Chapter 9 Photocatalysis with Nanoparticles for Environmental Applications: Reactor Design Issues

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9.1 Introduction

 3×10^{24} J is the annual solar irradiance, ca. 10,000 times the worldwide energy consumption [\[1](#page-24-0), [2](#page-24-1)]. This strongly supports the use of solar power to drive chemical processes, either for the photocatalytic/photochemical synthesis of fuels or chemicals or for the photoconversion of noxious compounds.

The photodegradation of various pollutants has been proposed through photocatalytic processes thanks to the in situ formation of powerful oxidising species upon irradiation. Different approaches can involve homogeneous processes, such as the Photo-Fenton reaction [[3\]](#page-24-2) or photooxidation through UVC irradiation of H_2O_2 or heterogeneous photocatalysis. The heterogeneous photocatalytic approach has the advantage of separating and recovering the solid catalyst or using it immobilised on a separate phase with respect to the reacting mixture, limiting its impact, especially when treating wastewaters [\[2](#page-24-1), [4](#page-24-3)[–6](#page-24-4)].

By far, the most used heterogeneous photocatalyst is $TiO₂$, which can be photoexcited upon absorption of a radiation with energy larger than its bandgap. An electron is thus promoted to the conduction band of the semiconductor, leaving a hole in the valence band,

$$
\text{TiO}_2 \overset{hv \geq E_g}{\rightarrow} h^+ + e^-.
$$

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M. Piumetti, S. Bensaid (eds.), *Nanostructured Catalysts for Environmental Applications*, [https://doi.org/10.1007/978-3-030-58934-9_9](https://doi.org/10.1007/978-3-030-58934-9_9#DOI)

The hole h⁺ is a strong oxidant, with a Relative Oxidation Power ROP = 2.35 [\[7](#page-25-0)] and can oxidise directly an organic pollutant in water or interact with water forming hydroxyl radicals, characterised by slightly lower $ROP = 1.25$,

$$
h^+ + R \rightarrow R^+,
$$

$$
h^+ + H_2O \rightarrow OH^+ + H^+,
$$

$$
OH^+ + RH \rightarrow R^+ + H_2O
$$

with R being a generic organic moiety. On the other hand, O_2 can act as electron scavenger forming a superoxide radical anion, which in turn acts as oxidant,

$$
e^{-} + O_{2} \rightarrow O_{2}^{-},
$$

\n
$$
O_{2}^{-} + H^{+} \rightleftarrows HO_{2},
$$

\n
$$
2 HO_{2} \rightarrow H_{2}O_{2},
$$

\n
$$
H_{2}O_{2} + e^{-} \rightarrow OH^{+} + OH^{-}.
$$

All the radicals so formed contribute to the oxidation of the target pollutant.

Various applications of photocatalytic oxidation processes can be found in either liquid or gas phase [\[8](#page-25-1)]. However, the research is mainly focussed on the development of photoactive materials, often neglecting the careful design of the photoreactor and of the relative process. This is a key point, since the reactor geometry and size, the radiation source and its distribution across the reacting medium hugely impact the photocatalytic performance, determining the success or failure of the process and posing important issues for the scale-up. Therefore, this chapter presents some different photoreactor confgurations, predominantly focussing on water treatment, but dealing also in some cases with different reactions, e.g. H_2 production or $CO₂$ photoreduction if they present interesting reactor layouts that may suggest improvements also for water decontamination. A further section deals with modelling of different photoreactors and presents the main descriptors and mathematical approaches for their design.

9.2 Photoreactor Confgurations

The main features that an efficient photoreactor should exhibit are: (1) high surface exposure of the photocatalyst per unit volume; (2) appropriate mixing and turbulence regime to ensure fast mass transfer of the reactants and effcient suspension of the catalyst in the case of slurry confguration; (3) good transparency of the solution/ suspension to radiation and (4) effective and continuous oxygen supply for oxidation reactions.

Different types of photoreactors are described in the literature. Some examples are reported in Fig. [9.1.](#page-3-0)

The most used confguration for photoreactors is the slurry one, at least in a research stage, where the catalyst is suspended in the aqueous medium containing the pollutants and irradiated either from the external surface of the reactor or through immersed lamps. This confguration has the advantage of exploiting the intrinsic activity of the catalytic material and providing optimal active site exposure to light and to the contaminant. Provided that effcient mixing conditions are achieved, it might ensure that mass transfer limitations are negligible and that statistically all the solid particles can be reached by radiation. However, a portion of the whole solid volume is shadowed by the absorption and scattering effects of the solid itself, depending on the solid particle concentration. To limit this drawback, typically small reactor volume is adopted, obtaining a more uniform radiation pattern penetrating the whole reacting volume. Furthermore, issues for scale-up and industrial exploitation arise from the need to separate and recover the catalyst after use.

A UV Free-Surface Reactor (UV-FSR) has been used for the degradation of dyes in a pilot 10 L scale [\[10](#page-25-2)]. It is basically a stirred tank reactor, working in continuous or batch mode, with part of its surface irradiated by one or several UV lamps. Intensive stirring throughout the process provides almost ideally turbulent conditions in the aqueous reactor content. This results in a constantly renewed surface, allowing water or contaminants to meet statistically the photon source, independent of colour and turbidity of the suspension or solution. It was tested with or without a heterogeneous photocatalyst as AOP for water treatment.

A fully new concept of high pressure slurry photoreactor, operating up to 20 bar pressure, has been developed. This prototype, with ca. 1.5 L capacity, has been used at the moment to improve the $CO₂$ solubility in water, boostening the reaction rate for CO_2 photoreduction and for the photoreforming of organic molecules for H_2 production $[11–18]$ $[11–18]$ $[11–18]$. However, the possibility to increase significantly O_2 concentration in water makes it very promising for also oxidation reactions in water treatment.

The choice between the use of suspended catalysts in a slurry arrangement or immobilised materials is hard. On one hand, a slurry confguration is preferable to maintain unaltered the intrinsic activity of the photocatalyst, since after deposition, some decrease of performance is often noticed. However, in order to achieve continuous operation when using slurry type reactors, as required for water treatment, catalyst separation, recovery and recycle are needed. A possible option is to use a selective membrane in a hybrid membrane photoreactor, as recently reviewed [\[5](#page-24-5), [19\]](#page-25-5). In this arrangement, a slurry type photoreactor is used, with suitable irradiation, and a selective membrane is immersed in it to collect the clarifed solution as permeate, retaining the photocatalyst in the reactor [[20\]](#page-25-6). A fouling problem may arise for the membrane, since, depending on the surface charge on the photocatalyst particles, agglomeration may occur, altering the membrane properties [\[21](#page-25-7)]. A possible solution consists in adapting the solution pH based on surface charge measurements.

Different photoreactor confgurations have been reviewed recently, for application to Advanced Oxidation processes (AOPs) [\[2](#page-24-1)]. A very simple option to overcome the limitations of slurry reactors and improve the throughput is to adopt a fxed bed confguration, which is possible by immobilisation of the active phase over a transparent material, such as glass beads [[1\]](#page-24-0). However, transparency is limited and the surface interaction with the liquid is less efficient than slurry.

The possibility to use photomicroreactors, where the catalyst is immobilised in different forms over a solid substrate, allows us to overcome the downstream separation problems of the catalyst powder and to irradiate a clear and transparent solution, but also implies some limits for the decrease of active site exposure with respect to slurries.

Photocatalytic microreactors have been developed with signifcant surface to volume ratio $(>10,000 \text{ m}^2/\text{m}^3)$ to overcome the latter limitation. Another very interesting feature that underpins the use of microreactors is that when optimised fow patterns are achieved, heat and mass transfer limitations are negligible and a uniform and predictable irradiation can be achieved [[2\]](#page-24-1). Different fow patterns in three-phase microreactors are reported for the fuid fow through the microchannels, such as bubble flow [\[22](#page-25-9)], slug flow [\[23](#page-25-10)] or annular flow [[24\]](#page-26-0). Microreactors can assume different shape, such as microcapillary [\[25](#page-26-1)], single- [\[26](#page-26-2)] or multimicrochannel reactors [\[27](#page-26-3)] or planar structure [\[28](#page-26-4)].

An example of capillary microreactor has been proposed by Eskandarloo et al. [\[26](#page-26-2)] by coating a stainless steel plate including microchannels with the photocatalyst to accomplish the oxidation of terephthalic acid under UV irradiation. Additionally, Castedo et al. [\[29](#page-26-5)] reported the photoproduction of H_2 with Au/TiO₂ immobilised in a silicon microreactor, which guarantees higher versatility.

Monolith-type microreactors are an interesting confguration, which consists of a regular array of channels, usually straight and coated with a thin layer of catalyst, irradiated effciently though optical fbres located inside the channels. The fbres should not be coated to distribute uniformly the radiation across the monolith (Fig. [9.2](#page-5-0)) [\[30](#page-26-6)].

Different confgurations have been developed. For instance, a photomicroreactor containing TiO₂ nanoparticles coated over ZnO nanorod arrays has been obtained by growing the active phase in the capillaries [\[31](#page-26-7)]. The assembly was used to photodegrade methylene blue. In addition, the mass and photon transfer have been improved for the reduction of Cr(VI) to Cr(III) comparing two photoreactors: a monolithic tubular photoreactor and a micro-meso-structured photoreactor, achieving in both cases a satisfactory irradiated surface to volume ratio [\[32](#page-26-8)]. To reduce photon transfer limitations, the tubular photoreactor was packed with transparent cellulose acetate monoliths loaded with the catalyst by dip-coating. For the other layout, a thin flm was uniformly deposited on a glass slab or on a network of channels and chambers printed in the back stainless steel slab, achieving much better performance than the monolith reactor.

A completely different confguration is achieved in a twin reactor, used, e.g. to achieve the photoreduction of $CO₂$ by CO , which is a stronger reducing agent than water [[33\]](#page-26-9). It consists of two compartments separated by an ion exchange membrane,

Fig. 9.2 Example of honeycomb-type micro-photoreactor, irradiated through optical fbres. (Reproduced from [\[30\]](#page-26-6) by kind permission. Copyright (2016) Elsevier)

in the specific case allowing H^+ transfer between the two reacting sections, e.g. Nafion. The same concept has been applied to the $CO₂$ photoreduction assisted by $H₂$ evolution, where the latter is accomplished by photoreforming organic pollutants in water [\[1](#page-24-0)]. The same twin reactor confguration has also been applied by Baniasadi et al. [[34\]](#page-26-10) for the photoproduction of H_2 under visible light. Two different catalysts, one for H_2 evolution (Pt/SrTiO₃:Rh) and the other for oxygen reaction (WO₃), were separately loaded in two compartments, separated by a Nafon membrane, and using the Fe(II)–Fe(III) couple as electron transfer agent. A Compound Parabolic Concentrator (CPC) is also included. The reactor performance is not only dependent on the fow regime, needing high turbulence, but also sensitive to the high pressure drop under that conditions.

The conversion of the Direct Red 23 (DR23) dye has been successfully achieved using solar light with $Fe₃O₄/TiO₂$, coated on a glass tube and inserted in an annular shaped photoreactor. This assembly is also mounted in a CPC confguration and allows almost quantitative abatement of the dye [\[35](#page-26-11)].

A solar Offset Multi Tubular Photoreactor (OMTP) is a variant of the CPC photoreactor layout, including additional tubes in the space occupied by the axes of intersection of the CPC refective involutes. This allows a considerable increase of the irradiated reactor volume (by 79%) and of the fuid residence time (up to 1.8 factor) with respect to a conventional CPC with the same footprint. With this assem-bly, the degradation efficiencies significantly increased for various pollutants [[36\]](#page-26-12).

A combined solar-electro-Fenton approach has been applied for the pilot scale removal of different pesticides and emerging contaminants [[37\]](#page-26-13). The selected reaction is based on a redox cycle mediated by the Fe^{2+}/Fe^{3+} couple, where the ferrous ions react with H_2O_2 to provide 'OH radicals as Reactive Oxygen Species (ROS). The electrochemical formation of H_2O_2 in situ avoids the handling, storage and transportation of this reactant. Furthermore, the irradiation with UVA radiation, or

better, with solar light, promotes the regeneration of $Fe²⁺$ with concomitant production of additional 'OH through photoreduction of Fe³⁺, the photodegradation of possible Fe(III)-carboxylate complexes formed as intermediates and the additional direct photolysis of the pollutants or intermediates. The pilot reactor was supplied with a CPC concentrating the solar radiation fux to the solution. The pilot scale reactor (Fig. [9.3\)](#page-6-0), the largest available for this application, allowed 100 L capacity with natural solar irradiation in the facility of the Plataforma Solar de Almería (Spain). The electrochemical section of the device was constituted of four plateand-frame electrochemical reactors, while the irradiated area was 2 m^2 for a 23 L volume of solution.

A similar reactor confguration on a pilot scale has been used for the photo-Fenton degradation of a solution of tetracyclines, which are the second largest group of antibiotics worldwide, largely excreted in water reservoirs through hospital or

Fig. 9.3 Front view of (**a**) the four flter-press type electrochemical cells of the pilot unit and (**b**) the CPC photoreactor. In (**c**), schematic diagram of the pilot unit equipped with one cell (examined in this work), showing (*1*) CPC photoreactor, (*2*) valve, (*3*) feed tank, (*4*) power supply, (*5*) electrochemical reactor, (*6*) liquid fowmeter, (*7*) air compressor and (*8*) magnetic pump. (Reproduced from [\[37\]](#page-26-13) by kind permission. Copyright (2019) by Elsevier)

domestic wastes, besides through industrial residue. Given the ineffective removal through conventional biological treatments, to avoid the insurgence of bacterial resistance, a pilot-scale photoreactor constituted of a CPC and 6 borosilicate glass tubes in series with a batch recirculating circuit was used to treat 35 L of a tetracycline solution, achieving almost 90% conversion in 2 h under different operating conditions [\[38](#page-26-14)]. Also, a CPC system for hydrogen production under solar light irradiation has been designed and modelled by Cao et al. [\[39](#page-26-15)].

A further different confguration is constituted by spinning disc reactors, which are proposed as tools for process intensifcation [[40–](#page-27-0)[43\]](#page-27-1). The reactor is composed of a spinning disc placed in the horizontal plane, which is let to rotate at different spinning velocity. The reacting mixture is injected from a central nozzle and moves towards the periphery due to centrifugal draft. The photocatalyst may be immobilised on the disc and irradiation can be achieved from top. The intensifcation of the process mainly occurs through a more uniform and effective penetration of light through the thin liquid layer formed over the disc and by improvement of the mass transfer coeffcient. This device has been applied to the photodegradation of methylene blue and of dehydroabietic acid, a resin acid found in pulp and paper wastewater, which acts as endocrine disruptor, with adverse ecological effects even at very low concentration. The overall reaction rate has been measured at different feeding fow rates and spinning regimes, obtaining widely different results and even a change of the apparent reaction order. The latter parameter is expected as frst order for most photocatalytic reactions, but it turned out to be second order for the selected applications in very specifc cases. The mass balance of a Continuously Stirred Tank Reactor, perfectly mixed but unsteady due to recirculation, has been coupled with a characteristic parameter of the spinning disc device, i.e. the liquid flm thickness over the disc, calculated though the Nusselt model [\[40](#page-27-0)]. The kinetics was interpreted in correlation with the fow pattern of the liquid on the disc. Under specifc rotation speed and fow rate combinations, large standing waves were observed through a camera, disrupting in an array of irregular waves at large rotational speeds. Overall, the change in the fow structure affected the reaction rate, with the biggest intensifcation of mass transfer rate with the highest non-linearity of the wave pattern. However, an additional detrimental effect is due to possible scattering and modifcation of the light penetration in the case of large waves. The comparison of performance with an annular reactor layout revealed a signifcant gain of conversion of methylene blue, though at the expense of a more complex system from the mechanical point of view, due to moving parts, with consequent heavier mainte-nance issues [\[41](#page-27-2)].

In order to improve mass transfer limitations, a different confguration has also been proposed, in form of a multiblade impeller, on which the photocatalyst has been coated. This system has been named Stacked Frame Photoreactor (SFPR) and is sketched in Fig. [9.4](#page-8-0) [[44\]](#page-27-3). In this photoreactor, the propeller rotation helps mixing and thus a more effcient mass transfer with respect to immobilised photocatalytic surfaces.

An impinging jet stream photoreactor was set up with a $TiO₂$ -coated disc for the abatement of phenol in water. Different parameters were optimised, such as liquid

Fig. 9.4 Schematic of the SFPR with the rotating titanium propeller housed within the reactor. (Reprinted from [[44](#page-27-3)] by kind permission. Copyright (2019) of American Chemical Society)

Fig. 9.5 Schematic drawing of TiO₂ "water-bell" photoreactor. (Reproduced from [[48](#page-27-6)] by kind permission. Copyright (2018) by Elsevier)

fow rate, disc diameter, nozzle-to-disc distance and initial pollutant concentration. A newly developed hydrodynamic model, Langmuir-Hinshelwood kinetics and mass transfer phenomena were used to rate the system and it was compared with a spinning disc reactor, showing better performance [[45–](#page-27-4)[47\]](#page-27-5).

A "water-bell" TiO_2 -based photoreactor (Fig. [9.5\)](#page-8-1) has been designed to achieve a thin water/photocatalyst film efficiently irradiated through UVA lamps, envisaging a future application with solar light [\[48](#page-27-6)]. The thin flm has the advantage of exchanging oxygen efficiently with surrounding air and of allowing high photocatalyst concentration, since high light penetration is not needed due to small flm thickness. This also allows loading the catalyst on coarse supporting particles, to ensure easy settling and recovery of the catalyst in a so-called clarifying zone. The system is designed at a large scale as a modular set of nozzles that spray the water/photocatalyst suspension as a flm to a basin for recirculation. The catalyst loading proved to be a more sensitive parameter than irradiance and the presence of an effcient electron scavenger also demonstrated important, as well as the pH.

A falling flm reactor has been tested for the photodegradation of methylene blue. Ca. 4 g of TiO₂ were deposited in a four step configuration exposing ca. 0.2 m² surface to UVA irradiation. The mechanism of oxidation has been investigated, together with the optimisation of the reaction conditions [[49\]](#page-27-7). Furthermore, falling flm microreactors irradiated with LED lamps were constituted by parallel microchannels for photocatalytic-assisted synthesis [[2,](#page-24-1) [50\]](#page-27-8).

Similarly, Bahbani Ghatar et al. [[51\]](#page-27-9) investigated the same reaction in a UVAirradiated plane falling flm photoreactor based on ZnO. Also, in this case, the optimisation of the operating conditions was the objective of the work, using an experimental design over different variables: the slope of the plate, the number of lamps, their distance from the plate and the fow rate of wastewater. The combined effect of all these parameters over the observable result (dye conversion) was expressed as a function of the single parameters or combined in couples. Only some of the variables revealed particularly signifcant, as determined by ANOVA analysis. For instance, the distance of the lamps from the plate was not signifcant by itself, but the combination of distance/number of lamps and distance/water fow rate revealed important. Furthermore, the water fow rate was signifcant by itself and in combination with other parameters, since it can modify the flm thickness of the solution and turbulence can affect the light absorbance.

The FluHelik photoreactor was used for a photochemical UVC/H₂O₂ degradation of various emerging contaminants [[52\]](#page-27-10). Tests have been carried out in both ultrapure water matrix and a real urban wastewater collected after secondary treatment. The FluHelik reactor showed 1.3 times higher mineralisation with respect to a reference reactor. Furthermore, the residual toxicity of the photodegraded products was evaluated through embryo toxicity bioassays. Furthermore, a miliphotoreactor (NETmix) was used for the degradation of oxytetracycline from urban wastewater [[53\]](#page-27-11). It is constituted of an array of small cylindrical chambers and prismatic transport channels and the system is assembled with a quartz slab. The optimisation of the radiation source, ensuring uniform light distribution, and the H_2O_2 concentration was performed. The same type of reactor was also used with a heterogeneous catalytic system, based on $TiO₂$ irradiated with a UVA source [\[54](#page-27-12)] and for bromate reduction [[55\]](#page-28-0).

The photodegradation of VOCs in air has been tested on a semi-pilot scale in a 420 L photoreactor consisting of a photocatalytic chamber that included a set of 9 W UVA fuorescent lamps and a photocatalytic flter. The latter was assembled on a cellulose-based material on which a titania-silica coating has been immobilised.

The air flow containing the pollutant was fed and recirculated to a first larger chamber, where sampling points allowed us to measure the concentration of the target molecules vs. time. The flow rate was measured before entering the photocatalytic module and air circulation was provided by a fan [[56\]](#page-28-1).

A summary of the different possible confgurations is reported in Table [9.1.](#page-10-0)

Reactor type	Features	Advantages	Disadvantages
Slurry	Catalyst in suspension; intense mixing needed	Good surface exposure; limited mass transfer if efficient mixing is achieved	Scale-up issues; inefficient light penetration due to scattering or opaque solutions; recovery of catalyst powder
Fixed bed	Catalyst immobilised over beads	Immobilised catalyst; no recovery issues	Light penetration is inefficient
Microreactors/ capillary reactors	Catalyst immobilised in microchannels exposed to light	Good catalyst exposure to light and reactants; limited size of the channels and so limited mass transfer issues	Needed high surface to volume ratio
Monoliths	Catalyst immobilised in microchannels of a monolith exposed to light	Good surface to volume ratio; effective irradiation with LED or optical fibres in each channel	Need of complex arrangement for irradiation; complex coating procedures
Twin reactor	Separate compartments for different half reactions	Separate production of different products	Applications limited by needs of transport through a membrane; rates possibly limited for the same reason
Single or multitubular with compound parabolic collectors (CPC)	Effective collection and concentration of radiation on the tubes	Enhanced irradiation; relatively compact and simple design for multitubular systems	Sizing dependent on reaction rate
Spinning disc	Solution/suspension spread in thin film irradiated on top	Short optical path; limited mass transfer	Possible irregular flow shape; mechanically more complex due to moving parts
Stacked frame/ rotating impeller	Catalyst immobilised over rotating blades	Improved mass transfer	Mechanically more complex due to moving parts; irradiation spreading and averaging over the surface
Water bell/falling film	Solution/suspension recirculated in thin layers flowing	Improved mass transfer; effective irradiation	Need of high recirculation; intrinsically unsteady operation

Table 9.1 Main reactor configurations

9.3 Photoreactor Scale-Up and Modelling Issues

One of the key points for modelling photochemical or photocatalytic reactors is the need to include the radiation as a "pseudo-reactant" in the kinetics of the process. This allows us to integrate the characteristic equations for the reactor used (mass, momentum and heat balances) to size or rate the photoreactor. Computational methods in photocatalysis have been very recently summarised and include methods for the solution of the radiation pattern [[57\]](#page-28-2).

The frst issue is to estimate the Local Volumetric Rate of Photon Absorption (LVRPA, *ea*), i.e. a sort of "concentration of photons". LVRPA expresses the useful photon concentration as they were reactants and thus the useful number of photons absorbed by the surface of the catalyst in the unit volume of the reactor [\[58](#page-28-3)].

The radiative transport equation (RTE) [\[59](#page-28-4), [60\]](#page-28-5) represents the most rigorous approach,

$$
\frac{dI_{\lambda}(s,\Omega)}{ds} = -k_{\lambda}I_{\lambda}(s,\Omega) - \sigma_{\lambda}I_{\lambda}(s,\Omega) + \frac{1}{4\pi}\sigma_{\lambda}\int_{0}^{4\pi}p(\Omega'\to\Omega)I_{\lambda}(s,\Omega')d\Omega'.
$$

*I*_λ represents the irradiance at wavelength $λ$, *s* is a spatial coordinate and $Ω$ represents a directional solid angle. The properties of the photocatalytic material are instead represented by an absorption coefficient k_{λ} and a scattering coefficient σ_{λ} , both depending on the wavelength. The term to be integrated $p(\Omega' \rightarrow \Omega)$ represents a scattering phase function, i.e. the probability that around the position *s*, a photon can be redirected from the direction Ω' to Ω . Therefore, the first term on the right represents the absorption, the second the out-scattering and the third the in-scattering of photons.

Different approaches were implemented to solve the RTE, where a rigorous Discrete Ordinate Method (DO) discretises the volume coupling and conventional mass and energy balances of Computational Fluid Dynamic (CFD) methods, with this additional radiation balance. More precisely, the knowledge of the hydrodynamics is needed to calculate the catalyst distribution in slurry reactors. CFD simulation typically solves the fow pattern, from which the time-averaged or steady state results of catalyst distribution are used to calculate the radiation feld through the RTE [[9\]](#page-25-8).

In the case of homogeneous photoreactors, the Navier-Stokes equations are used, while for multiphase fow, an Eulerian-Eulerian (E-E) or Eulerian-Lagrangian (E-L) approach can be applied. The E-E model solves the mass and momentum balance equations for the different phases, considered as interpenetrated continua. Additional relations are needed in the case of solid phases. The E-E approach has been more frequently implemented for multiphase fows in this application [[61–](#page-28-6)[66\]](#page-28-7). The E-L approach simulates the trajectory of the dispersed phase particles through the equations of motion, by applying a force balance, while the continuous phase is solved through an Eulerian approach. Different turbulence models have been used for bubble columns and photocatalytic reactors, among which the κ-ε one, in case accounting for bubble induced turbulence, is a good compromise between accuracy and computational demand [[9\]](#page-25-8). On the contrary, for immobilised catalysts, the nearwall zone is computed by discretisation under laminar conditions.

The Discrete Ordinate (DO) model transforms the RTE, which is an integrodifferential equation, into a system of algebraic equations [\[67](#page-28-8), [68](#page-28-9)]. The radiation feld is divided into a number of discrete directions for which the RTE is solved. Alternatively, a Finite Volume (FV) method is more fexible with respect to geometries, while the DO can span the whole optical thickness, has dedicated non-grey models and has been successfully applied to both immobilised and slurry photoreactors. A very nice example of application of CFD modelling to photoreactor design is proposed by Boyjoo et al. [[66\]](#page-28-7), where a DO approach is used, also implemented as a FV method in a commercial software for CFD, such as Fluent®. The reaction order with respect to the LVRPA showed variable depending on light intensity, i.e. 0.5 at high irradiance (due to more frequent parasitic electron-hole recombination events) and 1 for low irradiance. Following an E-E model, the phases are modelled as interpenetrating continua for the solution of the mass and momentum conservation balances. The rate of reaction is expressed as

$$
-r_{\rm p} = kf(c_{\rm p}) \text{LVRPA}^m,
$$

where $m = 0.5-1$ depending on the regimes and $f(c_n)$ was a function of the concentration of pollutant to be treated corresponding to a frst order kinetics or a Langmuir-Hinshelwood expression. The radiation modelling was carried out through DO modelling, by dividing the radiation feld into discrete directions and solving the RTE separately for each of them. The FV tool in Fluent® was also used, by discretising each control volume (e.g. 72 solid angles). However, the simultaneous solution of the radiation feld and the fuid dynamics is not currently possible in the commercial tool. So, at frst, the simulation was made on the phases, getting the mean catalyst concentration and velocity vectors. Then, the radiation was modelled, accessing the catalyst local concentration data previously stored, to calculate the optical properties of the medium and LVRPA. The latter was then used to model the reaction rate and, consequently, the reactor balances. Different photoreactor layouts were tested, holding 1, 2 or 4 lamps, with optimised, different grids for the fuid dynamic simulation and the radiation modelling [[66\]](#page-28-7).

CFD modelling was also applied to the SFPR reactor depicted in Fig. [9.4](#page-8-0) using a rotating machinery turbulent fow κ-ε model [[44\]](#page-27-3). The rays emitted by a LED lamp were modelled as deriving from a point source and so, depending on the distance from the $TiO₂$ -coated propeller, the amount of radiation reaching the surface was calculated. As reasonable, the latter parameter decreased upon increasing the distance from the lamp. However, another important feature was highlighted through the model. It was previously reported that rapidly switching LED sources can improve the activity due to more controlled excitation, which would decrease the recombination of the photogenerated charges [\[44](#page-27-3), [69–](#page-28-10)[71\]](#page-28-11). This effect can be achieved through the rotation of the propeller, which alternates the surfaces exposed to light. The calculated irradiance at different lamp-reactor distance and with different positions of the propeller is exemplifed in Fig. [9.6](#page-13-0). The effect of the lamp distance from the photocatalyst layer is immediately evident between 4 cm (Fig. [9.6a,](#page-13-0) [b\)](#page-13-0) and 7 cm (Fig. [9.6c, d\)](#page-13-0). The shadowing zones are also evidenced during the rotation of the photocatalyst.

Similarly, a CFD approach allowed us to optimise the confguration of a UV LED array, tested for the photodegradation of cinnamic acid as model for agroindustrial waste water [[72\]](#page-28-12). LEDs may be usually represented as point sources and can now reach high intensity irradiance, which makes the reaction rate order with respect to radiation intensity non-linear. Thus, the accurate prediction of light distribution (again with fuid dynamic assessment of catalyst and reactant distribution across the reactor) is a fundamental task for optimisation. Different reactor geometries were also tested [\[73](#page-29-0)].

Fig. 9.6 3D surface plots of the irradiance on the catalyst surface at 4 cm lamp/reactor distance in both the initial confguration (**a**) and the rotated propeller position (**b**) or at 7 cm distance in the initial static position (**c**) and the rotated propeller one (**d**). (Reproduced from [\[44\]](#page-27-3) by kind permission. Copyright (2019) by American Chemical Society)

The best piping arrangement in a pilot annular photoreactor has been modelled, considering a FluHelik reactor and a conventional Jet reactor. The residence time and radiation distribution have been calculated by CFD for the degradation of 3-amino-5-methylisoxazole with H_2O_2/UVC and UVC processes [[74\]](#page-29-1).

Various reactor layouts were modelled through CFD. For instance, as a nonexhaustive list, a channelled optical fbre reactor was applied to ethylene photodegradation in gas phase [\[75](#page-29-2)] and for its oxidation in non-ideal fxed bed fow reactors [\[76](#page-29-3)]; a baffed fat plate photoreactor irradiated with UVA-LED was used for the degradation of $Fe(CN)_6$ ⁻, a recalcitrant compound present in mining wastewaters [\[77](#page-29-4)]; the FV module implemented in Fluent® was used to model a single phase annular photoreactor [[78\]](#page-29-5); an annular photocatalytic reactor to degrade Rhodamine B was simulated using Fluent[®] [\[79](#page-29-6)]; the photodegradation of acetaldehyde in air has been modelled in a multi-tubular reactor [[80\]](#page-29-7).

Alternatively, a Monte Carlo (MC) stochastic approach can be used to solve the RTE. For both the DO and MC approaches, the rigorous solution is obtained at the expense of computational cost, so that typically, these methods are employed for constant irradiance sources, such as lamps, avoiding solar photoreactors for which the variation of the input source would be scarcely manageable.

The MC method [\[58](#page-28-3)] is based on the random scattering and absorption phenomena (Fig. [9.7,](#page-14-0) right), which can be modelled in a slab reactor through the defnition of a grid of thickness d*x* and by correlating the trajectory of each photon in the slab through an array R of numbers, which assumes random values between 0 and 1 . A mean free path (*l*) of the photon is defned as the free pathway between two consecutive absorption or scattering events. This is correlated with the extinction coefficient of the suspension (β) as

Fig. 9.7 *Left:* Scattering directions considered by the SFM in a differential volume of liquid containing a particle. *Right:* schematic diagram of the slab photoreactor with absorbing and scattering particles suspended in the inside. I_0 , I_R and I_T are the incident, reflected and transmitted intensity of light, *L* is the reactor length and *l* the free path of photons. (Reproduced from [[58](#page-28-3)] by kind permission. Copyright (2016) by Elsevier)

$$
l=-\frac{1}{\beta}\ln R,
$$

l and the direction θ _I where the photon is emitted allow the calculation of the new position (x_i) of the photon, which should lie between 0 and L; otherwise, the photon exits the slab and it is no more useful,

$$
x_i = x_{i-1} + l \cos \theta_i.
$$

When the free path is passed, the photon is absorbed or scattered based on the value of *R*: the photon is absorbed and its location recorded if $R \ge \omega$, where ω is defned as the scattering albedo, i.e. the ratio between the scattering and the extinction coefficients, defining the scattering probability. If $R < \omega$, the photon is scattered towards a new random direction, whose longitude and latitude angles (azimuth *α* and zenith φ , respectively) are found randomly in the range $[0; 2\pi]$. φ is expressed as a function of *R* and of the phase function, while the azimuth as

$$
\alpha_{i+1} = 2\pi R.
$$

Finally, the scattering angle is

$$
\cos\theta_{i+1} = \cos\varphi_{i+1}\cos\theta_i + \sin\varphi_{i+1}\cos\alpha_{i+1}\sin\theta_i.
$$

The radiative heat transfer in a multi-channel solar reactor was performed through MC on a single channel level [\[81](#page-29-8)]. A coating of ZnFe_2O_4 was considered to optimise the channel geometry and flm thickness.

Another alternative approach is the Six Flux Absorption-Scattering Model (SFM), which provides a simplifed solution adequate for most practical purposes. The latter has been proposed for homogeneous photoreactors by Brucato et al. [[82\]](#page-29-9), with pioneering simplifed models to predict radiation distribution in various reactor layouts [[83–](#page-29-10)[87\]](#page-29-11). The SFM model admits scattering in the six directions of a Cartesian space, defned as forward (f), backward (b) and sideward (s), with relative probability p_i (i = f,b,s) (Fig. [9.7\)](#page-14-0). p_i depends on the phase function for scattering and has been variously calculated depending on the type of scattering that is assumed [\[58](#page-28-3), [82,](#page-29-9) [88,](#page-29-12) [89\]](#page-30-0). For heterogeneous photocatalysts, the determination of the absorption and scattering coeffcients is a fundamental step for the reliability of the model and some estimates of the relevant parameters for $TiO₂$ is reported by Acosta-Herazo et al. [\[58](#page-28-3)], together with a very effective comparison between the MC and SFM approaches.

A SFM approach has been used to evaluate the LVRPA in a planar photoreactor, distinguishing the radiation fraction in UVA and Vis regions when coupling materials able to absorb different spectral ranges and therefore, they contribute differently to absorption and scattering. In the case of polychromatic sources, such as solar radiation, the photoactive materials should be carefully analysed to predict their optical properties and to calculate the relevant parameters. In the example reported by Acosta-Herazo et al. $[90]$ $[90]$, TiO₂ absorbs and scatters light in the UVA region, while goethite absorbs both UVA and part of the Visible spectrum. The SFM is therefore applied to polychromatic radiation and the net radiation absorbed by the goethite photocatalyst was calculated summing the contributions from the UVA and the Visible radiation. So the net radiation absorbed by goethite was calculated summing both contributions,

$$
e_{\text{TiO}_2}^a(x) = e_{\text{UVA}}^a(x) \left(\frac{W}{m^3}\right),
$$

$$
e_{\text{goethite}}^a(x) = e_{\text{UVA}}^a(x) + e_{\text{Vis}}^a(x) \left(\frac{W}{m^3}\right).
$$

The optical properties of both materials were averaged in the relevant spectral range and depend on catalyst concentration, e.g.

$$
\kappa^* = \frac{\int_{\frac{\lambda_{\min}}{\lambda_{\min}}}^{\lambda_{\max}} \kappa_{\lambda}^* I_{\lambda} d\lambda}{\int_{\frac{\lambda_{\min}}{\lambda_{\min}}}^{\lambda_{\max}} I_{\lambda} d\lambda} \left(\frac{m^2}{kg}\right),
$$

$$
\sigma^* = \frac{\int_{\frac{\lambda_{\min}}{\lambda_{\min}}}^{\lambda_{\max}} I_{\lambda} d\lambda}{\int_{\frac{\lambda_{\min}}{\lambda_{\min}}}^{\lambda_{\max}} I_{\lambda} d\lambda} \left(\frac{m^2}{kg}\right),
$$

$$
g = \frac{\int_{\frac{\lambda_{\min}}{\lambda_{\min}}}^{\lambda_{\max}} I_{\lambda} d\lambda}{\int_{\frac{\lambda_{\min}}{\lambda_{\min}}}^{\lambda_{\max}} I_{\lambda} d\lambda} \left(-\right),
$$

where κ is the absorption coefficient, σ is the scattering coefficient, while *g* is a parameter used to compute the phase function (e.g. defned according to Henyey-Greenstein, but different approaches are also reviewed in [\[9](#page-25-8)]). The limits for integration depend on the absorption edge of each material.

The optical thickness of an absorbing medium can also be calculated, assuming a characteristic length *L* of the system (e.g. slab length or annulus thickness), as [\[9](#page-25-8)]

$$
\tau_{\rm abs} = \kappa_{\lambda} L.
$$

The particle agglomeration, if occurring near the isoelectric point of the material, can also be an issue, which is accounted for by correcting the extinction coeffcient (*β*) with respect to aggregation,

$$
\beta_{\lambda}^* = \kappa_{\lambda}^* + \sigma_{\lambda}^*,
$$

$$
\beta_{\text{ager},\lambda}^* = \beta_{\lambda}^* \frac{\ln T(t, W_{\text{cat}})}{\ln T(t_1, W_{\text{cat}})},
$$

$$
\sigma_{\text{ager},\lambda}^* = \beta_{\text{ager},\lambda}^* - \kappa_{\lambda}^*,
$$

where t_1 is the time in which the suspension is formed and T is the transmission, depending on catalyst concentration (W_{cat}) and, in the case of particles aggregation, on time *t*,

$$
T(t, W_{\text{cat}}) = \frac{\text{Transmitted radiation flux at time } t \text{ and catalyst loading } W_{\text{cat}}}{\text{Transmitted radiation flux at time } t \text{ and pure water}}.
$$

Different parameters have been introduced to analyse and compare the performance of photoreactors, such as the Initial Rate of Photon Absorption (IRPA),

$$
IRPA = -\left[\frac{d}{dx}\left(e^{a}(x)\right)\right]_{x=0},
$$

which may result in the following equation:

$$
IRPA = -I_0 c_{cat}^2 \psi \quad \left(\frac{W}{m^4}\right),
$$

$$
\psi = \left(\beta^*\right)^2 \frac{a^2 \left(1 - \omega_{con}^2\right)}{\left(1 - \gamma\right) \omega_{con}} \left[\gamma \left(\omega_{con} - 1 - \sqrt{1 - \omega_{con}^2}\right) - \left(\omega_{con} - 1 + \sqrt{1 - \omega_{con}^2}\right)\right]
$$

with $\beta^* = \sigma^* + \kappa^*$ being specific extinction coefficient. The IRPA is very sensitive to catalyst concentration and thus, the optimal catalyst amount has to be calculated depending on its optical properties and be in turn related to the reactor thickness *L* to maximise the absorption of the incoming radiation.

Another useful indicator is the Total Rate of Photon Absorption (TRPA), which quantifes the total radiation absorbed in the whole volume,

$$
\frac{\text{TPRA}}{A} = \int_0^L e^a(x) \, \mathrm{d}x.
$$

Also, in this case, useful parameters relative to $TiO₂$ are reported by Acosta-Herazo et al. [\[90](#page-30-1)].

The SFM approach has also been used to model the LVRPA in a CPC photoreactor for the abatement of acetaminophen, a hugely used drug [[91\]](#page-30-2), by using a Langmuir-Hinshelwood kinetic expression. The reactor mass balance has been solved as a series of plug-fow and continuous stirred reactors to treat 5000 L of wastewater (daily amount of a medium-size Colombian hospital).

Different lamp models were proposed, as line source, surface or volume emitters [\[9](#page-25-8)], whose equations are summarised in the cited reference.

In the case of immobilised photocatalysts (planar coatings, monoliths, etc.), there is no scattering and absorption at difference with slurry systems (if a clear and transparent medium is treated). In this case, the RTE is calculated through the Beer-Lambert equation where the incident light on the photocatalyst surface (I_0) is calculated by subtracting the reflected part from the inlet radiation intensity $(I_0 = I_{in} - I_r)$. Possibly, some scattering may be due to air bubbles, if present [\[9](#page-25-8)].

Once the model for the source and the integration of the RTE gives quantitative indications on the useful radiation amount available as a "pseudo-reactant", the kinetic equation can be derived. At frst, intrinsic kinetics is considered, i.e. neglecting the possible physical limitations due to mass transport in the fuid phase or in the porous solid particles. Subsequently, these factors should be added.

In general, the kinetic equation can be written as a function of the concentration of the reactants, usually the pollutant concentration in the case of water treatment. The function is commonly frst order or expressed in Langmuir-Hinshelwood form,

$$
-r = kf(c_i),
$$

$$
k = f(pH, T, c_{O_2}, W_{cat}, LVRPA).
$$

The dependence on the radiation intensity is often expressed through a "reaction order" of the LVRPA, which depends on different parameters, as reviewed else-where [\[9](#page-25-8)]. Furthermore, photon absorption with varying catalyst concentration has been modelled by Cao et al. [\[92](#page-30-3)].

A generalised model was developed to rate solar CPCs with different arrangement for degradation of water pollutants. A modifed Langmuir–Hinshelwood kinetics has been developed, together with a new model to quantify the "effective" quantum yield. It is computed from the optical properties of the catalyst and the incident photon fux, without dependence on the composition, the operating conditions, the geometry and the scale of the reactor [[93\]](#page-30-4).

The RTE has been integrated with a Dirichlet boundary condition on the flm surface exposed to irradiation, [[60\]](#page-28-5)

$$
\frac{dI_{\lambda}(s)}{ds} = -k_{\lambda}I_{\lambda}(s).
$$

$$
I_{\lambda,s} = I_{\lambda,0}.
$$

It can be integrated to yield

$$
I_{\lambda,s}=I_{\lambda,0}e^{-k_{\lambda}s}.
$$

The mass transport of each species has been modelled under steady state conditions according to Fick's frst law, as diffusive fux *Ni*,

$$
N_{i} = -D_{\text{eff}} \nabla c_{i},
$$

$$
D_{\text{eff}} = \frac{D_{ij} \varepsilon}{\tau}.
$$

Take into account the effective diffusivity D_{eff} as a function of the diffusion coefficient of species *i* in the reaction medium $j(D_{i,j})$, the porosity (ε) and the tortuosity factor $(τ)$ of pores constituting the photocatalyst.

Accordingly, the rate of conversion of species *i* was defned assuming a frst order dependence on concentration and an order *γ* on light intensity, as

$$
r_i = -\eta k_{r, \text{app}} c_i I_{\lambda, \text{s}}^{\gamma} (1 - \varepsilon).
$$

The effectiveness of the catalyst (η) has been conventionally defined as

$$
\eta = \frac{\tanh(\phi)}{\phi},
$$

where the Thiele modulus for a frst order reaction was formulated as

$$
\phi = (\delta - y) \sqrt{\frac{k_{r, \text{app}} I'_{\lambda, 0}}{D_{\text{eff}}}},
$$

representing as usual the ratio between the potential rate of reaction and that of diffusion inside the pores of the solid. A discussion on the suggested values for the various parameters can be found in [[60\]](#page-28-5) and references therein. The adopted model neglects diffusional limitations in the liquid phase, which is reasonable for monolith or microchannel photoreactors operating at high Reynolds number. The normalised reaction rate (r/r_{max}) was modelled depending on the thickness of the photocatalyst coating for the two front side (FSI) and back side (BSI) irradiation (Fig. [9.8b, c\)](#page-20-0). The study returned the optimal value of the thickness for both cases and the attenuation of light intensity depending on the latter parameter. A correlation for flm thickness optimisation depending on the various operating parameters was also proposed $[60]$ $[60]$.

Additionally, $CO₂$ photoreduction has been investigated in the presence of a more powerful reducing agent than H_2O , studying the effect of different operating parameters in a twin reactor to predict the productivity of methanol under natural irradiance exposure [[33\]](#page-26-9). A model for the reactor has been derived considering the reaction rate as a function of irradiance, where higher *m* exponents mean a better utilisation of photons in the given reactor geometry,

Fig. 9.8 (**a**) Top, Back Side (BSI) and bottom, Front Side (FSI) Irradiation of photocatalyst coating. (**b**, **c**) Normalised reaction rate vs. photocatalyst flm thickness (*δ*) for BSI and FSI arrangement, respectively. (Reproduced from [\[60\]](#page-28-5) by kind permission. Copyright (2017) by Elsevier)

$$
r_j = I^m\bigg(k_j \prod_{i=1}^N c_i^{v_i}\bigg).
$$

Mass transfer was computed according to the two-flm theory and the solubility following Henry's coeffcient, calculated keeping into account the effect of the ionic strength of the solution. Sunlight was modelled using dedicated software. It was concluded that temperature and pressure were determinant factors to improve the mass transfer coefficients and gas solubility in water, respectively [[13,](#page-25-11) [16,](#page-25-12) [17\]](#page-25-13).

The optimal photocatalyst layer thickness has been calculated for the degradation of methylene blue as model reaction in microphotoreactors [\[60](#page-28-5)]. The comparison between two irradiation modes in immobilised catalyst layers, i.e. BSI or FSI, was compared, developing a radiation model to be included in the reaction kinetics. The fuid dynamic properties of the fuid in the channel were also considered. In the BSI arrangement, the maximum incident radiation and the maximum pollutant concentration occur on the opposite side of the assembly, while in FSI, they occur on the same side of the photocatalytic flm (as they were in a counter- or co-current arrangement, respectively) (Fig. [9.8\)](#page-20-0).

A similar comparison was done for the gas phase photodegradation of *n*-decane as model compound for VOCs. In such a case, however, no signifcant difference between FSI or BSI illumination was noticed, thanks to the high porosity of the catalytic layer and high diffusivity of the reactant [[94\]](#page-30-5).

Kinetic modelling has been considered by Charles et al. [[95\]](#page-30-6) in a microreactor assuming a Langmuir-Hinshelwood expression to represent the degradation of salicylic acid, with K being the adsorption constant of the compound to oxidise and C_s its surface concentration,

$$
r = \frac{kKC_s}{1+KC_s}.
$$

The residence time distribution was evaluated in the microreactor with the auxilium of an inert tracer and compared with a plug-fow reactor. The mass transfer coefficient k_{max} has been calculated through the Sherwood number (Sh), thanks to an existing correlation developed by Ergu et al. [[95,](#page-30-6) [96](#page-30-7)] as a function of the Reynolds (Re) and Schmidt (Sc) numbers,

$$
Sh = \frac{k_{m,ex}d_h}{D},
$$

\n
$$
Re = \frac{\rho ud_h}{\mu},
$$

\n
$$
Sc = \frac{\mu}{\rho D},
$$

\n
$$
Sh = 2.076 \times Re^{0.28} \left(\frac{L}{d_h}\right)^{-0.12} \times Sc^{0.33},
$$

where μ is the fluid viscosity, ρ its density, u the linear velocity, D the diffusion coefficient and d_h the characteristic length of the system, e.g. the channel hydraulic diameter.

Different arrangements of micro- and meso-photoreactors have been proposed, either for scale-up or numbering-up purposes [[97\]](#page-30-8) for the degradation of dyes or organic pollutants. The fow systems have been characterised according to the fow rate, Reynolds, Damköhler and Peclet numbers, also identifying the defnition of pertinent efficiency quantifiers. At first, the irradiance (I) has been simply defined as dependent on the optical absorbance of the medium, obtained multiplying the extinction coefficient ε , the concentration of the absorbing species c and the optical path (l) . The incident radiation I_0 can be measured through a radiometer,

$$
I = I_0 e^{-\varepsilon c l}.
$$

The authors considered strictly laminar fow regimes. The second Damköhler (Da_{II}) number was calculated considering the ratio between the reaction and the mass transfer rates as

$$
Da_{II} = \frac{\text{reaction rate}}{\text{diffusive mass transfer rate}} = \frac{k''H}{D}
$$

with *H* being the channel height and $k'' = k(V/A)$ the reaction rate constant corrected by the ratio between the reactor volume and surface area of the catalyst. In addition, the Peclet (Pe) number is introduced to quantify the relative roles of advection and diffusive mass transport, with *L* being the characteristic length for advection [\[97](#page-30-8)],

$$
\text{Pe} = \frac{\text{diffusion time}}{\text{advection time}} = \frac{H^2}{L/2}.
$$

To quantify the efficiency of the process, different concepts can be used. According to IUPAC, the correct statement of quantum yield would be the ratio between the number of reactive events and the number of photons absorbed, which are impractical to determine. A more direct parameter is defned as apparent quantum yield, which compares the rate of conversion of the target molecule with the rate of photons incident on the catalyst surface. However, other indexes of efficiency may be related to the energy consumption to degrade a given amount of pollutant. For instance, IUPAC defined the Electrical Energy per Order (E_{E0}) as the electrical energy required for the degradation of a contaminant by one order of magnitude, in a unit reactor volume of wastewater [[97,](#page-30-8) [98\]](#page-30-9) for both batch and continuous flow operation (P is the power consumption, t the time, V the reactor volume and *F* the volumetric flow rate),

$$
E_{\text{EO-Batch}} = \frac{1000 \times P \times t}{V \times \log \left(\frac{c_0}{c_t}\right)} \quad \text{or} \quad E_{\text{EO-Continuous}} = \frac{P}{F \times \log \left(\frac{c_0}{c_t}\right)}.
$$

Alternatively, a Reactor Characterisation (R_c) number has been introduced as a factor for selecting the best UV source from an economic point of view during dye degradation in a pilot scale photoreactor [\[10](#page-25-2)],

$$
R_{\rm c} = \frac{\text{input power}(W) \times \text{UV} - \text{C efficiency}(\%)}{\text{Reactor Volume}(L)}.
$$

As said, homogeneous Fenton or Photo-Fenton approach can be used to abate pollutants from waste water [[3,](#page-24-2) [99,](#page-30-10) [100](#page-30-11)]. The kinetic model of this reaction, applied to the oxidation of paracetamol, has been proposed including a form of LVRPA. The kinetic expression was expressed as.

$$
\big[\boldsymbol{R}(x,t)\big]\big[\boldsymbol{R}^{T}(x,t)\big]\!+\!\Phi\!\sum_{\lambda}e_{\lambda}^{a}(x,t)\tau(x,t),
$$

where $R(x,t)$ is the vector of formation/conversion rates of each product/reactant, $\mathbf{R}^T(x,t)$ is the contribution to the rate of thermal (dark) reactivity, Φ is the quantum yield averaged over the wavelength and the summation computes the LVRPA in polychromatic conditions. $\tau(x,t)$ is the term that completes the expression of the reaction rates, as detailed in [\[101](#page-30-12)].

This expression of the rate is therefore included in the mass balance of the reactor, setting $C = C⁰$ at $t₀ = 0$ as initial condition and considering the volume averaged quantity indicated between 〈 〉,

I. Rossetti et al.

$$
\frac{\mathrm{d} \mathbf{C}}{\mathrm{d} t} = \mathbf{R}^T(t) + \frac{V_{\text{IRR}}}{V_{\text{TOT}}} \Phi \sum_{\lambda} e_{\lambda}^a(x,t) \tau(x,t).
$$

To compute the spatial distribution of the radiation, a line source radiation model was introduced, assuming a spherical and isotropic emission,

$$
e_{\lambda}^{a}(x,t) = \kappa_{\lambda}(x,t) \frac{P_{\lambda s}}{2\pi L_{L}} \int_{\theta_{i}}^{\theta_{2}} \exp \left[-\frac{\kappa_{T,\lambda}(x,t)(r_{i} - r_{int})}{\cos \theta}\right] d\theta,
$$

where $P_{\lambda s}$ is the lamp emission power, $\kappa_{\lambda}(x,t)$ and $\kappa_{\lambda}(x,t)$ are the volumetric absorption coeffcient of the reacting species and of the medium, respectively, *r* the radius and L_L the useful length of the lamp [\[101](#page-30-12)].

The overall quantum yield is a parameter that can allow the comparison of different catalysts and to assess the overall feasibility of the photocatalytic process. It may be calculated as the ratio between the pollutant degradation reaction rate and the LVRPA (both volume averaged) [[9\]](#page-25-8).

An interesting defnition of energy and exergy balances has been proposed [[34\]](#page-26-10), though for the photoproduction of H_2 and not for the abatement of pollutants. At first, the irradiance of solar light was determined as follows: K_λ measures monochromatic intensity of the radiation and *n* represents the numbering of intervals of wavelengths selected,

$$
j_s = 4.329 \times 10^{-5} \pi \sum_{n} (K_{\lambda} \Delta \lambda)_{n},
$$

$$
q_e = \alpha_s \varepsilon_b \left(T^4 - T_0^4 \right),
$$

where q_e is the energy exchanged between the reactor and the surroundings, α_s is the absorptivity of the reactor surface and ε_b is the Boltzmann constant. The energy effciency of the process has been calculated as

$$
\eta = \frac{n_{\text{H}_2} \text{HHV}_{\text{H}_2}}{\gamma \left[\alpha_{\text{v}} j_{\text{v}} + \alpha_{\text{nv}} \left(j_{\text{s}} - j_{\text{v}} \right) \right] + n_{\text{H}_2 \text{O}} h_{\text{H}_2 \text{O}} + 0.5 V_{\text{H}_2 \text{O,in}}^2}
$$

with *γ* being radiation weakening factor, α _v and α _{nv} the absorptivity of the catalyst in the visible region and the non-visible one, *n* the moles of each compound, HHV the higher heating value, *V* the volume and j_y the irradiance in the visible range. The results revealed the highest effciency at the minimum visible light absorptivity and a signifcant energy and exergy loss when increasing the temperature, suggesting the need of an efficient heat exchange system.

The modelling of falling flm reactors has been accomplished under turbulent flow regime considering a pilot scale photoreactor, $1.25 \text{ m} \times 1 \text{ m}$ in size [[102\]](#page-30-13). Different papers considered a laminar regime, but in order to keep the photocatalyst in suspension, turbulence was needed, with $Re > 15,000$. This induces waves to be statistically described since they determine the falling flm thickness and, hence, the absorption of radiation by the suspension. The velocity profles were then estimated by momentum balances and the LVRPA was solved according to a SFM approach.

9.4 Conclusions

The scale-up of photochemical or photocatalytic reactors is a required step to exploit advanced oxidation processes and similar remediation processes to treat water effuents. Besides searching for active materials, it is important to describe reactor geometry, fow patterns and solid properties in order to understand the interaction between the reactants, the catalyst and the third actor, i.e. radiation.

Different reactor layouts are available, yet conceptualised at a pilot scale, to maximise the exposure to solar light and to limit mass transfer and radiation distribution issues. Various modelling tools are also available to describe the reaction rate as dependent on reactants, catalyst and radiation distribution across the photoreactor.

Broader efforts are needed for radiation modelling and its coupling to reactor design. Specifc attention should be put on microreactor devices and on those confgurations, which reached a pilot scale development to allow validation and further improvement.

Tools which couple CFD with radiation modelling should be implemented to provide rigorous methods to scale-up photoreactors.

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