0.2299695869	0.3825201644
0.1	0.2266810832	0.3754454339
0.2	0.2180411628	0.3600500656
0.3	0.2062671368	0.3434211029
0.4	0.1900081364	0.3158370163
0.5	0.1697602524	0.2874089253
0.6	0.1451770620	0.2475168041
0.7	0.1161364572	0.2007161568
0.8	0.0824312374	$0.1443\overline{651864}$
0.9	0.0438131825	0.0777826865
1.0	-8.95×10^{-10}	-3.54×10^{-11}

Table 1. Approximate solutions by reproducing kernel method for $\gamma = \frac{19}{10}$, $\beta = \frac{9}{10}$, $\alpha = \frac{1}{2}$ and different values of Ha^2 .

Table 2. Approximate solutions by reproducing kernel method for $\gamma = \frac{19}{10}$, $\beta = \frac{9}{10}$, $Ha^2 = 1$ and different values of α .

r	α = 1	α = 2
0.0	0.2259726183	0.2093112589
0.1	0.2226475274	0.2058336031
0.2	0.2141458034	0.1979096087
0.3	0.2029104831	0.1888516297
0.4	0.1870087144	0.1744727611
0.5	0.1673526319	0.1579851722
0.6	0.1433018034	0.1361693674
0.7	0.1147939073	0.1098299577
0.8	0.0815892945	0.0785261212
0.9	0.0434202449	0.0419967019
1.0	-9.79×10^{-10}	-3.97×10^{-10}

r	α = 1	α = 2
0.0	0.4454618088	0.3386406399
0.1	0.4284599310	0.3204729467
0.2	0.3272191714	0.2609343009
0.3	0.2498402621	0.2849602586
0.4	0.2020451400	0.3044655878
0.5	0.1890369644	0.2413493934
0.6	0.2063732634	0.1921955815
0.7	0.1160557551	0.1449551875
0.8	0.0896312152	0.1026976873
0.9	0.0478590068	0.0537369610
1.0	2.77×10^{-6}	-8.159×10^{-7}

Table 3. Approximate solutions by reproducing kernel method for $\gamma = \frac{13}{10}$, $\beta = \frac{3}{10}$, $Ha^2 = 1$ and different values of α .

4 Conclusions

In this work, we acquired the solutions of fractional version of a singular boundary value problem occurring in the electrohydrodynamic flow in a circular cylindrical conduit based on the reproducing kernel method. We demonstrated our effective results by some tables. We investigated the effect of the Hartmann electric number and the fractional order of the problem. We concluded that the reproducing kernel method can be applied to much more complicated fractional differential equations.

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An Efficient Approach for the Model of Thrombin Receptor Activation Mechanism with Mittag-Leffler Function

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Abstract. In the present work, we haired an efficient technique called, q-homotopy analysis transform method (q-HATM) in order to find the solution for the model of thrombin receptor activation mechanism (TRAM) and examine the nature of *q*-HATM solution with distinct fractional order. The considered model elucidates the TRA mechanism in calcium signalling, and this mechanism plays a vital role in the human body. We defined fractional derivative defined with Atangana-Baleanu (AB) operator and the projected scheme is an amalgamation of Laplace transform with q-homotopy analysis scheme. For the achieved results, to present the existence and uniqueness we hired the fixed point hypothesis. To validate and illustrate the effectiveness of the considered scheme, we examined the projected model with arbitrary order. The behaviour of the achieved results is captured in terms of plots and also showed the importance of the parameters offered by the considered solution procedure. The attained results illuminate, the projected scheme is easy to employ and more effective in order to analyse the behaviour of fractional order differential systems exemplifying real word problems associated with science and technology.

Keywords: *q*-Homotopy analysis method \cdot Laplace transform \cdot Thrombin receptor activation mechanism \cdot Fixed point theorem \cdot Atangana-Baleanu derivative

1 Introduction

The human body is mainly the composition of six elements of about 99%, and those are namely, phosphorus, calcium, nitrogen, hydrogen, carbon and oxygen. In our body, the most generous mineral is calcium (*Ca*) and it is about 1.5%. *Ca* is the most play a vital rule in muscle contraction and protein regulation and also it's very essential in the processes of contractions bones and their protection. Most of the phenomena including cell death and fertilization are achieved with the help of calcium oscillations. With the

exploit of the inositol phospholipid cascade by raising the cytosolic calcium levels produced, most of the pathways of signal transduction are arbitrated [1, 2]. Ca acts as emissary in information processing. For the analysis of enzyme phospholipase C (PLC), G protein is playing an important role.

In the present scenario, the study of most of the phenomena related to the human body like diseases and their behaviour, the essential components of our body and their functions; magnetize the attention of mathematicians and researchers associated to mathematics in order to model and analyse as well as predict its essential behaviours. In connection with this, authors in [3] nurtured the mathematical model in order to illustrate the mediated activation of human platelets, researchers in [4] analyse the cytosolic calcium dynamics by the aid of mathematical model and later by the help of fractional calculus (FC), authors in [5] present their viewpoint in order to understand the importance of FC while analysing the mathematical model stimulating the abovecited phenomenon.

The seed of fractional calculus (FC) is planted before 324 years, however lately become an essential tool for the distinct discipline of science and engineering, and hence fascinated the attention of authors. It was shortly discovered that fractional calculus is more appropriate for modelling the phenomena describing nature in a systematic manner as associated with integer order calculus. The calculus of arbitrary order turned out one of the most essential tools to describe biological phenomena. The human diseases which are modelled through derivative having fractional-order help us to incorporate the information about its present and past states. Diverse pioneering notions and fundamentals are prescribed by many senior researchers [6-11]. Recently, due to diverse applications and favourable properties, the concept of FC is widely hired to investigate real world problems [12–20]. Particularly, authors in [21] analysed the fractional order system exemplifying the fish farm model within the frame of new fractional operator and also the captured some simulating consequences associated to the model using efficient scheme, authors in [22, 23] investigated the numerical solution for the fractional order coupled special cases of KdV equations and presented some interesting results with respect to different fractional order. The epidemic model of childhood disease is analysed by the authors in [24] within the frame of fractional calculus and they presented the nature of the corresponding results for distinct arbitrary order. Authors in [19] analyse the evolution of 2019-nCoV and its dynamic structures with help of nonlocal operator and presented some numerical surface using efficient scheme.

The activated form of phospholipase C(PLC) hydrolyzes the diacylglycerol (*DAG*), 5-trisphosphate [*Ins*(*I*, 4, 5)*P*₃], 5-bisphosphate [*PtdIns*(4, 5)*P*₂] to inositol and phosphatidylinositol 4. From the endoplasmic reticulum, the 5-trisphosphate is helps to stimulate the let out of endogenous calcium. The number of activated cell surface receptor proportional to the rate of generates of the 5-trisphosphate. The thrombin is a multiprocessing serine protease aids from the endothelial cell to take calcium transient and it acts as a ligand for the present model. Here, we consider the system of the equation which described the TRS mechanism. In endothelial cells, this model provides incite of calcium arbitrated signal transduction. The release of calcium is determined by the 5-trisphosphate cytosolic level in the calcium homeostasis and the number of active surface receptors (*S*) aid to generate the 5-trisphosphate. The receptor-ligand complex (*C*) formed due to ligand binding with surface receptors and on cleavage outcomes in activated receptors (A). The above-cited phenomenon is illustrated with the aid of the system of three differential equations and concentration of thrombin (ε) as follows [4, 5]

$$\frac{dS(t)}{dt} = -\delta\varepsilon S(t) + \beta C(t)$$

$$\frac{dC(t)}{dt} = \delta\varepsilon S(t) - (\beta + \lambda)C(t), \qquad (1)$$

$$\frac{dA(t)}{dt} = \lambda C(t)$$

where δ and β respectively symbolise the on and off rate constant of thrombin binding.

Many nonlinear and important models are effectively and methodically examined with the assist of FC. Many senior pioneers proposed distinct definitions including, Riemann, Liouville, Caputo and Fabrizio. Soon after the invention of each notion, many researchers identify some limitations while examining specific problems. Including physical meaning of the initial conditions, kernel associated to singularity, non-locality and others associated with complex phenomena. With the assist of Mittag–Leffler function, Atangana and Baleanu [25] proposed a new fractional-order operator and overcome all the above-cited consequences which play a vital role while investigating properties of the models.

Authors in [5] presented the simulation for the fractional system with Caputo-Fabrizio derivative using perturbation iterative scheme, which poses interesting consequences. In the present framework, we consider with AB derivative and which as follows

$$\begin{split} & {}^{ABC}_{a} D^{\alpha}_{t} S(t) = -\delta \varepsilon S(t) + \beta C(t) \\ & {}^{ABC}_{a} D^{\alpha}_{t} C(t) = \delta \varepsilon S(t) - (\beta + \lambda) C(t), \\ & {}^{ABC}_{a} D^{\alpha}_{t} A(t) = \lambda C(t), \end{split}$$

$$\end{split}$$

$$(2)$$

where α is fractional order of the system.

As much as impartment of modelling real-world problems, finding the solution for the corresponding system is also vital and difficult. Most of the complex and nonlinear problems don't have an analytical solution. In this connection, researchers preferred for semi-analytical or numerical schemes. One of the efficient and most widely hired methods to solve nonlinear problems is the homotopy analysis method (HAM) and which natured by *Liao Shijun* [26, 27]. This solution procedure overcomes most of the limitation arise while solving nonlinear problems with dissertation and perturbation. However, a few limitations have been pointed out by researchers in order to reduce computational work and time. The presented method is the mixture of LT with q-HAM and nurtured by Singh et al. [28]. Clearly, q-HATM is an enhanced algorithm of HAM; it does not require linearization, perturbation or discretization. Recently, many researchers hired the considered method due to its efficacy and reliability to understand physical behaviour numerous classes of nonlinear problems [29–34]. The considered scheme gives more freedom to choose problems associated with distinct initial conditions and it proposed with axillary and homotopy corresponding phenomena [35, 36].

2 Preliminaries

Here, the basic notions and definitions of FC and LT are presented [25, 37-41].

Definition 1. In Caputo and Riemann-Liouville sense, for a function $f \in H^1(a, b)$ the fractional Atangana-Baleanu-derivative are presented respectively as follows [25]:

$${}^{ABC}_{a}D^{\alpha}_{t}(f(t)) = \frac{\mathcal{B}[\alpha]}{1-\alpha} \int_{a}^{t} f'(\vartheta) E_{\alpha} \left[\alpha \frac{(t-\vartheta)^{\alpha}}{\alpha-1} \right] d\vartheta.$$
(3)

$${}^{ABR}_{a}D^{\alpha}_{t}(f(t)) = \frac{\mathcal{B}[\alpha]}{1-\alpha}\frac{d}{dt}\int_{a}^{t}f(\vartheta)E_{\alpha}\left[\alpha\frac{(t-\vartheta)^{\alpha}}{\alpha-1}\right]d\vartheta,\tag{4}$$

where $\mathcal{B}[\alpha]$ is a normalization function such that $\mathcal{B}(0) = \mathcal{B}(1) = 1$.

Definition 2. The AB integral with fractional order is presented [25] as

$${}^{AB}_{a}I^{\alpha}_{t}(f(t)) = \frac{1-\alpha}{\mathcal{B}[\alpha]}f(t) + \frac{\alpha}{\mathcal{B}[\alpha]\Gamma(\alpha)}\int_{a}^{t}f(\vartheta)(t-\vartheta)^{\alpha-1}d\vartheta.$$
(5)

Definition 3. The Laplace transform (LT) Associated to AB operator is defined as

$$L\begin{bmatrix} ABR\\ 0 \end{bmatrix} = \frac{\mathcal{B}[\alpha]}{1-\alpha} \frac{s^{\alpha} L[f(t)] - s^{\alpha-1} f(0)}{s^{\alpha} + (\alpha/(1-\alpha))}.$$
(6)

Theorem 1. The following Lipschitz conditions satisfy respectively for the Riemann-Liouville and AB derivatives [25]

$$\left\|_{a}^{ABC} D_{t}^{\alpha} f_{1}(t) - {}^{ABC}_{a} D_{t}^{\alpha} f_{2}(t)\right\| < K_{1} \|f_{1}(x) - f_{2}(x)\|,$$
(7)

and

$$\left\|_{a}^{ABR} D_{t}^{\alpha} f_{1}(t) - \frac{ABR}{a} D_{t}^{\alpha} f_{2}(t)\right\| < K_{2} \|f_{1}(x) - f_{2}(x)\|.$$
(8)

Theorem 2. The fractional differential equation ${}^{ABC}_{a}D^{\mu}_{t}f(t) = s(t)$ has a unique solution is given by [25]

$$f(t) = \frac{1-\alpha}{\mathcal{B}[\alpha]}s(t) + \frac{\alpha}{\mathcal{B}[\alpha]\Gamma(\alpha)}\int_{0}^{t}s(\varsigma)(t-\varsigma)^{\alpha-1}d\varsigma.$$
(9)

3 Basic idea of *q*-HATM

Here, we consider the differential equation of fractional order with respectively linear \mathcal{R} and nonlinear \mathcal{N} differential operator form

$${}^{ABC}_{a}D^{\alpha}_{t}v(x,t) + \mathcal{R}v(x,t) + \mathcal{N}v(x,t) = f(x,t), n-1 < \alpha \le n,$$
(10)

with the initial condition

$$v(x,0) = g(x), \tag{11}$$

where ${}^{ABC}_{a}D^{\alpha}_{t}v(x,t)$ symbolise the AB derivative of v(x,t), f(x,t) denotes the source term. Using LT, Eq. (10) gives

$$\mathcal{L}[\nu(x,t)] - \frac{g(x)}{s} + \frac{1}{B[\alpha]} \left(1 - \alpha + \frac{\alpha}{s^{\alpha}} \right) \{ \mathcal{L}[\mathcal{R}\nu(x,t)] + \mathcal{L}[\mathcal{N}\nu(x,t)] - \mathcal{L}[f(x,t)] \} = 0.$$
(12)

By the assist of HAM, \mathcal{N} is projected as

$$\mathcal{N}[\varphi(x,t;q)] = \mathcal{L}[\varphi(x,t;q)] - \frac{\varphi(x)}{s}$$

$$+ \frac{1}{\mathcal{B}[\alpha]} \left(1 - \alpha + \frac{\alpha}{s^{\alpha}}\right) \{\mathcal{L}[\mathcal{R} \ \varphi(x,t;q)] + L[\mathcal{N}\varphi(x,t;q)] - L[f(x,t)]\}.$$
(13)

Here, $\varphi(x,t;q)$ is the real-valued function. Now, we have

$$(1 - nq)\mathcal{L}[\varphi(x, t; q) - v_0(x, t)] = \hbar q \mathcal{N}[\varphi(x, t; q)],$$
(14)

where *L* is signifying *LT*, $q \in [0, \frac{1}{n}]$ $(n \ge 1)$ is the embedding parameter and $\hbar \ne 0$ is an auxiliary parameter. For q = 0 and $q = \frac{1}{n}$, we have

$$\varphi(x,t;0) = v_0(x,t), \varphi\left(x,t;\frac{1}{n}\right) = v(x,t).$$
(15)

Thus, by intensifying q from 0 to $\frac{1}{n}$, then $\varphi(x,t;q)$ changes from $v_0(x,t)$ to v(x,t). Using Taylor theorem near to q, we defining $\varphi(x,t;q)$ in series form and then we get

$$\varphi(x,t;q) = v_0(x,t) + \sum_{m=1}^{\infty} v_m(x,t)q^m,$$
(16)

where

$$v_m(x,t) = \frac{1}{m!} \frac{\partial^m \varphi(x,t;q)}{\partial q^m} \Big|_{q=0}.$$
(17)

For the proper chaise of $v_0(x,t)$, n and \hbar , the series (14) converges at $q = \frac{1}{n}$. By simplifying Eq. (14), we achieved

$$\mathcal{L}[v_m(x,t) - \mathbf{k}_m v_{m-1}(x,t)] = \hbar \Re_m(\vec{v}_{m-1}), \tag{18}$$

where the vectors are defined as

$$\vec{v}_m = \{v_0(x,t), v_1(x,t), \dots, v_m(x,t)\}.$$
(19)

On employing inverse LT on Eq. (18), we get

$$v_m(x,t) = k_m v_{m-1}(x,t) + \hbar \mathcal{L}^{-1}[\Re_m(\vec{v}_{m-1})], \qquad (20)$$

where

$$\Re_{m}(\vec{v}_{m-1}) = L[v_{m-1}(x,t)] - \left(1 - \frac{k_{m}}{n}\right) \left(\frac{\mathcal{G}(x)}{s} + \frac{1}{\mathcal{B}[\alpha]} \left(1 - \alpha + \frac{\alpha}{s^{\alpha}}\right) L[f(x,t)]\right) + \frac{1}{\mathcal{B}[\alpha]} \left(1 - \alpha + \frac{\alpha}{s^{\alpha}}\right) L[Rv_{m-1} + \mathcal{H}_{m-1}],$$
(21)

and

$$\mathbf{k}_m = \begin{cases} 0, m \le 1, \\ n, m > 1. \end{cases}$$
(22)

In Eq. (21), \mathcal{H}_m signifies homotopy polynomial and which is defined as

$$\mathcal{H}_m = \frac{1}{m!} \left[\frac{\partial^m \varphi(x,t;q)}{\partial q^m} \right]_{q=0} \text{ and } \varphi(x,t;q) = \varphi_0 + q\varphi_1 + q^2\varphi_2 + \dots$$
(23)

By the aid of Eqs. (20) and (21), one can get

$$v_m(x,t) = (k_m + \hbar)v_{m-1}(x,t) - \left(1 - \frac{k_m}{n}\right)\mathcal{L}^{-1}\left(\frac{\mathcal{G}(x)}{s} + \frac{1}{\mathcal{B}[\alpha]}\left(1 - \alpha + \frac{\alpha}{s^{\alpha}}\right)L[f(x,t)]\right) + \hbar\mathcal{L}^{-1}\left\{\frac{1}{\mathcal{B}[\alpha]}\left(1 - \alpha + \frac{\alpha}{s^{\alpha}}\right)L[Rv_{m-1} + \mathcal{H}_{m-1}]\right\}.$$
(24)

The q-HATM solution is presented as

$$v(x,t) = v_0(x,t) + \sum_{m=1}^{\infty} v_m(x,t) \left(\frac{1}{n}\right)^m.$$
 (25)

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4 Solution for Proposed Model

To demonstrate the efficiency and solution procedure of the projected method, in here we consider system describing considered a model with arbitrary order. By the assist of Eq. (2), one can get

$${}^{ABC}_{a}D^{\alpha}_{t}S(t) + \delta\varepsilon S(t) - \beta C(t) = 0,$$

$${}^{ABC}_{a}D^{\alpha}_{t}C(t) - \delta\varepsilon S(t) + (\beta + \lambda)C(t) = 0, \qquad 0 < \alpha \le 1,$$

$${}^{ABC}_{a}D^{\alpha}_{t}A(t) - \lambda C(t) = 0$$
(26)

with initial conditions

$$S(0) = S_0(t), C(0) = C_0(t), A(0) = A_0(t).$$
(27)

Taking LT on Eq. (26) and then using the Eq. (27), we get

$$L[S(t)] = \frac{1}{s} (S_0(t)) + \frac{1}{\mathcal{B}[\alpha]} \left(1 - \alpha + \frac{\alpha}{s^{\alpha}} \right) L\{\delta \varepsilon S(t) - \beta C(t)\},$$

$$L[C(t)] = \frac{1}{s} (C_0(t)) - \frac{1}{\mathcal{B}[\alpha]} \left(1 - \alpha + \frac{\alpha}{s^{\alpha}} \right) L\{\delta \varepsilon S(t) - (\beta + \lambda)C(t)\},$$

$$L[A(t)] = \frac{1}{s} (A_0(t)) - \frac{1}{\mathcal{B}[\alpha]} \left(1 - \alpha + \frac{\alpha}{s^{\alpha}} \right) L\{\lambda C(t)\}.$$
(28)

Now, we define N as below

$$N^{1}[\varphi_{1}(t;q),\varphi_{2}(t;q),\varphi_{3}(t;q)] = L[\varphi_{1}(t;q)] - \frac{1}{s}(S_{0}(t)) + \frac{1}{\mathcal{B}[\alpha]} \left(1 - \alpha + \frac{\alpha}{s^{\alpha}}\right) L\{\delta\varepsilon\varphi_{1}(t;q) + \beta\varphi_{2}(t;q)\}, N^{2}[\varphi_{1}(t;q),\varphi_{2}(t;q),\varphi_{3}(t;q)] = L[\varphi_{2}(t;q)] - \frac{1}{s}(C_{0}(t)) - \frac{1}{\mathcal{B}[\alpha]} \left(1 - \alpha + \frac{\alpha}{s^{\alpha}}\right) L\{\delta\varepsilon\varphi_{1}(t;q) + (\beta + \lambda)\varphi_{2}(t;q)\}, N^{3}[\varphi_{1}(t;q),\varphi_{2}(t;q),\varphi_{3}(t;q)] = L[\varphi_{3}(t;q)] - \frac{1}{s}(A_{0}(t)) - \frac{1}{\mathcal{B}[\alpha]} \left(1 - \alpha + \frac{\alpha}{s^{\alpha}}\right) L\{\lambda\varphi_{2}(t;q)\}.$$
(29)

The deformation equation of *m*-th order at $\mathcal{H}(x, t) = 1$ is defined as

$$L[S_{m}(t) - \mathbf{k}_{m}S_{m-1}(t)] = \hbar \Re_{1,m} \Big[\vec{S}_{m-1}, \vec{C}_{m-1}, \vec{A}_{m-1} \Big],$$

$$L[C_{m}(t) - \mathbf{k}_{m}C_{m-1}(t)] = \hbar \Re_{2,m} \Big[\vec{S}_{m-1}, \vec{C}_{m-1}, \vec{A}_{m-1} \Big],$$

$$L[A_{m}(t) - \mathbf{k}_{m}A_{m-1}(t)] = \hbar \Re_{3,m} \Big[\vec{S}_{m-1}, \vec{C}_{m-1}, \vec{A}_{m-1} \Big],$$
(30)

where

$$\begin{aligned} \Re_{1,m} \Big[\vec{S}_{m-1}, \vec{C}_{m-1}, \vec{A}_{m-1} \Big] &= L[S_{m-1}(t)] - \left(1 - \frac{k_m}{n}\right) \Big\{ \frac{1}{s} \left(S_0(t)\right) \Big\} \\ &+ \frac{1}{\mathcal{B}[\alpha]} \Big(1 - \alpha + \frac{\alpha}{s^{\alpha}} \Big) L\{\delta \varepsilon S_{m-1}(t) - \beta C_{m-1}(t)\}, \\ \Re_{2,m} \Big[\vec{S}_{m-1}, \vec{C}_{m-1}, \vec{A}_{m-1} \Big] &= L[C_{m-1}(t)] + \left(1 - \frac{k_m}{n}\right) \Big\{ \frac{1}{s} \left(C_0(t)\right) \Big\} \\ &+ \frac{1}{\mathcal{B}[\alpha]} \Big(1 - \alpha + \frac{\alpha}{s^{\alpha}} \Big) L\{\delta \varepsilon S_{m-1}(t) - (\beta + \lambda) C_{m-1}(t)\}, \\ \Re_{3,m} \Big[\vec{S}_{m-1}, \vec{C}_{m-1}, \vec{A}_{m-1} \Big] &= L[A_{m-1}(t)] + \left(1 - \frac{k_m}{n}\right) \Big\{ \frac{1}{s} \left(A_0(t)\right) \Big\} \\ &- \frac{1}{\mathcal{B}[\alpha]} \Big(1 - \alpha + \frac{\alpha}{s^{\alpha}} \Big) L\{\lambda C_{m-1}(t)\}. \end{aligned}$$
(31)

Eq. (31) reduces after employing inverse LT, as follows

$$S_{m}(t) = \mathbf{k}_{m}S_{m-1}(t) + \hbar L^{-1} \left\{ \Re_{1,m} \left[\vec{S}_{m-1}, \vec{C}_{m-1}, \vec{A}_{m-1} \right] \right\},\$$

$$C_{m}(t) = \mathbf{k}_{m}C_{m-1}(t) + \hbar L^{-1} \left\{ \Re_{2,m} \left[\vec{S}_{m-1}, \vec{C}_{m-1}, \vec{A}_{m-1} \right] \right\},\$$

$$A_{m}(t) = \mathbf{k}_{m}A_{m-1}(t) + \hbar L^{-1} \left\{ \Re_{3,m} \left[\vec{S}_{m-1}, \vec{C}_{m-1}, \vec{A}_{m-1} \right] \right\}.$$
(32)

Using $S_0(t) = R_T$, $C_0 = 0$ and $A_0(t) = 0$ we can obtain the terms of the series solution with the help of the above system

$$S(t) = S_0(t) + \sum_{m=1}^{\infty} S_m(t) \left(\frac{1}{n}\right)^m,$$

$$C(t) = C_0(t) + \sum_{m=1}^{\infty} C_m(t) \left(\frac{1}{n}\right)^m,$$

$$A(t) = A_0(t) + \sum_{m=1}^{\infty} A_m(t) \left(\frac{1}{n}\right)^m,$$

(33)

5 Existence of Solutions for the Proposed Problem

Here, to present the existence of the solution, we considered the fixed-point theorem. Now, the system (27) is considered as

$$\begin{cases} {}^{ABC}_{0}D^{x}_{t}[S(t)] = \mathcal{G}_{1}(t,S), \\ {}^{ABC}_{0}D^{x}_{t}[C(t)] = \mathcal{G}_{2}(t,C), \\ {}^{ABC}_{0}D^{x}_{t}[A(t)] = \mathcal{G}_{3}(t,A). \end{cases}$$
(34)

Using the Theorem 2, Eq. (35) is transformed to the Volterra integral equation and defined as

$$\begin{cases} S(t) - S(0) = \frac{(1-\alpha)}{B(\alpha)} \mathcal{G}_1(t, S) + \frac{\alpha}{B(\alpha)\Gamma(\alpha)} \int_0^t \mathcal{G}_1(\zeta, S)(t-\zeta)^{\alpha-1} d\zeta, \\ C(t) - C(0) = \frac{(1-\alpha)}{B(\alpha)} \mathcal{G}_2(t, C) + \frac{\alpha}{B(\alpha)\Gamma(\alpha)} \int_0^t \mathcal{G}_2(\zeta, C)(t-\zeta)^{\alpha-1} d\zeta, \\ A(t) - A(0) = \frac{(1-\alpha)}{B(\alpha)} \mathcal{G}_3(t, A) + \frac{\alpha}{B(\alpha)\Gamma(\alpha)} \int_0^t \mathcal{G}_3(\zeta, A)(t-\zeta)^{\alpha-1} d\zeta. \end{cases}$$
(35)

Theorem 3. The kernel G_1 satisfies the Lipschitz condition and contraction if $0 \le (\delta \varepsilon - \beta \lambda_2) < 1$ holds.

Proof. We consider the two functions u and u_1 to prove the required result, as follows

$$\begin{aligned} \|\mathcal{G}_{1}(t,S) - \mathcal{G}_{1}(t,S_{1})\| &= \|(\delta c[S(t) - S(t_{1})] - \beta C(t))\| \\ &\leq \|\delta \varepsilon - \beta C(t)\| \|S(t) - S(t_{1})\| \\ &\leq (\delta \varepsilon - \beta \lambda_{2})\|S(t) - S(t_{1})\|, \end{aligned} (36)$$

where $|| C(t) || \le \lambda_2$ be the bounded function. Putting $\eta_1 = \delta \varepsilon - \beta \lambda_2$ in Eq. (37), we have

$$\|\mathcal{G}_1(t,S) - \mathcal{G}_1(t,S_1)\| \le \eta_1 \|S(t) - S(t_1)\|.$$
(37)

Equation (38) signifies Lipschitz condition for G_1 . If $0 \le (\delta \varepsilon - \beta \lambda_2) < 1$, then it gives the contraction. Similarly, we have

$$\begin{cases} \|\mathcal{G}_{2}(t,C) - \mathcal{G}_{2}(t,C_{1})\| \le \eta_{2} \|C(t) - C(t_{1})\|, \\ \|\mathcal{G}_{3}(t,A) - \mathcal{G}_{3}(t,A_{1})\| \le \eta_{3} \|A(t) - A(t_{1})\|. \end{cases}$$
(38)

Now, we define the recursive form of Eq. (36) as with initial conditions

$$\begin{cases} S_n(t) = \frac{(1-\alpha)}{\mathcal{B}(\alpha)} \mathcal{G}_1(t, S_{n-1}) + \frac{\alpha}{\mathcal{B}(\alpha)\Gamma(\alpha)} \int_0^t \mathcal{G}_1(\zeta, S_{n-1})(t-\zeta)^{\alpha-1} d\zeta, \\ C_n(t) = \frac{(1-\alpha)}{\mathcal{B}(\alpha)} \mathcal{G}_2(t, C_{n-1}) + \frac{\alpha}{\mathcal{B}(\alpha)\Gamma(\alpha)} \int_0^t \mathcal{G}_2(\zeta, C_{n-1})(t-\zeta)^{\alpha-1} d\zeta, \\ A_n(t) = \frac{(1-\alpha)}{\mathcal{B}(\alpha)} \mathcal{G}_3(t, A_{n-1}) + \frac{\alpha}{\mathcal{B}(\alpha)\Gamma(\alpha)} \int_0^t \mathcal{G}_3(\zeta, A_{n-1})(t-\zeta)^{\alpha-1} d\zeta, \end{cases}$$
(39)

and

$$S(0) = S_0(t), C(0) = C_0(t)$$
 and $A(0) = A_0(t).$ (40)

The successive difference between the terms presented as

$$\begin{aligned}
\begin{pmatrix}
\phi_{1n}(t) &= S_n(t) - S_{n-1}(t) \\
&= \frac{(1-\alpha)}{\mathcal{B}(\alpha)} (\mathcal{G}_1(t, S_{n-1}) - \mathcal{G}_1(t, S_{n-2})) + \frac{\alpha}{\mathcal{B}(\alpha)\Gamma(\alpha)} \int_0^t \mathcal{G}_1(\zeta, S_{n-1})(t-\zeta)^{\alpha-1} d\zeta, \\
\phi_{2n}(t) &= C_n(t) - C_{n-1}(t) \\
&= \frac{(1-\alpha)}{\mathcal{B}(\alpha)} (\mathcal{G}_2(t, C_{n-1}) - \mathcal{G}_2(t, C_{n-2})) + \frac{\alpha}{\mathcal{B}(\alpha)\Gamma(\alpha)} \int_0^t \mathcal{G}_2(\zeta, C_{n-1})(t-\zeta)^{\alpha-1} d\zeta, \\
\phi_{3n}(t) &= A_n(t) - A_{n-1}(t) \\
&= \frac{(1-\alpha)}{\mathcal{B}(\alpha)} (\mathcal{G}_3(t, A_{n-1}) - \mathcal{G}_3(t, A_{n-2})) + \frac{\alpha}{\mathcal{B}(\alpha)\Gamma(\alpha)} \int_0^t \mathcal{G}_3(\zeta, A_{n-1})(t-\zeta)^{\alpha-1} d\zeta.
\end{aligned}$$
(41)

Notice that

$$\begin{cases} S_n(t) = \sum_{i=1}^n \phi_{1i}(t), \\ C_n(t) = \sum_{i=1}^n \phi_{2i}(t), \\ A_n(t) = \sum_{i=1}^n \phi_{3i}(t). \end{cases}$$
(42)

By using Eq. (39) and applying the norm on the first term of Eq. (42), we have

$$\|\phi_{1n}(t)\| \leq \frac{(1-\alpha)}{\mathcal{B}(\alpha)}\eta_1 \left\|\phi_{1(n-1)}(t)\right\| + \frac{\alpha}{\mathcal{B}(\alpha)\Gamma(\alpha)}\eta_1 \int_0^t \left\|\phi_{1(n-1)}(\zeta)\right\| d\zeta.$$
(43)

Similarly, we have

$$\begin{cases} \left\|\phi_{2n}(t)\right\| \leq \frac{(1-\alpha)}{\mathcal{B}(\alpha)}\eta_{2}\left\|\phi_{2(n-1)}(t)\right\| + \frac{\alpha}{\mathcal{B}(\alpha)\Gamma(\alpha)}\eta_{2}\int_{0}^{t}\left\|\phi_{2(n-1)}(\zeta)\right\|d\zeta,\\ \left\|\phi_{3n}(t)\right\| \leq \frac{(1-\alpha)}{\mathcal{B}(\alpha)}\eta_{3}\left\|\phi_{3(n-1)}(t)\right\| + \frac{\alpha}{\mathcal{B}(\alpha)\Gamma(\alpha)}\eta_{3}\int_{0}^{t}\left\|\phi_{3(n-1)}(\zeta)\right\|d\zeta. \end{cases}$$
(44)

Next using the above result, we have following results.

Theorem 4. The solution for Eq. (27) will exist and unique if we have t_0 then

$$\frac{(1-\alpha)}{\mathcal{B}(\alpha)}\eta_i + \frac{\alpha}{\mathcal{B}(\alpha)\Gamma(\alpha)}\eta_i < 1,$$

for i = 1, 2 and 3.

Proof. Let S(t), C(t) and A(t) be the bounded functions admitting the Lipschitz condition. Now, we have by Eqs. (43) and (45)

$$\begin{aligned} \|\phi_{1i}(t)\| &\leq \|S_n(0)\| \left[\frac{(1-\alpha)}{\mathcal{B}(\alpha)} \eta_1 + \frac{\alpha}{\mathcal{B}(\alpha)\Gamma(\alpha)} \eta_1 \right]^n, \\ \|\phi_{2i}(t)\| &\leq \|C_n(0)\| \left[\frac{(1-\alpha)}{\mathcal{B}(\alpha)} \eta_2 + \frac{\alpha}{\mathcal{B}(\alpha)\Gamma(\alpha)} \eta_2 \right]^n, \\ \|\phi_{3i}(t)\| &\leq \|A_n(0)\| \left[\frac{(1-\alpha)}{\mathcal{B}(\alpha)} \eta_3 + \frac{\alpha}{\mathcal{B}(\alpha)\Gamma(\alpha)} \eta_3 \right]^n. \end{aligned}$$
(45)

This proves the continuity as well as existence. Now, we consider showing Eq. (46) is a solution for the system (27)

$$S(t) - S(0) = S_n(t) - \mathcal{K}_{1n}(t),$$

$$C(t) - C(0) = C_n(t) - \mathcal{K}_{2n}(t),$$

$$A(t) - A(0) = A_n(t) - \mathcal{K}_{3n}(t).$$
(46)

To achieve the required result, we consider

$$\begin{aligned} \|\mathcal{K}_{1n}(t)\| &= \left\| \frac{(1-\alpha)}{\mathcal{B}(\alpha)} (\mathcal{G}_{1}(t,S) - \mathcal{G}_{1}(t,S_{n-1})) + \frac{\alpha}{\mathcal{B}(\alpha)\Gamma(\alpha)} \int_{0}^{t} (t-\zeta)^{\mu-1} (\mathcal{G}_{1}(\zeta,S) - \mathcal{G}_{1}(\zeta,S_{n-1})) d\zeta \right\| \\ &\leq \frac{(1-\alpha)}{\mathcal{B}(\alpha)} \| (\mathcal{G}_{1}(t,S) - \mathcal{G}_{1}(t,S_{n-1})) \| \\ &+ \frac{\alpha}{\mathcal{B}(\alpha)\Gamma(\alpha)} \int_{0}^{t} \| (\mathcal{G}_{1}(\zeta,S) - \mathcal{G}_{1}(\zeta,S_{n-1})) \| d\zeta \\ &\leq \frac{(1-\alpha)}{\mathcal{B}(\alpha)} \eta_{1} \| S - S_{n-1} \| + \frac{\alpha}{\mathcal{B}(\alpha)\Gamma(\alpha)} \eta_{1} \| S - S_{n-1} \| t. \end{aligned}$$

In the same way at t_0 , we can obtain

$$\|\mathcal{K}_{1n}(t)\| \le \left(\frac{(1-\alpha)}{\mathcal{B}(\alpha)} + \frac{\alpha t_0}{\mathcal{B}(\alpha)\Gamma(\alpha)}\right)^{n+1} \eta_1^{n+1} M.$$
(48)

We can see that form Eq. (49), when *n* approaches to ∞ , $\|\mathcal{K}_{1n}(t)\|$ tends to 0. We can verify similarly for $\|\mathcal{K}_{2n}(t)\|$ and $\|\mathcal{K}_{3n}(t)\|$.

Now, we present the uniqueness. Suppose $S^*(t)$, $C^*(t)$ and $A^*(t)$ be the set of other solutions, then we have

$$S(t) - S^{*}(t) = \frac{(1-\alpha)}{\mathcal{B}(\alpha)} (\mathcal{G}_{1}(t,S) - \mathcal{G}_{1}(t,S^{*})) + \frac{\alpha}{\mathcal{B}(\alpha)\Gamma(\alpha)} \int_{0}^{t} (\mathcal{G}_{1}(\zeta,S) - \mathcal{G}_{1}(\zeta,S^{*})) d\zeta.$$

$$(49)$$

By employing norm on Eq. (51), we get

$$\begin{split} \|S(t) - S^*(t)\| &= \left\| \frac{(1-\alpha)}{\mathcal{B}(\alpha)} (\mathcal{G}_1(t,S) - \mathcal{G}_1(t,S^*)) + \frac{\alpha}{\mathcal{B}(\alpha)\Gamma(\alpha)} \int_0^t (\mathcal{G}_1(\zeta,S) - \mathcal{G}_1(\zeta,S^*)) d\zeta \right\| \\ &\leq \frac{(1-\alpha)}{\mathcal{B}(\alpha)} \eta_1 \|S(t) - S^*(t)\| + \frac{\alpha}{\mathcal{B}(\alpha)\Gamma(\alpha)} \eta_1 t \|S(t) - S^*(t)\|. \end{split}$$

$$(50)$$

On simplification

$$\|S(t) - S^*(t)\| \left(1 - \frac{(1-\alpha)}{\mathcal{B}(\alpha)}\eta_1 - \frac{\alpha}{\mathcal{B}(\alpha)\Gamma(\alpha)}\eta_1 t\right) \le 0.$$
(51)

From the above condition, it is clear that $S(t) = S^*(t)$, if

$$\left(1 - \frac{(1-\alpha)}{\mathcal{B}(\alpha)}\eta_1 - \frac{\alpha}{\mathcal{B}(\alpha)\Gamma(\alpha)}\eta_1 t\right) \ge 0.$$
(52)

Therefore Eq. (52) proves our result.

6 Numerical Results and Discussion

Here, we illustrated the nature of q-HATM solution for different α . The initial conditions for the proposed model is defined as

$$S(0) = S_0(t) = N, C(0) = C_0(t) = 0, A(0) = A_0(t) = 0.$$

where N is the total number of receptors and which is 4.4×10^4 No./cell. In order to capture the behaviour, the value of the parameters cited in Eq. (2) are considered as follows

$$\delta = 0.0005 M^{-1} s^{-1}, \beta = 142.8 s^{-1}, \varepsilon = 1 unit/mL, \lambda = 0.12 s^{-1}.$$

The nature of results obtained by *q*-HATM for a considered model with different α is dissipated in Fig. 1 with different fractional order. To analyse the behaviour of archived results associated with \hbar , the \hbar -curves are plotted for distinct α is captured in Fig. 2. These help us to adjust and control the convergence region of the obtained results. For a suitable \hbar , the obtained results rapidly tend to an analytical solution. Moreover, in the plots the convergence region is denoted by the horizontal line. The captured figures show the degree of freedom and more simulating consequences about the hired model with different arbitrary order and also it signifies the novelty of the fractional operator employed. Further, from all plots one can observer that the projected solution procedure is and very effective and more accurate to examine the considered nonlinear problem.



Fig. 1. Behaviour of the obtained results for (a) S(t), (b) C(t) and (c) A(t) with different α at n = 1 and $\hbar = -1$.



Fig. 2. *h*-curves for (a) S(t), (b) C(t) and (c) A(t) with distinct α at t = 0.01 and n = 1.

7 Conclusion

The *q*-HATM is employed efficiently in the present framework to find the solution for the system of equation with arbitrary order and illustrating the model of TRA mechanism in calcium signalling. Since, generalized Mittag-Leffler function is hired to define fractional-order AB integrals and derivatives, these operators help us to capture more simulating consequences and also it incorporate most essential behaviours of the models, and hence the current study exemplifies the effeteness of the projected derivative. Further, for the obtained results we presented the existence and uniqueness within the frame of fixed point hypothesis. As associated to consequences available in the literature, the results obtained by the help of projected method are more stimulating. The graphical representations show the dependence of the considered nonlinear model on parameters offered by the considered scheme and fractional order, and also it exemplifies the degree of freedom when we incorporate the fractional operator in the systems. We can be observed by the present study, the projected model is remarkably associated with the time instant and time history-based consequences, and which can be efficiently examined by the help of fractional calculus. Lastly, we can conclude that the present study can aid the researchers to analyse the nature system corresponded to very useful and interesting and consequences.

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Stability Analysis of Bifurcated Limit Cycles in a Labor Force Evolution Model

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Abstract. This article focuses on the fluctuations observed in the labor markets. We divide the total population into three categories: employed, unemployed and inactive, then we describe the entry-exit flows between these different categories by two delay differential equations. Our contribution is to compute an indicator for determining the behavior of the model variables, in a neighborhood of the critical delay. Our findings show that the model can undergo a Hopf bifurcation and the bifurcated limit cycles is stable (or unstable), according to the crossing direction of critical delay.

Keywords: Differential equations \cdot Delay \cdot Periodic solutions \cdot Hopf bifurcation \cdot Limit cycles \cdot Labor market \cdot Employed persons \cdot Unemployed

1 Introduction

The study of economic fluctuations and cycles has attracted much attention in macroeconomic theory and applied mathematics. On one hand, the resulting mathematical models have been involved solving many problems related to the explanation, identification and measurement or estimation of these phenomena, and on the other hand, the obtained results align well with proposed results in econometric studies [6,7,10,12,14]. Within this framework, we developed a mathematical model to study the fluctuations observed empirically in the three aggregates of labor markets namely, employment, unemployment and inactivity. The idea is to divide the population into three distinct categories of individuals: the employed, the unemployed and the inactive and describe the flows between these different categories by the following system (For more details, see Fig. 1):

$$\begin{cases} \frac{dL}{dt} = \gamma U(t) - (s + \alpha + m)L(t),\\ \frac{dU}{dt} = \rho \left(1 - \frac{L(t-r) + U(t-r)}{N_c}\right)L(t-r) + sL(t) - (m+\gamma+\beta)U(t), \end{cases}$$
(1)

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Fig. 1. Labor market flows

with the initial condition:

$$(L(\xi), U(\xi)) = (\varphi_1(\xi), \varphi_2(\xi)), \quad \forall \xi \in [-r, 0],$$
 (2)

where the variable L is the number of the employed population, U is the number of the unemployed population, γ denotes the employment level, s indicates the job loss rate, α is the rate of workers who have withdrawn from labor market due to retirement or disability, β is the rate of unemployed people who are no longer able to work, m is the mortality rate, ρ is the maximum population growth rate, N_c is the maximum load capacity, r is the time needed for a new person who has found a job to contribute to the reproductive process and $\varphi_i \in \mathcal{C}([-r, 0], \mathcal{R}^+)$, i = 1, 2. Here $\mathcal{C}([-r, 0], \mathbb{R}^+)$ is the Banach space of continuous functions from the interval [-r, 0] to the set of positive real numbers \mathbb{R}^+ .

Model (1) is composed of two differential equations which model the inflowsoutflows in the three categories of the total population (category U of the unemployed, category L of the employed and category I of inactive people). On the one hand, the second equation translates the feeding of category U by people who have reached working age and who are looking for a job. This flow is indicated by a logistic growth rate ρ . People in this category transform to category I with a disability rate β , or category L with a recruitment rate γ or exit with a death rate m. On the other hand, the first equation describes the evolution of category L by the difference between the newly employed, noted γU and those leaving this category by job loss, sL, by disability αL or by death, mL.

There have been several attempts in this area. We cite for example the model of labor force evolution, proposed by Farkas (in 1995 [16]), Only a handful of studies have been found to examine its dynamics systematically. Papers [9,15,20] study the local and global stability of labor force evolution model using linearization technique and Lyapunov method. The resulting numerical results usually give the local and global stability. In empirical studies, however, the observed data of the employed and the unemployed have oscillatory behavior. In this work, we prove the existence of a Hopf bifurcation point and we also study the direction and stability of the periodic branches (limit cycles) that evolve around this point.

The study of the oscillatory behavior of dynamical systems can be done by fixed point methods [8,17] or by Hopf bifurcation theorem [13,25]. In the latter case, several researchers have proposed different techniques to investigate the behavior of dynamical systems in the neighborhoods of the critical delays. The first is to look for a normal form in a central manifold [5,22,23]. This method entails a long calculation. The second is the singular perturbation approach. This technique is preferable for its computational efficiency (multiple-scale analysis [4,21], Krylov-Bogoliubov-Mitropolsky method [2], Poincaré-Lindstedt method [1,19], harmonic balance method [3,18] and pseudo-oscillator analysis [27]). In this work we chose to work with the Kuznetsov method [11,26]. The method only requires a computation of the first Lyapunov coefficient to determine the behavior in the neighborhoods of the critical delays.

This work is structured as follows. In Sect. 2, we study the local stability and the Hopf bifurcation of the nontrivial equilibrium position of the System (1). In Sect. 3, we first give the essential calculations of the central manifold and the reduction of our model to a normal form. Then, we use the Kuzentsov method to determine the direction and stability of the periodic orbit resulting from the Hopf bifurcation. Numerical simulations are given in Sect. 4 to support the main aspects of our study. Finally, in Sect. 5, we summarize the main findings, our conclusion, the gaps we encountered and some perspectives on this study.

2 Hopf Bifurcation Analysis

2.1 Equilibria

In the following, we study the existence of equilibrium points for the system (1).

Proposition 1. If $\frac{\gamma\rho}{m(\gamma + s + \alpha + \alpha_2 + m) + \gamma\alpha + \beta(s + \alpha)} > 1$, then system (1) admits two equilibria: $E_0 = (0,0)$ and a unique positive equilibrium $E_* = (L_*, U_*)$, where

$$U_* = \frac{(s + \alpha + m)L_*}{\gamma}$$

and

$$L_* = \frac{\gamma N_c}{\gamma + s + \alpha + m} \left(1 - \frac{m(\gamma + s + \alpha + \beta + m) + \gamma \alpha + \beta(s + \alpha)}{\gamma \rho}\right).$$

Proof. Suppose (U, L) is an equilibrium point, that is,

$$\begin{cases} \gamma U - (s + \alpha + m)L = 0, \\ \rho (1 - \frac{L+U}{N_c})L + sL - (\gamma + \beta + m)U = 0. \end{cases}$$
(3)

It is clear that (0,0) is a solution of the system (3). This gives that E_0 is a trivial equilibrium of system (1). Moreover, if (U, L) is an equilibrium such that U > 0 and L > 0 then we have

$$\begin{cases} U = \frac{(s+\alpha+m)L}{\gamma},\\ \rho(1 - \frac{(\gamma+s+\alpha+m)L}{\gamma N_c}) - \frac{(s+\alpha+m)(\gamma+\beta+m)}{\gamma} = 0. \end{cases}$$
(4)

Under the condition $\frac{\gamma\rho}{m(\gamma+s+\alpha+\beta+m)+\gamma\alpha+\beta(s+\alpha)} > 1$, the system of linear equations (4) has a unique non-trivial solution $E_* = (L_*, U_*)$, where

$$U_* = \frac{(s + \alpha + m)L_*}{\gamma}$$

and

$$L_* = \frac{\gamma N_c}{(\gamma + s + \alpha + m)} \left(1 - \frac{m(\gamma + s + \alpha + \beta + m) + \gamma \alpha + \beta(s + \alpha)}{\gamma \rho}\right).$$

This completes the proof.

2.2 Local Stability

The Linearized system of Eqs. (1) at the positive equilibrium E_* is

$$\begin{cases} \frac{dx}{dt} = -(s+\alpha+m)x + \gamma y,\\ \frac{dy}{dt} = sx + \left(\frac{m(\gamma+s+\alpha+\beta+m)+\gamma\alpha+\beta(s+\alpha)}{\gamma} - \frac{\rho L_*}{N_c}\right)x_r - (\gamma+\beta+m)y - \frac{\rho L_*}{N_c}y_r. \end{cases}$$
(5)

For System (5) the characteristic equation is:

$$\lambda^2 + \theta_1 \lambda + \theta_2 \lambda e^{-\lambda r} + \theta_3 + \theta_4 e^{-\lambda r} = 0, \tag{6}$$

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where

$$\theta_1 = \gamma + s + \alpha + \beta + 2m, \quad \theta_2 = \frac{\rho L_*}{N_c}$$
$$\theta_3 = m(\gamma + s + \alpha + \beta + m) + \gamma \alpha + \beta(s + \alpha)$$

and

$$\theta_4 = \rho(\gamma + s + \alpha + m) \frac{L_*}{N_c} - (m(\gamma + s + \alpha + \beta + m) + \gamma\alpha + \beta(s + \alpha)).$$

Using the Routh-Hurwitz criterion and Kuang's results [25] for Eq. (6), we prove the following results:

Proposition 2. If r = 0, then the positive equilibrium E_* is locally asymptotically stable.

Proof. When r = 0, the Eq. (6) becomes

$$\lambda^2 + (\gamma + s + \alpha + \beta + 2m + \frac{\rho L_*}{N_c})\lambda + \rho(\gamma + s + \alpha + m)\frac{L_*}{N_c} = 0, \qquad (7)$$

Since $\gamma + s + \alpha + \beta + 2m + \frac{\rho L_*}{N_c} > 0$ and $\rho(\gamma + s + \alpha + m) \frac{L_*}{N_c} > 0$, then, by the Routh-Hurwitz criterion, all the roots of Eq. (6) have non-negative real parts, and therefore the positive equilibrium E_* is locally asymptotically stable.

Let (H_1) : $\frac{\gamma \rho}{m(\gamma + s + \alpha + \beta + m)} > 3.$

Proposition 3. If (H_1) is valid. Then there exists $r_0 > 0$ such that,

(i) for $0 \le r < r_0$, E_* is locally asymptotically stable; (ii) for $r > r_0$, E_* is unstable; (iii) for $r = r_0$, Eq. (5) admits two purely imaginary roots;

where

$$r_0 = \frac{1}{\omega_0} \arccos(\frac{-\theta_1 \theta_2 \omega_0^2 - \theta_4 (\theta_3 - \omega_0^2)}{\theta_2 \omega_0^2 + \theta_4^2}),$$

and

$$\omega_0 = \sqrt{\frac{1}{2} \{ (\theta_2^2 + 2\theta_3 - \theta_1^2) + \sqrt{(\theta_2^2 + 2\theta_3 - \theta_1^2)^2 + 4(\theta_3^2 - \theta_4^2)} \}}$$

Proof. If $\frac{\gamma\rho}{m(\gamma+s+\alpha+\beta+m)} \ge 3$, then

$$\theta_3^2 - \theta_4^2 = \Lambda_1 \times \Lambda_2 \times \Lambda_3 < 0,$$

where

$$\Lambda_1 = [m(\gamma + s + \alpha + \beta + m) + \gamma \alpha + \beta(s + \alpha)]^2,$$

$$\Lambda_2 = \frac{\gamma \rho}{m(\gamma + s + \alpha + \beta + m) + \gamma \alpha + \beta(s + \alpha)} - 1$$

and

$$\Lambda_3 = 3 - \frac{\gamma \rho}{m(\gamma + s + \alpha + \beta + m) + \gamma \alpha + \beta(s + \alpha)}.$$

Consequently, Eq. (6) has only one purely imaginary solution,

$$i\omega_0 = i\sqrt{\frac{1}{2}\{(\theta_2^2 + 2\theta_3 - \theta_1^2) + \sqrt{(\theta_2^2 + 2\theta_1^2)^2 + 4(\theta_3^2 - \theta_4^2)}\}},$$

with $\omega_0 > 0$. By Theorem 2.7 in ([25], p. 77), we conclude that there exists $r_0 > 0$ which satisfies the three statements (i), (ii) and (iii) of Proposition 3.
2.3 Local Hopf Bifurcation

From (*iii*) of Proposition 3, we have proved that Eq. (5) has a pair of purely imaginary roots $\pm i\omega_0$, $\omega_0 > 0$, when the delay crosses the critical value r_0 . In the following result we show the birth of the Hopf bifurcation, in a small vicinity of r_0 .

Theorem 1. Under hypothesis (H_1) , the system (1) loses its stability through a Hopf bifurcation when $r = r_0$, i.e., a limit cycle appears out of the equilibrium E_* .

Proof. From Proposition 3, the characteristic equation (6) has a pair of imaginary roots $\pm i\omega_0$ at $r = r_0$. It's easy to show that this root is simple. Thus it suffices to show that

$$\frac{dRe(\lambda)}{dr}(r_0) > 0$$

(see, for example [13]). We have:

$$Sign\frac{dRe(\lambda)}{dr}|r_{0} = Sign\{\theta_{2}^{2} + 2\theta_{3} - \theta_{1}^{2} - 4(\theta_{3}^{2} - \theta_{4}^{2})\}$$

After some calculations, we get:

$$\theta_2^2 + 2\theta_3 - \theta_1^2 - 4(\theta_3^2 - \theta_4^2) = \Gamma_1 - \Gamma_2(\Gamma_3 - 1)(3 - \Gamma_3), \tag{8}$$

where

$$\Gamma_1 = \frac{\gamma^2 \rho^2 (1 - m^2) + [\gamma - (\gamma + s + \alpha + \beta + m)^2]^2}{(\gamma + s + \alpha + \beta + m)^2},$$

$$\Gamma_2 = 4(m(\gamma + s + \alpha + \beta + m) + \gamma \alpha + \beta(s + \alpha))^2,$$

and

$$\Gamma_3 = \frac{\gamma \rho}{m(\gamma + s + \alpha + \beta + m) + \gamma \alpha + \beta(s + \alpha)}.$$

If the hypothesis (H_1) is verified, then $\Gamma_2(\Gamma_3 - 1)(3 - \Gamma_3) < 0$. Moreover, we have 0 < m < 1. Consequently

$$\frac{dRe(\lambda)}{dr}(r_0) > 0.$$

3 Direction and Stability of the Hopf Bifurcation

In this section, we use Kuznetsov's method [24] to calculate an indicator of the direction and stability of the bifurcated branches (limit cycles) from E_* .

By the change of variables: $x(t) = L(rt) - L_*$, $y(t) = U(rt) - L_*$ and $r = r_0 + \varepsilon$, where $\varepsilon \in \mathbb{R}$ is the bifurcation parameter, the system (1) becomes

$$\begin{cases} \frac{dx}{dt} = (r_0 + \varepsilon)[a_1 x + a_2 y] \\ \frac{dy}{dt} = (r_0 + \varepsilon)[b_{10} x + b_{01} y + b'_{10} x_1 + b'_{11} (y_1 + x_1 y_1 + x_1^2)] \end{cases}$$
(9)

with

$$\begin{aligned} x_1 &:= x(t-1), \quad y_1 &:= y(t-1), \\ a_1 &= -(s+\alpha+m), \quad a_2 &= \gamma, \\ b_{10} &= s, \ b_{01} &= -(\gamma+\beta+m), \end{aligned}$$
$$b_{10}' &= (\frac{m(\gamma+s+\alpha+\beta+m)+\gamma\alpha+\beta(s+\alpha)}{\gamma} - \frac{\rho L_*}{N_c}) \end{aligned}$$

and

$$b_{01}^{'} = -\frac{\rho L_{*}}{N_{c}}, \quad b_{11}^{'} = -\frac{\rho}{N_{c}}$$

Hence, system (9) is transformed into a functional differential equation in $\mathcal{C} := C([-1,0], \mathbb{R}^2)$ as follows,

$$\dot{x}(t) = \mathcal{L}_{\varepsilon}(x_t) + f_{\varepsilon}(x_t), \tag{10}$$

where $x = (x, y)^T \in \mathcal{C}$, $x_t \in \mathcal{C}$ is defined by $x_t(\theta) = x(t + \theta)$ for any $\theta \in [-1, 0]$, $\varepsilon \in \mathbf{R}$ is the bifurcation parameter, $\mathcal{L}_{\varepsilon} : \mathcal{C} \to \mathbb{R}^2$, is a bounded linear operator and $f : \mathbb{R} \times \mathcal{C} \to \mathcal{C}$ is the nonlinear operator. $\mathcal{L}_{\varepsilon}$ and f are given respectively by:

$$\mathcal{L}_{\varepsilon}(\psi) := (r_0 + \varepsilon) \left(A_1(\psi(0)) + A_2(\psi(-1)) \right)$$
(11)

and

$$f_{\varepsilon}(\psi) = (r_0 + \varepsilon) \begin{pmatrix} 0 \\ b'_{11}\psi_1(-1)\psi_2(-1) + b'_{20}\psi_1^2(-1) \end{pmatrix}.$$
 (12)

where

$$A_1 = \begin{pmatrix} -(s+\alpha+m) & \gamma \\ s & -(\gamma+\beta+m) \end{pmatrix}, \tag{13}$$

$$A_2 = \begin{pmatrix} 0 & 0\\ \Lambda & \frac{-\rho L_*}{N_c} \end{pmatrix},\tag{14}$$

with $\Lambda = \frac{m(\gamma + s + \alpha + \beta + m) + \gamma \alpha + \beta(s + \alpha)}{\gamma} - \frac{\rho L_*}{N_c}.$

Using the Riesz's representation (see [13]), we get

$$\mathcal{L} := \mathcal{L}_{\varepsilon}(\psi) = \int_{-1}^{0} d\eta(\theta, \varepsilon) \psi(\theta).$$
(15)

where,

$$\eta(\theta,\varepsilon) = (r_0 + \varepsilon) \left(A_1 \delta(\theta) + A_2 \delta(\theta + 1) \right)$$
(16)

The solution operator of Eq. (10) generates a \mathcal{C}_0 -semigroup with the infinitesimal generator $\mathcal{A}_{\varepsilon}$ defined by

$$\mathcal{A}_{\varepsilon}\psi(\theta) = \begin{cases} \frac{d\psi}{d\theta}(\theta) \text{ for } \theta \in [-1,0)\\ \int_{-1}^{0} d\eta(\theta,\varepsilon)\psi(\theta) \text{ for } \theta = 0 \end{cases}$$
(17)

We rewrite Eq. (10) as an abstract ordinary differential equation

$$\frac{dx_t}{dt} = \mathcal{A}_{\varepsilon}(x_t) + R_{\varepsilon}(x_t), \qquad (18)$$

with the nonlinear term

$$R_{\varepsilon}(x_t) = \begin{cases} 0 \text{ for } \theta \in [-1,0), \\ f_{\varepsilon}(x_t) \text{ for } \theta = 0. \end{cases}$$
(19)

We denote by \mathcal{A}^* the adjoint operator of $\mathcal{A}_{\varepsilon}$

$$\mathcal{A}^*\psi(s) = \begin{cases} -\frac{d\psi}{ds}(s), \text{ for } s \in (0,1] \\ \int_{-1}^0 \psi^T(-s)d\eta(s), \text{ for } s = 0, \end{cases}$$
(20)

where η^T is the transposed matrix of η .

In order to normalize the eigenvectors of operator A and A^{*}, we define the bilinear form

$$\langle \psi, \phi \rangle = \bar{\psi}(0)\phi(0) - \int_{-1}^{0} \int_{0}^{\theta} \bar{\psi}(\xi - \theta)d\eta(\theta, 0)\phi(\xi)d\xi$$

where $\phi \in \mathcal{C}$ and $\psi \in \mathcal{C}^* = C([0,1), (\mathbb{R}^{2*})).$

Assume that \mathcal{L} has two eigenvalues on the imaginary axis. Let $p(\theta)$ and $p^*(s)$ are eigenvectors of \mathcal{A}_0 and \mathcal{A}^* . In order to determine the Poincare normal form of operator A, we needs to calculate the eigenvector $p(\theta)$ and $p^*(s)$ corresponding to $i\omega_0r_0$ and $-i\omega_0r_0$, respectively, with $\langle p^*, p \rangle = 1$ and $\langle p^*, \bar{p} \rangle = 0$. Let P be the generalized eigenspace spanned by $p(\theta)$ and $\bar{p}(\theta)$ defined as

$$P = \{zp + \bar{z}\bar{p}, z \in \mathbb{C}\}.$$

Then the orthogonal complement of P in \mathcal{C} is

$$Q = \{ \psi \in \mathcal{C}, = 0, \ <\bar{p^*}, \psi > = 0 \}.$$

Therefore, we get a decomposition of \mathcal{C} as follows

$$\mathcal{C} = P \oplus Q. \tag{21}$$

A straightforward calculation gives

$$p(\theta) = (p_1, 1)^T e^{i\theta\omega_0 r_0}$$

and

$$p^*(s) = \kappa(p_2, 1)^T e^{is\omega_0 r_0},$$

where $p_1 = \frac{i\omega_0 - b_{01} - b'_{01}}{b_{10} + b'_{10}}$, and $p_2 = \frac{-(i\omega_0 + b_{01} + b_{01})}{a_2}$.

Using the normalization condition $\langle p^*, p \rangle = 1$, we get

$$\kappa = \frac{a_2(b_{10} + b_{10})}{\Upsilon},$$

where

$$\Upsilon = a_2(b_{10}^{'} + b_{10})(1 + b_{01}r_0e^{i\omega_0r_0}) + (i\omega_0 + b_{01} + b_{01}^{'})((i\omega_0 + b_{01} + b_{01}^{'}) - a_2b_{10}).$$

From (21), the state variable x_t of Eq. (10) could be decomposed by

$$x_t = \Phi z + w(z, \bar{z}, \theta)$$

= $-zp(\theta) - \bar{z}\bar{p}(\theta) + w(z, \bar{z}, \theta)$ (22)

where $w(z, \overline{z}, \theta) \in Q$. On the center manifold at $r = r_0$, we define

$$z(t) = \langle p^*, x_t \rangle,$$

$$w(z, \bar{z}, \theta) = x_t(\theta) - Re \{ z(t)p(\theta) \}, \quad w(z, \bar{z}) = w(z, \bar{z}, \theta).$$

Then

$$\dot{z}(t) = \langle p^*, \dot{u}_t \rangle
= \langle p^*, \mathcal{A}_0 x_t + R_0 x_t \rangle
= \langle \mathcal{A}^* p^*, x_t \rangle + \langle p^*, R_0 x_t \rangle$$
(23)

On the invariant manifold, system (1) can be written as

$$\dot{z}(t) = i\omega_0 r_0 z(t) + g(z, \bar{z}) \tag{24}$$

where

$$f_0(z, \bar{z}) = f(0, w(z, \bar{z}) + Re(z(t)p(\theta))$$
 (25)

$$g(z,\bar{z}) = p^*(0)f_0(z,\bar{z})$$

= $g_{20}\frac{z^2}{2} + g_{11}z\bar{z} + g_{02}\frac{\bar{z}^2}{2} + \dots$ (26)

and

$$w(z,\bar{z}) = w_{20}(\theta)\frac{z^2}{2} + w_{11}(\theta)z\bar{z} + w_{02}(\theta)\frac{\bar{z}^2}{2} + \dots$$

where z and \bar{z} are local coordinates for center manifold in the direction of p and $\bar{p}^*.$

Thus, from (10) and (24), we have

$$\dot{w} = \dot{u}_t - p\dot{z} - \bar{p}\dot{\bar{z}},$$

which leads to

$$\dot{w} = \mathcal{A}_0 w + H(z, \bar{z}, \theta), \tag{27}$$

where

$$H(z,\bar{z},\theta) = H_{20}(\theta)\frac{z^2}{2} + H_{11}(\theta)\frac{z\bar{z}}{2} + H_{02}(\theta)\frac{\bar{z}^2}{2} + \dots$$
(28)

By expanding (28) and identifying its coefficients, we get

$$H_{20} = -(\mathcal{A}_0 - 2i\omega_0 r_0) w_{20}(\theta),$$

$$H_{11} = -\mathcal{A}_0 w_{11}(\theta),$$

$$H_{02} = -(\mathcal{A}_0 + 2i\omega_0 r_0) w_{02}(\theta).$$

(29)

By

$$\dot{x}_t = w(z, \bar{z}) + zp(\theta) + \bar{z}\bar{p}(\theta)$$

and

$$q(\theta) = (1, p_1)^T e^{i\theta\omega_0 r_0},$$

we get

$$x_{1}(t) = p_{1}z + \bar{p}_{1}\bar{z} + \frac{1}{2}w_{20}(0)z^{2} + w_{11}(0)z\bar{z} + \frac{1}{2}w_{02}(0)\bar{z}^{2} + \dots$$

$$x_{1}(t-1) = p_{1}ze^{-i\theta\omega_{0}r_{0}} + \bar{p}_{1}\bar{z}e^{i\theta\omega_{0}r_{0}} + \frac{1}{2}w_{20}(-1)z^{2} + w_{11}(-1)z\bar{z} + \frac{1}{2}w_{02}(0)\bar{z}^{2} + \dots$$

$$x_{2}(t) = z + \bar{z} + \frac{1}{2}w_{20}(0)z^{2} + w_{11}(0)z\bar{z} + \frac{1}{2}w_{02}(0)\bar{z}^{2} + \dots$$

$$x_{2}(t-1) = ze^{-i\theta\omega_{0}r_{0}} + \bar{z}e^{i\theta\omega_{0}r_{0}} + \frac{1}{2}w_{20}(-1)z^{2} + \frac{1}{2}w_{11}(-1)z\bar{z} + \frac{1}{2}w_{02}(-1)\bar{z}^{2} + \dots$$
(30)

Comparing with (27), we have the coefficients of (26):

$$g_{20} = -2r_0 \bar{\kappa} \frac{\rho}{N_c} p_1(p_1 + 1) e^{2ir_0 \omega_0},$$

$$g_{02} = -2 \frac{\rho \bar{\kappa} r_0}{N_c} \bar{p}_1(\bar{p}_1 + 1) e^{-2ir_0 \omega_0},$$

$$g_{11} = -\frac{\rho \bar{\kappa} r_0}{N_c} (2\bar{p}_1 p_1 + \bar{p}_1 + p_1),$$

$$g_{21} = -\frac{\rho \bar{\kappa} r_0}{N_c} (e^{-i\omega_0 r_0} w_{1_{11}}(-1)(2p_1 + 1) + \frac{1}{2} \bar{p}_1 e^{-i\omega_0 r_0} w_{2_{20}}(-1) + p_1 e^{i\omega_0 r_0} w_{2_{11}}(-1) + (\frac{1}{2} + \bar{p}_1) e^{-i\omega_0 r_0} w_{1_{20}}(-1)).$$

(31)

Next, we calculate $w_{11}(\theta)$, $w_{20}(\theta)$ and g_{21} . For $\theta \in [-1,0)$, we have

$$H(z, \bar{z}, \theta) = \bar{p}^{*}(0) f_{0} p(0) - p^{*}(0) \bar{f}_{0} \bar{p}(0)$$

$$= -g(z, \bar{z}) p(\theta) - \bar{g}(z, \bar{z}) \bar{p}(\theta)$$

$$= -(g_{20} p(\theta) + \bar{g}_{02} \bar{p}(\theta)) \frac{z^{2}}{2} - (g_{11} p(\theta) + \bar{g}_{11} \bar{p}(\theta)) z \bar{z} + \dots$$
(32)

Using formula (27), we find

$$H_{20}(\theta) = -(g_{20}p(\theta) + \bar{g}_{02}\bar{p}(\theta)), \qquad (33)$$

$$H_{11}(\theta) = -(g_{11}p(\theta) + \bar{g}_{11}\bar{p}(\theta)).$$
(34)

Substituting (34) into (25) and (33) into (25), respectively, we get

$$\dot{w}_{20} = 2ir_0\omega_0 w_{20}(\theta) + g_{20}p(\theta) + \bar{g}_{02}\bar{p}(\theta).$$
(35)

 So

$$w_{20} = \frac{ig_{20}e^{ir_0\omega_0\theta}}{r_0\omega_0}p(0) + \frac{i\bar{g}_{02}e^{-ir_0\omega_0\theta}}{3r_0\omega_0}\bar{p}(0) + E_1e^{2ir_0\omega_0\theta},$$
(36)

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$$w_{11} = \frac{-ig_{11}e^{ir_0\omega_0\theta}}{r_0\omega_0}p(0) + \frac{i\bar{g}_{11}e^{-ir_0\omega_0\theta}}{r_0\omega_0}\bar{p}(0) + E_2.$$
 (37)

In the sequel, we determine E_1 and E_2 .

By (27) and the operator \mathcal{A}_0 , we have

$$\int_{-1}^{0} d\eta(\theta) w_{20}(\theta) = 2i\omega_0 r_0 w_{20}(0) - H_{20}(0), \qquad (38)$$

$$\int_{-1}^{0} d\eta(\theta) w_{11}(\theta) = -H_{11}(0).$$
(39)

From (27) and (28), we obtain

$$H_{20}(\theta) = -(g_{20}q(0) + \bar{g}_{02}\bar{p}(0)) + \frac{\rho r_0 e^{2ir_0\omega_0}}{N_c} \begin{pmatrix} 0\\ p_1(p_1+1) \end{pmatrix},$$
(40)

$$H_{11}(\theta) = -(g_{11}q(0) + \bar{g}_{11}\bar{p}(0)) + \frac{\rho r_0}{N_c} \begin{pmatrix} 0\\ 2\bar{p}_1p_1 + \bar{p}_1 + p_1 \end{pmatrix},$$
(41)

substituting (36) and (40) into (38), and noticing that

$$\left(i\omega_0 r_0 I - \int_{-1}^0 d\eta(\theta) e^{ir_0\omega_0\theta}\right) p(0) = 0, \tag{42}$$

$$\left(-i\omega_0 r_0 - \int_{-1}^0 d\eta(\theta) e^{-ir_0\omega_0\theta}\right)\bar{p}(0) = 0,$$
(43)

•

we get

$$\left(2i\omega_0 r_0 - \int_{-1}^0 d\eta(\theta) e^{2ir_0\omega_0\theta}\right) E_1 = -\frac{\rho r_0 e^{2ir_0\omega_0}}{N_c} \begin{pmatrix} 0\\ p_1(p_1+1) \end{pmatrix}, \quad (44)$$

that is

$$\begin{pmatrix} a_1 - 2ir_0\omega_0 & a_2 \\ b_{10}e^{-2ir_0\omega_0} + b'_{10} & b_{01} + b'_{01}e^{-2ir_0\omega_0} - 2ir_0\omega_0 \end{pmatrix} E_1 = -\frac{\rho r_0 e^{2ir_0\omega_0}}{N_c} \begin{pmatrix} 0 \\ p_1(p_1+1) \end{pmatrix},$$
(45)

where $E_1 = (E_{1_1}, E_{1_2})^T$, with

$$E_{1_1} = \frac{\rho r_0 a_2 p_1(p_1+1) e^{2ir_0\omega_0}}{N_c((a_1 - 2i\omega_0)(b_{01} + b'_{01}e^{-2ir_0\omega_0} - 2i\omega_0) - a_2(b_{10} + b'_{10}e^{-2ir_0\omega_0})},$$

and

$$E_{1_2} = -\frac{\rho r_0 (a_1 - 2i\omega_0) E_{1_1}}{a_2 N_c}$$

Similarly, substituting (37) and (41) into (39), we get

$$\int_{-1}^{0} d\eta(\theta) E_2 = -\frac{\rho r_0}{N_c} \begin{pmatrix} 0\\ 2\bar{p_1}p_1 + \bar{p_1} + p_1 \end{pmatrix},\tag{46}$$

that is

$$\begin{pmatrix} a_1 & a_2 \\ \\ b_{10} + b'_{10} & b_{01} + b'_{01} \end{pmatrix} E_2 = -\frac{\rho r_0}{N_c} \begin{pmatrix} 0 \\ 2\bar{p_1}p_1 + \bar{p_1} + p_1 \end{pmatrix},$$
(47)

where $E_2 = (E_{2_1}, E_{2_2})^T$, with

$$E_{2_1} = \frac{\rho r_0 a_2 p_1(p_1 + 1)(a_1(b_{01} + b'_{01}) - a_2(b_{10} + b'_{10}))}{N_c},$$

and

$$E_{2_2} = -\frac{\rho r_0 a_1 E_{1_1}}{a_2 N_c}.$$

To give our main result, we recall the definition of an indicator of direction and stability of limit cycles.

Definition 1. The first Lyapunov coefficient is given by [24]

$$l_1(r) = \frac{Re(c_1)}{\omega r} + Re(\lambda) \frac{Im(c_1)}{\omega^2 r^2},$$

where

$$c_1 = \frac{g_{21}}{2} + \frac{|g_{11}|^2}{\lambda} + \frac{|g_{02}|^2}{2(2\lambda - \bar{\lambda})} + \frac{g_{20}g_{11}(2\lambda + \bar{\lambda})}{2|\lambda|^2}.$$

For $r = r_0$, we have $\lambda = \lambda_0 = i\omega_0$, and consequently, we obtain the following result.

Theorem 2. [24] Suppose that hypothesis (H_1) holds. Then, for λ in neighborhood of λ_0 , the Eq. (10) is locally topologically equivalent to the following equation:

$$\dot{z} = (\sigma + i)z + sign(l_1(r_0))z \mid z \mid^2 + O(\mid z \mid^4),$$
(48)

with $\sigma = \frac{Re(\lambda)}{Im(\lambda)} |r_0.$

Theorem 3. [24] Suppose that hypothesis (H_1) holds. Then

- (a) if $l_1(r_0) < 0$, then a stable limit cycle appears out of the equilibrium E_* , for $r > r_0$ (supercritical Hopf bifurcation).
- (b) if $l_1(r_0) > 0$, then an unstable limit cycle appears out of the equilibrium E_* , for $r < r_0$ (sub-critical Hopf bifurcation).



Fig. 2. Stable solutions of Model (1) for a delay, r smaller than the critical value, r_0 : r = 3.1068 and $r_0 = 4.1068$



Fig. 3. Periodic solutions have bifurcated from the positive equilibrium of Model (1) for a delay closer to the critical value, $r_0 = 4.1068$

4 Numerical Simulations

4.1 Qualitative Behavior of Solutions

We consider the following hypothetical numerical parameters:

 $\gamma = 0.7, \ s = 0.01, \ \rho = 0.4, \ m = 0.005, \ \alpha = 0.01, \ \beta = 0.03, \ N_c = 11000000.$

The positive equilibrium $E^* = (1.0137 \times 10^7, 3.5561 \times 10^5)$. The first Lyapunov coefficient $l_1(r_0) = 1.106447511 \times 10^{-7}$, then the subcritical Hopf bifurcation exist and an unstable limit cycle appears out of E^* .

According to Fig. 2, Fig. 3 and Fig. 4, we observe three oscillatory regimes on the labor market: convergent oscillations towards the equilibrium (see, Fig. 2),



Fig. 4. Unstable solutions of Model (1) for a delay, r greater than the critical value: r = 5.1068 and $r_0 = 4.1068$



Fig. 5. The variation curve of the critical delay, r_0 as a function of: (a) employment level, γ , (b) the job loss rate, s and (c) the maximum population growth rate ρ



Fig. 6. The effect of the simultaneous variation of two parameters on the critical delay, r_0 : (i) (β, ρ) , (ii) (ρ, γ) and (iii) (γ, s)

periodic oscillations (see, Fig. 3) or divergent oscillations (see, Fig. 4). In summary, the number of the employed persons and the number of the unemployed oscillate around the labor market equilibrium, under the effect of the delay.

4.2 Effect of Parameters on Critical Delay

In this section, we examine the effect of the parameters on the critical delay. First, we vary one parameter and find that the critical value is a monotonic function, see Fig. 5. Next, we vary two parameters simultaneously and find a similar result, Fig. 6.

5 Conclusion

In this document, we have proposed a delayed labor model. We have studied the effect of lagging on job market fluctuations. From this analysis, we concluded that

this time lag can destabilize the Model (1) via the Hopf bifurcation phenomenon. Using the Kuznetsov method [11], we have also shown that the proposed model undergo a subcritical Hopf bifurcation and that the bifurcated limit cycles are unstable, in the vicinity of the critical lag. These results can help to control the functioning of the labor market by rationalizing the reproduction process. The difficulties we encountered are essentially related to the non-linearity of our model and the presence of the temporal deviation. In order to develop our conclusions, we plan in our next work to study the effect of two delays.

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Existence and Uniqueness Results of Fractional Differential Equations with Fuzzy Data

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Abstract. In this paper, we are going to study the existence and uniqueness solutions of fractional differential equations with fuzzy data, involving the fuzzy fractional differential operators of the order $\gamma \in \mathbb{R}_+$. The aid method of successive approximation is provided with adequate conditions for the existence and uniqueness solution. Examples are given to explain the theory obtained.

1 Introduction

Fractional differential equations (FDEs) is a generalization and integration of ordinary differential equations into arbitrary non-integer orders. This is commonly and effectively used to explain many phenomena that occur in specific scientific fields and engineering. Indeed, many applications can be found in viscoelasticity, electrochemistry, power, porous media, electromagnetic, etc. (see [1-3]). And they got a lot of attention. For the most recent work on the existence and uniqueness of solutions of initial and boundary value problems for fractional differential equations, we list [4,5,10,23,24].

Agarwal et al. [6] have taken an initiative to incorporate the idea of a solution for fuzzy fractional differential equations in order $\gamma > 0$ to get a more practical model than (FDEs). This contribution has inspired other writers to draw some results about the solution's existence and uniqueness. (see [8,10,15,16,20–23,25]. In this paper, we will study the fuzzy fractional differential equation

$$\begin{cases} D^{\gamma}y(t) = F(t, y(t)), \ t \in [0, a], \ \gamma \in \mathbf{R}_{+}, \\ D^{j}y(t)|_{t=0} = y_{j}(0), \ j = 0, 1, 2, \dots, \ k = [\gamma], \end{cases}$$
(1)

where $F: [0, a] \times \mathbb{R}_{\mathscr{F}} \to \mathbb{R}_{\mathscr{F}}$ is continuous, we shall consider this equation with some appropriate initial condition for a given equation $y_j(0) \in \mathbb{R}_{\mathscr{F}}$ ($\mathbb{R}_{\mathscr{F}}$ be the set of fuzzy real numbers [20,21]) and D^{γ} is the fuzzy fractional differential operator (γ be a positive real number with $j = [\gamma]$ ([γ] is the smallest integer greater than or equal to γ)). For earlier works concerning the crisp problem 1, the first author studied it in [11,12] when $F \in [0, a] \times C([0, a]) \to C([0, a])$ (C([0, a]))

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is the set of continuous functions defined on [0, a]) and [13] when $F \in [0, a] \times X \to X$, (X is the Banach space). Here we generalize this work for fuzzy set $\mathbb{R}_{\mathscr{F}}$.

The paper is structured as follows: In Sect. 2, we remember some basic knowledge of fuzzy calculus. Several basic principles and properties of fuzzy fractional calculus are introduced in Sect. 3 and in Sect. 4 we prove some results on the existence and uniqueness of solutions of fuzzy fractional differential equations. We denote to some examples, finally.

2 Preliminaries

We now recall some definitions needed in throughout the paper. Let us denote by $\mathbb{R}_{\mathscr{F}}$ the class of fuzzy subsets of the real axis $y : \mathbb{R} \to [0, 1]$ satisfying the following properties:

- (i) y is normal: there exists $x_0 \in \mathbb{R}$ with $y(x_0) = 1$,
- (*ii*) y is convex fuzzy set: for all $x, t \in \mathbb{R}$ and $0 < \lambda \leq 1$, it holds that

 $y(\lambda x + (1 - \lambda)t) \ge \min\{y(x), y(t)\},\$

(*iii*) y is upper semicontinuous: for any $x_0 \in \mathbb{R}$, it holds that

$$y(x_0) \ge \lim_{x \to x_0} y(x) \; ,$$

(*iv*) $[y]^0 = cl\{x \in \mathbb{R} | y(x) > 0\}$ is compact.

Then $\mathbb{R}_{\mathscr{F}}$ is called the space of fuzzy numbers see [27]. Obviously, $\mathbb{R} \subset \mathbb{R}_{\mathscr{F}}$. If y is a fuzzy set, we define $[y]^{\alpha} = \{x \in \mathbb{R} | y(x) \geq \alpha\}$ the α -level (cut) sets of y, with $0 < \alpha \leq 1$. Also, if $y \in \mathbb{R}_{\mathscr{F}}$ then α -cut of y denoted by $[y]^{\alpha} = [y_1^{\alpha}, y_2^{\alpha}]$.

Lemma 1. See ([14]) Let $y, z : \mathbb{R}_{\mathscr{F}} \to [0, 1]$ be the fuzzy sets. Then y = z if and only if $[y]^{\alpha} = [z]^{\alpha}$ for all $\alpha \in [0, 1]$.

For $y, z \in \mathbb{R}_{\mathscr{F}}$ and $\lambda \in \mathbb{R}$ the sum y + z and the product λy are defined by

$$[y+z]^{\alpha} = [y_1^{\alpha} + z_1^{\alpha}, y_2^{\alpha} + z_2^{\alpha}],$$
$$[\lambda y]^{\alpha} = \lambda [y]^{\alpha} = \begin{cases} [\lambda y_1^{\alpha}, \lambda y_2^{\alpha}], \ \lambda \ge 0;\\ [\lambda y_2^{\alpha}, \lambda y_1^{\alpha}], \ \lambda < 0, \end{cases}$$

 $\forall \alpha \in [0,1]$. Additionally if we denote $\hat{0} = \chi_{\{0\}}$, then $\hat{0} \in \mathbb{R}_{\mathscr{F}}$ is a neutral element with respect to +.

Let $d: \mathbb{R}_{\mathrm{F}} \times \mathbb{R}_{\mathscr{F}} \to \mathbb{R} + \cup \{0\}$ by the following equation:

$$d(y,z) = \sup_{\alpha \in [0,1]} d_H([y]^{\alpha}, \ [z]^{\alpha}), \text{ for ally, } z \in \mathbf{R}_{\mathscr{F}},$$

where d_H is the Hausdorff metric defined as:

$$d_H([y]^{\alpha}, [z]^{\alpha}) = \max\{|y_1^{\alpha} - z_1^{\alpha}|, |y_2^{\alpha} - z_2^{\alpha}|\}$$

The following properties are well-known see [26]:

$$\begin{aligned} d(y+w,z+w) &= d(y,z) \quad and \quad d(y,z) = d(z,y), \quad \forall \ y,z,w \in \mathbb{R}_{\mathscr{F}}, \\ d(ky,kz) &= |k|d(y,z), \quad \forall k \in \mathbb{R}, \ y,z \in \mathbb{R}_{\mathscr{F}} \\ d(y+z,w+e) &\leq d(y,w) + d(z,e), \quad \forall \ y,z,w,e \in \mathbb{R}_{\mathscr{F}}, \end{aligned}$$
(2)

and $(\mathbf{R}_{\mathscr{F}}, d)$ is a complete metric space.

Remark 1. We denote by $C([0, a], \mathbb{R}_{\mathscr{F}})$ the space of all continuous fuzzy functions on [0, a] and is a complete metric space with respect to the metric

$$h(y,z) = \sup_{t \in [0,a]} d\big(y(t), z(t)\big).$$

We denote by $L^1([0,a], \mathbb{R}_{\mathscr{F}})$ the space of all fuzzy functions $F : [0,a] \to \mathbb{R}_{\mathscr{F}}$ which are Lebesgue integrable on the bounded interval [0,a].

Definition 1. The mapping $y : [0, a] \to \mathbb{R}_{\mathscr{F}}$ for some interval [0, a] is called a fuzzy process. Therefore, its α -level set can be written as follows:

$$[y(t)]^{\alpha} = [y_1^{\alpha}(t), y_2^{\alpha}(t)], \ t \in [0, a], \ \alpha \in [0, 1].$$

Theorem 1. [8] Let $y : [0, a] \to \mathbb{R}_{\mathscr{F}}$ be Seikkala differentiable and denote $[y(t)]^{\alpha} = [y_1^{\alpha}(t), y_2^{\alpha}(t)]$. Then, the boundary function $y_1^{\alpha}(t)$ and $y_2^{\alpha}(t)$ are differentiable and

$$(y'(t))^{\alpha} = [(y_1^{\alpha})'(t), (y_2^{\alpha})'(t)], \ \alpha \in [0, 1].$$

Definition 2. [9] Let $y : [0, a] \to \mathbb{R}_{\mathscr{F}}$. The fuzzy integral, denoted by $\int_{b}^{c} y(t) dt, b, c \in [0, a]$, is defined levelwise by the following equation:

$$\left[\int_{b}^{c} y(t)dt\right]^{\alpha} = \left[\int_{b}^{c} y_{1}^{\alpha}(t)dt, \int_{b}^{c} y_{2}^{\alpha}(t)dt\right],$$

for all $0 \le \alpha \le 1$. In [9], if $y : [0, a] \to \mathbb{R}_{\mathscr{F}}$ is continuous, it is fuzzy integrable.

Theorem 2. [7] If $y \in \mathbb{R}_{\mathscr{F}}$, then the following properties hold:

- (*i*) $[y]^{\alpha_2} \subset [y]^{\alpha_1}$, *if* $0 \le \alpha_1 \le \alpha_2 \le 1$;
- (ii) $\{\alpha_k\} \subset [0,1]$ is a nondecreasing sequence which converges to α then

$$[y]^{\alpha} = \bigcap_{k \ge 1} [y]^{\alpha_k}.$$

Conversely if $A_{\alpha} = \{[y_1^{\alpha}, y_2^{\alpha}]; \alpha \in (0, 1]\}$ is a family of closed real intervals verifying (i) and (ii), then $\{A_{\alpha}\}$ defined a fuzzy number $y \in \mathbb{R}_{\mathscr{F}}$ such that $[y]^{\alpha} = A_{\alpha}$.

3 Fuzzy Fractional Integral and Fuzzy Fractional Derivative

Let $\gamma \in \mathbf{R}_+$ and $y : [0, a] \to \mathbf{R}_{\mathscr{F}}$ be such that $[y(t)]^{\alpha} = [y_1^{\alpha}(t), y_2^{\alpha}(t)]$ for all $t \in [0, a]$. Suppose that $y_1^{\alpha}, y_2^{\alpha} \in C([0, a], \mathbf{R}) \cap L^1([0, a], \mathbf{R})$ for all $\alpha \in [0, 1]$ and let

$$A_{\alpha} := \frac{1}{\Gamma(\gamma)} \Big[\int_{0}^{t} (t-s)^{\gamma-1} y_{1}^{\alpha}(s) ds, \ \int_{0}^{t} (t-s)^{\gamma-1} y_{2}^{\alpha}(s) ds \Big],$$

$$:= \Big[\Psi_{\gamma}(t) * y_{1}^{\alpha}(t), \ \Psi_{\gamma}(t) * y_{2}^{\alpha}(t) \Big].$$
(3)

Lemma 2. See ([10]) The family $\{A_{\alpha}; \alpha \in [0,1]\}$ given by 3, defined a fuzzy number $y \in \mathbb{R}_{\mathscr{F}}$ such that $[y]^{\alpha} = A_{\alpha}$.

Now for any positive real number $\gamma > 0$, we define

$$\Psi_{\gamma}(t) = \begin{cases} \frac{t^{\gamma-1}}{\Gamma(\gamma)}, t > 0, \\ 0, \quad t \le 0, \end{cases}$$

and

$$\Psi_{-\gamma}(t) = \Psi_{1+k-\gamma}(t) * \delta^{1+k}(t) , \quad k = [\gamma], \Psi_{-n}(t) = \delta^{n}(t), \quad n = 0, 1, 2, \dots.$$

with the property $\Psi_{\gamma}(t) * \Psi_{p}(t) = \Psi_{\gamma+p}(t)$ for p > 0, where $\delta^{n}(t)$ is the nth derivative of the delta function and Γ is the gamma function (for the properties of $\Psi_{\gamma}(t)$ see [17,18]).

Definition 3. Let $y \in C([0, a], \mathbb{R}_{\mathscr{F}}) \cap L^1([0, a], \mathbb{R}_{\mathscr{F}})$. The fuzzy fractional primitive of order $\gamma > 0$ of y, is defined by

$$I^{\gamma}y(t) = \frac{1}{\Gamma(\gamma)} \int_0^t (t-s)^{\gamma-1}y(s)ds,$$

by

$$\begin{bmatrix} I^{\gamma}y(t) \end{bmatrix}^{\alpha} = \frac{1}{\Gamma(\gamma)} \begin{bmatrix} \int_{0}^{t} (t-s)^{\gamma-1}y_{1}^{\alpha}(s)ds, & \int_{0}^{t} (t-s)^{\gamma-1}y_{2}^{\alpha}(s)ds \end{bmatrix}, \qquad (4)$$
$$= \begin{bmatrix} y_{1}^{\alpha}(t) * \Psi_{\gamma}(t), y_{2}^{\alpha}(t) * \Psi_{\gamma}(t) \end{bmatrix},$$

For $\gamma = 1$, we obtain $I^1 y(t) = \int_0^t y(s) ds$, $t \in [0, a]$, that is, the integral operator. Also, Subsequent properties are evident.

- (i) $I^{\gamma}(\lambda y)(t) = \lambda I^{\gamma}(y)(t)$ for each constant $\lambda \in \mathbf{R}_{\mathscr{F}}$,
- (*ii*) $I^{\gamma}(y+z)(t) = I^{\gamma}(y)(t) + I^{\gamma}(z)(t)$.

Proposition 1. [10] If $y \in C([0, a], \mathbb{R}_{\mathscr{F}}) \cap L^1([0, a], \mathbb{R}_{\mathscr{F}})$ and $p, \gamma > 0$, then we have

$$I^p I^\gamma y = I^{p+\gamma} y.$$

Definition 4. Let $y \in C^{1+k}([0, a], \mathbb{R}_{\mathscr{F}}) \cap L^1([0, a], \mathbb{R}_{\mathscr{F}})$ be a given function such that $[y]^{\alpha} = [y_1^{\alpha}, y_2^{\alpha}]$ for all $t \in [0, a]$ and $\alpha \in [0, 1]$ the fuzzy fractional differential operator is defined

$$D^{\gamma}y(t) = \frac{1}{\Gamma(1+k-\gamma)} \int_0^t (t-s)^{k-\gamma} D^{1+k}y(s)ds,$$
(5)
= $D^{1+k}y(t) * \Psi_{1+k-\gamma}(t),$

by

$$\begin{split} \left[D^{\gamma}y(t)\right]^{\alpha} &= \frac{1}{\Gamma(1+k-\gamma)} \left[\int_{0}^{t} (t-s)^{k-\gamma} D^{1+k} y_{1}^{\alpha}(s) ds, \int_{0}^{t} (t-s)^{k-\gamma} D^{1+k} y_{2}^{\alpha}(s) ds\right] \\ &= \left[D^{1+k} y_{1}^{\alpha}(t) * \Psi_{1+k-\gamma}(t), D^{1+k} y_{2}^{\alpha}(t) * \Psi_{1+k-\gamma(t)}\right]. \end{split}$$

For k = 0, we obtain

$$\left[D^{\gamma}y(t)\right]^{\alpha} = \frac{1}{\Gamma(1-\gamma)} \left[\int_0^t (t-s)^{-\gamma} \frac{d}{ds} y_1^{\alpha}(t) ds, \int_0^t (t-s)^{-\gamma} \frac{d}{ds} y_2^{\alpha}(t) ds\right],$$

provided that the equation defines a fuzzy number $D^{\gamma}y(t) \in \mathbb{R}_{\mathscr{F}}$. In fact $\left[D^{\gamma}y(t)\right]^{\alpha} = \left[D^{\gamma}y_{1}^{\alpha}(t), D^{\gamma}y_{2}^{\alpha}(t)\right]$ for all $t \in [0, a]$ and $\alpha \in [0, 1]$.

4 Existence and Uniqueness of the Fuzzy Solution

We now consider the fuzzy fractional differential equation

$$\begin{cases} D^{\gamma}y(t) = F(t, y(t)) , & t \in [0, a], \\ D^{j}y(t)|_{t=0} = y_{j}(0) \in \mathbb{R}_{\mathscr{F}}, & j = 0, 1, 2, \dots, k. \end{cases}$$
(6)

where $\gamma \in \mathbf{R}_+$ and $F \in [0, a] \times \mathbf{R}_{\mathscr{F}} \to \mathbf{R}_{\mathscr{F}}$ is a continuous function on $(0, a] \times \mathbf{R}_{\mathscr{F}}$. We call $y : [0, a] \to \mathbf{R}_{\mathscr{F}}$ a fuzzy solution of 6, if

$$D^{\gamma}y_{1}^{\alpha}(t) = f_{1}(t, y(t)), \quad D^{j}y_{1}^{\alpha}(t)|_{t=0} = y_{1j}^{\alpha}(0)$$

$$D^{\gamma}y_{2}^{\alpha}(t) = f_{2}(t, y(t)), \quad D^{j}y_{1}^{\alpha}(t)|_{t=0} = y_{2j}^{\alpha}(0)$$
(7)

for $t \in [0, a]$ and $0 < \alpha \le 1$, where

$$\begin{bmatrix} F(t,y) \end{bmatrix}^{\alpha} = \begin{bmatrix} f_1(t,y), f_2(t,y) \end{bmatrix}$$

= $\begin{bmatrix} \min\{F(t,x) : x \in [y_1^{\alpha}, y_2^{\alpha}]\}, \max\{F(t,x) : x \in [y_1^{\alpha}, y_2^{\alpha}]\} \end{bmatrix}.$

If we can solve it (uniquely), we have only to verify that the intervals $[y_1^{\alpha}(t), y_2^{\alpha}(t)]$, for all $\alpha \in (0, 1]$, define a fuzzy number $y(t) \in \mathbb{R}_{\mathscr{F}}$.

Definition 5. A mapping $y: [0, a] \to \mathbb{R}_{\mathscr{F}}$ is a solution to the problem 6 if it is continuous and satisfies the integral equation

$$y(t) = \sum_{j=0}^{k} \frac{t^{j}}{j!} y_{j}(0) + \frac{1}{\Gamma(\gamma)} \int_{0}^{t} (t-s)^{\gamma-1} F(s,y(s)) ds.$$
(8)

According to the method of successive approximation, let us consider the sequence $\{y_n(t)\}$ such that $y_0: [0, a] \to \mathbb{R}_{\mathscr{F}}$ be continuous,

$$y_n(t) = \sum_{j=0}^k \frac{t^j}{j!} y_j(0) + \frac{1}{\Gamma(\gamma)} \int_0^t (t-s)^{\gamma-1} F(s, y_{n-1}(s)) ds,$$
(9)

where n = 1, 2, 3, ...

Now we are proving the following theorem on equivalence.

Theorem 3. Let $F : [0, a] \times \mathbb{R}_{\mathscr{F}} \to \mathbb{R}_{\mathscr{F}}$ be continuous on $[0, a] \times \mathbb{R}_{\mathscr{F}}$. And suppose $\exists \eta > 0$, such that

$$d\big(F(t,y(t)),F(t,z(t))\big) \le \eta d(y(t),z(t)),\tag{10}$$

for every y(t), $z(t) \in \mathbb{R}_{\mathscr{F}}$, $t \in [0, a]$. If $\left|\frac{\eta a^{\gamma}}{\Gamma(\gamma + 1)}\right| < 1$ then the problem 6 has a unique solution $y(t) \in C([0, a], \mathbb{R}_{\mathscr{F}})$.

Proof. By using the definition 4, we can write 6 in the form

$$\left[D^{1+k}y_1^{\alpha}(t) * \Psi_{1+k-\gamma}(t), D^{1+k}y_2^{\alpha}(t) * \Psi_{1+k-\gamma}(t)\right] = \left[F(t,y)\right]^{\alpha}$$

from lemma 1

$$D^{1+k}y_1^{\alpha}(t) * \Psi_{1+k-\gamma}(t) = F(t, y_1^{\alpha}, y_2^{\alpha}),$$

$$D^{1+k}y_2^{\alpha}(t) * \Psi_{1+k-\gamma}(t) = F(t, y_1^{\alpha}, y_2^{\alpha}),$$
 (11)

where $\overline{F} = (f_1, f_2)$, operating with the convolution of $\Psi_{\gamma}(t)$, we get

$$\begin{split} D^{1+k}y_1^{\alpha}(t) * \Psi_{1+k}(t) &= \overline{F}(t, y_1^{\alpha}(t), y_2^{\alpha}(t)) * \Psi_{\gamma}(t), \\ D^{1+k}y_2^{\alpha}(t) * \Psi_{1+k}(t) &= \overline{F}(t, y_1^{\alpha}(t), y_2^{\alpha}(t)) * \Psi_{\gamma}(t), \end{split}$$

and taking into consideration the initial values 6 by choosing $y_n(0) = [y_{1n}^{\alpha}(0), y_{2n}^{\alpha}(0)]$ we obtain

$$y(t) = \sum_{j=0}^{k} \frac{t^{j}}{j!} y_{j}(0) + \frac{1}{\Gamma(\gamma)} \int_{0}^{t} (t-s)^{\gamma-1} \overline{F}(s,y(s)) ds,$$
(12)

where $y(t) = (y_1^{\alpha}(t), y_2^{\alpha}(t))$ to 7 for all $\alpha \in [0, 1]$.

We will prove that the intervals $[y_1^{\alpha}(t), y_2^{\alpha}(t)]$, for $0 < \alpha \leq 1$, define a fuzzy number. $y(t) \in \mathbb{R}_{\mathscr{F}}$ for each $t \geq 0$; Means that y is a fuzzy solution to 6.

The successive approximation $y_0 \in \mathbf{R}_{\mathscr{F}}$,

$$y_n(t) = \sum_{j=0}^k \frac{t^j}{j!} y_j(0) + \frac{1}{\Gamma(\gamma)} \int_0^t (t-s)^{\gamma-1} \overline{F}(s, y_{n-1}(s)) ds,$$

where n = 1, 2, 3...,. And the integral is the fuzzy integral, define a sequence of fuzzy numbers $y_n(t) \in \mathbb{R}_{\mathscr{F}}$. Let us show that there exists a fuzzy set-valued mapping $y : [0, a] \to \mathbb{R}_{\mathscr{F}}$ such that $d(y_n(t), y(t)) \to 0$ uniformly on $t \in [0, a]$ as $n \to \infty$.

Let $t \in [0, a]$, from 9, it follows that, for n = 1

$$y_1(t) = \sum_{j=0}^k \frac{t^j}{j!} y_j(0) + \frac{1}{\Gamma(\gamma)} \int_0^t (t-s)^{\gamma-1} \overline{F}(s, y_0(s)) ds,$$
(13)

and for n = 2 from 9

$$y_2(t) = \sum_{j=0}^k \frac{t^j}{j!} y_j(0) + \frac{1}{\Gamma(\gamma)} \int_0^t (t-s)^{\gamma-1} \overline{F}(s, y_1(s)) ds.$$
(14)

From 13 and 14, we have

$$d_{H}\left([y_{2}(t)]^{\alpha}, [y_{1}(t)]^{\alpha}\right) = d_{H}\left(\left[\frac{1}{\Gamma(\gamma)} \int_{0}^{t} (t-s)^{\gamma-1} F(s, y_{1}(s)) ds\right]^{\alpha}, \left[\frac{1}{\Gamma(\gamma)} \int_{0}^{t} (t-s)^{\gamma-1} F(s, y_{0}(s)) ds\right]^{\alpha}\right) \\ \leq \frac{1}{\Gamma(\gamma)} \int_{0}^{t} (t-s)^{\gamma-1} d_{H}\left(\left[F(s, y_{1}(s))\right]^{\alpha}, \left[F(s, y_{0}(s))\right]^{\alpha}\right) ds,$$
(15)

for any $\alpha \in [0, 1]$.

According to the condition 10 and using proprieties 2, we get

$$d(y_{2}(t), y_{1}(t)) \leq \frac{\eta}{\Gamma(\gamma)} \int_{0}^{t} (t-s)^{\gamma-1} d(y_{1}(s), y_{0}(s)) ds$$
$$\leq \frac{\eta}{\Gamma(\gamma)} \int_{0}^{t} (t-s)^{\gamma-1} \sup_{s \in [0,a]} d(y_{1}(s), y_{0}(s)) ds$$
(16)

Now, we can apply 16 to get

$$d(y_2(t), y_1(t)) \leq \frac{\eta a^{\gamma}}{\Gamma(\gamma+1)} h(y_1, y_0).$$
(17)

Starting from 16 and 17, we assume that

$$d(y_n(t), y_{n-1}(t)) \le \left(\frac{\eta a^{\gamma}}{\Gamma(\gamma+1)}\right)^{n-1} h(y_1, y_0),$$
(18)

and we will show that inequality holds for $d(y_{n+1}(t), y_n(t))$.

Indeed, from 9 and condition 10, so

$$d_{H}\left([y_{n+1}(t)]^{\alpha}, [y_{n}(t)]^{\alpha}\right) = d_{H}\left(\left[\frac{1}{\Gamma(\gamma)} \int_{0}^{t} (t-s)^{\gamma-1} F(s, y_{n}(s)) ds\right]^{\alpha}, \\ \left[\frac{1}{\Gamma(\gamma)} \int_{0}^{t} (t-s)^{\gamma-1} F(s, y_{n-1}(s)) ds\right]^{\alpha}\right) \\ \leq \frac{1}{\Gamma(\gamma)} \int_{0}^{t} (t-s)^{\gamma-1} d_{H}\left(\left[F(s, y_{n}(s))\right]^{\alpha}, \\ \left[f(s, y_{n-1}(s))\right]^{\alpha}\right) ds, \tag{19}$$

for any $\alpha \in [0, 1]$. And by properties 2, we obtain

$$d(y_{n+1}(t), y_n(t)) \leq \frac{\eta}{\Gamma(\gamma)} \int_0^t (t-s)^{\gamma-1} d(y_n(s), y_{n-1}(s)) ds$$

$$\leq \frac{\eta}{\Gamma(\gamma)} \int_0^t (t-s)^{\gamma-1} \left(\frac{\eta a^{\gamma}}{\Gamma(\gamma+1)}\right)^{n-1} h(y_1, y_0) ds \quad (20)$$

$$\leq \frac{\eta}{\Gamma(\gamma)} \left(\frac{\eta a^{\gamma}}{\Gamma(\gamma+1)}\right)^{n-1} h(y_1, y_0) \int_0^t (t-s)^{\gamma-1} ds.$$

Considering 18 we have

$$d(y_{n+1}(t), y_n(t)) \le \left(\frac{\eta a^{\gamma}}{\Gamma(\gamma+1)}\right)^n h(y_1, y_0).$$
(21)

Consequently, inequality 18 holds for n = 1, 2, ... We can also write

$$d(y_n(t), y_{n-1}(t)) \le \left(\frac{\eta a^{\gamma}}{\Gamma(\gamma+1)}\right)^{n-1} h(y_1, y_0).$$

$$(22)$$

From 22, with to the convergence, it follows that the suite having the general

 $\left(\frac{\eta a^{\gamma}}{\Gamma(\gamma+1)}\right)^{n-1} \to 0, \text{ so } d(y_n(t), y_{n-1}(t)) \to 0 \text{ uniformly on } 0 \le t \le a \text{ as } n \to \infty.$

Hence, there exists a fuzzy set-valued mapping $y: [0, a] \to \mathbb{R}_{\mathscr{F}}$ such that $d(y_n(t), y(t)) \to 0$ uniformly on $0 \le t \le a$ as $n \to \infty$. From 10 and by 2, we get

$$d\big(F(t, y_n(t), F(t, y(t))\big) \le \eta d\big(y_n(t), y(t)\big) \to 0,$$
(23)

uniformly on $0 \le t \le a$ as $n \to \infty$.

With 23 into account, from 9, we obtain, for $n \to \infty$

$$y(t) = \sum_{j=0}^{k} \frac{t^{j}}{j!} y_{j}(0) + \frac{1}{\Gamma(\gamma)} \int_{0}^{t} (t-s)^{\gamma-1} \mathbf{F}(s, y(s)) ds,$$
(24)

by the convergence of sequence 9, the end points of $\left[y_n(t)\right]^{\alpha}$ converge to $y_1^{\alpha}(t)$ and

 $y_2^\alpha(t)$ respectively. Therefore at least one continuous solution exists 6. Now, we prove that this solution is unique that, is from

$$z(t) = \sum_{j=0}^{k} \frac{t^{j}}{j!} y_{j}(0) + \frac{1}{\Gamma(\gamma)} \int_{0}^{t} (t-s)^{\gamma-1} F(s, z(s)) ds,$$
(25)

it follows that $d(y(t), z(t)) \equiv 0$ Indeed, from 9 and 25, we obtain

$$d_{H}\left([z(t)]^{\alpha}, [y_{n}(t)]^{\alpha}\right) = d_{H}\left(\left[\frac{1}{\Gamma(\gamma)}\int_{0}^{t}(t-s)^{\gamma-1}F(s, z(s))ds\right]^{\alpha}, \\ \left[\frac{1}{\Gamma(\gamma)}\int_{0}^{t}(t-s)^{\gamma-1}F(s, y_{n-1}(s))ds\right]^{\alpha}\right) \\ \leq \frac{1}{\Gamma(\gamma)}\int_{0}^{t}(t-s)^{\gamma-1}d_{H}\left(\left[F(s, z(s))\right]^{\alpha}, \left[F(s, y_{n-1}(s))\right]^{\alpha}\right)ds,$$
(26)

for any $\alpha \in [0, 1]$. And by 2, we obtain

$$d(z(t), y_n(t)) \le \frac{\eta}{\Gamma(\gamma)} \int_0^t (t-s)^{\gamma-1} d(z(s), y_{n-1}(s)) ds, \quad n = 1, 2, \dots, \quad (27)$$

but $\sup_{t\in[0,a]} d(z(t), y_0(t)) < \infty$ being a solution of 25. It follows from 27 that

$$d(z(t), y_1(t)) \le \eta \frac{a^{\gamma}}{\Gamma(\gamma+1)} h(z, y_0), \quad t \in [0, a].$$

$$(28)$$

Assume that

$$d(z(t), y_n(t)) \le \left(\eta \frac{a^{\gamma}}{\Gamma(\gamma+1)}\right)^n h(z, y_0), \quad t \in [0, a].$$
⁽²⁹⁾

From

$$d(z(t), y_{n+1}(t)) \le \frac{\eta}{\Gamma(\gamma)} \int_0^t (t-s)^{\gamma-1} d(z(s), y_n(s)) ds, \quad t \in [0, a],$$
(30)

and 29, one obtains

$$d(z(t), y_{n+1}(t)) \le \left(\eta \frac{a^{\gamma}}{\Gamma(\gamma+1)}\right)^{n+1} h(z, y_0), \quad t \in [0, a].$$
(31)

Consequently, (29) holds for any n, therefore we have

$$d(z(t), y_n(t)) = d(y(t), y_n(t)) \to 0$$
(32)

on $t \in [0, a]$ as $n \to \infty$. This proves the uniqueness of the solution for 6. Now write $\gamma = 1 + k - p$, $0 , <math>k = [\gamma]$ and consider the problem

$$\begin{cases} D^{1+k}z(t) = F(t, z(t)), & t \in [0, a], \\ D^{j}z(t)|_{t=0} = z_{j}(0) \in \mathbf{R}_{\mathscr{F}}, \ j = 0, 1, 2, \dots, k. \end{cases}$$
(33)

Then, we get the following result.

Theorem 4. Let $F \in [0, a] \times \mathbb{R}_{\mathscr{F}} \to \mathbb{R}_{\mathscr{F}}$ be continuous on $[0, a] \times \mathbb{R}_{\mathscr{F}}$. and satisfy the Lipschitz condition 10. If $p \to 0$ (i.e $\gamma \to 1+k$) then the solution of 6 coincides with the solution of 33.

Proof. Suppose that y(t) is a solution of 6 and z(t) is a solution of 33, then by the equivalence between 6 and the integral Eq. 8 and matching equivalence between 33 and the integral equation

$$z(t) = \sum_{j=0}^{k} \frac{t^{j}}{j!} y_{j}(0) + \frac{1}{\Gamma(1+k)} \int_{0}^{t} (t-s)^{k} F(s, z(s)) ds,$$
(34)

we have

$$\begin{split} d(y(t), z(t)) &\leq d\Big(\frac{1}{\Gamma(\gamma)} \int_0^t (t-s)^{\gamma-1} F(s, y(s)) ds, \frac{1}{\Gamma(1+k)} \int_0^t (t-s)^k F(s, z(s)) ds\Big) \\ &= d\Big(\frac{1}{\Gamma(\gamma)} \int_0^t (t-s)^{\gamma-1} F(s, y(s)) ds + \frac{1}{\Gamma(\gamma)} \int_0^t (t-s)^{\gamma-1} F(s, z(s)) ds, \\ &\frac{1}{\Gamma(1+k)} \int_0^t (t-s)^k F(s, z(s)) ds + \frac{1}{\Gamma(\gamma)} \int_0^t (t-s)^{\gamma-1} F(s, z(s)) ds\Big) \\ &\leq \frac{1}{\Gamma(\gamma)} d\Big(\int_0^t (t-s)^{\gamma-1} F(s, y(s)) ds, \int_0^t (t-s)^{\gamma-1} F(s, z(s)) ds\Big) \\ &+ \frac{1}{\Gamma(1+k)} d\Big(\int_0^t (t-s)^k F(s, z(s)) ds + \frac{\Gamma(1+k)}{\Gamma(\gamma)} \int_0^t (t-s)^{\gamma-1} F(s, z(s)) ds\Big) \\ &\leq \frac{1}{\Gamma(\gamma)} \int_0^t (t-s)^{\gamma-1} d(F(s, y(s)), F(s, z(s)) ds \\ &+ \frac{1}{\Gamma(1+k)} \int_0^t (t-s)^k (d(F(s, z(s)), \hat{0})|1 - \frac{\Gamma(1+k)}{\Gamma(\gamma)} (t-s)^{\gamma-1-k}|) ds. \end{split}$$

Therefore

$$h(y,z) \le \frac{\eta a^{\gamma}}{\Gamma(\gamma+1)} h(y,z) + \kappa_p, \qquad (35)$$

and hence

$$h(y,z) \le \frac{\kappa_p}{\left(1 - \frac{\eta a^{\gamma}}{\Gamma(\gamma+1)}\right)},\tag{36}$$

where

$$\kappa_p = \frac{1}{\Gamma(1+k)} \int_0^t (t-s)^k \Big(d\big(F(s,z(s)),0\big) |1 - \frac{\Gamma(1+k)}{\Gamma(\gamma)} (t-s)^p| \Big) ds.$$
(37)

Now, since

$$(t-s)^k d\big(F(s,z(s)),\hat{0}\big)|1 - \frac{\Gamma(1+k)}{\Gamma(1+k-p)}(t-s)^p| \le (t-s)^k d\big(F(s,z(s)),\hat{0}\big)$$

and

$$(t-s)^k d\Big(F\big(s,z(s)\big),\hat{0}\Big)\Big|1 - \frac{\Gamma(1+k)}{\Gamma(1+k-p)}(t-s)^p\Big| \to 0 \text{ as } p \to 0,$$

It follows from a theorem dominated by Lebesgue [19] that $\kappa_p \to 0$ as $p \to 0$ which proves that $h(y, z) \to 0$ as $\gamma \to 1 + k$

$$\lim_{\gamma \to 1+k} y(t) = z(t) \text{ in } C\Big([0,a], \mathbf{R}_{\mathscr{F}}\Big)$$

Remark 2. If the assumptions of Theorem 3 are satisfied, then

$$\lim_{\gamma \to 1+k} D^{\gamma} y(t) = D^{1+k} z(t) \text{ in } C\Big([0,a], \mathbf{R}_{\mathscr{F}}\Big).$$

From 6 and 33 we have

$$\lim_{\gamma \to 1+k} d(D^{\gamma}y(t), D^{1+k}z(t)) = \lim_{\gamma \to 1+k} d\left(F(t, y(t)), F(t, z(t))\right)$$
$$\leq \eta \lim_{\gamma \to 1+k} d\left(y(t), z(t)\right),$$

then

$$\lim_{\gamma \to 1+k} h(D^{\gamma}y, D^{1+k}z) \le \eta \lim_{\gamma \to 1+k} h(y, z) = 0$$

which proves the result.

5 Examples

In order to illustrate the previous results, we give here two examples.

Example 1. Let $t \in [0, a]$, so the function F(t, y(t)) = t + y(t) is continuous on $[0, a] \times \mathbb{R}_{\mathscr{F}}$ and Lipschitzians

$$d\Big(F(t, y(t)), F(t, z(t))\Big) \le d(y(t), z(t)),$$

for all $y, z \in \mathbb{R}_{\mathscr{F}}$ and $t \in [0, a]$ it follows that

 $d(F(t, y(t)), F(t, z(t))) \le \eta d(y(t), z(t)), with \eta = 1$

Hence, we can apply our theorems to the initial value problem

$$\begin{cases} D^{\gamma} y(t) = t + y(t), & t \in [0, a], \ \gamma \in \mathbf{R}_{+}, \\ D^{j} y(t)|_{t=0} = y_{j}(0) \in \mathbf{R}_{\mathscr{F}}, & j = 0, 1, 2 \dots, k. \end{cases}$$

Example 2. Let $t \in [0, a]$, then the function $F(t, z(t)) = z^2(t)$ is continuous on $[0, a] \times \mathbb{R}_{\mathscr{F}}$ and Lipschitzians

$$\begin{aligned} d_{H}\Big([F(t,z(t))]^{\alpha}, [F(t,y(t))]^{\alpha}\Big) &= d_{H}\left(F(t, [z(t)]^{\alpha}), F(t, [y(t)]^{\alpha})\right) \\ &= d_{H}\left([z^{2}(t)]^{\alpha}, [y^{2}(t)]^{\alpha}\right) \\ &= \max\Big\{\Big|\Big(z^{\alpha}_{1}(t))^{2} - (y^{\alpha}_{1}(t)\Big)^{2}\Big|, \left|\Big(z^{\alpha}_{2}(t))^{2} - (y^{\alpha}_{2}(t)\Big)^{2}\Big|\Big\} \\ &\leq \max\Big\{\Big|\Big(z^{\alpha}_{1}(t) - y^{\alpha}_{1}(t)\Big)\Big(z^{\alpha}_{1}(t) + y^{\alpha}_{1}(t)\Big)\Big|, \\ &\left|\Big(z^{\alpha}_{2}(t) - y^{\alpha}_{2}(t)\Big)\Big(z^{\alpha}_{2}(t) + y^{\alpha}_{2}(t)\Big)\Big|\Big\} \\ &\leq \max\Big\{\Big|\Big(z^{\alpha}_{1}(t) + y^{\alpha}_{2}(t)\Big)\Big|\Big(\Big|(z^{\alpha}_{1}(t) - y^{\alpha}_{1}(t))\Big|, \left|(z^{\alpha}_{2}(t) - y^{\alpha}_{2}(t))\Big|\Big)\Big\} \\ &\leq \Big|\Big(z^{\alpha}_{2}(t) + y^{\alpha}_{2}(t)\Big)\Big|d_{H}\Big([z(t)]^{\alpha}, [y(t)]^{\alpha}\Big) \end{aligned}$$

$$d\Big(F(t,z(t)),F(t,y(t))\Big) \leq \sup_{\alpha\in[0,1]} \left| \Big(z_2^{\alpha}(t)+y_2^{\alpha}(t)\Big) \right| d(z(t),y(t)) + \sum_{\alpha\in[0,1]} \left| \int_{-\infty}^{\infty} |f(t,y(t))|^2 d(z(t),y(t)) \right| d(z(t),y(t)) + \sum_{\alpha\in[0,1]} |f(t,y(t))|^2 d(z(t),y(t)) + \sum_{\alpha$$

for all $z, y \in \mathbb{R}_{\mathscr{F}}$ and $t \in [0, a]$ it follows that $d(F(t, z(t)), F(t, y(t))) \leq \eta d(z(t), y(t))$

with $\eta = \sup_{\alpha \in [0,1]} \left| \left(z_2^{\alpha}(t) + y_2^{\alpha}(t) \right) \right|.$

Hence, using our results to the initial value problem

$$\begin{cases} D^{\gamma} z(t) = z^{2}(t) , & t \in [0, a], \ \gamma \in \mathbf{R}_{+}, \\ D^{j} z(t)|_{t=0} = z_{j}(0) \in \mathbf{R}_{\mathscr{F}}, & j = 0, 1, 2, \dots, k. \end{cases}$$

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Approximate Efficient Solutions of Nonsmooth Vector Optimization Problems via Approximate Vector Variational Inequalities

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Abstract. In this work, we demonstrate the connection between the solutions of approximate vector variational inequalities and approximate efficient solutions of corresponding nonsmooth vector optimization problems via generalized approximate invex functions. The underlying variational inequalities are stated under the Clarke's generalized Jacobian.

1 Introduction

Various significant applications in engineering and economics can only be stated as a multiobjective optimization problem [1]. Nowadays, the connection of these problems to vector variational inequalities is well-established for differentiable convex functions [2]. In particular, results in this direction were developed under various assumptions of generalized convexity [3–7] and nonsmooth invexity [8– 11]. On the other hand, relationships between a vector variational inequality and a nonsmooth vector optimization problem (NVOP) were established under the generalized approximate convexity assumption [12–14].

This paper is devoted to the case of NVOP involving generalized approximate invex multiobjective functions, which we have introduced in [15]. Our aim is to use approximate vector variational inequalities (AVVIs) of Stampacchia and Minty type in terms of Clarke's generalized Jacobian to characterize approximate efficient solutions. It is worth mentioning that, as generalized approximate invexity is an extension of generalized approximate convexity, the results obtained in our work are improvements and generalizations of the main results in [14].

The paper is organized as follows: in Sect. 2, we give some preliminary definitions, notation, and auxiliary results. In Sect. 3, we introduce the concept of approximate efficiency for NVOPs, and derive their relationships to AVVIs using the assumption of approximate invex functions. In Sect. 4, we give an example to illustrate our main results. Finally, we conclude our paper in Sect. 5.

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2 Preliminaries

Let \mathbb{R}^n be the n-dimensional Euclidean spaces, $S \subseteq \mathbb{R}^n$ be a given nonempty set and $C \subseteq \mathbb{R}^m$ be a solid pointed convex cone. We use the following partial ordering relations:

$$u \ge_C v \Leftrightarrow u - v \in C;$$

$$u >_C v \Leftrightarrow u - v \in intC.$$

Definition 1 ([16]). Let $F: S \to \mathbb{R}^m$ be a vector-valued function. F is locally Lipschitz if for each $w \in S$ there is k > 0 and $\rho > 0$ such that, for all $u, v \in B(w; \rho)$

$$||F(u) - F(v)|| \le k ||u - v||.$$

Throughout this paper, we let $F := (F_1, ..., F_m) : S \to \mathbb{R}^m$ be a locally lipschitz function, $\theta : S \times S \to \mathbb{R}^n$ be a mapping and $\tau >_C 0$ be a vector.

Definition 2 ([16]). The Clarke's generalized Jacobian of F at $u \in S$ is given by

$$\partial F(u) = conv \{ \lim_{i \to +\infty} JF(u^{(i)}) : u^{(i)} \to u, u^{(i)} \in D \},$$

where *conv* denotes the convex hull, $JF(u^{(i)})$ indicates the Jacobian of F at $u^{(i)}$, and D is the differentiability set of F.

We note that the Clarke's generalized Jacobian is not equal to the cartesian product of the components' Clarke subdifferentials. Nevertheless, one has

$$\partial F(u) \subseteq \partial F_1(u) \times \ldots \times \partial F_m(u).$$

Note also that $\partial(-F)(u) = -\partial F(u)$.

We recall some definitions given in [15] which are a generalization of the concepts of generalized approximate convexity provided in [12, 14, 17].

Definition 3. F is called approximate (θ, τ) -invex $(A(\theta, \tau)I)$ at $w \in S$, if there is $\rho > 0$ satisfying

$$F(u) - F(v) \ge_C A_v \theta(u, v) - \tau \|\theta(u, v)\|, \quad \text{for each} \quad u, v \in B(w, \rho), \, A_v \in \partial F(v).$$

If F is $A(\theta, \tau)I$ at each $w \in S$, we say that F is $A(\theta, \tau)I$ on S.

Taking $\theta(u, v) = u - v$, approximate invexity reduces to approximate convexity [18]. The counter-example given in [15, Example 2.2] shows that approximate invexity is still more general.

Definition 4. • F is approximate pseudo (θ, τ) -invex of type 1 (AP (θ, τ) I-1) at $w \in S$ if there is $\rho > 0$ such that, whenever $u, v \in B(w, \rho)$ and if

$$F(u) - F(v) <_C -\tau \|\theta(u, v)\|,$$

then

$$A_v \theta(u, v) <_C 0$$
 for each $A_v \in \partial F(v)$

• F is approximate pseudo (θ, τ) -invex of type 2 (AP (θ, τ) I-2) at $w \in S$ if there is $\rho > 0$ such that, whenever $u, v \in B(w, \rho)$ and if

$$F(u) - F(v) <_C 0,$$

then

$$A_v \theta(u, v) + \tau \|\theta(u, v)\| <_C 0 \text{ for all } A_v \in \partial F(v).$$

Proposition 1. If F is $AP(\theta, \tau)I$ -2 at $w \in S$, then F is $AP(\theta, \tau)I$ -1 at w.

Proof. Assume that there is $\overline{\rho} > 0$ satisfying for each $u, v \in B(w, \overline{\rho})$

$$F(u) - F(v) <_C -\tau \|\theta(u, v)\|,$$

then

$$F(u) - F(v) <_C 0.$$

Since F is AP(θ, τ)I-2 at w, then there is $\rho > 0$, $\rho < \overline{\rho}$, satisfying for each $u, v \in B(w, \rho)$

$$A_v \theta(u, v) + \tau \|\theta(u, v)\| <_C 0$$
 for each $A_v \in \partial F(v)$,

which further implies that

$$A_v \theta(u, v) <_C 0$$
 for each $A_v \in \partial F(v)$.

Hence F is $AP(\theta, \tau)$ I-1 at $w \in S$.

Definition 5. • F is approximate quasi (θ, τ) -invex of type 1 (AQ (θ, τ) I-1) at $w \in S$ if there is $\rho > 0$ such that for each $u, v \in B(w, \rho)$

$$A_v \theta(u, v) - \tau \|\theta(u, v)\| >_C 0$$
, for some $A_v \in \partial F(v)$,

implies

$$F(u) >_C F(v).$$

• F is approximate quasi (θ, τ) -invex of type 2 (AQ (θ, τ) I-2) at $w \in S$ if there is $\rho > 0$ such that, for each $u, v \in B(w, \rho)$

$$A_v \theta(u, v) >_C 0$$
, for some $A_v \in \partial F(v)$,

implies

$$F(u) - F(v) >_C \tau \|\theta(u, v)\|.$$

The next proposition can be easily proven.

Proposition 2. If F is $AQ(\theta, \tau)I$ -2 at $v \in S$, then F is $AQ(\theta, \tau)I$ -1 at v.

Remark 1. •
$$A(\theta, \tau)I \Rightarrow \left[AP(\theta, \tau)I-1 \text{ and } AQ(\theta, \tau)I-1\right].$$

• There is no relation between AP(θ, τ)I-2 and AQ(θ, τ)I-2 and A(θ, τ)I (for examples, see [14]).

Now, we consider the following NVOP:

$$(NVOP)$$
 min $F(u) := (F_1(u), \dots, F_m(u))$ subject to $u \in S$,

where each $F_i: S \to \mathbb{R}$ are real-valued functions for any $i \in \{1, \dots, m\}$.

Definition 6. Let $\zeta \in S$.

(i) ζ is an efficient solution of (NVOP) iff there is no vector $u \in S$ such that

$$F(u) \leq_C F(\zeta).$$

(ii) ζ is an τ -approximate efficient solution (τ -AES) of (NVOP) iff there is no $\rho > 0$ such that, for each $u \in B(\zeta; \rho) \setminus \{\zeta\}$

$$F(u) - F(\zeta) \leq_C -\tau \|\theta(u,\zeta)\|.$$

3 Relationships Between NVOP and AVVI

Consider the following AVVI of Stampacchia and Minty type in terms of Clarke subdifferentials as follows:

(ASVVI). To find $\zeta \in S$ such that, there is no $\rho > 0$ satisfying for each $u \in B(\zeta, \rho)$ and $A_{\zeta} \in \partial F(\zeta)$

$$A_{\zeta}\theta(u,\zeta) \leq_C -\tau \|\theta(u,\zeta)\|.$$

(AMVVI). To find $\zeta \in S$ such that, there is no $\rho > 0$ satisfying for each $u \in B(\zeta, \rho)$ and $A_u \in \partial F(u)$

$$A_u\theta(u,\zeta) \leq_C -\tau \|\theta(u,\zeta)\|.$$

The following theorems describe relations between AVVI and NVOP.

Theorem 1. Let F be $A(\theta, \tau)I$ at $\zeta \in S$. If ζ solves (ASVVI) w.r.t. τ , then ζ is a 2τ -AES of (NVOP).

Proof. Assume ζ fails to be a 2τ -AES of (NVOP). It means that there is $\overline{\rho} > 0$ satisfying for each $u \in B(\zeta, \overline{\rho})$

$$F(u) - F(\zeta) \leq_C -2\tau \|\theta(u,\zeta)\|.$$
(1)

As F is $A(\theta, \tau)I$ at ζ , it follows that there is $\tilde{\rho} > 0$, satisfying

$$F(u) - F(\zeta) \ge_C A_{\zeta} \theta(u, \zeta) - \tau \| \theta(u, \zeta) \| \quad \forall \ u \in B(\zeta, \widetilde{\rho}), \ A_{\zeta} \in \partial F(\zeta).$$

By using (1) and the definition of approximate (θ, τ) – invexity, and by taking $\rho := \min(\overline{\rho}, \widetilde{\rho})$, we obtain

$$A_{\zeta}\theta(u,\zeta) - \tau \|\theta(u,\zeta)\| \leq_C -2\tau \|\theta(u,\zeta)\|.$$

Hence

$$A_{\zeta}\theta(u,\zeta) \leq_C -\tau \|\theta(u,\zeta)\|.$$

This means ζ does not solve (ASVVI) w.r.t τ .

Theorem 2. Let -F be $A(\theta, \tau)I$ at $\zeta \in S$. If $\zeta \in S$ is a τ -AES for (NVOP), then ζ solves (ASVVI) w.r.t 2τ .

Proof. Assume ζ fails to be a solution of (ASVVI) w.r.t 2τ . It means that there is $\overline{\rho} > 0$ such that, for each $u \in B(\zeta, \overline{\rho}), A_{\zeta} \in \partial F(\zeta)$, we have

$$A_{\zeta}\theta(u,\zeta) \leq_C -2\tau \|\theta(u,\zeta)\|.$$

Then

$$-A_{\zeta}\theta(u,\zeta) \ge_C 2\tau \|\theta(u,\zeta)\|. \tag{2}$$

By $\partial(-F)(\zeta) = -\partial F(\zeta)$ we deduce that $-A_{\zeta} \in \partial(-F)(\zeta)$. As -F is $A(\theta, \tau)I$ at ζ , it yields that there is $\tilde{\rho} > 0$ satisfying

$$(-F)(u) - (-F)(\zeta) \ge_C -A_{\zeta}\theta(u,\zeta) - \tau \|\theta(u,\zeta)\| \qquad \forall u \in B(\zeta,\widetilde{\rho}).$$

By using (2) and by taking $\rho := \min(\overline{\rho}, \widetilde{\rho})$, we obtain

$$-F(u)+F(\zeta)+\tau \|\theta(u,\zeta)\| \ge_C -A_{\zeta}\theta(u,\zeta) \ge_C 2\tau \|\theta(u,\zeta)\| \qquad \forall u \in B(\zeta,\rho) \setminus \{\zeta\},$$

which implies

$$F(u) - F(\zeta) \leq_C -\tau \|\theta(u,\zeta)\|.$$

Therefore ζ cannot be a τ -AES of (NVOP).

Theorem 3. Let F be $A(\theta, \tau)I$ at $\zeta \in S$ and $\theta(u, \zeta) + \theta(\zeta, u) = 0$ for any $u \in S$. If ζ solves (AMVVI) w.r.t τ , then ζ is a 2τ -AES of (NVOP).

Proof. Assume ζ fails to be a 2τ -AES of (NVOP). It means that there is $\overline{\rho} > 0$ satisfying for each $u \in B(\zeta, \overline{\rho})$

$$F(u) - F(\zeta) \leq_C -2\tau \|\theta(u,\zeta)\|.$$
(3)

As -F is $A(\theta, \tau)I$ at ζ , it yields that there is $\tilde{\rho} > 0$ satisfying

$$(-F)(\zeta) - (-F)(u) \ge_C A_v \theta(\zeta, u) - \tau \|\theta(\zeta, u)\| \qquad \forall u \in B(\zeta, \widetilde{\rho}), A_v \in \partial(-F)(u),$$

then

$$F(u) - F(\zeta) \ge_C A_v \theta(\zeta, u) - \tau \|\theta(\zeta, u)\|.$$

By using (3) and by taking $\rho := \min(\overline{\rho}, \widetilde{\rho})$, we obtain

$$A_v\theta(\zeta, u) - \tau \|\theta(\zeta, u)\| \leq_C -2\tau \|\theta(u, \zeta)\| \qquad \forall u \in B(\zeta, \rho) \setminus \{\zeta\}.$$

From $\partial(-F)(u) = -\partial F(u)$, there is $A_u = -A_v \in \partial F(u)$. Consequently, using $\theta(u,\zeta) + \theta(\zeta,u) = 0$ together with the above inequality, we deduce

$$A_u\theta(u,\zeta) \leq_C -\tau \|\theta(u,\zeta)\|.$$

This means ζ does not solve (AMVVI) w.r.t τ .

Theorem 4. Let -F be $A(\theta, \tau)I$ at $\zeta \in S$ and $\theta(u, \zeta) + \theta(\zeta, u) = 0$ for all $u \in S$. If $\zeta \in S$ is a τ -AES for (NVOP), then ζ solves (AMVVI) w.r.t 2τ .

Proof. Assume ζ fails to be a solution of (AMVVI) w.r.t 2τ . Thus, there is $\overline{\rho} > 0$ satisfying for any $u \in B(\zeta, \overline{\rho}), A_u \in \partial F(u)$

$$A_u \theta(u,\zeta) \le_C -2\tau \|\theta(u,\zeta)\|.$$
(4)

As F is $A(\theta, \tau)I$ at ζ , it yields that there is $\tilde{\rho} > 0$, such that

 $F(\zeta) - F(u) \geq_C A_u \theta(\zeta, u) - \tau \| \theta(\zeta, u) \| \quad \forall u \in B(\zeta, \widetilde{\rho}), A_u \in \partial F(u).$

Since $\theta(\zeta, u) = -\theta(u, \zeta)$, then

$$F(u) - F(\zeta) - \tau \|\theta(u,\zeta)\| \leq_C A_u \theta(\zeta, u).$$

By using (3) and by taking $\rho := \min(\overline{\rho}, \widetilde{\rho})$, we obtain

$$F(u) - F(\zeta) \leq_C -\tau \|\theta(u,\zeta)\|.$$

We conclude that ζ cannot be a τ -AES of (NVOP).

Theorem 5. Let F be $AP(\theta, \tau)I$ -2 at $\zeta \in S$. If ζ solves (ASVVI) w.r.t. τ , then ζ is a τ -AES of (NVOP).

Proof. Assume ζ fails to be a τ -AES of (NVOP). It means that there is $\overline{\rho} > 0$ satisfying for all $u \in B(\zeta, \overline{\rho})$

$$F(u) - F(\zeta) \leq_C -\tau \|\theta(u,\zeta)\| <_C 0.$$
(5)

As F is AP (θ, τ) I-2 at ζ , it yields that there is $\tilde{\rho} > 0$, such that, whenever $u \in B(\zeta, \tilde{\rho})$

$$F(u) - F(\zeta) <_C 0 \Rightarrow A_{\zeta} \theta(u, \zeta) <_C -\tau \| \theta(u, \zeta) \|, \quad \forall A_{\zeta} \in \partial F(\zeta).$$

By using (5) and the definition of approximate quasi (θ, τ) -invexity type 2, and by taking $\rho := \min(\overline{\rho}, \widetilde{\rho})$, we obtain

$$A_{\zeta}\theta(u,\zeta) \leq_C -\tau \|\theta(u,\zeta)\|.$$

This means ζ does not solve (ASVVI) w.r.t. τ .

Theorem 6. Let -F be $AQ(\theta, \tau)I$ -2 at $\zeta \in S$. If ζ is a τ -AES of (NVOP), then ζ solves (ASVVI) w.r.t. τ .

Proof. Assume ζ fails to be a solution of (ASVVI) w.r.t. τ , then, there is $\rho > 0$ satisfying for each $A_{\zeta} \in \partial F(\zeta)$ and $u \in B(\zeta, \overline{\rho})$

$$A_{\zeta}\theta(u,\zeta) \leq_C -\tau \|\theta(u,\zeta)\|.$$

Then

$$-A_{\zeta}\theta(u,\zeta) \ge_C \tau \|\theta(u,\zeta)\| >_C 0.$$
(6)

As $\partial(-F)(\zeta) = -\partial F(\zeta)$ it yields that $-A_{\zeta} \in \partial(-F)(\zeta)$. Since -F is AQ (θ, τ) I-2 at ζ , it follows that there is $\tilde{\rho} > 0$ such that, whenever $u \in B(\zeta, \tilde{\rho})$

$$-A_{\zeta}\theta(u,\zeta) >_C 0 \Rightarrow -F(u) - (-F(\zeta)) >_C \tau \|\theta(u,\zeta)\|.$$

By using (6) and the definition of approximate pseudo (θ, τ) -invexity type 2, and by taking $\rho := \min(\overline{\rho}, \widetilde{\rho})$, we get

$$F(u) - F(\zeta) \leq_C -\tau \|\theta(u,\zeta)\|.$$

Consequently ζ cannot be a τ -AES of (NVOP).

The following corollary can be deduced from Theorems 5 and 6.

Corollary 1. Let F be $AP(\theta, \tau)I$ -2 at $\zeta \in S$ and -F be $AQ(\theta, \tau)I$ -2 at ζ . ζ is a τ -AES of (NVOP) if and only if ζ solves (ASVVI) w.r.t. τ .

Theorem 7. Let F be $AQ(\theta, \tau)I$ -2 at ζ and $\theta(u, \zeta) + \theta(\zeta, u) = 0$, $\forall u \in S$. If ζ is a τ -AES of (NVOP), then ζ solves (AMVVI) w.r.t. τ .

Proof. Assume ζ fails to be a solution of (AMVVI) w.r.t. τ . Then, there is $\overline{\rho} > 0$ satisfying for each $A_u \in \partial F(u)$ and $u \in B(\zeta, \overline{\rho})$

$$A_u\theta(u,\zeta) \leq_C -\tau \|\theta(u,\zeta)\|.$$

From $\theta(u,\zeta) + \theta(\zeta,u) = 0$, we obtain

$$A_u \theta(\zeta, u) \ge_C \tau \|\theta(\zeta, u)\| >_C 0.$$
(7)

As F is AQ(θ, τ)I-2 at ζ , it yields that, there is $\tilde{\rho} > 0$ such that, whenever $u \in B(\zeta, \tilde{\rho})$

$$A_u\theta(\zeta, u) >_C 0 \Rightarrow F(\zeta) - F(u) >_C \tau \|\theta(\zeta, u)\|.$$

By using (7) and the definition of approximate quasi (θ, τ) -invexity type 2, and by taking $\rho := \min(\overline{\rho}, \widetilde{\rho})$, we deduce

$$F(u) - F(\zeta) \leq_C -\tau \|\theta(u,\zeta)\|.$$

This means that ζ is not a τ -AES of (NVOP).

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Theorem 8. Let -F be $AP(\theta, \tau)I$ -2 at ζ and $\theta(u, \zeta) + \theta(\zeta, u) = 0$, $\forall u \in S$. If $\zeta \in S$ solves (AMVVI) w.r.t. τ , then ζ is a τ -AES of (NVOP).

Proof. Assume ζ fails to be a τ -AES of (NVOP). It means that there is $\overline{\rho} > 0$ satisfying for any $u \in B(\zeta, \overline{\rho})$

$$F(u) - F(\zeta) \le_C -\tau \|\theta(u,\zeta)\|.$$

Thus

$$-F(\zeta) - (-F)(u) \le_C -\tau \|\theta(u,\zeta)\| <_C 0.$$
(8)

As -F is AP (θ, τ) I-2 at ζ , it yields that there is $\tilde{\rho} > 0$, such that, whenever $u \in B(\zeta, \tilde{\rho})$

$$-F(\zeta) - (-F)(u) <_C 0 \Rightarrow A_v \theta(u,\zeta) <_C -\tau \|\theta(u,\zeta)\|, \quad \forall A_v \in \partial(-F)(u).$$

By using (8) and the definition of approximate pseudo (θ, τ) -invexity type 2, and by taking $\rho := \min(\overline{\rho}, \widetilde{\rho})$, we obtain

$$A_v \theta(u,\zeta) \leq_C -\tau \|\theta(u,\zeta)\|, \quad \forall A_v \in \partial(-F)(u), \ u \in B(\zeta,\rho).$$

Using $\partial(-F)(u) = -\partial F(u)$, there is $A_u = -A_v \in \partial F(u)$, then we have

$$-A_u\theta(\zeta, u) \leq_C -\tau \|\theta(u, \zeta)\|.$$

Since $\theta(u, \zeta) + \theta(\zeta, u) = 0$, therefore,

$$A_u\theta(u,\zeta) \leq_C -\tau \|\theta(u,\zeta)\|.$$

This means ζ does not solve (AMVVI) w.r.t. τ .

The following corollary can be deduced from Theorems 7 and 8.

Corollary 2. Let F be $AQ(\theta, \tau)I$ -2 at $\zeta \in S$ and -F be $AP(\theta, \tau)I$ -2 at ζ and $\theta(u, \zeta) + \theta(\zeta, u) = 0$, $\forall u \in S$. ζ is a τ -AES of (NVOP) if and only if ζ solves (AMVVI) w.r.t. τ .

4 Example

Consider the following NVOP as an example to illustrate the obtained results.

$$\min_{u \in S} F(u) = \begin{cases} u^2 + 3u, & u \ge 0\\ -u^2 + 4u, & u < 0, \end{cases}$$

where $S = \mathbb{R}$, $C = \mathbb{R}^+$ and $\theta(u, v) = (u - v)^3$ for each $u, v \in S$.

The Clarke subdifferential of F at $u \in S$ is defined by

$$\partial F(u) = \begin{cases} 2u+3, & u > 0; \\ [3,4], & u = 0; \\ -2u+4, & u < 0. \end{cases}$$

For $1<\tau<2$, there is $\rho=\frac{1}{2}>0$ such that, for each $u,v\in B(\zeta,\rho),\,\zeta=0,\,A_v\in\partial F(v),$ we have

$$F(u) - F(v) = \begin{cases} (u - v)(u + v + 3) > 0, & \text{if } v > 0, u > 0, u - v > 0; \\ (u - v)(u + v + 3) < 0, & \text{if } v > 0, u > 0, u - v < 0; \\ -u^2 + 4u - v^2 - 3v < 0, & \text{if } v > 0, u \ge 0, u - v < 0; \\ u^2 + 3u + v(v - 4) > 0, & \text{if } v < 0, u \ge 0; \\ (u - v)(4 - u - v) > 0, & \text{if } v < 0, u < 0, u - v > 0; \\ (u - v)(4 - u - v) < 0, & \text{if } v < 0, u < 0, u - v < 0; \\ u^2 + 3u > 0, & \text{if } v = 0, u > 0; \\ -u^2 + 4u < 0, & \text{if } v = 0, u < 0. \end{cases}$$

Also,

$$A_{v}\theta(u,v) + \tau \|\theta(u,v)\| = \begin{cases} (2v+3-\tau)(u-v)^{3} < 0, & \text{if } v > 0, u > 0, u - v < 0; \\ (2v+3-\tau)(u-v)^{3} < 0, & \text{if } v > 0, u \le 0; \\ (-2v+4-\tau)(u-v)^{3} < 0, & \text{if } v < 0, u < 0, u - v < 0; \\ ku^{3} < 0, & \text{if } v = 0, u < 0, \end{cases}$$

where $k \in [3, 4]$. Hence, F is AP (θ, τ) I-2 at $\zeta = 0$.

Since for any u > 0, one has

$$A_{\zeta}\theta(u,\zeta) + \tau \|\theta(u,\zeta)\| = ku^3 + \tau u^3 > 0, \ k \in [2,3].$$

Hence, there is no $\rho > 0$ satisfying for each $u \in B(\zeta, \rho)$ and $A_{\zeta} \in \partial F(\zeta)$

$$A_{\zeta}\theta(u,\zeta) \leq_C -\tau \|\theta(u,\zeta)\|_{\mathcal{L}}$$

Thus, $\zeta = 0$ solves (ASVVI) w.r.t. τ .

Finally, as F is AP(θ, τ)I-2 at $\zeta = 0$, then, from Theorem 5, $\zeta = 0$ should be a τ -AES of (NVOP). Indeed, for all u > 0 we have

$$F(u) - F(\zeta) + \tau \|\theta(u,\zeta)\| = u^2 + 3u + \tau u^3 > 0.$$

Hence, there is no $\rho > 0$ such that, for each $u \in B(\zeta; \rho) \setminus \{\zeta\}$

$$F(u) - F(\zeta) \leq_C -\tau \|\theta(u,\zeta)\|.$$

Therefore, $\zeta = 0$ is a τ -AES of (NVOP).

Remark 2. In the above example, the function -F is AQ (θ, τ) I-2 at $\zeta = 0$ and $\theta(u, \zeta) + \theta(\zeta, u) = 0, \forall u \in S$. We can easily show that it verifies the conditions of Theorem 6.

5 Conclusions

We have shown the relationships between AVVI in terms of Clarke's generalized Jacobian and NVOP using the concepts of approximate efficiency and generalized approximate invexity. Our work improves that of Gupta and Mishra [14] with respect to two aspects:

- If the generalized approximate invexity assumption is replaced by generalized approximate convexity assumption, then the proof arguments remain the same. Consequently, our theorems are more general since the concept of invexity includes that of convexity as a special case.
- In addition to necessary conditions of approximate efficient solutions of NVOP, we have also provided sufficient conditions using the generalized approximate invexity of -F.

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Existence of Entropy Solutions for Anisotropic Elliptic Nonlinear Problem in Weighted Sobolev Space

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Abstract. In this paper, we will study the existence of an entropy solution to the unilateral problem for a class of nonlinear anisotropic elliptic equation, with second term being an element of $L^1(\Omega)$. Our technical approach is based on a monotony method and the truncation techniques in the framework of the weighted anisotropic Sobolev space.

1 Introduction

The unilateral elliptic problems in weighted anisotropic Sobolev space have recently attracted the attention of many authors (see [5,8,12]), who used different methods to solve the question of the existence of solutions in the framework of weighted anisotropic Sobolev space (we refer to [1,2,12,13] for more details). One of the motivations for studying the unilateral elliptic problems comes from applications of mathematical modeling of physical and mechanical processes in anisotropic continuous medium.

The purpose of this paper is to study the unilateral problem for a class of nonlinear anisotropic elliptic equation of type:

$$\begin{cases} Au - div(\phi(u)) = f & \text{in } \Omega, \\ u = 0 & \text{on } \partial \Omega, \end{cases}$$
(1)

where $\Omega \subset \mathbb{R}^N$ $(N \ge 2)$ is a bounded open subset with smooth boundary $\partial \Omega$, $1 < p_1, \dots, p_N < +\infty$, \overrightarrow{p} and \overrightarrow{w} are respectively the exponent and weight function vectors, which will be specified in the following. The term $\phi = (\phi_1, \dots, \phi_N)$ belongs to $C^0(\mathbb{R}, \mathbb{R}^N)$, $Au = -div(a(x, u, \nabla u))$ is the Leray-Lions operator defined on $W_0^{1,\overrightarrow{p}}(\Omega, \overrightarrow{w})$, with $a(x, u, \nabla u)$ is a Carathéodory's function satisfying some hypotheses which will be stated later. Finally, we mention that the second member f belongs to $L^1(\Omega)$.

In the non weighted case $w_i \equiv 1$ for any $i \in \{1,...,N\}$, by using monotony method and the truncation techniques, the authors in [4] has established the existence of an entropy solutions for anisotropic elliptic unilateral problem like (1). we refer the reader to the papers [8,15] and the references therein. Moreover, Boccardo et al. [10] studied the existence of weak solutions for nonlinear elliptic problem with

 $Au = -\sum_{i=1}^{N} \frac{\partial}{\partial x_i} \left(\left| \frac{\partial u}{\partial x_i} \right|^{p_i - 2} \frac{\partial u}{\partial x_i} \right), \text{ when } \phi_i(u) = 0 \text{ for } i = 1, \cdots, N \text{ and the right-hand side is a bounded Radon measure on } \Omega.$

In general the function ϕ_i does not belongs to $L^1_{loc}(\Omega)$. Then, the problem (1) does not admit weak solution. To avoid this situation, we use entropy solutions in this paper, this concept of entropy solution was first proposed by Benilan et al. see [7].

Motivated by the above cited papers and the results in [4], we show the existence result for the anisotropic unilateral nonlinear elliptic problem related to the equation in the problem (1). Specifically, we show the existence result of an entropy solutions for the following unilateral anisotropic problem,

$$\begin{cases} u \geq \psi \text{a.e. in} \Omega, T_k(u) \in W_0^{1, \overrightarrow{p}}(\Omega, \overrightarrow{\omega}), \\ \sum_{i=1}^N \int_{\Omega} a_i(x, u, \nabla u) \partial_i T_k(u-v) dx + \sum_{i=1}^N \int_{\Omega} \phi_i(u) \partial_i T_k(u-v) dx \\ \leq \int_{\Omega} f T_k(u-v) dx, \ \forall v \in K_{\Psi} \cap L^{\infty}(\Omega), \ \forall k > 0, \end{cases}$$
(2)

in the convex class $K_{\psi} := \{ u \in W_0^{1, \overrightarrow{p}}(\Omega, \overrightarrow{\omega}), u \ge \psi \quad \text{a.e in } \Omega \}$, where ψ is a measurable function on Ω such that

$$\psi^{+} \in W_{0}^{1,\overrightarrow{p}}(\Omega,\overrightarrow{\omega}) \cap L^{\infty}(\Omega).$$
(3)

Note that the uniqueness result being a rather delicate one, due to a counter-example by Serrin (see [16]), we also mention some works [11,14] for further remarks.

The paper is outlined as follow: In the next section, we will give a brief discussion of the weighted Lebesgue space and the weighted anisotropic Sobolev space. The Sect. 3 is dedicated to some necessary lemmas and basic assumptions of our problem. In the last section, we present the main result and proofs.

2 Preliminaries

In this section, we recall some basic properties of the weighted Lebesgue-Sobolev spaces needed to study problem (1), and we give the fundamental definitions and lemmas which will be used in the following pages.

Let Ω be a bounded open subset of $\mathbb{R}^N (N \ge 2)$ with smooth boundary $\partial \Omega$. Let p_1, \ldots, p_N be N real numbers and $\vec{p} = \{p_1, \ldots, p_N\}$ be a vector of exponent, the following vector $\vec{w} = \{w_1, \ldots, w_N\}$ be a vector of weight functions, i.e., every component w_i is a measurable function which is positive a.e. in Ω . Further, we suppose in all our considerations that

(H₁)
$$w_i \in L^1_{loc}(\Omega)$$
 and $w_i^{\frac{-1}{p_i-1}} \in L^1_{loc}(\Omega).$

for any $i = 1, \ldots, N$, we denote

$$\partial_i u = \frac{\partial u}{\partial x_i}$$
 for $i = 1, \dots, N$,

$$p^- = \min\{p_1, \ldots, p_N\}, \quad p^+ = \max\{p_1, \ldots, p_N\}.$$

We define the weighted Lebesgue space $L^p(\Omega, \gamma)$ with weight γ in Ω as, the set of all measurable functions u on Ω .

we endow it

$$\|u\|_{L^p(\Omega,\gamma)} \equiv \|u\|_{p,\gamma} = \left(\int_{\Omega} |u|^p \gamma(x) dx\right)^{\frac{1}{p}} \qquad 1 \le p < \infty.$$
(4)

We denote by $W^{1,\overrightarrow{p}}(\Omega,\overrightarrow{w})$ the weighted anisotropic Sobolev space of all functions $u \in L^{1}_{loc}(\Omega)$ such that the derivatives $\partial_{i}u$ are in $L^{p_{i}}(\Omega,w_{i})$ for any i = 1, ..., N.

This set of functions is a Banach space with respect to norm (see [12])

$$\|u\|_{1,\overrightarrow{p},\overrightarrow{w}} = \|u\|_{L^{1}(\Omega)} + \sum_{i=1}^{N} \|\partial_{i}u\|_{p_{i},w_{i}}.$$
(5)

In the following to study the Dirichlet problem, we use the functional space $W_0^{1,\overrightarrow{p}}(\Omega,\overrightarrow{w})$ defined as the closure of $C_0^{\infty}(\Omega)$ in $W^{1,\overrightarrow{p}}(\Omega,\overrightarrow{w})$ with respect to the norm (5).

Let us remark that $C_0^{\infty}(\Omega)$ is dense in $W_0^{1,\vec{p}}(\Omega,\vec{w})$ and $(W_0^{1,\vec{p}}(\Omega,\vec{w}), \|.\|_{1,\vec{p},\vec{w}})$ is a reflexive Banach space, for all i = 1, ..., N such that $1 < p_i < \infty$, (see [2] for more details).

We next recall that the dual of the weighted anisotropic Sobolev space $W_0^{1,\overrightarrow{p}}(\Omega,\overrightarrow{w})$ is equivalent to $W^{-1,\overrightarrow{p'}}(\Omega,\overrightarrow{w^*})$, where $\overrightarrow{p'}$ is the conjugate of \overrightarrow{p} , i.e. $p'_i = \frac{p_i}{p_i - 1}$ and $\overrightarrow{w^*} = \{w_i^* = w_i^{1-p'_i}, i = 1, ..., N\}.$

Remark 1. suppose there is $s_i \in]\frac{N}{p_i}, +\infty[\cap]\frac{1}{p_i-1}, +\infty[$ such that

$$w_i^{-s_i} \in L^1(\Omega), \quad \text{for all } i = 1, \cdots, N.$$
 (6)

Then, the expression

$$\|u\|_{W_0^{1,\vec{p}}(\Omega,\vec{w})} = \sum_{i=1}^N \|\partial_i u\|_{p_i,w_i}$$
(7)

is a norm defined on $W_0^{1,\vec{p}}(\Omega,\vec{w})$ which is equivalent to (5).

Note that (6) is stronger than the second integrability condition in (H_1) .

Let us consider the following exponent vector $\overrightarrow{p_s} = \{ p_{s_i} = \frac{p_i s_i}{s_i + 1}, i = 1, ..., N \}.$

Lemma 1. Suppose that (H_1) and (6) hold, we have

• If
$$p^- < N$$
, then $W_0^{1,\overrightarrow{p}}(\Omega,\overrightarrow{w}) \subset L^q(\Omega)$ for all $q \in [p^-, p^*[$, with $p^* = \frac{Np^-}{N-p^-}$.

• If $p^- = N$, then $W_0^{1, p'}(\Omega, \vec{w}) \subset L^q(\Omega)$ for all $q \in [p^-, +\infty[$.

Furthermore, the above embeddings are compacts.

The proof of this lemma comes from the fact that the following embedding (see [9] for more details)

$$W_0^{1,\overrightarrow{p}}(\Omega,\overrightarrow{w})\subset W_0^{1,\overrightarrow{p_s}}(\Omega)\subset W_0^{1,p^-}(\Omega)$$

We consider the space

$$\mathscr{T}_{0}^{1,\overrightarrow{p}}(\Omega,\overrightarrow{w}) := \left\{ u \text{ measurable in } \Omega, T_{k}(u) \in W_{0}^{1,\overrightarrow{p}}(\Omega,\overrightarrow{w}), \text{ for any } k > 0 \right\},\$$

where

$$T_k(z) := \begin{cases} z & \text{if } |z| \le k, \\ k \frac{z}{|z|} & \text{if } |z| > k. \end{cases}$$

3 Basic Assumptions and Notion of Solutions

In this section, we recall some useful technical lemmas to show our aim, and we give the assumptions of our problem.

We suppose that $a_i : \Omega \times \mathbb{R} \times \mathbb{R}^N \mapsto \mathbb{R}$ are Carathéodory functions, for i = 1, ..., Nwhich satisfies the following assumptions, for every $\xi, \xi' \in \mathbb{R}^N, \theta \in \mathbb{R}$ and a.e. in $x \in \Omega$,

$$a_i(x,\theta,\xi) \cdot \xi_i \ge \alpha \, w_i \, \big| \xi_i \big|^{p_i},\tag{8}$$

$$|a_i(x,\theta,\xi)| \le \beta w_i^{1/p_i'} \big(R_i(x) + \sigma^{1/p_i'} |\theta|^{p_i/p_i'} + w_i^{1/p_i'} |\xi_i|^{p_i-1} \big), \tag{9}$$

$$\left(a_i(x,\theta,\xi) - a_i(x,\theta,\xi')\right) \cdot \left(\xi_i - \xi_i'\right) > 0 \quad \text{for} \quad \xi_i \neq \xi_i', \tag{10}$$

where $R_i(\cdot)$ is a nonnegative function lying in $L^{p'_i}(\Omega)$ and $\alpha, \beta > 0$.

Moreover, we suppose that

$$\phi_i \in C^0(\mathbb{R}, \mathbb{R}) \quad \text{for} \quad i = 1, \dots, N, \tag{11}$$

and

$$f \in L^1(\Omega). \tag{12}$$

Lemma 2. [1] Let $g \in L^r(\Omega, \gamma)$ and $g_n \subset L^r(\Omega, \gamma)$ such that $||g_n||_{r,\gamma} \leq C$, $1 < r < \infty$, If $g_n(x) \to g(x)$ a.e. in Ω then $g_n \to g$ weakly in $L^r(\Omega, \gamma)$.

Lemma 3. [3] Suppose that (8)–(10) hold, let $(u_n)_n$ a sequence in $W_0^{1,\vec{p}}(\Omega,\vec{w})$ such that $u_n \rightarrow u$ weakly in $W_0^{1,\vec{p}}(\Omega,\vec{w})$ and

$$\sum_{i=1}^{N} \int_{\Omega} (a_i(x, u_n, \nabla u_n) - a_i(x, u, \nabla u)) \partial_i(u_n - u) dx \to 0,$$

then $u_n \longrightarrow u$ strongly in $W_0^{1,\overrightarrow{p}}(\Omega,\overrightarrow{w})$.

Lemma 4. Let $(u_n)_n$ be a sequence from $W_0^{1,\overrightarrow{p}}(\Omega,\overrightarrow{w})$, if $u_n \rightharpoonup u$ weakly in $W_0^{1,\overrightarrow{p}}(\Omega,\overrightarrow{w})$. Then $T_k(u_n)$ weakly converges to $T_k(u)$ in $W_0^{1,\overrightarrow{p}}(\Omega,\overrightarrow{w})$.

Proof. We have $u_n \rightarrow u$ weakly in $W_0^{1,\overrightarrow{p}}(\Omega,\overrightarrow{w})$ and $W_0^{1,\overrightarrow{p}}(\Omega,\overrightarrow{w}) \rightarrow L^q(\Omega)$, we obtain $u_n \rightarrow u$ strongly in $L^q(\Omega)$ and a.e. in Ω , thus $T_k(u_n) \rightarrow T_k(u)$ a.e. in Ω . On the other hand

$$\begin{aligned} \|T_{k}(u_{n})\|_{W_{0}^{1,\overrightarrow{p}}(\Omega,\overrightarrow{w})} &= \sum_{i=1}^{N} \|\partial_{i}T_{k}(u_{n})\|_{p_{i},w_{i}} \\ &\leq \sum_{i=1}^{N} \left(\int_{\Omega} |T_{k}'(u_{n})\partial_{i}u_{n}|^{p_{i}}w_{i}(x)dx \right)^{1/p_{i}} \\ &\leq \sum_{i=1}^{N} \left(\int_{\Omega} |\partial_{i}u_{n}|^{p_{i}}w_{i}(x)dx \right)^{1/p_{i}} = \|u_{n}\|_{W_{0}^{1,\overrightarrow{p}}(\Omega,\overrightarrow{w})} \end{aligned}$$

Thus $(T_k(u_n))_n$ is bounded in $W_0^{1,\overrightarrow{p}}(\Omega,\overrightarrow{w})$, consequently $T_k(u_n) \rightharpoonup T_k(u)$ weakly in $W_0^{1,\overrightarrow{p}}(\Omega,\overrightarrow{w})$.

Lemma 5. [3] If $u \in W_0^{1, \overrightarrow{p}}(\Omega, \overrightarrow{w})$ then $\sum_{i=1}^N \int_{\Omega} \partial_i u \, dx = 0$.

Proof. Since $u \in W_0^{1,\vec{p}}(\Omega, \vec{w})$ there exists $u_k \in C_0^{\infty}(\Omega)$ such that $u_k \to u$ strongly in $W_0^{1,\vec{p}}(\Omega, \vec{w})$

Moreover, since $u_k \in C_0^{\infty}(\Omega)$ by Green's Formula, we have

$$\sum_{i=1}^{N} \int_{\Omega} \partial_{i} u_{k} \, dx = \int_{\partial \Omega} u_{k} \cdot \overrightarrow{n} \, ds = 0$$

Since $\partial_i u_k \to \partial_i u$ strongly in $L^{p_i}(\Omega, w_i)$ we have $\partial_i u_k \to \partial_i u$ strongly in $L^1(\Omega)$

Passing to the limit in (3), we conclude that $\sum_{i=1}^{N} \int_{\Omega} \partial_i u \, dx = 0.$

4 Main Results

In this section we state and show the main result of our article.

The definition of an entropy solution for problem (1) can be defined as follows.

Definition 1. A function $u \in \mathscr{T}_0^{1,\overrightarrow{p}}(\Omega,\overrightarrow{w})$ such that $u \ge \psi$ a.e. in Ω is said to be an entropy solution for the unilateral problem (1), if

$$\sum_{i=1}^{N} \int_{\Omega} \left[a_i(x, u, \nabla u) \, \partial_i T_k(u - \varphi) + \phi_i(u) \, \partial_i T_k(u - \varphi) \right] dx \le \int_{\Omega} f \, T_k(u - \varphi) dx$$

for all $\varphi \in K_{\psi} \cap L^{\infty}(\Omega)$.

Theorem 1. Under the Assumptions (8)–(12), then the problem (1) admits at least one entropy solution.

Proof:

Step 1: Approximate problems.

let us consider the following approximate problems

$$\begin{cases} u_n \in K_{\psi} \\ \sum_{i=1}^N \int_{\Omega} a_i(x, u_n, \nabla u_n) \partial_i(u_n - v) dx + \sum_{i=1}^N \int_{\Omega} \phi_i^n(u_n) \partial_i(u_n - v) dx \le \int_{\Omega} f_n(u_n - v) dx \\ \forall v \in K_{\psi} \quad \text{and } \forall k > 0, \end{cases}$$

$$(13)$$

where $f_n = T_n(f)$ and $\phi_i^n(s) = \phi_i(T_n(s))$.

We define the operators Φ_n of K_{ψ} to $W_0^{-1,\vec{p}'}(\Omega,\vec{w}^*)$ by:

$$\langle \Phi_n u, v \rangle = \sum_{i=1}^N \int_{\Omega} \phi_i(T_n(u)) \partial_i v \, dx \quad \text{for all } u \in K_{\Psi} \text{ and } v \in W_0^{1, \overrightarrow{p}}(\Omega, \overrightarrow{\omega}).$$

Lemma 6. The operator $B_n = A + \Phi_n$ is pseudomonotone. Furthermore, B_n is coercive in the following sense: there exists $v_0 \in K_{\Psi}$ such that

$$\frac{\langle B_n v, v - v_0 \rangle}{\|v\|_{1,\overrightarrow{p},\overrightarrow{w}}} \longrightarrow +\infty \qquad if \quad \|v\|_{1,\overrightarrow{p},\overrightarrow{w}} \to +\infty \qquad for \quad v \in K_{\Psi}.$$

Proof. In light of the Hölder's type inequality, we get for every $u, v \in W_0^{1, \overrightarrow{p}}(\Omega, \overrightarrow{w})$,

$$\begin{split} \left| \left\langle \boldsymbol{\Phi}_{n} u, v \right\rangle \right| &\leq \sum_{i=1}^{N} \int_{\Omega} \phi_{i}(T_{n}(u)) \,\partial_{i} v \, w_{i}^{\frac{-1}{p_{i}}} \, w_{i}^{\frac{1}{p_{i}}} \, dx \\ &\leq \sum_{i=1}^{N} \left(\int_{\Omega} \left| \phi_{i}(T_{n}(u)) w_{i}^{\frac{-1}{p_{i}}} \right|^{p_{i}'} dx \right)^{\frac{1}{p_{i}'}} \left(\int_{\Omega} \left| \partial_{i} v \, w_{i}^{\frac{1}{p_{i}}} \right|^{p_{i}} dx \right)^{\frac{1}{p_{i}'}} \\ &\leq \sum_{i=1}^{N} \left(\int_{\Omega} \sup_{|s| \leq n} \left| \phi_{i}(s) \right|^{p_{i}'} w_{i}^{\frac{-p_{i}'}{p_{i}}} \, dx \right)^{\frac{1}{p_{i}'}} \left(\int_{\Omega} \left| \partial_{i} v \right|^{p_{i}} w_{i} dx \right)^{\frac{1}{p_{i}'}} \\ &\leq \sum_{i=1}^{N} \left(\int_{\Omega} (\sup_{|s| \leq n} \left| \phi_{i}(s) \right| + 1)^{p_{i}'} w_{i}^{\frac{-1}{p_{i}-1}} \, dx \right)^{\frac{1}{p_{i}'}} \left(\int_{\Omega} \left| \partial_{i} v \right|^{p_{i}} w_{i} dx \right)^{\frac{1}{p_{i}'}} \\ &\leq \sum_{i=1}^{N} \left(\sup_{|s| \leq n} \left| \phi_{i}(s) \right| + 1 \right) \left(\int_{\Omega} w_{i}^{\frac{-1}{p_{i}-1}} \, dx \right)^{\frac{1}{p_{i}'}} \left(\int_{\Omega} \left| \partial_{i} v \right|^{p_{i}} w_{i} dx \right)^{\frac{1}{p_{i}'}} \\ &\leq C(n) \| v \|_{W_{0}^{1,\overline{p}'}(\Omega,\overline{w})}, \end{split}$$

which implies that $\frac{| < \Phi_n u, v > |}{||v||_{1,\overrightarrow{p},\overrightarrow{w}}} \le C(n).$

Let $v_0 \in K_{\Psi}$, thanks to Hölder's inequality and (9), by using the following continuous embeddings $W_0^{1,p_i}(\Omega, w_i) \hookrightarrow L^{p_i}(\Omega, w_i)$, we obtain

$$\begin{split} | < Av, v_{0} > | \leq \sum_{i=1}^{N} \int_{\Omega} |a_{i}(x, v, \nabla v) \partial_{i} v_{0} w_{i}^{\frac{1}{p_{i}}} w_{i}^{\frac{1}{p_{i}}} |dx \\ \leq \sum_{i=1}^{N} \left(\int_{\Omega} |a_{i}(x, v, \nabla v) w_{i}^{\frac{1}{p_{i}}} |^{p_{i}'} dx \right)^{\frac{1}{p_{i}'}} \left(\int_{\Omega} |\partial_{i} v_{0} w_{i}^{\frac{1}{p_{i}}} |^{p_{i}} dx \right)^{\frac{1}{p_{i}}} \\ \leq \beta \sum_{i=1}^{N} \left(\int_{\Omega} \mathcal{R}_{i}^{p_{i}'}(x) + \sigma |v|^{p_{i}} + |\partial_{i} v|^{p_{i}} w_{i} dx \right)^{\frac{1}{p_{i}'}} \left(\int_{\Omega} |\partial_{i} v_{0}|^{p_{i}} w_{i} dx \right)^{\frac{1}{p_{i}}} \\ \leq \beta \sum_{i=1}^{N} \left(C_{1} + C_{2} \int_{\Omega} |\partial_{i} v|^{p_{i}} w_{i} dx + \int_{\Omega} |\partial_{i} v|^{p_{i}} w_{i} dx \right)^{\frac{1}{p_{i}'}} \left(\int_{\Omega} |\partial_{i} v_{0}|^{p_{i}} w_{i} dx \right)^{\frac{1}{p_{i}'}} \\ \leq \beta \sum_{i=1}^{N} C_{1}^{\frac{1}{p_{i}'}} \left(1 + \frac{C_{2} + 1}{C_{1}} \sum_{i=1}^{N} \int_{\Omega} |\partial_{i} v|^{p_{i}} w_{i} dx \right)^{\frac{1}{p_{i}'}} \left(\int_{\Omega} |\partial_{i} v_{0}|^{p_{i}} w_{i} dx \right)^{\frac{1}{p_{i}'}} \\ \leq \beta C_{4} \sum_{i=1}^{N} \left(1 + C_{3} \left(\sum_{i=1}^{N} \int_{\Omega} |\partial_{i} v|^{p_{i}} w_{i} dx \right)^{\frac{1}{p_{i}'}} \right) \left(\int_{\Omega} |\partial_{i} v_{0}|^{p_{i}} w_{i} dx \right)^{\frac{1}{p_{i}'}} \\ \leq \beta C_{4} \left(1 + C_{3} \left(\sum_{i=1}^{N} \int_{\Omega} |\partial_{i} v|^{p_{i}} w_{i} dx \right)^{\frac{1}{p_{i}'}} \right) \sum_{i=1}^{N} \left(\int_{\Omega} |\partial_{i} v_{0}|^{p_{i}} w_{i} dx \right)^{\frac{1}{p_{i}'}} \\ \leq \beta C_{4} \left(1 + C_{3} \left(\sum_{i=1}^{N} \int_{\Omega} |\partial_{i} v|^{p_{i}} w_{i} dx \right)^{\frac{1}{p_{i}'}} \right) \|v_{0}\|_{W_{0}^{1,\overline{p}'}(\Omega,\overline{w})}. \end{split}$$

Therefore

$$\frac{\left| < Av, v - v_0 > \right|}{\left\| v \right\|_{W_0^{1, \overrightarrow{p}}(\Omega, \overrightarrow{w})}} \ge \alpha \frac{\sum_{i=1}^N \int_{\Omega} |\partial_i v|^{p_i} w_i dx}{\left\| v \right\|_{W_0^{1, \overrightarrow{p}}(\Omega, \overrightarrow{w})}} - \frac{\beta C_4 \| v_0 \|_{W_0^{1, \overrightarrow{p}}(\Omega, \overrightarrow{w})}}{\left\| v \right\|_{W_0^{1, \overrightarrow{p}}(\Omega, \overrightarrow{w})}} - \frac{\beta C_4 C_3}{\left\| v \right\|_{W_0^{1, \overrightarrow{p}}(\Omega, \overrightarrow{w})}} \left(\sum_{i=1}^N \int_{\Omega} |\partial_i v|^{p_i} w_i dx \right)^{\frac{1}{p_-}} \| v_0 \|_{W_0^{1, \overrightarrow{p}}(\Omega, \overrightarrow{w})}.$$

Then,

$$\frac{|\langle Av, v-v_0 \rangle|}{\|v\|_{W_0^{1,\vec{p}}(\Omega,\vec{w})}} \ge \alpha \frac{\sum_{i=1}^N \int_{\Omega} |\partial_i v|^{p_i} w_i dx}{\|v\|_{W_0^{1,\vec{p}}(\Omega,\vec{w})}} \Big[1 - \frac{\beta}{\alpha} C_4 C_3 \Big(\sum_{i=1}^N \int_{\Omega} |\partial_i v|^{p_i} dx \Big)^{\frac{1}{p_-} - 1} \\ \|v_0\|_{W_0^{1,\vec{p}}(\Omega,\vec{w})} \Big] - \frac{\beta C_4 \|v_0\|_{W_0^{1,\vec{p}}(\Omega,\vec{w})}}{\|v\|_{W_0^{1,\vec{p}}(\Omega,\vec{w})}}.$$

According to Jensen's inequality, we obtain

$$\begin{split} \|v\|_{W_0^{1,\overrightarrow{p}}(\Omega,\overrightarrow{w})}^{p^+_-} &= \Big(\sum_{i=1}^N \Big(\int_{\Omega} |\partial_i v|^{p_i} w_i dx\Big)^{\frac{1}{p_i}}\Big)^{p^+_-} \\ &\leq \Big(\sum_{i=1}^N \Big(\int_{\Omega} |\partial_i v|^{p_i} w_i dx\Big)^{\frac{1}{p^+_-}}\Big)^{p^+_-} \\ &\leq C \sum_{i=1}^N \int_{\Omega} |\partial_i v|^{p_i} w_i dx, \end{split}$$

where

$$p_{-}^{+} = \begin{cases} p^{-} & \text{if } \|\partial_{i}v\|_{L^{p_{i}}(\Omega,w_{i})} \geq 1\\ p^{+} & \text{if } \|\partial_{i}v\|_{L^{p_{i}}(\Omega,w_{i})} < 1. \end{cases}$$

Then

$$\frac{\sum_{i=1}^{N} \int_{\Omega} |\partial_{i}v|^{p_{i}} w_{i} dx}{\|v\|_{W_{0}^{1,\overrightarrow{p}}(\Omega,\overrightarrow{w})}} \to +\infty \text{ and } \sum_{i=1}^{N} \int_{\Omega} |\partial_{i}v|^{p_{i}} w_{i} dx \to +\infty \text{ as } \|v\|_{W_{0}^{1,\overrightarrow{p}}(\Omega,\overrightarrow{w})} \to +\infty.$$

Using (4), we obtain $\frac{|\langle Av, v - v_0 \rangle|}{\|v\|_{W_0^{1,\vec{p}}(\Omega,\vec{w})}} \to +\infty \text{ as } \|v\|_{1,\vec{p},\vec{w}} \to +\infty.$ Since $\frac{\langle \Phi_n v, v \rangle}{\|v\|_{W_0^{1,\vec{p}}(\Omega,\vec{w})}}$ and $\frac{\langle \Phi_n v, v_0 \rangle}{\|v\|_{W_0^{1,\vec{p}}(\Omega,\vec{w})}}$ are bounded, then we get

$$\frac{\langle B_n v, v, -v_0 \rangle}{\|v\|_{W_0^{1,\overrightarrow{p}}(\Omega,\overrightarrow{w})}} = \frac{\langle Av, v - v_0 \rangle}{\|v\|_{W_0^{1,\overrightarrow{p}}(\Omega,\overrightarrow{w})}} + \frac{\langle \Phi_n v, v, -v_0 \rangle}{\|v\|_{W_0^{1,\overrightarrow{p}}(\Omega,\overrightarrow{w})}} \to +\infty \text{ as } \|v\|_{W_0^{1,\overrightarrow{p}}(\Omega,\overrightarrow{w})} \to +\infty.$$

We conclude that $B_n = A + \Phi_n$ is coercive.

It remains to show that B_n is pseudomonotone.

Let $(u_k)_k$ be a sequence in $W_0^{1, \vec{p}}(\Omega, \vec{w})$ such that

$$\begin{cases} u_k \rightharpoonup u & \text{weakly in } W_0^{1, \overrightarrow{p}'}(\Omega, \overrightarrow{w}) \\ B_n u_k \rightharpoonup \chi & \text{weakly in } W_0^{-1, \overrightarrow{p}'}(\Omega, \overrightarrow{w}^*) \\ \limsup_{k \to +\infty} < B_n u_k, u_k > \le < \chi, u > . \end{cases}$$

We will show that $\chi = B_n u$ and $\langle B_n u_k, u_k \rangle \longrightarrow \langle \chi, u \rangle$ as $k \to +\infty$. Since $W_0^{1,\overrightarrow{p}}(\Omega, \overrightarrow{w}) \hookrightarrow L^{p^-}(\Omega)$, then $u_k \to u$ strongly in $L^{p^-}(\Omega)$ and a.e. in Ω for a subsequence denoted again $(u_k)_k$. Since $(u_k)_k$ is bounded in $W_0^{1,\overrightarrow{p}}(\Omega, \overrightarrow{w})$. By using (9) we have $(a_i(x, u_k, \nabla u_k))_k$ is bounded in $L^{p'_i}(\Omega, w_i^*)$, then there exists a function $\varphi_i \in L^{p'_i}(\Omega, w_i^*)$ such that

$$a_i(x, u_k, \nabla u_k) \rightharpoonup \varphi_i \quad \text{as} \quad k \to +\infty$$
 (14)

Moreover, since $(\phi_i^n(u_k))_k$ is bounded in $L^{p'_i}(\Omega, w_i^*)$ and $\phi_i^n(u_k) \to \phi_i^n(u)$ a.e. in Ω , we obtain

$$\phi_i^n(u_k) \to \phi_i^n(u)$$
 strongly in $L^{\rho_i'}(\Omega, w_i)$ as $k \to +\infty.$ (15)

For all $v \in W_0^{1,\overrightarrow{p}}(\Omega,\overrightarrow{w})$ combining (14) and (15), we have

$$<\chi, v> = \lim_{k \to +\infty} < B_n u_k, v>$$

= $\lim_{k \to +\infty} \sum_{i=1}^N \int_{\Omega} a_i(x, u_k, \nabla u_k) \partial_i v dx + \lim_{k \to +\infty} \sum_{i=1}^N \int_{\Omega} \phi_i^n(u_k) \partial_i v dx$
= $\sum_{i=1}^N \int_{\Omega} \varphi_i \partial_i v dx + \sum_{i=1}^N \int_{\Omega} \phi_i^n(u) \partial_i v dx.$

Hence, we obtain

$$\begin{split} \limsup_{k \to +\infty} \langle B_n u_k, u_k \rangle &= \limsup_{k \to +\infty} \left[\sum_{i=1}^N \int_\Omega a_i(x, u_k, \nabla u_k) \partial_i u_k dx + \sum_{i=1}^N \int_\Omega \phi_i^n(u_k) \partial_i u_k dx \right] \\ &= \limsup_{k \to +\infty} \sum_{i=1}^N \int_\Omega a_i(x, u_k, \nabla u_k) \partial_i u_k dx + \sum_{i=1}^N \int_\Omega \phi_i^n(u) \partial_i u dx \\ &\leq \langle \chi, u \rangle \\ &= \sum_{i=1}^N \int_\Omega \phi_i \partial_i u dx + \sum_{i=1}^N \int_\Omega \phi_i^n(u) \partial_i u dx \end{split}$$

as a result

$$\limsup_{k \to +\infty} \sum_{i=1}^{N} \int_{\Omega} a_i(x, u_k, \nabla u_k) \partial_i u_k dx \le \sum_{i=1}^{N} \int_{\Omega} \varphi_i \partial_i u dx.$$
(16)

Using (10), we get $\sum_{i=1}^{N} \int_{\Omega} (a_i(x, u_k, \nabla u_k) - a_i(x, u_k, \nabla u)) (\partial_i u_k - \partial_i u) dx > 0.$ Then

$$\begin{split} \sum_{i=1}^{N} \int_{\Omega} a_{i}(x, u_{k}, \nabla u_{k}) \partial_{i} u_{k} dx &\geq -\sum_{i=1}^{N} \int_{\Omega} a_{i}(x, u_{k}, nablau) \partial_{i} u dx \\ &+ \sum_{i=1}^{N} \int_{\Omega} a_{i}(x, u_{k}, \nabla u_{k}) \partial_{i} u dx + \sum_{i=1}^{N} \int_{\Omega} a_{i}(x, u_{k}, \nabla u) \partial_{i} u_{k} dx. \end{split}$$

By (14), we have

$$\liminf_{k \to +\infty} \sum_{i=1}^{N} \int_{\Omega} a_i(x, u_k, \nabla u_k) \partial_i u_k dx \ge \sum_{i=1}^{N} \int_{\Omega} \varphi_i \partial_i u dx.$$
(17)

Using (16) and (17), we get

$$\lim_{k \to +\infty} \sum_{i=1}^{N} \int_{\Omega} a_i(x, u_k, \nabla u_k) \partial_i u_k dx = \sum_{i=1}^{N} \int_{\Omega} \varphi_i \partial_i u dx \tag{18}$$

$$\lim_{k \to +\infty} \langle B_n u_k, u_k \rangle = \lim_{k \to +\infty} \sum_{i=1}^N \int_{\Omega} a_i(x, u_k, \nabla u_k) \partial_i u_k dx + \lim_{k \to +\infty} \sum_{i=1}^N \int_{\Omega} \phi_i^n(u_k) \partial_i u_k dx$$
$$= \sum_{i=1}^N \int_{\Omega} \phi_i \partial_i u dx + \sum_{i=1}^N \int_{\Omega} \phi_i^n(u) \partial_i u dx$$
$$= \langle \chi, u \rangle.$$

Moreover, since $a_i(x, u_k, \nabla u) \longrightarrow a_i(x, u, \nabla u)$ strongly in $L^{p'_i}(\Omega, w_i)$, by using (18) we have

$$\sum_{i=1}^{N} \int_{\Omega} (a_i(x, u_k, \nabla u_k) - a_i(x, u_k, \nabla u)) (\partial_i u_k - \partial_i u) dx = 0.$$

Using Lemma 3, we obtain u_k converges to u strongly in $W_0^{1,\vec{p}}(\Omega, \vec{w})$ and a.e. in Ω , then $a_i(x, u_k, \nabla u)$ converges to $a_i(x, u, \nabla u)$ weakly in $L^{p'_i}(\Omega, w_i)$ and $\phi_i^n(u)$ converges to $\phi_i^n(u)$ strongly in $L^{p'_i}(\Omega, w_i)$. Then for all $v \in W_0^{1,\vec{p}}(\Omega, \vec{w})$, we get

$$\langle \chi, v \rangle = \lim_{k \to +\infty} \langle B_n u_k, v \rangle$$

$$= \lim_{k \to +\infty} \sum_{i=1}^N \int_{\Omega} a_i(x, u_k, \nabla u_k) \partial_i v dx + \lim_{k \to +\infty} \sum_{i=1}^N \int_{\Omega} \phi_i(u_k) \partial_i v dx$$

$$= \sum_{i=1}^N \int_{\Omega} a_i(x, u, \nabla u) \partial_i v dx + \sum_{i=1}^N \int_{\Omega} \phi_i(u) \partial_i v dx$$

$$= \langle B_n u, v \rangle$$

Therefore $B_n u = \chi$.

Proposition 1. Assume that (8)–(12) hold, then the problem (13) admits at least one solution.

Proof. From Lemma 6 and Theorem 8.2 chapter 2 in [13], then the problem (13) admits at least one solution. \Box

Step 2: A priori estimate.

Proposition 2. Under the assumptions (8)–(12) and if u_n is a solution of the approximate problem (13). Then the following assertion is valid:

$$\sum_{i=1}^{N} \int_{\Omega} |\partial_i T_k(u_n)|^{p_i} w_i dx \le C(k+1) \qquad \text{for all } k > 0,$$

where C is a constant.

Proof. Let $v = u_n - \eta T_k(u_n^+ - \psi^+)$ where $\eta \ge 0$. Since $v \in W_0^{1,\overrightarrow{p}}(\Omega, \overrightarrow{w})$ and for all η small enough, we get $v \in K_{\psi}$. We take v as test function in problem (13), we obtain

$$\begin{split} \sum_{i=1}^N \int_{\Omega} a_i(x, u_n, \nabla u_n) \partial_i T_k(u_n^+ - \psi^+) dx &+ \sum_{i=1}^N \int_{\Omega} \phi_i^n(u_n) \partial_i T_k(u_n^+ - \psi^+) dx \\ &\leq \int_{\Omega} f_n T_k(u_n^+ - \psi^+) dx. \end{split}$$

As result

$$\begin{split} \sum_{i=1}^N \int_{\Omega} a_i(x, u_n, \nabla u_n) \partial_i T_k(u_n^+ - \psi^+) dx &\leq \int_{\Omega} f_n T_k(u_n^+ - \psi^+) dx \\ &+ \sum_{i=1}^N \int_{\Omega} |\phi_i^n(u_n)| |\partial_i T_k(u_n^+ - \psi^+)| dx. \end{split}$$

Since $\partial_i T_k(u_n^+ - \psi^+) = 0$ on the set $\{u_n^+ - \psi^+ > k\}$, we get

$$\sum_{i=1}^{N} \int_{\{u_{n}^{+}-\psi^{+}\leq k\}} a_{i}(x,u_{n},\nabla u_{n})\partial_{i}(u_{n}^{+}-\psi^{+})dx \leq \int_{\Omega} f_{n}T_{k}(u_{n}^{+}-\psi^{+})dx + \sum_{i=1}^{N} \int_{\{u_{n}^{+}-\psi^{+}\leq k\}} |\phi_{i}^{n}(u_{n})||\partial_{i}(u_{n}^{+}-\psi^{+})|dx,$$

thus, we have

$$\begin{split} \sum_{i=1}^{N} \int_{\{u_{n}^{+}-\psi^{+}\leq k\}} a_{i}(x,u_{n}^{+},\nabla u_{n}^{+})\partial_{i}u_{n}^{+}dx &\leq \int_{\Omega} f_{n}T_{k}(u_{n}^{+}-\psi^{+})dx \\ + \sum_{i=1}^{N} \int_{\{u_{n}^{+}-\psi^{+}\leq k\}} |\phi_{i}^{n}(u_{n})||\partial_{i}u_{n}^{+}|w_{i}^{\frac{-1}{p_{i}}}w_{i}^{\frac{1}{p_{i}}}dx + \sum_{i=1}^{N} \int_{\{u_{n}^{+}-\psi^{+}\leq k\}} |\phi_{i}^{n}(u_{n})||\partial_{i}\psi^{+}|dx \\ + \sum_{i=1}^{N} \int_{\{u_{n}^{+}-\psi^{+}\leq k\}} |a_{i}(x,u_{n}^{+},\nabla u_{n}^{+})\partial_{i}\psi^{+}|dx \end{split}$$

According to Young's inequalities, we have for a positive constant λ

$$\begin{split} &\sum_{i=1}^{N} \int_{\{u_{n}^{+}-\psi^{+}\leq k\}} a_{i}(x,u_{n}^{+},\nabla u_{n}^{+})\partial_{i}u_{n}^{+}dx \leq \int_{\Omega} f_{n}T_{k}(u_{n}^{+}-\psi^{+})dx \\ &+ C_{1}(\alpha)\sum_{i=1}^{N} \int_{\{u_{n}^{+}-\psi^{+}\leq k\}} |\phi_{i}^{n}(T_{k+||\psi||_{\infty}}(u_{n}))|^{p_{i}'}w_{i}^{\frac{-1}{p_{i}'-1}}dx + \frac{\alpha}{6}\sum_{i=1}^{N} \int_{\{u_{n}^{+}-\psi^{+}\leq k\}} |\partial_{i}u_{n}^{+}|^{p_{i}}w_{i}dx \\ &+ \sum_{i=1}^{N} \int_{\{u_{n}^{+}-\psi^{+}\leq k\}} |\phi_{i}^{n}(T_{k+||\psi||_{\infty}}(u_{n}))||\partial_{i}\psi^{+}|dx \\ &+ \sum_{i=1}^{N} \frac{\lambda^{p_{i}'}}{p_{i}} \int_{\{u_{n}^{+}-\psi^{+}\leq k\}} |a_{i}(x,u_{n},\nabla u_{n})|^{p_{i}'}w_{i}^{1-p_{i}'}dx + \sum_{i=1}^{N} \frac{1}{p_{i}\lambda^{p_{i}}} \int_{\{u_{n}^{+}-\psi^{+}\leq k\}} |\partial_{i}\psi^{+}|^{p_{i}}w_{i}dx. \end{split}$$

Using to (9) and taking
$$\lambda = \left(\frac{p'_i \alpha}{6\beta}\right)^{\frac{1}{p'_i}}$$
, we have

$$\begin{split} &\sum_{i=1}^{N} \int_{\{u_{n}^{+}-\psi^{+}\leq k\}} a_{i}(x,u_{n},\nabla u_{n})\partial_{i}u_{n}^{+}dx \leq \int_{\Omega} f_{n}T_{k}(u_{n}^{+}-\psi^{+})dx \\ &+ C_{1}(\alpha)\sum_{i=1}^{N} \int_{\{u_{n}^{+}-\psi^{+}\leq k\}} |\phi_{i}^{n}(T_{k+\|\psi\|_{\infty}}(u_{n}))|^{p_{i}'}w_{i}^{\frac{-1}{p_{i}'-1}}dx + \frac{\alpha}{6}\sum_{i=1}^{N} \int_{\{u_{n}^{+}-\psi^{+}\leq k\}} |\partial_{i}u_{n}^{+}|^{p_{i}}w_{i}dx \\ &+ \sum_{i=1}^{N} \int_{\{u_{n}^{+}-\psi^{+}\leq k\}} |\phi_{i}^{n}(T_{k+\|\psi\|_{\infty}}(u_{n}))||\partial_{i}\psi^{+}|dx + \sum_{i=1}^{N} \frac{\alpha}{6} \int_{\{u_{n}^{+}-\psi^{+}\leq k\}} R_{i}(x)|^{p_{i}'}dx \\ &+ \sum_{i=1}^{N} \frac{\alpha}{6} \int_{\{u_{n}^{+}-\psi^{+}\leq k\}} |u_{n}^{+}|^{p_{i}}w_{i}dx + \sum_{i=1}^{N} \frac{\alpha}{6} \int_{\{u_{n}^{+}-\psi^{+}\leq k\}} |\partial_{i}u_{n}^{+}|^{p_{i}}w_{i}dx \\ &+ \sum_{i=1}^{N} \frac{(6\beta)^{p_{i}-1}}{p_{i}(p_{i}'\alpha)^{p_{i}-1}} \int_{\{u_{n}^{+}-\psi^{+}\leq k\}} |\partial_{i}\psi^{+}|^{p_{i}}w_{i}dx. \end{split}$$

Combining (3), (8), (9), (10) and (H_1) , we have

$$\sum_{i=1}^{N} \int_{\{u_{n}^{+} - \psi^{+} \le k\}} |\partial_{i}u_{n}^{+}|^{p_{i}} w_{i} dx \le Ck + C'$$
(19)

As $\{x \in \Omega, u^+ \le k\} \subset \{x \in \Omega, u^+ - \psi^+ \le k + \|\psi^+\|_{\infty}\}$, then

$$\sum_{i=1}^{N} \int_{\Omega} |\partial_i T_k(u_n^+)|^{p_i} w_i dx = \sum_{i=1}^{N} \int_{\{u^+ \le k\}} |\partial_i u_n^+|^{p_i} w_i dx \le \sum_{i=1}^{N} \int_{\{u^+ - \psi^+ \le k + \|\psi^+\|_{\infty}\}} |\partial_i u_n^+|^{p_i} w_i dx.$$

Hence, thanks to (19), we get

$$\sum_{i=1}^{N} \int_{\Omega} |\partial_i T_k(u_n^+)|^{p_i} w_i dx \le (k + \|\psi^+\|_{\infty})C + C' \quad \forall k > 0.$$
⁽²⁰⁾

Similarly taking $v = u_n + T_k(u_n^-)$ as test function in approximate problem (13), we have

$$\sum_{i=1}^{N} \int_{\Omega} |\partial_i T_k(u_n)|^{p_i} w_i dx \le C''(k+1).$$

$$\tag{21}$$

By (20) and (21), we obtain

$$\sum_{i=1}^{N} \int_{\Omega} |\partial T_k(u_n)|^{p_i} w_i(x) dx \le (k + \|\psi^+\|_{\infty} + 1)C' \quad \text{for all } k > 0.$$

Step 3: Strong convergence of truncations.

Proposition 3. If u_n is a solution of approximate problem (13). Then there is a function u and a subsequence of u_n such that

$$T_k(u_n) \to T(u)$$
 strongly in $W_0^{1, \overrightarrow{p}}(\Omega, \overrightarrow{w})$

Proof. According to Proposition 2, we obtain

$$|T_k(u_n)||_{W_0^{1,\vec{p}}(\Omega,\vec{w})} \le C(k + ||\psi^+||_{\infty} + 1)^{\frac{1}{p^-}}.$$
(22)

Firstly, we shall demonstrate that $(u_n)_n$ is a Cauchy sequence in measure in Ω . For every $\lambda > 0$, we obtain $\{|u_n - u_m| > \lambda\} \subset \{|u_n| > k\} \cup \{|u_m| > k\} \cup \{|T_k(u_n) - T_k(u_m)| > \lambda\}$, thus $meas\{|u_n - u_m| > \lambda\} < meas\{|u_n| > k\} + meas\{|u_m| > k\}$

$$neas\{|u_n - u_m| > \lambda\} \le meas\{|u_n| > k\} + meas\{|u_m| > k\} + meas\{|T_k(u_n) - T_k(u_m)| > \lambda\}.$$
(23)

Using Hölder's inequality, Lemma 1 and (22), we have

$$\begin{aligned} \mathsf{k}.meas\{|u_n| > k\} &= \int_{\{|u_n| > k\}} |T_k(u_n)| dx \leq \int_{\Omega} |T(u_n) dx \\ &\leq (meas(\Omega))^{\frac{1}{p^-}} \|T_k(u_n)\|_{L^{p^-}(\Omega)} \\ &\leq C(meas(\Omega))^{\frac{1}{p^-}} \|T_k(u_n)\|_{W_0^{1,\overrightarrow{p}}(\Omega,\overrightarrow{w})} \\ &\leq C(k + \|\psi^+\|_{\infty} + 1)^{\frac{1}{p^-}}. \end{aligned}$$

Thus, $meas\{|u_n| > k\} \le C\left(\frac{1}{k^{-1+p^-}} + \frac{1+\|\psi^+\|_{\infty}}{kp^-}\right)^{\frac{1}{p^-}} \to 0$ as $k \to +\infty$. Which means that, for each $\varepsilon > 0$, there exists k_0 such that for all $k > k_0$, we get

$$meas\{|u_n| > k\} \le \frac{\varepsilon}{3} \quad \text{and} \quad meas\{|u_m| > k\} \le \frac{\varepsilon}{3}.$$
 (24)

As the sequence $(T_k(u_n))_n$ is bounded in $W_0^{1,\overrightarrow{p}}(\Omega,\overrightarrow{w})$, then there exists a subsequence $(T_k(u_n))_n$ such that $T(u_n)$ converges to v_k a.e. in Ω , weakly in $W_0^{1,\overrightarrow{p}}(\Omega,\overrightarrow{w})$ and strongly in $L^{p^-}(\Omega)$ as *n* goes to $+\infty$. Which implies that the sequence $(T_k(u_n))_n$ is a Cauchy sequence in measure in Ω , then for all $\lambda > 0$, there is n_0 such that

$$meas\{|T_k(u_n) - T_k(u_m)| > \lambda\} \le \frac{\varepsilon}{3}, \qquad \forall n, m \ge n_0.$$
(25)

Using (23), (24) and (25), then $\forall \lambda, \varepsilon > 0$, we have

$$meas\{|u_n - u_m| > \lambda\} \le \varepsilon$$
 for all $n, m \ge n_0$.

Hence $(u_n)_n$ is a Cauchy sequence in measure in Ω , then there exists a subsequence denoted again by $(u_n)_n$ such that u_n converges to a measurable function u a.e. in Ω and

$$T_k(u_n) \rightharpoonup T(u)$$
 weakly in $W_0^{1,\overrightarrow{p}}(\Omega,\overrightarrow{w})$ and a.e. in Ω for all $k > 0$. (26)

Now, we will prove that

$$\lim_{n \to \infty} \sum_{i=1}^{N} \int_{\Omega} \left[a_i(x, T_k(u_n), \nabla T_k(u_n)) - a_i(x, T_k(u_n), \nabla T_k(u)) \right] (\partial_i T_k(u_n) - \partial_i T_k(u)) dx = 0.$$
(27)

Let us consider $v = u_n + T_1(u_n - T_m(u_n))^-$ as test function in approximate problem (13), we obtain

$$-\sum_{i=1}^{N}\int_{\Omega}a_{i}(x,u_{n},\nabla u_{n})\partial_{i}T_{1}(u_{n}-T_{m}(u_{n}))^{-}dx$$
$$-\sum_{i=1}^{N}\int_{\Omega}\phi_{i}^{n}(u_{n})\partial_{i}T_{1}(u_{n}-T_{m}(u_{n}))^{-}dx\leq-\int_{\Omega}f_{n}T_{1}(u_{n}-T_{m}(u_{n}))^{-}dx.$$

Then

$$\sum_{i=1}^{N} \int_{\{-(m+1) \le u_n \le -m\}} a_i(x, u_n, \nabla u_n) \partial_i u_n dx + \sum_{i=1}^{N} \int_{\{-(m+1) \le u_n \le -m\}} \phi_i(u_n) \partial_i u_n dx \le -\int_{\Omega} f_n T_1(u_n - T_m(u_n))^- dx.$$

We pose $\Phi_i^n(s) = \int_0^s \phi_i^n(t) \chi_{\{-(m+1) \le t \le -m\}} dt$. By using the Green's formula, we obtain

$$\sum_{i=1}^N \int_{\{-(m+1)\leq u_n\leq -m\}} \phi_i(u_n)\partial_i u_n dx = \sum_{i=1}^N \int_{\Omega} \partial_i \Phi_i^n(u_n) dx = 0.$$

Then, we have

$$\sum_{i=1}^{N} \int_{\{-(m+1)\leq u_n\leq -m\}} a_i(x,u_n,\nabla u_n)\partial_i u_n dx \leq -\int_{\Omega} f_n T_1(u_n-T_m(u_n))^- dx$$

According to Lebesgue's theorem, we have

$$\lim_{m \to +\infty} \limsup_{n \to +\infty} \int_{\Omega} f_n T_1(u_n - T_m(u_n))^- dx = 0$$

Then, we get

$$\lim_{m \to +\infty} \limsup_{n \to +\infty} \sum_{i=1}^{N} \int_{\{-(m+1) \le u_n \le -m\}} a_i(x, u_n, \nabla u_n) \partial_i u_n dx = 0.$$
(28)

Similarly, we choose $v = u_n - \eta T_1(u_n - T_m(u_n))^+$ as test function in approximate problem (13), we have

$$\lim_{m \to \infty} \limsup_{n \to \infty} \sum_{i=1}^{N} \int_{\{m \le u_n \le m+1\}} a_i(x, u_n, \nabla u_n) \partial_i u_n dx = 0.$$
(29)

We define the following function for each m > k:

$$h_m(z) = \begin{cases} 1 & \text{if } |z| \le m \\ 0 & \text{if } |z| \ge m+1 \\ m+1-|z| & \text{if } m \le |z| \le m+1, \end{cases}$$

By using in (13) the test function $\varphi = u_n - \eta (T_k(u_n) - T(u))^+ h_m(u_n)$, we obtain

$$\sum_{i=1}^{N} \int_{\Omega} a_{i}(x, u_{n}, \nabla u_{n}) \partial_{i}(T_{k}(u_{n}) - T(u))^{+} h_{m}(u_{n}) dx$$

$$+ \sum_{i=1}^{N} \int_{\Omega} a_{i}(x, u_{n} \nabla u_{n})(T_{k}(u_{n}) - T_{k}(u))^{+} \partial_{i}u_{n}h'_{m}(u_{n}) dx$$

$$+ \sum_{i=1}^{N} \int_{\Omega} \phi_{i}^{n}(u_{n}) \partial_{i}(T_{k}(u_{n}) - T_{k}(u))^{+} h_{m}(u_{n}) dx$$

$$+ \sum_{i=1}^{N} \int_{\Omega} \phi_{i}^{n}(u_{n}) \partial_{i}u_{n}(T_{k}(u_{n}) - T_{k}(u))^{+} h'_{m}(u_{n}) dx$$

$$\leq \int_{\Omega} f_{n}(T_{k}(u_{n}) - T_{k}(u))^{+} h_{m}(u_{n}) dx.$$
(30)

Using (28) and (29), we get the second integral in (30) converges to 0 when *n* and *m* tend to $+\infty$.

As $h_m(u_n) = 0$ if $|u_n| > m + 1$. Then, we obtain

$$\begin{split} \sum_{i=1}^{N} \int_{\Omega} \phi_{i}^{n}(u_{n}) \partial_{i}(T_{k}(u_{n}) - T_{k}(u))^{+} h_{m}(u_{n}) dx \\ &= \sum_{i=1}^{N} \int_{\Omega} \phi_{i}(T_{m+1}(u_{n})) h_{m}(u_{n}) \partial_{i}(T_{k}(u_{n}) - T_{k}(u))^{+} dx. \end{split}$$

By Lebesgue's theorem, we get $\phi_i^n(T_{m+1}(u_n))h_m(u_n) \to \phi_i(T(u))h_m(u)$ in $L^{p'_i}(\Omega, w_i^*)$ and $\partial_i T_k(u_n) \rightharpoonup \partial_i T(u)$ weakly in $L^{p_i}(\Omega, w_i)$ as *n* goes to $+\infty$, then the third integral in (30) converges to 0 when *n* and *m* tend to $+\infty$.

Combining (8), (28), (29) and Lebesgue's theorem, we get

$$\lim_{m \to +\infty} \lim_{n \to +\infty} \sum_{i=1}^{N} \int_{\{-(m+1) \le u_n \le -m\}} |\partial_i u_n|^{p_i} (T_k(u_n) - T_k(u))^+ w_i dx = 0,$$

and

$$\lim_{m \to +\infty} \lim_{n \to +\infty} \sum_{i=1}^{N} \int_{\{m \le u_n \le m+1\}} |\partial_i u_n|^{p_i} (T_k(u_n) - T_k(u))^+ w_i dx = 0.$$

We conclude that

$$\lim_{m\to+\infty}\lim_{n\to+\infty}\sum_{i=1}^N\int_{\Omega}a_i(x,u_n,\nabla u_n)\partial_i(T_k(u_n)-T_k(u))^+h_m(u_n)dx\leq 0,$$

which implies that

$$\lim_{m \to +\infty} \lim_{n \to +\infty} \sum_{i=1}^{N} \int_{\{T_k(u_n) - T_k(u) \ge 0, |u_n| \le k\}} a_i(x, u_n, \nabla u_n) \partial_i(T_k(u_n) - T_k(u)) h_m(u_n) dx$$
$$- \lim_{m \to +\infty} \lim_{n \to +\infty} \sum_{i=1}^{N} \int_{\{T_k(u_n) - T_k(u) \ge 0, |u_n| > k\}} a_i(x, u_n, \nabla u_n) \partial_i T_k(u) h_m(u_n) dx \le 0.$$

As $h_m(u_n) = 0$ in $\{|u_n| > m+1\}$, then we obtain

$$\sum_{i=1}^{N} \int_{\{T_{k}(u_{n})-T_{k}(u)\geq 0, |u_{n}|>k\}} a_{i}(x,u_{n},\nabla u_{n})\partial_{i}T_{k}(u)h_{m}(u_{n})dx$$
$$= \sum_{i=1}^{N} \int_{\{T_{k}(u_{n})-T_{k}(u)\geq 0, |u_{n}|>k\}} a_{i}(x,T_{m+1}(u_{n}),\nabla T_{m+1}(u_{n}))\partial_{i}T_{k}(u)h_{m}(u_{n})dx.$$

Since $(a_i(x, T_{m+1}(u_n), \nabla T_{m+1}(u_n)))_{n \ge 0}$ is bounded in $L^{p'_i}(\Omega, w^*_i)$.

We have $a_i(x, T_{m+1}(u_n), \nabla T(u))$ converges to Y_m^i weakly in $L^{p'_i}(\Omega, w_i^*)$. Hence

$$\begin{split} \lim_{m \to +\infty} \lim_{n \to +\infty} \sum_{i=1}^{N} \int_{\{T_{k}(u_{n}) - T_{k}(u) \ge 0, |u_{n}| > k\}} a_{i}(x, T_{m+1}(u_{n}), \nabla T_{m+1}(u_{n})) \partial_{i} T_{k}(u) h_{m}(u_{n}) dx \\ &= \lim_{m \to +\infty} \sum_{i=1}^{N} \int_{\{|u| > k\}} Y_{m}^{i} \partial_{i} T_{k}(u) h_{m}(u) dx = 0, \end{split}$$

as results

$$\lim_{m \to +\infty} \lim_{n \to +\infty} \sum_{i=1}^{N} \int_{\{T_{k}(u_{n}) - T_{k}(u) \ge 0\}} a_{i}(x, T_{k}(u_{n}), \nabla T_{k}(u_{n})) \\ \partial_{i}(T_{k}(u_{n}) - T_{k}(u))h_{m}(u_{n})dx \le 0.$$
(31)

Moreover, we have $a_i(x, T_k(u_n), \nabla T_k(u))h_m(u_n) \rightarrow a_i(x, T_k(u), \nabla T_k(u))h_m(u)$ in $L^{p'_i}(\Omega, w_i^*)$ and $\partial_i(T_k(u_n) - T_k(u))$ converges to 0 weakly in $L^{p_i}(\Omega, w_i)$, then

$$\lim_{m \to +\infty} \lim_{n \to +\infty} \sum_{i=1}^{N} \int_{\{T_{k}(u_{n}) - T_{k}(u) \ge 0\}} a_{i}(x, T_{k}(u_{n}), \nabla T_{k}(u)) \\ \partial_{i}(T_{k}(u_{n}) - T_{k}(u))h_{m}(u_{n})dx = 0.$$
(32)

According to (10), (31) and (32), we deduce

$$\lim_{m \to +\infty} \lim_{n \to +\infty} \sum_{i=1}^{N} \int_{\{T_{k}(u_{n}) - T_{k}(u) \ge 0\}} \left[a_{i}(x, T_{k}(u_{n}), \nabla T_{k}(u_{n})) - a_{i}(x, T_{k}(u_{n}), \nabla T_{k}(u)) \right] \\ \partial_{i}(T_{k}(u_{n}) - T_{k}(u)) h_{m}(u_{n}) dx = 0.$$
(33)

Similarly, we choose $\varphi = u_n + (T_k(u_n) - T_k(u))^- h_m(u_n)$ as test function in (13), we obtain

$$\lim_{m \to +\infty} \lim_{n \to +\infty} \sum_{i=1}^{N} \int_{\{T_{k}(u_{n}) - T_{k}(u) \le 0\}} \left[a_{i}(x, T_{k}(u_{n}), \nabla T_{k}(u_{n})) - a_{i}(x, T_{k}(u_{n}), \nabla T_{k}(u)) \right] \\ \partial_{i}(T_{k}(u_{n}) - T_{k}(u)) h_{m}(u_{n}) dx = 0.$$
(34)

Using (33) and (34), we have

$$\lim_{m \to +\infty} \lim_{n \to +\infty} \sum_{i=1}^{N} \int_{\Omega} \left[a_i(x, T_k(u_n), \nabla T_k(u_n)) - a_i(x, T_k(u_n), \nabla T_k(u)) \right]$$
$$\partial_i(T_k(u_n) - T_k(u)) h_m(u_n) dx = 0.$$
(35)

Now, we show

$$\lim_{m \to +\infty} \lim_{n \to +\infty} \sum_{i=1}^{N} \int_{\Omega} \left[a_i(x, T_k(u_n), \nabla T_k(u_n)) - a_i(x, T_k(u_n), \nabla T_k(u)) \right]$$
$$\partial_i(T_k(u_n) - T_k(u))(1 - h_m(u_n)) dx = 0.$$

Let $\varphi = u_n + T_k(u_n)^-(1 - h_m(u_n))$ as test function in approximate problem (1), we obtain

$$-\sum_{i=1}^{N} \int_{\Omega} a_{i}(x, u_{n}, \nabla u_{n}) \partial_{i} T_{k}(u_{n})^{-} (1 - h_{m}(u_{n})) dx$$

+
$$\sum_{i=1}^{N} \int_{\Omega} a_{i}(x, u_{n}, \nabla u_{n}) \partial_{i} u_{n} T_{k}(u_{n})^{-} h_{m}'(u_{n}) dx - \sum_{i=1}^{N} \int_{\Omega} \phi_{i}(u_{n}) \partial_{i} T_{k}(u_{n})^{-} (1 - h_{m}(u_{n})) dx$$

+
$$\sum_{i=1}^{N} \int_{\Omega} \phi_{i}(u_{n}) \partial_{i} u_{n} T_{k}(u_{n})^{-} h_{m}'(u_{n}) dx \leq -\int_{\Omega} f_{n} T_{k}(u_{n})^{-} (1 - h_{m}(u_{n})) dx.$$
(36)

Thanks to (28) and (29), we have

$$\lim_{m \to +\infty} \lim_{n \to +\infty} \sum_{i=1}^{N} \int_{\Omega} a_i(x, u_n, \nabla u_n) \partial_i u_n T_k(u_n)^{-} h'_m(u_n) dx = 0.$$

Thus, the second integral in (36) converges to 0 when *n* and *m* goes to $+\infty$. As $\partial_i T_k(u_n)^- \rightharpoonup \partial_i T_k(u)^-$ in $L^{p_i}(\Omega, w_i)$ and $\phi_i(T_k(u_n))(1 - h_m(u_n)) \rightarrow \phi_i(T_k(u))(1 - h_m(u))$ strongly in $L^{p'_i}(\Omega, w_i^*)$, we get

$$\begin{split} \lim_{m \to +\infty} \lim_{n \to +\infty} \sum_{i=1}^{N} \int_{\Omega} \phi_i(u_n) \partial_i T_k(u_n)^- (1 - h_m(u_n)) dx \\ = \lim_{m \to +\infty} \sum_{i=1}^{N} \int_{\Omega} \phi_i(T_k(u)) \partial_i T_k(u)^- (1 - h_m(u)) dx. \end{split}$$

In view to Lebesgue's theorem, we get

$$\lim_{m\to+\infty}\sum_{i=1}^N\int_{\Omega}\phi_i(T_k(u))\partial_iT_k(u)^-(1-h_m(u))dx=0.$$

Hence, the third integral in (36) converges to 0 when *m* and *n* tends to $+\infty$.

We take $\Phi_i^n(t) = \int_0^t \phi_i(s) T_k(s)^- h'_m(s) ds$, in light of Green's Formula, we obtain

$$\sum_{i=1}^N \int_{\Omega} \phi_i^n(u_n) \partial_i u_n T_k(u_n)^- h'_m(u_n) dx = \sum_{i=1}^N \int_{\Omega} \partial_i \Phi_i^n(u_n) dx = 0.$$

Then the last integral of the left-hand side of (36) converges to 0 when *n* and *m* tend to $+\infty$. By using to Lebesgue dominated convergence theorem, we get the term of the right-hand side of (36) converges to 0 as *m* and *n* goes to $+\infty$. We Conclude

$$\lim_{m \to +\infty} \lim_{n \to +\infty} \sum_{i=1}^{N} \int_{\{u_n \le 0\}} a_i(x, u_n, \nabla u_n) \partial_i T_k(u_n) (1 - h_m(u_n)) dx = 0.$$
(37)

Following this, for η small enough, we choose $\varphi = u_n - \eta T_k(u_n^+ - \psi^+)(1 - h_m(u_n))$ as test function in (13), we obtain

$$\sum_{i=1}^{N} \int_{\Omega} a_{i}(x, u_{n}, \nabla u_{n}) \partial_{i} T_{k}(u_{n}^{+} - \psi^{+}) (1 - h_{m}(u_{n})) dx$$

$$- \sum_{i=1}^{N} \int_{\Omega} a_{i}(x, u_{n}, \nabla u_{n}) \partial_{i} u_{n} T_{k}(u_{n}^{+} - \psi^{+}) h'_{m}(u_{n}) dx$$

$$+ \sum_{i=1}^{N} \int_{\Omega} \phi_{i}^{n}(u_{n}) \partial_{i} T_{k}(u_{n}^{+} - \psi^{+}) (1 - h_{m}(u_{n})) dx$$

$$- \sum_{i=1}^{N} \int_{\Omega} \phi_{i}^{n}(u_{n}) \partial_{i} u_{n} T_{k}(u_{n}^{+} - \psi^{+}) h'_{m}(u_{n}) dx$$

$$\leq \int_{\Omega} f_{n} T_{k}(u_{n}^{+} - \psi^{+}) (1 - h_{m}(u_{n})) dx. \quad (38)$$

From the Hölder inequality, (8), (28) and (29), we get

$$\lim_{m \to +\infty} \lim_{n \to +\infty} \sum_{i=1}^{N} \int_{\Omega} \phi_i^n(u_n) \partial_i u_n T_k(u_n^+ - \psi^+) h'_m(u_n) dx = 0.$$

By the Young inequality, we obtain

$$\begin{split} &\sum_{i=1}^{N} \int_{\Omega} a_{i}(x, u_{n}, \nabla u_{n}) \partial_{i} T_{k}(u_{n}^{+} - \psi^{+}) (1 - h_{m}(u_{n})) dx \\ &\leq \sum_{i=1}^{N} \int_{\{-(m+1) \leq u_{n} \leq -m\}} a_{i}(x, u_{n}, \nabla u_{n}) \partial_{i} u_{n} T_{k}(u_{n}^{+} - \psi^{+}) dx \\ &+ \int_{\Omega} f_{n} T_{k}(u_{n}^{+} - \psi^{+}) (1 - h_{m}(u_{n})) dx \\ &+ \sum_{i=1}^{N} \int_{\{u_{n}^{+} - \psi^{+} \leq k\}} \phi_{i}^{n}(u_{n}) \partial_{i} u_{n}^{+} (1 - h_{m}(u_{n})) dx \\ &+ \sum_{i=1}^{N} \int_{\{u_{n}^{+} - \psi^{+} \leq k\}} \phi_{i}^{n}(u_{n}) \partial_{i} \psi^{+} (1 - h_{m}(u_{n})) dx \end{split}$$
(39)

Thank to (28), we get the first term on the right-hand converges to 0 when *n* and *m* tend to $+\infty$. By the Lebesgue dominated convergence theorem, we obtain the second part in the right-hand converges to 0 when *m* and *n* tend to $+\infty$.

As

$$\sum_{i=1}^{N} \int_{\{u_{n}^{+}-\psi^{+}\leq k\}} \phi_{i}^{n}(u_{n})\partial_{i}u_{n}^{+}(1-h_{m}(u_{n}))dx$$

$$= \sum_{i=1}^{N} \int_{\Omega} \phi_{i}^{n}(T_{\{k+\parallel\psi^{+}\parallel_{L^{\infty}(\mathcal{O})}\}}(u_{n}))\partial_{i}T_{\{k+\parallel\psi^{+}\parallel_{L^{\infty}(\mathcal{O})}\}}(u_{n}^{+})(1-h_{m}(u_{n}))dx.$$
(40)

Since $\partial_i T_{\{k+\|\psi^+\|_{L^{\infty}(\Omega)}\}}(u_n^+) \rightharpoonup \partial_i T_{\{k+\|\psi^+\|_{L^{\infty}(\Omega)}\}}(u^+)$ weakly in $L^{p_i}(\Omega, w_i)$ and $\phi_i^n(T_{\{k+\|\psi^+\|_{L^{\infty}(\Omega)}\}}(u_n))(1-h_m(u_n)) \rightarrow \phi_i(T_{\{k+\|\psi^+\|_{L^{\infty}(\Omega)}\}}(u))(1-h_m(u))$ strongly in $L^{p'_i}(\Omega, w_i^*)$, we obtain

$$\begin{split} &\sum_{i=1}^{N} \int_{\Omega} \phi_{i}^{n} (T_{\{k+\|\psi^{+}\|_{L^{\infty}(\Omega)}\}}(u_{n})) \partial_{i} T_{\{k+\|\psi^{+}\|_{L^{\infty}(\Omega)}\}}(u_{n}^{+})(1-h_{m}(u_{n})) dx \\ &= \sum_{i=1}^{N} \int_{\Omega} \phi_{i} (T_{\{k+\|\psi^{+}\|_{L^{\infty}(\Omega)}\}}(u)) \partial_{i} T_{\{k+\|\psi^{+}\|_{L^{\infty}(\Omega)}\}}(u)(1-h_{m}(u)) dx + \varepsilon(n). \end{split}$$

Using the Lebesgue dominated convergencetheorem, we obtain

$$\lim_{m\to\infty}\sum_{i=1}^N\int_{\Omega}\phi_i(T_{\{k+\|\psi^+\|_{L^{\infty}(\Omega)}\}}(u))\partial_iT_{\{k+\|\psi^+\|_{L^{\infty}(\Omega)}\}}(u)(1-h_m(u))dx=0.$$

Hence, we get the third integral converges to 0 as m and n tend to $+\infty$. Similarly as (37), we have

$$\lim_{m \to +\infty} \lim_{n \to +\infty} \sum_{i=1}^{N} \int_{\{u_n > 0\}} a_i(x, u_n, \nabla u_n) \partial_i T_k(u_n) (1 - h_m(u_n)) dx = 0.$$
(41)

According to (37) and (41), we obtain

$$\lim_{m \to +\infty} \lim_{n \to +\infty} \sum_{i=1}^{N} \int_{\Omega} a_i(x, u_n, \nabla u_n) \partial_i T_k(u_n) (1 - h_m(u_n)) dx = 0.$$
(42)

Furthermore, we have

$$\begin{split} &\sum_{i=1}^{N} \int_{\Omega} (a_i(x, T_k(u_n), \nabla T_k(u_n)) - a_i(x, T_k(u_n), \nabla T_k(u))) (\partial_i T_k(u_n) - \partial_i T_k(u)) dx \\ &= \sum_{i=1}^{N} \int_{\Omega} (a_i(x, T_k(u_n), \nabla T_k(u_n)) - a_i(x, T_k(u_n), \nabla T_k(u))) (\partial_i T_k(u_n) - \partial_i T_k(u)) h(u_n) dx \\ &+ \sum_{i=1}^{N} \int_{\Omega} (a_i(x, T_k(u_n), \nabla T_k(u_n))) \partial_i T_k(u_n) (1 - h_m(u_n)) dx \\ &- \sum_{i=1}^{N} \int_{\Omega} (a_i(x, T_k(u_n), \nabla T_k(u_n))) \partial_i T_k(u) (1 - h_m(u_n)) dx \end{split}$$

$$-\sum_{i=1}^{N}\int_{\Omega}(a_i(x,T_k(u_n),\nabla T_k(u)))(\partial_iT_k(u_n)-\partial_iT_k(u))(1-h_m(u_n))dx.$$

Combining (35) and (42), the first and the second integrals on the right-hand converge to 0 when *m* and *n* goes to ∞ .

As $(a_i(x, T_k(u_n), \nabla T_k(u_n)))_n$ is bounded in $L^{p'_i}(\Omega, w_i^*)$ and $\partial_i T_k(u)(1 - h_m(u_n)) \longrightarrow$ 0 in $L^{p_i}(\Omega, w_i)$ when m and n goes to $+\infty$, hence the third term on the right-hand side converge to 0 as *m* and *n* goes to $+\infty$.

Where

$$a_i(x, T_k(u_n), \nabla T_k(u_n))(1 - h_m(u_n)) \longrightarrow a_i(x, T_k(u), \nabla T_k(u))(1 - h_m(u))$$

strongly in $L^{p'_i}(\Omega, w_i^*)$ and $\partial_i T_k(u_n) \rightharpoonup \partial_i T(u)$ weakly in $L^{p_i}(\Omega, w_i)$, we get the last integral on the right-hand side converge to 0 as m and n goes to $+\infty$. Then, we obtain (27).

Thanks to (26), (27) and Lemma 3, we have

$$T_k(u_n) \to T(u)$$
 strongly in $W_0^{1,\overrightarrow{p}}(\Omega,\overrightarrow{w})$ and a. e. in Ω for all $k > 0$.

Step 4: Passing to the limit.

Let $\varphi \in K_{\Psi} \cap L^{\infty}(\Omega)$, we choose $v = u_n - T_k(u_n - \varphi)$ as test function in approximate problem (13), we have

$$\sum_{i=1}^{N} \int_{\Omega} a_i(x, u_n, \nabla u_n) \partial_i T_k(u_n - \varphi) dx + \sum_{i=1}^{N} \int_{\Omega} \phi_i^n(u_n) \partial_i T_k(u_n - \varphi) dx$$

$$\leq \int_{\Omega} f_n T_k(u_n - \varphi) dx,$$
(43)

which implies that,

$$\sum_{i=1}^{N} \int_{\Omega} a_i(x, T_{k+\|\varphi\|_{\infty}}(u_n), \nabla T_{k+\|\varphi\|_{\infty}}(u_n)) \partial_i T_k(u_n - \varphi) dx + \sum_{i=1}^{N} \int_{\Omega} \phi_i(T_{k+\|\varphi\|_{\infty}}(u_n)) \partial_i T_k(u_n - \varphi) dx \le \int_{\Omega} f_n T_k(u_n - \varphi) dx.$$

As $T_k(u_n) \to T(u)$ strongly in $W_0^{1, \overrightarrow{p}}(\Omega, \overrightarrow{w})$ and a.e. in Ω for all k > 0, we obtain

$$a_{i}(x, T_{k+\|\varphi\|_{\infty}}(u_{n}), \nabla T_{k+\|\varphi\|_{\infty}}(u_{n})) \rightharpoonup a_{i}(x, T_{k+\|\varphi\|_{\infty}}(u), \nabla T_{k+\|\varphi\|_{\infty}}(u))$$
weakly in $L^{p'_{i}}(\Omega, w_{i}^{*})$
$$\phi_{i}(T_{k+\|\varphi\|_{\infty}}(u_{n})) \rightarrow \phi_{i}(T_{k+\|\varphi\|_{\infty}}(u))$$
strongly in $L^{p'_{i}}(\Omega, w_{i}^{*})$

$$\phi_i(T_{k+\|\varphi\|_{\infty}}(u_n)) \to \phi_i(T_{k+\|\varphi\|_{\infty}}(u)) \quad \text{strongly in } L^{p_i'}(\Omega, w_i^*)$$

and

$$\partial_i T_k(u_n - \varphi) \to \partial_i T_k(u - \varphi)$$
 strongly in $L^{p_i}(\Omega, w_i)$.

Passing to the limit in (43) and this completes the proof of theorem 1.

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Well-Posedness and Stability for the Viscous Primitive Equations of Geophysics in Critical Fourier-Besov-Morrey Spaces

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Abstract. In this paper we study the Cauchy problem of the viscous primitive equations of geophysics in critical Fourier-Besov-Morrey spaces. By using the Fourier localization argument and the Littlewood-Paley theory, we prove that the Cauchy problem with Prankster number P = 1 is local well-posedness and global well-posedness when the initial data (u_0, θ_0) are small and we give a stability result for global solutions.

Keywords: Navier-Stokes equations \cdot Global well-posedness \cdot Analytic solutions \cdot Coriolis force \cdot Fourier-Besov-Morrey space

1 Introduction

In this paper, we study the initial value problem of the viscous primitive equations of geophysics in \mathbb{R}^3 , which is a fundamental mathematical model in the field of fluid geophysics. The model reads as follows:

$$\begin{cases} \partial_t u + \nu \Delta u + \Omega e_3 \times u + (u.\nabla)u + \nabla p = g\theta e_3 & (t,x) \in \mathbb{R}^+ \times \mathbb{R}^3, \\ \partial_t \theta + \mu \Delta \theta + (u.\nabla)\theta = -\mathcal{N}^2 u_3 & (t,x) \in \mathbb{R}^+ \times \mathbb{R}^3, \\ \nabla .u = 0, & (1.1) \\ u(0,x) = u_0(x) & x \in \mathbb{R}^3, \end{cases}$$

where $u = u(t, x) = (u^1(t, x), u^2(t, x), u^3(t, x))$ and p = p(t, x) denotes the unknown velocity field and the unknown pressure of the fluid at the point $(t, x) \in \mathbb{R}^+ \times \mathbb{R}^3$, respectively and θ is a scalar function representing the density fluctuation in the fluid (in the case of the ocean it depends on the temperature and the salinity, and in the case of the atmosphere it depends on the temperature), while $u_0 = u_0(x) = (u_0^1(x), u_0^2(x), u_0^3(x))$ denote the given initial velocity flied satisfying the compatibility condition $\nabla u = 0$. ν , μ and g are positive

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Z. Hammouch et al. (Eds.): SM2A 2019, LNNS 168, pp. 123–140, 2021. https://doi.org/10.1007/978-3-030-62299-2_9 constants related to viscosity, diffusivity and gravity, respectively, $\Omega \in \mathbb{R}$ represents the speed of rotation around the vertical unit vector $e_3 = (0, 0, 1)$, which is called the Coriolis parameter, and "×" represents the outer product, hence, $-\Omega e_3 \times u = (\Omega u_2, -\Omega u_1, 0)$. We recall that the Coriolis term has an another expression $-\Omega e_3 \times u = -\Omega J u$, where the skew-symmetric matrix J defined by

$$\mathbf{J} = \begin{pmatrix} 0 & -1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}.$$

 \mathcal{N} is the stratification parameter, a nonnegative constant representing the Brunt-Visala wave frequency. The ratio $P := \frac{\nu}{\mu}$ is known as the Prandtl number and $B := \frac{\Omega}{\mathcal{N}}$ is essentially the "Burger" number of geophysics. When $\theta \equiv 0$, $\mathcal{N} = 0$ and $\Omega = 0$, the problem (1.1) become the classical

When $\theta \equiv 0$, $\mathcal{N} = 0$ and $\Omega = 0$, the problem (1.1) become the classical Navier-Stokes equation:

$$\begin{cases} u_t - \nu \Delta u + (u \cdot \nabla) u + \nabla p = 0 \quad (t, x) \in \mathbb{R}^+ \times \mathbb{R}^3, \\ \nabla \cdot u = 0, \\ u(0, x) = u_0(x) \quad x \in \mathbb{R}^3. \end{cases}$$

The existence of mild solutions and the regularity have been established locally in time and global for small initial data in various functional spaces, for example [5–8,28,29,34,36].

If only $\theta \equiv 0$, $\mathcal{N} = 0$ but $\Omega \neq 0$ the problem (1.1) corresponds to the usual Navier-Stokes equation with Coriolis force,

$$\begin{cases} u_t - \nu \Delta u + \Omega e_3 \times u + (u \cdot \nabla) u + \nabla p = 0 \quad (t, x) \in \mathbb{R}^+ \times \mathbb{R}^3, \\ \nabla \cdot u = 0, \\ u(0, x) = u_0(x) \quad x \in \mathbb{R}^3. \end{cases}$$

Hieber and Shibata [22] obtained the uniform global well-posedness for the Navier-Stokes equations with Coriolis force for small initial data in the Sobolev space $H^{\frac{1}{2}}(\mathbb{R}^3)$. Iwabuchi and Takada [26] proved the existence of global solutions for the Navier-Stokes equations with Coriolis force in Sobolev spaces $\dot{H}^s(\mathbb{R}^3)$ with 1/2 < s < 3/4 if the speed of rotation Ω is large enough compared with the norm of initial data $||u_0||_{\dot{H}^s}$, they also obtained the global existence and the uniqueness of the mild solution for small initial data in the Fourier-Besov spaces $\dot{\mathrm{FB}}_{1,2}^{-1}$ and proved the ill-posedness in the space $\dot{\mathrm{FB}}_{1,q}^{-1}$, $2 < q \leq \infty$ for all $\Omega \in \mathbb{R}$ see [27]. El Baraka and Toumllin [10] got global well posedness result with small initial data in $\mathcal{FN}_{p,\lambda,q}^{1-2\alpha+\frac{3}{p'}+\frac{\lambda}{p}}$ for $\alpha \neq 1$ and $\Omega = 0$, moreover, in [12] they generalize this result for $\alpha \neq 1$ and $\Omega \neq 0$ where they proved local well-posedness results and global well-posedness results with small initial data in fourier-Besov-Morrey spaces.

When $\theta \neq 0$, $\mathcal{N} \neq 0$ and $\Omega \neq 0$, Babin, Maholov and Nicolaenko [3] proved the existance of global solution for problem (1.1) in $[H^s(\mathbb{T}^3)]^4$ with $s \geq 3/4$ for small initial data when the stratification parameter \mathcal{N} is sufficiently large. Charve [18,19] obtained the global well-posedness of problem (1.1) in $[\dot{H}^{\frac{1}{2}}(\mathbb{R}^3) \cap$ $\dot{H}^{1}(\mathbb{R}^{3})]^{4}$ under the assumptions that both Ω and \mathcal{N} are sufficiently large for arbitrary initial data, moreover we get the global well-posedness of (1.1) in less regular initial value spaces. [24] J.Sun and S.Cui proved that the Cauchy problem (1.1) with P = 1 is locally well-posed and globally well-posed when the initial data (u_0, θ_0) are small in Fourier-Bessov spaces $FB_{p,r}^{2-\frac{3}{p}}$ for 1 , $<math>1 \leq r < \infty$ and $FB_{1,r}^{-1}$ for $1 \leq r \leq 2$, they also proved that such problem is ill-posed in $FB_{1,r}^{-1}$ for $2 < r \leq \infty$.

We refer to [14,23,25,32] for rich literature about global-in-time well-posedness for fluid dynamics PDEs.

We first transform the Cauchy problem in to an equivalent Cauchy probleme. By setting $N := \mathcal{N}\sqrt{g}, v := (v^1, v^2, v^3, v^4) := (u^1, u^2, u^3, \frac{\sqrt{g}\theta}{\mathcal{N}}), v_0 := (v_0^1, v_0^2, v_0^3, v_0^4) := (u_0^1, u_0^2, u_0^3, \frac{\sqrt{g}\theta_0}{\mathcal{N}})$ and $\tilde{\nabla} := (\partial_1, \partial_2, \partial_3, 0), (1.1)$ can be rewritten into the following problem:

$$\begin{cases} v_t + \mathcal{A}v + \mathcal{B}v + \tilde{\nabla}p = -(v.\tilde{\nabla})v & (t,x) \in \mathbb{R}^+ \times \mathbb{R}^3, \\ \tilde{\nabla}v = 0, & \\ v(0,x) = v_0(x) & x \in \mathbb{R}^3. \end{cases}$$
(1.2)

Where

$$\mathcal{A} = \begin{pmatrix} -\nu \bigtriangleup & 0 & 0 & 0 \\ 0 & -\nu \bigtriangleup & 0 & 0 \\ 0 & 0 & -\nu \bigtriangleup & 0 \\ 0 & 0 & 0 & -\mu \bigtriangleup \end{pmatrix} \text{ and } \mathcal{B} = \begin{pmatrix} 0 & -\Omega & 0 & 0 \\ \Omega & 0 & 0 & 0 \\ 0 & 0 & 0 & -N \\ 0 & 0 & N & 0 \end{pmatrix}.$$

To solve the original problem (1.1), we may consider the following integral equation:

$$v(t) = T_{\Omega,N}(t)v_0 - \int_0^t T_{\Omega,N}(t-\tau)\tilde{\mathbb{P}}\tilde{\nabla} \cdot (v \otimes v)d\tau, \qquad (1.3)$$

where, $\tilde{\mathbb{P}} = (\tilde{\mathbb{P}}_{ij})_{4 \times 4}$ the Helmholtz projection onto the divergence-free vector fields defined by:

$$\tilde{\mathbb{P}}_{ij} = \begin{cases} \delta_{ij} + R_i R_j & 1 \le i, j \le 3\\ \delta_{ij} & \text{otherwise,} \end{cases}$$

and $T_{\Omega,N}(.)$ denotes Stokes-Coriolis Stratification to the linear probleme of (1.2) via Fourier transform, which is given explicitly by

$$T_{\Omega,N}(t)f = \mathcal{F}^{-1}[\cos(\frac{|\xi|'}{|\xi|}t)M_1 + \sin(\frac{|\xi|'}{|\xi|}t)M_2 + M_3] * (e^{-\nu\Delta t}f).$$

Where

$$|\xi| := \sqrt{\xi_1^2 + \xi_2^2 + \xi_3^2}$$
 and $|\xi|' := |\xi|'_{\Omega,N} := \sqrt{N^2 \xi_1^2 + N^2 \xi_2^2 + \Omega^2 \xi_3^2}$

and

$$M_{1} = \begin{pmatrix} \frac{\Omega^{2}\xi_{3}^{2}}{|\xi|^{\prime 2}} & 0 & -\frac{N^{2}\xi_{1}\xi_{3}}{|\xi|^{\prime 2}} & \frac{\Omega N\xi_{2}\xi_{3}}{|\xi|^{\prime 2}} \\ 0 & \frac{\Omega^{2}\xi_{3}^{2}}{|\xi|^{\prime 2}} & -\frac{N^{2}\xi_{2}\xi_{3}}{|\xi|^{\prime 2}} & -\frac{\Omega N\xi_{1}\xi_{3}}{|\xi|^{\prime 2}} \\ -\frac{\Omega^{2}\xi_{1}\xi_{3}}{|\xi|^{\prime 2}} & -\frac{\Omega^{2}\xi_{2}\xi_{3}}{|\xi|^{\prime 2}} & \frac{N^{2}(\xi_{1}^{2}+\xi_{2}^{2})}{|\xi|^{\prime 2}} & 0 \\ \frac{\Omega N\xi_{2}\xi_{3}}{|\xi|^{\prime 2}} & -\frac{\Omega N\xi_{1}\xi_{3}}{|\xi||\xi|^{\prime}} & 0 & \frac{N^{2}(\xi_{1}^{2}+\xi_{2}^{2})}{|\xi|^{\prime 2}} \end{pmatrix}.$$

$$M_{2} = \begin{pmatrix} 0 & -\frac{\Omega\xi_{3}^{2}}{|\xi|^{\prime 2}} & -\frac{\Omega\xi_{2}\xi_{3}}{|\xi||\xi|^{\prime}} & \frac{N\xi_{1}\xi_{3}}{|\xi||\xi|^{\prime}} \\ \frac{\Omega\xi_{3}^{2}}{|\xi||\xi|^{\prime}} & 0 & -\frac{\Omega\xi_{1}\xi_{3}}{|\xi||\xi|^{\prime}} & \frac{N\xi_{2}\xi_{3}}{|\xi||\xi|^{\prime}} \\ -\frac{\Omega^{2}\xi_{2}\xi_{3}}{|\xi||\xi|^{\prime}} & \frac{\Omega^{2}\xi_{1}\xi_{3}}{|\xi||\xi|^{\prime}} & 0 & -\frac{N(\xi_{1}^{2}+\xi_{3}^{2})}{|\xi||\xi|^{\prime}} \\ -\frac{N\xi_{1}\xi_{3}}{|\xi||\xi|^{\prime}} & -\frac{N\xi_{2}\xi_{3}}{|\xi||\xi|^{\prime}} & \frac{N(\xi_{1}^{2}+\xi_{3}^{2})}{|\xi||\xi|^{\prime}} & 0 \end{pmatrix}.$$

$$M_{3} = \begin{pmatrix} \frac{N^{2}\xi_{2}}{|\xi|^{\prime 2}} & -\frac{N^{2}\xi_{1}\xi_{3}}{|\xi||\xi|^{\prime}} & 0 & -\frac{\Omega N\xi_{1}\xi_{2}}{|\xi|^{\prime 2}} \\ -\frac{N^{2}\xi_{1}\xi_{2}}{|\xi|^{\prime 2}} & -\frac{N^{2}\xi_{1}\xi_{3}}{|\xi||\xi|^{\prime}} & 0 & -\frac{\Omega N\xi_{1}\xi_{2}}{|\xi|^{\prime 2}} \\ 0 & 0 & 0 & 0 \\ -\frac{\Omega N\xi_{2}\xi_{3}}{|\xi|^{\prime 2}} & \frac{\Omega N\xi_{1}\xi_{3}}{|\xi|^{\prime 2}} & 0 & \frac{\Omega^{2}\xi_{3}}{|\xi|^{\prime 2}} \\ 0 & 0 & 0 & 0 \\ -\frac{\Omega N\xi_{2}\xi_{3}}{|\xi|^{\prime 2}} & \frac{\Omega N\xi_{1}\xi_{3}}{|\xi|^{\prime 2}} & 0 & \frac{\Omega^{2}\xi_{3}}{|\xi|^{\prime 2}} \end{pmatrix}.$$

Note that, denoting by M_{jk}^l -th component of the matrix $M_l(\xi)$, it is obvious that non-vanishing M_{jk}^l satisfies

$$|M_{ik}^l| \le 2$$
 for $\xi \in \mathbb{R}^3, j, k = 1, 2, 3, 4, l = 1, 2, 3$

Inspired by the work [10, 24], the aim of this paper is to prove the global existence and the stability of the global solution of the viscous primitive equations of geophysics in critical Fourier-Besov-Morrey spaces, using abstract lemma on the existence of fixed point solutions.

Lemma 1.1. Let X be a Banach space with norm $\|.\|_X$ and $B: X \times X \longrightarrow X$ be a bounded bilinear operator satisfying

$$||B(u,v)||_X \le \eta ||u||_X ||v||_X$$

for all $u, v \in X$ and a constant $\eta > 0$. Then, if $0 < \varepsilon < \frac{1}{4\eta}$ and if $y \in X$ such that $\|y\|_X \leq \varepsilon$, the equation x := y + B(x, x) has a solution \overline{x} in X such that $\|\overline{x}\|_X \leq 2\varepsilon$. This solution is the only one in the ball $\overline{B}(0, 2\varepsilon)$. Moreover, the solution depends continuously on y in the sense: if $\|y'\|_X < \varepsilon$, x' = y' + B(x', x'), and $\|x'\|_X \leq 2\varepsilon$, then

$$\|\overline{x} - x'\|_X \le \frac{1}{1 - 4\varepsilon\eta} \|y - y'\|_X.$$

2 Preliminaries and Main Results

To give the precise statements of our main results, we first recall the definitions of the Morrey space $M_p^{\lambda}(\mathbb{R}^n)$, Besov space $B_{p,q}^s(\mathbb{R}^n)$ and Fourier-Besov-Morrey space $\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^s(\mathbb{R}^n)$ were introduced by Ferreira and Lima [16] in order to analyze a class of active scalar equations. As usual we denote by the space of Schwartz functions on \mathbb{R}^3 , and by the space of tempred distributions on \mathbb{R}^3 . Choose two nonnegative smooth radial functions χ , φ satisfying

$$\begin{aligned} \sup \varphi &\subset \{\xi \in \mathbb{R}^n : \frac{3}{4} \le |\xi| \le \frac{8}{3}\}, \quad \sum_{j \in \mathbb{Z}} \varphi(2^{-j}\xi) = 1, \quad \xi \in \mathbb{R}^n \setminus \{0\}, \\ \sup \chi &\subset \{\xi \in \mathbb{R}^n : |\xi| \le \frac{4}{3}\}, \quad \chi(\xi) + \sum_{j \ge 0} \varphi(2^{-j}\xi) = 1, \quad \xi \in \mathbb{R}^n. \end{aligned}$$

We denote $\varphi_j(\xi) = \varphi(2^{-j}\xi)$ and \mathcal{P} the set of all polynomials. The space of tempered distributions is denoted by S'. The homogeneous dyadic blocks $\dot{\Delta}_j$ and \dot{S}_j are defined for all $j \in \mathbb{Z}$ by

$$\dot{\Delta}_j u = \varphi(2^{-j}D)u = 2^{jn} \int h(2^j y)u(x-y) \, dy,$$
$$\dot{S}_j u = \sum_{k \le j-1} \dot{\Delta}_k u = \chi(2^{-j}D)u = 2^{jn} \int \tilde{h}(2^j y)u(x-y) \, dy,$$

where $h = \mathcal{F}^{-1}\varphi$ and $\tilde{h} = \mathcal{F}^{-1}\chi$. We defined the function spaces $\mathrm{M}_{p}^{\lambda}(\mathbb{R}^{n})$.

Definition 2.1. [28,34]. For $1 \le p < \infty$, $0 \le \lambda < n$, the Morrey spaces $M_p^{\lambda} = M_p^{\lambda}(\mathbb{R}^n)$ is defined by

$$\mathbf{M}_{p}^{\lambda}(\mathbb{R}^{n}) = \{ f \in L_{loc}^{p}(\mathbb{R}^{n}); \|f\|_{\mathbf{M}_{p}^{\lambda}} < \infty \},\$$

where

$$\|f\|_{\mathcal{M}_{p}^{\lambda}} = \sup_{x_{0} \in \mathbb{R}^{n}} \sup_{r > 0} r^{-\frac{\lambda}{p}} \|f\|_{L^{p}(B(x_{0}, r))},$$

with $B(x_0, r)$ the ball in \mathbb{R}^n with center x_0 and radius r. The space M_p^{λ} endowed with the norm $||f||_{M_p^{\lambda}}$ is a Banach space. If $1 \leq p_1, p_2, p_3 < \infty$ and $0 \leq \lambda_1, \lambda_2, \lambda_3 < n$ with $\frac{1}{p_3} = \frac{1}{p_1} + \frac{1}{p_2}$ and $\frac{\lambda_3}{p_3} = \frac{\lambda_1}{p_1} + \frac{\lambda_2}{p_2}$, then we have the Hölder inequality

$$\|fg\|_{\mathcal{M}_{p_3}^{\lambda_3}} \le \|f\|_{\mathcal{M}_{p_1}^{\lambda_1}} \|g\|_{\mathcal{M}_{p_2}^{\lambda_2}}.$$

Also, for $1 \le p < \infty$ and $0 \le \lambda < n$,

$$\|\varphi * g\|_{\mathbf{M}_{p}^{\lambda}} \le \|\varphi\|_{L^{1}} \|g\|_{\mathbf{M}_{p}^{\lambda}}, \tag{2.1}$$

for all $\varphi \in L^1$ and $g \in \mathcal{M}_p^{\lambda}$.

Bernstein type lemma in Fourier variables in Morrey spaces.

Lemma 2.2. [16]. Let $1 \le q \le p < \infty$, $0 \le \lambda_1, \lambda_2 < n$, $\frac{n-\lambda_1}{p} \le \frac{n-\lambda_2}{q}$ and let γ be a multi-index. If $supp(\widehat{f}) \subset \{|\xi| \le A2^j\}$, then there is a constant C > 0 independent of f and j such that

$$\|(i\xi)^{\gamma}\widehat{f}\|_{\mathcal{M}^{\lambda_{2}}_{q}} \leq C2^{j|\gamma|+j(\frac{n-\lambda_{2}}{q}-\frac{n-\lambda_{1}}{p})}\|\widehat{f}\|_{\mathcal{M}^{\lambda_{1}}_{p}}.$$
(2.2)

Then, we define the function spaces $\mathcal{F}\dot{\mathcal{N}}^{s}_{p,\lambda,q}(\mathbb{R}^{n})$.

Definition 2.3. (Homogeneous Besov-Morrey spaces) Let $s \in \mathbb{R}$, $1 \le p < \infty$, $1 \le q \le \infty$ and $0 \le \lambda < n$, the space $\dot{\mathcal{N}}^s_{p,\lambda,q}(\mathbb{R}^n)$ is defined by

$$\dot{\mathcal{N}}^{s}_{p,\lambda,q}(\mathbb{R}^{n}) = \left\{ u \in \mathcal{Z}'(\mathbb{R}^{n}); \quad \|u\|_{\dot{\mathcal{N}}^{s}_{p,\lambda,q}(\mathbb{R}^{n})} < \infty \right\}.$$

Here

$$\|u\|_{\dot{\mathcal{N}}^{s}_{p,\lambda,q}(\mathbb{R}^{n})} = \begin{cases} \left\{ \sum_{j \in \mathbb{Z}} 2^{jqs} \|\dot{\Delta}_{j}u\|_{\mathcal{M}^{\lambda}_{p}}^{q} \right\}^{1/q} & for \ q < \infty, \\ \sup_{j \in \mathbb{Z}} 2^{js} \|\dot{\Delta}_{j}u\|_{\mathcal{M}^{\lambda}_{p}} & for \ q = \infty. \end{cases}$$

The space $\mathcal{Z}'(\mathbb{R}^n)$ denotes the topological dual of the space $\mathcal{Z}(\mathbb{R}^n) = \{f \in \mathcal{S}(\mathbb{R}^n); \partial^{\alpha} \widehat{f}(0) = 0 \text{ for every multi-index } \alpha \}$ and can be identified to the quotient space $\mathcal{S}'(\mathbb{R}^n)/\mathcal{P}$, where \mathcal{P} represents the set of all polynomials on \mathbb{R}^n . We refer to [37, chap. 8] and [15] for more details.

Definition 2.4. (Homogeneous Fourier-Besov-Morrey spaces)

Let $s \in \mathbb{R}$, $0 \leq \lambda < n$, $1 \leq p < \infty$ and $1 \leq q \leq \infty$. The space $\mathcal{F}\dot{\mathcal{N}}^{s}_{p,\lambda,q}(\mathbb{R}^{n})$ denotes the set of all $u \in \mathcal{Z}'(\mathbb{R}^{n})$ such that

$$\|u\|_{\mathcal{F}\dot{\mathcal{N}}^{s}_{p,\lambda,q}(\mathbb{R}^{n})} = \left\{ \sum_{j\in\mathbb{Z}} 2^{jqs} \|\widehat{\dot{\Delta}_{j}u}\|_{\mathcal{M}^{\lambda}_{p}}^{q} \right\}^{1/q} < \infty,$$
(2.3)

with appropriate modifications made when $q = \infty$.

Note that the space $\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{s}(\mathbb{R}^{n})$ equipped with the norm (2.3) is a Banach space. Since $\mathcal{M}_{p}^{0} = L^{p}$, we have $\mathcal{F}\dot{\mathcal{N}}_{p,0,q}^{s} = F\dot{B}_{p,q}^{s}$, $\mathcal{F}\dot{\mathcal{N}}_{1,0,q}^{s} = F\dot{B}_{1,q}^{s} = \dot{\mathcal{B}}_{q}^{s}$ and $\mathcal{F}\dot{\mathcal{N}}_{1,0,1}^{-1} = \chi^{-1}$ where $\dot{\mathcal{B}}_{q}^{s}$ is the Fourier-Herz space and χ^{-1} is the Lei-Lin space [4,11].

Now, we give the definition of the mixed space-time spaces.

Definition 2.5. Let $s \in \mathbb{R}$, $1 \leq p < \infty$, $1 \leq q, \rho \leq \infty$, $0 \leq \lambda < n$, and $I = [0,T), T \in (0,\infty]$. The space-time norm is defined on u(t,x) by

$$\|u(t,x)\|_{\mathcal{L}^{\rho}(I,\mathcal{F}\dot{\mathcal{N}}^{s}_{p,\lambda,q})} = \Big\{\sum_{j\in\mathbb{Z}} 2^{jqs} \|\widehat{\dot{\Delta}_{j}u}\|_{L^{\rho}(I,\mathcal{M}^{\lambda}_{p})}^{q}\Big\}^{1/q}$$

and denote by $\mathcal{L}^{\rho}(I, \mathcal{F}\dot{\mathcal{N}}^{s}_{p,\lambda,q})$ the set of distributions in $S'(\mathbb{R} \times \mathbb{R}^{n})/\mathcal{P}$ with finite $\|.\|_{\mathcal{L}^{\rho}(I,\mathcal{F}\dot{\mathcal{N}}^{s}_{p,\lambda,q})}$ norm.

Theorem 2.6. Let Prandtl number P = 1, i.e., $\mu = \nu$, $\Omega \in \mathbb{R}$, $0 \le \lambda < 3$, $1 \le q \le 2$. For $\max\{1, \frac{3-\lambda}{2}\} \le p < \infty$, there exists a positive time T such that for $v_0 = (u_0, \theta_0) \in \mathcal{FN}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}$ and $\nabla .u_0 = 0$, the problem (1.1) admits a unique local solution $(u, \theta) \in \mathcal{L}^4([0, T), \mathcal{FN}_{p,\lambda,q}^{-\frac{1}{2}+\frac{3}{p'}+\frac{\lambda}{p}})$.

Furthermore $1 \le p < \infty$ there exists a constant $C_0(p,q)$ such that for any $v_0 = (u_0, \theta_0) \in \mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}$ satisfying $\nabla .u_0 = 0$ and $\|(u_0, \frac{\sqrt{g}\theta_0}{\mathcal{N}})\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}} < C_0\mu$, the problem (1.1) admits a unique global solution

 $(u,\theta) \in \mathcal{L}^{\infty}\left([0,\infty); \mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}\right) \cap \mathcal{L}^{1}\left([0,\infty), \mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{1+\frac{3}{p'}+\frac{\lambda}{p}}\right),$

and it satisfies

$$\begin{aligned} \|(u,\frac{\sqrt{g}\theta}{\mathcal{N}})\|_{\mathcal{L}^{\infty}([0,\infty);\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}})} + \mu\|(u,\frac{\sqrt{g}\theta}{\mathcal{N}})\|_{\mathcal{L}^{1}([0,\infty),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{1+\frac{3}{p'}+\frac{\lambda}{p}})} \\ &\leq 2C\|(u_{0},\frac{\sqrt{g}\theta_{0}}{\mathcal{N}})\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}}, \end{aligned}$$

where C is a positive constant.

Theorem 2.7. Let T^* denote the maximal time of existence of a solution $v = (u, \theta)$ in $\mathcal{L}^{\infty}\left([0, T^*); \mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}\right) \cap \mathcal{L}^1\left([0, T^*), \mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{1+\frac{3}{p'}+\frac{\lambda}{p}}\right)$. If $T^* < \infty$, then $\|v\|_{\mathcal{L}^1([0, T^*), \mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{1+\frac{3}{p'}+\frac{\lambda}{p}})} = \infty.$

Besides; if $v' = (u', \theta') \in C(\mathbb{R}^+, \mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}})$ is a global solution of (1.1), and for all $v'_0 \in \mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}$ satisfying

$$\|v_0' - v_0\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}} < C_0 \frac{\mu}{8} \exp\left\{-\int_0^\infty \frac{1}{C_0} (||\mathcal{B}|| + \|v\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{1+\frac{3}{p'}+\frac{\lambda}{p}}})\right\}$$
(2.4)

for some constant C_0 sufficiently small and $||\mathcal{B}||$ is matrix norm, then the viscous primitive equations starting from v_0 has a global solution v fulfilling the inequality

$$\begin{aligned} \|v'(t) - v(t)\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}} + \frac{\mu}{2} \|v'(s) - v(s)\|_{\mathcal{L}^{1}([0,t),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{1+\frac{3}{p'}+\frac{\lambda}{p}})} \\ < C \|v'_{0} - v_{0}\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}} \exp\Big\{\int_{0}^{\infty} C(||\mathcal{B}|| + \|v\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{1+\frac{3}{p'}+\frac{\lambda}{p}}})\Big\} \end{aligned}$$

where C is a positive constant.

3 Well-Posedness

In this section we present the proof of Theorem 2.6. To this end, we establish some basic estimates.

Lemma 3.1. Let T > 0, $s \in \mathbb{R}$, $0 \le \lambda < 3, 1 \le p < \infty, 1 \le q, \rho, r \le \infty$ and $f \in \mathcal{L}^r([0,T), \mathcal{F}\dot{\mathcal{N}}^s_{p,\lambda,q})$. There exists a constant C > 0 such that

$$\|\int_0^t T_{\Omega,N}(t-\tau)f(\tau)d\tau\|_{\mathcal{L}^\rho([0,T),\mathcal{F}\dot{\mathcal{N}}^s_{p,\lambda,q})} \le C\|f\|_{\mathcal{L}^r\left([0,T),\mathcal{F}\dot{\mathcal{N}}^{s-2-\frac{2}{\rho}+\frac{2}{r}}_{p,\lambda,q}\right)}$$

Proof: Set $1 + \frac{1}{\rho} = \frac{1}{\bar{\rho}} + \frac{1}{r}$. The definition of the space-time norm of $\mathcal{L}^{\rho}([0,T), \mathcal{F}\dot{\mathcal{N}}^{s}_{p,\lambda,q})$ and Young's inequality give

$$\begin{split} & \left\| \int_{0}^{t} T_{\Omega,N}(t-\tau) f(\tau) d\tau \right\|_{\mathcal{L}^{\rho}([0,T),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{s})} \\ &= \Big\{ \sum_{j \in \mathbb{Z}} 2^{jqs} \Big(\int_{0}^{T} \|\varphi_{j} \int_{0}^{t} \mathcal{F}(T_{\Omega,N}(t-\tau)f)(\tau) d\tau \|_{\mathcal{M}_{p}^{\lambda}}^{\rho} dt \Big)^{\frac{q}{\rho}} \Big\}^{1/q} \\ &\leq C \Big\{ \sum_{j \in \mathbb{Z}} 2^{jqs} \Big(\int_{0}^{T} \|\varphi_{j} \int_{0}^{t} e^{-\mu |\xi|^{2}(t-\tau)} \hat{f}(\tau) d\tau \|_{\mathcal{M}_{p}^{\lambda}}^{\rho} dt \Big)^{\frac{q}{\rho}} \Big\}^{1/q} \\ &\leq C \Big\{ \sum_{j \in \mathbb{Z}} 2^{jqs} \Big(\int_{0}^{T} \|\varphi_{j} \int_{0}^{t} e^{-\mu 2^{2j}(t-\tau)} \hat{f}(\tau) d\tau \|_{\mathcal{M}_{p}^{\lambda}}^{\rho} dt \Big)^{\frac{q}{\rho}} \Big\}^{1/q} \\ &\leq C \Big\{ \sum_{j \in \mathbb{Z}} 2^{jqs} \Big(\int_{0}^{T} e^{-t\mu \tilde{\rho} 2^{2j}} dt \Big)^{\frac{q}{\rho}} \|\varphi_{j} \hat{f}(\tau)\|_{L^{r}([0,T),\mathcal{M}_{p}^{\lambda})}^{q} \Big\}^{1/q} \\ &\leq C \Big\{ \sum_{j \in \mathbb{Z}} 2^{jq(s-2-\frac{2}{\rho}+\frac{2}{\tau})} \|\varphi_{j} \hat{f}(\tau)\|_{L^{r}([0,T),\mathcal{M}_{p}^{\lambda})}^{q} \Big\}^{1/q} \\ &\leq C \|f\|_{\mathcal{L}^{r}([0,T),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{s-2-\frac{2}{\rho}+\frac{2}{\tau})} \,. \end{split}$$

Lemma 3.2. Let T > 0, $0 \le \lambda < 3$, $1 \le p < \infty$, $1 \le q \le \infty$, $s \in \mathbb{R}$ and $u_0 \in \mathcal{FN}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}(\mathbb{R}^3)$. Then there exists a constant C > 0 such that

$$\|T_{\Omega,N}(.)v_0\|_{\mathcal{L}^{\rho}\left([0,T),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{s+\frac{2}{\rho}}\right)} \leq C\|v_0\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{s}},\tag{3.1}$$

Proof: Since Supp $\psi_j \subset \{\xi \in \mathbb{R}^3 : 2^{j-1} \le |\xi| \le 2^{j+1}\}$, one has

$$||\Delta_j \widehat{T_{\Omega,N}(.)} v_0||_{M_p^{\lambda}} \le C e^{-\mu 2^{2j} t} ||\widehat{\psi_j} \widehat{v_0}||_{M_p^{\lambda}}.$$

for all $t \ge 0$, which yields that

$$||\Delta_j \widehat{T_{\Omega,N}(.)} v_0||_{L^{\rho}([0,T),M_p^{\lambda})} \le C \Big(\frac{1-e^{-\mu 2^{2j}\rho T}}{\mu 2^{2j}\rho}\Big)^{\frac{1}{\rho}} ||\widehat{\psi}_j \widehat{v}_0||_{M_p^{\lambda}}.$$

Thus, we have

$$\|T_{\Omega,N}(.)v_0\|_{\mathcal{L}^{\rho}\left([0,T),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{s+\frac{2}{\rho}}\right)} \leq C\|v_0\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{s}}$$

Proposition 3.3. [10] Let $1 \le p < \infty$, $1 \le q \le 2$, $\frac{1}{2} \le \alpha \le 1 + \frac{3}{2p'} + \frac{\lambda}{2p}$ and $0 \le \lambda < 3$. Set

$$X = \mathcal{L}^{\infty}\left([0,\infty), \mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{1-2\alpha+\frac{3}{p'}+\frac{\lambda}{p}}\right) \cap \mathcal{L}^{1}\left([0,\infty), \mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{1-2\alpha+\frac{3}{p'}+\frac{2\alpha}{\rho}+\frac{\lambda}{p}}\right).$$

with the norm

$$\|u\|_{X} = \|u\|_{\mathcal{L}^{\infty}\left([0,\infty),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{1-2\alpha+\frac{3}{p'}+\frac{\lambda}{p}}\right)} + \mu\|u\|_{\mathcal{L}^{1}\left([0,\infty),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{1-2\alpha+\frac{3}{p'}+\frac{2\alpha}{\rho}+\frac{\lambda}{p}}\right)}.$$

There exists a constant C = C(p,q) > 0 depending on α, p, q such that

$$\|\nabla (u \otimes v)\|_{\mathcal{L}^1\left([0,\infty),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{1-4\alpha+\frac{3}{p'}+\frac{2\alpha}{\rho}+\frac{\lambda}{p}\right)} \leq C\mu^{-1}\|u\|_X\|v\|_X.$$
 (3.2)

Proposition 3.4. [10] Let $0 \le \lambda < 3$, $\max\{1, \frac{3-\lambda}{2}\} \le p < \infty$, $1 \le q \le 2$, $I = [0,T), 0 < T \le \infty$ and $\frac{2}{3} < \alpha \le \frac{2}{3} + \frac{1}{p'} + \frac{\lambda}{3p}$. Set

$$Y = \mathcal{L}^4(I, \mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{1-\frac{3}{2}\alpha+\frac{3}{p'}+\frac{\lambda}{p}}),$$

there exists a constant C = C(p,q) > 0 depending on p,q such that

$$\|uv\|_{\mathcal{L}^{2}(I,\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{2-3\alpha+\frac{3}{p'}+\frac{\lambda}{p}})} \leq C\|u\|_{Y}\|v\|_{Y}.$$
(3.3)

Proof of Theorem 2.6. For the local existence, we set

$$Y = \mathcal{L}^4(I, \mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-\frac{1}{2} + \frac{3}{p'} + \frac{\lambda}{p}}), I = [0,T)$$

Here, as usual, we begin with the mild integral equation

$$v(t) = T_{\Omega,N}(t)v_0 - \int_0^t T_{\Omega,N}(t-\tau)\tilde{\mathbb{P}}\tilde{\nabla} \cdot (v \otimes v)d\tau, \qquad (3.4)$$

and we consider the bilinear operator B given by

$$B(v,v') = \int_0^t T_{\Omega,N}(t-\tau) \tilde{\mathbb{P}} \tilde{\nabla} \cdot (v \otimes v') d\tau.$$

According to Lemma 3.1 and Proposition 3.4 with $\alpha = 1$, we obtain

$$\begin{split} \|B(v,v')\|_{\mathcal{L}^{4}(I,\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-\frac{1}{2}+\frac{3}{p'}+\frac{\lambda}{p}})} \\ &= \|\int_{0}^{t} T_{\Omega,N}(t-\tau)\tilde{\mathbb{P}}\tilde{\nabla}\cdot(v\otimes v')d\tau\|_{\mathcal{L}^{4}(I,\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-\frac{1}{2}+\frac{3}{p'}+\frac{\lambda}{p}})} \\ &\leq C\|\tilde{\nabla}.(v\otimes v')\|_{\mathcal{L}^{2}(I,\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-2+\frac{3}{p'}+\frac{\lambda}{p}})} \\ &\leq C\|vv'\|_{\mathcal{L}^{2}(I,\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}})} \\ &\leq C\|v\|_{Y}\|v'\|_{Y} \,. \end{split}$$

Lemma 3.2 yields

$$\|T_{\Omega,N}(t)v_0\|_{Y} \le C \|v_0\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}}.$$
(3.5)

Now, we shall decompose the initial data u_0 into two terms

$$v_0 = \mathcal{F}^{-1}(\chi_{B(0,\delta)}\hat{v_0}) + \mathcal{F}^{-1}(\chi_{B^C(0,\delta)}\hat{v_0}) := v_{0,1} + v_{0,2},$$

where $\delta = \delta(v_0) > 0$ is a real number. Since $v_{0,2}$ converge to 0 in $\mathcal{FN}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}$ as $\delta \to +\infty$, by (3.5) there exists δ large enough such that

$$||T_{\Omega,N}(t)v_{0,2}||_Y \le \frac{1}{8C}.$$

For the first term $v_{0,1}$,

$$\begin{split} \left\| T_{\Omega,N}(t)v_{0,1} \right\|_{Y} &\leq \left\| 2^{j(-\frac{1}{2}+\frac{3}{p'}+\frac{\lambda}{p})} \|\varphi_{j}e^{-\mu t|\xi|^{2}}\chi_{B(0,\delta)}\hat{v_{0}} \|_{L^{4}(I,\mathcal{M}_{p}^{\lambda})} \right\|_{\ell^{q}} \\ &\leq \left\| 2^{j(-\frac{1}{2}+\frac{3}{p'}+\frac{\lambda}{p})} \right\| \sup_{\xi\in B(0,\delta)} e^{-\mu t|\xi|^{2}} |\xi|^{\frac{1}{2}} \|_{L^{4}([0,T))} \|\varphi_{j}|\xi|^{-\frac{1}{2}} \hat{v_{0}} \|_{\mathcal{M}_{p}^{\lambda}} \right\|_{\ell^{q}} \\ &\leq C\delta^{\frac{1}{2}}T^{\frac{1}{4}} \left\| v_{0} \right\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}} \,. \end{split}$$

Thus for arbitrary v_0 in $\mathcal{FN}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}$, (3.4) has a unique local solution in Y on [0,T) where

$$T \le \left(\frac{1}{8C^2 \delta^{\frac{1}{2}} \|u_0\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}}}\right)^4.$$

For the global existence, we will again use Lemma 1.1 to ensure the existence of global mild solution with small initial data in the Banach space X given by

$$X = \mathcal{L}^{\infty}([0,\infty), \mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}) \cap \mathcal{L}^{1}([0,\infty), \mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{1+\frac{3}{p'}+\frac{\lambda}{p}}).$$

According to Lemma 3.1 and Proposition 3.3, we obtain

$$\begin{split} \|B(v,v')\|_{\mathcal{L}^{1}([0,\infty),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{1+\frac{3}{p'}+\frac{\lambda}{p}})} \\ &= \left\|\int_{0}^{t} T_{\Omega,N}(t-\tau)\tilde{\mathbb{P}}\tilde{\nabla}\cdot(v\otimes v')d\tau\right\|_{\mathcal{L}^{1}([0,\infty),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{1+\frac{3}{p'}+\frac{\lambda}{p}})} \\ &\leq C\|\tilde{\nabla}.(v\otimes v')\|_{\mathcal{L}^{1}([0,\infty),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{1-2\alpha+\frac{3}{p'}+\frac{\lambda}{p}})} \\ &\leq C\mu^{-1}\|v'\|_{X}\|v\|_{X} \,. \end{split}$$

Similarly,

$$\begin{split} \|B(v,v')\|_{\mathcal{L}^{\infty}([0,\infty),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{1-2\alpha+\frac{3}{p'}+\frac{\lambda}{p}})} \\ &= \left\|\int_{0}^{t} T_{\Omega,N}(t-\tau)\tilde{\mathbb{P}}\tilde{\nabla}\cdot(v\otimes v')d\tau\right\|_{\mathcal{L}^{\infty}([0,\infty),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}})} \\ &\leq C\|\tilde{\nabla}.(v\otimes v')\|_{\mathcal{L}^{1}([0,\infty),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}})} \\ &\leq C\mu^{-1}\|v\|_{X}\|v'\|_{X} \,. \end{split}$$

Finally,

$$||B(v,v')||_X \le C\mu^{-1} ||v||_X ||v'||_X.$$

Lemma 3.2 yields

$$||T_{\Omega,N}(t)v_0||_X \le C ||v_0||_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{1-2\alpha+\frac{3}{p'}+\frac{\lambda}{p}}}$$

If $\|v_0\|_{\mathcal{FN}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}} < C_0\mu$ with $C_0 = \frac{1}{4C^2}$, then (1.1) has a unique global solution $u \in X$ satisfying

$$\|v\|_{\mathcal{L}^{\infty}([0,\infty);\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}})} + \mu\|v\|_{\mathcal{L}^{1}([0,\infty),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{1+\frac{3}{p'}+\frac{\lambda}{p}})} \le 2C\|v_{0}\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}}.$$

Stability of Global Solutions 4

In this section we prove Theorem 2.7. Let T^* be the maximal existence time of

In this section we prove Theorem 2.1. Let Γ be the maximal existence time of a solution u of (1.1) in $\mathcal{L}^{\infty}\left([0,T^*); \mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}\right) \cap \mathcal{L}^{1}\left([0,T^*), \mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{1+\frac{3}{p'}+\frac{\lambda}{p}}\right)$. In order to prove a blow-up criterion of the solution given by Theorem 2.6, assume that $T^* < \infty$ and $\|v\|_{\mathcal{L}^{1}([0,T^*),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{1+\frac{3}{p'}+\frac{\lambda}{p}})} < \infty$, then we can find $0 < T_0 < T^*$ satisfying

$$\|v\|_{\mathcal{L}^1([T_0,T^*),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{1+\frac{3}{p'}+\frac{\lambda}{p}})} < \frac{1}{2}$$

For $t \in [T_0, T^*)$, we explicitly consider the integral equation

$$v(t) = T_{\Omega,N}(t)v(T_0) - \int_{T_0}^t T_{\Omega,N}(t-\tau)\tilde{\mathbb{P}}\tilde{\nabla} \cdot (v \otimes v)d\tau,$$

we obtain

$$|\widehat{v}(t,\xi)| \leq e^{-\mu|\xi|^2 t} |\widehat{v}(T_0,\xi)| + \int_{T_0}^t e^{-\mu(t-s)|\xi|^2} |\widetilde{\mathbb{P}}\widetilde{\nabla} \cdot (v \otimes v)(s,\xi)| \,\mathrm{d}s \,.$$

The same reasoning as in the proof of Proposition 3.3 gives

$$\begin{split} \|v\|_{\mathcal{L}^{\infty}([T_0,t),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}})} &\lesssim \|v(T_0)\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}} \\ &+ \|v\|_{\mathcal{L}^{\infty}([T_0,t),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}})} \|v\|_{\mathcal{L}^{1}([T_0,t),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{1+\frac{3}{p'}+\frac{\lambda}{p}})}. \end{split}$$

It follows that

$$\|v\|_{\mathcal{L}^{\infty}([T_{0},t),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}})} \lesssim \|v(T_{0})\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}} + \frac{1}{2}\|v\|_{\mathcal{L}^{\infty}([T_{0},t),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}})}.$$

We can deduce that

$$\sup_{T_0 \le s \le t} \|v\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}} \lesssim 2\|v(T_0)\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}}, \forall t \in [T_0, T^*).$$

Setting

$$M = \max(2\|v(T_0)\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}}, \max_{t \in [0,T_0]} \|v\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}}),$$

we have

$$\|v(t)\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}} \lesssim M, \ \forall t \in [0,T^*).$$

On the other side

$$v(t) = e^{-t\mu(-\Delta)}u_0 - \Omega \int_0^t e^{-\mu(t-\tau)(-\Delta)} \tilde{\mathbb{P}}\mathcal{B}v(\tau)d\tau - \int_0^t e^{-\mu(t-\tau)(-\Delta)} \tilde{\mathbb{P}}\tilde{\nabla} \cdot (v \otimes v)(\tau)d\tau.$$

Then,

$$\begin{split} v(t') - v(t) &= (e^{-\mu t'(-\Delta)}v_0 - e^{-\mu t(-\Delta)}v_0) \\ &- \left(\int_0^{t'} e^{-\mu (t'-\tau)(-\Delta)}\tilde{\mathbb{P}}\tilde{\nabla} \cdot (v\otimes v)(\tau)d\tau - \int_0^t e^{-\mu (t-\tau)(-\Delta)}\tilde{\mathbb{P}}\tilde{\nabla} \cdot (v\otimes v)(\tau)d\tau\right) \\ &- \Omega \Big(\int_0^{t'} e^{-\mu (t'-\tau)(-\Delta)}\tilde{\mathbb{P}}\mathcal{B}v(\tau)d\tau - \int_0^t e^{-\mu (t-\tau)(-\Delta)}\tilde{\mathbb{P}}\mathcal{B}v(\tau)d\tau\Big) \\ &= [e^{-\mu t'(-\Delta)}v_0 - e^{-\mu t(-\Delta)}v_0] - \Big[\int_t^{t'} e^{-\mu (t'-\tau)(-\Delta)}\tilde{\mathbb{P}}\tilde{\nabla} \cdot (v\otimes v)(\tau)d\tau\Big] \\ &- \Big[\int_0^t e^{-\mu (t-\tau)(-\Delta)}(e^{-\mu (t'-t)(-\Delta)} - 1)\tilde{\mathbb{P}}\tilde{\nabla} \cdot (v\otimes v)(\tau)d\tau\Big] \\ &- \Omega\Big[\int_t^{t'} e^{-\mu (t'-\tau)(-\Delta)}\tilde{\mathbb{P}}\mathcal{B}v(\tau)d\tau\Big] \\ &- \Omega\Big[\int_0^t e^{-\mu (t-\tau)(-\Delta)}(e^{-\mu (t'-t)(-\Delta)} - 1)\tilde{\mathbb{P}}\mathcal{B}v(\tau)d\tau\Big] \\ &:= J_1 + J_2 + J_3 + J_4 + J_5 \,. \end{split}$$

We will estimate J_1 , J_2 , J_3 , J_4 and J_5 ;

$$\begin{split} \|J_1\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}} &= \left\|2^{j(-1+\frac{3}{p'}+\frac{\lambda}{p})}\|\varphi_j(e^{-\mu t'|\xi|^2} - e^{-\mu t|\xi|^2})\hat{u}_0\|_{\mathcal{M}_p^{\lambda}}\right\|_{\ell^q} \\ &\leq \left\|2^{j(-1+\frac{3}{p'}+\frac{\lambda}{p})}\|\varphi_j(e^{-\mu (t'-t)|\xi|^2} - 1)\hat{u}_0\|_{\mathcal{M}_p^{\lambda}}\right\|_{\ell^q}, \end{split}$$

$$\begin{aligned} \|J_2\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}} &\leq \left\|2^{j(-1+\frac{3}{p'}+\frac{\lambda}{p})} \int_t^{t'} \|\varphi_j e^{-\mu(t'-\tau)|\xi|^2} \mathcal{F}(\tilde{\nabla} \cdot v \otimes v)(\tau)\|_{\mathcal{M}_p^{\lambda}} d\tau \right\|_{\ell^q} \\ &\leq \left\|2^{j(\frac{3}{p'}+\frac{\lambda}{p})} \int_t^{t'} \|\varphi_j \mathcal{F}(v \otimes v)(\tau)\|_{\mathcal{M}_p^{\lambda}} d\tau \right\|_{\ell^q}, \end{aligned}$$

$$\begin{split} \|J_{3}\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}} &\leq \\ & \left\|2^{j(-1+\frac{3}{p'}+\frac{\lambda}{p})} \times \int_{0}^{t} \|\varphi_{j}e^{-\mu(t'-\tau)|\xi|^{2}}(1-e^{-\mu(t'-t)|\xi|^{2}})\mathcal{F}(\tilde{\nabla}\cdot v\otimes v)(\tau)\|_{M_{p}^{\lambda}}d\tau\right\|_{\ell^{q}} \\ & \leq \left\|2^{j(\frac{3}{p'}+\frac{\lambda}{p})} \int_{0}^{t} \|\varphi_{j}(e^{-\mu(t'-t)|\xi|^{2}}-1)\mathcal{F}(v\otimes v)(\tau)\|_{M_{p}^{\lambda}}d\tau\right\|_{\ell^{q}}, \end{split}$$

$$\begin{aligned} \|J_4\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}} &\lesssim \left\|2^{j(-1+\frac{3}{p'}+\frac{\lambda}{p})} \int_{t}^{t'} \|\varphi_{j}e^{-\mu(t'-\tau)|\xi|^{2}}\mathcal{F}(\mathcal{B}v)(\tau)\|_{\mathcal{M}_{p}^{\lambda}}d\tau\right\|_{\ell^{q}} \\ &\lesssim \left\|2^{j(-1+\frac{3}{p'}+\frac{\lambda}{p})} \int_{t}^{t'} \|\varphi_{j}\mathcal{F}(\mathcal{B}v)(\tau)\|_{\mathcal{M}_{p}^{\lambda}}d\tau\right\|_{\ell^{q}}, \end{aligned}$$

and

$$\begin{split} \|J_{5}\|_{\mathcal{FN}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}} &\lesssim \\ & \left\|2^{j(-1+\frac{3}{p'}+\frac{\lambda}{p})} \times \int_{0}^{t} \|\varphi_{j}e^{-\mu(t'-\tau)|\xi|^{2}}(1-e^{-\mu(t'-t)|\xi|^{2\alpha}})\mathcal{F}(\mathcal{B}v)(\tau)\|_{\mathbf{M}_{p}^{\lambda}}d\tau\right\|_{\ell^{q}} \\ & \lesssim \left\|2^{j(-1+\frac{3}{p'}+\frac{\lambda}{p})} \int_{0}^{t} \|\varphi_{j}(e^{-\mu(t'-t)|\xi|^{2\alpha}}-1)\mathcal{F}(\mathcal{B}v)(\tau)\|_{\mathbf{M}_{p}^{\lambda}}d\tau\right\|_{\ell^{q}}. \end{split}$$

The dominated convergence theorem gives

$$\limsup_{t,t' \nearrow T^*, t \le t'} \|v(t) - v(t')\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}} = 0.$$

This means that v(t) satisfies the Cauchy criterion at T^* . As $\mathcal{FN}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}$ is a Banach space, then there exists an element v^* in $\mathcal{FN}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}$ such that $v(t) \to v^*$ in $\mathcal{FN}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}$ as $t \to T^*$. Set $v(T^*) = v^*$ and consider the probleme (1.2) starting by v^* . By the well-posedness we obtain a solution existing on a larger time interval than $[0,T^*)$, which is a contradiction. Now, let $v \in \mathcal{C}([0,T^*); \mathcal{FN}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}) \cap \mathcal{L}^1([0,T^*), \mathcal{FN}_{p,\lambda,q}^{1+\frac{3}{p'}+\frac{\lambda}{p}})$ be the maximal solution of (1.1) corresponding to the initial condition v_0 . We want to prove $T^* = \infty$. Put w = v' - v and $w_0 = v_0' - v_0$. We have

$$w_t + \mathcal{A}w + \mathcal{B}w + w \cdot \tilde{\nabla}w + v \cdot \tilde{\nabla}w + w \cdot \tilde{\nabla}v = -\tilde{\nabla}p$$

We first apply $\tilde{\mathbb{P}}$ to the above equation, then we have

$$w_t + \mathcal{A}w = -\tilde{\mathbb{P}}\mathcal{B}w - \tilde{\mathbb{P}}\tilde{\nabla}.(w \otimes w) - \tilde{\mathbb{P}}\tilde{\nabla}.(v \otimes w) - \tilde{\mathbb{P}}\tilde{\nabla}.(w \otimes v) + \tilde{\mathbb{P}}\tilde{\nabla}.$$

Due to Duhamel's formula, we write

$$\begin{split} |\widehat{w}(t,\xi)| &\leq e^{-\mu|\xi|^2 t} |\widehat{w}(0,\xi)| + \int_0^t e^{-\mu(t-s)|\xi|^2} |\mathcal{F}(\tilde{\mathbb{P}}\tilde{\nabla}.(w\otimes w))(s,\xi)| \,\mathrm{d}s \\ &+ \int_0^t e^{-\mu(t-s)|\xi|^2} |\mathcal{F}(\tilde{\mathbb{P}}\tilde{\nabla}.(w\otimes w))(s,\xi)| \,\mathrm{d}s \\ &+ \int_0^t e^{-\mu(t-s)|\xi|^2} |\mathcal{F}(\tilde{\mathbb{P}}\tilde{\nabla}.(w\otimes v))(s,\xi)| \,\mathrm{d}s \\ &+ \int_0^t e^{-\mu(t-s)|\xi|^2} |\mathcal{F}(\tilde{\mathbb{P}}\mathcal{B}w)(s,\xi)| \,\mathrm{d}s \,. \end{split}$$

Then, for $t \in [0, T^*)$ we get

$$\begin{split} \mu \|w\|_{\mathcal{L}^{1}([0,t),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{1+\frac{3}{p'}+\frac{\lambda}{p}})} &\leq C \Big\{ \|w_{0}\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}} + \|\tilde{\nabla}.(w\otimes w)\|_{\mathcal{L}^{1}([0,t),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}})} \\ &+ \|\tilde{\nabla}.(w\otimes w)\|_{\mathcal{L}^{1}([0,t),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}})} \\ &+ \|\tilde{\nabla}.(w\otimes v)\|_{\mathcal{L}^{1}([0,t),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}})} \\ &+ \|\mathcal{B}w\|_{\mathcal{L}^{1}([0,t),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}})} \Big\} \,. \end{split}$$

Similarly,

$$\begin{split} \|w\|_{\mathcal{L}^{\infty}([0,t),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}})} &\leq \|w_0\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}} + \|\tilde{\nabla}.(w\otimes w)\|_{\mathcal{L}^{1}([0,t),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}})} \\ &+ \|\tilde{\nabla}.(v\otimes w)\|_{\mathcal{L}^{1}([0,t),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}})} \\ &+ \|\tilde{\nabla}.(w\otimes v)\|_{\mathcal{L}^{1}([0,t),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}})} \\ &+ \|\mathcal{B}w\|_{\mathcal{L}^{1}([0,t),\mathcal{F}\mathcal{N}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}})} \cdot \end{split}$$

Consequently, for $t \in [0, T^*)$ we get

$$\begin{split} \|w(t)\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}} + \mu\|w\|_{\mathcal{L}^{1}([0,t),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{1+\frac{3}{p'}+\frac{\lambda}{p}})} \\ &\leq C\Big\{\|w_{0}\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}} + \|\tilde{\nabla}.(w\otimes w)\|_{\mathcal{L}^{1}([0,t),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}})} \\ &+ \|\tilde{\nabla}.(v\otimes w)\|_{\mathcal{L}^{1}([0,t),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}})} \\ &+ \|\tilde{\nabla}.(w\otimes v)\|_{\mathcal{L}^{1}([0,t),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}})} \\ &+ \|\mathcal{B}w\|_{\mathcal{L}^{1}([0,t),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}})}\Big\} \\ &\lesssim \|w_{0}\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}} + L_{1} + L_{2} + L_{3} \,. \end{split}$$

Where

$$L_{1} = \|\tilde{\nabla}.(w \otimes w)\|_{\mathcal{L}^{1}([0,t),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}})},$$

$$L_{2} = \|\tilde{\nabla}.(v \otimes w)\|_{\mathcal{L}^{1}([0,t),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}})} + \|\tilde{\nabla}.(w \otimes v)\|_{\mathcal{L}^{1}([0,t),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}})}$$

and $L_3 = \|\mathcal{B}w\|_{\mathcal{L}^1([0,t),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}})}$. The same calculus in the proof of Proposition 3.3 gives

$$\begin{split} L_{1} &\lesssim \|w\|_{\mathcal{L}^{\infty}([0,t),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}})} \|w\|_{\mathcal{L}^{1}([0,t),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{1+\frac{3}{p'}+\frac{\lambda}{p}})},\\ L_{2} &\lesssim \int_{0}^{t} \|w\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}} \|v\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{1+\frac{3}{p'}+\frac{\lambda}{p}}},\\ L_{3} &\lesssim ||\mathcal{B}|| \|w\|_{\mathcal{L}^{1}([0,t),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}})}. \end{split}$$
Then

$$\begin{split} \|w(t)\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}} + \mu \|w\|_{\mathcal{L}^{1}([0,t),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{1+\frac{3}{p'}+\frac{\lambda}{p}})} \\ &\leq C \Big\{ \|w_{0}\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}} + \|w\|_{\mathcal{L}^{\infty}([0,t),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}})} \|w\|_{\mathcal{L}^{1}([0,t),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{1+\frac{3}{p'}+\frac{\lambda}{p}})} \\ &+ \int_{0}^{t} \|w\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}} \|v\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{1+\frac{3}{p'}+\frac{\lambda}{p}}} + \|\mathcal{B}\|\|w\|_{\mathcal{L}^{1}([0,t),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}})} \Big\}. \end{split}$$

Put

$$T = \sup\{t \in [0, T^*), \|w\|_{\mathcal{L}^{\infty}([0,t], \mathcal{F} \dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}})} < \frac{\mu}{4C}\}.$$
(4.1)

For $t \in [0, T)$, we have

$$\begin{split} \|w(t)\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}} &+ \frac{\mu}{2} \|w\|_{\mathcal{L}^{1}([0,t),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}})} \\ &\leq C \left\{ \|w_{0}\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}} + \int_{0}^{t} \|w\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}} (\|\mathcal{B}\| + \|u\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{1+\frac{3}{p'}+\frac{\lambda}{p}}}) \right\}. \end{split}$$

Gronwall's Lemma yields

$$\begin{split} \|w(t)\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}} &+ \frac{\mu}{2} \int_{0}^{t} \|w\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{1+\frac{3}{p'}+\frac{\lambda}{p}}} \\ &\leq C \|w_{0}\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}} \exp\left\{ \int_{0}^{t} C(\|\mathcal{B}\|+\|u\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{1+\frac{3}{p'}+\frac{\lambda}{p}}}) \right\} \\ &\leq C \|w_{0}\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}} \exp\left\{ \int_{0}^{\infty} C(\|\mathcal{B}\|+\|v\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{1+\frac{3}{p'}+\frac{\lambda}{p}}}) \right\}. \end{split}$$

Thus if we take C_0 sufficiently small in (2.4), we have

$$\|w(t)\|_{\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}}} + \frac{\mu}{2} \|w\|_{\mathcal{L}^{1}([0,t),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{-1+\frac{3}{p'}+\frac{\lambda}{p}})} < \frac{\mu}{8C} \,,$$

which contradicts the Definition (4.1).

Then $T = T^*$ and $||w||_{\mathcal{L}^1([0,T^*),\mathcal{F}\dot{\mathcal{N}}_{p,\lambda,q}^{1+\frac{3}{p'}+\frac{\lambda}{p}})} < \infty$, therefore $T^* = \infty$. This completes the proof of Theorem 2.7.

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Regional Controllability of a Class of Time-Fractional Systems

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Abstract. The main purpose of this paper is to develop the concept of regional controllability for an important class of Caputo time-fractional semi-linear systems using the analytical approach, where the dynamic of the considered system is generates by an analytical semigroup. This approach use the fixed point techniques and semigroup theory. Finally, we present some numerical simulations to approve our theoretical results.

1 Introduction

Fractional Calculus has received a considerable amount of interest in the last years, its main purpose is the investigation of the notions of derivation and integration of real or complex order. Many problems in physics, chemistry, engineering and control theory are represented by fractional equations (see [4,14] and [9]), which are being used in modeling the anomalous behavior of problems occurring in the real world . Fractional operators (integration and differentiation) have an important advantage, which is the nonlocal property, where the current state, of a fractional system, depends on historical and past states. Many researchers worked on the existence of solutions for initial and boundary value fractional differential equations (see [13,18]), Zhou and Jiao discussed the existence of mild solution for fractional evolution and neutral evolution equations in Banach spaces based on a probability density function and semigroup theory (see [20] and [19]), seeing this big interest on fractional order systems, it is natural to study and analyze these kinds of systems as an extension or a general case of classical dynamical systems (ie. systems with integer order derivatives).

The analysis of dynamical systems consists of many branches and various concepts, Controllability being one amongst others. The concept of controllability consists of steering a system into a desired state (exactly or approximately) at time T from an arbitrary initial state. The concept in hand has a very vast literature for various type of systems (linear, Semi-linear, nonlinear...), for more informations (see [3,6,8,11,17,21]). In many practical applications their exists states which are not reachable, also sometimes we only need to control the system on a particular region, in this cases the regional controllability concept should be considered (see [12,15,16] and references therein).

Regional controllability's purpose is to steer a system into a desired state only in a subregion of the whole evolution domain, this notion is a general case of

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'global' Controllability. This notion is developed by several researchers to cover various types of systems. In particular, recently Ge,Chen and Kou discussed the regional controllability for time-fractional sub-diffusion systems with Caputo and Riemann-Liouville fractional derivatives (see [7]).

The concept in hand, namely regional controllability for nonlinear fractional systems, is in an initial stage and needs some more research, thus the motivation for this work, is to develop this theory for semi-linear time-fractional systems with Caputo derivative by using the analytical approach, which is based on the fixed point techniques and semigroup theory.

This paper is presented as follows, in Sect. 2, we introduce some preliminaries, definitions and results which will be used throughout this work. In Sect. 3, by using some properties of analytical semigroup and under suitable assumptions we show that the considered system is regionally controllable by a control that will be given later. In Sect. 4, we provide an algorithm, which is based on the steps of the used approach. The Sect. 5 is devoted to present successful numerical results illustrating the theoretical ones. Finally a conclusion shall be giving.

2 Preliminaries and Considered System

In this section, we recall some basic definitions and properties used throughout this paper.

Definition 1 [10]. The (left) Caputo fractional derivative of a function y at a point t of order $\alpha \in [0, 1]$ is defined as follows :

$${}^{C}\mathrm{D}_{0^{+}}^{\alpha}y(t) = \frac{1}{\Gamma(1-\alpha)} \int_{0}^{t} (t-s)^{-\alpha} \frac{d}{ds}(y(s))ds, \qquad 0 \le t < T.$$
(1)

We have the following two propositions.

Proposition 1 [1]. Let X and Y be two Banach spaces. Let's consider $f \in L^1_{loc}(0,T;X)$ and $\mathscr{T}:[0,T] \to \mathscr{L}(X,Y)$ be a strongly continuous function. Then the convolution

$$(\mathscr{T}*f)(t) := \int_0^t \mathscr{T}(t-s)f(s)ds,$$

exists in the bochner sense and defines a continuous function from [0,T] into Y

$$\begin{split} & \textbf{Proposition 2 } [1](Young's \ Inequality). \ Let's \ consider \ p,q,s \geq 1 \ such \ that \ \frac{1}{q} + \\ & \frac{1}{p} = 1 + \frac{1}{s}. \\ & If \ \mathscr{T} \in L^p(0,T;\mathscr{L}(X,Y)) \ and \ f \in L^q(0,T;X), \ then \\ & \mathscr{T} * f \in L^s(0,T;Y) \qquad and \quad \|\mathscr{T} * f\|_{L^s(0,T;Y)} \leq \|f\|_{L^q(0,T;X)}. \|\mathscr{T}\|_{L^p(0,T;\mathscr{L}(X,Y))}. \end{split}$$

Let w be a measurable function defined from [0,T] to \mathbb{R}^+ . The Weighted Lebesgue space ([2]) associate to w is defined by:

$$L_{w}^{p}[0,T] := \left\{ f \in L_{loc}^{p}[0,T] \mid \int_{0}^{T} w(t) |f(t)|^{p} dt < +\infty \right\}, \qquad p \ge 1$$

which is a Banach space endowed with the norm :

$$\|f\|_{L^p_w[0,T]} = \left[\int_0^T w(t)|f(t)|^p dt\right]^{\frac{1}{p}}$$

For $0 < \alpha \leq 1$, let $w(t) = t^{\alpha-1}$, we denote $L_w^p[0,T] := L_{\alpha-1}^p(0,T)$ and we have the following inclusion $L_{\alpha-1}^p(0,T) \subset L^p[0,T]$.

Let's consider $n \in \mathbb{N}^*$, Ω an open bounded subset of \mathbb{R}^n with smooth enough boundary $\partial \Omega$ and let $\alpha \in]0,1]$. For a time T > 0, set $Q = \Omega \times]0,T]$ and $\Sigma = \partial \Omega \times]0,T]$. Let's consider the following fractional semi-linear evolution equation:

$$\begin{cases} {}^{C}D_{0^{+}}^{\alpha}y_{u}(x,t) + Ay_{u}(x,t) = Ny_{u}(x,t) + Bu(t) \text{ in } Q, \\ y_{u}(\xi,t) = 0 & \text{ on } \Sigma, \\ y_{u}(x,0) = y_{0}(x) & \text{ in } \Omega, \end{cases}$$
(2)

where

- -A is the infinitesimal generator of a C_0 semi-group $\{\mathscr{S}(t)\}_{t\geq 0}$ on the Hilbert space $\mathbf{X} = L^2(\Omega)$.
- N a nonlinear operator.
- B is the control operator from ${\mathscr U}$ into X which is linear.
- u is given in $U = L^2(0, T, \mathscr{U})$ and $y_0 \in X$.

Without loss of generality, we denote $y_u(.,t) := y_u(t)$.

Definition 2 [5,19]. A mild solution of the system (2) is any function y_u in C(0,T;X) satisfying the following integral equation :

$$y_{u}(t) = S_{\alpha}(t)y_{0} + \int_{0}^{t} (t-\tau)^{\alpha-1} \mathbf{K}_{\alpha}(t-\tau)Ny(\tau)d\tau + \int_{0}^{t} (t-\tau)^{\alpha-1} \mathbf{K}_{\alpha}(t-\tau)Bu(\tau)d\tau,$$
(3)

where

$$S_{\alpha}(t) = \int_{0}^{\infty} \phi_{\alpha}(\theta) \mathscr{S}(t^{\alpha}\theta) d\theta,$$
$$K_{\alpha}(t) = \alpha \int_{0}^{\infty} \theta \phi_{\alpha}(\theta) \mathscr{S}(t^{\alpha}\theta) d\theta,$$

$$\phi_{\alpha}(\theta) = \frac{1}{\alpha} \theta^{-1-\frac{1}{\alpha}} W_{\alpha}(\theta^{-\frac{1}{\alpha}}) \quad \text{for all } \theta \quad \text{positive,}$$

and

$$W_{\alpha}(\theta) = \frac{1}{\pi} \sum_{n=1}^{\infty} (-1)^{n-1} \theta^{-n\alpha-1} \frac{\Gamma(n\alpha+1)}{n!} sin(n\pi\alpha).$$

we have the following proposition.

Proposition 3 [20]. For all $\beta \geq -1$, we have

$$\int_0^\infty \theta^\beta \phi_\alpha(\theta) d\theta = \frac{\Gamma(1+\beta)}{\Gamma(1+\alpha\beta)},$$

then we have the following remark.

Remark 1. If $\beta = 0$, we can see that ϕ_{α} is a probability density.

Let $\omega \subset \Omega$ be a subregion with positive Lebesgue measure. The restriction operator in ω is defined as follows:

$$\chi_{\omega}: L^2(\Omega) \longrightarrow L^2(\omega)$$
$$y \longmapsto y_{|_{\omega}}$$

and we denote its adjoint by χ^*_{ω} .

The mild solution defined by (3) can be written :

$$y_u(t) = S_\alpha(t)y_0 + \mathcal{L}_\alpha(t)Ny_u(.) + \mathcal{L}_\alpha(t)Bu(.), \qquad (4)$$

where

$$\mathcal{L}_{\alpha}(t)y(.) = \int_{0}^{t} (t-\tau)^{\alpha-1} \mathcal{K}_{\alpha}(t-\tau)y(\tau)d\tau.$$

We also define the restriction of the controllability operator in ω by:

$$\begin{aligned} \mathbf{H}^{\alpha}_{\omega} &: U \longrightarrow L^{2}(\omega) \\ u &\longmapsto \chi_{\omega} \mathbf{L}_{\alpha}(T) B u. \end{aligned}$$

Definition 3. The system (2) is said to be exactly (respectively, approximately) ω -controllable if for all $y_d \in L^2(\omega)$ (respectively, for all $\epsilon > 0$ and for all $y_d \in L^2(\omega)$), we can find a control $u \in U$ such that $\chi_{\omega} y_u(T) = y_d$ (respectively, $||\chi_{\omega} y_u(T) - y_d||_{L^2(\omega)} \leq \epsilon$).

Problem: For any state y_d in $L^2(\omega)$, is it possible to find a control u^* that steer the system (2) in a finite time T to y_d only in the subregion ω ?

We consider the following linear system associate to the nonlinear system (2):

$$\begin{cases} {}^{C}D_{0^{+}}^{\alpha}y(x,t) + Ay(x,t) = Bu(t) \text{ in } Q, \\ y(\xi,t) = 0 & \text{ on } \Sigma, \\ y(x,0) = y_{0}(x) & \text{ in } \Omega, \end{cases}$$
(5)

which we suppose, for the rest of this work, to be approximately ω -controllable. Then we give the following proposition.

If the following hypotheses hold Proposition 4.

• $[y_d - \chi_\omega S_\alpha(T)y_0 - \chi_\omega \mathcal{L}_\alpha(T)Ny_{u^*}(.)] \in Im(H^\alpha_\omega),$ • $Im(H^\alpha_\omega)$ a closed subset.

Then the system (2) is exactly ω - controllable by the control

$$u^*(.) = \mathbf{H}_{\omega}^{\alpha^{\dagger}}[y_d - \chi_{\omega}S_{\alpha}(T)y_0 - \chi_{\omega}\mathbf{L}_{\alpha}(T)Ny_{u^*}(.)].$$

Where

$$H^{\alpha^{\dagger}}_{\omega} := H^{\alpha^{*}}_{\omega} \left(H^{\alpha}_{\omega} H^{\alpha^{*}}_{\omega} \right)^{-1}$$
 is the Pseudo-inverse operator of H^{α}_{ω} .

Proof. Using the expression (3), the solution of system (2) controlled by u^* is giving by the following formula

$$y_{u^*}(t) = S_{\alpha}(t)y_0 + L_{\alpha}(t)Ny_{u^*}(.) + L_{\alpha}(t)BH_{\omega}^{\alpha'}[y_d - \chi_{\omega}S_{\alpha}(T)y_0 - \chi_{\omega}L_{\alpha}(T)Ny_{u^*}(.)],$$

hence

$$\chi_{\omega}y_{u^*}(T) = \chi_{\omega}S_{\alpha}(T)y_0 + \chi_{\omega}\mathcal{L}_{\alpha}(T)Ny_{u^*}(.) + \mathcal{H}_{\omega}^{\alpha}\mathcal{H}_{\omega}^{\alpha^{\dagger}}[y_d - \chi_{\omega}S_{\alpha}(T)y_0 - \chi_{\omega}\mathcal{L}_{\alpha}(T)Ny_{u^*}(.)],$$

since

$$[y_d - \chi_{\omega} S_{\alpha}(T) y_0 - \chi_{\omega} \mathcal{L}_{\alpha}(T) N y_{u^*}(.)] \in \operatorname{Im}(H_{\omega}^{\alpha}),$$

 $H^{\alpha}_{\omega}H^{\alpha^{\dagger}}_{\omega}$ is the orthogonal projection on $Im(H^{\alpha}_{\omega})$. and Then

$$\chi_{\omega}y_{u^*}(T) = \chi_{\omega}S_{\alpha}(T)y_0 + \chi_{\omega}\mathcal{L}_{\alpha}(T)Ny_{u^*}(.) + y_d - \chi_{\omega}S_{\alpha}(T)y_0 - \chi_{\omega}\mathcal{L}_{\alpha}(T)Ny_{u^*}(.) = y_d.$$

In the next section, we will study the regional controllability of the system (2)in $\operatorname{Im}(H^{\alpha}_{\omega})$ endowed with the norm

$$||y_d||_{\operatorname{Im}(H^{\alpha}_{\omega})} = ||H^{\alpha^{\dagger}}_{\omega} y_d||_{\omega}$$

Remark 2. $||.||_{\operatorname{Im}(H^{\alpha}_{\omega})}$ defines a semi-norm on $\operatorname{Im}(H^{\alpha}_{\omega})$ but it becomes a norm if the linear system (5) is approximately ω -controllable.

Proof. It is sufficient to show that

$$||y_d||_{\operatorname{Im}(H^{\alpha}_{\omega})} = 0 \Longrightarrow y_d = 0$$

We have

$$||y_d||_{\operatorname{Im}(H^{\alpha}_{\omega})} = 0 \Longrightarrow ||H^{\alpha^{\dagger}}_{\omega} y_d||_{U} = 0$$
$$\Longrightarrow H^{\alpha^{\dagger}}_{\omega} y_d = 0$$
$$\Longrightarrow (H^{\alpha*}_{\omega} H^{\alpha}_{\omega}) H^{\alpha^{\dagger}}_{\omega} y_d = 0$$
$$\Longrightarrow H^{\alpha*}_{\omega} y_d = 0.$$

Since the linear system (5) is approximately ω -controllable, then ker $(H_{\omega}^{\alpha*}) = \{0\}$ by [7], therefore $y_d = 0$.

3 Analytical Approach

We consider the system (2) with $y_0 = 0$, moreover, let -A the infinitesimal generator of an analytic semigroup of bounded linear operator $(T(t))_{t>0}$ on X.

Let 0 be an element of the resolvent set of -A, then it is possible to define the fractional power A^{ν} for any ν belongs to the interval]0,1]. $X^{\nu} := D(A^{\nu})$ is a Banach space, which is dense in X, endowed with the graph norm: $||.||_{X^{\nu}} =$ $||A^{\nu}(.)||_{X}$.

Remark 3. For the sake of simplification, we choose the order of fractional power of A to be the same as the order of fractional derivative.

We have the following proposition.

Proposition 5. [12] For all $\alpha \in [0,1]$, the following properties are satisfied

- (i) $\exists C_{\alpha} > 0$ such that $||A^{\alpha}T(t)||_{\mathscr{L}(X,X)} \leq C_{\alpha}t^{-\alpha} \qquad 0 < t \leq T.$
- (ii) $\forall t \in [0, T]$, we have

$$||\mathbf{K}_{\alpha}(t)||_{\mathscr{L}(X,X^{\alpha})} \leq \frac{\alpha C_{\alpha}}{t^{\alpha^{2}}} \times \frac{\Gamma(2-\alpha)}{\Gamma(1+\alpha(1-\alpha))} := f_{\alpha}(t).$$

Corollary 1. Let's consider $H(t) = t^{\alpha-1}K_{\alpha}(t)$ and $q \ge 1$. If $f_{\alpha} \in L^{q}_{\alpha-1}(0,T)$, then

$$H \in L^{q}(0,T;\mathscr{L}(X,X^{\alpha}))$$
 and $||H(.)||_{L^{q}(0,T;\mathscr{L}(X,X^{\alpha}))} \leq ||f_{\alpha}(.)||_{\mathrm{L}^{q}_{\alpha-1}(0,T)}.$

Hypotheses :

We assume that the following conditions hold.

(i) For all $p, s \ge 1$, there exists $q \ge 1$ such that

$$\frac{1}{q} = 1 + \frac{1}{p} - \frac{1}{s}$$
 and $f_{\alpha} \in L^{q}_{\alpha-1}(0,T).$ (6)

(ii) Let $N: L^p(0,T;X^{\alpha}) \longrightarrow L^s(0,T;X)$ be the nonlinear operator satisfying

$$\begin{cases} N(0) = 0, \\ ||Nx - Ny||_{L^{s}(0,T;X)} \le k(||x||, ||y||)||x - y||_{L^{p}(0,T;X^{\alpha})}, \end{cases}$$
(7)

where $k : \mathbb{R}^+ \times \mathbb{R}^+ \longrightarrow \mathbb{R}^+$ is such that $\lim_{(\theta_1, \theta_2) \to (0, 0)} k(\theta_1, \theta_2) = 0.$

We define the operator

$$\Psi(y_d, u) = \mathbf{H}_{\omega}^{\alpha^{\dagger}}(y_d - \chi_{\omega} \mathbf{L}_{\alpha}(T) N y_u).$$

Then the regional controllability problem becomes a fixed point problem of the function $\Psi(y_d, .)$, where y_d is an element of $\text{Im}(H_{\omega}^{\alpha})$

Theorem 1. If the hypotheses (i) and (ii) hold and

(iii)

$$|\mathcal{L}_{\alpha}(.)Bu||_{L^{p}(0,T;X^{\alpha})} \leq \beta ||u||_{U}, \qquad \beta > 0,$$
(8)

(iv)

$$||\chi_{\omega} \mathbf{K}_{\alpha}(.)||_{\mathscr{L}(X,Im(H_{\omega}^{\alpha}))} = g_{\alpha} \in L^{r}_{\alpha-1}(0,T), \quad \frac{1}{r} + \frac{1}{s} = 1,$$
(9)

are satisfied, then the following assertions hold.

- 1. There exists a > 0, $\rho = \rho(a) > 0$ and m = m(a) > 0 such that for any state y_d in $\mathscr{B}(0,\rho) \subset Im(H^{\alpha}_{\omega})$ there exists u^* in $\mathscr{B}(0,m)$ that steers the system (2) to y_d in ω . Where $\mathscr{B}(0,k)$ is a ball with center 0 and radius k.
- 2. The mapping

$$F: \mathscr{B}(0,\rho) \longrightarrow U$$

$$y_d \longrightarrow u^*$$

 $is \ a \ lipschitz \ mapping.$

Proof. 1- Based on hypothesis (ii), we have

$$\lim_{(\theta_1,\theta_2)\to(0,0)} k(\theta_1,\theta_2) = 0,$$

then $\exists a > 0, \exists \nu > 0$ such that

$$k(\theta_1, \theta_2) < \nu < \frac{1}{\beta ||g_{\alpha}||_{L^r_{\alpha-1}(0,T)} + ||f_{\alpha}||_{L^q_{\alpha-1}(0,T)}} \qquad \forall \theta_1, \theta_2 \le a_1$$

which gives

$$\sup_{\theta_i \leq a} k(\theta_1, \theta_2) \leq \nu < \frac{1}{\beta ||g_{\alpha}||_{L^r_{\alpha-1}(0,T)} + ||f_{\alpha}||_{L^q_{\alpha-1}(0,T)}}.$$

Let's consider $A_1 = \beta ||g_{\alpha}||_{L^r_{\alpha-1}(0,T)} \sup_{\theta_i \leq a} k(\theta_1, \theta_2)$ and $A_2 = ||f_{\alpha}||_{L^q_{\alpha-1}(0,T)}$

 $\sup_{\substack{\theta_i \leq a}} k(\theta_1, \theta_2).$ We have $A_1 < 1$ and $A_2 < 1$. If we set $m = \frac{a}{2}(1 + \frac{1}{2})$

$$m = \frac{a}{\beta} (1 - ||f_{\alpha}||_{L^{q}_{\alpha-1}(0,T)} \sup_{\theta \le a} k(\theta, 0)),$$

then m is positive. In fact,

$$||f_{\alpha}||_{L^{q}_{\alpha-1}(0,T)} \sup_{\theta \le a} k(\theta,0) \le A_{2} < 1.$$

Moreover, the following function

$$\begin{array}{ccc} f:\mathscr{B}(0,m) \longrightarrow \mathscr{B}(0,a) \\ \\ u & \longmapsto y_u \end{array}$$

is a Lipschitz mapping with constant $\frac{\beta}{1-A_2}$. For that, by the equation (4) and corollary (1), for all $u, v \in B(0, m)$, we have

$$\begin{aligned} ||y_u - y_v||_{L^p(0,T;X^{\alpha})} &= ||\mathcal{L}_{\alpha}(.)N(y_u - y_v) + \mathcal{L}_{\alpha}(.)B(u - v)||_{L^p(0,T;X^{\alpha})} \\ &\leq ||(H * N(y_u - y_v))(.)||_{L^p(0,T;X^{\alpha})} + ||\mathcal{L}_{\alpha}(.)B(u - v)||_{L^p(0,T;X^{\alpha})}, \end{aligned}$$

using hypotheses (ii) and (iii) we get

$$||y_u - y_v||_{L^p(0,T;X^{\alpha})} \le A_2||y_u - y_v||_{L^p(0,T;X^{\alpha})} + \beta||u - v||_U$$

hence f is a lipschitz mapping with constant $\frac{\beta}{1-A_2}$. Next we show that $\Psi(y_d, .)$ has a unique fixed point in $\mathscr{B}(0, m)$. Let's consider $y_d \in \text{Im}(H^{\alpha}_{\omega})$ and $u, v \in \mathscr{B}(0, m)$, we have

$$\begin{split} ||\Psi(y_d, u) - \Psi(y_d, v)||_U &= ||\chi_{\omega} \mathcal{L}_{\alpha}(T)(Ny_u - Ny_v)||_{\operatorname{Im} H^{\alpha}_{\omega}} \\ &\leq ||g_{\alpha}||_{L^{r}_{\alpha-1}(0,T)} ||Ny_u - Ny_v||_{L^{s}(0,T;X)} \\ &\leq ||g_{\alpha}||_{L^{r}_{\alpha-1}(0,T)} \sup_{(\theta_i \leq a)} k(\theta_1, \theta_2)||y_u - y_v||_{L^{p}(0,T;X^{\alpha})}, \end{split}$$

since f is lipschitz, then

$$||\Psi(y_d, u) - \Psi(y_d, v)||_U \le \frac{\beta A_1}{1 - A_2} ||u - v||_U.$$
(10)

If we denote $A_3 := \frac{\beta A_1}{1 - A_2}$, we have $A_3 < 1$, thus $\Psi(y_d, .)$ is a strict contraction mapping. For $u \in \mathscr{B}(0, m)$, we have $y_u \in \mathscr{B}(0, a)$ and

$$\begin{split} ||\Psi(y_d, u)||_U &= ||y_d - \chi_{\omega} \mathcal{L}_{\alpha}(T) N y_u||_{\operatorname{Im}(H^{\alpha}_{\omega})} \\ &\leq ||y_d||_{\operatorname{Im}(H^{\alpha}_{\omega})} + ||\chi_{\omega} \mathcal{L}_{\alpha}(T) N y_u||_{\operatorname{Im}(H^{\alpha}_{\omega})} \\ &\leq ||y_d||_{\operatorname{Im}(H^{\alpha}_{\omega})} + ||g_{\alpha}||_{L^{r}_{\alpha-1}(0,T)} a \sup_{(\theta \leq \alpha)} k(\theta, 0), \end{split}$$

therefore, if

$$||y_d||_{\operatorname{Im}(H^{\alpha}_{\omega})} \le m - ||g_{\alpha}||_{L^r_{\alpha-1}(0,T)} a \sup_{(\theta \le a)} k(\theta, 0),$$

then $\Psi(y_d, u) \in \mathscr{B}(0, m)$. We set $\rho = \frac{a}{\beta}(1 - (||f_{\alpha}||_{L^q_{\alpha-1}(0,T)} + \beta||g_{\alpha}||_{L^r_{\alpha-1}(0,T)}) \sup_{\theta \leq a} k(\theta, 0))$, hence, if $y_d \in \mathscr{B}(0, \rho) \subset \text{Im}(H^{\alpha}_{\omega})$, we deduce from the Picard fixed point theorem that $\Psi(y_d, .)$ admits a unique fixed point $u^* \in \mathscr{B}(0, m)$. We remark that u^* obtained is solution of the exact regional controllability problem.

2- Let z_d and y_d in $\mathscr{B}(0,\rho)$, we have

$$F(z_d) - F(y_d) = \Psi(z_d, F(z_d)) - \Psi(z_d, F(y_d)) + \Psi(z_d, F(y_d)) - \Psi(y_d, F(y_d)),$$

since

$$||\Psi(z_d, F(z_d)) - \Psi(z_d, F(y_d))||_U \le A_3 ||F(z_d) - F(y_d)||_U,$$

$$||\Psi(z_d, F(y_d)) - \Psi(y_d, F(y_d))||_U = ||z_d - y_d||_{\text{Im}(H^{\alpha}_{\omega})}.$$

Then

$$||F(z_d) - F(y_d)||_U \le \frac{1}{1 - A_3} ||z_d - y_d||_{\operatorname{Im}(H^{\alpha}_{\omega})},$$

therefore, F satisfies the Lipschitz condition.

Moreover, the Picard fixed point theorem gives also the existence of a sequence which converges to the control u^* .

We give the following proposition.

Proposition 6. The sequence

$$\begin{cases} u_0 = 0 \\ u_{n+1} = \mathbf{H}_{\omega}^{\alpha^{\dagger}} (y_d - \chi_{\omega} \mathbf{L}_{\alpha}(T) N y_{u_n}), \end{cases}$$
(11)

converges to u^* in $\mathscr{B}(0,m) \subset U$.

Proof. Let's consider $n, k \in \mathbb{N}^*$ we have

$$||u_{n+k} - u_n||_U \le \sum_{l=n}^{n+k-1} ||u_{l+1} - u_l||_U.$$

By the Inequality (10) we obtain

$$||u_{l+1} - u_l||_U = ||\Psi(y_d, u_l) - \Psi(y_d, u_{l-1})||_U \le A_3 ||u_l - u_{l-1}||_U \le A_3^l ||u_1||_U,$$

which yields

$$||u_{n+k} - u_n||_U \le \sum_{l=n}^{n+k-1} A_3^l ||u_1||_U \le \frac{1 - A_3^k}{1 - A_3} A_3^n ||u_1||_U,$$

hence since $A_3^n \to 0$, we conclude that $\lim_{n \to +\infty} ||u_{n+k} - u_n||_U = 0.$

The sequence $(u_n)_n$ is a Cauchy sequence on $\mathscr{B}(0,m)$, then $(u_n)_n$ converges to u in $\mathscr{B}(0,m)$.

Passing to the limit in (11), we have $u = \Psi(y_d, u)$, since $\Psi(y_d, .)$ has a unique fixed point in B(0, m), then $u = u^*$.

4 Algorithm

In this section, we present an algorithm which has as objective, finding a control that steering the considered system to the desired state only in ω , this leads to some numerical simulations which will be presented in the next section.

Algorithm 1

```
Initialization:

Fractional order of derivative \alpha.

The region \omega.

Actuator (D, f).

r_1 = y_d.

Error estimate \varepsilon.

Calculation of u_1 = \operatorname{H}_{\omega}^{\dagger \alpha} r_1 and obtain y_{u_1}(T).

repeat

r_n = r_{n-1} + (y_d - \chi_{\omega} y_{u_{n-1}}(T)), \quad n \ge 2.

Calculation of u_n = \operatorname{H}_{\omega}^{\omega^{\dagger}} r_n.

Solve the semi-linear system (2) controlled by u_n.

until

||\chi_{\omega} y_{u_n}(T) - y_d||_{Im(\operatorname{H}^{\alpha})} < \varepsilon.
```

5 Numerical Results

In this section, we present two numerical simulations illustrating our theoretical result where the first one is done by using zonal actuator and the second example is giving by using a pointwise actuator.

5.1 Case of Zonal Actuator

Let's consider the following sub-diffusion one-dimensional system with order $\alpha = 0.7$:

$$\begin{cases} {}^{\scriptscriptstyle C}D_{0^+}^{0.7}z(x,t) - \frac{\partial^{2z(x,t)}}{\partial x^2} = \chi_D u(t) + \sum_{j=1}^{\infty} (\langle z, \varphi_j \rangle)^2 \varphi_j(x) & \text{in } [0,1] \times]0,3] \\ z(x,t) = 0 & \text{on}\{0,1\} \times]0,3] \\ z(x,0) = 0 & \text{in } [0,1], \end{cases}$$
(12)

where $\varphi_j(x) = \sqrt{2}\sin(j\pi x)$.

The control operator in the system (12) is given by a zonal actuator (D, f) where D = [0.2, 0.3] and f = 1.

We consider the region $\omega = [0.4, 0.68]$ and the desired state $z_d(x) = 5.3 x^2 (x-1)^2 (x-0.4)$.

Using the previous algorithm, we obtain the following results:



Fig. 1. Desired state and estimate final one in $\omega = [0.4, 0.68]$

In the subregion]0.4, 0.68], we can see that the regional final state and the desired state z_d are very close with error $|| \chi_{\omega} z_u(x,t) - z_d ||_{L^2(\omega)} = 7.05 \times 10^{-4}$. Here we have the evolution of the control function.

0.5 0.4 0.3 **Control Function** 0.2 0.1 0 -0.1 0 0.5 1 1.5 2 2.5 3 Times

Fig. 2. Control function.

with a transfer cost $|| u^* ||_{L^2(0,T)}^2 = 0.13$.

5.2 Case of Pointwise Actuator

We will treat the same kind of system with $\alpha = 0.8$ and a pointwise actuator located in b = 0.4, which amounts to consider the following system:

$$\begin{cases} {}^{\scriptscriptstyle C}D_{0^+}^{^{0.8}}z(x,t) - \frac{\partial^2}{\partial x^2}z(x,t) = \delta_b u(t) + \sum_{j=1}^{\infty} (\langle z, \varphi_j \rangle)^2 \varphi_j(x) & \text{in } [0,1] \times]0,2] \\ z(\xi,t) = 0 & \text{on } \{0,1\} \times]0,2] \\ z(x,0) = 0 & \text{in } [0,1], \end{cases}$$

where $\varphi_j(x) = \sqrt{2}\sin(j\pi x)$.

The subregion under consideration is $\omega =]0.45, 0.7]$. Let's consider the following desired state:

$$z_d(x) = x(x-1)(3.5x-0.2)(0.6-x)(x-0.1).$$

Using the proposed algorithm, we obtain the following figure:



Fig. 3. Desired and estimate final states in ω .

We can see that the final state is very close to the desired state in the subregion]0.45, 0.7] with an error of

$$|| \chi_{\omega} z_u(t) - z_d ||_{L^2(\omega)} = 6.04 \times 10^{-4}$$

The following figure shows the evolution of the control function. with a transfer cost $||\;u^*\;||^2_{L^2(0,T)}=0.03.$



Fig. 4. Control function.

6 Conclusion

In this work, we have studied the regional controllability of Caputo timefractional sub-diffusion system with analytical approach, which is a technical one, based on fixed point techniques and semigroup theory, we also presented an algorithm based on our theoretical results, which leads to successful numerical results. As a future work, we are working on the concept of boundary and gradient regional controllability for the same type of systems.

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Quadratic Optimal Control for Bilinear Systems

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Abstract. In this work, we will investigate the quadratic optimal control for bilinear systems. We will first study the existence of a solution for the considered optimal control. Then, we will focus on a special class of bilinear systems for which the quadratic optimal control can be expressed in a feedback law form. The approach relies on the conditions of optimality and linear semi-group theory.

Keywords: Quadratic cost \cdot Optimal control \cdot Bilinear systems

1 Introduction

The subject of this paper is to study the quadratic optimal control for infinite dimensional bilinear systems. The importance of bilinear systems lies in the fact that they represent a theoretical model for real processes (natural or industrial) and that they represent a first generalization of linear systems. Optimal control theory consists in seeking the best control strategy among others from the set of admissible controls, that is the one that enables us to reach a precise objective (reach a desired state, minimize a cost or energy, etc.). Quadratic cost functions have a wide use in differential geometry, statistics, special relativity, solid mechanics, etc. Optimal control problems have been the subject of several works. In [7], Kalman has studied the problem of quadratic optimal control for linear finite-dimensional systems. He characterized the control in term of the system's state and the solution of the Riccati equation. These results have been generalized by Banks and Yu [4] to a class of infinite dimensional semi-linear systems. The case of distributed bilinear systems has been considered by Alami [2]. Another approach, based on the Pontryagin's maximum principle, has been developed by Pontryagin [10] for finite-dimensional systems. These results have been generalized by different authors (e.g. [1, 5, 8, 11, 12]) for a class of semi-linear systems with a variety of cost functions. The main limit of the approach based on the adjoint equation is that such an equation involves the unknown optimal control. This together with the nonlinear dependence of the state w.r.t to control becomes somehow embarrassing when one looks for the explicit expression of optimal control.

In this work, we consider the problem of optimal control that minimizes a given quadratic cost. In Sect. 2, we will first state our quadratic optimal problem

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for bilinear systems and then study the existence of an optimal control. In Sect. 3, we will focus on a special class of bilinear system for which the quadratic optimal control can be expressed in a feedback law form.

2 Quadratic Optimal Control

2.1 Problem Statement

Let us consider the following bilinear system

$$\begin{cases} \dot{y}(t) = Ay(t) + u(t)By(t)\\ y(0) = y_0 \in X \end{cases}$$

$$\tag{1}$$

where

- $A: D(A) \subset X \mapsto X$ is the infinitesimal generator of a linear C_0 semigroup S(t) of isometries on a real Hilbert space X whose inner product and the corresponding norm are denoted respectively by $\langle ., . \rangle$ and $\|.\|$,
- $B: X \mapsto X$ is a linear bounded operator,
- $u(\cdot)$ is a scalar valued control that belongs to the control space $L^2(0,T)$ and $y(\cdot)$ is the corresponding mild solution with initial state $y_0 \in X$.

The quadratic cost function J to be minimized is defined by

$$J(u) = 2||y(T)|| + \int_0^T ||y(t)||^2 dt + \int_0^T u(t)^2 dt,$$
(2)

for any admissible control u, i.e. for which the corresponding solution y exists and J(u) makes sense.

It is well known that for any $u \in L^2(0,T)$ the system (1) admits a unique mild solution (see [3]) and we have $J(u) < +\infty$.

In the sequel we take $U_{ad} := L^2(0, T)$ (the set of admissible control).

The optimal control problem may be stated as follows

$$\begin{cases} \min J(u) \\ u \in U_{ad} \end{cases}$$
(3)

Let us introduce the following time varying cost function

$$J(u)(t) = 2||y(t)|| + \int_0^t ||y(s)||^2 ds + \int_0^t u(s)^2 ds, \quad t \in [0, T].$$
(4)

2.2 Existence of the Optimal Control

In this subsection, we prove the existence of a square integrable control u that minimizes the cost (2).

Theorem 1. There exists an optimal control solution of the problem (3).

Proof: Since the set $\{J(u)/u \in U_{ad}\} \subset \mathbb{R}^+$ is not empty and bounded from below, it admits a lower bound J^* . Let $(u_n)_{n \in \mathbb{N}}$ be a minimizing sequence such that $J(u_n) \to J^*$.

By the coercivity of the mapping $R: u \mapsto \int_0^T ||u(t)||^2 dt$, we deduce that the sequence (u_n) is bounded, so it admits a subsequence still denoted by (u_n) as well, which weakly converges to $u^* \in U_{ad}$.

Let y_n and y^* be the solutions of (1) respectively corresponding to u_n and u^* . From Theorem 3.6 of [3] we have

$$\lim_{n \to +\infty} \|y_n(t) - y^*(t)\| = 0, \ \forall t \in [0, T].$$
(5)

Since the norm $\|.\|$ is lower semi-continuous, it follows from (5) that for all $t \in [0, T]$

$$||y^*(t)||^2 \le \lim_{n \to +\infty} \inf ||y_n(t)||^2$$

Applying Fatou's lemma we obtain

$$\int_{0}^{T} \|y^{*}(t)\|^{2} dt \leq \lim_{n \to +\infty} \inf \int_{0}^{T} \|y_{n}(t)\|^{2} dt.$$
(6)

Taking into account that R is convex and lower semi-continuous with respect to the weak topology, we get (see Corollary III.8 of [6])

$$R(u^*) \le \lim_{n \to +\infty} \inf R(u_n).$$
(7)

Combining the formulas (5), (6) and (7) we deduce that

$$\begin{aligned} J(u^*) &= 2\|y^*(T)\| + \int_0^T \|y(t)\|^2 dt + \int_0^T u^*(t)^2 dt \\ &\leq 2 \lim_{n \to +\infty} \inf \|y_n(T)\| + \lim_{n \to +\infty} \inf \int_0^T \|y_n(t)\|^2 dt + \lim_{n \to +\infty} \inf \int_0^T u_n(t)^2 dt \\ &\leq \lim_{n \to +\infty} \inf J(u_n) \\ &< J^*. \end{aligned}$$

So we conclude that $J(u^*) = J^*$ and hence u^* is a solution of the problem (3).

3 Feedback Optimal Control

In the previous section, we have established an existence result of the quadratic optimal control for the problem (3). However, this result does not provide any information about the expression of the optimal control. This section consists of expressing the quadratic optimal control as an explicit feedback for a class of bilinear systems with B = I. Thus the system (1) becomes

$$\begin{cases} \dot{y}(t) = Ay(t) + u(t)y(t) \\ y(0) = y_0 \in X \end{cases}$$
(8)

Theorem 2. The feedback control defined by

$$u^*(t) = -\|y^*(t)\|,\tag{9}$$

is an optimal control for the problem (3).

Proof: Observing that the mapping $y \to -||y||y$ is locally Lipschitz, we deduce that the system (8), controlled by (9), has a unique mild solution defined on a maximal sub-interval [0, T], which is given by the following variation of constants formula

$$y^{*}(t) = S(t)y_{0} - \int_{0}^{t} S(t-s) \|y^{*}(s)\|y^{*}(s)ds$$

Moreover, the control (9) results in decreasing norm state which implied that the optimal solution is global (see [9] p.185).

Let $v \in U_{ad}$ and let y_v be the respective solution to system (8). Here, we will show that $J(u^*) \leq J(v)$. To this end, two cases will be discussed.

Case 1: $y^*(t) \neq 0$, $\forall t \in [0, T]$. Case 1.1: $y_v(t) \neq 0$, $\forall t \in [0, T]$.

Let $A_n = nA(nI - A)^{-1}$ be the Yoshida approximation of the operator A, and let y_{v_n} be the respective solution to (8) with A_n instead of A. Since the operator A_n is bounded, it follows that $y_{v_n} \in H^1(0,T)$.

Multiplying the system (8) by the state y_{v_n} , and using that A generates an isometric semi-group we deduce that

$$\|y_{v_n}(t)\|\frac{d}{dt}\|y_{v_n}(t)\| = \frac{1}{2}\frac{d}{dt}\|y_{v_n}(t)\|^2 = v(t)\|y_{v_n}(t)\|^2, \quad \forall t \in (0,T).$$

Taking into account that $||y_v(t)|| \neq 0$, for all $t \in [0,T]$ and that $y_{v_n} \longrightarrow y_v$ (strongly) as $n \to +\infty$, we deduce that

$$\exists N \in \mathbb{N}, \forall n \ge N, \ \|y_{v_n}(t)\| \neq 0, \forall t \in [0, T]$$

so that

$$\frac{d}{dt}\|y_{v_n}(t)\| = v(t)\|y_{v_n}(t)\|, \quad \forall t \in (0,T).$$
(10)

Integrating (10) over [0, T], we get

$$2(\|y_{v_n}(T)\| - \|y_0\|) = \int_0^T 2v(t)\|y_{v_n}(t)\|dt$$

= $\int_0^T \left((v(t) + \|y_{v_n}(t)\|)^2 - (v^2(t) + \|y_{v_n}(t)\|^2) \right) dt.$ (11)

Using the fact that $y_{v_n} \longrightarrow y_v$ (strongly) as $n \to +\infty$, we deduce by taking the limit in the relation (11) that

$$J(v) - 2\|y_0\| = \int_0^T (v(t) + \|y_v(t)\|)^2 dt \ge 0.$$
(12)

In particular for $v = u^*$ we get

$$J(u^*) - 2\|y_0\| = \int_0^T (u^*(t) + \|y^*(t)\|)^2 dt.$$
 (13)

This, together with the expression (9), gives

$$J(u^*) - 2\|y_0\| = 0.$$
(14)

Combining (12) and (14) we conclude that

$$J(v) \ge J(u^*).$$

Case 1.2: $\exists t_1 \in (0,T) / y_v(t_1) = 0$.

In this case we have $y_v(t) = 0$, $\forall t \in [t_1, T]$.

Let us define t' by

$$t' = \inf\{t \in [0, T] | y_v(t) = 0\}.$$

By the continuity of the state y_v we deduce that $y_v(t') = 0$ and that $y_v(t) \neq 0$, $\forall t \in [0, t')$.

Following the same method as in the previous case, we conclude that

$$J(v)(t') \ge J(u^*)(t').$$
 (15)

Moreover, it comes from the expression of J(v) that

$$J(v)(T) = \int_0^T (v^2(t) + \|y_v(t)\|^2) dt = J(v)(t') + \int_{t'}^T v^2(t) dt \ge J(v)(t')$$
(16)

Since $y^*(t) \neq 0$ for all $t \in [0, t']$ we deduce, according to Case 1.1, that

$$J(u^*)(T) = J(u^*)(t') = 2||y_0||.$$
(17)

Combining (15), (16) and (17) we conclude that

$$J(v)(T) \ge J(v)(t') \ge J(u^*)(t') = J(u^*)(T).$$

Case 2: $\exists t_1 \in [0,T] \ /y^*(t_1) = 0$. Here we have $y^*(t) = 0, \ \forall t \in [t_1,T]$.

Then, taking $t' = \inf\{t \in [0,T] | y^*(t) = 0\}$, we deduce that $y^*(t') = 0$ and $y^*(t) \neq 0$, $\forall t \in [0,t')$.

According to Case 1 we have

$$J(v)(t') \ge J(u^*)(t'), \forall v \in U_{ad}.$$
(18)

Using the fact that $y^*(t) = 0$, for all $t \in [t', T]$, we derive from the expression of $J(u^*)$ that

$$J(u^*)(T) = \int_0^T (u^{*2}(t) + \|y^*(t)\|^2) dt = \int_0^{t'} (u^{*2}(t) + \|y^*(t)\|^2) dt = J(u^*)(t').$$
(19)

In the sequel, we will show that $J(v)(T) \ge J(v)(t')$, which amounts to showing that

$$\int_{t'}^{t} (v^2(t) + \|y_v(t)\|^2) dt + 2\|y_v(T)\| - 2\|y_v(t')\| \ge 0.$$

Case 2.1: $y_v(t) \neq 0$, $\forall t \in [t', T]$. According to Case 1.1 we have

$$\int_{t'}^{T} (v^2(t) + \|y_v(t)\|^2) dt + 2\|y_v(T)\| - 2\|y_v(t')\| = \int_{t'}^{T} (v(t) + \|y_v(t)\|)^2 dt \ge 0.$$

It follows that

$$J(v)(T) \ge J(v)(t'). \tag{20}$$

Combining (18), (19) and (20) we obtain that

$$J(v)(T) \ge J(u^*)(T).$$

Case 2.2: $\exists t_1 \in [t', T] / y_v(t_1) = 0$.

In this case we have $y_v(t) = 0$, $\forall t \in [t_1, T]$.

Then, letting $t'' = \inf\{t \in [t', T] | y_v(t) = 0\}$, we deduce that $y_v(t'') = 0$ and that $y_v(t) \neq 0, \forall t \in [t', t'')$. Here again we get from Case 1.1

$$J(v)(t'') \ge J(v)(t').$$

Since $y_v(t) = 0$ for all $t \in [t'', T]$, we can see that

$$J(v)(T) = \int_0^T (v^2(t) + \|y_v(t)\|^2) dt = J(v)(t'') + \int_{t''}^T v^2(t) dt \ge J(v)(t'') \cdot$$

We conclude that

$$J(v)(T) \ge J(v)(t'') \ge J(v)(t').$$
 (21)

Combining (18), (19) and (21) we get that

$$J(v)(T) \ge J(u^*)(T).$$

Then, we conclude that u^* is an optimal control for the problem (3).

4 Examples

4.1 Wave Equation

Let us consider the following wave equation

$$\begin{cases} \frac{\partial^2}{\partial t^2} z(t,x) = \Delta z(t,x) + u(t)a(x)z(t,x), & t \in [0,T] \text{ and } x \in \Omega = (0,1) \\ z(t,0) &= z(t,1) = 0, & t \in [0,T] \\ z(0,x) &= z_0(x), & x \in \Omega \end{cases}$$

where $a(.) \in L^{\infty}(\Omega)$ and $u \in L^{2}(0,T)$. This system has the form of the system (1) if we take $X = H_{0}^{1}(\Omega) \times L^{2}(\Omega)$ with $\langle (y_{1}, z_{1}), (y_{2}, z_{2}) \rangle_{X} = \langle y_{1}, y_{2} \rangle_{H_{0}^{1}(\Omega)} + \langle z_{1}, z_{2} \rangle_{L^{2}(\Omega)}$ and

$$A = \begin{pmatrix} 0 & I \\ \Delta & 0 \end{pmatrix} \text{ with } D(A) = H_0^1(\Omega) \cap H^2(\Omega) \times H_0^1(\Omega) \text{ and } B = \begin{pmatrix} 0 & 0 \\ a & 0 \end{pmatrix}$$

Here B is a linear bounded operator on X and A is the infinitesimal generator of a linear C_0 - semi-group S(t) of isometries and $y(t) = (z(t), \dot{z}(t))$.

The quadratic cost function is given by

$$J(u) = 2(\|z(T)\|_{H_0^1(\Omega)}^2 + \|\dot{z}(T)\|_{L^2(\Omega)}^2)^{\frac{1}{2}} + \int_0^T (\|z(t)\|_{H_0^1(\Omega)}^2 + \|\dot{z}(t)\|_{L^2(\Omega)}^2)dt + \int_0^T u(t)^2 dt \cdot (22)$$

According to Theorem 1, there exists an optimal control u^* that minimizes the quadratic cost (22).

4.2 The Transport Equation

Let us consider the following transport problem

$$\begin{cases} \frac{\partial}{\partial t}y(t,x) = -\frac{\partial}{\partial x}y(t,x) + u(t)y(t,x) & t \in [0,T] \text{ and } x \in \Omega = (0,+\infty) \\ y(t,0) &= 0, \ t \in [0,T] \\ y(0,x) &= y_0(x), \ x \in \Omega \end{cases}$$

where $u \in L^2(0,T)$ is the control and $y(t) = y(t,.) \in L^2(\Omega)$ is the state. The operator $A = -\frac{\partial}{\partial x}$ with domain $D(A) = H_0^1(\Omega)$ generates a C_0 semi-group S(t) of isometries on $X = L^2(\Omega)$.

According to Theorem 2, the feedback control u(t) = -||y(t)|| minimizes the following functional cost

$$J(u) = \left(\int_0^{+\infty} y(T, x)^2 dx\right)^{\frac{1}{2}} + \int_0^T \left(\int_0^{+\infty} y(t, x)^2 dx\right) dt + \int_0^T u(t)^2 dt$$

5 Conclusion

In this paper, we investigated the quadratic optimal control problem for a class of bilinear infinite dimensional systems . We formulated optimality conditions in the general case, then we showed that the optimal control can be expressed as a feedback law for a class of bilinear systems. The established results are applied to wave and transport equations. As a natural continuation of the present work is to extend the obtained results to a larger class of bilinear systems, and to study the problems of controllability and stability of bilinear systems using a quadratic optimal control.

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Regional Observability of Linear Fractional Systems Involving Riemann-Liouville Fractional Derivative

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Abstract. In this paper, we study the concept of regional observability, more precisely the regional reconstruction of the initial state of a linear fractional system on a subregion ω of the evolution domain Ω . We use the Hilbert uniqueness method in order to reconstruct the initial state of the given system, which consists of transforming the reconstruction problem into a solvability one. After presenting an algorithm that allows us to reconstruct the regional initial state, we give, at the end, two successful numerical results, in order to backup our theoretical work, each with a different type of sensor and with a reasonable value of error.

1 Introduction

Let Ω be a bounded domain in \mathbb{R}^n , $n \ge 1$, with smooth enough boundary $\partial \Omega$, [0,T] a time interval, $\alpha \in [0,1]$ and A a second order, linear, differential operator. We consider the following fractional system :

$$\begin{cases} {}^{RL} \mathscr{D}_{0^{+}}^{\alpha} y(x,t) = Ay(x,t) & \text{in } \Omega \times [0,T], \\ y(\xi,t) = 0 & \text{on } \partial \Omega \times [0,T], \\ \lim_{t \to 0^{+}} \mathscr{I}_{0^{+}}^{1-\alpha} y(x,t) = y_{0}(x) \text{ in } \Omega, \end{cases}$$
(1)

where ${}^{RL} \mathscr{D}_{0^+}^{\alpha}$ is the left sided Riemann-Liouville fractional derivative of order α . This kind of systems are called fractional diffusion equations or processes, which mean some kind of a diffusion phenomena governed by evolution equations involving fractional derivatives with respect to time and whose solution is given by means of a probability density function [21]. These systems were and are still being widely investigated because, as the theory of continuous time random walks (CTRW) states, they provide a better characterization of anomalous diffusion processes, they also give a better performance compared with conventional diffusion systems [13].

Not only for this kind of systems, Fractional Calculus is a valuable and useful tool especially in modeling real world phenomena in the fields of physics, engineering, aerospace, visco-elasticity, electricity, chemistry, control theory and so forth. For more details see [3,18,23,25].

For instance, in [19], a time-fractional diffusion system for signal smoothing is used, it was mentioned that the fractional model has another adjustable time-fractional

derivative order to control the diffusion process. In the same work, already simulated signals were used in order to compare between the classical diffusion equation and the fractional one. It is claimed that the fractional diffusion filtering, which was applied to nuclear magnetic resonance (NMR) spectrum smoothing, has more advantageous results than those of the classical diffusion filtering, it was also stated that its performance is higher than that of the classical smoothing methods.

Fractional Calculus (FC) is a wide discipline of mathematics which has been around since 300 years ago, the first traces of this subject goes back to Leibniz and L'hospital in their discussion about the meaning of $\frac{d^n}{dx^n}$ if $n = \frac{1}{2}$. The first attempt to give a logical definition is due to Liouville in 1832, and since then a lot of researchers came up with new and different definitions of fractional derivatives, we mention : Riemann-Liouville, Caputo, Riez, Caputo-Fabrizo, Atangana-Baleanu and many other ones. We refer the reader seeking more information about fractional calculus and its properties to see the following books and the references therein [15, 17, 23].

One thing that some find hard to grasp is the initial conditions in a Riemann-Liouville type time-fractional system, which are given as a limit of an integral, we refer any one wondering about this issue to the work of Heymans and Podlubny [14] where they demonstrated that, in fact, one can give a significant physical meaning to such types of initial conditions, and it is very much possible to attribute some values to those kinds of conditions by using appropriate measurements and observations. They also gave a series of concrete examples where these initial conditions make complete sense.

A very important discipline of mathematics that we are dealing with in this paper, is control theory, this field of study plays a serious role in linking between mathematics and technology and it includes several notions such as controllability, stability, observability, stabilization and many more.

In this paper, we deal with observability, precisely the regional observability of a time fractional diffusion system written in terms of Riemann-Liouville time-fractional derivative. The concept of observability was introduced for the first time by the Hungarian-American engineer Rudolf Kalman [16]. This notion has as goal the possibility of finding and reconstructing the initial state of the considered system in a finite time using only the outputs (measurements). This concept has been thoroughly investigated and it also possesses a large literature for various types of system (linear, semilinear...). For more information see [9, 24, 26, 27] and the references therein.

We shall point out the fact that in case of distributed or diffusion systems not all states are observable, hence the necessity of introducing a more weaker notion to cut back the losses for non observable systems, we are speaking about regional observability which also consists of finding and reconstructing the initial state of a system but only in a desired subregion of the evolution domain $\omega \subset \Omega$. Regional observability had seen light for the first time in the nineties with professors El Jai and Afifi for discrete systems see [1], and El Jai and Zerrik for continuous systems see [2, 12]. Afterwards this concept started to be developed by Badraoui, Boutoulout, El Alaoui, Bourray, Zouiten and Torres to cover various types of systems and cases see [4–8, 10, 28, 30].

Lately, regional observability for time-fractional diffusion systems was being studied, see [13], which does not only cover the regional observability but regional analysis in general (regional controllability, regional stability, regional detectability...) for time fractional diffusion processes.

The main goal of this work is to reconstruct the initial state of the considered system using an extension of the Hilbert uniqueness method (HUM), which was firstly introduced by Lions in [20]. This approach relies on the concept of duality for integer order distributed parameter systems, this duality comes from Green's formula. This property fails to work for non-integer order systems, yet we can derive a similar property of duality where the adjoint or dual system is given in terms of Caputo fractional derivative. This relation is obtained with the help of fractional green's formula see [22].

This paper will be organized as follows : After this introduction, we give, in Sect. 2, some preliminary results to be used along this work also as a quick follow up of the considered system. In Sect. 3, we show the formulation and steps of the HUM approach and in Sect. 4, we propose an algorithm that reconstructs the initial state of our system in a desired subregion. As for Sect. 5, we present two successful numerical simulations to back up our work, the first is given with a pointwise sensor and second with a zonal sensor.

2 Considered System and Preliminaries

In this section, we shall introduce some basic definitions needed to present our main result. We give a quick reminder of some necessary notions and properties of fractional calculus followed by other tool of control theory.

We recall the following definitions.

Definition 1 [17]. We call the left sided fractional integral of order $\alpha \in [0, 1]$ of a function y(x,t) at $t \in [0,T]$ for all $x \in \Omega$, the following integral formula :

$$\left(\mathscr{I}_{0^+}^{\alpha}y(x,.)\right)(t) = \frac{1}{\Gamma(\alpha)}\int_0^t (t-s)^{\alpha-1}y(x,s)ds,$$

where $\Gamma(\alpha) := \int_0^{+\infty} t^{\alpha-1} e^{-t} dt$ is the Euler's gamma function.

Definition 2 [17]. We define the left sided Riemann-Liouville fractional derivative of order $\alpha \in [0, 1]$ of a function y(x, t) in $t \in [0, T]$ for all $x \in \Omega$, by :

$$\left({}^{RL}\mathscr{D}^{\alpha}_{0^+}y(x,.)\right)(t):=\frac{d}{dt}\left(\mathscr{I}^{1-\alpha}_{0^+}y(x,.)\right)(t)=\frac{1}{\Gamma(1-\alpha)}\frac{d}{dt}\int_0^t(t-s)^{-\alpha}y(x,s)ds.$$

Definition 3 [17]. The right sided Caputo fractional derivative of order $\alpha \in [0, 1]$ of a function y(x,t) in $t \in [0, 1]$ for all $x \in \Omega$ is given as follows :

$$\binom{C}{\mathscr{D}_{T^{-}}^{\alpha}}y(x,.)(t) = \frac{-1}{\Gamma(1-\alpha)}\int_{t}^{T}(s-t)^{-\alpha}\frac{\partial}{\partial s}y(x,s)ds$$

Let's denote, if there is no confusion,

$$\left(\mathscr{I}_{0^+}^{\alpha}y(x,.)\right)(t) := \mathscr{I}_{0^+}^{\alpha}y(x,t), \quad \left({}^{RL}\mathscr{D}_{0^+}^{\alpha}y(x,.)\right)(t) := {}^{RL}\mathscr{D}_{0^+}^{\alpha}y(x,t) \text{ and} \\ \left({}^{C}\mathscr{D}_{T^-}^{\alpha}y(x,.)\right)(t) := {}^{C}\mathscr{D}_{T^-}^{\alpha}y(x,t).$$

Let $A: D(A) \subseteq L^2(\Omega) \to L^2(\Omega)$ be a second order, linear, differential operator, which generates a C_0 -semigroup $\{S(t)\}_{t\geq 0}$ on $L^2(\Omega)$ and $C: D(C) \subseteq L^2(\Omega) \longrightarrow \mathcal{O}$ a linear, possibly unbounded, operator called the observation operator, where \mathcal{O} is the observation space.

We consider the system (1) augmented with the output equation,

$$z(t) = Cy(.,t), \quad t \in [0,T].$$
 (2)

We give now the definition of the mild solution for the above system.

Definition 4 [29]. We say that a function $y \in C(0,T;L^2(\Omega))$ is a mild solution of (1) if the following formula is satisfied:

$$\mathbf{y}(\mathbf{x},t) = t^{\alpha-1} \mathbf{R}_{\alpha}(t) \mathbf{y}_0(\mathbf{x}), \quad \forall (\mathbf{x},t) \in \Omega \times [0,T],$$
(3)

where

$$egin{aligned} R_{lpha}(t) &= lpha \int_{0}^{+\infty} heta \xi_{lpha}(heta) S(t^{lpha} heta) d heta, \quad t \in [0,T], \ \xi_{lpha}(heta) &= rac{1}{lpha} heta^{-1-rac{1}{lpha}} \, arpi_{lpha}\left(heta^{-rac{1}{lpha}}
ight), \quad heta \in]0, +\infty[. \end{aligned}$$

and

$$\varpi_{\alpha}(\theta) = \frac{1}{\pi} \sum_{n=1}^{+\infty} (-1)^{n-1} \theta^{-n\alpha-1} \frac{n\alpha+1}{n!} \sin(n\pi\alpha), \quad \theta \in]0, +\infty[.$$

Note that ξ_{α} is probability density, that is,

$$\xi_{\alpha}(\theta) \ge 0, \ \forall \theta \in]0, +\infty[$$
 and $\int_{0}^{+\infty} \xi_{\alpha}(\theta) d\theta = 1,$

and that the output function can be written as follows,

$$z(t) = t^{\alpha - 1} C R_{\alpha}(t) y_0(.) = K_{\alpha}(t) y_0.$$

The operator $K_{\alpha}(.): L^{2}(\Omega) \longrightarrow L^{2}([0,T], \mathcal{O})$ appears to be a linear operator, it is called the observability operator and it plays an important role in the characterization of observability.

Remark 1. The operator $K_{\alpha}(.)$ is bounded if *C* is bounded.

In order to study the concept of regional observability we need to use the adjoint operator of $K_{\alpha}(.)$ which is not always defined, precisely in the case when the operator *C* is not bounded, so for $K_{\alpha}^{*}(.)$ to be well defined we need to assume that *C* is an admissible observation operator, then we have the following definition. **Definition 5** [31]. We say that the operator *C* is an admissible observation operator for R_{α} if,

$$\exists M > 0, \text{ such that } \int_0^T \|CR_\alpha(t)x\|_{\mathscr{O}}^2 dt \le M \|x\|_{L^2(\Omega)}^2 \quad \forall x \in D(A).$$

Remark 2. If C is bounded then it is also an admissible observation operator.

For the rest of this work we suppose that *C* is an admissible observation operator, in this case the adjoint operator $K_{\alpha}(.)$ is given as follows,

$$K^*_{\alpha}(.): L^2(0,T; \mathscr{O}) \longrightarrow L^2(\Omega)$$

$$q\longmapsto \int_0^T t^{\alpha-1} R^*_{\alpha}(t) C^*q(t) dt.$$

Let $\omega \subset \Omega$ be a subregion with positive Lebesgue measure, we define the restriction operator in ω by,

$$\chi_{\omega}: L^{2}(\Omega) \longrightarrow L^{2}(\omega),$$
$$y \longmapsto y_{|\omega},$$

and χ^*_{ω} denotes its adjoint, and we have the following definition.

Definition 6 [13]. The system (1) together with the output (2) is said to be approximately regionally observable in ω (or approximately ω -observable) if

$$\mathscr{I}m(\chi_{\omega}K_{\alpha}^{*}(.))=L^{2}(\omega),$$

equivalently

$$\mathscr{K}er(K_{\alpha}(.)\chi_{\omega}^{*}) = \{0\}.$$

We now introduce the notion of sensors which plays an important role in the domain of observability, their main role is to collect data on the studied phenomenon.

Definition 7 [11]. A sensor is a couple (Σ, f) , where Σ is a non empty subset of the evolution domain Ω , it is called the spatial support of the sensor, and f is the spatial distribution.

A sensor is called zonal if $\Sigma \subset \Omega$ is a subset with strictly positive Lebesgue measure, in this case $f \in L^2(\Sigma)$, $\mathscr{O} = \mathbb{R}$ and $z(t) = Cy(.,t) = \langle f, y(.,t) \rangle_{L^2(\Sigma)}$.

A sensor is called pointwise if $\Sigma = \{b\} \in \Omega$, in this case $f = \delta_b$, where δ_b is the Dirac delta function centered at *b*, and the output function is written

$$z(t) = \left\langle \delta_b, y(.,t) \right\rangle_{L^2(\Omega)} = y(b,t).$$

Definition 8 [13]. A sensor (Σ, f) is called strategic if the system (1) augmented with the output function (2), which is given by means of the sensor (Σ, f) is approximately ω -observable. It is called non strategic if not.

We now present one version of the fractional green's formula [22]. $\forall w \in C^{\infty}(0, T \cdot L^2(\Omega))$

$$\int_{0}^{T} \int_{\Omega} \left[{}^{RL} \mathscr{D}_{0^{+}}^{\alpha} y(x,t) - Ay(x,t) \right] \psi(x,t) ds dt = \int_{\Omega} \psi(x,T) \mathscr{I}_{0^{+}}^{1-\alpha} y(x,T) dx + \int_{0}^{T} \int_{\Omega} \left[{}^{C} \mathscr{D}_{T^{-}}^{\alpha} \psi(x,t) - A^{*} \psi(x,t) \right] y(x,t) dx dt - \int_{\Omega} \psi(x,0) \lim_{t \to 0^{+}} \mathscr{I}_{0^{+}}^{1-\alpha} y(x,t) dx$$
(4)
$$+ \int_{0}^{T} \int_{\partial\Omega} \frac{\partial y(\varsigma,t)}{\partial v_{A}} \psi(\varsigma,t) d\varsigma dt - \int_{0}^{T} \int_{\partial\Omega} y(\varsigma,t) \frac{\partial \psi(\varsigma,t)}{\partial v_{A^{*}}} d\varsigma dt.$$

HUM Approach 3

Firstly we define the following set,

$$G = \Big\{ g \in L^2(\Omega) \mid g_{\mid_{\Omega \setminus \omega}} = 0 \Big\},$$

this choice of G is not arbitrary, in fact, we are searching for the value of the initial state in a subregion ω without taking into account the residual part (the value of the initial *state in* $\Omega \setminus \omega$), hence it is natural to consider it null as in the definition of G.

For all $\varphi_0 \in G$, we consider the following linear system,

$$\begin{cases} {}^{RL}\mathscr{D}^{\alpha}_{0^{+}}\varphi(x,t) = A\varphi(x,t) & \text{in } \Omega \times [0,T], \\ \varphi(\xi,t) = 0 & \text{on } \partial \Omega \times [0,T], \\ \lim_{t \mapsto 0^{+}} \mathscr{I}^{1-\alpha}_{0^{+}}\varphi(x,t) = \varphi_{0}(x) \text{ in } \Omega, \end{cases}$$
(5)

the unique mild solution of this system is written,

$$\varphi(x,t) = t^{\alpha-1} R_{\alpha}(t) \varphi_0(x), \quad \forall (x,t) \in \Omega \times [0,T].$$
(6)

We define on $G \times G$ the following bilinear form,

<

$$\langle ,, . \rangle_G : G \times G \longrightarrow \mathbb{C}$$

 $(\varphi_0, g_0) \longmapsto \int_0^T \langle t^{\alpha - 1} C R_\alpha(t) \varphi_0, t^{\alpha - 1} C R_\alpha(t) g_0 \rangle_{\mathscr{O}} dt$

This form satisfies the properties of conjugate symmetry $\left(\langle \varphi_0, g_0 \rangle_G = \overline{\langle g_0, \varphi_0 \rangle_G}\right)$ and positiveness $(\langle \varphi_0, \varphi_0 \rangle_G \ge 0).$

We give the following proposition.

Proposition 1. If the system (5) together with (2) is approximately ω -observable then the form $\langle .,. \rangle_G$ defines a scalar product on G.

Proof. All that remains is to prove the definiteness of $\langle ., . \rangle_G$, (*i.e* : $\langle \varphi_0, \varphi_0 \rangle_G = 0 \implies \varphi_0 = 0$), in fact

$$\langle \varphi_0, \varphi_0 \rangle_G = \int_0^T \langle t^{\alpha - 1} C R_\alpha(t) \varphi_0, t^{\alpha - 1} C R_\alpha(t) \varphi_0 \rangle_{\mathscr{O}} dt = \int_0^T \| t^{\alpha - 1} C R_\alpha(t) \varphi_0 \|_{\mathscr{O}}^2 dt,$$

hence

$$\langle \varphi_0, \varphi_0 \rangle_G = 0 \implies \| t^{\alpha - 1} C R_{\alpha}(t) \varphi_0 \|_{\mathscr{O}}^2 = 0, \ \forall t \in [0, T],$$

then

$$t^{\alpha-1}CR_{\alpha}(t)\varphi_0=0, \ \forall t\in[0,T],$$

which gives

$$t^{\alpha-1}CR_{\alpha}(t)\chi_{\omega}^{*}\chi_{\omega}\varphi_{0}=K_{\alpha}(t)\chi_{\omega}^{*}(\chi_{\omega}\varphi_{0})=0, \forall t\in[0,T],$$

and since the system (5) is approximately ω -observable we have that

$$\chi_{\omega}\varphi_0=0,$$

thus

 $\varphi_0 = 0, \quad \text{in } \omega,$

and by definition of G, we have

$$\varphi_0=0, \quad \text{ in } \Omega \setminus \omega,$$

finally

$$arphi_0=0, \quad ext{ in } \Omega.$$

For the rest of this work, we assume that the system (5) – (2) is approximately ω -observable, hence $\langle ., . \rangle_G$ is a scalar product and we denote by $\|.\|_G := \sqrt{\langle ., . \rangle_G}$ the natural norm on *G* based upon the scalar product of *G*.

With the help of the fractional green's formula, see equation(4), we derive the following retrograded system,

$$\begin{cases} {}^{C}D_{T^{-}}^{\alpha}\phi(x,t) = A^{*}\phi(x,t) + C^{*}C\phi(.,t) & \text{in } \Omega \times [0,T], \\ \phi(\xi,t) = 0 & \text{on } \partial\Omega \times [0,T], \\ \phi(x,T) = 0 & \text{in } \Omega, \end{cases}$$
(7)

which has a unique mild solution $\phi \in C(0,T;L^2(\Omega))$, given by the following integral formula,

$$\phi(t) = \int_t^T (\tau - t)^{\alpha - 1} R^*_\alpha(\tau - t) C^* C \phi(\tau) d\tau.$$
(8)

if φ_0 is chosen in *G* such that $C\varphi(t) = z(t)$, then the following system

$$\begin{cases} {}^{C}D_{T^{-}}^{\alpha}\psi(x,t) = A^{*}\psi(x,t) + C^{*}z(t) \quad \text{in } \Omega \times [0,T], \\ \psi(\xi,t) = 0 \qquad \qquad \text{on } \partial\Omega \times [0,T], \\ \psi(x,T) = 0 \qquad \qquad \text{in } \Omega, \end{cases}$$
(9)

can be seen as the adjoint system of (5).

We define the mapping

$$egin{aligned} \Lambda: G \longrightarrow G \ & arphi_0 \longmapsto \chi^*_\omega \chi_\omega \left(\psi(0)
ight), \end{aligned}$$

hence the problem of regional reconstruction is reduced to solving the following equation,

$$\Lambda \varphi_0 = \chi_\omega^* \chi_\omega \left(\psi(0) \right). \tag{10}$$

Remark 3. $\chi_{\omega}^* \chi_{\omega}$ is projection operator on *G*.

We have the following theorem,

Theorem 1. If the system (5) together with (2) is approximately ω -observable then the equation (10) has a unique solution which corresponds with the initial state in ω .

Proof. Let's consider $\varphi_0 \in G$, we have :

$$\begin{split} \langle \Lambda \, \varphi_0, \varphi_0 \rangle_G &= \langle \chi^*_{\omega} \chi_{\omega}(\psi(0)), \varphi_0 \rangle_G, \\ &= \langle \psi(0), \varphi_0 \rangle_G, \\ &= \langle \int_0^T t^{2\alpha - 2} R_{\alpha}(t)^* C^* C R_{\alpha}(t) \varphi_0 dt, \varphi_0 \rangle_G, \\ &= \int_0^T \langle t^{\alpha - 1} C R_{\alpha}(t) \varphi_0, t^{\alpha - 1} C R_{\alpha}(t) \varphi_0 \rangle_{\mathscr{O}} dt, \\ &= \langle \varphi_0, \varphi_0 \rangle_G \\ &= \| \varphi_0 \|_G^2. \end{split}$$

Thus Λ is an isomorphism.

4 Algorithm

This section is reserved to the proposed algorithm for the reconstruction of the initial state in ω , this leads to some numerical results which will be presented in the next section.

In order to give an algorithm that reconstructs the initial state, we assume that the operator *A* generates a complete system of eigenfunctions $\{\varphi_i\}_{i \in \mathbb{N}^*}$ on the state space $L^2(\Omega)$ associated with the eigenvalues $\{\lambda_i\}_{i \in \mathbb{N}^*}$.

Note that the family $\{\varphi_i\}_{i \in \mathbb{N}^*}$ is an orthonormal basis of $L^2(\Omega)$. In this case, R_{α} and φ are respectively expressed as follows, $\forall w \in L^2(\Omega), \forall (x,t) \in \Omega \times [0,T]$:

$$\begin{split} R_{\alpha}(t)w &= \sum_{i=1}^{+\infty} E_{\alpha,\alpha}\left(\lambda_{i}t^{\alpha}\right) \left\langle w, \varphi_{i} \right\rangle_{L^{2}(\Omega)} \varphi_{i}(.).\\ \varphi(x,t) &= \sum_{i=1}^{+\infty} t^{\alpha-1} E_{\alpha,\alpha}\left(\lambda_{i}t^{\alpha}\right) \left\langle \varphi_{0}, \varphi_{i} \right\rangle_{L^{2}(\Omega)} \varphi_{i}(x), \end{split}$$

where, $E_{\alpha,\beta}(t) := \sum_{n=0}^{+\infty} \frac{t^n}{\Gamma(n\alpha + \beta)}$, is the two parameter Mittag Leffler function. The solution of (9) at t = 0 can be written as follows $\forall r \in \Omega$

The solution of (9) at t = 0 can be written as follows, $\forall x \in \Omega$,

$$\psi(x,0) = \sum_{i=1}^{+\infty} \int_0^T \tau^{\alpha-1} E_{\alpha,\alpha} \left(\lambda_i \tau^{\alpha} \right) \left\langle C^* z(\tau), \varphi_i \right\rangle_{L^2(\Omega)} d\tau \varphi_i(x)$$

From theorem 1, we have

$$\left\langle \Lambda \varphi_0, \varphi_0 \right\rangle_G = \left\| \varphi_0 \right\|_G^2 = \int_0^T \left\| C \varphi(.,t) \right\|_{\mathscr{O}}^2 dt$$

Case 1 : *I*f the measurements are given by mean of a pointwise sensor (b, δ_b) :

$$\begin{split} \left\langle \Lambda \varphi_{0}, \varphi_{0} \right\rangle_{G} &= \sum_{i,j=1}^{+\infty} \int_{0}^{T} t^{2\alpha-2} E_{\alpha,\alpha} \left(\lambda_{i} t^{\alpha} \right) E_{\alpha,\alpha} \left(\lambda_{j} t^{\alpha} \right) dt \varphi_{i}(b) \varphi_{j}(b) \\ &\times \left\langle \varphi_{0}, \varphi_{i} \right\rangle_{L^{2}(\Omega)} \left\langle \varphi_{0}, \varphi_{j} \right\rangle_{L^{2}(\Omega)}. \end{split}$$

We set

$$\Lambda_{ij} = \int_0^T t^{2\alpha - 2} E_{\alpha, \alpha} \left(\lambda_i t^{\alpha} \right) E_{\alpha, \alpha} \left(\lambda_j t^{\alpha} \right) dt \varphi_i(b) \varphi_j(b), \quad \forall i, j = 1, ..., \infty$$

Case 2 : If the measurements are given by mean of a zonal sensor $(D, f) : \langle \Lambda \varphi_0, \varphi_0 \rangle_G = \sum_{i,j=1}^{+\infty} \int_0^T t^{2\alpha-2} E_{\alpha,\alpha} \left(\lambda_i t^{\alpha} \right) E_{\alpha,\alpha} \left(\lambda_j t^{\alpha} \right) dt \langle f, \varphi_i \rangle_{L^2(D)} \times \langle f, \varphi_j \rangle_{L^2(D)} \langle \varphi_0, \varphi_i \rangle_{L^2(\Omega)} \langle \varphi_0, \varphi_j \rangle_{L^2(\Omega)}.$

We set

$$\Lambda_{ij} = \int_0^T t^{2\alpha-2} E_{\alpha,\alpha} \left(\lambda_i t^{\alpha}\right) E_{\alpha,\alpha} \left(\lambda_j t^{\alpha}\right) dt \langle f, \varphi_i \rangle_{L^2(D)} \langle f, \varphi_j \rangle_{L^2(D)}, \quad \forall i, j = 1, ..., \infty.$$

The problem (10) can be written now as

$$AX = b, \quad \text{where} \quad A \in \mathscr{M}_{N,N}(\mathbb{C}), \ X \in \mathscr{M}_{N,1}(\mathbb{C}) \text{ and } b \in \mathscr{M}_{N,1}(\mathbb{C}), \qquad (11)$$

such that,

$$A_{ij} = \Lambda_{ij}, \quad X_i = \langle \varphi_0, \varphi_i \rangle_{L^2(\Omega)} \quad \text{ and } \ b_j = \langle \chi_{\omega}^* \chi_{\omega}(\psi(0)), \varphi_j \rangle_{L^2(\Omega)}$$

After resolving the system (11), we obtain the reconstructed initial state. Then we give the following algorithm.

Algorithm

- Initialization of : ε , α , ω , Sensors, y_0 .
- Repeat
 - \rightarrow Solve (9), and get ψ
 - ightarrow Calculate the components of Λ
 - \rightarrow Solve the system (11) and get φ_0 .

• Until
$$\|y_0 - \varphi_0\|_{L^2(\omega)} \leq \varepsilon.$$

5 Numerical Results

In this section, we show some numerical illustrations of our result. We will present two examples with the same system but with different output functions, the first one will be given with a pointwise senor whereas for the second one, measurements are given by a zonal sensor.

Pointwise Sensor

Let us consider the following time-fractional system,

$$\begin{cases} {}^{RL}\mathscr{D}_{0^{+}}^{0.5}y(x,t) = \frac{\partial^{2}}{\partial x^{2}}y(x,t) \text{ in } [0,1] \times [0,2], \\ y(0,t) = y(1,t) = 0 \quad \text{ in } [0,2], \\ \lim_{t \mapsto 0^{+}} \mathscr{I}_{0^{+}}^{0.5}y(x,t) = y_{0}(x) \quad \text{ in } [0,1], \end{cases}$$
(12)

The operator $\frac{\partial^2}{\partial x^2}$ has a complete set of eigenfunctions $\varphi_i(x) = \sqrt{2}\sin(i\pi x)$ with the corresponding eigenvalues $\lambda_i = -i^2\pi^2$.
The system (12) is augmented by the output equation given by means of a pointwise sensor localized at b = 0.72,

$$z(t) = y(b,t), \quad t \in [0,T],$$

we consider the region $\omega = [0.35, 0.65]$ and the initial state (Supposed to be unknown)

$$y_0(x) = 2x(x-1)(2x-1).$$

By applying the proposed algorithm, we obtain Fig. 1.



Fig. 1. Initial and estimated initial state in $\omega = [0.35, 0.65]$.

The reconstruction error is $\left\|y_0 - \varphi_0\right\|_{L^2(\omega)} = 1.0410 \times 10^{-4}.$

The Fig. 1 shows that the estimated and real initial state are near one another in the subregion ω .

Figure 2, shows the evolution of the reconstruction error in terms of the sensor's location.

Zonal Sensor

We consider, in this example, the same system (12) but with $\alpha = 0.4$ and measurements given by a zonal sensor (D, f) where D = [0.25, 0.35] and $f \equiv 1$. The output function is given as follows

$$z(t) = \langle f, y(.,t) \rangle_{L^2(D)}.$$

The considered subregion is $\omega = [0.2, 0.6]$ and the initial state (Supposedly unknown) is $y_0(x) = (e^x - 1)ln(2 - x)$.



Fig. 2. Evolution of the reconstruction error in function of the sensor's location.



Fig. 3. Initial and estimated initial state in $\omega = [0.2, 0.6]$.

The reconstruction error is $\|y_0 - \varphi_0\|_{L^2(\omega)} = 3.2 \times 10^{-3}$. We also see, in Fig. 3, that the reconstructed initial state and the initial one are besides each other in the desired subregion ω .

The following table shows the evolution of the reconstruction error in function of the geometric domain of the sensor.

Geometric domain of the sensor	Error
[0.15, 0.25]	1.8×10^{-3}
[0.25, 0.35]	$3.2 imes 10^{-3}$
[0.35, 0.45]	1.9×10^{-3}
[0.45, 0.55]	4.77×10^{11}
[0.55, 0.65]	$3.5 imes 10^{-3}$
[0.65, 0.75]	$2.9 imes 10^{-3}$
[0.75, 0.85]	$2.6 imes 10^{-3}$
[0.85, 0.95]	$3.4 imes 10^{-3}$

Table 1. Evolution of the reconstruction error in function of the geometric domain of the sensor.

The same remark about the position of the sensor applies here. We see that if the sensor is placed in [0.45, 0.55] then the sensor is not strategic.

6 Conclusion

In this work, we studied the regional reconstruction of the initial state for a Riemann-Liouville type time-fractional diffusion system, and for that we adopted an extension of the Hilbert uniqueness method, which was introduced by the french mathematicians Jacques-Louis Lions. We supposed the admissibility condition on the observation operator, which is necessary when it is unbounded, and that the considered system is approximately regionally observable, which is also necessary, because we can't reconstruct the initial state, regionally, if the system isn't at least approximately regionally observable. We proposed an algorithm that helps us achieve our goal and at the end we gave two successful numerical results, with satisfying errors of reconstruction, to valid our algorithm. For the future, we are working on the concept of regional observability for semilinear fractional diffusion systems, we also plan on making the algorithm better in order to get a lesser error.

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Stability Analysis of Fractional Differential Systems Involving Riemann–Liouville Derivative

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Abstract. We introduce the stability notion of the fractional differential systems under Riemann-Liouville time derivative of order $\alpha \in (0, 1)$, evolving on a spatial domain Ω . Then, we characterize the asymptotic behavior of the state. Also, we present sufficient and necessary conditions to achieve the exponential stability of this important class of systems. Hence, we study the state stabilization of fractional differential systems by means of decomposition method. Several examples and simulations are given to show the applicability of our presented results.

1 Introduction

The fractional calculus history dates back to the Seventeenth century (more precisely to 1695), when the possible meaning of the half order differentiation was discussed by Marquis de L'Hospital and Gottfried Wilhelm Leibniz. Since then, this question has been studied by many well-known mathematicians over the years, such as: Euler, Feller, Fourier, Laurent, Letnikov, Liouville, Grünewald, Riemann and many others. However, the fractional calculus theory has been evolved speedily since the Nineteenth century, mainly as a foundation of several mathematical branches such as, fractional differential systems and fractional geometry. Moreover, during last decades, the investigation, especially, of the fractional calculus one. Hence, nowadays, the fractional differential systems have been proved as a powerful tool to characterize various dynamical models and many real world systems. For example, it can be mentioned viscoelastic systems [1], diffusion and some heat transfer process [5], amongst others.

Stability analysis is an interesting concept in differential systems and control theories, as well as in their applications. Indeed, the stability of integer order differential systems was widely studied ([2, 13, 16]). Many approaches were used to establish several degrees of their stability and stabilization: the asymptotic and exponential stability have been considered using Lyapunov equation [13]. Also, the strong stabilization has been treated by means of Riccati equation [4]. Moreover, the exponential stabilization has been developed using a specific state space and system decomposition [16]. Recently, the stability was introduced

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to fractional calculus [11,14,18–20]. Furthermore, many studies investigated the stability of fractional order systems [17]. For example Qian et al have been developed some stability analytical results for Riemann-Liouville fractional systems of order $\alpha \in (0, 1)$, including perturbed systems, linear systems and time-delayed systems [14]. Also, a fractional Lyapunov direct method has been proposed, by Li et al, to study the power law stability and the exponential stability [11]. Moreover, in [6], Ge et al introduced the regional stability notion for Riemann–Liouville linear fractional differential systems, where they charcterized the strong stability using the spectrum properties of the system dynamic and the strong stabilization via decomposition method.

In this work, we shall present some new stability and stabilization theorems for Riemann–Liouville linear fractional differential systems of order $\alpha \in (0, 1)$. In details, we formulate the problem and we characterize the asymptotic and exponential stability of such class of systems in Sect. 2. In Sect. 3, we derive the fractional differential systems stabilization using especially decomposition method. Finally, we give a conclusion in the last section which contains a synthesis and some perspectives.

2 Stability of Fractional Differential Systems

In this section we consider, in $\Omega \subset \mathbb{R}^n$ (n = 1, 2, 3, ...), an open bounded subset with a regular boundary $\partial \Omega$, the following fractional diffusion system, defined as

$$\begin{cases} {}^{RL}_{0}D^{\alpha}_{t}z(x,t) = Az(x,t), & x \in \Omega, \ t \in]0, +\infty[\\ z(\eta,t) = 0, & \eta \in \partial\Omega, \ t \in]0, +\infty[\\ \lim_{t \longrightarrow 0^{+}} {}^{0}I^{1-\alpha}_{t}z(x,t) = z_{0}(x), \ x \in \Omega, \end{cases}$$
(1)

where $A : D(A) \subset L^2(\Omega) \longrightarrow L^2(\Omega)$ is a linear operator that generates a C_0 -semi-group $(S(t))_{t\geq 0}$ [3] on $L^2(\Omega)$, ${}_0^{RL}D_t^{\alpha}$ and ${}_0I_t^{\alpha}$ denote, respectively, the Riemann-Liouville derivative and integral of order $\alpha \in (0,1)$ [10], that are given by

$${}_{0}I_{t}^{\alpha}z(.,t) = \Gamma(\alpha)^{-1} \int_{0}^{t} (t-v)^{\alpha-1}z(.,v) \,\mathrm{d}v$$

and

$${}_{0}^{RL}D_{t}^{\alpha}z(.,t) = \frac{d}{dt}{}_{0}I_{t}^{1-\alpha}z(.,t),$$

with $\Gamma(\alpha) = \int_0^{+\infty} y^{\alpha-1} e^{-y} \, \mathrm{d}y$ is the Gamma function. The mild solution $z \in C(0, T, L^2(\Omega))$ of system (1) [7] is a

The mild solution $z \in C(0, T, L^2(\Omega))$ of system (1) [7] is defined as

$$z(.,t) = H_{\alpha}(t)z_{0}(.) = t^{\alpha-1}K_{\alpha}(t)z_{0}(.), \qquad (2)$$

where

$$K_{\alpha}(t) = \alpha \int_{0}^{+\infty} \xi \phi_{\alpha}(\xi) S(t^{\alpha}\xi) \,\mathrm{d}\xi \tag{3}$$

and

$$\phi_{\alpha}(\xi) = \alpha^{-1} \xi^{-1 - \frac{1}{\alpha}} P_{\alpha}(\xi^{-\frac{1}{\alpha}}), \tag{4}$$

with

$$P_{\alpha}(\xi) = \pi^{-1} \sum_{n=1}^{+\infty} (-1)^n \frac{\Gamma(n\alpha+1)}{n!} \xi^{\alpha n-1} \sin(n\pi\alpha), \ \xi \in (0,\infty).$$

Remark 1. $P_{\alpha}(.)$ represents the function of the probability density.

First, we state some stability definitions.

Definition 1. System (1) is said to be

• Exponentially stable, if for all $z_0 \in L^2(\Omega)$ there exist Q and $\sigma > 0$ satisfying

$$||z(.,t)|| \le Qe^{-\sigma t} ||z_0||, \ \forall t \ge 0.$$

• Strongly stable, if for all $z_0 \in L^2(\Omega)$ the corresponding solution z(.,t) of (1) fulfills

$$\lim_{t \longrightarrow +\infty} \|z(.,t)\| = 0.$$

Remark 2. The exponential stability implies the strong one.

In the following theorem we present the link between the strong stability of the fractional differential system (1) and the spectrum properties of its dynamic A.

Let us introduce the sets

$$\sigma^{1}(A) = \left\{ \lambda \in \sigma(A) : |arg(\lambda)| \le \frac{\alpha \pi}{2} \right\}$$

and

$$\sigma^2(A) = \left\{ \lambda \in \sigma(A) : |arg(\lambda)| > \frac{lpha \pi}{2} \right\},$$

where $\sigma(A)$ indicates the points spectrum of the operator A.

Theorem 1. Let $(\lambda_n)_{n\geq 1}$ and $(\chi_n)_{n\geq 1}$ be the eigenvalues and the corresponding eigenfunctions of the operator A, with $(\chi_n)_{n\geq 1}$ form an orthonormal basis on $L^2(\Omega)$. If $\sigma^1(A) = \emptyset$ and $\forall \lambda_n \in \sigma^2(A)$, n = 1, 2, ..., there exists $\varepsilon > 0$ satisfying $\lambda_n \leq -\varepsilon$, then the system (1) is strongly stable in Ω .

Proof. For $z_0 \in L^2(\Omega)$, the solution of (1) [7] can be written as

$$z(.,t) = t^{\alpha-1} \sum_{n=1}^{+\infty} E_{\alpha,\alpha}(\lambda_n t^{\alpha}) \langle z_0, \chi_n \rangle \chi_n(.), \ \forall z_0 \in L^2(\Omega),$$
(5)

where

$$E_{\alpha,\alpha}(\lambda_n t^{\alpha}) = \sum_{k=0}^{+\infty} \frac{(\lambda_n t^{\alpha})^k}{\Gamma(\alpha k + \alpha)}$$

From (5), one has

$$||z(.,t)||^{2} = t^{2(\alpha-1)} \sum_{n=1}^{+\infty} (E_{\alpha,\alpha}(\lambda_{n}t^{\alpha}))^{2} \langle z_{0}, \chi_{n} \rangle^{2}.$$

Also, using the fact that $\sigma^1(A) = \emptyset$ and $\lambda_n \leq -\varepsilon$ for all $\lambda_n \in \sigma^2(A)$, and the Mittag-Leffler function $E_{\alpha,\alpha}(-x)$, $x \geq 0$, is completely monotonic [15], yields

$$||z(.,t)|| \le t^{\alpha-1} E_{\alpha,\alpha}(-\varepsilon t^{\alpha}) ||z_0||.$$

It follows, since $|E_{\alpha,\alpha}(-\varepsilon t^{\alpha})| \leq 1$ for $\alpha \in (0,1)$ [8], that

$$||z(.,t)|| \longrightarrow 0 \quad as \quad t \longrightarrow +\infty.$$

Example 1. Let's consider the sub-diffusion system

$$\begin{cases} {}_{0}^{RL}D_{t}^{0.6}z(x,t) = \frac{\partial^{2}}{\partial x^{2}}z(x,t), & x \in \Omega, t \in]0, +\infty[\\ z(\eta,t) = 0, & \eta \in \partial\Omega, t \in]0, +\infty[\\ \lim_{t \longrightarrow 0^{+}} I_{t}^{0.4}z(x,t) = x^{3}(x-1), x \in \Omega, \end{cases}$$
(6)

with $\Omega =]0,1[$. According to (6), we get that the dynamic $A = \frac{\partial^2}{\partial x^2}$, with the eigenvalues being

$$\lambda_n = -n^2 \pi^2, \ n \ge 1 \tag{7}$$

and the corresponding eigenfunctions being

$$\chi_n(x) = \sqrt{2}\sin(n\pi x), \ n \ge 1.$$

The solution of system (6) is defined by

$$z(x,t) = t^{0.4} \sum_{n=1}^{+\infty} E_{0.6}(\lambda_n t^{0.6}) \langle z_0, \chi_n \rangle \chi_n(x).$$

One has, for all $n \ge 1$, that

$$|arg(\lambda_n)| = \pi > \frac{\alpha\pi}{2} = \frac{3\pi}{10},$$

which implies that $\sigma^1(A) = \emptyset$ and $\sigma^2(A) = \{-n^2 \pi^2, n \ge 1\}.$

Also, from (7), one has that

$$\lambda_n \leq -\pi^2, \ n \geq 1.$$

so, for all $\lambda_n \in \sigma^2(A)$, there exist $\varepsilon = \pi^2 > 0$ such that

$$\lambda_n \leq -\varepsilon, \ n \geq 1.$$

Hence, by applying the above Theorem, we get the strong stability of system (6) as it is illustrated by Fig. 1.



Fig. 1. The state z(x, t) behavior of system (6) at t = 0, t = 10, t = 16.

We shall use the next lemma to study the exponential stability for system (1).

Lemma 1. Suppose there exists a function $R(.) \in L^2(0, +\infty; \mathbb{R}^+)$ fulfilling

$$\|H_{\alpha}(t+\tau)z\| \le R(t)\|H_{\alpha}(\tau)z\|, \forall t, \tau \ge 0,$$
(8)

for all $z \in L^2(\Omega)$, then the operators $(H_{\alpha}(t))_{t\geq 0}$ are uniformly bounded. Proof. To show the boundedness of $(H_{\alpha}(t))_{t\geq 0}$, we prove that

$$\sup_{t\geq 0} \|H_{\alpha}(t)z\| < \infty, \ \forall z \in L^{2}(\Omega).$$

Otherwise, there exists a sequence $(t_1 + r_n), t_1 > 0$ and $r_n \longrightarrow +\infty$ with

$$||H_{\alpha}(t_1 + r_n)z|| \longrightarrow +\infty \ as \ n \longrightarrow +\infty.$$
(9)

From the following relation

$$\int_0^{+\infty} \|H_\alpha(\tau+r_n)z\|^2 \,\mathrm{d}\tau = \int_{r_n}^{+\infty} \|H_\alpha(\tau)z\|^2 \,\mathrm{d}\tau, \quad \underset{n \longrightarrow +\infty}{\longrightarrow} 0$$

and by Fatou's Lemma, it follows that

$$\lim_{n \to +\infty} \inf \|H_{\alpha}(\tau + r_n)z\| = 0,$$

almost everywhere $0 \leq \tau < +\infty$.

Thus, for some $\tau_0 < t_1$ we can find a subsequence r_{n_p} such that

$$\lim_{p \to +\infty} \|H_{\alpha}(\tau_0 + r_{n_p})z\| = 0.$$
 (10)

Moreover, by virtue of (8), one has

$$\|H_{\alpha}(t_1 + r_{n_p})z\| \le R(t_1 - \tau_0) \|H_{\alpha}(\tau_0 + r_{n_p})z\|.$$
(11)

Then, combining (11) and (10), one obtains

$$||H_{\alpha}(t_1 + r_{n_p})z|| \longrightarrow 0 \ as \ p \longrightarrow +\infty, \tag{12}$$

which is absurd. Hence, using the principale of the uniform boundedness, we get the stated result.

Theorem 2. Assume that the operator $(H_{\alpha}(t))_{t\geq 0}$ satisfies the condition (8) and the inequality

$$||H_{\alpha}(t+\tau)z|| \le ||H_{\alpha}(t)z||.||H_{\alpha}(\tau)z||, \quad \forall t, \tau \ge 0,$$
(13)

holds for all $z \in L^2(\Omega)$, then the system (1) is exponentially stable, if and only if

$$\int_{0}^{+\infty} \|H_{\alpha}(t)z\|^{2} \,\mathrm{d}t < \infty, \ \forall z \in L^{2}(\Omega).$$

$$(14)$$

Proof. Let us show that

$$\sigma_0 = \lim_{t \longrightarrow +\infty} \frac{\ln \|H_\alpha(t)\|}{t}.$$
(15)

We have, for all $t \ge 0$, the following relation

$$t \|H_{\alpha}(t)z\|^{2} = \int_{0}^{t} \|H_{\alpha}(t)z\|^{2} d\tau$$
$$= \int_{0}^{t} \|H_{\alpha}(\tau + t - \tau)z\|^{2} d\tau$$

Using (13), yields

$$t \|H_{\alpha}(t)z\|^{2} \leq \int_{0}^{t} \|H_{\alpha}(t-\tau)z\|^{2} \|H_{\alpha}(\tau)z\|^{2} \,\mathrm{d}\tau.$$

Since the operator $H_{\alpha}(t)$ is bounded for all $t \geq 0$, and by virtue of (14), one gets

 $t \|H_{\alpha}(t)z\|^2 \le \xi \|z\|^2$, for some $\xi > 0$,

moreover, for t sufficiently large, yields

$$\|H_{\alpha}(t)\| < 1,$$

hence, there exists $t_1 > 0$ satisfying

$$\ln \|H_{\alpha}(t)\| < 0,$$

for all $t \geq t_1$. Then,

$$\sigma_0 = \inf_{t \ge 0} \frac{\ln \|H_{\alpha}(t)\|}{t} < 0.$$

Furthermore, let $S_p = \sup_{t \in [0,t_1]} ||H_{\alpha}(t)||$ with $t_1 > 0$ is fixed. Thus, for every $t > t_1$ we may find an integer $\beta \ge 0$ such that $\beta t_1 \le t \le (\beta + 1)t_1$.

From (13), yields

$$|H_{\alpha}(t)|| = ||H_{\alpha}(\beta t_1 + (t - \beta t_1))||$$

$$\leq ||H_{\alpha}(\beta t_1)|||H_{\alpha}(t - \beta t_1)||$$

which implies that

$$\frac{\ln \|H_{\alpha}(t)\|}{t} \le \frac{\ln \|H_{\alpha}(\beta t_1)\|}{t} + \frac{\ln \|H_{\alpha}(t-\beta t_1)\|}{t}$$

using again condition (13), we obtain that

$$\frac{\ln \|H_{\alpha}(t)\|}{t} \le \frac{\beta t_1}{t} \frac{\ln \|H_{\alpha}(t_1)\|}{t_1} + \frac{\ln \|S_p\|}{t},$$

taking into acount that t_1 is arbitrary and $\frac{\beta t_1}{t} \leq 1$, it follows

$$\lim_{t \longrightarrow +\infty} \sup_{t \to \infty} \frac{\ln \|H_{\alpha}(t)\|}{t} \le \inf_{t > 0} \frac{\ln \|H_{\alpha}(t)\|}{t} \le \liminf_{t \longrightarrow +\infty} \frac{\ln \|H_{\alpha}(t)\|}{t},$$

which implies that (15) is satisfied.

Then, for all $\sigma \in [0, -\sigma_0]$, there exists Q > 0 such that

$$\|H_{\alpha}(t)z\| \le Qe^{-\sigma t}\|z\|,$$

for all $z \in L^2(\Omega)$ and $t \ge 0$.

The converse implication of the theorem is immediate.

Remark 3. When $\alpha = 1$, we retrieve the exponential stability result established in [2].

3 Stabilization of Fractional Differential Systems

In this section we investigate the strong stabilization of time fractional differential systems under Riemann-Liouville derivative of order $\alpha \in (0, 1)$, described by

$$\begin{cases} R^{L} D^{\alpha}_{t} z(x,t) = A z(x,t) + B u(x,t), x \in \Omega, \ t \in]0, +\infty[\\ z(\eta,t) = 0, & \eta \in \partial\Omega, \ t \in]0, +\infty[\\ \lim_{t \longrightarrow 0^{+}} 0I^{1-\alpha}_{t} z(x,t) = z_{0}(x), & x \in \Omega, \end{cases}$$
(16)

where the operator A is defined as in system (1), the operator B is linear and bounded from X into $L^2(\Omega)$, with $\Omega \subset \mathbb{R}^n$ is an open bounded subset and X is a Hilbert space of controls, and $u \in L^2(0, +\infty, X)$. **Definition 2.** The system (16) is said to be strongly stabilizable if there exists a bounded operator $K \in \mathcal{L}(L^2(\Omega), X)$ such that the system

$$\begin{cases} {}^{RL}_{0}D^{\alpha}_{t}z(x,t) = (A+BK)z(x,t), x \in \Omega, t \in]0, +\infty[\\ z(\eta,t) = 0, & \eta \in \partial\Omega, t \in]0, +\infty[\\ \lim_{t \longrightarrow 0^{+}} 0I^{1-\alpha}_{t}z(x,t) = z_{0}(x), & x \in \Omega \end{cases}$$
(17)

is strongly stable in Ω .

The solution of system (17) is defined by

$$z(.,t) = t^{\alpha-1} K_{\alpha}^{K}(t) z_{0}(.),$$

with

$$K_{\alpha}^{K}(t) = \alpha \int_{0}^{+\infty} \xi \phi_{\alpha}(\xi) S^{K}(t^{\alpha}\xi) \,\mathrm{d}\xi,$$

where $\phi_{\alpha}(.)$ is given by (4) and $(S^{K}(t))_{t\geq 0}$ is the semi-group generated by A + BK.

3.1 Characterization of Stabilization

We have the following theorem.

Theorem 3. Let $(\lambda_n^K)_{n\geq 1}$ and $(\chi_n^K)_{n\geq 1}$ be the eigenvalues and the corresponding eigenfunctions of the operator A + BK, with $(\chi_n^K)_{n\geq 1}$ form an orthonormal basis on $L^2(\Omega)$. If $\sigma^1(A + BK) = \emptyset$ and $\forall \lambda_n^K \in \sigma^2(A + BK)$, n = 1, 2, ..., there exists $\varepsilon > 0$ satisfying $\lambda_n^K \leq -\varepsilon$, then the system (16) is strongly stabilizable in Ω by the control

$$u(x,t) = Kz(x,t).$$
(18)

Proof. The system (1) admits a unique mild solution [7] given by

$$z(.,t) = t^{\alpha-1} \sum_{n=1}^{+\infty} E_{\alpha,\alpha}(\lambda_n^K t^\alpha) \langle z_0, \chi_n^K \rangle \chi_n^K(.), \ \forall z_0 \in L^2(\Omega),$$

It follows

$$||z(.,t)||^{2} = t^{2(\alpha-1)} \sum_{n=1}^{+\infty} (E_{\alpha,\alpha}(\lambda_{n}^{K}t^{\alpha}))^{2} \langle z_{0}, \chi_{n} \rangle^{2}.$$

Also, Using the fact that $\sigma^1(A) = \emptyset$ and $\lambda_n^K \leq -\varepsilon$ for all $\lambda_n^K \in \sigma^2(A)$, and the Mittag-Leffler function $E_{\alpha,\alpha}(-x)$, $x \geq 0$, is completely monotonic [15], yields

$$||z(.,t)|| \le t^{\alpha-1} E_{\alpha,\alpha}(-\varepsilon t^{\alpha}) ||z_0||.$$

Using the fact that $|E_{\alpha,\alpha}(-\varepsilon t^{\alpha})| \leq 1$ for $\alpha \in (0,1)$ [8], it follows that

$$\|z(.,t)\| \longrightarrow 0 \ as \ t \longrightarrow +\infty,$$

which means that the system (16) is strongly stabilizable by the feedback control u(x,t) = Kz(x,t).

Example 2. Let us consider the fractional diffusion system

$$\begin{cases} {}_{0}^{RL} D_{t}^{0.8} z(x,t) = A z(x,t) + B u(t), \ x \in \Omega, \ t \in]0, +\infty[\\ z(0,t) = z(\pi,t) = 0, \ t \in]0, +\infty[\\ \lim_{t \longrightarrow 0^{+}} 0 I_{t}^{0.2} z(x,t) = x sin(x), \ x \in \Omega, \end{cases}$$
(19)

where $\Omega =]0, \pi[$, the operator $Az = z + \frac{1}{4\pi^2} \frac{\partial^2 z}{\partial x^2}$ and the control operator B = I. The eigenvalues of A are defined by

$$\lambda_n = 1 - \frac{n^2}{4}, \ n \ge 1 \tag{20}$$

and the corresponding eigenfunctions are given by

$$\chi_n(x) = \sqrt{2}\sin(n\pi x), \ n \ge 1.$$

System (19) is unstable since $\lambda_1, \lambda_2 \ge 0$.

Applying the control (18), with K = -I, to system (19). One has, the operator $A + BK = \frac{1}{4\pi^2} \frac{\partial^2}{\partial x^2}$, with the eigenvalues being

$$\lambda_n^K = -\frac{n^2}{4}, \ n \ge 1 \tag{21}$$

and the corresponding eigenfunctions being $\chi_n^K(x) = \chi_n(x), n \ge 1$.

Moreover, one has

$$z(x,t) = t^{0.2} \sum_{n=1}^{+\infty} E_{0.8,0.8}(\lambda_n^K t^{0.8}) \langle z_0, \chi_n^K \rangle \chi_n^K(x).$$

From (21), we have

$$|arg(\lambda_n^K)| = \pi > \frac{\alpha\pi}{2} = \frac{2\pi}{5},$$

yields $\sigma^1(A + BK) = \emptyset$ and $\sigma^2(A + BK) = \left\{-\frac{n^2}{4}, n \ge 1\right\}.$

Also, from (21), one has that

$$\lambda_n^K \le -\frac{1}{4}, \ n \ge 1.$$

Consequently, from Theorem 3, we conclude that the system (19) is strongly stabilizable in Ω by the beedback control (18). Numerical illustration is given in Fig. 2.



Fig. 2. The state z(x,t) behavior of system (19) at t = 0, t = 3, t = 8 and t = 19.

3.2 Decomposition Approach

In the following, we propose an approach characterizing a feedback control that guarantees the stabilization of system (16). We suppose that A is a self adjoint operator with compact resolvent on $H = L^2(\Omega)$. So, the eigenvalues $(\lambda_n)_{n\geq 1}$ of A are real (which can be numbered in decreasing order in such away that $\lambda_n \xrightarrow[n \to +\infty]{} \infty$) and there are at most finitely many nonnegative eigenvalues $(\lambda_n)_{1\leq n\leq l}$, each with finite-dimensional eigenspace, such that $\lambda_n \geq -\delta$, for some $\delta > 0$. Yields $\sigma(A)$ can be decomposed as

$$\sigma(A) = \sigma_s(A) \cup \sigma_u(A), \tag{22}$$

with $\sigma_s(A)$ and $\sigma_u(A)$ defined as

$$\sigma_s(A) = \{\lambda_n \le -\delta, \quad n = l+1, l+2...\},$$

$$\sigma_u(A) = \{\lambda_n \ge \delta, \quad n = 1, 2, ..., l\}.$$

Since the eigenvectors $(\chi_n)_{n\geq 1}$ associated to the eigenvalues $(\lambda_n)_{n\geq 1}$ forms a complete and orthonormal basis in H [16], then one has the following decomposition of the state space

$$H = H_s \oplus H_u,$$

with

$$H_s = (I - P)H = Vect\{\chi_{l+1}, \chi_{l+2}, ...\}$$

and

$$H_u = PH = Vect\{\chi_1, \chi_2, ..., \chi_l\},\$$

where the operator $P \in L(H)$ represents the projection one [9].

Furthermore, the decomposition of system (16) may be described by

$$\begin{cases} R^{L}D_{t}^{\alpha}z_{s}(x,t) = A_{s}z_{s}(x,t) + (I-P)Bu(x,t), & x \in \Omega, \ t \in]0, +\infty[\\ z_{s}(\eta,t) = 0, & \eta \in \partial\Omega, \ t \in]0, +\infty[\\ \lim_{t \longrightarrow 0^{+}} z_{s}(x,t) = z_{0s}(x) = (I-P)z_{0}(x), & x \in \Omega\\ z_{s} = (I-P)z, & z \in L^{2}(\Omega) \end{cases}$$
(23)

and

$$\begin{cases} R^{L}D_{t}^{\alpha}z_{u}(x,t) = A_{u}z_{u}(x,t) + PBu(x,t), & x \in \Omega, \ t \in]0, +\infty[\\ z_{u}(\eta,t) = 0, & \eta \in \partial\Omega, \ t \in]0, +\infty[\\ \lim_{t \longrightarrow 0^{+}} z_{u}(x,t) = z_{0u}(x) = Pz_{0}(x), & x \in \Omega\\ z_{u} = Pz, & z \in L^{2}(\Omega), \end{cases}$$

$$(24)$$

where A_s and A_u define the restrictions of A on H_s and H_u respectively, with

$$\begin{cases} \sigma(A_s) = \sigma_s(A) \\ \sigma(A_u) = \sigma_u(A) \end{cases}$$

and the operator A_u is bounded on H_u .

For $\alpha = 1$ case, in [16], it has been shown that if A_s satisfies the following spectrum growth condition

$$\lim_{t \to +\infty} \frac{\|S_s(t)\|}{t} = \sup(\operatorname{Re}(\sigma(A_s))),$$

then the stabilization of system (16) boils down to the stabilization of (24).

The following theorem gives an extension of this result to $\alpha \in (0, 1)$ case.

Theorem 4. Let the spectrum $\sigma(A)$ of A satisfies the above spectrum decomposition assumption (22) and $\sigma(A_s) \subset \sigma^2(A)$. If the system (24) is strongly stabilizable by the control

$$u(x,t) = K_u z_u(x,t), \tag{25}$$

where $K_u \in L(H, U)$ with

$$||z_u(.,t)|| \le M \ t^{-\mu},\tag{26}$$

for some $\mu > 0$ and M > 0, then the system (16) is strongly stabilizable using the feedback control (25).

Proof. One has that the system (23) admits a unique mild solution [7] given by

$$z_{s}(.,t) = t^{\alpha-1} \sum_{n=l+1}^{+\infty} E_{\alpha,\alpha}(\lambda_{n}t^{\alpha}) \langle z_{0s}, \chi_{n} \rangle \chi_{n}(.)$$

+
$$\sum_{n=l+1}^{+\infty} \int_{0}^{t} (t-\tau)^{\alpha-1} E_{\alpha,\alpha}(\lambda_{n}(t-\tau)^{\alpha}) \langle (I-P)Bu(.,\tau), \chi_{n} \rangle \chi_{n}(.) \, \mathrm{d}\tau.$$

$$(27)$$

From the spectrum decomposition relation (22), we have that $\lambda_n \leq -\delta$, for all $n \geq l+1$, then using the completely monotonic property [15] of the Mittag-Leffler function $E_{\alpha,\alpha}(-x)$, $x \geq 0$, yields

$$E_{\alpha,\alpha}(\lambda_n t^{\alpha}) \le E_{\alpha,\alpha}(-\delta t^{\alpha}), \ \forall \ n \ge l+1,$$
(28)

and

$$E_{\alpha,\alpha}(\lambda_n(t-\tau)^{\alpha}) \le E_{\alpha,\alpha}(-\delta(t-\tau)^{\alpha}), \ \forall \ n \ge l+1.$$
(29)

Replacing (28) and (29) in (27) and applying the control $u(x,t) = K_u z_u(x,t)$ into (23), one obtains

$$\begin{aligned} \|z_s(.,t)\| &\leq t^{\alpha-1} E_{\alpha,\alpha}(-\delta t^{\alpha}) \|z_{0s}\| \\ &+ C_p \int_0^t (t-\tau)^{\alpha-1} \tau^{-\mu} E_{\alpha,\alpha}(-\delta (t-\tau)^{\alpha}) \,\mathrm{d}\tau, \end{aligned}$$

with $C_p = M\Gamma(1-\mu) ||K_u|| ||I - P|| ||B||$. It implies

$$\begin{aligned} \|z_s(.,t)\| &\leq \frac{E_{\alpha,\alpha}(-\delta t^{\alpha})}{t^{1-\alpha}} \|z_{0s}\| + C_p \sum_{k=1}^{+\infty} \int_0^t \frac{(-\delta)^k (t-s)^{\alpha k+\alpha-1} s^{-\mu} \, \mathrm{d}s}{\Gamma(\alpha k+\alpha)} \\ &\leq \frac{E_{\alpha,\alpha}(-\delta t^{\alpha})}{t^{1-\alpha}} \|z_{0s}\| + C_p \sum_{k=1}^{+\infty} \frac{(-\delta)^k t^{\alpha k+\alpha-\mu}}{\Gamma(\alpha k+\alpha-\mu-1)} \Gamma(1-\mu) \\ &\leq \frac{E_{\alpha,\alpha}(-\delta t^{\alpha})}{t^{1-\alpha}} \|z_{0s}\| + C_p t^{\alpha-\mu} E_{\alpha,\alpha-\mu+1}(-\delta t^{\alpha}). \end{aligned}$$

Then, since $\sigma(A_s) \subset \sigma^2(A)$, it follows

$$||z_s(.,t)|| \le \frac{\omega_1}{t^{1-\alpha}(1+\delta t^{\alpha})} ||z_{s0}|| + C_p \frac{\omega_2 t^{\alpha-\mu}}{1+\delta t^{\alpha}}, \omega_1, \ \omega_2 > 0,$$

which leads to

$$\lim_{t \to +\infty} \|z_s(.,t)\| = 0.$$
(30)

On the other hand, taking into acount (26) and that control (25) strongly stabilizes system (24), one gets

$$\lim_{t \to +\infty} \|z_u(.,t)\| = 0.$$
(31)

Hence, from the relation

$$||z(.,t)|| \le ||z_s(.,t)|| + ||z_u(.,t)||,$$
(32)

it follows, by using (31) and (30), that $\lim_{t \to +\infty} ||z(.,t)|| = 0$, which achieves the proof.

Example 3. Let's consider $\Omega =]0, 2[$ and the fractional diffusion system

$$\begin{cases} {}^{RL}_{0} D^{0.4}_{t} z(x,t) = \frac{2}{9\pi^2} \frac{\partial^2 z}{\partial x^2} z(x,t) + \delta z(x,t) + Bu(t), \ x \in \Omega, \ t \in]0, +\infty[\\ z(0,t) = z(2,t) = 0, \\ \lim_{t \to 0^+} 0 I^{0.6}_{t} z(x,t) = z_0(x), \\ x \in \Omega, \end{cases}$$
(33)

where the operator $Az = \frac{2}{9\pi^2} \frac{\partial^2 z}{\partial x^2} + \delta z$, with $\delta = 2$ and

$$D(A) = \{ z \in L^2(0,2), z(0,t) = z(2,t) = 0, (\forall t > 0) \}$$

is self-adjoint and the control operator $B = \delta I$.

The eigenvalues and the eigenfunctions of A are given by

$$\begin{cases} \lambda_n = 2 - \frac{2n^2}{9}, \quad n \ge 1\\ \chi_n(x) = \sqrt{2}\sin(n\pi x), \quad n \ge 1. \end{cases}$$

One has $\sigma(A) = \{\frac{16}{9}, \frac{4}{3}, 0\} \cup \{\lambda_n, n = 4, 5, ...\}$ which satisfies the spectrum decomposition assumption (22) with

$$\sigma_s(A) = \{2 - \frac{2n^2}{9}, \quad n = 4, 5, \dots\},\$$
$$\sigma_u(A) = \{\frac{16}{9}, \frac{4}{3}, 0\},\$$

and $\sigma_s(A) \subset \sigma^2(A)$ because of $|arg(2 - \frac{2n^2}{9})| = \pi > \frac{0.4\pi}{2}$, for all $n \ge 4$.

The eigenvectors $(\chi_n)_{n\geq 1}$ associated to the eigenvalues $(\lambda_n)_{n\geq 1}$ forms a complete basis on $L^2(\Omega)$, thus the system (33) may be decomposed into sub-systems (23) and (24) with

$$A_s z_s = \sum_{n=4}^{+\infty} \lambda_n \langle z_s, \chi_n \rangle \chi_n, \forall \ z_s \in H_s = Vect\{\chi_n, n \ge 4\}$$
$$A_u z_u = \sum_{n=1}^{3} \lambda_n \langle z_u, \chi_n \rangle \chi_n, \forall \ z_u \in H_u = Vect\{\chi_1, \chi_2, \chi_3\}.$$
(34)

On the other hand, system (33) is unstable since $\lambda_1, \lambda_2, \lambda_3 \ge 0$.

So, applying the control (18), with K = -I, to the unstable part of system (33) with (34). It follows that the operator $A_u + B_u K = \frac{2}{9\pi^2} \frac{\partial^2}{\partial x^2}$, with the eigenvalues being

$$\lambda_n^K = -\frac{2n^2}{9}, \ n \ge 1$$

and the corresponding eigenfunctions being $\chi_n^K(x) = \chi_n(x), \ n \ge 1.$

Yields (33) is strongly stabilizable. Indeed, for all $\lambda_n^K \in \sigma_u(A)$, one has that

$$|arg(\lambda_n^K)| = |arg(-\frac{2n^2}{9})| = \pi > \frac{0.4\pi}{2}, \ 1 \le n \le 3$$

and

$$\lambda_n^K < -\frac{2}{9}, \ 1 \le n \le 3.$$

Moreover, one has

$$z_u(x,t) = t^{0.6} \sum_{n=4}^{+\infty} E_{0.4,0,4}(\lambda_n^K t^{0.4}) \langle z_0, \chi_n^K \rangle \chi_n^K(x).$$

Using the fact that the Mittag-Leffler function $E_{\alpha,\alpha}(-x)$, $x \ge 0$, is completely monotonic [15], it follows that

$$||z_u(.,t)|| \le ||t^{0.6} \sum_{n=1}^{+\infty} E_{0.4,0,4}(-\frac{2}{9}t^{0.4})\langle z_{u0}, \chi_n^K \rangle \chi_n^K ||.$$

Also, by considering the fact that $|E_{\alpha,\alpha}(-\varepsilon t^{\alpha})| \leq 1$ for $\alpha \in (0,1)$ [8] yields

$$||z_u(.,t)|| \le t^{0.6} ||z_{u0}||,$$

which means that (26) holds with $M = ||z_{u0}||$ and $\mu = 0.6$.

Hence, all the contitions of Theorem 4 are satisfied. Thus system (33) is strongly stabilizable by u(x,t) = -2z(x,t).

4 Conclusion

The present paper deals with the concepts of the stability and stabilization of the state for Riemann–Liouville time fractional differential system of order $\alpha \in (0, 1)$. We investigated several interesting strong stability criterion's. Also, we explored the exponential stability. Furthermore, the decomposition method is utilized to derive the stabilization of fractional differential systems. Hence, we presented different examples with some simultations to illustrate the applicability of the estabilished theorems. We claim that our developed results can be useful to analyse and control the behaviour of several real world phenomena such as some heat transfer processes and anomalous sub-diffusion ones.

The problem of the state gradient stability of fractional time differential systems of order $\alpha \in (0,1)$ could be considered as our future work. Various questions are still open, for example, extending the presented results here to a class of complex linear fractional systems and studying the stabilization of fractional semilinear systems as well as nonlinear ones in general, which are closer to real applications.

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Deformed Joint Free Distributions of Semicircular Elements Induced by Multi Orthogonal Projections

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Abstract. In this paper, we consider (i) how to establish semicircular elements $\{U_k\}_{k=1}^N$ induced by *N*-many mutually orthogonal projections $\{q_k\}_{k=1}^N$, for $N \in (\mathbb{N} \setminus \{1\}) \cup \{\infty\}$, and the corresponding free product Banach *-probability space $\mathbb{L}_Q^{(N)}$ generated by $\{U_k\}_{k=1}^N$, (ii) the free-distributional data on $\mathbb{L}_Q^{(N)}$, (iii) certain *-homomorphisms on $\mathbb{L}_Q^{(N)}$, and (iv) how the *-homomorphisms of (iii) deform the original free-distributional data of (ii).

Keywords: Free probability \cdot Projections \cdot (Weighted-)semicircular elements \cdot Banach *-probability spaces \cdot Integer-shifts \cdot Restricted-integer-shifts

1 Introduction

In this paper, we study certain *-homomorphisms acting on a free product Banach *-algebra $\mathbb{L}_Q^{(N)}$ generated by mutually free, N-many semicircular elements $\mathcal{S}^{(N)} = \{U_k\}_{k=1}^N$, induced by mutually orthogonal N-many projections $\mathbf{Q}_o = \{q_k\}_{k=1}^N$, for

$$N \in \mathbb{N}_{>1}^{\infty} \stackrel{def}{=} (\mathbb{N} \setminus \{1\}) \cup \{\infty\},\$$

where $\infty = |\mathbb{N}|$. Especially, we consider the cases where such *-homomorphisms are constructed by some shifting processes on the index set

 $\{1, ..., N\}$

of $\mathcal{S}^{(N)}$, or of \mathbf{Q}_o . The main results show how our *-homomorphisms deform the free probability on $\mathbb{L}_O^{(N)}$.

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1.1 Motivations

In earlier works (e.g., [1,4,6,12,15,19-21]), semicircular elements are constructed-and-studied in topological *-probability spaces (e.g., C^* -probability spaces, or W^* -probability spaces, or Banach *-probability spaces, etc.). Different from them, the construction of semicircular elements here is motivated by that of weighted-semicircular elements in the Banach *-probability spaces of [5] and [8] from an analysis on the p-adic number fields \mathbb{Q}_p , for primes p.

By mimicking the weighted-semicircularity of [5] and [8], a construction of (weighted-)semicircular elements from arbitrary mutually orthogonal $|\mathbb{Z}|$ -many projections in fixed C^* -probability spaces is introduced, and the corresponding (weighted-)semicircular law(s) is (are) considered in [6] (See short Sects. 3 through Sect. 5 below). Independently, free distributions of free reduced words in mutually free, multi semicircular elements were characterized, estimated, and asymptotically estimated by their joint free moments in [7] (See Sect. 6.2 below).

Based on the main results of [6] and [7], the free product Banach *-algebra $\mathbb{L}_Q^{(N)}$ generated by the free semicircular family $\mathcal{S}^{(N)} = \{U_k\}_{k=1}^N$, which is induced by the family $\mathbf{Q}_o = \{q_k\}_{k=1}^N$ of mutually orthogonal projections $q_1, ..., q_N$ of an arbitrary C^* -probability space (A_o, ψ_o) , is considered as a Banach *-subalgebra of the Banach *-probability space \mathbb{L}_Q , generated by mutually free, $|\mathbb{Z}|$ -many semicircular elements. Then the free-distributional data on $\mathbb{L}_Q^{(N)}$ would be characterized naturally (e.g., [7]). And then, we define-and-study a certain type of *-homomorphisms on $\mathbb{L}_Q^{(N)}$. In particular, we are interested in how these morphisms on $\mathbb{L}_Q^{(N)}$ affect the original free-probabilistic information on $\mathbb{L}_Q^{(N)}$.

1.2 Overview

In short Sects. 2, 3, 4 and 5, we introduce backgrounds of our works briefly. In Sect. 6, we construct an operator algebra \mathbb{L}_Q generated by our semicircular elements under free product, and free-distributional data on \mathbb{L}_Q are considered.

In Sect. 7, certain shifting processes on \mathbb{Z} are defined, and the corresponding *-*isomorphisms* are established on \mathbb{L}_Q . It is shown that such *-isomorphisms form a subgroup \mathfrak{B} of the *automorphism group* $Aut(\mathbb{L}_Q)$ of \mathbb{L}_Q ; and, it is isomorphic to the infinite cyclic abelian group $(\mathbb{Z}, +)$ as groups. Interestingly, our (weighted-)semicircularity on \mathbb{L}_Q is preserved by the action of \mathfrak{B} , implying that the action of \mathfrak{B} preserves the free probability on \mathbb{L}_Q .

In Sect. 8, arbitrarily given N-many mutually orthogonal projections of a C^* algebra are fixed for $N \in \mathbb{N}_{\geq 1}^{\infty}$, and we study how they induce the corresponding free semicircular family $\mathcal{S}^{(N)}$, and show this family generates the Banach *probability space $\mathbb{L}_Q^{(N)}$. Especially, $\mathbb{L}_Q^{(N)}$ can be understood as a free-probabilistic sub-structure of \mathbb{L}_Q of Sect. 7. By restricting the action of \mathfrak{B} on \mathbb{L}_Q to that on $\mathbb{L}_Q^{(N)}$, it is proven that this restricted action of \mathfrak{B} distorts the original free probability "on $\mathbb{L}_Q^{(N)}$." Such distortions are characterized.

2 Preliminaries

For fundamental free probability theory, e.g., see [17, 19], and the citations therein. *Free probability* is the noncommutative operator-algebraic analogue of classical *measure theory* (including *probability theory*) and *statistical analysis*. It is not only an important branch of functional analysis (e.g., [2-4, 7, 12, 14, 15]), but also an interesting application in related fields (e.g., [5, 6, 8, 13, 16, 20, 21]).

We here use combinatorial approach [17] of Speicher. *Joint free moments* and *joint free cumulants* of operators will be computed, and the (free-probabilistic) free product (of [17] and [19]) will be used without detailed definitions and backgrounds.

3 The Banach *-Algebra \mathfrak{L}_Q

Let (\mathcal{B}, φ) be a topological *-probability space (a C^* -probability space, or a W^* -probability space, or a Banach *-probability space, etc.), where \mathcal{B} is a topological *-algebra (a C^* -algebra, resp., a W^* -algebra, resp., a Banach *-algebra, etc.), and φ is a bounded linear functional on \mathcal{B} .

An operator $a \in \mathcal{B}$ is said to be a *free random variable*, if we understand it as an element of (\mathcal{B}, φ) . A free random variable $a \in (\mathcal{B}, \varphi)$ is said to be *self-adjoint*, if the operator a is self-adjoint in \mathcal{B} in the sense that $a^* = a$ in \mathcal{B} , where a^* is the *adjoint of a* (e.g., [11]).

Definition 3.1. A self-adjoint free random variable a is weighted-semicircular in (\mathcal{B}, φ) with the weight $t_0 \in \mathbb{C}^{\times} = \mathbb{C} \setminus \{0\}$ (or, in short, t_0 -semicircular), if

$$k_n^{\mathcal{B}}(a,...,a) = \begin{cases} k_2^{\mathcal{B}}(a, a) = t_0 & if \ n = 2\\ 0 & otherwise, \end{cases}$$
(3.1)

for all $n \in \mathbb{N}$, where $k_{\bullet}^{\mathcal{B}}(...)$ is the free cumulant on \mathcal{B} in terms of φ under the Möbius inversion of [17].

If $t_0 = 1$ in (3.1), the 1-semicircular element *a* is said to be semicircular in (\mathcal{B}, φ) . i.e., *a* is semicircular in (\mathcal{B}, φ) , if

$$k_n(a,...,a) = \begin{cases} 1 & \text{if } n = 2\\ 0 & \text{otherwise,} \end{cases}$$
(3.2)

for all $n \in \mathbb{N}$.

By the *Möbius inversion* of [17], the weighted-semicircularity (3.1) is recharacterized as follows: a self-adjoint operator a is t_0 -semicircular in (\mathcal{B}, φ) , if and only if

$$\varphi(a^n) = \omega_n \left(t_0^{\frac{n}{2}} c_{\frac{n}{2}} \right), \qquad (3.3)$$

where

$$\omega_n \stackrel{def}{=} \begin{cases} 1 & \text{if } n \text{ is even} \\ 0 & \text{if } n \text{ is odd,} \end{cases}$$

for all $n \in \mathbb{N}$, and c_k are the k-th Catalan numbers,

$$c_k = \frac{1}{k+1} \binom{2k}{k} = \frac{1}{k+1} \frac{(2k)!}{k!(2k-k)!} = \frac{(2k)!}{k!(k+1)!},$$

for all $k \in \mathbb{N}_0 = \mathbb{N} \cup \{0\}$.

So, by (3.3), a free random variable *a* is semicircular in (\mathcal{B}, φ) , if and only if *a* is 1-semicircular in (\mathcal{B}, φ) , if and only if

$$\varphi(a^n) = \omega_n \, c_{\frac{n}{2}},\tag{3.4}$$

for all $n \in \mathbb{N}$.

From below, we use the t_0 -semicircularity (3.1) (or the semicircularity (3.2)) and its characterization (3.3) (resp., (3.4)) alternatively.

If a is a self-adjoint free random variable of (\mathcal{B}, φ) , then

the free moments $\{\varphi(a^n)\}_{n=1}^{\infty}$,

and

the free cumulants
$$\{k_n^{\mathcal{B}}(a,...,a)\}_{n=1}^{\infty}$$

provide equivalent free-distributional data of a in (\mathcal{B}, φ) (e.g., [17]). Indeed, the *Möbius inversion* makes us have

$$\varphi(a^n) = \sum_{\pi \in NC(n)} \left(\prod_{V \in \pi} k^{\mathcal{B}}_{|V|}(a, \ ..., \ a) \right),$$

and

$$k_n^{\mathcal{B}}(a,...,a) = \sum_{\pi \in NC(n)} \left(\prod_{V \in \theta} \varphi(a^{|V|}) \right) \mu(\pi,1_n),$$

where NC(n) is the *lattice* of all *noncrossing partitions* over $\{1, ..., n\}$, and " $V \in \pi$ " means "V is a *block of* π ," and where $\mu(\pi, 1_n)$ is the Möbius functional values at the interval $[\pi, 1_n]$ in NC(n), where

$$1_n = \{(1, ..., n)\}$$

is the maximal element of the lattice NC(n) having a single block (1, ..., n).

We now fix a C^{*}-probability space (A, ψ) , where A is a C^{*}-algebra, and assume that A contains mutually orthogonal, $|\mathbb{Z}|$ -many projections $\{q_j\}_{j\in\mathbb{Z}}$, i.e.,

$$q_j^* = q_j = q_j^2 \text{ in } A, \text{ for all } j \in \mathbb{Z},$$
(3.5)

and

$$q_i q_j = \delta_{i,j} q_j$$
 in A, for all $i, j \in \mathbb{Z}$,

where δ is the *Kronecker delta*. Remark that there do exist such C^* -probabilistic structures naturally (e.g., [5,8,11]), or artificially (e.g., [6]).

Fix the family,

$$\mathbf{Q} = \{q_j : j \in \mathbb{Z}\} \text{ in } A, \tag{3.6}$$

of mutually orthogonal projections q_j 's of (3.5).

And let Q be the C^{*}-subalgebra of A generated by the family \mathbf{Q} of (3.6),

$$Q \stackrel{def}{=} C^* \left(\mathbf{Q} \right) \subseteq A, \tag{3.7}$$

where $C^*(Y)$ are the C^* -subalgebras generated by the subsets $Y \cup Y^*$ of A, where

$$Y^* = \{y^* : y \in Y\}$$
 in A.

Proposition 3.1. If Q is the C^* -subalgebra (3.7) of A, then

$$Q \stackrel{*-iso}{=} \bigoplus_{j \in \mathbb{Z}} (\mathbb{C} \cdot q_j) \stackrel{*-iso}{=} \mathbb{C}^{\oplus |\mathbb{Z}|}, \qquad (3.8)$$

in A, where \oplus is the direct product of C^{*}-algebras.

Proof. The characterization (3.8) is proven by the mutual-orthogonality (3.5) of **Q**.

Define the linear functionals ψ_i on the C^{*}-algebra Q by

$$\psi_j(q_i) = \delta_{ij}\psi(q_j), \text{ for all } i \in \mathbb{Z},$$
(3.9)

for all $j \in \mathbb{Z}$, where ψ is the linear functional of (A, ψ) . These linear functionals $\{\psi_j\}_{j\in\mathbb{Z}}$ of (3.9) are well-defined on Q by (3.8).

Assumption. Let (A, ψ) be a fixed C^* -probability space, and let Q be the C^* -subalgebra (3.7) of A. From below, we assume

$$\psi(q_j) \neq 0$$
 in $\mathbb{C}^{\times} = \mathbb{C} \setminus \{1\}$, for all $j \in \mathbb{Z}$,

where q_j are projections in the generating family **Q** of (3.6).

Then, as an independent C^* -algebra, the C^* -subalgebra Q of A forms C^* -probability spaces (Q, ψ_j) , where ψ_j are the linear functionals (3.9) on Q, for all $j \in \mathbb{Z}$. We call them, the *j*-th C^* -probability spaces of Q in (A, ψ) , for all $j \in \mathbb{Z}$.

Now, define bounded linear transformations **c** and **a** acting on the C^* -algebra Q, by linear morphisms satisfying

$$\mathbf{c}(q_j) = q_{j+1}, \text{ and } \mathbf{a}(q_j) = q_{j-1},$$
 (3.10)

for all $j \in \mathbb{Z}$. Then **c** and **a** are well-defined operators "acting on Q" by (3.8). These are understood to be *Banach-space operators* in the *operator space* B(Q), consisting of all bounded linear transformations on Q, by understanding Q as a Banach space under its C^* -norm topology (e.g., [9]). **Definition 3.2.** The Banach-space operators \mathbf{c} and \mathbf{a} of (3.10) are said to be the creation, respectively, the annihilation on Q. Define

$$\mathbf{l} = \mathbf{c} + \mathbf{a} \ on \ Q. \tag{3.11}$$

We call this Banach-space operator \mathbf{l} of (3.11), the radial operator on Q.

Now, define a subspace \mathfrak{L} of B(Q) by

$$\mathfrak{L} \stackrel{def}{=} \overline{\mathbb{C}[\{l\}]}^{\|\cdot\|}, \qquad (3.12)$$

equipped with the operator norm,

$$||T|| = \sup\{||Tq||_Q : ||q||_Q = 1\},$$
(3.13)

where $\|.\|_Q$ is the C^* -norm on Q, where $\overline{Z}^{\|.\|}$ are the operator-norm closures of subsets $Z \subseteq B(Q)$ (e.g., [9]). By (3.12), this subspace \mathfrak{L} forms a Banach algebra in the vector space B(Q).

On this Banach algebra \mathfrak{L} of (3.12), define an operation (*) by

$$\left(\sum_{n=0}^{\infty} t_n \mathbf{l}^n\right)^* = \sum_{n=0}^{\infty} \overline{t_n} \mathbf{l}^n \operatorname{in} \mathfrak{L}, \qquad (3.14)$$

where \overline{z} are the conjugates of $z \in \mathbb{C}$.

Then the operation (3.13) is a well-defined *adjoint on* \mathfrak{L} (See [6]), and hence, every element of \mathfrak{L} is *adjointable* (in the sense of [9]) in B(Q). So, the Banach algebra \mathfrak{L} of (3.12) forms a *Banach* *-algebra with the adjoint (3.13) in B(Q). We call this Banach *-algebra \mathfrak{L} , the *radial* (*Banach* *-)algebra on Q.

Construct now the tensor product Banach *-algebra \mathfrak{L}_Q ,

$$\mathfrak{L}_Q = \mathfrak{L} \otimes_{\mathbb{C}} Q, \tag{3.15}$$

where $\otimes_{\mathbb{C}}$ is the tensor product of Banach *-algebras, where \mathfrak{L} is the radial algebra (3.12).

Definition 3.3. The Banach *-algebra \mathfrak{L}_Q of (3.14) is called the radial projection (Banach *-)algebra on Q.

4 Weighted-Semicircular Elements

In this section, we construct weighted-semicircular elements induced by the family \mathbf{Q} of (3.6) in the radial projection algebra \mathfrak{L}_Q of (3.14). Let (Q, ψ_j) be the *j*-th C^* -probability spaces of Q in (A, ψ) , where ψ_j are the linear functionals of (3.9), for all $j \in \mathbb{Z}$.

Note that, if

$$u_j \stackrel{def}{=} \mathbf{l} \otimes q_j \in \mathfrak{L}_Q, \text{ for all } j \in \mathbb{Z},$$

$$(4.1)$$

then

$$u_j^n = (\mathbf{l} \otimes q_j)^n = \mathbf{l}^n \otimes q_j, \text{ for all } n \in \mathbb{N},$$

for $j \in \mathbb{Z}$. i.e., such operators $\{u_j\}_{j\in\mathbb{Z}}$ generate \mathfrak{L}_Q , by (3.8), (3.12) and (3.14).

By (4.1), one can define a linear functional φ_j on \mathfrak{L}_Q by a morphism satisfying that

$$\varphi_j\left((\mathbf{l}\otimes q_i)^n\right) \stackrel{def}{=} \psi_j\left(\mathbf{l}^n(q_i)\right) \tag{4.2}$$

for all $n \in \mathbb{N}$, for all $i, j \in \mathbb{Z}$.

By the well-defindness of the linear functionals $\{\varphi_j\}_{j\in\mathbb{Z}}$ of (4.2), the Banach *-algebra \mathfrak{L}_Q forms well-defined Banach *-probability spaces,

$$(\mathfrak{L}_Q, \varphi_j), \text{ for all } j \in \mathbb{Z}.$$
 (4.3)

If \mathbf{c} and \mathbf{a} are the creation, respectively, the annihilation on Q of (3.10), then

 $\mathbf{ca} = \mathbf{1}_Q = \mathbf{ac}$, the identity operator on Q,

in B(Q). So, one has

$$\mathbf{c}^{n_1}\mathbf{a}^{n_2} = \mathbf{a}^{n_2}\mathbf{c}^{n_1}, \forall n_1, n_2 \in \mathbb{N}.$$
(4.4)

By (4.4), we have

$$\mathbf{l}^{n} = (\mathbf{c} + \mathbf{a})^{n} = \sum_{k=0}^{n} \binom{n}{k} \mathbf{c}^{k} \mathbf{a}^{n-k}, \qquad (4.5)$$

for all $n \in \mathbb{N}$, where

$$\binom{n}{k} = \frac{n!}{k!(n-k)!}, \text{ for all } k \le n \in \mathbb{N}_0.$$

By (4.5), for any $n \in \mathbb{N}$,

$$\mathbf{l}^{2n-1} = \sum_{k=0}^{2n-1} \binom{2n-1}{k} \mathbf{c}^k \mathbf{a}^{n-k},$$
(4.6)

and

$$\mathbf{l}^{2n} = \sum_{k=0}^{2n} \binom{2n}{k} \mathbf{c}^k \mathbf{a}^{n-k} = \binom{2n}{n} \mathbf{c}^n \mathbf{a}^n + [\text{Rest terms}]$$
(4.7)

(e.g., see [6] for details).

Proposition 4.1. Let l be the radial operator on Q. Then

(4.8) l^{2n-1} does not contain the nonzero 1_Q -summand.

(4.9)
$$\mathbf{l}^{2n}$$
 contains the nonzero $\mathbf{1}_Q$ -summand, $\binom{2n}{n} \cdot \mathbf{1}_Q$.

Proof. The statements (4.8) and (4.9) are shown by (4.6) and (4.7), respectively.

Since

one has

$$u_j^n = (\mathbf{l} \otimes q_j)^n = \mathbf{l}^n \otimes q_j,$$

$$\varphi_j \left(u_j^{2n-1} \right) = \psi_j \left(\mathbf{l}^{2n-1} \left(q_j \right) \right) = 0,$$
(4.10)

for all $n \in \mathbb{N}$, by (4.8).

Also, we have

$$\varphi_j \left(u_j^{2n} \right) = \psi_j \left(\mathbf{l}^{2n} \left(q_j \right) \right) = \psi_j \left(\begin{pmatrix} 2n \\ n \end{pmatrix} q_j + [\text{Rest terms}] \right)$$
$$= \begin{pmatrix} 2n \\ n \end{pmatrix} \psi_j \left(q_j \right) = \begin{pmatrix} 2n \\ n \end{pmatrix} \psi \left(q_j \right),$$

by (4.7) and (4.9). i.e.,

$$\varphi_j\left(u_j^{2n}\right) = {\binom{2n}{n}}\psi\left(q_j\right), \text{ for all } n \in \mathbb{N}.$$
 (4.11)

Thus, by (4.10) and (4.11), the following free-distributional data are obtained.

Proposition 4.2. Fix $j \in \mathbb{Z}$, and let $u_k = \mathbf{l} \otimes q_k$ be the k-th generating operators of $(\mathfrak{L}_Q, \varphi_j)$, for all $k \in \mathbb{Z}$. Then

$$\varphi_j\left(u_k^n\right) = \delta_{j,k}\omega_n\left(\left(\frac{n}{2}+1\right)\psi\left(q_j\right)\right)c_{\frac{n}{2}},\tag{4.12}$$

where ω_n and $c_{\frac{n}{2}}$ are in the sense of (3.3) for all $k \in \mathbb{Z}$, and $n \in \mathbb{N}$.

Proof. By (4.10) and (4.11), one can get that: if u_j is the *j*-th generating operator of \mathfrak{L}_Q , then

$$\varphi_j\left(u_j^{2n-1}\right) = 0,$$

and

$$\varphi_j \left(u_j^{2n} \right) = \binom{2n}{n} \psi \left(q_j \right) = \binom{n+1}{n+1} \binom{2n}{n} \psi \left(q_j \right)$$
$$= \left((n+1)\psi \left(q_j \right) \right) \left(\frac{1}{n+1} \binom{2n}{n} \right)$$
$$= \left((n+1)\psi \left(q_j \right) \right) c_n,$$

for all $n \in \mathbb{N}$.

If $k \neq j$ in \mathbb{Z} , and u_k is the k-th generating operator of \mathfrak{L}_Q , then

$$\varphi_j(u_k^n) = 0$$
, for all $n \in \mathbb{N}$,

by (3.9) and (4.2).

Based on (4.12), define the linear morphisms,

$$E_{j,Q}: \mathfrak{L}_Q \to \mathfrak{L}_Q,$$

by linear transformations satisfying

$$E_{j,Q}\left(u_{i}^{n}\right) \stackrel{def}{=} \begin{cases} \frac{\psi(q_{j})^{n-1}}{\left(\left[\frac{n}{2}\right]+1\right)} u_{j}^{n} & \text{if } i=j\\ 0_{\mathfrak{L}_{Q}}, \text{ the zero operator of } \mathfrak{L}_{Q} & \text{otherwise,} \end{cases}$$
(4.13)

for all $n \in \mathbb{N}$, $i, j \in \mathbb{Z}$, where $[\frac{n}{2}]$ means the *minimal integer* greater than or equal to $\frac{n}{2}$.

The linear transformations $E_{j,Q}$ of (4.13) are well-defined on \mathfrak{L}_Q by the cyclicity (3.12) of a tensor factor \mathfrak{L} of \mathfrak{L}_Q , and by the structure theorem (3.8) of the other tensor factor Q of \mathfrak{L}_Q , by (3.14).

Now, define the linear functionals τ_j on \mathfrak{L}_Q by

$$\tau_j \stackrel{def}{=} \varphi_j \circ E_{j,Q} \text{ on } \mathfrak{L}_Q, \text{ for all } j \in \mathbb{Z},$$
(4.14)

where $E_{j,Q}$ are in the sense of (4.13).

Definition 4.1. The Banach *-probability spaces,

$$\mathfrak{L}_Q(j) \stackrel{denote}{=} (\mathfrak{L}_Q, \ \tau_j), \qquad (4.15)$$

are called the *j*-th (free) filter of \mathfrak{L}_Q , for all $j \in \mathbb{Z}$.

Observe on the *j*-th filter $\mathfrak{L}_Q(j)$ of (4.15) that:

$$\tau_j \left(u_j^n \right) = \varphi_j \left(E_{j,Q} \left(u_j^n \right) \right)$$
$$= \varphi_j \left(\frac{\psi(q_j)^{n-1}}{\left([\frac{n}{2}] + 1 \right)} \left(u_j^n \right) \right) = \frac{\psi(q_j)^{n-1}}{\left([\frac{n}{2}] + 1 \right)} \varphi_j \left(u_j^n \right)$$
$$= \frac{\psi(q_j)^{n-1}}{\left([\frac{n}{2}] + 1 \right)} \omega_n \left(\left(\frac{n}{2} + 1 \right) \psi \left(q_j \right) \right) c_{\frac{n}{2}},$$

by (4.12), i.e.,

$$\tau_j \left(u_j^n \right) = \omega_n \psi(q_j)^n c_{\frac{n}{2}}, \tag{4.16}$$

where ω_n and $c_{\frac{n}{2}}$ are in the sense of (3.3), for all $n \in \mathbb{N}$, for $j \in \mathbb{Z}$.

Lemma 4.3. Let $\mathfrak{L}_Q(j) = (\mathfrak{L}_Q, \tau_j)$ be the *j*-th filter of \mathfrak{L}_Q , for $j \in \mathbb{Z}$. Then

$$\tau_j\left(u_i^n\right) = \delta_{j,i}\left(\omega_n \psi(q_j)^n c_{\frac{n}{2}}\right),\tag{4.17}$$

for all $n \in \mathbb{N}$, for all $i \in \mathbb{Z}$.

Proof. If i = j in \mathbb{Z} , then the formula (4.17) holds by (4.16). Meanwhile, if $i \neq j$ in \mathbb{Z} , then $\tau_j(u_i^n) = 0$, by (4.2) and (4.13). Therefore, the formula (4.17) holds for all $i \in \mathbb{Z}$.

The following theorem is proven by (4.17).

Theorem 4.4. Let $\mathfrak{L}_Q(j)$ be the *j*-th filter of \mathfrak{L}_Q , for $j \in \mathbb{Z}$. Then the *j*-th generating operator (4.1) of \mathfrak{L}_Q is $\psi(q_j)^2$ -semicircular in $\mathfrak{L}_Q(j)$. Meanwhile, all other *k*-th generating operators u_k of \mathfrak{L}_Q have the zero free distribution on $\mathfrak{L}_Q(j)$, for all $k \neq j$ in \mathbb{Z} .

Proof. Note that the generating operators u_k are self-adjoint in \mathfrak{L}_Q , since

$$u_k^* = (\mathbf{l} \otimes q_k)^* = \mathbf{l} \otimes q_k = u_k$$

for all $k \in \mathbb{Z}$, by (3.13).

If u_j is the *j*-th generating operator of \mathfrak{L}_Q , then

$$\tau_j\left(u_j^n\right) = \omega_n\left(\psi\left(q_j\right)^2\right)^{\frac{n}{2}}c_{\frac{n}{2}},$$

for all $n \in \mathbb{N}$, by (4.17). Therefore, by (3.3), u_j is $\psi(q_j)^2$ -semicircular in $\mathfrak{L}_Q(j)$.

Now, suppose $k \neq j$ in \mathbb{Z} , and take the generating operator u_k of $\mathfrak{L}_Q(j)$. By the self-adjointness of u_k , the free distribution of u_k is characterized by the free-moment sequence,

$$(\tau (u_k^n))_{n=1}^{\infty} = (0, 0, 0, 0, ...),$$

by (4.17). So, u_k has the zero free distribution on $\mathfrak{L}_Q(j)$, whenever $k \neq j$.

By using the Möbius inversion of [17], one can obtain that: if $k^j_{\bullet}(...)$ is the free cumulant on \mathfrak{L}_Q in terms of a linear functional τ_j , then

$$k_n^j\left(\underbrace{u_k, u_k, \dots, u_k}_{n\text{-times}}\right) = \begin{cases} \delta_{j,k}\psi(q_j)^2 & \text{if } n = 2\\ 0 & \text{otherwise,} \end{cases}$$
(4.18)

for all $n \in \mathbb{N}$, for all $j, k \in \mathbb{Z}$ (e.g., see [6] for details).

5 Semicircular Elements

Let $\mathfrak{L}_Q(j)$ be the *j*-th filter for $j \in \mathbb{Z}$. Then, for a fixed $j \in \mathbb{Z}$, the *j*-th generating operator u_j is $\psi(q_j)^2$ -semicircular in $\mathfrak{L}_Q(j)$, since

$$\tau_j\left(u_j^n\right) = \omega_n \psi(q_j)^n c_{\frac{n}{2}},$$

equivalently,

$$k_n^j(u_j, ..., u_j) = \begin{cases} \psi(q_j)^2 & \text{if } n = 2\\ 0 & \text{otherwise,} \end{cases}$$
(5.1)

for all $n \in \mathbb{N}$, by (4.17) and (4.18).

Theorem 5.1. Let $U_j = \frac{1}{\psi(q_j)} u_j$ in the *j*-th filter $\mathfrak{L}_Q(j)$ for $j \in \mathbb{Z}$, where u_j is the *j*-th generating operator of \mathfrak{L}_Q . If

$$\psi(q_j) \in \mathbb{R}^{\times} = \mathbb{R} \setminus \{0\} \ in \ \mathbb{C}^{\times}, \tag{5.2}$$

then U_j is semicircular in $\mathfrak{L}_Q(j)$, for $j \in \mathbb{Z}$.

Proof. Fix $j \in \mathbb{Z}$, and assume the condition (5.2) holds. Then

$$U_j^* = \left(\frac{1}{\psi(q_j)}u_j\right)^* = U_j,$$

in \mathfrak{L}_Q , because u_j is self-adjoint. Consider now that

$$\tau_j \left(U_j^n \right) = \frac{1}{\psi(q_j)^n} \tau \left(u_j^n \right)$$

$$= \frac{1}{\psi(q_j)^n} \left(\omega_n \psi(q_j)^n c_{\frac{n}{2}} \right) = \omega_n c_{\frac{n}{2}},$$

(5.3)

for all $n \in \mathbb{N}$.

Therefore, under (5.2), the self-adjoint free random variable U_j is semicircular in $\mathfrak{L}_Q(j)$, by (3.4) and (5.3).

Assumption. For convenience, we assume from below that

$$\psi(q_j) \in \mathbb{R}^{\times} \text{ in } \mathbb{C}, \text{ for } q_j \in \mathbf{Q},$$

for all $j \in \mathbb{Z}$.

6 The Free Filterization $\mathfrak{L}_Q(\mathbb{Z})$

In this section, we construct the free product Banach *-probability space $\mathfrak{L}_Q(\mathbb{Z})$ of the free filters $\{\mathfrak{L}_Q(j)\}_{j\in\mathbb{Z}}$, and the corresponding sub-structure $\mathbb{L}_Q = (\mathbb{L}_Q, \tau)$ generated by a free semicircular family

$$\{U_j \in \mathfrak{L}_Q(j) : j \in \mathbb{Z}\},\$$

and study free-distributional information on \mathbb{L}_Q .

6.1 The Semicircular Filterization \mathbb{L}_Q

As before, let (A, ψ) be the fixed C^{*}-probability space containing a family $\mathbf{Q} = \{q_j\}_{j \in \mathbb{Z}}$ of mutually orthogonal projections, satisfying

$$\psi(q_i) \in \mathbb{R}^{\times}$$
, for all $j \in \mathbb{Z}$,

and let $\mathfrak{L}_Q(j)$ be the *j*-th filters of Q, for all $j \in \mathbb{Z}$. From the system, $\{\mathfrak{L}_Q(j): j \in \mathbb{Z}\},\$

define the free product Banach *-probability space $\mathfrak{L}_Q(\mathbb{Z})$ by

$$\mathfrak{L}_Q(\mathbb{Z}) \stackrel{denote}{=} (\mathfrak{L}_Q(\mathbb{Z}), \tau) \\
\stackrel{def}{=} \underset{j \in \mathbb{Z}}{\star} \mathfrak{L}_Q(j) = \begin{pmatrix} \star \\ j \in \mathbb{Z}} \mathfrak{L}_{Q,j}, & \star \\ j \in \mathbb{Z}} \tau_j \end{pmatrix},$$
(6.1.1)

with

$$\mathfrak{L}_Q(\mathbb{Z}) = \underset{j \in \mathbb{Z}}{\star} \mathfrak{L}_{Q,j}$$
, with $\mathfrak{L}_{Q,j} = \mathfrak{L}_Q, \forall j \in \mathbb{Z}$,

and

$$\tau = \underset{j \in \mathbb{Z}}{\star} \tau_j \text{ on } \mathfrak{L}_Q(\mathbb{Z})$$

For more about free-probabilistic free product, see [17] and [19].

Definition 6.1. The free product Banach *-probability space $\mathfrak{L}_Q(\mathbb{Z})$ of (6.1.1) is said to be the free filterization of $Q \subset (A, \psi)$.

Define now two subsets \mathcal{X} and \mathcal{S} of $\mathfrak{L}_Q(\mathbb{Z})$ by

$$\mathcal{X} = \{ u_j \in \mathfrak{L}_Q(j) : j \in \mathbb{Z} \}, \tag{6.1.2}$$

and

$$\mathcal{S} = \{ U_j = \frac{1}{\psi(q_j)} u_j \in \mathfrak{L}_Q(j) : j \in \mathbb{Z} \},\$$

where u_j are the operators (4.1), for all $j \in \mathbb{Z}$.

A subset \mathcal{Y} of a topological *-probability space (\mathcal{B}, φ) is said to be a *free* family, if all elements of \mathcal{Y} are mutually free in (\mathcal{B}, φ) . And, a free family \mathcal{Y} is called a *free* (weighted-)semicircular family in (\mathcal{B}, φ) , if all elements of \mathcal{Y} are (weighted-)semicircular in (\mathcal{B}, φ) . (e.g., [5] and [19]).

Theorem 6.1. Let \mathcal{X} and \mathcal{S} be the families of (6.1.2) in $\mathfrak{L}_Q(\mathbb{Z})$.

(6.1.3) \mathcal{X} is a free weighted-semicircular family in $\mathfrak{L}_Q(\mathbb{Z})$. (6.1.4) \mathcal{S} is a free semicircular family in $\mathfrak{L}_Q(\mathbb{Z})$.

Proof. Let \mathcal{X} be in the family of (6.1.2) in $\mathfrak{L}_Q(\mathbb{Z})$. By (6.1.1), all elements u_j of \mathcal{X} are from mutually distinct free blocks $\mathfrak{L}_Q(j)$ for all $j \in \mathbb{Z}$, and hence, they are mutually free in $\mathfrak{L}_Q(\mathbb{Z})$. Thus, the subset \mathcal{X} forms a free family in $\mathfrak{L}_Q(\mathbb{Z})$. Moreover, the powers $u_j^n \in \mathfrak{L}_Q(\mathbb{Z})$ of $u_j \in \mathcal{X}$ are contained in the same free block $\mathfrak{L}_Q(j)$ as free reduced words with their lengths-1 of $\mathfrak{L}_Q(\mathbb{Z})$, for all $n \in \mathbb{N}$, implying that

$$\tau\left(u_{j}^{n}\right)=\tau_{j}\left(u_{j}^{n}\right)=\omega_{n}\psi(q_{j})^{n}c_{\frac{n}{2}},\forall n\in\mathbb{N},$$

by (5.1). Therefore, the statement (6.1.3) holds. Similarly, the family S of (6.1.2) is a free family in $\mathfrak{L}_Q(\mathbb{Z})$, because

$$U_j = \frac{1}{\psi(q_j)} u_j = U_j^*, \text{ for all } j \in \mathbb{Z},$$

and the family \mathcal{X} is a free family in $\mathfrak{L}_Q(\mathbb{Z})$. So, the semicircularity (5.4) of U_j 's shows that the statement (6.1.4) holds.

By (6.1.3) and (6.1.4), the "*j*-th" generating operators u_j of the free blocks $\mathfrak{L}_Q(j)$, and their powers $u_j^n \in \mathfrak{L}_Q(j)$ provide nonzero free-distributional data on the free filterization $\mathfrak{L}_Q(\mathbb{Z})$. In particular, the free (reduced) words in $\mathcal{X} \cup \mathcal{S}$ (under operator-multiplication on $\mathfrak{L}_Q(\mathbb{Z})$) have non-vanishing free distributions on $\mathfrak{L}_Q(\mathbb{Z})$.

Definition 6.2. In the free filterization $\mathfrak{L}_Q(\mathbb{Z})$, define a Banach * -subalgebra \mathbb{L}_Q of $\mathfrak{L}_Q(\mathbb{Z})$ by

$$\mathbb{L}_Q \stackrel{def}{=} \overline{\mathbb{C}\left[\mathcal{X}\right]},\tag{6.1.5}$$

where \mathcal{X} is the free weighted-semicircular family (6.1.3) of $\mathfrak{L}_Q(\mathbb{Z})$, and \overline{Y} are the topological closures of the subsets Y of $\mathfrak{L}_Q(\mathbb{Z})$. Construct the Banach *-probability space,

$$\mathbb{L}_{Q} \stackrel{denote}{=} \left(\mathbb{L}_{Q}, \ \tau = \tau \mid_{\mathbb{L}_{Q}} \right), \tag{6.1.6}$$

in $\mathfrak{L}_Q(\mathbb{Z}) = (\mathfrak{L}_Q(\mathbb{Z}), \tau)$. We call \mathbb{L}_Q of (6.1.5) or (6.1.6), the semicircular (free-sub-)filterization of $\mathfrak{L}_Q(\mathbb{Z})$.

By the definitions (6.1.5) and (6.1.6), one obtains the following structure theorem.

Theorem 6.2. Let \mathbb{L}_Q be the semicircular filterization (6.1.5). Then

$$\mathbb{L}_Q = \overline{\mathbb{C}[S]} \stackrel{*-iso}{=} \underbrace{\star}_{j \in \mathbb{Z}} \overline{\mathbb{C}[\{u_j\}]} \stackrel{*-iso}{=} \mathbb{C} \left[\underbrace{\star}_{j \in \mathbb{Z}} \{u_j\} \right], \tag{6.1.7}$$

in $\mathfrak{L}_Q(\mathbb{Z})$, where $\overset{**-iso}{=}$ means "being Banach-*-isomorphic," and where (\star) in the first *-isomorphic relation of (6.1.7) means the free-probabilistic free product of [17] and [19], and (\star) in the second *-isomorphic relation of (6.1.7) is the pure-algebraic free product inducing noncommutative free words in \mathcal{X} .

Proof. Set-theoretically, one has

$$\mathcal{X} = \{\psi(q_j)U_j \in \mathfrak{L}_Q(j) : j \in \mathbb{Z}\}$$

in the free filterization $\mathfrak{L}_Q(\mathbb{Z})$, where $U_j \in \mathcal{S}$ are the semicircular elements of (6.1.4). Therefore,

$$\overline{\mathbb{C}[\mathcal{X}]} = \overline{\mathbb{C}[\mathcal{S}]} \operatorname{in} \mathfrak{L}_Q(\mathbb{Z}),$$

i.e., the equality (=) of (6.1.7) holds.

By the definition (6.1.5) of \mathbb{L}_Q , it is generated by the free family \mathcal{X} , and hence, the first *-isomorphic relation of (6.1.7) holds in $\mathfrak{L}_Q(\mathbb{Z})$ by (6.1.1) and (6.1.6).

Since

$$\mathbb{L}_Q \stackrel{*-\mathrm{iso}}{=} \underset{j \in \mathbb{Z}}{\star} \overline{\mathbb{C}[\{u_j\}]} \text{ in } \mathfrak{L}_Q(\mathbb{Z}),$$

every element T of \mathbb{L}_Q is a limit of linear combinations of free reduced words in \mathcal{X} . Note that all (pure-algebraic) free words in \mathcal{X} have their unique freereduced-word forms as their operator-product in $\mathcal{L}_Q(\mathbb{Z})$ (e.g., [17] and [19]). Therefore, the second *-isomorphic relation of (6.1.7) holds.

6.2 Free-Distributional Data Induced by Semicircular Elements

In this section, we consider general free-distributional data on \mathbb{L}_Q . In particular, we are interested in joint free moments of S in \mathbb{L}_Q . Throughout this section, let (B, φ) be an arbitrarily fixed topological *-probability space, and suppose there are N-many semicircular elements $x_1, ..., x_N$ in (B, φ) , for $N \in \mathbb{N} \setminus \{1\}$. Assume further that they are free from each other in (B, φ) .

By the self-adjointness of these semicircular elements $x_1, ..., x_N \in (B, \varphi)$, the free distribution, say

$$\rho \stackrel{denote}{=} \rho_{x_1,\dots,x_N},$$

of them are characterized by the joint free-moments

$$\bigcup_{n=1}^{\infty} \left(\bigcup_{(i_1,\dots,i_n)\in\{1,\dots,N\}^n} \left\{ \varphi\left(x_{i_1}x_{i_2}\dots x_{i_n}\right) \right\} \right)$$
(6.2.1)

(e.g., [17]). i.e., the free distribution ρ of (6.2.1), is characterized by the freemoments,

$$\bigcup_{l=1}^{N} \left\{ \varphi(x_l^n) \right\}_{n=1}^{\infty}, \qquad (6.2.2)$$

and the "mixed" free-moments,

$$\bigcup_{\substack{s=2\\s=2}}^{\infty} \left\{ \varphi \left(x_{i_1}^{n_1} x_{i_2}^{n_2} \dots x_{i_s}^{n_s} \right) \middle| \begin{array}{c} (i_1, \dots, i_s) \in \{1, \dots, N\}^s \\ \text{are mixed in } \{1, \dots, N\}, \\ \text{for all } n_1, \dots, n_s \in \mathbb{N} \end{array} \right\},$$
(6.2.3)

by (6.2.1). To characterize the free distribution ρ , we consider the freedistributional data (6.2.2) and (6.2.3), independently. **Corollary 6.3.** The free-distributional data (6.2.2) of ρ are determined by the semicircularity. *i.e.*,

$$\varphi\left(x_{l}^{n}\right) = \omega_{n} c_{\frac{n}{2}}, \text{ for all } n \in \mathbb{N},$$

$$(6.2.4)$$

for all l = 1, ..., N.

Proof. The formula (6.2.4) is obtained by the semicircularity (3.4).

Let's concentrate on the free-distributional data (6.2.3) of the free distribution ρ . For any $s \in \mathbb{N} \setminus \{1\}$, we fix an s-tuple I_s ,

$$I_s \stackrel{denote}{=} (i_1, ..., i_s) \in \{1, ..., N\}^s,$$
(6.2.5)

which is mixed in $\{1, ..., N\}$ in the sense that there exists at least one entry i_{k_0} in I_s such that $i_{k_0} \neq i_l$, for some $l \neq k_0$ in $\{1, ..., s\}$.

For example,

$$I_8 = (1, 1, 3, 2, 4, 2, 2, 1),$$

in $\{1, 2, 3, 4, 5\}^8$. From the sequence I_s of (6.2.5), define a set

$$[I_s] = \{i_1, i_2, \dots, i_s\},\tag{6.2.6}$$

without considering repetition. For instance, if I_8 is as above, then

$$[I_8] = \{i_1, i_2, \dots, i_8\},\$$

with

$$i_1 = i_2 = i_8 = 1,$$

 $i_4 = i_6 = i_7 = 2,$
 $i_3 = 3, \text{ and } i_5 = 4.$

i.e., all 1's in I_8 are regarded as distinct elements i_1, i_2 and i_8 in the set $[I_8]$.

From the set $[I_s]$ of (6.2.6), define a unique "noncrossing" partition $\pi_{(I_s)}$ of the non-crossing-partition lattice $NC([I_s])$ over $[I_s]$, such that (i) starting from the very first entry i_1 , construct the block V_1 of $\pi_{(I_s)}$, satisfying

$$V_1 = \left(i_{j_1} = i_1, \ i_{j_2}, \ \dots, \ i_{j_{|V_1|}}\right) \in \pi_{(I_s)},\tag{6.2.7}$$

$$\Leftrightarrow$$

$$\exists k \in \{1, ..., N\}, \text{ s.t., } i_{j_1} = i_{j_2} = ... = i_{j_{|V|}} = k,$$

and then do the same process to the very next entry other than i_{j_1} , ..., $i_{j_{|V_1|}}$, step-by-step, until such processes end, (ii) such a partition $\pi_{(I_s)}$ of (i) is "maximal" in $NC([I_s])$ (e.g., [17]).

For example, if I_8 and $[I_8]$ are as above, then there exists a noncrossing partition

$$\pi_{(I_8)} = \{ (i_1, i_2, i_8), (i_3), (i_4, i_6, i_7), (i_5) \} \\ = \{ (1, 1, 1), (3), (2, 2, 2), (4) \},\$$

in $NC([I_8])$, satisfying the above conditions (i) and (ii). In this case,

$$V_1 = \{i_1, i_2, i_8\} = \{1, 1, 1\}, \text{ as in } (6.2.7).$$

Denote the noncrossing partition $\pi_{(I_s)} \in NC([I_s])$ of (6.2.6) by

$$\pi_{(I_s)} = \{V_1, \dots, V_t\},\$$

where $t \leq s$ and $V_k \in \pi_{(I_s)}$ are the blocks of (ii), satisfying (i), for k = 1, ..., t. Then the partition $\pi_{(I_s)}$ is the joint partition,

$$\pi_{(I_s)} = \mathbf{1}_{|V_1|} \vee \mathbf{1}_{|V_2|} \vee \dots \vee \mathbf{1}_{|V_t|}, \tag{6.2.8}$$

where $1_{|V_k|}$ are the maximal partitions of $NC(V_k)$, for all k = 1, ..., t, by regarding V_k as discrete sets.

Let I_s be in the sense of (6.2.5), and let $x_{i_1}, ..., x_{i_s}$ be the corresponding semicircular elements of (B, φ) induced by I_s , without considering repetition in the set $\{x_1, ..., x_N\}$ of our fixed mutually free, N-many semicircular elements of (B, φ) . And then, define a free random variable $X[I_s]$ by

$$X[I_s] \stackrel{def}{=} \prod_{l=1}^s x_{i_l} \in (B,\varphi).$$
(6.2.9)

If $X[I_s]$ is a free random variable (6.2.9), then

$$\varphi\left(X[I_s]\right) = \sum_{\pi \in NC([I_s])} k_{\pi}$$

by the Möbius inversion of [17], where k_{π} are the partition-depending free cumulants of [17],

$$k_{\pi} = \prod_{V \in \pi} k_V,$$

where k_V is the block-depending free cumulants of [17], and hence, it goes to

$$= \sum_{\pi \in NC([I_s]), \ \pi \le \pi_{(I_s)}} k_{\pi}$$

by the mutual-freeness of $x_1, ..., x_N$ in (B, φ)

$$= \sum_{(\theta_1, \dots, \theta_t) \in NC(V_1) \times \dots \times NC(V_t)} k_{\theta_1 \vee \dots \vee \theta_t}$$
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by (6.2.8)

$$= \sum_{(\theta_1,\dots,\theta_t)\in NC_2(V_1)\times\dots\times NC_2(V_t)} k_{\theta_1\vee\dots\vee\theta_t}$$
$$= \sum_{(\theta_1,\dots,\theta_t)\in NC_2(V_1)\times\dots\times NC_2(V_t)} \left(\prod_{l=1}^t k_{\theta_l}\right), \qquad (6.2.10)$$

by the semicircularity (3.2) of $x_{i_1}, ..., x_{i_s}$ in (B, φ) , where $NC_2(X)$ is the subset,

$$NC_2(X) = \{ \pi \in NC(X) : \forall V \in \pi, |V| = 2 \},$$
 (6.2.11)

of the noncrossing-partition lattice NC(X) over sets X.

By (6.2.10), (6.2.11) and (3.2), if there is at least one $k_0 \in \{1, ..., t\}$, such that $|V_{k_0}|$ is odd in \mathbb{N} (or equivalently, if s is odd), then

$$\varphi\left(X[I_s]\right) = 0,$$

where $X[I_s]$ is in the sense of (6.2.9). Meanwhile, if

$$|V_k| \in 2\mathbb{N}, \text{ for all } k = 1, ..., t,$$
 (6.2.12)

where $2\mathbb{N} = \{2n : n \in \mathbb{N}\}\$, then the formula (6.2.10) is nonzero.

More precisely, if the condition (6.2.12) is satisfied, then the summands $k_{\theta_1 \vee \ldots \vee \theta_t}$ of (6.2.10) satisfy that

$$k_{\theta_1 \vee \ldots \vee \theta_t} = \prod_{V \in \theta_1 \vee \ldots \vee \theta_t} k_V = \prod_{V \in \theta_1 \vee \ldots \vee \theta_t} \left(\prod_{i=1}^t 1^{\#(\theta_i)} \right) = 1,$$
(6.2.13)

by the semicircularity (3.2), where $\#(\theta_i)$ are the number of blocks of θ_i , for all i = 1, ..., t. Therefore, if the condition (6.2.12) holds, then

$$\varphi\left(X[I_s]\right) = \sum_{\left(\theta_1, \dots, \theta_t\right) \in NC_2(V_1) \times \dots \times NC_2(V_t)} 1$$

= $\left|NC_2\left(V_1\right) \times \dots \times NC_2\left(V_t\right)\right|,$ (6.2.14)

by (6.2.10) and (6.2.13), where |Y| are the cardinalities of sets Y.

Theorem 6.4. Let I_s be a mixed s-tuple (6.2.5), and let $X[I_s] = \prod_{l=1}^{n} x_{i_l}$ be the corresponding free random variable (6.2.9) of (B, φ) . If

$$\pi_{(I_s)} = 1_{|V_1|} \lor \dots \lor 1_{|V_t|},$$

in the sense of (6.2.7) and (6.2.8), then

$$\varphi\left(X[I_s]\right) = \begin{cases} t & \text{if } |V_k| \in 2\mathbb{N}, \\ \prod_{i=1}^{l} c_{\frac{|V_i|}{2}} & \text{for all } k = 1, ..., t \\ 0 & \text{otherwise,} \end{cases}$$
(6.2.15)

where c_k are the k-th Catalan numbers for all $k \in \mathbb{N}_0$.

Proof. Under hypothesis, $\varphi(X[I_s])$

$$= \begin{cases} |NC_{2}(V_{1}) \times ... \times NC_{2}(V_{t})| & \text{ if } |V_{k}| \in 2\mathbb{N}, \\ & \text{ for all } k = 1, ..., t \\ 0 & \text{ otherwise,} \end{cases}$$

by (6.2.14).

Recall that, for every countable set X, with $|X| \in 2\mathbb{N}$, the set

$$NC_{2}(X) = \{\theta \in NC(X) : \forall V \in \theta, |V| = 2\}$$

is equipotent (or bijective) to the noncrossing-partition lattice $NC\left(\frac{|X|}{2}\right)$ over $\{1, ..., \frac{|X|}{2}\}$ (e.g., [5] and [8]). i.e., if $|V_k| \in 2\mathbb{N}$, then

$$|NC_2(V_k)| = \left|NC\left(\frac{|V_k|}{2}\right)\right|, \qquad (6.2.16)$$

for all k = 1, ..., t. So, we have $\varphi(X[I_s])$

$$= \begin{cases} \left| NC\left(\frac{|V_{1}|}{2}\right) \times ... \times NC\left(\frac{|V_{t}|}{2}\right) \right| & \text{if } |V_{k}| \in 2\mathbb{N}, \\ \text{for all } k = 1, ..., t \\ 0 & \text{otherwise}, \end{cases}$$
$$= \begin{cases} \prod_{l=1}^{t} c_{\frac{|V_{l}|}{2}} & \text{if } |V_{l}| \in 2\mathbb{N}, \text{ for all } l = 1, ..., t \\ 0 & \text{otherwise}, \end{cases}$$
(6.2.17)

by (6.2.16), because $|NC(X)| = c_{|X|}$, for all finite sets X (e.g., [7,8,14,17]). Therefore, the formula (6.2.15) holds by (6.2.17).

Remark 6.1. The more combinatorial computational techniques, and the refined results of (6.2.15) are considered "analytically" in [7], including direct estimations, and asymptotic estimations of (6.2.15). However, in this paper, the free-distributional data (6.2.15) is enough for our purposes. The importance here is that the free-distributional data induced by mutually free, multi semicircular elements are dictated by the semicircularity (3.2), by (6.2.4) and (6.2.15).

We provide some examples before finishing this section.

Example 6.1. (1) Let x_1 , x_2 , x_3 , x_4 be mutually free semicircular elements of (B, φ) , and let

$$W = x_1^2 x_2^4 x_1^2 x_3^2 \in (B, \varphi)$$

be a free reduced word with its length-4. Then one can take

$$I_W = (1, 1, 2, 2, 2, 2, 1, 1, 3, 3) \stackrel{let}{=} (i_1, ..., i_{10}),$$

and

$$\pi_{(I_W)} = \{(i_1, i_2, i_7, i_8), (i_3, i_4, i_5, i_6), (i_9, i_{10})\},\$$

with

$$V_1 = \{i_1, i_2, i_7, i_8\} = \{1, 1, 1, 1\}, V_2 = \{i_3, i_4, i_5, i_6\} = \{2, 2, 2, 2\},$$

and

$$V_3 = \{i_7, \, i_8\} = \{3, \, 3\}$$

Therefore,

$$\varphi(W) = c_{\frac{4}{2}}c_{\frac{4}{2}}c_{\frac{2}{2}} = c_{2}^{2}c_{1} = 4.$$
(2) Meanwhile, if $W = x_{1}^{2}x_{3}x_{2}x_{4}x_{2}^{2}x_{1} \in (B, \varphi)$, then

$$I_W = (1, 1, 3, 2, 4, 2, 2, 1)$$

and

$$\pi_{(I_8)} = \{ (i_1, i_2, i_8), (i_3), (i_4, i_6, i_7), (i_5) \} \\ = \{ (1, 1, 1), (3), (2, 2, 2), (4) \}$$

satisfying that

$$\varphi\left(X[I_8]\right) = \varphi\left(x_1^2 x_3 x_2 x_4 x_2^2 x_1\right) = c_{\frac{3}{2}} c_{\frac{1}{2}} c_{\frac{3}{2}} c_{\frac{1}{2}} = 0,$$

by (6.2.15). (3) Now, let $W = x_1^4 x_2^4 x_1^6 x_2^4 \in (B, \varphi)$. Then one can take

$$I_W = (i_1, i_2, \dots, i_{18}),$$

having

$$\pi_{(I_{18})} = \{V_1, V_2, V_3\},\$$

with

$$V_1 = \{i_1, i_2, i_3, i_4, i_9, i_{10}, i_{11}, i_{12}, i_{13}, i_{14}\},\$$

and

$$V_2 = \{i_5, i_6, i_7, i_8\}, V_3 = \{i_{15}, i_{16}, i_{17}, i_{18}\}.$$

(Here, since all entries of V_2 and V_3 are identical to 2, one may/can be tempted to make a block

$$\{i_5, i_6, i_7, i_8, i_{15}, i_{16}, i_{17}, i_{18}\},\$$

but, in such a case, this block has crossing with V_1 , disobeying the conditions (i) and (ii)!)

Therefore,

$$\varphi\left(x_1^4 x_2^4 x_1^6 x_3^4\right) = c_{\frac{10}{2}} c_{\frac{4}{2}} c_{\frac{4}{2}} = 168,$$

by (6.2.15).

6.3 Free-Distributional Data on \mathbb{L}_Q

Let \mathbb{L}_Q be our semicircular filterization (6.1.5) of the free filterization $\mathfrak{L}_Q(\mathbb{Z})$, generated by the free semicircular family S of (6.1.4). By the structure theorem (6.1.7), all free random variables of \mathbb{L}_Q are the limits of linear combinations of free reduced words, formed by

$$W = \prod_{l=1}^{N} U_{j_{l}}^{n_{l}}, \text{ for } U_{j_{l}} \in \mathcal{S}, \forall l = 1, ..., N,$$
(6.3.1)

in S, for all $N \in \mathbb{N}$, where $n_1, ..., n_N \in \mathbb{N}$, and the N-tuple $(j_1, ..., j_N)$ is alternating in \mathbb{Z} .

Theorem 6.5. Let W be a free reduced word (6.3.1) of \mathbb{L}_Q in S.

(6.3.2) If
$$j_1 = j_2 = ... = j_N$$
 in (6.3.1), then $\tau(W)$ is characterized by (6.2.4).

(6.3.3) If $(j_1, ..., j_N)$ is mixed in (6.3.1), then $\tau(W)$ is determined by (6.2.15).

Proof. Note that all semicircular elements of any topological *-probability spaces have the same free distribution, "the" semicircular law, characterized by the free-moment sequence

 $(0, c_1, 0, c_2, 0, c_3, 0, c_4, \ldots),$

equivalently, the free-cumulant sequence

$$(0, 1, 0, 0, 0, 0, \ldots),$$

by (3.2) and (3.4), where c_k are the k-th Catalan numbers for all $k \in \mathbb{N}$.

By this universality of the semicircular law (or, by the identically-freedistributedness of all semicircular elements in terms of [19]), the statements (6.3.2) and (6.3.3) are shown by (6.2.4) and (6.2.15), respectively.

The above theorem characterizes the free-distributional data on the semicircular filterization \mathbb{L}_Q , in terms of joint free moments of generating semicircular elements of \mathcal{S} , by (6.3.2) and (6.3.3).

7 Integer-Shifts on \mathbb{L}_Q

In this section, let (A, ψ) be the fixed C^{*}-probability space containing a family $\mathbf{Q} = \{q_j\}_{j \in \mathbb{Z}}$ of mutually-orthogonal projections q_j 's having

$$\psi(q_j) \in \mathbb{R}^{\times}$$
, for all $j \in \mathbb{Z}$,

and let \mathbb{L}_{Q} be the semicircular filterization.

7.1 (±)-Shifts on \mathbb{Z}

Let \mathbb{Z} be the set of all integers. Define functions h_+ and h_- on \mathbb{Z} by

$$h_+(j) = j + 1, \tag{7.1.1}$$

and

$$h_{-}(j) = j - 1,$$

for all $j \in \mathbb{Z}$. By the definition (7.1.1), these two functions h_{\pm} are well-defined bijections on \mathbb{Z} , satisfying $h_{+}^{-1} = h_{-}$, where f^{-1} mean the functional inverses of invertible functions f.

Then, for these bijections h_{\pm} of (7.1.1), one can construct the bijections $h_{\pm}^{(n)}$ on \mathbb{Z} ,

$$h_{\pm}^{(n)} = \underbrace{h_{\pm} \circ h_{\pm} \circ \cdots \circ h_{\pm}}_{n\text{-times}}, \tag{7.1.2}$$

for all $n \in \mathbb{N}$, with identities, $h_{\pm}^{(1)} = h_{\pm}$, where (\circ) is the usual functional composition. It is easy to check that

$$h^{(n)}_{\pm}(j) = j \pm n$$
, for all $j \in \mathbb{Z}$,

for all $n \in \mathbb{N}$.

Definition 7.1. We call the functions $h_{\pm}^{(n)}$ of (7.1.2), the n-(\pm)-shifts on \mathbb{Z} , for $n \in \mathbb{N}$.

7.2 Integer-Shifts on \mathbb{L}_Q

Let $h_{\pm}^{(n)}$ be n-(\pm)-shifts of (7.1.2) on \mathbb{Z} , for all $n \in \mathbb{N}$. Define now a multiplicative bounded linear transformation β_{\pm} on \mathbb{L}_Q by the morphisms satisfying that:

$$\beta_{\pm} (U_j) = U_{h_{\pm}(j)}, \tag{7.2.1}$$

for $U_j \in S$, for all $j \in \mathbb{Z}$, where S is our free semicircular family (6.1.4) of $\mathfrak{L}_Q(\mathbb{Z})$, generating \mathbb{L}_Q .

By (6.1.6) and (6.1.7), the above multiplicative linear transformation β_{\pm} of (7.2.1) is well-defined on \mathbb{L}_Q .

Lemma 7.1. Let $Y = \prod_{l=1}^{N} U_{j_l}^{n_l} \in \mathbb{L}_Q$, for $U_{j_1}, ..., U_{j_N} \in S$, and $n_1, ..., n_N \in \mathbb{N}$, for $N \in \mathbb{N}$. Then

$$\beta_{\pm}(Y) = \prod_{l=1}^{N} U_{h\pm(j_l)}^{n_l}.$$
(7.2.2)

Proof. If $Y \in \mathbb{L}_Q$ be as above, then, by the multiplicativity of β_{\pm} , one has that

$$\beta_{\pm}(Y) = \prod_{l=1}^{N} \beta_{\pm} \left(U_{j_{l}}^{n_{l}} \right) = \prod_{l=1}^{N} \left(\beta_{\pm} \left(U_{j_{l}} \right) \right)^{n_{l}} = \prod_{l=1}^{N} U_{h_{\pm}(j_{l})}^{n_{l}}.$$

So, the formula (7.2.2) holds.

Let $u_{j_1}, ..., u_{j_N} \in \mathcal{X}$ be weighted-semicircular elements of \mathbb{L}_Q , for $N \in \mathbb{N}$, where \mathcal{X} is the free weighted-semicircular family (6.1.3), generating \mathbb{L}_Q , and let

$$X = \prod_{l=1}^{N} u_{j_l}^{n_l}, \text{ for } n_1, ..., n_N \in \mathbb{N}.$$

Then one has

$$\beta_{\pm}(X) = \beta_{\pm} \left(\left(\prod_{l=1}^{N} \psi(q_{j_l})^{n_l} \right) \left(\prod_{l=1}^{N} U_{j_l}^{n_l} \right) \right)$$

since

$$U_{j_l} = \frac{1}{\psi(q_{j_l})} u_{j_l} \in \mathcal{S} \iff u_{j_l} = \psi(q_{j_l}) U_{j_l} \in \mathcal{X},$$

so, the above formula goes to

$$= \left(\prod_{l=1}^{N} \psi(q_{j_l})^{n_l}\right) \beta_{\pm} \left(\prod_{l=1}^{N} U_{j_l}^{n_l}\right)$$
$$= \left(\prod_{l=1}^{N} \psi(q_{j_l})^{n_l}\right) \left(\prod_{l=1}^{N} U_{h_{\pm}(j_l)}^{n_l}\right), \tag{7.2.2}'$$
by (7.2.2)

By
$$(7.2.2)'$$
, one can see that the freeness on \mathbb{L}_Q is preserved to that on \mathbb{L}_Q .

Theorem 7.2. The multiplicative linear transformations β_{\pm} of (7.2.1) are *isomorphisms on \mathbb{L}_Q .

Proof. By (6.1.5) and (6.1.6), each element of the semicircular filterization \mathbb{L}_Q is a limit of linear combinations of free reduced words in the generating free semicircular family \mathcal{S} . So, we focus on free reduced words of \mathbb{L}_Q in \mathcal{S} .

Let $(j_1, ..., j_N)$ be an alternating N-tuple in \mathbb{Z} for $N \in \mathbb{N}$, and let

$$Y = \prod_{l=1}^{N} U_{j_l}^{n_l}, \text{ for } n_1, ..., n_N \in \mathbb{N}.$$

be a free reduced word with its length-N in \mathbb{L}_Q by (6.1.7). Then, by (7.2.2),

$$\beta_{\pm}(Y) = \prod_{l=1}^{N} U_{h\pm(j_l)}^{n_l}, \qquad (7.2.3)$$

are free reduced words with their lengths-N in \mathbb{L}_Q , where h_{\pm} are the (\pm) -shifts (7.1.1) on \mathbb{Z} , since

$$(h_{\pm}(j_1), ..., h_{\pm}(j_N)) = (j_1 \pm 1, ..., j_N \pm 1)$$

are alternating in \mathbb{Z} , too. Also, if Y is as above, then

$$\beta_{\pm}(Y^*) = \beta_{\pm} \left(\prod_{l=1}^{N} U_{j_{N-l+1}}^{n_{N-l+1}} \right)$$

by the self-adjointness of U_{j_1} , ..., U_{j_N}

$$=\prod_{l=1}^{N} U_{h_{\pm}(j_{N-l+1})}^{n_{N-l+1}}$$

by

$$= \left(\prod_{l=1}^{N} U_{h_{\pm}(j_{l})}^{n_{l}}\right)^{*} = \left(\beta_{\pm}(Y)\right)^{*},$$

showing that

$$\beta_{\pm}(S^*) = \left(\beta_{\pm}(S)\right)^*, \text{ for all } S \in \mathbb{L}_Q.$$
(7.2.4)

By (7.2.4), the bijective bounded multiplicative linear transformations β_{\pm} of (7.2.1) are well-defined *-isomorphisms on \mathbb{L}_Q .

The above theorem illustrates that the (±)-shifts h_{\pm} of (7.1.1) on \mathbb{Z} induce the *-isomorphisms β_{\pm} of (7.2.2) on \mathbb{L}_Q .

Definition 7.2. The *-isomorphisms β_{\pm} of (7.2.1) are said to be (\pm) -integershift(-*-isomorphism)s on \mathbb{L}_Q .

These two *-isomorphisms β_{\pm} satisfy the following identity relation on \mathbb{L}_Q .

Lemma 7.3. The (\pm) -integer shifts β_{\pm} satisfy

$$\beta_{+}\beta_{-} = 1_{\mathbb{L}_{Q}} = \beta_{-}\beta_{+} on \ \mathbb{L}_{Q}, \tag{7.2.5}$$

where $1_{\mathbb{L}_Q}$ is the identity map on \mathbb{L}_Q .

Proof. It is enough to consider the cases where we have free reduced words formed by

$$Y = \prod_{l=1}^{N} U_{j_l}^{n_l} \text{ of } \mathbb{L}_Q, \text{ for } n_1, ..., n_N \in \mathbb{N},$$

for $N \in \mathbb{N}$, where $(j_1, ..., j_N)$ is alternating in \mathbb{Z} . Observe that

$$\beta_{+}\beta_{-}(Y) = \beta_{+} \left(\prod_{l=1}^{N} U_{h_{-}(j_{l})}^{n_{l}}\right) = \beta_{+} \left(\prod_{l=1}^{N} U_{j_{l}-1}^{n_{l}}\right)$$
$$= \prod_{l=1}^{N} U_{h_{+}(j_{l}-1)}^{n_{l}} = \prod_{l=1}^{N} U_{j_{l}-1+1}^{n_{l}} = Y.$$

Similarly,

$$\beta_{-}\beta_{+}(Y) = Y$$
, in \mathbb{L}_Q .

So, for any arbitrary operators $S \in \mathbb{L}_Q$,

$$\beta_+\beta_-(S) = S = \beta_-\beta_+(S), \text{ in } \mathbb{L}_Q.$$

Thus, the relation (7.2.5) holds.

From the (±)-shifts β_{\pm} on \mathbb{L}_Q , construct the *-isomorphisms β_{\pm}^n ,

$$\beta_{\pm}^{n} = \underbrace{\beta_{\pm}\beta_{\pm}\cdots\cdots\beta_{\pm}}_{n\text{-times}} \text{ on } \mathbb{L}_{Q}, \qquad (7.2.6)$$

for all $n \in \mathbb{N}_0 = \mathbb{N} \cup \{0\}$, with axiomatization:

$$\beta^0_+ = \mathbb{1}_{\mathbb{L}_O} = \beta^0_-$$

Since β_{\pm} and $1_{\mathbb{L}_Q}$ are well-defined *-isomorphisms, the morphisms β_{\pm}^n of (7.2.9) are well-defined *-isomorphisms on \mathbb{L}_Q , too, for all $n \in \mathbb{N}_0$.

Definition 7.3. Let β_{\pm}^n be the *-isomorphisms (7.2.6) on the semicircular filterization \mathbb{L}_Q , for all $n \in \mathbb{N}_0$, with axiomatization $\beta_{\pm}^0 = 1_{\mathbb{L}_Q}$. Then we call them the n- (\pm) -(integer-)shifts on \mathbb{L}_Q , for all $n \in \mathbb{N}_0$.

By (7.2.5) and (7.2.6), we obtain the following relation on the set $\{\beta_{\pm}^{n} : n \in \mathbb{N}_{0}\}$.

Lemma 7.4. Let β_{\pm}^n be the n- (\pm) -shifts on \mathbb{L}_Q , for $n \in \mathbb{N}_0$. Then

$$\beta_{+}^{n_{1}}\beta_{-}^{n_{2}} = \beta_{-}^{n_{2}}\beta_{+}^{n_{1}} = \begin{cases} 1_{\mathbb{L}_{Q}} & \text{if } n_{1} = n_{2} \\ \beta_{+}^{n_{1}-n_{2}} & \text{if } n_{1} > n_{2} \\ \beta_{-}^{n_{2}-n_{1}} & \text{if } n_{1} < n_{2}, \end{cases}$$
(7.2.7)

on \mathbb{L}_Q , for all $n_1, n_2 \in \mathbb{N}_0$. And

$$\beta_{+}^{n_{1}}\beta_{+}^{n_{2}} = \beta_{+}^{n_{1}+n_{2}}, and \ \beta_{-}^{n_{1}}\beta_{-}^{n_{2}} = \beta_{-}^{n_{1}+n_{2}}, \tag{7.2.8}$$

on \mathbb{L}_Q , for all $n_1, n_2 \in \mathbb{N}_0$.

Proof. The formulas (7.2.7) and (7.2.8) are proven by the straightforward computations.

The above relations (7.2.7) and (7.2.8) can be re-expressed as follows;

$$\beta_{e_1}^{n_1}\beta_{e_2}^{n_2} = \beta_{e_2}^{n_2}\beta_{e_1}^{n_1} = \beta_{sgn(e_1n_1+e_2n_2)}^{|e_1n_1+e_2n_2|} \text{ on } \mathbb{L}_Q,$$
(7.2.9)

with

$$sgn(e_1n_1 + e_2n_2) = \begin{cases} + \text{ if } e_1n_1 + e_2n_2 \ge 0\\ - \text{ if } e_1n_1 + e_2n_2 \le 0, \end{cases}$$

for all $e_1, e_2 \in \{\pm\}$, and $n_1, n_2 \in \mathbb{N}_0$, where sgn is the sign map,

$$sgn(j) = \begin{cases} + & \text{if } j \ge 0 \\ - & \text{if } j < 0, \end{cases}$$

for all $j \in \mathbb{Z}$, and |.| is the *absolute value* on \mathbb{Z} . Now, let

$$\mathfrak{B} = \{\beta^n_\pm\}_{n \in \mathbb{N}_0}.\tag{7.2.10}$$

Then it is a subset of the *automorphism group*,

$$Aut\left(\mathbb{L}_Q\right) = \left(\left\{ \alpha : \mathbb{L}_Q \to \mathbb{L}_Q \middle| \begin{array}{c} \alpha \text{ are} \\ * \text{ -isomorphisms} \\ \text{ on } \mathbb{L}_Q \end{array} \right\}, \cdot \right), \quad (7.2.11)$$

of all *-isomorphisms on \mathbb{L}_Q , where the operation (.) is the product of *-isomorphisms.

By (7.2.10), this system \mathfrak{B} is clearly a "subset" of the automorphism group $Aut(\mathbb{L}_Q)$ of (7.2.11). Note that, by (7.2.9),

$$\left(\beta_{e_1}^{n_1}\beta_{e_2}^{n_2}\right)\beta_{e_3}^{n_3} = \beta_{sgn(e_1n_1e_2n_2)}^{|e_1n_1e_2n_2|}\beta_{e_3}^{n_3} = \beta_{sgn(e_1n_1e_2n_2e_3n_3)}^{|e_1n_1e_2n_2e_3n_3|} = \beta_{e_1}^{n_1}\beta_{sng(e_2n_2e_3n_3)}^{|e_2n_2e_3n_3|} = \beta_{e_1}^{n_1}\left(\beta_{e_2}^{n_2}\beta_{e_3}^{n_3}\right),$$
(7.2.12)

on \mathbb{L}_Q , for all $e_1, e_2, e_3 \in \{\pm\}, n_1, n_2, n_3 \in \mathbb{N}_0$.

Theorem 7.5. Let \mathfrak{B} be the subset (7.2.10) of $Aut(\mathbb{L}_Q)$. Then

(7.2.13)
$$\mathfrak{B}$$
 is a subgroup of $Aut(\mathbb{L}_Q)$

Proof. Let \mathfrak{B} be the set (7.2.10). By (7.2.9), the operation (·) is closed on \mathfrak{B} . By (7.2.12), this operation is associative on \mathfrak{B} .

Since

$$\beta^0_+ = 1_{\mathbb{L}_Q} = \beta^0_- \in \mathfrak{B},$$

and

$$\beta_e^n \cdot 1_{\mathbb{L}_Q} = \beta_e^n = 1_{\mathbb{L}_Q} \cdot \beta_e^n \text{ on } \mathbb{L}_Q,$$

for all $e \in \{\pm\}$, and $n \in \mathbb{N}_0$, the set \mathfrak{B} contains the group-identity $1_{\mathbb{L}_Q}$ of $Aut(\mathbb{L}_Q)$.

By (7.2.7), all elements $\beta_{\pm}^n \in \mathfrak{B}$ have their unique (·)-inverses $\beta_{\mp}^n \in \mathfrak{B}$, such that

$$\beta_+^n \beta_-^n = 1_{\mathbb{L}_Q} = \beta_-^n \beta_+^n \text{ on } \mathbb{L}_Q,$$

for all $n \in \mathbb{N}_0$,

Therefore, the system \mathfrak{B} is a subgroup of $Aut(\mathbb{L}_Q)$.

As a subgroup of $Aut(\mathbb{L}_Q)$, the group \mathfrak{B} satisfies the following algebraic property.

Theorem 7.6. Let \mathfrak{B} be the subgroup (7.2.10) of $Aut(\mathbb{L}_Q)$. Then

$$\mathfrak{B} \stackrel{Group}{=} (\mathbb{Z}, +), \tag{7.14}$$

where " $\overset{Group}{=}$ " means "being group-isomorphic," where $(\mathbb{Z}, +)$ is the infinite cyclic abelian group.

Proof. Define a function $\Phi : \mathbb{Z} \to \mathfrak{B}$ by

$$\Phi: j \in \mathbb{Z} \longmapsto \beta_{sgn(j)}^{|j|} \in \mathfrak{B}, \tag{7.2.15}$$

with correspondence

$$0 \in \mathbb{Z} \longmapsto 1_{\mathbb{L}_Q} = \beta_{\pm}^0 \in \mathfrak{B}.$$

Then this function Φ of (7.2.15) is a well-defined bijection from \mathbb{Z} onto \mathfrak{B} , by (7.2.6). And it satisfies that

$$\Phi(j_1 + j_2) = \beta_{sgn(j_1 + j_2)}^{|j_1 + j_2|} = \beta_{sgn(j_1)}^{|j_1|} \beta_{sgn(j_2)}^{|j_2|}
= \Phi(j_1) \Phi(j_2),$$
(7.2.16)

in \mathfrak{B} , for all $j_1, j_2 \in \mathbb{Z}$.

So, the bijection Φ is a group-homomorphism by (7.2.16), i.e., the relation (7.14) holds.

The above theorem characterizes the algebraic structure of the group $\mathfrak{B} = \{\beta_{\pm}^n\}_{n\in\mathbb{N}_0}$ in the automorphism group $Aut(\mathbb{L}_Q)$.

Definition 7.4. The subgroup \mathfrak{B} of the automorphism group $Aut(\mathbb{L}_Q)$ is called the integer-shift (sub)group (of $Aut(\mathbb{L}_Q)$ acting) on \mathbb{L}_Q .

7.3 Free Distributions on \mathbb{L}_Q Affected by \mathfrak{B}

Let \mathfrak{B} be the integer-shift group (7.2.10) acting on the semicircular filterization \mathbb{L}_Q of Q. We here consider how the action of our *-isomorphisms $\beta_{\pm}^n \in \mathfrak{B}$ affects the original free-distributional data on the semicircular filterization \mathbb{L}_Q .

Take an arbitrary free reduced word,

$$Y = \prod_{l=1}^{N} U_{j_l}^{n_l} in \mathbb{L}_Q \tag{7.3.1}$$

where $U_{j_l} = \frac{1}{\psi(q_{j_l})} u_{j_l} \in \mathcal{S}$ are the generating semicircular elements of \mathbb{L}_Q , where $u_{j_l} \in \mathcal{X}$ are the $\psi(q_j)^2$ -semicircular elements, for all l = 1, ..., N, for $N \in \mathbb{N}$, and where the N-tuple $(j_1, ..., j_N)$ is alternating in \mathbb{Z} , and $n_1, ..., n_N \in \mathbb{N}$. **Theorem 7.7.** Let Y be a free reduced word (7.3.1) of \mathbb{L}_Q in S. Then

$$\tau\left(\beta_e^k(Y)\right) = \tau(Y),\tag{7.3.2}$$

for all $e \in \{\pm\}$ and $k \in \mathbb{N}_0$.

Proof. First assume that N = 1, and hence, $Y = U_{j_1}^{n_1}$ in \mathbb{L}_Q . Then, by the semicircularity of $U_{j_1}, U_{j_1ek} \in S$ in \mathbb{L}_Q , one has that

$$\tau\left(\beta_e^k(Y)\right) = \tau\left(U_{jek}^{n_1}\right) = \omega_{n_1} c_{\frac{n_1}{2}} = \tau\left(U_{j_1}^{n_1}\right),\tag{7.3.3}$$

for all $\beta_e^k \in \mathfrak{B}$.

Assume now that N > 1 in \mathbb{N} , and the free reduced word $Y \in \mathbb{L}_Q$ with its length-N is in the sense of (7.3.1). Note that the image

$$\beta_e^k(Y) = \prod_{l=1}^N U_{j_lek}^{n_k}$$

is again a free reduced word with the same length-N in \mathbb{L}_Q , for all $\beta_e^k \in \mathfrak{B}$. Now, let I_s be the s-tuple of (6.2.5), satisfying

$$Y = X[I_s]$$
 in \mathbb{L}_Q , for some $s \ge N$,

where $X[I_s]$ is in the sense of (6.2.9). Similarly, let $I_{s'}$ be the s'-tuple of (6.2.5), satisfying

$$\beta_e^k(Y) = X[I_{s'}] \text{ in } \mathbb{L}_Q,$$

for $\beta_e^k \in \mathfrak{B}$, where $X[I_{s'}]$ is in the sense of (6.2.9).

Then, since Y and $\beta_e^k(Y)$ are the same-length free reduced words having the same free-ness structure, one has not only

$$s = s' \text{ in } \mathbb{N},$$

but also

$$\pi(I_s) = \pi(I_{s'}) \text{ in } NC\left([I_s]\right) \stackrel{\text{lattice}}{=} NC\left([I_{s'}]\right),$$

where $\pi(I_s)$ and $\pi(I_{s'})$ are the noncrossing partitions of (6.2.7), and " $\stackrel{\text{lattice,m}}{=}$ " means "being lattice-isomorphic." (Recall that if X and Y are finite discrete sets, then NC(X) and NC(Y) are lattice-isomorphic, if and only if |X| = |Y| in \mathbb{N}).

By the semicircularity (7.3.3) of $U_{j_1}, ..., U_{j_N}, U_{j_1ek}, ..., U_{j_Nek} \in S$ in \mathbb{L}_Q , we have

$$\tau\left(\beta_{e}^{k}(Y)\right) = \tau\left(X[I_{s'}]\right) = \tau\left(X[I_{s}]\right) = \tau\left(Y\right),$$

by (6.2.15) or (6.3.3), for all $\beta_e^k \in \mathfrak{B}$.

Therefore, the statement (7.3.2) holds.

The above theorem shows how the original free distributional data on the semicircular filterization \mathbb{L}_Q is affected by the group-action of the integer-shifts of \mathfrak{B} on \mathbb{L}_Q . i.e., \mathfrak{B} preserves the free probability on \mathbb{L}_Q , by (7.3.2).

8 Semicircular Elements Induced by Multi Projections

Now, we have all ingredients for studying our main interests. In this section, we show that if there are N-many mutually orthogonal projections in an arbitrary C^* -probability space, then there exists a corresponding free semicircular family $\mathcal{S}^{(N)}$, induced by the projections in a certain free product Banach *-probability space $\mathbb{L}_{O}^{(N)}$, for any

$$N \in \mathbb{N}_{>1}^{\infty} = (\mathbb{N} \setminus \{1\}) \cup \{\infty\}.$$

And then, consider certain *-homomorphisms acting on $\mathbb{L}_Q^{(N)}$, and study how they deform the free probability on $\mathbb{L}_Q^{(N)}$.

8.1 A Free Semicircular Family $S^{(N)}$ Induced by N-many Projections

Let (A_o, ψ_o) be a C^* -probability space containing its N-many mutually orthogonal projections

$$\mathbf{Q}_o = \{q_1^o, ..., q_N^o\}$$
(8.1.1)

for $N \in \mathbb{N}_{>1}^{\infty}$, and let

$$Q_o = C^* \left(\mathbf{Q}_o \right) \subseteq A_o \tag{8.1.2}$$

be the C^* -subalgebra of A generated by the family \mathbf{Q}_o of (8.1.1). Suppose

$$\psi_o(q_k^o) \in \mathbb{R}^{\times} \text{ in } \mathbb{C}, \forall k = 1, ..., N.$$
 (8.1.3)

Proposition 8.1. Let Q_o be the C^{*}-subalgebra (8.1.2). Then

$$Q_o \stackrel{*-iso}{=} \stackrel{N}{\bigoplus}_{l=1} (\mathbb{C} \cdot q_l^o) \stackrel{*-iso}{=} \mathbb{C}^{\oplus N}.$$
(8.1.4)

Proof. Since the generating set \mathbf{Q}_o of Q_o consists of mutually orthogonal *N*-many projections q_1^o , ..., q_N^o , the structure theorem (8.1.4) is immediately proven.

Suppose there is a C^* -probability space (A, ψ) containing a family $\mathbf{Q} = \{q_j\}_{j \in \mathbb{Z}}$ of mutually orthogonal $|\mathbb{Z}|$ -many projections q_j 's, satisfying

 $\psi(q_j) \in \mathbb{R}^{\times}$ in \mathbb{C} , for all $j \in \mathbb{Z}$.

Assume further that there exist projections $q_{j_1}, ..., q_{j_N} \in \mathbf{Q}$, such that

$$\psi\left(q_{j_l}\right) = \psi_o\left(q_l^o\right) \text{ in } \mathbb{R}^\times, \tag{8.1.5}$$

for all l = 1, ..., N, where ψ_o is the linear functional on the C^{*}-algebra A_o , satisfying the condition (8.1.3).

For convenience, we re-index the subfamily

$$\{q_{j_1}, ..., q_{j_N}\}$$
 of **Q** (8.1.6)

by

 $\{q_1, ..., q_N\}$ in **Q**,

without loss of generality, from below.

Theorem 8.2. Let Q_o be the C^{*}-subalgebra (8.1.2) of a fixed C^{*}-probability space (A_o, ψ_o) , generated by the family \mathbf{Q}_o of (8.1.1). Then, under (8.1.5) and (8.1.6), there exists a Banach *-subalgebra

$$\mathbb{L}_Q^{(N)} \stackrel{*\text{-}iso}{=} \stackrel{N}{\underset{l=1}{\overset{\times}{\leftarrow}}} \overline{\mathbb{C}[\{U_l\}]}$$

of the semicircular filterization $\mathbb{L}_{\mathcal{O}}$ of (6.1.5).

Proof. Let $Q = C^*(\mathbf{Q})$ be the C^{*}-subalgebra (3.7) of the C^{*}-probability space (A, ψ) satisfying (8.1.5), under the re-indexing process (8.1.6). First, define a linear morphism

by

$$\Psi\left(\sum_{l=1}^{N} t_l q_l^o\right) \stackrel{def}{=} \sum_{l=1}^{N} t_l q_l + \sum_{j \in \mathbb{Z} \setminus \{1, \dots, N\}} 0 \cdot q_j.$$

Then it is an injective *-homomorphism from Q_o to Q, by (3.8) and (8.1.4). Thus, one can construct the semicircular elements

$$U_l = \mathbf{l} \otimes q_l = \mathbf{l} \otimes \Psi(q_l^o) \in \mathbb{L}_Q, \qquad (8.1.7)$$

in the free semicircular family \mathcal{S} of (6.1.4), generating the semicircular filterization \mathbb{L}_Q , for all l = 1, ..., N.

By the structure theorem (6.1.7) of \mathbb{L}_Q , one can define the Banach *subalgebra

$$\mathbb{L}_{Q}^{(N)} \stackrel{def}{=} \overline{\mathbb{C}[\{U_{1}, ..., U_{N}\}]}$$

$$\stackrel{*\text{-iso}}{=} \overline{\mathbb{C}\left[\{\mathbf{l} \otimes \Psi(q_{l}^{o}) : l = 1, ..., N\}\right]}$$

$$\stackrel{*\text{-iso}}{=} \frac{N}{\star} \overline{\mathbb{C}[\{\mathbf{l} \otimes \Psi(q_{j}^{o})\}]} = \overset{N}{\underset{l=1}{\overset{N}{\star}} \overline{\mathbb{C}}[\{U_{l}\}]}$$

$$(8.1.8)$$

of \mathbb{L}_Q , by (8.1.7).

$$\Psi: Q_o \to Q$$

$$\Psi: Q_o \to \zeta$$

It shows that: if the condition (8.1.5) is satisfied under (8.1.6), then the family \mathbf{Q}_o of (8.1.1) induces Banach *-probability space,

$$\mathbb{L}_Q^{(N)} = \left(\mathbb{L}_Q^{(N)}, \ \tau = \tau \mid_{\mathbb{L}_Q^{(N)}}\right),$$

generated by the free semicircular family

$$\mathcal{S}^{(N)} = \{ U_l = \mathbf{l} \otimes \Psi(q_l^o) \}_{l=1}^N, \tag{8.1.9}$$

as a free-probabilistic sub-structure of the semicircular filterization \mathbb{L}_Q .

The above theorem shows that if the conditions (8.1.3) and (8.1.5) are satisfied under the re-indexing process (8.1.6), then the family \mathbf{Q}_o of (8.1.1) induces a Banach *-probability space $\mathbb{L}_Q^{(N)}$ of (8.1.8) generated by the free semicircular family $\mathcal{S}^{(N)}$ of (8.1.9), as an embedded free-probabilistic sub-structure of the semicircular filterization \mathbb{L}_Q of (6.1.5).

Remark 8.1. As we briefly discussed in [6], whenever such a family \mathbf{Q}_o of (8.1.1) in a C^{*}-probability space (A_o, ψ_o) is fixed, in fact, one can construct the corresponding C^{*}-probability space (A, ψ) , having its family \mathbf{Q} of mutually orthogonal $|\mathbb{Z}|$ -many projections, artificially-but-naturally.

Assume first that $N < \infty$ in $\mathbb{N}_{>1}^{\infty}$. If $Q_o = C^*(\mathbf{Q}_o)$ is the C^{*}-subalgebra (8.1.2) of A_o , satisfying (8.1.4), then one can construct the direct product C^{*}-algebra Q,

$$Q = \bigoplus_{k \in \mathbb{Z}} Q_{o,k} \stackrel{*-iso}{=} \mathbb{C}^{\oplus |\mathbb{Z}|}, \text{ with } Q_{o,k} = Q_o,$$

equipped with its linear functional ψ ,

$$\psi = \bigoplus_{k \in \mathbb{Z}} \psi_{o,k} \text{ on } Q, \text{ with } \psi_{o,k} = \psi_o.$$

So, in the above C^* -probability space (Q, ψ) (or a C^* -probability space (A, ψ) with $A \supseteq Q$), there exist infinitely many projections q, such that

 $\psi(q) = \psi_o(q_l), \text{ for some } l \in \{1, ..., N\}.$

Assume now that $N = \infty$ in $\mathbb{N}_{>1}^{\infty}$, and let

$$Q_o = C^* \left(\mathbf{Q}_o \right) = C^* \left(\{ q_1^o, q_2^o, q_3^0, \ldots \} \right).$$

Then, for convenience, by the canonical re-indexing, one can let

 $Q_o = C^* (\{q_0^o, q_1^o, q_2^o, \ldots\}),$

by identifying q_i^o with q_{i-1}^o , for all $i \in \mathbb{N}$. Now, consider a subfamily \mathbf{Q}'_o of \mathbf{Q}_o ,

$$\mathbf{Q}_o' = \{q_{-1}^o, q_{-2}^o, \ldots\} = \{q_0^o, q_1^o, q_2^o, \ldots\} \setminus \{q_0^o\},\$$

with identity:

 $q_{-k}^o = q_k^o$, for all $k \in \mathbb{N}$,

of the re-indexed family \mathbf{Q}_o , and construct

 $Q' = C^* \left(\mathbf{Q}_o' \right).$

Then we have the direct product C^* -algebra

$$Q = Q_o \oplus Q'_o \stackrel{*\text{-}iso}{=} \mathbb{C}^{\oplus |\mathbb{Z}|},$$

equipped with its linear functional

$$\psi = \psi_o \oplus \psi'_o, \text{ with } \psi'_o = \psi_o \mid_{Q'_o},$$

satisfying

$$\psi(q_0^o) = \psi_o(q_0^o), \text{ and } \psi(q_{\pm n}^o) = \psi_o(q_n^o),$$

for all $n \in \mathbb{N}$.

Therefore, in fact, whenever such a family \mathbf{Q}_o of (8.1.1) is fixed in a C^* -probability space (A_o, ψ_o) , there does exist a family \mathbf{Q} of mutually orthogonal $|\mathbb{Z}|$ -many projections in a C^* -probability space (Q, ψ) (or (A, ψ) with $A \supseteq Q$), such that Q automatically satisfies (8.1.5) (and (8.1.6)).

By the above remark, the following corollary is regarded as a re-statement of the above theorem.

Corollary 8.3. Let (A_o, ψ_o) be an arbitrary C^* -probability space containing mutually orthogonal N-many projections $q_1, ..., q_N$, satisfying (8.1.3), for $N \in \mathbb{N}_{>1}^{\infty}$. Then there exists a free semicircular family $\mathcal{S}^{(N)}$ induced by $\{q_k\}_{k=1}^N$ in a certain Banach *-probability space $\mathbb{L}_O^{(N)}$.

Proof. The proof is done by (8.1.7), (8.1.8), (8.1.9), and the very above remark.

8.2 Restricted Integer-Shifts on $\mathbb{L}_{O}^{(N)}$

Let (A_o, ψ_o) be a C^* -probability space containing the family \mathbf{Q}_o of (8.1.1), satisfying (8.1.3). In Sect. 8.1, we showed that such a family \mathbf{Q}_o of mutually orthogonal N-many projections induces the free semicircular family $\mathcal{S}^{(N)}$ of (8.1.9) in a free product Banach *-probability space

$$\mathbb{L}_Q^{(N)} = (\mathbb{L}_Q^{(N)}, \tau), \text{ with } \tau = \tau \mid_{\mathbb{L}_Q^{(N)}},$$

of (8.1.8), for $N \in \mathbb{N}_{>1}^{\infty}$. Moreover, this Banach *-probability space is understood as a free-probabilistic sub-structure of the semicircular filterization \mathbb{L}_Q of (6.1.5).

Since the integer-shift group

$$\mathfrak{B} = \{\beta_e^k \in Aut(\mathbb{L}_Q) : e \in \{\pm\}, k \in \mathbb{N}_0\}$$

of (7.2.13) acts on \mathbb{L}_Q naturally, as an infinite cyclic abelian group, one can restrict the action of \mathfrak{B} on \mathbb{L}_Q to that on $\mathbb{L}_Q^{(N)}$.

Lemma 8.4. Let $\beta_e^k \in \mathfrak{B}$ be an integer-shift on \mathbb{L}_Q , and $U_l \in \mathcal{S}^{(N)}$, a semicircular element, generating $\mathbb{L}_Q^{(N)}$, for l = 1, ..., N, and suppose $\beta_e^k \mid_{\mathbb{L}_Q^{(N)}}$ is the restriction of β_e^k on $\mathbb{L}_Q^{(N)}$, also denoted simply by β_e^k . If

$$N \in \mathbb{N}_{>1} = \mathbb{N}_{>1}^{\infty} \setminus \{\infty\},\$$

then

$$\beta_e^k \left(U_l \right) = \begin{cases} U_{lek} & \text{if } 1 \le lek \le N \\ O & otherwise, \end{cases}$$
(8.2.1)

in $\mathbb{L}_Q^{(N)}$, where O is the zero element of $\mathbb{L}_Q^{(N)}$. Meanwhile, if $N = \infty$ in $\mathbb{N}_{\geq 1}^{\infty}$, then

$$\beta_{e}^{k}(U_{l}) = \begin{cases} U_{l+k} & \text{if } e = + \\ U_{l-k} & \text{if } e = -, \text{ and } l > k \\ O & \text{if } e = -, \text{ and } l \le k, \end{cases}$$
(8.2.2)

in $\mathbb{L}_Q^{(N)}$.

Proof. First, assume that $N < \infty$ in $\mathbb{N}_{>1}^{\infty}$, i.e., $N \in \mathbb{N}_{>1}$, and fix $l \in \{1, ..., N\}$ arbitrarily, and let $\beta_e^k \in \mathfrak{B}$ be an integer-shift on \mathbb{L}_Q . Let's restrict β_e^k on $\mathbb{L}_Q^{(N)}$, i.e.,

$$\beta_e^k \stackrel{denote}{=} \beta_e^k \mid_{\mathbb{L}_Q^{(N)}} \text{ on } \mathbb{L}_Q^{(N)}.$$

Then, for a semicircular elements $U_l \in \mathcal{S}^{(N)}$, generating $\mathbb{L}_Q^{(N)}$, one has that: if e = +, then

$$\beta_e^k(U_l) = \begin{cases} U_{l+k} & \text{if } l+k \le N\\ O & \text{if } l+k > N; \end{cases}$$
(8.2.3)

and if e = -, then

$$\beta_{e}^{k}(U_{l}) = \begin{cases} U_{l-k} & \text{if } l-k \ge 1\\ O & \text{if } l-k < 1, \end{cases}$$
(8.2.4)

in $\mathbb{L}_Q^{(N)}$. By (8.2.3) and (8.2.4), one obtains that

$$\beta_e^k(U_l) = \begin{cases} U_{lek} & \text{if } 1 \le lek \le N\\ O & \text{otherwise,} \end{cases}$$
(8.2.5)

in $\mathbb{L}_{Q}^{(N)}$. Therefore, the formula (8.2.1) holds by (8.2.5).

Now, assume that $N = \infty$ in $\mathbb{N}_{>1}^{\infty}$. Then, the restricted action of $\beta_e^k \in \mathfrak{B}$ on $\mathbb{L}_O^{(N)}$ satisfies that: if e = +, then

$$\beta_e^k\left(U_l\right) = U_{l+k};\tag{8.2.6}$$

if e = -, then

$$\beta_e^k(U_l) = \begin{cases} U_{l-k} & \text{if } l-k \ge 1\\ O & \text{if } l-k < 1, \end{cases}$$
(8.2.7)

in $\mathbb{L}_Q^{(N)}$. Therefore, the formula (8.2.2) is shown by (8.2.6) and (8.2.7).

The above lemma not only show how the restricted action of the integershift group \mathfrak{B} on $\mathbb{L}_Q^{(N)}$ acts on the free generator set $\mathcal{S}^{(N)}$, but also demonstrates that the restrictions of integer-shifts are no longer *-isomorphisms on $\mathbb{L}_Q^{(N)}$, in general.

Let B be an arbitrary topological *-algebra. Then the (*-)homomorphismsemigroup Hom(B) is defined to be the semigroup (under composition)

 $Hom(B) = \{f : f \text{ is a } * \text{ -homomorphism on } B\}.$

Since the zero map on B is contained in Hom(B), it cannot be a group (under composition), however, it forms a well-defined semigroup.

Notation. From below, we denote the family of restricted integer-shifts on $\mathbb{L}_Q^{(N)}$ by $\mathfrak{B}^{(N)}$, i.e.,

$$\mathfrak{B}^{(N)} = \left\{ \beta_e^k \mid_{\mathbb{L}_Q^{(N)}} \left| \begin{array}{c} \beta_e^k \in \mathfrak{B}, \text{ with} \\ e \in \{\pm\}, \ k \in \mathbb{N}_0 \end{array} \right\}.$$
(8.2.8)

Also, for convenience, we denote the restrictions $\beta_e^k \mid_{\mathbb{L}_Q^{(N)}} \in \mathfrak{B}^{(N)}$, simply by β_e^k , as above.

Lemma 8.5. Let $\mathfrak{B}^{(N)}$ be the set (8.2.8) of the restricted integer-shifts on $\mathbb{L}_Q^{(N)}$. Then

$$\mathfrak{B}^{(N)} \subseteq Hom\left(\mathbb{L}_Q^{(N)}\right),\tag{8.2.9}$$

equivalently, every element $\beta_e^k \in \mathfrak{B}^{(N)}$ is a *-homomorphism on $\mathbb{L}_Q^{(N)}$.

Proof. First, assume that $N \in \mathbb{N}_{>1}$ in $\mathbb{N}_{>1}^{\infty}$. If $\beta_e^k \in \mathfrak{B}^{(N)}$ satisfies

$$lek < 1, \text{ or } lek > N, \forall l = 1, ..., N,$$
 (8.2.10)

then such a restricted integer-shift β_e^k is identified with the zero *-homomorphism $0_Q^{(N)}$ on $\mathbb{L}_Q^{(N)}$, i.e.,

$$\beta_e^k(T) = 0_Q^{(N)}(T) = O \text{ in } \mathbb{L}_Q^{(N)},$$

for all $T \in \mathbb{L}_Q^{(N)}$, by (8.1.8), (8.2.1) and (8.2.10). And hence, all elements β_e^k of $\mathfrak{B}^{(N)}$ satisfying (8.2.10) are identified with $0_Q^{(N)}$, i.e.,

$$\beta_e^k = 0_Q^{(N)} \in Hom\left(\mathbb{L}_Q^{(N)}\right). \tag{8.2.11}$$

Suppose that $\beta_e^k \in \mathfrak{B}^{(N)}$ satisfies

$$1 \le lek \le N,\tag{8.2.12}$$

for some $l \in \{1, ..., N\}$. Then, by (8.2.1), the morphism β_e^k is a well-defined *-homomorphism on $\mathbb{L}_{Q}^{(N)}$, since it is the restriction of a *-isomorphism on $\mathbb{L}_{Q} \supseteq$ $\mathbb{L}_{Q}^{(N)}$. i.e.,

$$\beta_e^k \in Hom(\mathbb{L}_Q^{(N)}), \tag{8.2.13}$$

under (8.2.12). So, if $N \in \mathbb{N}_{>1}$, then

$$\mathfrak{B}^{(N)} \subseteq Hom\left(\mathbb{L}_Q^{(N)}\right),\tag{8.2.14}$$

by (8.2.11) and (8.2.13).

Assume now that $N = \infty$ in $\mathbb{N}_{>1}^{\infty}$. If $\beta_{+}^{k} \in \mathfrak{B}^{(N)}$, then

$$\beta_{+}^{k} \in Hom\left(\mathbb{L}_{Q}^{(N)}\right); \qquad (8.2.15)$$

if $\beta_{-}^{k} \in \mathfrak{B}^{(N)}$, then

$$\beta_{-}^{k} \in Hom\left(\mathbb{L}_{Q}^{(N)}\right),\tag{8.2.16}$$

by (8.2.2), because there are infinitely many semicircular elements $\{U_l\}_{l=1}^{\infty}$ generating $\mathbb{L}_Q^{(N)}$. However, in this case,

$$\beta_e^k \neq 0_Q^{(N)} inHom\left(\mathbb{L}_Q^{(N)}\right), \forall k \in \mathbb{N}_0, e \in \{\pm\},$$
(8.2.17)

different from (8.2.11). Indeed, for any arbitrarily fixed $\beta_{-}^{k} \in \mathfrak{B}^{(N)}$, there always exists n > k in \mathbb{N} , such that

$$\beta_{-}^{k}(U_{n}) = U_{n-k} \neq O \text{ in } \mathbb{L}_{Q}.$$

Therefore, if $N = \infty$ in $\mathbb{N}_{>1}^{\infty}$, then

$$\mathfrak{B}^{(N)} \subseteq Hom\left(\mathbb{L}_Q^{(N)}\right) \setminus \{\mathbf{0}_Q^{(N)}\},\tag{8.2.18}$$

by (8.2.15), (8.2.16) and (8.2.17).

In conclusion, if $\mathfrak{B}^{(N)}$ is the family (8.2.8) of the restricted integer-shifts on $\mathbb{L}_Q^{(N)}$, then

$$\mathfrak{B}^{(N)} \subseteq Hom\left(\mathbb{L}_Q^{(N)}\right),$$

by (8.2.14) and (8.2.18), for all $N \in \mathbb{N}_{>1}^{\infty}$.

The relation (8.2.9) shows that all restricted integer-shifts of the family $\mathfrak{B}^{(N)}$ of (8.2.8) are well-defined *-homomorphisms on $\mathbb{L}_Q^{(N)}$. However, they cannot be *-isomorphisms in general on $\mathbb{L}_Q^{(N)}$. Also, from the proof of (8.2.9), one can realize that the size of $\mathfrak{B}^{(N)}$ can be much smaller than the original integer-shift group \mathfrak{B} , especially when $N < \infty$ in $\mathbb{N}_{\geq 1}^{\infty}$. Also, the proof shows that

$$N < \infty \Longleftrightarrow 0_Q^{(N)} \in \mathfrak{B}^{(N)}$$

Definition 8.1. Let $\mathfrak{B}^{(N)}$ be in the sense of (8.2.8). We call $\mathfrak{B}^{(N)}$, the restricted(-integer)-shift family on $\mathbb{L}_Q^{(N)}$, for $N \in \mathbb{N}_{>1}^{\infty}$.

Now, consider an algebraic property of the restricted-shift family $\mathfrak{B}^{(N)}$ in the homomorphism semigroup $Hom\left(\mathbb{L}_Q^{(N)}\right)$. Recall that the integer-shift group \mathfrak{B} is a group acting on \mathbb{L}_Q in $Aut\left(\mathbb{L}_Q\right)$. How about the restricted-shift family $\mathfrak{B}^{(N)}$ in $Hom\left(\mathbb{L}_Q^{(N)}\right)$? This question can be answered by (8.2.1) and (8.2.2).

If the subset $\mathfrak{B}^{(N)}$ were an algebraic structure embedded in $Hom\left(\mathbb{L}_Q^{(N)}\right)$, then the following relation should hold;

$$\beta_{e_1}^{k_1}, \beta_{e_2}^{k_2} \in \mathfrak{B}^{(N)} \Longrightarrow \beta_{e_1}^{k_1} \beta_{e_2}^{k_2} \in \mathfrak{B}^{(N)}.$$
(8.2.19)

For instance, if $\mathfrak{B}^{(N)}$ were an algebraic structure, and if $\beta_e^k \in \mathfrak{B}^{(N)}$, then

 $\beta_+^k \beta_-^k, \beta_-^k \beta_+^k$ should be in $\mathfrak{B}^{(N)}$.

Observe now that, for a generating semicircular element $U_k \in \mathcal{S}^{(N)}$ of $\mathbb{L}_Q^{(N)}$, with

$$k \ge 1$$
, with $1 \le 2k \le N$,

in $\{1, ..., N\}$, one has that

$$\beta_{-}^{k}\beta_{+}^{k}(U_{k}) = \beta_{-}^{k}\left(\beta_{+}^{k}(U_{k})\right) = \beta_{-}^{k}(U_{k+k})$$

$$= \beta_{-}^{k}(U_{2k}) = U_{2k-k} = U_{k},$$
(8.2.20)

meanwhile, (8.2.20)

$$\beta_{+}^{k}\beta_{-}^{k}(U_{k}) = \beta_{+}^{k}(U_{k-k}) = \beta_{+}^{k}(O) = O,$$

in $\mathbb{L}_{O}^{(N)}$.

The relation (8.2.20) illustrates that two *-homomorphisms $\beta_{+}^{k}\beta_{-}^{k}$, and $\beta_{-}^{k}\beta_{+}^{k}$ are distinct *-homomorphisms in $Hom\left(\mathbb{L}_{Q}^{(N)}\right)$, i.e.,

$$\beta_{+}^{k}\beta_{-}^{k} \neq \beta_{-}^{k}\beta_{+}^{k} \text{ in } Hom\left(\mathbb{L}_{Q}^{(N)}\right).$$

$$(8.2.21)$$

The relation (8.2.21) shows that there does "not" exist $\beta_e^n = \beta_e^n \mid_{\mathbb{L}_Q^{(N)}} \in \mathfrak{B}^{(N)}$, such that either

$$\beta_e^n = \beta_+^k \beta_-^k$$
, or $\beta_e^n = \beta_-^k \beta_+^k$, in $\mathfrak{B}^{(N)}$,

by (8.2.8), i.e.,

$$\beta_{+}^{k}\beta_{-}^{k} \neq \beta_{-}^{k}\beta_{+}^{k} \notin \mathfrak{B}^{(N)}, \qquad (8.2.22)$$

 \Leftrightarrow

$$\beta_{+}^{k}\beta_{-}^{k} \neq \beta_{-}^{k}\beta_{+}^{k} \in Hom\left(\mathbb{L}_{Q}^{(N)}\right) \setminus \mathfrak{B}^{(N)}$$

i.e., the relation (8.2.19) does not hold on $\mathfrak{B}^{(N)}$, equivalently, the multiplication on *-homomorphisms is not closed (or, well-defined) on $\mathfrak{B}^{(N)}$.

Theorem 8.6. The restricted-shift family $\mathfrak{B}^{(N)}$ of (8.2.8) is a subset of $Hom\left(\mathbb{L}_Q^{(N)}\right)$, but it is not an algebraic sub-structure of $Hom\left(\mathbb{L}_Q^{(N)}\right)$.

Proof. By (8.2.9), the restricted-shift family $\mathfrak{B}^{(N)}$ is a well-defined subset of the homomorphism semigroup $Hom(\mathbb{L}_Q^{(N)})$. However, by (8.2.22), it cannot be an algebraic sub-structure of $Hom(\mathbb{L}_Q^{(N)})$.

The above theorem shows that, different from the integer-shift group \mathfrak{B} , a subgroup of the automorphism group $Aut(\mathbb{L}_Q)$, our restricted-shift families $\mathfrak{B}^{(N)}$ have no nice algebraic properties as a subset of $Hom\left(\mathbb{L}_Q^{(N)}\right)$, for $N \in \mathbb{N}_{\geq 1}^{\infty}$. However, every restricted shift $\beta_e^k \in \mathfrak{B}^{(N)}$ acts nicely on $\mathbb{L}_Q^{(N)}$, as a \ast -homomorphism.

8.3 Free Probability on $\mathbb{L}_{O}^{(N)}$ Affected by $\mathfrak{B}^{(N)}$

In this section, we consider how the restricted-shift family $\mathfrak{B}^{(N)}$ deform the original free-distributional data on the free product Banach *-probability space $\mathbb{L}_{O}^{(N)}$ of (8.1.8).

Theorem 8.7. Let $\beta_e^k \in \mathfrak{B}^{(N)}$ be a restricted shift on $\mathbb{L}_Q^{(N)}$, and let $U_l \in \mathcal{S}^{(N)}$ be a semicircular element of $\mathbb{L}_Q^{(N)}$, for $l \in \{1, ..., N\}$, for $N \in \mathbb{N}_{>1}^{\infty}$. Then the free distribution of $W_l = \beta_e^k (U_l)$ is either the semicircular law, or the zero free distribution in $\mathbb{L}_Q^{(N)}$. *Proof.* Under hypothesis, for any $N \in \mathbb{N}_{\geq 1}^{\infty}$,

$$W_l = \begin{cases} U_{lek} \in \mathcal{S}^{(N)} & \text{if } 1 \le lek \le N \\ O, & \text{otherwise,} \end{cases}$$

in $\mathbb{L}_Q^{(N)}$, by (8.2.1), (8.2.2), (8.2.9) and (8.2.19). So, if $W_l = U_{lek} \in \mathcal{S}^{(N)}$, then it is semicircular, while, if $W_l = O$, then it follows the zero free distribution in $\mathbb{L}_Q^{(N)}.$

The above theorem characterizes how the action of the restricted-shift family $\mathfrak{B}^{(N)}$ deform the semicircular law induced by $\mathcal{S}^{(N)}$; the semicircularity is deformed to be either the semicircular law, or the zero free distribution. So, one can have the following generalized result.

Theorem 8.8. Let U_{l_1} , ..., U_{l_s} be semicircular elements of $\mathcal{S}^{(N)}$, for

$$I_s = (l_1, ..., l_s) \in \{1, ..., N\}^n,$$

for $n \in \mathbb{N}$, in $\mathbb{L}_Q^{(N)}$, and let $\beta_e^k \in \mathfrak{B}^{(N)}$ be a restricted shift on $\mathbb{L}_Q^{(N)}$. Define a free (non-reduced, or reduced) word $X[I_s]$ by

$$X[I_s] = \prod_{t=1}^{s} U_{l_t} \in \mathbb{L}_Q^{(N)}.$$
(8.3.1)

Then one has either

$$\tau\left(\beta_e^k\left(X[I_s]\right)\right) = \tau\left(X[I_s]\right), \text{ satisfying } (6.2.15)$$
(8.3.2)

or

$$\tau\left(\beta_e^k\left(X[I_s]\right)\right) = 0.$$

Proof. Let $X[I_s]$ be in the sense of (8.3.1). Then it is a well-defined free random variable of $\mathbb{L}_Q^{(N)}$, as a free (non-reduced, or reduced) word in $\mathcal{S}^{(N)}$, by (8.1.8). Assume first that there exists at least one entry l_p in the *s*-tuple I_s such that

$$\beta_e^k\left(U_{l_p}\right) = O,$$

up to (8.2.1), or (8.2.2). Then, by the multiplicativity of $\beta_e^k \in \mathfrak{B}^{(N)}$,

$$\beta_e^k\left(X[I_s]\right) = \beta_e^k\left(U_{t_1}\right) \cdots \beta_e^k\left(U_{l_p}\right) \cdots \beta_e^k\left(U_{l_s}\right) = O$$

in $\mathbb{L}_{O}^{(N)}$, implying that

$$\tau\left(\beta_e^k\left(X[I_s]\right)\right) = 0.$$

Meanwhile, if

$$\beta_{e}^{k}\left(U_{l_{t}}\right)\neq O \text{ in } \mathbb{L}_{Q}^{\left(N\right)}, \text{ for all } t=1,...,s,$$

then

$$\tau\left(\beta_e^k\left(X[I_s]\right)\right) = \tau\left(X[I_s]\right),$$

by (6.3.2), (6.3.3), (7.3.2), (7.3.3), (8.2.1) and (8.2.2), because β_e^k preserves the free "reduced" word of $X[I_s]$ (as an operator) to the same-length free reduced word $\beta_e^k(X[I_s])$ with the same free-ness in $\mathbb{L}_Q^{(N)}$. Therefore, the freedistributional data (8.3.2) holds.

The above theorem generalizes Theorem 8.7 by (8.3.2). But, the proof of Theorem 8.8 illustrates that the free-distributional data (8.3.2) is dictated by the free-probabilistic information of Theorem 8.7, too.

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Several Explicit and Recurrent Formulas for Determinants of Tridiagonal Matrices via Generalized Continued Fractions

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Abstract. In the paper, by the aid of mathematical induction and some properties of determinants, the authors present several explicit and recurrent formulas of evaluations for determinants of general tridiagonal matrices in terms of finite generalized continued fractions and apply these newly-established formulas to evaluations for determinants of the Sylvester matrix and two Sylvester type matrices.

Keywords: Determinant \cdot Tridiagonal matrix \cdot Induction \cdot Explicit formula \cdot Recurrent formula \cdot Finite generalized continued fraction \cdot Sylvester matrix \cdot Sylvester type matrix

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1 Introduction

A finite generalized continued fraction is of the form

$$p_0 + \frac{q_1}{p_1 + \frac{q_2}{p_2 + \frac{q_3}{\cdots \frac{q_3}{\cdots \frac{q_{m-1}}{p_{m-2} + \frac{q_{m-1}}{p_{m-1} + \frac{q_m}{p_m}}}}},$$

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where p_0, p_1, \ldots, p_m and q_1, q_2, \ldots, q_m can be any complex numbers or functions. It can also be written equivalently as

$$p_0 + \mathop{K}\limits_{\ell=1}^m \frac{q_\ell}{p_\ell} = p_0 + \sum_{\ell=1}^m \frac{q_\ell}{|p_\ell|} = p_0 + \frac{q_1}{p_1 + p_2} + \cdots + \frac{q_{m-1}}{p_{m-1} + p_m} \frac{q_m}{p_m}.$$

In this paper, we will use the second compact form above. For more information on the theory of continued fractions, please refer to the papers [7,14] and closely related references therein.

In general, a tridiagonal matrix of order n is defined for $n \in \mathbb{N}$ by

where

$$e_{i,j} = \begin{cases} \alpha_i, & 1 \le i = j \le n; \\ \beta_i, & 1 \le i = j - 1 \le n - 1; \\ \gamma_j, & 1 \le j = i - 1 \le n - 1; \\ 0, & \text{otherwise.} \end{cases}$$

In the papers [15,16,18], the determinant $|D_n|$ and some special cases were discussed, computed, and applied to several problems in analytic combinatorics and analytic number theory. In the papers [2,5,6,9,15,16,18], there are some computation of the inverse and determinant of the general tridiagonal matrix D_n . For more information about this topic, please refer to the papers [4,8,12,13] and closely related references therein.

Let $n \geq 2$ and

$$P_{n} = (p_{i,j})_{1 \le i,j \le n} = \begin{pmatrix} a_{1} & c_{1} & 0 & 0 & \cdots & 0 & 0 & 0 & 0 \\ a_{2} & b_{2} & c_{2} & 0 & \cdots & 0 & 0 & 0 & 0 \\ a_{3} & 0 & b_{3} & c_{3} & \cdots & 0 & 0 & 0 & 0 \\ a_{4} & 0 & 0 & b_{4} & \cdots & 0 & 0 & 0 & 0 \\ \vdots & \vdots & \vdots & \vdots & \ddots & \vdots & \vdots & \vdots & \vdots & \vdots \\ a_{n-3} & 0 & 0 & 0 & \cdots & b_{n-3} & c_{n-3} & 0 & 0 \\ a_{n-2} & 0 & 0 & 0 & \cdots & 0 & b_{n-2} & c_{n-2} & 0 \\ a_{n-1} & 0 & 0 & 0 & \cdots & 0 & 0 & b_{n-1} & c_{n-1} \\ a_{n} & 0 & 0 & 0 & \cdots & 0 & 0 & 0 & b_{n} \end{pmatrix},$$
(1.2)

where

$$p_{i,j} = \begin{cases} a_i, & 1 \le i \le n, j = 1; \\ b_i, & 2 \le i = j \le n; \\ c_i, & 1 \le i = j - 1 \le n - 1; \\ 0, & \text{otherwise.} \end{cases}$$

In this paper, by the help of mathematical induction and some properties of determinants, we will present several explicit and recurrent formulas for evaluations of two determinants $|P_n|$ and $|D_n|$ and will apply these newlyestablished formulas to evaluations for determinants of the Sylvester matrix and two Sylvester type matrices.

2 Explicit and Recurrent Formulas for $|P_n|$

Right now we start off to present explicit and recurrent formulas for $|P_n|$.

Theorem 2.1. Let $n \ge 2$ and $b_k \ne 0$ for $2 \le k \le n$. Then the determinant $|P_n|$ can be computed recurrently by

$$|P_n| = \lambda_{1,n} \prod_{k=2}^n b_k,$$
 (2.1)

where

$$\lambda_{k,n} = a_k - \frac{c_k}{b_{k+1}} \lambda_{k+1,n}, \quad 1 \le k \le n-1$$
 (2.2)

and $\lambda_{n,n} = a_n$.

Proof. When n = 2, it is easy to see that

$$|P_2| = \begin{vmatrix} a_1 & c_1 \\ a_2 & b_2 \end{vmatrix} = a_1 b_2 - a_2 c_1$$

and

$$\lambda_{1,2} \prod_{k=2}^{2} b_k = \lambda_{1,2} b_2 = \left(a_1 - \frac{c_1 \lambda_{2,2}}{b_2}\right) b_2 = \left(a_1 - \frac{c_1 a_2}{b_2}\right) b_2 = a_1 b_2 - a_2 c_1 = |P_2|.$$

This means that the formula (2.1) is valid for n = 2.

Assume that the formula (2.1) validates for n = m - 1, equivalently speaking,

$$|P_{m-1}| = \lambda_{1,m-1} \prod_{k=2}^{m-1} b_k$$

When n = m, expanding the determinant $|P_m|$ according to the first rank and utilizing the assumption for n = m - 1 give

$$|P_{m}| = a_{1} \prod_{k=2}^{m} b_{k} - c_{1} \begin{vmatrix} a_{2} & c_{2} & 0 \cdots & 0 & 0 & 0 & 0 \\ a_{3} & b_{3} & c_{3} \cdots & 0 & 0 & 0 & 0 \\ a_{4} & 0 & b_{4} \cdots & 0 & 0 & 0 & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots & \vdots & \vdots & \vdots \\ a_{m-3} & 0 & 0 \cdots & b_{m-3} & c_{m-3} & 0 & 0 \\ a_{m-2} & 0 & 0 \cdots & 0 & b_{m-2} & c_{m-2} & 0 \\ a_{m-1} & 0 & 0 \cdots & 0 & 0 & b_{m-1} & c_{m-1} \\ a_{m} & 0 & 0 \cdots & 0 & 0 & 0 & b_{m} \end{vmatrix}$$
$$= a_{1} \prod_{k=2}^{m} b_{k} - c_{1} \lambda_{2,m} \prod_{k=3}^{m} b_{k} = \left(a_{1} - \frac{c_{1}}{b_{2}} \lambda_{2,m}\right) \prod_{k=2}^{m} b_{k} = \lambda_{1,m} \prod_{k=2}^{m} b_{k}.$$

By mathematical induction, we derive the formula (2.1). The proof of Theorem 2.1 is complete.

Theorem 2.2. For $n \ge 2$, the determinant $|P_n|$ can be computed explicitly by

$$|P_n| = a_1 \prod_{k=2}^n b_k - \sum_{k=2}^n (-1)^k \left(\prod_{\ell=1}^{k-1} c_\ell \prod_{m=k+1}^n b_m \right) a_k.$$
(2.3)

Proof. From the recurrent relation (2.2), it follows that

$$\begin{split} \lambda_{1,n} &= a_1 - \frac{c_1}{b_2} \lambda_{2,n} \\ &= a_1 - \frac{c_1}{b_2} \left(a_2 - \frac{c_2}{b_3} \lambda_{3,n} \right) \\ &= a_1 - \frac{c_1}{b_2} \left[a_2 - \frac{c_2}{b_3} \left(a_3 - \frac{c_3}{b_4} \lambda_{4,n} \right) \right] \\ &= \cdots \\ &= a_1 - \frac{c_1}{b_2} \left[a_2 - \frac{c_2}{b_3} \left(a_3 - \frac{c_3}{b_4} \left[a_4 - \cdots - \frac{c_{\ell-1}}{b_\ell} \left(a_\ell - \frac{c_\ell}{b_{\ell+1}} \lambda_{\ell+1,n} \right) \right] \right) \right] \right] \\ &= \cdots \\ &= a_1 - \frac{c_1}{b_2} \left[a_2 - \frac{c_2}{b_3} \left(a_3 - \frac{c_3}{b_4} \left[a_4 - \cdots - \frac{c_{\ell-1}}{b_\ell} \left(a_\ell - \frac{\cdots}{b_{\ell+1}} \lambda_{\ell+1,n} \right) \right] \right) \right] \right] \\ &= \cdots \\ &= a_1 - \frac{c_1}{b_2} \left[a_{2-\frac{c_2}{b_3}} \left(a_3 - \frac{c_3}{b_4} \left[a_4 - \cdots - \frac{c_{\ell-1}}{b_\ell} \left(a_\ell - \cdots \right) - \frac{c_{n-3}}{b_{n-2}} \left[a_{n-2} - \frac{c_{n-2}}{b_{n-1}} \left(a_{n-1} - \frac{c_{n-1}}{b_n} \lambda_{n,n} \right) \right] \right) \right] \right) \right] \\ &= a_1 - \frac{c_1}{b_2} \left[a_2 - \frac{c_2}{b_3} \left(a_3 - \frac{c_3}{b_4} \left[a_4 - \cdots - \frac{c_{\ell-1}}{b_\ell} \left(a_\ell - \cdots \right) - \frac{c_{n-3}}{b_\ell} \left[a_{2-\frac{c_2}{b_3}} \left(a_3 - \frac{c_3}{b_4} \left[a_4 - \cdots - \frac{c_{\ell-1}}{b_\ell} \left(a_\ell - \cdots \right) - \frac{c_{n-3}}{b_\ell} \left[a_{n-2} - \frac{c_{n-2}}{b_{n-1}} \left(a_{n-1} - \frac{c_{n-1}}{b_n} a_n \right) \right] \right) \right] \right) \right] \end{split}$$

$$= a_1 - \frac{c_1}{b_2} \left(a_2 - \frac{c_2}{b_3} \left[a_3 - \frac{c_3}{b_4} \left(a_4 - \dots - \frac{c_{\ell-1}}{b_\ell} \left[a_\ell - \dots - \frac{c_{n-3}}{b_\ell} \left(a_{n-2} - \frac{c_{n-2}}{b_{n-1}} a_{n-1} + \frac{c_{n-2}c_{n-1}}{b_{n-1}b_n} a_n \right) \right] \right) \right] \right)$$

$$= a_1 - \frac{c_1}{b_2} \left(a_2 - \frac{c_2}{b_3} \left[a_3 - \frac{c_3}{b_4} \left(a_4 - \dots - \frac{c_{\ell-1}}{b_\ell} \left[a_\ell - \dots - \left(\frac{c_{n-3}}{b_{n-2}} a_{n-2} - \frac{c_{n-3}c_{n-2}}{b_{n-2}b_{n-1}} a_{n-1} + \frac{c_{n-3}c_{n-2}c_{n-1}}{b_{n-2}b_{n-1}b_n} a_n \right) \right] \right) \right] \right)$$

$$= \dots$$

$$= a_1 - \sum_{k=2}^n (-1)^k \left(\prod_{\ell=2}^k \frac{c_{\ell-1}}{b_\ell} \right) a_k$$

for $n \ge 2$. Substituting this result into (2.1) and simplifying lead to (2.3). The proof of Theorem 2.2 is complete.

Remark 2.1. Applying $a_k = k$, $b_k = k$, and $c_k = k$ to the explicit formula (2.3) in Theorem 2.2 reveals

$$\begin{vmatrix} 1 & 1 & 0 & 0 & \cdots & 0 & 0 & 0 & 0 \\ 2 & 2 & 2 & 0 & \cdots & 0 & 0 & 0 & 0 \\ 3 & 0 & 3 & 3 & \cdots & 0 & 0 & 0 & 0 \\ 4 & 0 & 0 & 4 & \cdots & 0 & 0 & 0 & 0 \\ \vdots & \vdots & \vdots & \vdots & \ddots & \vdots & \vdots & \vdots & \vdots & \vdots \\ n - 3 & 0 & 0 & 0 & \cdots & n - 3 & n - 3 & 0 & 0 \\ n - 2 & 0 & 0 & 0 & n & -2 & n - 2 & 0 \\ n - 1 & 0 & 0 & 0 & 0 & n & -1 & n - 1 \\ n & 0 & 0 & 0 & \cdots & 0 & 0 & 0 & n \end{vmatrix} = \frac{1 - (-1)^n}{2} n!$$

3 Explicit and Recurrent Formulas for $|D_n|$

Now we are in a position to present explicit and recurrent formulas for $|D_n|$.

Theorem 3.1. For $n \in \mathbb{N}$, the determinant $|D_n|$ can be explicitly and recurrently computed by

$$|D_{n}| = \alpha_{1}\alpha_{2} + (\alpha_{1} - \beta_{1}\gamma_{1}) \prod_{m=3}^{n} \left[\alpha_{m} + \frac{m^{-2}}{K} \frac{(-\beta_{m-\ell}\gamma_{m-\ell})}{\alpha_{m-\ell}} \right]$$
(3.1)
$$- \sum_{k=3}^{n} \left[\prod_{\ell=1}^{k-1} (\beta_{\ell}\gamma_{\ell}) \right] \frac{\prod_{m=k+1}^{n} \left[\alpha_{m} + K_{\ell=1}^{m-2} \frac{(-\beta_{m-\ell}\gamma_{m-\ell})}{\alpha_{m-\ell}} \right]}{\prod_{m=2}^{k-1} \left[\alpha_{m} + K_{\ell=1}^{m-2} \frac{(-\beta_{m-\ell}\gamma_{m-\ell})}{\alpha_{m-\ell}} \right]}$$

and

$$|D_n| = \eta_{1,n} \left(\alpha_2 + \prod_{k=3}^n \left[\alpha_k + \frac{k-2}{K} \frac{(-\beta_{k-\ell}\gamma_{k-\ell})}{\alpha_{k-\ell}} \right] \right),$$
(3.2)

where $K_{\ell=q}^p$ for p < q is understood to be zero,

$$\eta_{1,n} = -1 - \frac{\beta_1}{\alpha_2} \eta_{2,n}, \quad \eta_{2,n} = \gamma_1 - \frac{\beta_2}{\alpha_3 - \frac{\beta_2 \gamma_2}{\alpha_2}} \eta_{3,n}, \quad \eta_{3,n} = -\frac{\gamma_1 \gamma_2}{\alpha_2} - \frac{\beta_3}{\alpha_4 - \frac{\beta_3 \gamma_3}{\alpha_3 - \frac{\beta_2 \gamma_2}{\alpha_2}}} \eta_{4,n},$$
$$\eta_{k,n} = (-1)^k \frac{\prod_{\ell=1}^{k-1} \gamma_\ell}{\alpha_2 + \prod_{\ell=3}^{k-1} \left[\alpha_\ell + K_{m=1}^{\ell-2} \frac{(-\beta_\ell - m\gamma_\ell - \ell)}{\alpha_{\ell-m}}\right]} - \frac{\beta_k}{\alpha_{k+1} + K_{\ell=1}^{k-1} \frac{(-\beta_k - \ell + 1\gamma_k - \ell + 1)}{\alpha_{k-\ell+1}}} \eta_{k+1,n}$$

for $4 \le k \le n-1$, and

$$\eta_{n,n} = (-1)^n \frac{\prod_{\ell=1}^{n-1} \gamma_\ell}{\alpha_2 + \prod_{k=3}^{n-1} \left[\alpha_k + K_{\ell=1}^{k-2} \frac{(-\beta_{k-\ell} \gamma_{k-\ell})}{\alpha_{k-\ell}} \right]}.$$

Proof. The determinant $|D_n|$ of the tridiagonal matrix D_n in (1.1) can be rewritten as

$$\begin{split} |D_n| &= \begin{vmatrix} \alpha_1 & \beta_1 & 0 & 0 & \cdots & 0 & 0 & 0 & 0 \\ \gamma_1 & \alpha_2 & \beta_2 & 0 & \cdots & 0 & 0 & 0 & 0 \\ -\frac{\gamma_1 \gamma_2}{\alpha_2} & 0 & \alpha_3 - \frac{\beta_2 \gamma_2}{\alpha_2} & \beta_3 & \cdots & 0 & 0 & 0 & 0 \\ 0 & 0 & \gamma_3 & \alpha_4 & \cdots & 0 & 0 & 0 & 0 \\ \vdots & \vdots & \vdots & \vdots & \ddots & \vdots & \vdots & \vdots & \vdots \\ 0 & 0 & 0 & 0 & \cdots & \gamma_{n-3} & \beta_{n-3} & 0 & 0 \\ 0 & 0 & 0 & 0 & \cdots & \gamma_{n-3} & \alpha_{n-2} & \beta_{n-2} & 0 \\ 0 & 0 & 0 & 0 & \cdots & 0 & \gamma_{n-2} & \alpha_{n-1} & \beta_{n-1} \\ 0 & 0 & 0 & 0 & \cdots & 0 & 0 & \gamma_{n-1} & \alpha_n \end{vmatrix} \\ \\ &= \begin{vmatrix} \alpha_1 & \beta_1 & 0 & 0 & \cdots & 0 & 0 & 0 & 0 \\ \gamma_1 & \alpha_2 & \beta_2 & 0 & \cdots & 0 & 0 & 0 & 0 \\ -\frac{\gamma_1 \gamma_2}{\alpha_2} & 0 & \alpha_3 - \frac{\beta_2 \gamma_2}{\alpha_2} & \beta_3 & \cdots & 0 & 0 & 0 & 0 \\ \frac{\gamma_1 \gamma_2 \gamma_3}{\alpha_2 \alpha_3 - \beta_2 \gamma_2} & 0 & 0 & \alpha_4 - \frac{\alpha_2 \beta_3 \gamma_3}{\alpha_2 \alpha_3 - \beta_2 \gamma_2} & \cdots & 0 & 0 & 0 & 0 \\ \vdots & \vdots & \vdots & \vdots & \ddots & \vdots & \vdots & \vdots & \vdots \\ 0 & 0 & 0 & 0 & 0 & \cdots & \alpha_{n-3} & \beta_{n-3} & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & \cdots & 0 & 0 & \gamma_{n-1} & \alpha_n \end{vmatrix} \\ = \cdots \\ \\ &= \begin{vmatrix} a_1 & \beta_1 & 0 & 0 & \cdots & 0 & 0 & 0 & 0 \\ a_2 & b_2 & \beta_2 & 0 & \cdots & 0 & 0 & 0 \\ a_3 & 0 & b_3 & \beta_3 & \cdots & 0 & 0 & 0 & 0 \\ a_4 & 0 & 0 & 4 & \cdots & 0 & 0 & 0 & 0 \\ a_{n-2} & 0 & 0 & \cdots & 0 & b_{n-2} & \beta_{n-2} & 0 \\ a_{n-1} & 0 & 0 & \cdots & 0 & 0 & b_{n-1} & \beta_{n-1} \\ a_n & 0 & 0 & \cdots & 0 & 0 & 0 & b_n \end{vmatrix} , \end{split}$$

where

$$b_{2} = \alpha_{2}, \quad b_{3} = \alpha_{3} - \frac{\beta_{2}\gamma_{2}}{b_{2}}, \quad b_{4} = \alpha_{4} - \frac{\beta_{3}\gamma_{3}}{b_{3}}, \quad \dots, \quad b_{n-3} = \alpha_{n-3} - \frac{\beta_{n-4}\gamma_{n-4}}{b_{n-4}},$$
$$b_{n-2} = \alpha_{n-2} - \frac{\beta_{n-3}\gamma_{n-3}}{b_{n-3}}, \quad b_{n-1} = \alpha_{n-1} - \frac{\beta_{n-2}\gamma_{n-2}}{b_{n-2}}, \quad b_{n} = \alpha_{n} - \frac{\beta_{n-1}\gamma_{n-1}}{b_{n-1}}$$

and

$$a_{1} = \alpha_{1}, \quad a_{2} = \gamma_{1}, \quad a_{3} = -\frac{\gamma_{2}}{b_{2}}a_{2}, \quad a_{4} = -\frac{\gamma_{3}}{b_{3}}a_{3}, \quad \dots, \quad a_{n-3} = -\frac{\gamma_{n-4}}{b_{n-4}}a_{n-4},$$
$$a_{n-2} = -\frac{\gamma_{n-3}}{b_{n-3}}a_{n-3}, \quad a_{n-1} = -\frac{\gamma_{n-2}}{b_{n-2}}a_{n-2}, \quad a_{n} = -\frac{\gamma_{n-1}}{b_{n-1}}a_{n-1}.$$

The sequences b_k and a_k for $k \geq 3$ can be formulated by finite generalized continued fractions

$$b_{k} = \alpha_{k} - \frac{\beta_{k-1}\gamma_{k-1}}{\alpha_{k-1} - \frac{\beta_{k-2}\gamma_{k-2}}{\alpha_{k-2} - \frac{\beta_{k-3}\gamma_{k-3}}{\alpha_{k-3} - \frac{\beta_{k-4}\gamma_{k-4}}{\alpha_{k-4} - \frac{\beta_{3}\gamma_{3}}{\alpha_{3} - \frac{\beta_{2}\gamma_{2}}{\alpha_{2}}}} = \alpha_{k} + \frac{k^{-2}}{K} \frac{(-\beta_{k-\ell}\gamma_{k-\ell})}{\alpha_{k-\ell}}$$

and

$$a_k = (-1)^k \frac{\prod_{\ell=1}^{k-1} \gamma_\ell}{\prod_{\ell=2}^{k-1} b_\ell}.$$

Using (2.3) results in

$$\begin{split} |D_n| &= \left(b_2 + \prod_{k=3}^n B_k\right) a_1 - \left(\beta_1 \prod_{m=3}^n B_m\right) a_2 - \sum_{k=3}^n (-1)^k \left(\prod_{\ell=1}^{k-1} \beta_\ell \prod_{m=k+1}^n B_m\right) a_k \\ &= \alpha_1 \left(\alpha_2 + \prod_{k=3}^n \left[\alpha_k + \frac{k-2}{\ell=1} \frac{(-\beta_{k-\ell}\gamma_{k-\ell})}{\alpha_{k-\ell}}\right]\right) - \beta_1 \gamma_1 \prod_{m=3}^n \left[\alpha_m + \frac{m-2}{K} \frac{(-\beta_{m-\ell}\gamma_{m-\ell})}{\alpha_{m-\ell}}\right] \\ &- \sum_{k=3}^n \prod_{\ell=1}^{k-1} \beta_\ell \prod_{m=k+1}^n \left[\alpha_m + \frac{m-2}{\ell=1} \frac{(-\beta_{m-\ell}\gamma_{m-\ell})}{\alpha_{m-\ell}}\right] \frac{\prod_{\ell=1}^{k-1} \gamma_\ell}{\prod_{\ell=2}^{k-1} \left[\alpha_\ell + K_{i=1}^{\ell-2} \frac{(-\beta_{\ell-i}\gamma_{\ell-i})}{\alpha_{\ell-i}}\right]} \end{split}$$

which can be rearranged as (3.1).

Making use of (2.1) and (2.2) yields (3.2). The proof of Theorem 3.1 is complete.

4 Discussions

In this section, we discuss our main results and related ones by several remarks.

Remark 4.1. In [3, p. 1018], it was stated that J. J. Sylvester found in 1854 that

$$|M_n(s)| = \begin{vmatrix} s & 1 & 0 & 0 \cdots 0 & 0 & 0 & 0 \\ n & s & 2 & 0 \cdots 0 & 0 & 0 & 0 \\ 0 & n-1 & s & 3 \cdots 0 & 0 & 0 & 0 \\ 0 & 0 & n-2 & s \cdots 0 & 0 & 0 & 0 \\ \vdots & \vdots & \vdots & \vdots & \ddots & \vdots & \vdots & \vdots & \vdots \\ 0 & 0 & 0 & 0 \cdots s & n-2 & 0 & 0 \\ 0 & 0 & 0 & 0 \cdots 0 & 2 & s & n \\ 0 & 0 & 0 & 0 \cdots 0 & 0 & 1 & s \end{vmatrix} = \prod_{k=0}^n (s+n-2k).$$

An application of (3.1) to $|M_n(s)|$ yields

$$\begin{split} |M_n(s)| &= s^2 + (s-n) \prod_{m=3}^n \left[s + \frac{m^{-2}}{K} \frac{-(m-\ell)(n-m+\ell+1)}{s} \right] \\ &- \sum_{k=3}^n \left[\prod_{\ell=1}^{k-1} \ell(n-\ell+1) \right] \frac{\prod_{m=k+1}^n \left[s + K_{\ell=1}^{m-2} \frac{-(m-\ell)(n-m+\ell+1)}{s} \right]}{\prod_{m=2}^{k-1} \left[s + K_{\ell=1}^{m-2} \frac{-(m-\ell)(n-m+\ell+1)}{s} \right]} \\ &= s^2 + (s-n) \prod_{m=3}^n \left[s + \frac{m^{-2}}{K} \frac{-(m-\ell)(n-m+\ell+1)}{s} \right] \\ &- n! \sum_{k=3}^n \frac{(k-1)!}{(n-k+1)!} \frac{\prod_{m=k+1}^n \left[s + K_{\ell=1}^{m-2} \frac{-(m-\ell)(n-m+\ell+1)}{s} \right]}{\prod_{m=2}^{k-1} \left[s + K_{\ell=1}^{m-2} \frac{-(m-\ell)(n-m+\ell+1)}{s} \right]} \\ &\triangleq s^2 + (s-n) \prod_{m=3}^n S(s;m,n) - n! \sum_{k=3}^n \frac{(k-1)!}{(n-k+1)!} \frac{\prod_{m=k+1}^n S(s;m,n)}{\prod_{m=2}^{k-1} S(s;m,n)} \end{split}$$

Now we try to explicitly compute

$$S(s;m,n) = s + \frac{K}{K} \frac{-(m-\ell)(n-m+\ell+1)}{s}.$$

When m = 3, it is easy to obtain that

$$S(s;3,n) = \frac{s^2 - 2(n-1)}{s} \triangleq \frac{\beta_1}{\alpha_1}.$$

When m = 4, employing the above result for S(s; 3, n), we can acquire

$$S(s;4,n) = \frac{\beta_1 s - 3(n-2)\alpha_1}{\beta_1} = \frac{s(s^2 - 5n + 8)}{s^2 - 2n + 2} \triangleq \frac{\beta_2}{\alpha_2}.$$

If assuming $S(s; k+1, n) = \frac{\beta_{k-1}}{\alpha_{k-1}}$, then, by mathematical induction, we have

$$S(s; k+2, m) = \frac{\beta_k}{\alpha_k} = \frac{\beta_{k-1}s - \alpha_{k-1}(k+1)(n-k)}{\beta_{k-1}}.$$

Note that $\alpha_{k-1} = \beta_{k-2}$. Then

$$\beta_k - \beta_{k-1}s + \beta_{k-2}(k+1)(n-k) = 0.$$

Further replacing k by k + 2 results in

$$\beta_{k+2} - \beta_{k+1}s + (k+3)(n-k-2)\beta_k = 0.$$

By the approach utilized in [15, Theorem 3.1], the characteristic equation is

$$t^{2} - st + (k+3)(n-k-2) = 0$$

which has solutions

$$t = \frac{s \pm \sqrt{s^2 - 4(k+3)(n-k-2)}}{2}$$

Consequently, it follows that

$$\beta_k = A \left(\frac{s + \sqrt{s^2 - 4(k+3)(n-k-2)}}{2} \right)^{k-1} + B \left(\frac{s - \sqrt{s^2 - 4(k+3)(n-k-2)}}{2} \right)^{k-1},$$

where

$$A = -\frac{2s^3 - 2(5n - 8) - (s^2 - 2n + 2)(s + \sqrt{s^2 - 20n + 80})}{2\sqrt{s^2 - 20n + 80}}$$

and

$$B = \frac{2s^3 - 2(5n - 8) - (s^2 - 2n + 2)(s - \sqrt{s^2 - 20n + 80})}{2\sqrt{s^2 - 20n + 80}}$$

.

In a word, we provide an alternative expression for the Sylvester determinant $|M_n(s)|$.

Remark 4.2. In [3], by virtue of left eigenvector method, the determinants

$$|M_n(s,t)| = \begin{cases} s & 1 & 0 & 0 & \cdots & 0 & 0 & 0 \\ n & s+t & 2 & 0 & \cdots & 0 & 0 & 0 \\ 0 & n-1 & s+2t & 3 & \cdots & 0 & 0 & 0 \\ 0 & 0 & n-2 & s+3t & \cdots & 0 & 0 & 0 \\ \vdots & \vdots & \vdots & \vdots & \ddots & \vdots & \vdots & \vdots \\ 0 & 0 & 0 & 0 & \cdots & s+(n-2)t & n-1 & 0 \\ 0 & 0 & 0 & 0 & \cdots & 2 & s+(n-1)t & n \\ 0 & 0 & 0 & 0 & \cdots & 0 & 1 & s+nt \end{cases}$$
$$= \prod_{k=0}^n \left(s + \frac{nt}{2} + \frac{n-2k}{2}\sqrt{t^4 + 4}\right),$$

and

$$|M_n(s,t;x,y)| = \begin{vmatrix} s & x & 0 & 0 & \cdots & 0 & 0 & 0 \\ nv & s+t & 2x & 0 & \cdots & 0 & 0 & 0 \\ 0 & (n-1)y & s+2t & 3x & \cdots & 0 & 0 & 0 \\ 0 & 0 & (n-2)y & s+3t & \cdots & 0 & 0 & 0 \\ \vdots & \vdots & \vdots & \vdots & \ddots & \vdots & \vdots & \vdots \\ 0 & 0 & 0 & 0 & \cdots & s+(n-2)t & (n-1)x & 0 \\ 0 & 0 & 0 & 0 & \cdots & 2y & s+(n-1)t & nx \\ 0 & 0 & 0 & 0 & \cdots & 0 & y & s+nt \end{vmatrix}$$
$$= \prod_{k=0}^n \left(s + \frac{nt}{2} + \frac{n-2k}{2} \sqrt{t^4 + 4xy} \right)$$

of tridiagonal matrices similar to the Sylvester matrix were collected and calculated. These evaluations can be computed alternatively by Theorem 3.1.

Remark 4.3. The condition $b_k \neq 0$ for $2 \leq k \leq n$ in Theorem 2.1 is removed off in Theorem 2.2. Therefore, the explicit formula (2.3) is better than the recurrent formulas (2.1) and (2.2).

Remark 4.4. The explicit formula (2.3) can be simply reformulated as

$$|P_n| = \sum_{k=1}^n (-1)^{k+1} \left(\prod_{\ell=1}^{k-1} c_\ell \prod_{m=k+1}^n b_m \right) a_k,$$

where the empty product is understood to be 1 as usual.

,

Remark 4.5. Let

$$U_n = (u_{i,j})_{1 \le i,j \le n} = \begin{pmatrix} \alpha_1 & \gamma_1 & 0 & 0 & \cdots & 0 & 0 & 0 & \tau_1 \\ \alpha_2 & \beta_2 & \gamma_2 & 0 & \cdots & 0 & 0 & 0 & \tau_2 \\ \alpha_3 & 0 & \beta_3 & \gamma_3 & \cdots & 0 & 0 & 0 & \tau_3 \\ \alpha_4 & 0 & 0 & \beta_4 & \cdots & 0 & 0 & 0 & \tau_4 \\ \vdots & \vdots & \vdots & \vdots & \ddots & \vdots & \vdots & \vdots & \vdots \\ \alpha_{n-3} & 0 & 0 & 0 & \cdots & \beta_{n-3} & \gamma_{n-3} & 0 & \tau_{n-3} \\ \alpha_{n-2} & 0 & 0 & 0 & \cdots & 0 & \beta_{n-2} & \gamma_{n-2} & \tau_{n-2} \\ \alpha_{n-1} & 0 & 0 & 0 & \cdots & 0 & 0 & \beta_n \end{pmatrix},$$

where

$$u_{i,j} = \begin{cases} \alpha_i, & 1 \le i \le n, j = 1; \\ \beta_i, & 2 \le i = j \le n; \\ \gamma_i, & 1 \le i = j - 1 \le n - 1; \\ \tau_i, & 1 \le i \le n - 2, j = n; \\ 0, & \text{otherwise.} \end{cases}$$

The determinant $|U_n|$ can be rewritten as

$$\begin{split} U_n| = \begin{vmatrix} \alpha_1 & \gamma_1 & 0 & 0 & \cdots & 0 & 0 & 0 & \tau_1 \\ \alpha_2 & \beta_2 & \gamma_2 & 0 & \cdots & 0 & 0 & 0 & \tau_2 \\ \alpha_3 & 0 & \beta_3 & \gamma_3 & \cdots & 0 & 0 & 0 & \tau_3 \\ \alpha_4 & 0 & 0 & \beta_4 & \cdots & 0 & 0 & 0 & \tau_4 \\ \vdots & \vdots & \vdots & \vdots & \ddots & \vdots & \vdots & \vdots & \vdots \\ \alpha_{n-3} & 0 & 0 & 0 & \cdots & \beta_{n-3} & \gamma_{n-3} & 0 & \tau_{n-3} \\ \alpha_{n-2} & 0 & 0 & 0 & \cdots & 0 & \beta_{n-2} & \gamma_{n-2} \\ \alpha_{n-1} & 0 & 0 & 0 & \cdots & 0 & 0 & \beta_n \end{vmatrix} \\ \\ = \begin{vmatrix} \alpha_1 - \frac{\alpha_n \tau_1}{\beta_n} & \gamma_1 & 0 & 0 & \cdots & 0 & 0 & 0 \\ \alpha_2 - \frac{\alpha_n \tau_2}{\beta_n} & \beta_2 & \gamma_2 & 0 & \cdots & 0 & 0 & 0 \\ \alpha_3 - \frac{\alpha_n \tau_3}{\beta_n} & 0 & \beta_3 & \gamma_3 & \cdots & 0 & 0 & 0 \\ \alpha_4 - \frac{\alpha_n \tau_4}{\beta_n} & 0 & 0 & \beta_4 & \cdots & 0 & 0 & 0 \\ \alpha_{n-2} - \frac{\alpha_n \tau_{n-3}}{\beta_n} & 0 & 0 & 0 & \cdots & \beta_{n-3} & \gamma_{n-3} & 0 & 0 \\ \alpha_{n-1} - \frac{\alpha_n \gamma_{n-1}}{\beta_n} & 0 & 0 & 0 & \cdots & 0 & 0 & \beta_{n-1} \\ \alpha_n & 0 & 0 & 0 & \cdots & 0 & 0 & 0 \\ \beta_{n-1} - \frac{\alpha_n \tau_1}{\beta_n} & \gamma_1 & 0 & 0 & \cdots & 0 & 0 & 0 \\ \alpha_3 - \frac{\alpha_n \tau_3}{\beta_n} & 0 & \beta_3 & \gamma_3 & \cdots & 0 & 0 & 0 \\ \alpha_3 - \frac{\alpha_n \tau_3}{\beta_n} & 0 & \beta_3 & \gamma_3 & \cdots & 0 & 0 & 0 \\ \alpha_4 - \frac{\alpha_n \tau_4}{\beta_n} & 0 & 0 & \beta_4 & \cdots & 0 & 0 & 0 \\ \alpha_{n-2} - \frac{\alpha_n \tau_{n-3}}{\beta_n} & 0 & 0 & \beta_4 & \cdots & 0 & 0 & 0 \\ \alpha_{n-2} - \frac{\alpha_n \tau_{n-3}}{\beta_n} & 0 & 0 & 0 & \cdots & 0 & \beta_{n-2} & \gamma_{n-2} \\ \alpha_{n-1} - \frac{\alpha_n \gamma_{n-1}}{\beta_n} & 0 & 0 & 0 & \cdots & 0 & 0 & \beta_{n-1} \end{vmatrix}$$

An application of Theorem 2.1 and 2.2 and Remark 4.4 straightforwardly yields

$$|U_n| = \sum_{k=1}^{n-2} (-1)^{k+1} (\alpha_k \beta_n - \tau_k \alpha_n) \prod_{\ell=1}^{k-1} \gamma_\ell \prod_{m=k+1}^{n-1} \beta_m + (-1)^n (\alpha_{n-1} \beta_n - \tau_{n-1} \alpha_n) \prod_{\ell=1}^{n-2} \gamma_\ell$$

and

$$|U_n| = \Lambda_{1,n-1} \prod_{k=2}^n \beta_k,$$

where

$$\Lambda_{k,n-1} = \alpha_k - \frac{\alpha_n}{\beta_n} \tau_k - \frac{\gamma_k}{\beta_{k+1}} \Lambda_{k+1,n-1}, \quad 1 \le k \le n-2$$

and $\Lambda_{n-1,n-1} = \alpha_{n-1} - \frac{\alpha_n}{\beta_n} \gamma_{n-1}.$

Remark 4.6. The determinant $|P_n|$ of P_n in (1.2) can be rearranged as

Therefore, by virtue of Theorems 2.1 and 2.2, we derive that the determinant $|Q_{n-1}|$ satisfies

$$|Q_{n-1}| = \frac{|P_n|}{a_1} = \frac{\lambda_{1,n}}{a_1} \prod_{k=2}^n b_k$$
(4.1)

and

$$|Q_{n-1}| = \frac{|P_n|}{a_1} = \prod_{k=2}^n b_k - \frac{1}{a_1} \sum_{k=2}^n (-1)^k \left(\prod_{\ell=1}^{k-1} c_\ell \prod_{m=k+1}^n b_m \right) a_k.$$
(4.2)

Further letting

$$\begin{cases} \alpha_k = b_{k+1} - \frac{a_{k+1}c_k}{a_k}, & 1 \le k \le n-1 \\ \beta_k = c_{k+1}, & 1 \le k \le n-2 \\ \gamma_k = -\frac{a_{k+2}b_{k+1}}{a_{k+1}}, & 1 \le k \le n-2 \end{cases}$$
(4.3)

in Eqs. (4.1) and (4.2) reveals

$$|D_{n-1}| = \frac{\lambda_{1,n}}{a_1} \prod_{k=2}^n b_k \tag{4.4}$$

and

$$|D_{n-1}| = \prod_{k=2}^{n} b_k - \frac{1}{a_1} \sum_{k=2}^{n} (-1)^k \left(\prod_{\ell=1}^{k-1} c_\ell \prod_{m=k+1}^{n} b_m \right) a_k.$$
(4.5)

From the second equality in (4.3), it is not difficult to see that $c_k = \beta_{k-1}$ for $2 \leq k \leq n-1$. If we can derive another relations from (4.3) to express a_k for $1 \leq k \leq n$ and b_k for $2 \leq k \leq n$ in terms of α_k for $1 \leq k \leq n-1$, β_k for $1 \leq k \leq n-2$, and γ_k for $1 \leq k \leq n-2$, then, by substituting these relations into (4.4) and (4.5), an alternative and explicit expression for evaluation of $|D_n|$ would be concluded. This is an open problem and we leave it to the interested readers.

Remark 4.7. In [1, Lemma 1.1] and [11, Lemma 2.1], it was acquired that

$$\begin{array}{c} \tau_{1} \tau_{2} \tau_{3} \tau_{4} \cdots \tau_{n-2} \tau_{n-1} \tau_{n} \\ \alpha \ \beta \ 0 \ 0 \cdots \ 0 \ 0 \ 0 \\ \gamma \ \alpha \ \beta \ 0 \cdots \ 0 \ 0 \ 0 \\ \gamma \ \alpha \ \beta \ 0 \cdots \ 0 \ 0 \ 0 \\ \vdots \ \vdots \ \vdots \ \vdots \ \ddots \ \vdots \ \vdots \ \vdots \\ 0 \ 0 \ 0 \ 0 \ \cdots \ \beta \ 0 \ 0 \\ 0 \ 0 \ 0 \ \cdots \ \alpha \ \beta \ 0 \\ 0 \ 0 \ 0 \ \cdots \ \gamma \ \alpha \ \beta \end{array} \right| = \sum_{k=1}^{n} (-1)^{k-1} \tau_{k} b^{n-k} (\beta \gamma)^{(k-1)/2} U_{k-1} \left(\frac{\alpha}{2\sqrt{\beta\gamma}}\right),$$

$$(4.6)$$

where $U_k(s)$ is the *k*th Chebyshev polynomials of the second kind, which can be generated [19,20,22] by

$$\frac{1}{1-2st+t^2} = \sum_{k=0}^{\infty} U_k(s)t^k, \quad |s| < 1, \quad |t| < 1.$$

Taking $\tau_1 = \tau_2 = \cdots = \tau_{n-1} = 0$ and $\tau_n = 1$ and reformulating, the formula (4.6) becomes

$$\begin{vmatrix} \alpha \ \beta \ 0 \cdots 0 \ 0 \\ \gamma \ \alpha \ \beta \cdots 0 \ 0 \\ 0 \ \gamma \ \alpha \cdots 0 \ 0 \\ \vdots \ \vdots \ \vdots \ \ddots \ \vdots \\ 0 \ 0 \ 0 \cdots \alpha \ \beta \\ 0 \ 0 \ 0 \cdots \gamma \ \alpha \end{vmatrix}_{n \times n} = (\beta \gamma)^{n/2} U_n \left(\frac{\alpha}{2\sqrt{\beta\gamma}}\right).$$
(4.7)
This is different from

$$\begin{vmatrix} \alpha \ \beta \ 0 \cdots 0 \ 0 \\ \gamma \ \alpha \ \beta \cdots 0 \ 0 \\ 0 \ \gamma \ \alpha \cdots 0 \ 0 \\ \vdots \ \vdots \ \ddots \ \vdots \ \vdots \\ 0 \ 0 \ 0 \cdots \alpha \ \beta \\ 0 \ 0 \ 0 \cdots \gamma \ \alpha \end{vmatrix}_{n \times n} = \begin{cases} \frac{\left(\alpha + \sqrt{\alpha^2 - 4\beta\gamma}\right)^{n+1} - \left(\alpha - \sqrt{\alpha^2 - 4\beta\gamma}\right)^{n+1}}{2^{n+1}\sqrt{\alpha^2 - 4\beta\gamma}}, & \alpha^2 \neq 4\beta\gamma \\ (n+1)\left(\frac{\alpha}{2}\right)^n, & \alpha^2 = 4\beta\gamma \end{cases}$$

and

$$\begin{vmatrix} \alpha \ \beta \ 0 \cdots 0 \ 0 \\ \gamma \ \alpha \ \beta \cdots 0 \ 0 \\ 0 \ \gamma \ \alpha \cdots 0 \ 0 \\ \vdots \ \vdots \ \vdots \ \ddots \ \vdots \ \vdots \\ 0 \ 0 \ 0 \cdots \alpha \ \beta \\ 0 \ 0 \ 0 \cdots \gamma \ \alpha \end{vmatrix}_{n \times n} = \prod_{j=1}^{n} \left(\beta + 2\alpha \sqrt{\frac{\gamma}{\alpha}} \cos \frac{j\pi}{n+1} \right)$$
(4.9)

(4.8)

which are established and collected in [15, pp. 130] and [18, Theorem 4].

Comparing (4.7) with (4.8) and (4.9), taking $\beta = \gamma = 1$ and $\alpha = 2s$, and simplifying yield

$$U_n(s) = \prod_{j=1}^n \left(1 + 2\sqrt{2s} \cos \frac{j\pi}{n+1} \right)$$
$$= \begin{cases} \frac{\left(s + \sqrt{s^2 - 1}\right)^{n+1} - \left(s - \sqrt{s^2 - 1}\right)^{n+1}}{2^{n+1}\sqrt{s^2 - 1}}, & s^2 \neq 1\\ (n+1)s^n, & s^2 = 1 \end{cases}$$

which are alternative explicit formulas for the Chebyshev polynomials of the second kind $U_n(s)$.

Remark 4.8. On 21 September 2019, we were reminded of the paper [10] in which an alternative explicit formula for elements of the inverse of a tridiagonal matrix and an efficient and fast computing method to obtain elements of the inverse of a tridiagonal matrix by backward continued fractions were investigated.

Remark 4.9. Theorem 2.2 in this paper has been applied in the proof of [17, Theorem 3.3].

Remark 4.10. This paper is a revised version of the preprint [21].

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